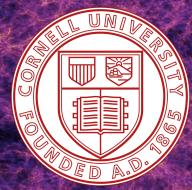
# Testing gravity with cosmology: efficient simulations, novel statistics and analytical approaches



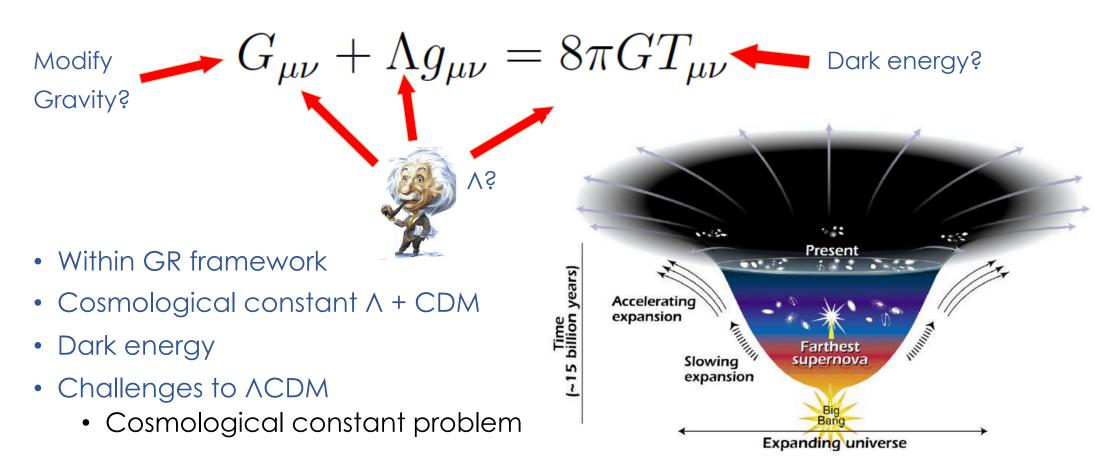
Georgios Valogiannis Cornell University

Berkeley Center for Cosmological Physics September 3<sup>rd</sup>, 2019

### Origins of cosmic acceleration



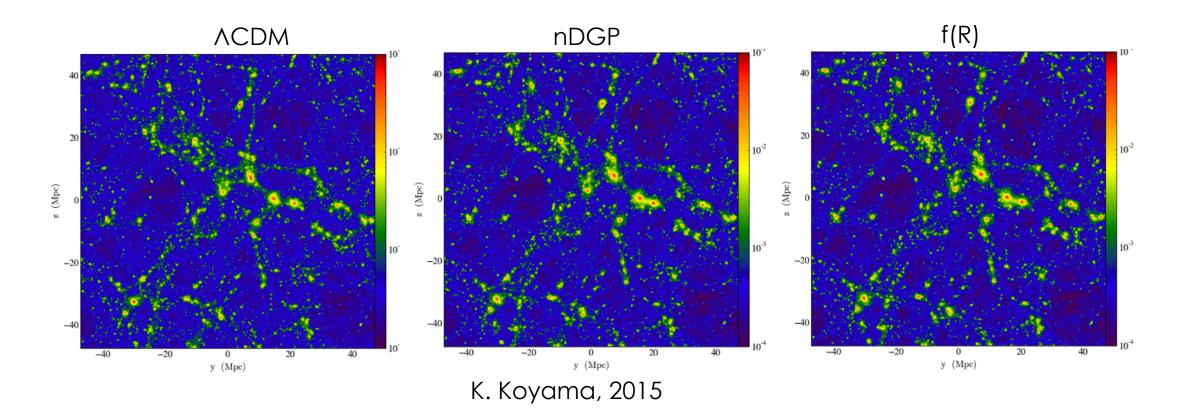
• Recent accelerative epoch poses great theoretical challenge



#### Cosmological scales as probes of viable MG models



- Alternative proposal modified gravity + screening
- MG-ACDM degeneracy broken at cosmological scales
- Upcoming LSS surveys can constrain MG models



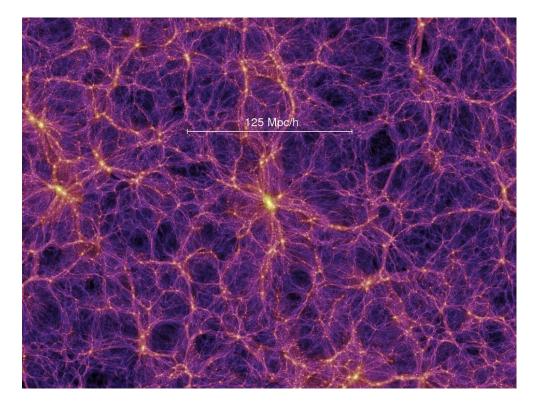
### Theoretical modeling of LSS in MG necessary



- Model-dependent cosmological tests of gravity from upcoming data
- Theoretically model structure formation in MG

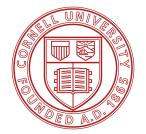
#### However:

- N-body simulations computationally expensive
  - Efficient simulation schemes
- Detection challenging
  - Novel statistics
- Surveys observe galaxies
  - Biased tracers
  - Redshift space distortions



Millennium Simulation, 2005

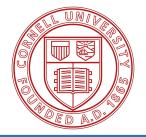
#### Efficient COLA hybrid scheme



#### Efficient simulation schemes

• Accurate but expensive N-body simulations vs Fast but approximate perturbation theory Why not combine? S. Tassev et al. 2013 Hybrid "COmoving Lagrangian Acceleration" COLA scheme  $m{x}_{
m res} \equiv m{x} - m{x}_{
m LPT}$   $\partial_t^2 m{x}_{
m res} = abla \Phi - \partial_t^2 m{x}_{
m LPT}$ 0.8 S. Tassev et al. 2013 COLA PM 0.2 2LPT 2LPT solution N-body 0.0 0.1 0.5 k[h/Mpc]

### **COLA simulations for MG: chameleons**



In MG: non-minimally coupled scalar  $\boldsymbol{\phi}$ 

N body component

$$\ddot{\mathbf{x}} + 2H\dot{\mathbf{x}} = -\frac{1}{a^2} \left( \nabla_{\mathbf{x}} \Phi_N + \frac{\beta}{M_{Pl}} \nabla_{\mathbf{x}} \phi \right)$$
  
Fifth force

 $\begin{aligned} & LPT \text{ component} \\ & \ddot{D}_1(\mathbf{k}, a) + 2H\dot{D}_1(\mathbf{k}, a) = \frac{3}{2}\Omega_m(a)H^2D_1(\mathbf{k}, a)\frac{G_{eff}}{G} \\ & \text{MG linear growth} \\ & \text{factor} \end{aligned}$ 

- MG introduces new fifth force
- Need to solve highly nonlinear Klein-Gordon equation
- Computationally expensive!

- LPT growth factor becomes scale-dependent
- Also computationally expensive!
- More about this part later

### Solutions: Effective screening implementation



#### Speed-up by 2 order of magnitude!

#### N body component

- Solve linearized KG
- Attach thin shell factor to KG fifth force

$$\ddot{\mathbf{x}} + 2H\dot{\mathbf{x}} = -\frac{1}{a^2} \left( \nabla_{\mathbf{x}} \Phi_N + \frac{\Delta \mathcal{R}_c}{\mathcal{R}_c} \frac{\beta}{M_{Pl}} \nabla_{\mathbf{x}} \phi \right)$$

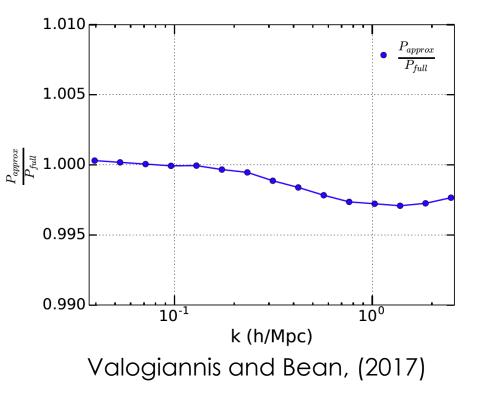
• Phenomenological factor

$$\frac{\Delta \mathcal{R}_c}{\mathcal{R}_c} = \begin{cases} \frac{\phi(a)}{2\beta(a)M_{Pl}|\Phi_N|} & \text{if } \frac{\phi(a)}{2\beta(a)M_{Pl}|\Phi_N|} < 1\\ 1 & \text{if } \frac{\phi(a)}{2\beta(a)M_{Pl}|\Phi_N|} > 1 \end{cases}$$

H. Winther & P. Ferreira, 2014

#### LPT component

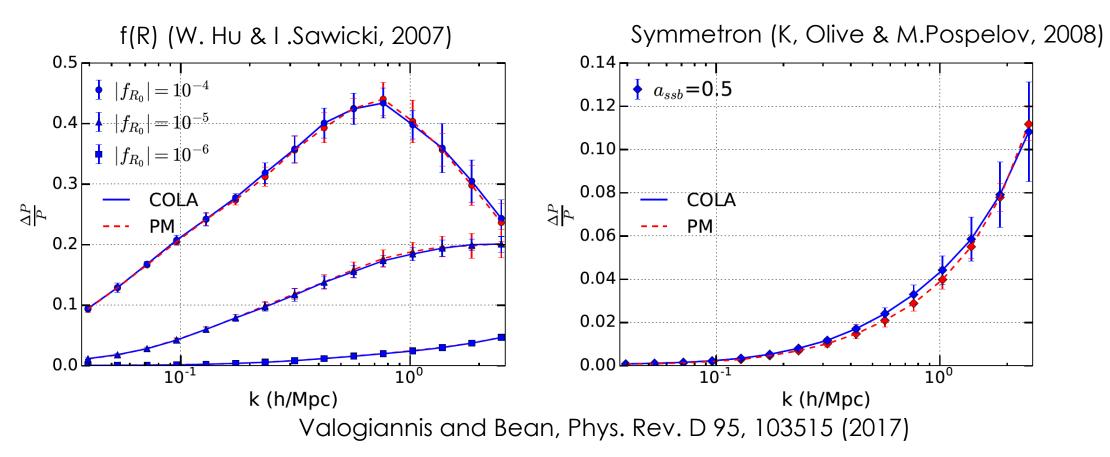
• Scale-dependent growth negligible!





#### Comparison of N-Body and COLA power spectra

- MG COLA hybrid agrees well with N-body results!
- ~100x faster!
- Emulators

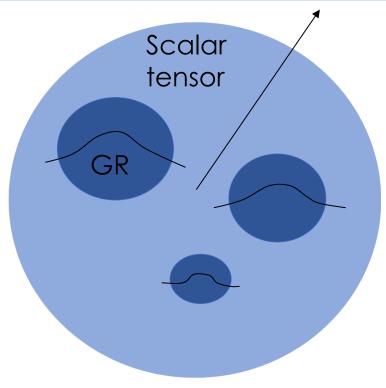




#### Detection challenging - need for a new statistic

Novel statistics

- "Screening" suppresses deviations in high densities
- Signals strongly suppressed by screening detection challenging
- Need for new statistic!



K. Koyama 2017

#### Marked density transformation



- Up-weighting low density, unscreened regions and down-weight highly screened regime can highlight MG signals in density fields
- Fundamental quantity of interest  $\delta(\mathbf{x}, a) = \frac{\rho_m(\mathbf{x}, a)}{-1}$

• Variety of density transformations in literature 
$$ar{
ho}_m$$

• Logarithmic re-mapping (M. Neyrinck et al. 2009)

 $\delta' = \ln{(\delta + 1)}$ 

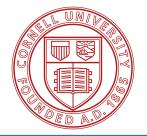
• Clipping density field (F. Simpson et al. 2011)

$$\delta' = \delta_c = \begin{cases} \delta & \text{if } \delta < \delta_0 \\ \delta_0 & \text{if } \delta > \delta_0 \end{cases}$$

• «Marked» transformation (M. White, 2016)

$$\delta' = m(\delta) = \left(\frac{\rho_* + 1}{\rho_* + \rho_m}\right)^p = \left(\frac{\rho_* + 1}{\rho_* + \bar{\rho}_m(\delta + 1)}\right)^p$$

#### Quantifying enhancement



- Dark matter N-body simulations using Particle-Mesh (PM) code (Valogiannis & Bean 2017)
- Simulation box side L=200 Mpc/h, 256<sup>3</sup> particles, resolved on 512<sup>3</sup> grid
- 40 density snapshots at z=0 for  $\Lambda CDM$ , f(R) and symmetron cosmologies
- 2D projection  $\rightarrow$  3x40=120 power spectra
- Covariance matrix

$$C_{ij} = \frac{1}{N_{seed} - 1} \sum_{r}^{N_{seed}} \left( P_r(k_i) - \bar{P}(k_i) \right) \left( P_r(k_j) - \bar{P}(k_j) \right)$$

• Fisher information in the parameter

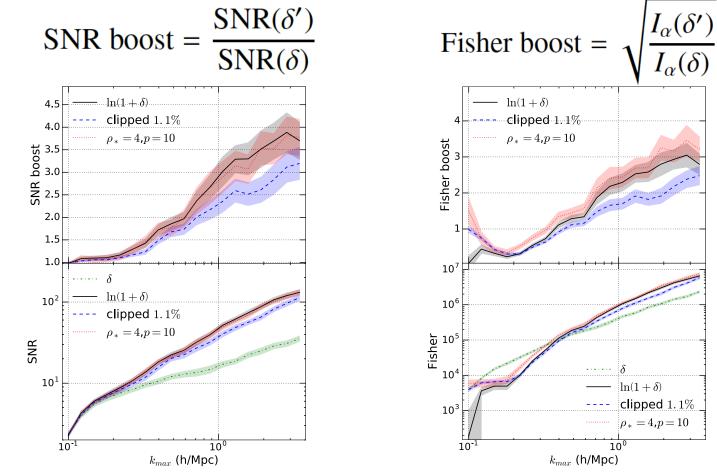
$$I_{\alpha} = \sum_{i,j}^{N_{bins}} \frac{\partial P(k_i)}{\partial \alpha} C_{ij}^{-1} \frac{\partial P(k_j)}{\partial \alpha}$$

• Signal-to-Noise Ratio (SNR)

$$SNR = \sqrt{\sum_{i,j}^{N_{bins}} \bar{P}(k_i) C_{ij}^{-1} \bar{P}(k_j)}$$

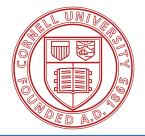


- Assess level of additional information encoded, in terms of "boost"
- Marked transformation increases information relative to standard  $\boldsymbol{\delta}$

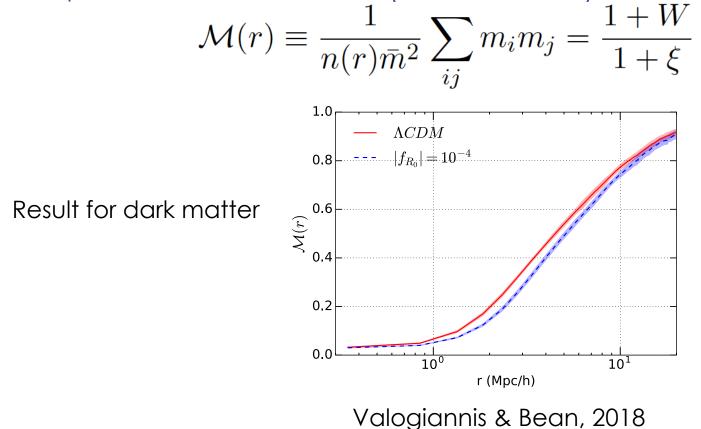


Valogiannis & Bean, 2018

### Marked correlation function in MG



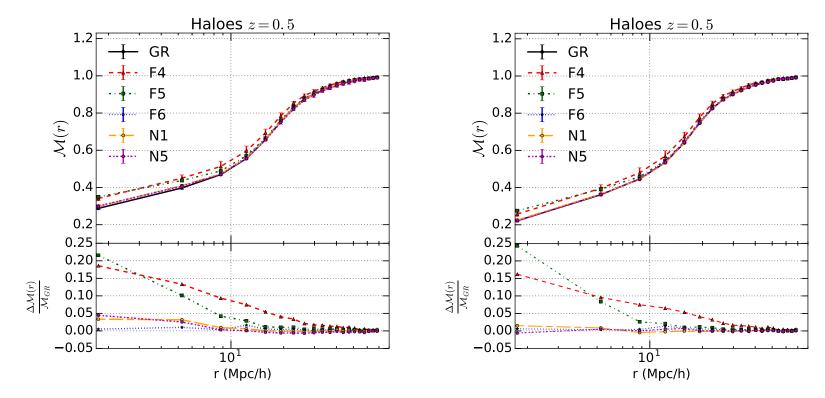
- Screening mechanism produces unique density dependent signature
- What other density-dependent statistics?
- Marked correlation function (Sheth, R.K., Connolly, A.J., & Skibba, R. 2005)
- Real space statistic to test MG (M. White, 2016)



#### Halo-marked correlation function



• Marked correlation functions for dark matter halos at z=0.5

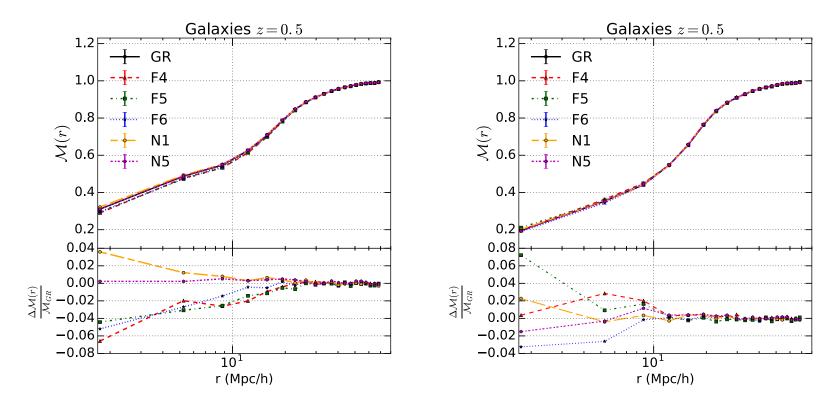


DESI MG white paper, in prep.

#### Galaxy-marked correlation function



• Marked correlation functions for galaxies at z=0.5

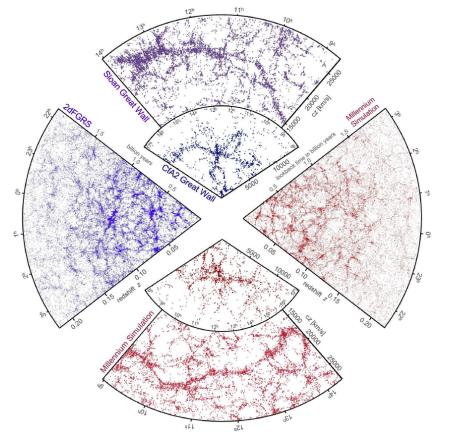


DESI MG white paper, in prep

#### Signatures of viable MG models using biased tracers

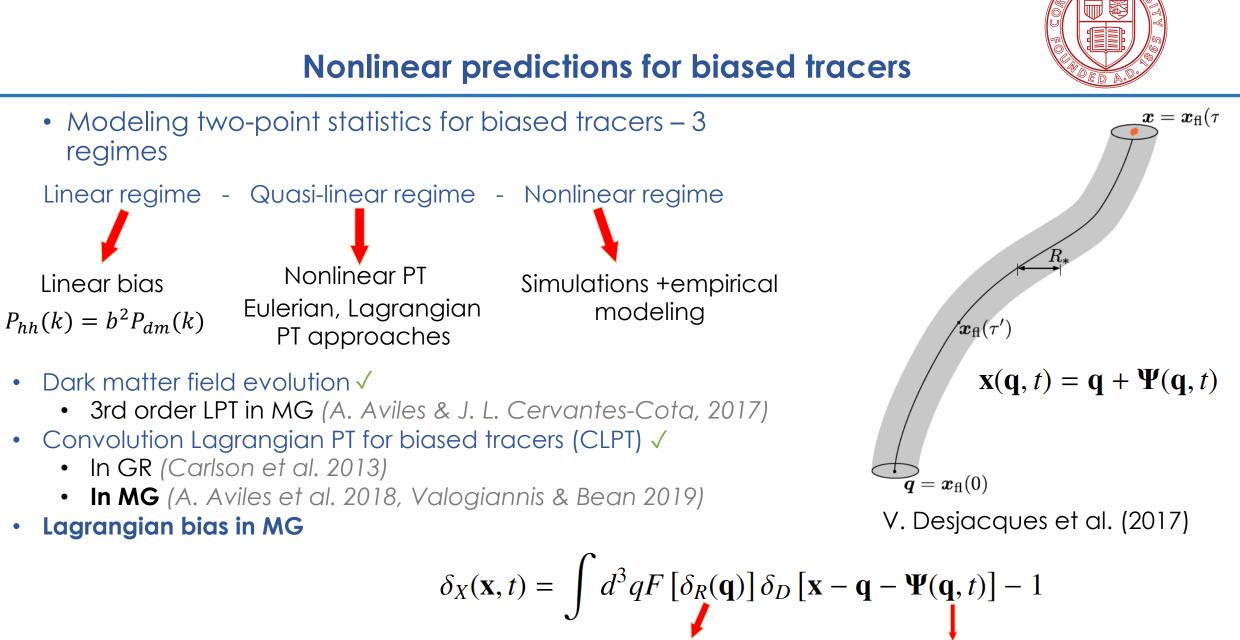
#### Biased tracers

- Spectroscopic surveys observe galaxies that do not trace dark matter field perfectly
  - Biased tracers (N. Kaiser 1987, G. Efstathiou 1988)
- Analytical predictions for 2-point statistics of biased tracers in gravity models (GR, MG) necessary



V. Springel et al. (2006)





Biased tracers

LPT displacement



### Peak-Background Split Lagrangian biases in GR

- Lagrangian bias factors rigorous definition
  - Response to a large wavelength density perturbation  $\boldsymbol{\Delta}$
- Unconditional/universal halo mass function

$$\bar{n}_h(M) = \frac{\partial^2 \bar{N}_h}{\partial V \partial \ln M} = \frac{\bar{\rho}_m}{M} v_c(M) f\left[v_c(M)\right] \frac{d \ln v_c(M)}{dM}$$

F. Schmidt et al. (2013)  

$$b_n^L(M) = \frac{1}{\bar{n}_h(M,0)} \frac{d^n \bar{n}_h(M,\Delta)}{d\Delta^n} \Big|_{\Delta=0}^{\text{Conditiond}}$$
Conditiond

Unconditional halo mass function

$$v_c(M) = \frac{\delta_{cr}}{\sigma(M, z)} = \frac{\delta_{cr}}{D(z)\sigma(M)}$$

- Peak-Background Split (PBS) formalism biases
- Calculation not valid in MG!

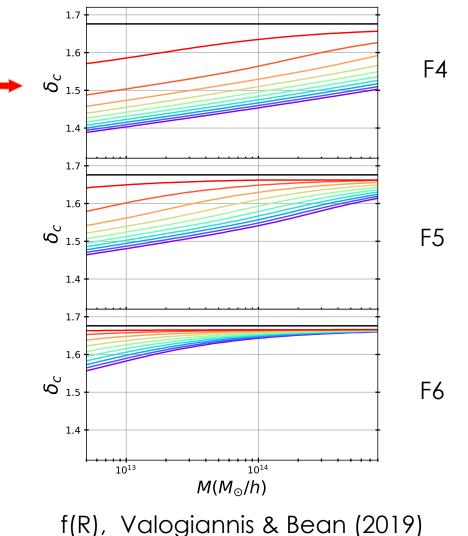
$$\begin{aligned} v_c f\left[v_c\right] &= \sqrt{\frac{2}{\pi}} A(p) \left[ 1 + \frac{1}{(qv_c^2)^p} \right] \sqrt{q} v_c e^{\frac{-qv_c^2}{2}} \\ b_1^L(M) &= \frac{-1}{\delta_{cr}} \left[ qv_c^2 - 1 + \frac{2p}{1 + \left(qv_c^2\right)^p} \right], \\ b_2^L(M) &= \frac{1}{\delta_{cr}^2} \left[ q^2 v_c^4 - 3qv_c^2 + \frac{2p \left(2qv_c^2 + 2p - 1\right)}{1 + \left(qv_c^2\right)^p} \right] \end{aligned}$$

#### Gravitational collapse in MG



[-1(purple), -0.72, -0.43, -0.15, 0.13, 0.42, 0.7, 0.98, 1.27, 1.55(red)]

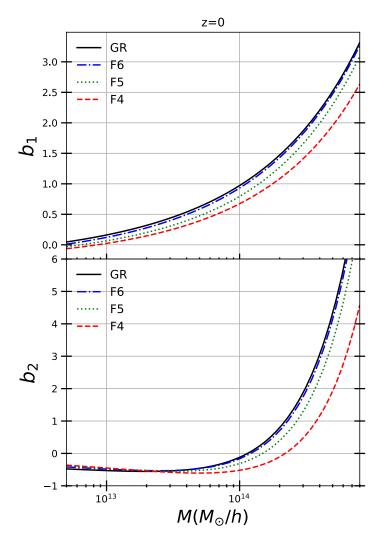
- Density collapse threshold now depends on mass & environment  $\delta_{cr} = \delta_{cr}(M, \delta_{env}, z)$ 
  - B. Li & G. Efstathiou (2012), L. Lombriser et al. (2013)
  - Birkhoff's theorem violated





#### Peak-Background Split Lagrangian biases in MG

- PBS formalism should be modified in MG
- Mass function model in MG  $v_{cMG}(z, M, \delta_{env}) = \frac{\delta_{cr}(z, M, \delta_{env})}{D^{(1)}(z)\sigma(M)}$
- PBS Lagrangian biases in MG
- $b_{1}^{MG} < b_{1}^{GR}$ , in agreement with Arnold et al. (2018)  $b_{MG}^{1}(M, \delta_{env}) = \frac{\frac{d\delta_{cr}(M, \delta_{env})}{d\delta_{env}} - 1}{\delta_{cr}(M, \delta_{env})} \left[ qv_{cMG}^{2} - 1 + \frac{2p}{1 + (qv_{cMG}^{2})^{p}} \right],$  $b_{MG}^{2}(M, \delta_{env}) = \frac{\left(\frac{d\delta_{cr}(M, \delta_{env})}{d\delta_{env}} - 1\right)^{2}}{\delta_{cr}^{2}(M, \delta_{env})} \left[ q^{2}v_{cMG}^{4} - 3qv_{cMG}^{2} + \frac{2p(2qv_{cMG}^{2} + 2p - 1)}{1 + (qv_{cMG}^{2})^{p}} \right] + \frac{d^{2}\delta_{cr}(M, \delta_{env})}{d\delta_{env}^{2}} \frac{1}{\delta_{cr}(M, \delta_{env})} \left[ qv_{cMG}^{2} - 1 + \frac{2p}{1 + (qv_{cMG}^{2})^{p}} \right].$  Valogiannis & Bean (2019)



f(R), Valogiannis & Bean (2019)

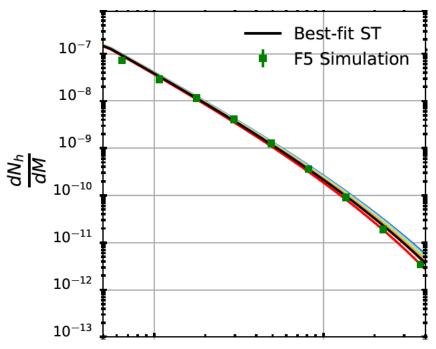
### Halo mass function fitting in MG



- In GR, standard fit (q,p)=(0.75,0.3)
  - (R. Sheth et al. 2001)
- In MG, ST values different absorb novel physics of gravitational collapse
- Best-fit (q,p) against MG simulations

	Best-fit ST		Predicted Biases	
Models	q	р	$b_1$	$b_2$
Group I : GR	0.726	0.345	0.301	-0.501
Group I: F4	0.671	0.361	0.120	-0.435
Group I: F5	0.765	0.321	0.211	-0.470
Group I: F6	0.670	0.362	0.230	-0.449
Group I: N1	0.701	0.369	0.224	-0.661
Group I : N5	0.702	0.357	0.268	-0.503
Group II : GR Low	0.674	0.362	0.345	-0.183
Group II : GR Mid.	0.728	0.342	0.925	-0.05
Group II : GR High	0.806	0.594	1.720	1.900
Group II : F5 Low	0.733	0.314	0.295	-0.170
Group II : F5 Mid.	0.788	0.282	0.909	-0.033
Group II : F5 High	0.746	0.305	1.491	0.416





Valogiannis & Bean (2019)

### 2-point correlation function results



Low mass bin

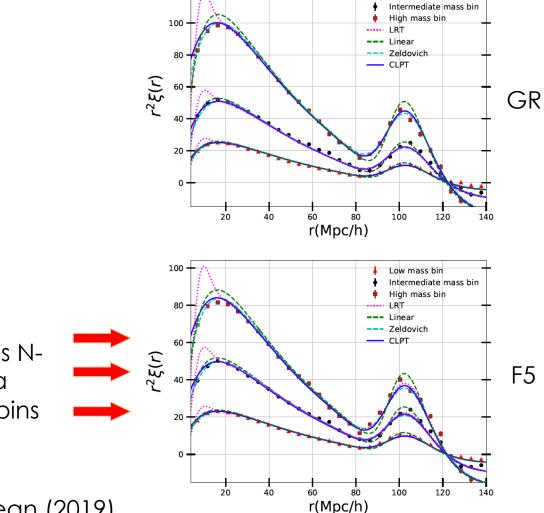


- f(R) (chameleon) Hu-Sawicki, 2007
- nDGP (Vainshtein) G. Dvali et al. 2000

#### N-body simulations:

- f(R) Lightcone project:
  - z=1, L=1536 Mpc/h
  - 3 mass bins
  - C. Arnold et al. (2018)
- f(R) & nDGP ELEPHANT sims:
  - z=0.5, L=1024 Mpc/h
  - 1 mass bin
  - M. Cautun et al. (2017)

CLPT matches Nbody data across mass bins

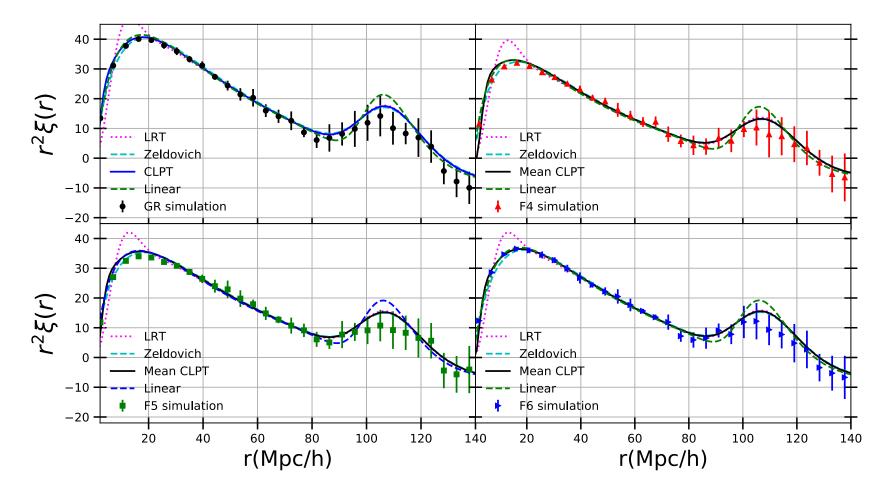


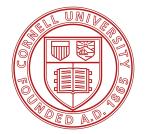
120

### 2-point correlation function results

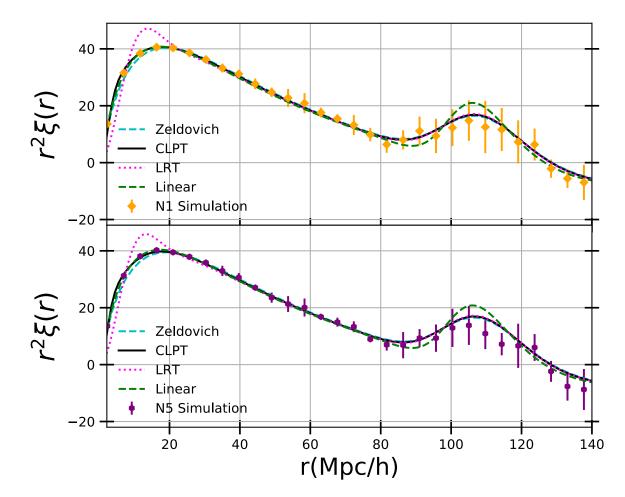


• Consistency for varying degrees of f(R) screening z=0.5 !





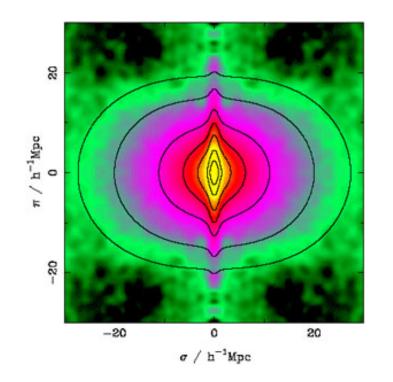
• Consistency for varying degrees of Vainshtein screening at z=0.5 !



## Testing MG models in redshift-space

#### Redshift space distortions

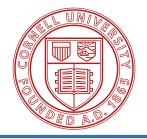
- Spectroscopic surveys observe galaxies:
  - **biased tracers** (galaxy clustering in different environments) = Lagrangian bias in MG (1901.03763)
  - **Redshift Space Distortions =** Gaussian streaming model (GSM) (paper in prep)





Peacock et al. 2001

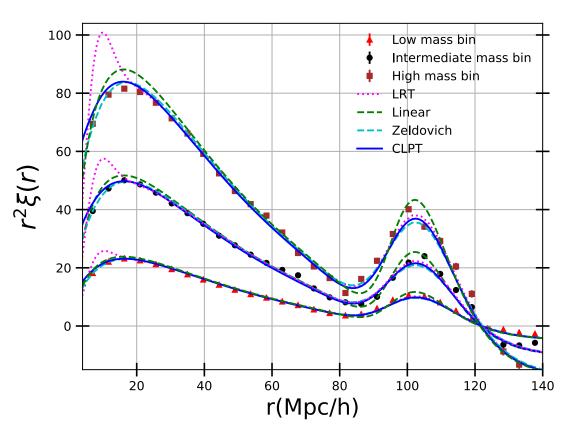
# Necessary ingredients to model RSD in MG



- Nonlinear evolution of DM field in MG
  - Lagrangian perturbation theory (LPT)  $\checkmark$ 
    - (A. Aviles & J. L. Cervantes-Cota, 2017)
- Halo bias in MG
  - Convolution LPT for biased tracers in MG  $\checkmark$

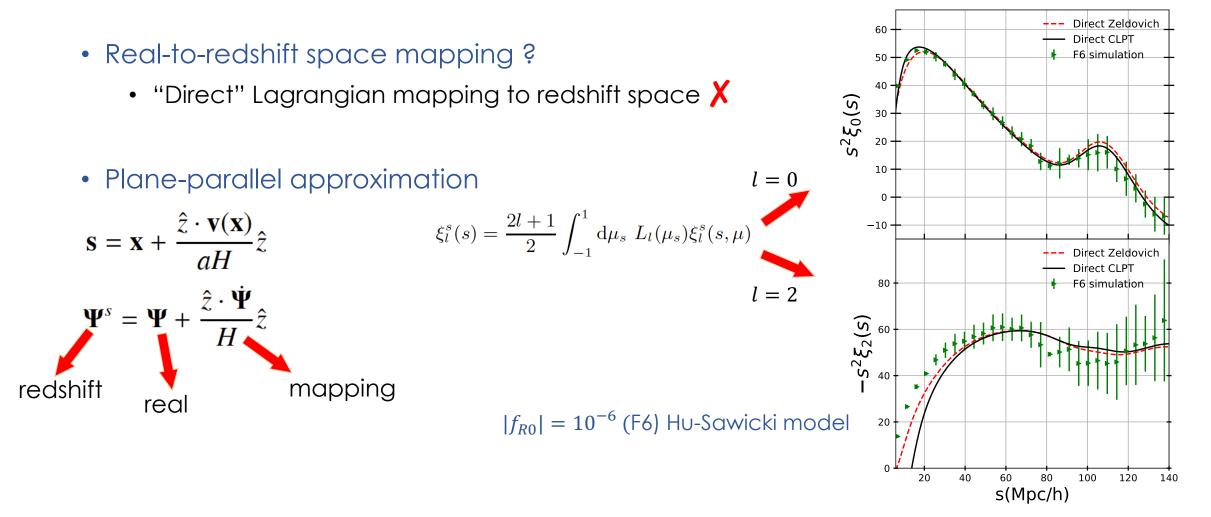
(A. Aviles et al. 2018, Valogiannis & Bean 2019)

 Sheth–Tormen (ST) halo bias in MG √ (G. Valogiannis & R. Bean 2019)





## Necessary ingredients to model RSD in MG



Valogiannis et al. (in prep.)

# Necessary ingredients to model RSD in MG



- Real-to-redshift space mapping
  - Scale-dependent Gaussian Streaming model (GSM)  $\checkmark$

(B. Reid & M. White, 2011, L. Wang et al. 2013)

$$1 + \xi^{s}(s_{\perp}, s_{\parallel}) = \int \frac{dy}{[2\pi]^{1/2}\sigma_{12}} [1 + \xi(r)] \exp\left\{-\frac{[s_{\parallel} - y - \mu v_{12}]^{2}}{2\sigma_{12}^{2}}\right\}$$
  
Redshift  
space

# **GSM ingredients in MG**



- Real-to-redshift space mapping
  - Scale-dependent Gaussian Streaming model (GSM)  $\checkmark$

(B. Reid & M. White, 2011, L. Wang et al. 2013)

$$1 + \xi^{s}(s_{\perp}, s_{\parallel}) = \int \frac{dy}{[2\pi]^{1/2} \sigma_{12}} [1 + \xi(r)] \exp\left\{-\frac{[s_{\parallel} - y - uv_{12}]^{2}}{2\sigma_{12}^{2}}\right\}$$

- GSM ingredients
  - Real-space  $\xi(r)$  for biased tracers in MG  $\checkmark$
  - Pairwise velocity  $v_{12}(r)$  for biased tracers in MG

$$v_{12}(r)\mathbf{\hat{r}} = \frac{\langle [1+\delta(\mathbf{x})][1+\delta(\mathbf{x}+\mathbf{r})][\mathbf{v}(\mathbf{x}+\mathbf{r})-\mathbf{v}(\mathbf{x})] \rangle}{\langle [1+\delta(\mathbf{x})][1+\delta(\mathbf{x}+\mathbf{r})] \rangle}$$

• Dispersion 
$$\sigma_{12}^2(r)$$
 for biased tracers in MG

- Model  $v_{12}(r)$  &  $\sigma_{12}^2(r)$  using LPT
- $\dot{\Psi} \neq f H \Psi$  in MG scale-dependence!!

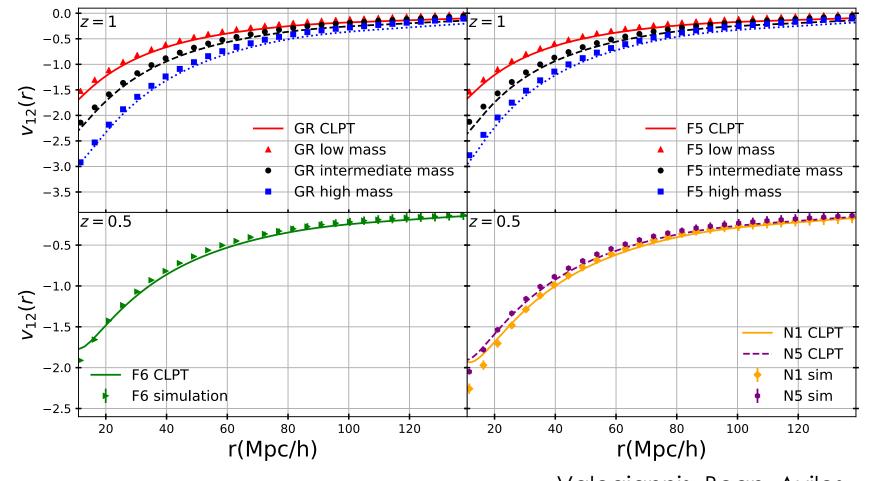
$$\sigma_{12}^2(r,\mu) = \frac{\left\langle [1+\delta(\mathbf{x})][1+\delta(\mathbf{x}+\mathbf{r})][v_z(\mathbf{x}+\mathbf{r})-v_z(\mathbf{x})]^2 \right\rangle}{\left\langle [1+\delta(\mathbf{x})][1+\delta(\mathbf{x}+\mathbf{r})] \right\rangle}$$



# Pairwise velocity (v) model matches sims well

Theory:

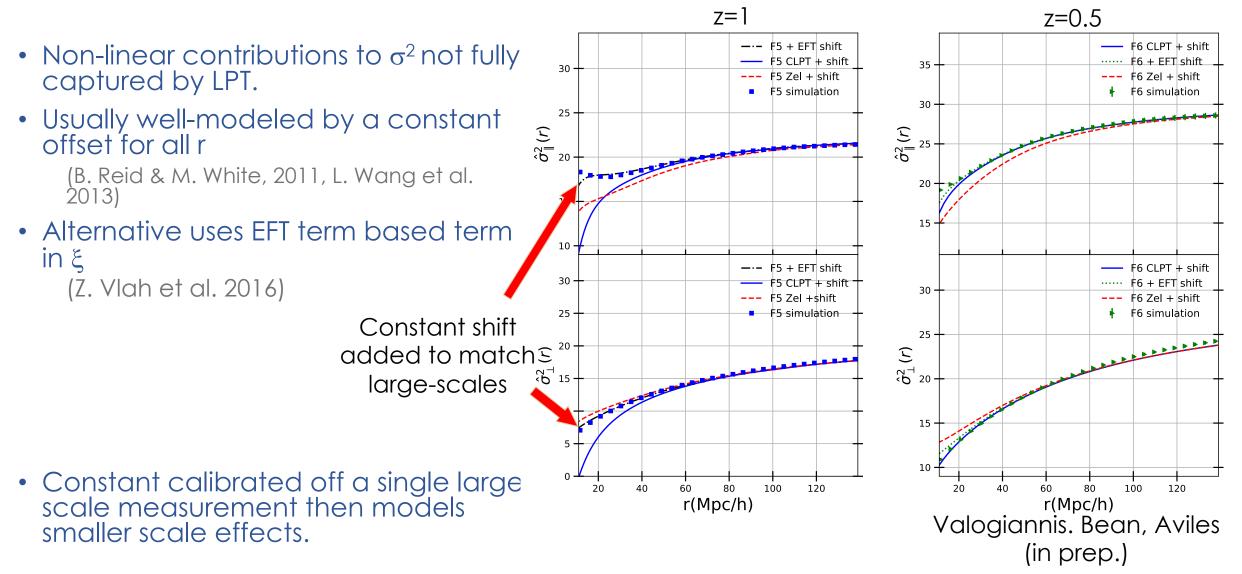
- f(R) (chameleon) Hu-Sawicki, 2007
- nDGP (Vainshtein) G. Dvali et al. 2000
- N-body simulations:
- f(R) Lightcone project:
  - z=1, L=1536 Mpc/h
  - 3 mass bins
  - C. Arnold et al. (2018)
- f(R) & nDGP ELEPHANT sims:
  - z=0.5, L=1024 Mpc/h
  - 1 mass bin
  - M. Cautun et al. (2017)



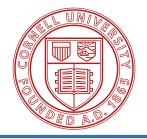
Valogiannis. Bean, Aviles (in prep.)

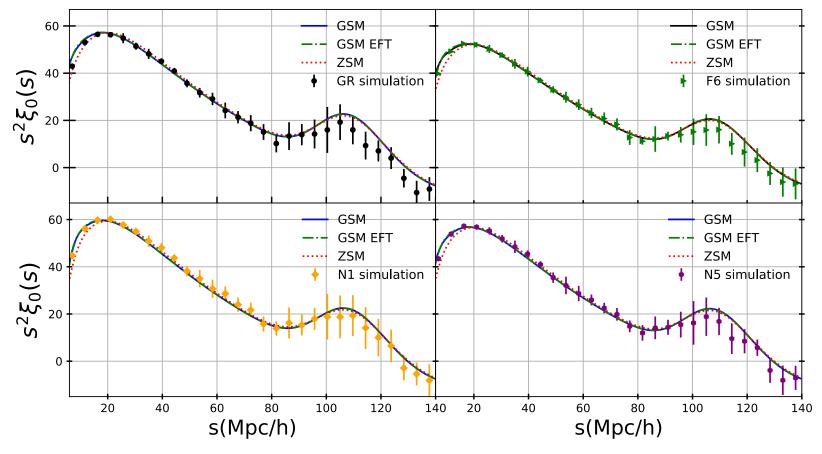


# Modeling velocity dispersion ( $\sigma^2$ )+ shift



### Accurate monopole predictions across MG models

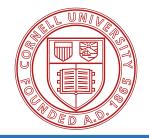


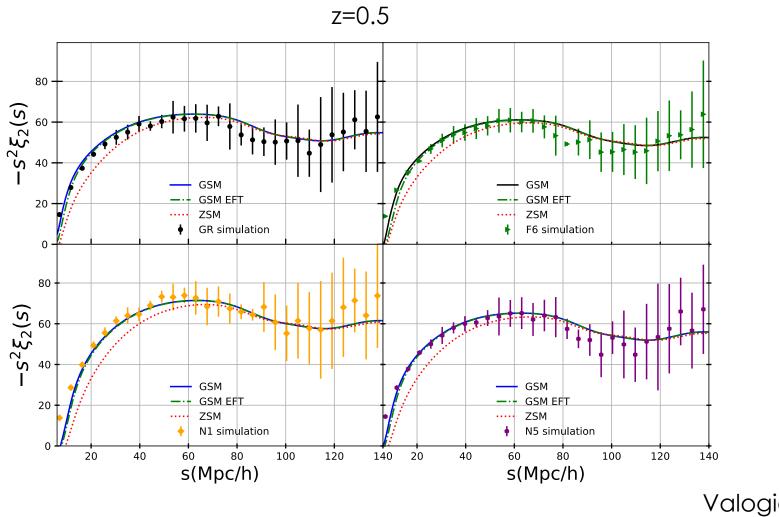


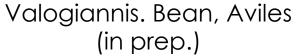
z=0.5

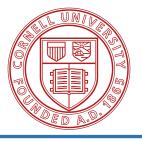
Valogiannis. Bean, Aviles (in prep.)

## Accurate quadrupole predictions across MG models

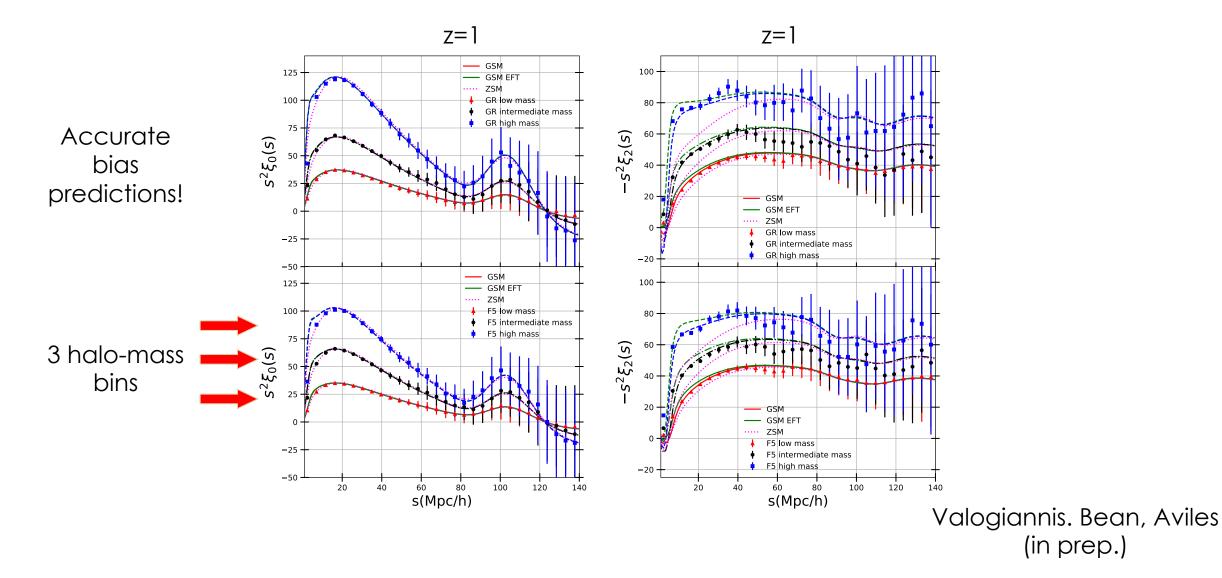








### Accurate predictions across different mass bins



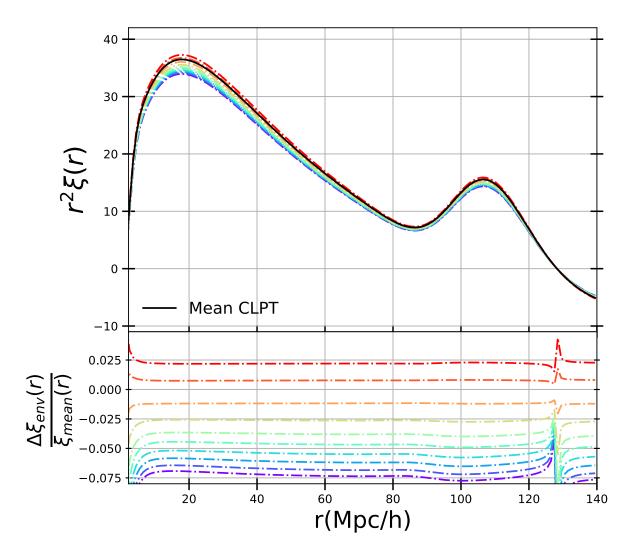
## Summary



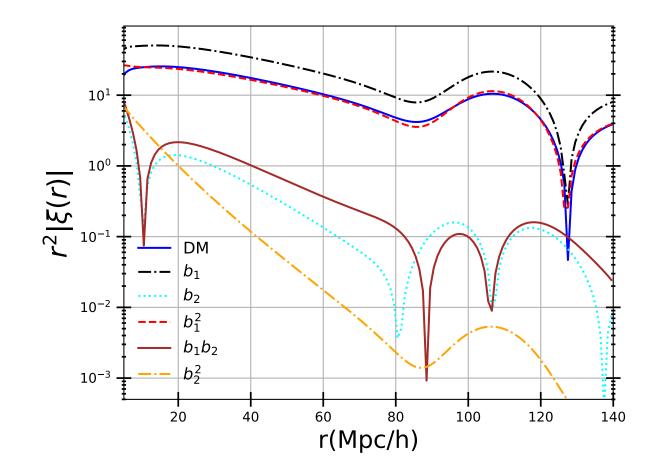
- Hybrid COLA scheme for efficient MG simulations
  - Valogiannis & Bean, 2017, arXiv:1612.06469, Phys. Rev. D 95, 103515
  - Ideal for emulators (in progress)
- Novel statistics to detect MG signals
  - Valogiannis & Bean, 2018 , arXiv:1708.05652, Phys. Rev. D 97, 023535
  - Testing the theory of gravity with DESI: estimators, predictions and simulation requirements (in progress)
- Modeling 2-point statistics for biased tracers in MG
  - Valogiannis & Bean, 2019, arXiv:1901.03763, Phys. Rev. D 99, 063526
- In redshift space: Gaussian Streaming Model in MG
  - Prepared for submission
  - Bias & RSD **simultaneously** modeled in MG for the first time!



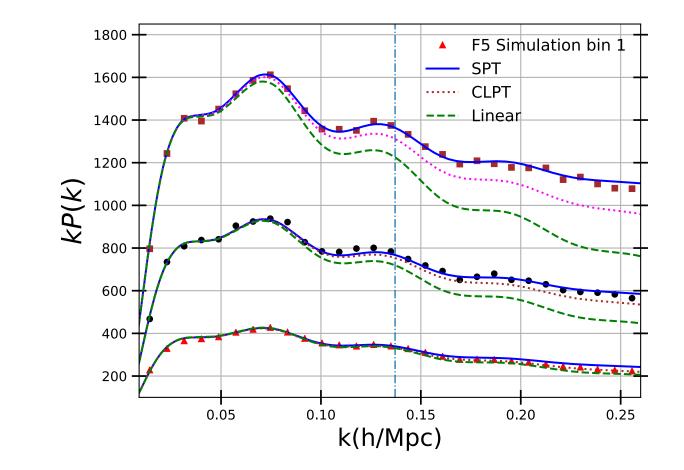








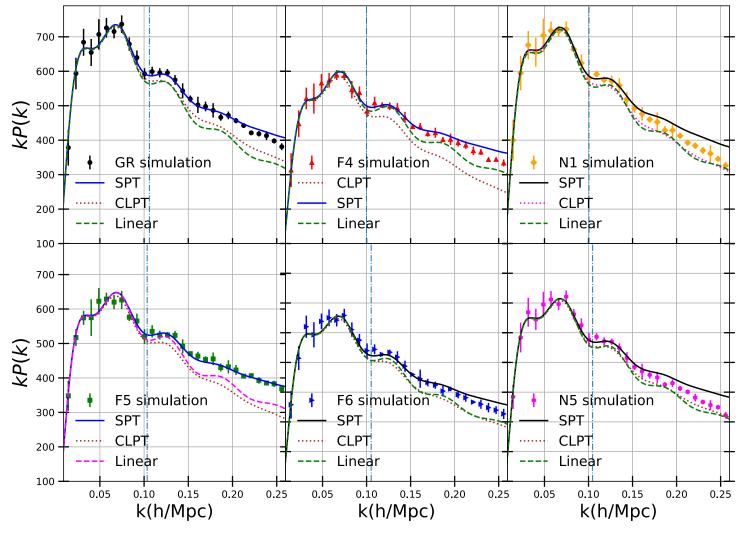




#### **Power spectrum results**



• Eulerian SPT expansion accurate



Valogiannis & Bean (2019)

# Gravitational collapse in MG



F4

F5

F6

[-1(purple), -0.72, -0.43, -0.15, 0.13, 0.42, 0.7, 0.98, 1.27, 1.55(red)]

1.7

 Density collapse threshold now depends on mass & environment  $\delta_{cr} = \delta_{cr}(M, \delta_{env}, z)$ 1.6  $\delta_c$ B. Li & G. Efstathiou (2012), L. Lombriser et al. (2013) 1.5 Birkhoff's theorem violated 1.4 • Collapsing halo-overdensity  $1 + \delta_h = y_h^{-3}$ 1.7 • In MG  $y_h'' - \left(2 - \frac{3}{2}\Omega(a)\right)y_h' + \frac{1}{2}\Omega(a)\frac{G_{eff}}{G}\left(y_h^{-3} - 1\right)y_h = 0$ 1.6  $\overset{\circ}{\wp}_{1.5}$  $G_{eff} = (1+E)G$ 1.4 1.7 1.6 Chameleons  $\overset{\circ}{\wp}_{1.5}$  $E = 2\beta^2 \left[ 3\frac{\Delta R}{R_{th}} - 3\left(\frac{\Delta R}{R_{th}}\right)^2 + \left(\frac{\Delta R}{R_{th}}\right)^3 \right]$ 1.4 1013 1014  $M(M_{\odot}/h)$ "Thin shell" f(R), Valogiannis & Bean (2019)



$$\frac{\Delta R}{R_{th}} = \frac{|f_{R0}|a^3}{\Omega_m y_h^{-3} H_0^2 R_{th}^2}$$

$$\times \left[ \left( \frac{1 + 4\frac{\Omega_\Lambda}{\Omega_m}}{(y_{env}a)^{-3} + 4\frac{\Omega_\Lambda}{\Omega_m}} \right)^{n+1} - \left( \frac{1 + 4\frac{\Omega_\Lambda}{\Omega_m}}{(y_ha)^{-3} + 4\frac{\Omega_\Lambda}{\Omega_m}} \right)^{n+1} \right]$$

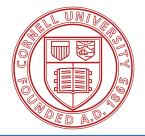
$$E = \frac{2}{3\beta(a)} \frac{\sqrt{1+\chi^3}-1}{\chi^{-3}} \qquad \qquad \chi^{-3} = \frac{\Omega_m n^2}{1.10894a^3\beta^2(a)} \frac{y_h^3-1}{y_h^3}$$

# **DESI-related work**

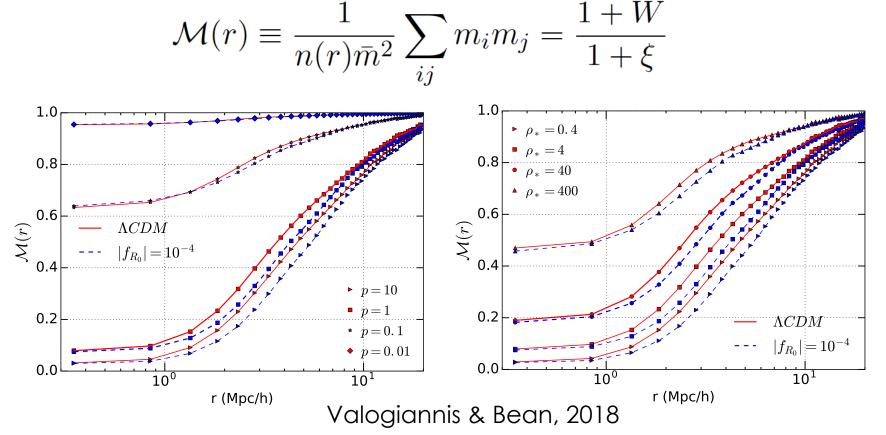


- Dark matter-only simulations shed light on underlying gravitational physics
- Realistic surveys do **not** trace dark matter
  - Dark matter halos
  - Biased tracers
- Test marked statistics on realistic applications DESI MG white paper (in prep)
- ELEPHANT simulations (M. Cautun et al., 2017): Simulation box side L=1024 Mpc/h, 1024<sup>3</sup> particles
- ACDM, f(R) and nDGP MG models
- Vainshtein Mechanism in nDGP (G. Dvali et al., 2000)
- Dark matter halo-finding
- Biased tracers (galaxies) Halo Occupation Distribution (HOD)

## Marked correlation function



- Screening mechanism produces unique density dependent signature
- What other density-dependent statistics?
- Marked correlation function (Sheth, R.K., Connolly, A.J., & Skibba, R. 2005)
- Real space statistic





- Hybrid COLA scheme for efficient MG chameleon simulations
  - Phys. Rev. D 95, 103515 (2017)
- Simple, "marked" density transformations serve as powerful tools for testing gravity
- Up-weight unscreened regions and down-weight high densities
  - Phys. Rev. D 97, 023535 (2018)
- "Testing the theory of gravity with DESI: estimators, predictions and simulation requirements" white paper in preparation
- C. H. Aguayo et al. (2018) & J. Armijo et al. (2018) on marks
- Further explore use of mark in the context of realistic observations
- Perturbation theory predictions for marked statistics

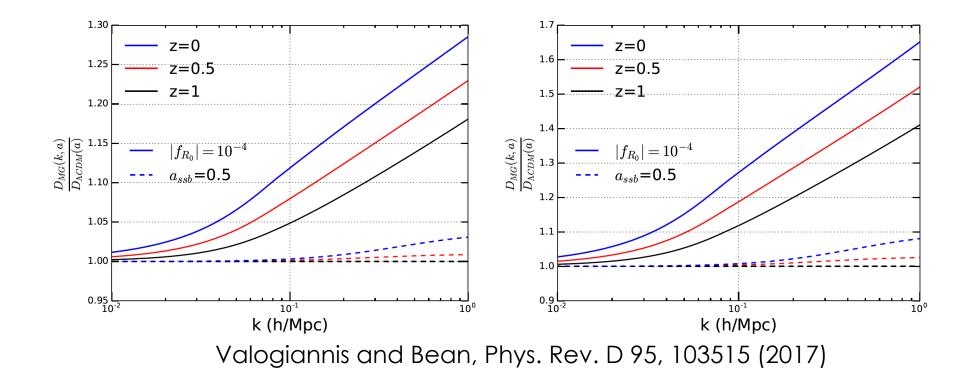


- Simple, "marked" density transformations serve as powerful tools for testing gravity
- Up-weight unscreened regions and down-weight high densities
- Enhance information encoded in 2-point statistics
- Marked correlation function for differentiating between MG and GR
- Phys. Rev. D 97, 023535 (2018)
- "Testing the theory of gravity with DESI: estimators, predictions and simulation requirements" white paper in preparation
- C. Hernandez Aguayo et al. (2018) & J. Armijo et al. (2018) on marks
- Further explore use of mark in the context of realistic observations
- Perturbation theory predictions for marked statistics

# **COLA simulations for MG: LPT component**



• LPT growth factor becomes scale dependent



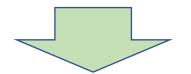


- In lower-density, unscreened regions, modifications significant
- In MG, dark matter particles in voids get pulled outwards faster
- Voids emptier in MG for **dark matter**  $\rightarrow M_{GR} > M_{MG}$
- However, due to stronger gravity, **more** halos/galaxies form in voids
- As a result, for **halo/galaxy** marks  $\rightarrow M_{GR} < M_{MG}$
- HOD parameters tuned to match observations deviations vanish
- Marks encoding additional information e.g. gravitational potential

$$\mathcal{M}(r) \equiv \frac{1}{n(r)\bar{m}^2} \sum_{ij} m_i m_j = \frac{1+W}{1+\xi}$$

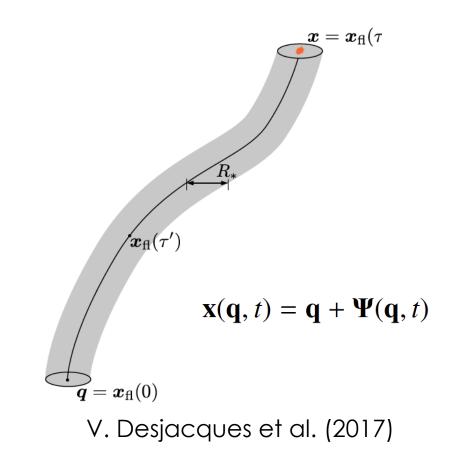
# Necessary ingredients to model RSD in MG

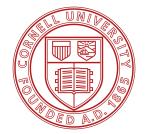
- Nonlinear evolution of Dark Matter field in MG
- Halo bias modeling in MG
- Real-to-redshift space mapping in MG



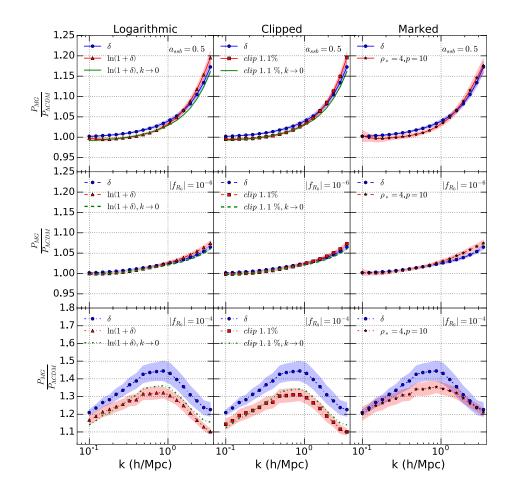
Need to include MG complications

- Modified, scale dependent growth rates
- Environment (density) dependent clustering
- Screening effects





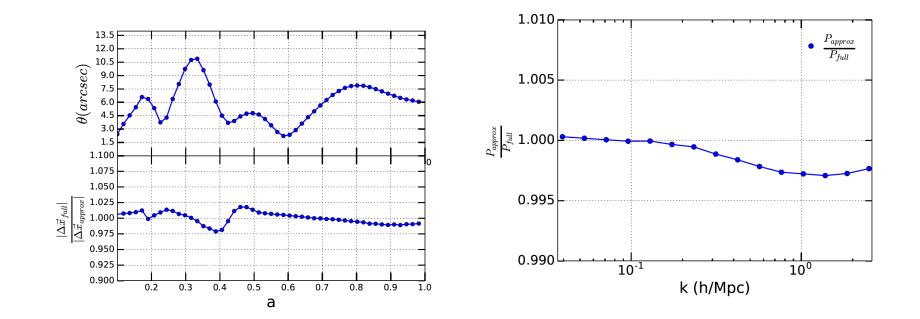
• MG deviations in power spectra enhanced in non-linear regime



Valogiannis & Bean, 2018



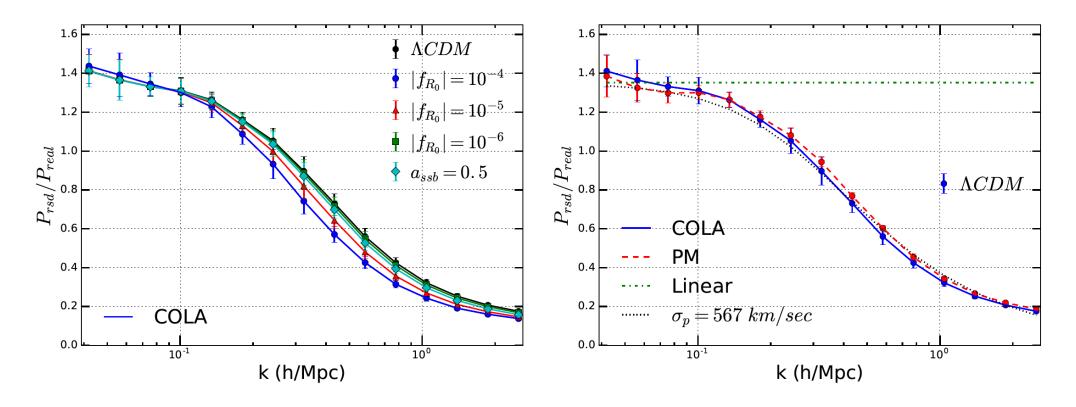
#### Solution: Scale dependent LPT growth negligible



Valogiannis and Bean, Phys. Rev. D 95, 103515 (2017)



#### **Redshift Space Distortions**



Valogiannis and Bean, Phys. Rev. D 95, 103515 (2017)

# UNIDA S

# Chameleon mechanism generates density-dependent mass

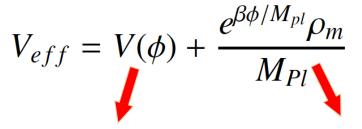
• Scalar-tensor chameleon

$$S = \int dx^4 \sqrt{-g} \left[ \frac{R}{2} M_{\rm Pl}^2 - \frac{1}{2} (\partial \phi)^2 - V(\phi) + S_m(g_{\mu\nu} A^2(\phi), \psi_m) \right]$$

- Klein-Gordon equation for  $\boldsymbol{\phi}$ 

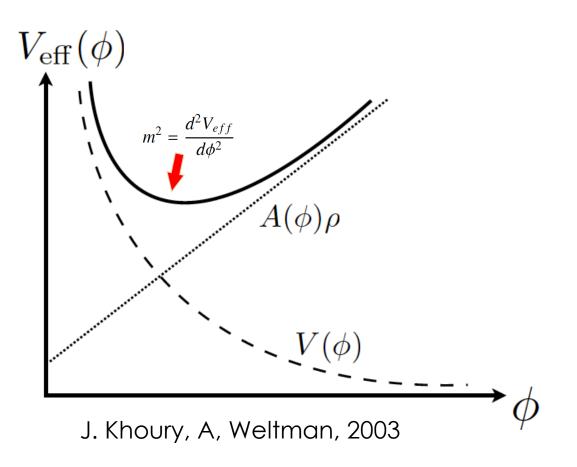
 $\Box \phi = V_{eff,\phi}$ 

• Effective potential



runaway potential

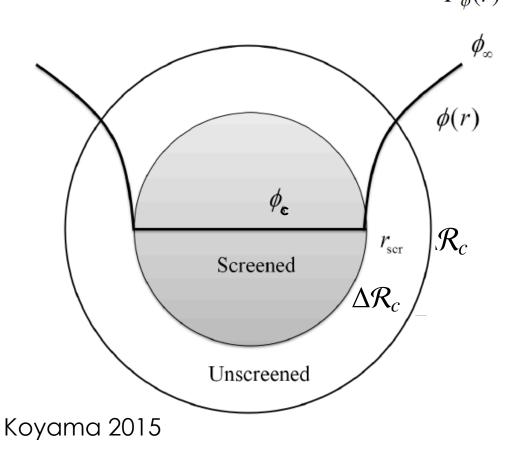
Coupling to matter



### Fifth force phenomenology - thin shell effect

 $\frac{\Delta \mathcal{R}_c}{\mathcal{R}_c} > \bar{1}$ 

- In linear regime, no screening
- Fifth force in full effect



 $F_{\phi}(r) \approx 2\beta_{\infty}^2 \frac{GM}{r^2}$  J. Khoury, A, Weltman, 2003

- In dense regime interior mass decouples due to chameleon (Yukawa suppression)
- Effective coupled mass confined to a "thin shell"

• Fifth force screened  $\frac{\Delta \mathcal{R}_c}{\mathcal{R}_c} = \frac{|\phi_{\infty} - \phi_c|}{2\beta_{\infty}M_{Pl}\Phi_N} \ll 1$ 

$$F_{\phi}(r) \approx 2\beta_{\infty}^2 \left(\frac{\Delta \mathcal{R}_c}{\mathcal{R}_c}\right) \frac{GM}{r^2}$$
  
Coupling "Thin Shell" Newtor

