

# The splashback boundary of galaxy clusters in mass and light and its implications for galaxy evolution



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(DES collaboration) (SPT collaboration) (ACT collaboration)

<https://arxiv.org/abs/1811.06081> (accepted to MNRAS); 2019 paper in preparation

Based on ~400 public cluster sample from ACT+SPT

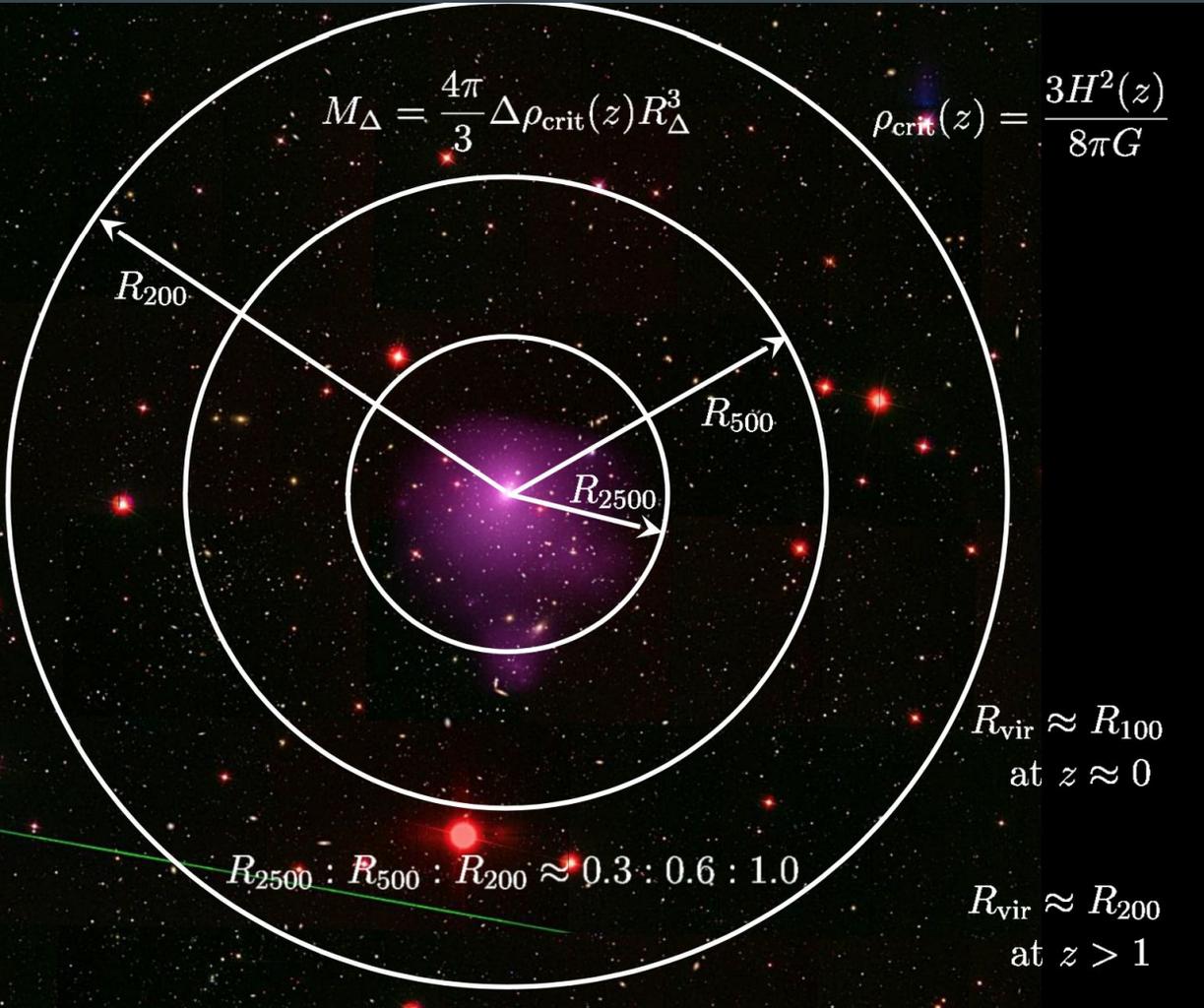
Ongoing analysis of 1000+ SZ-selected clusters from DES+ACT+SPT

# Background

# Mass and boundary of dark matter halos

$$M_{\Delta} = \frac{4\pi}{3} \Delta \rho_{\text{crit}}(z) R_{\Delta}^3$$

$$\rho_{\text{crit}}(z) = \frac{3H^2(z)}{8\pi G}$$



$$R_{2500} : R_{500} : R_{200} \approx 0.3 : 0.6 : 1.0$$

$$R_{\text{vir}} \approx R_{100} \\ \text{at } z \approx 0$$

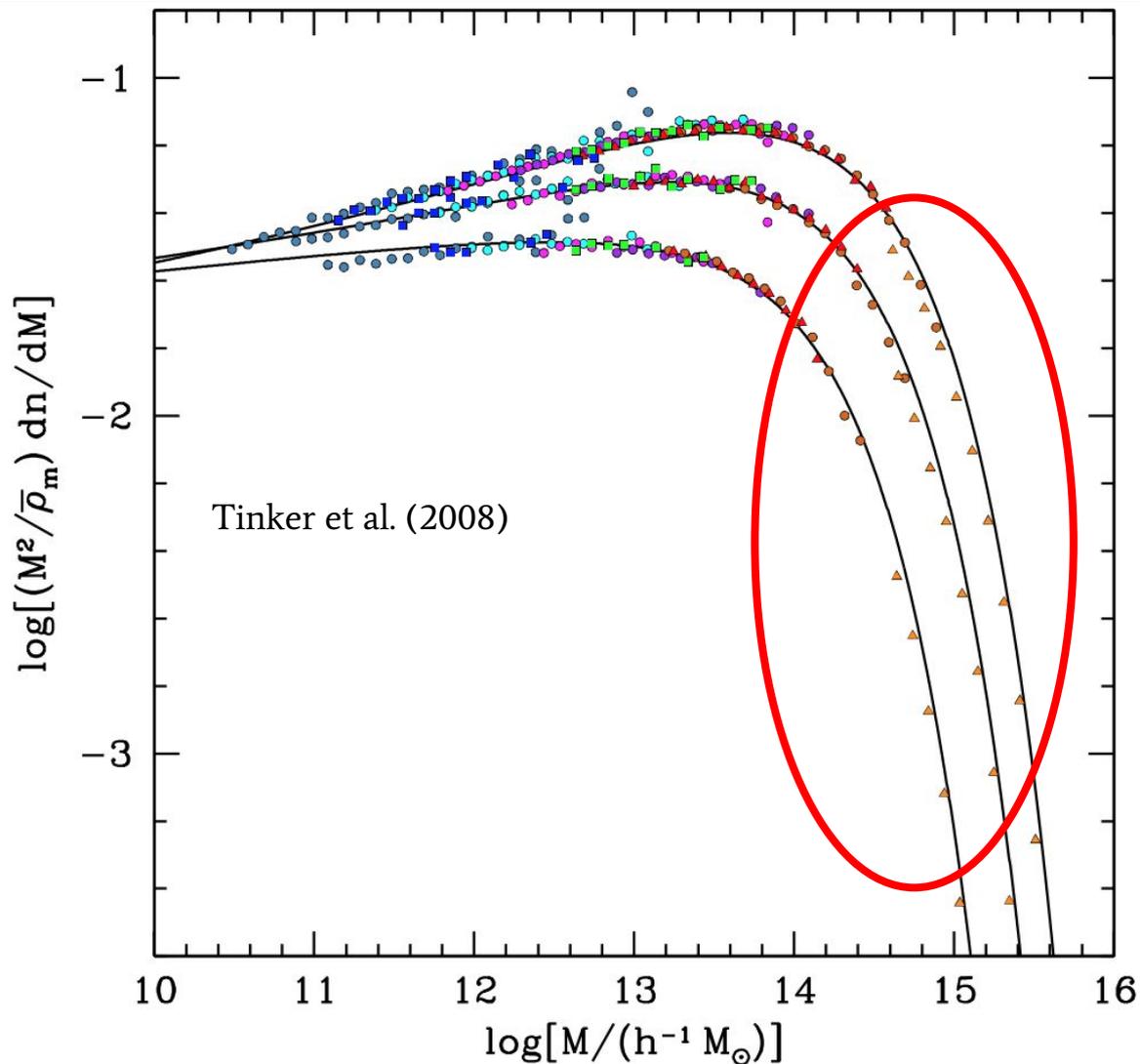
$$R_{\text{vir}} \approx R_{200} \\ \text{at } z > 1$$

However,  $M_{\Delta}$  and  $R_{\Delta}$  are subject to **pseudo-evolution** due to the decrease in the reference density ( $\rho_c$  or  $\rho_m$ )

Halos continuously accrete matter; there is no radius within which the matter is fully virialized

$\Rightarrow$  where is the **physical** boundary of the halos?

# Cosmology with galaxy clusters

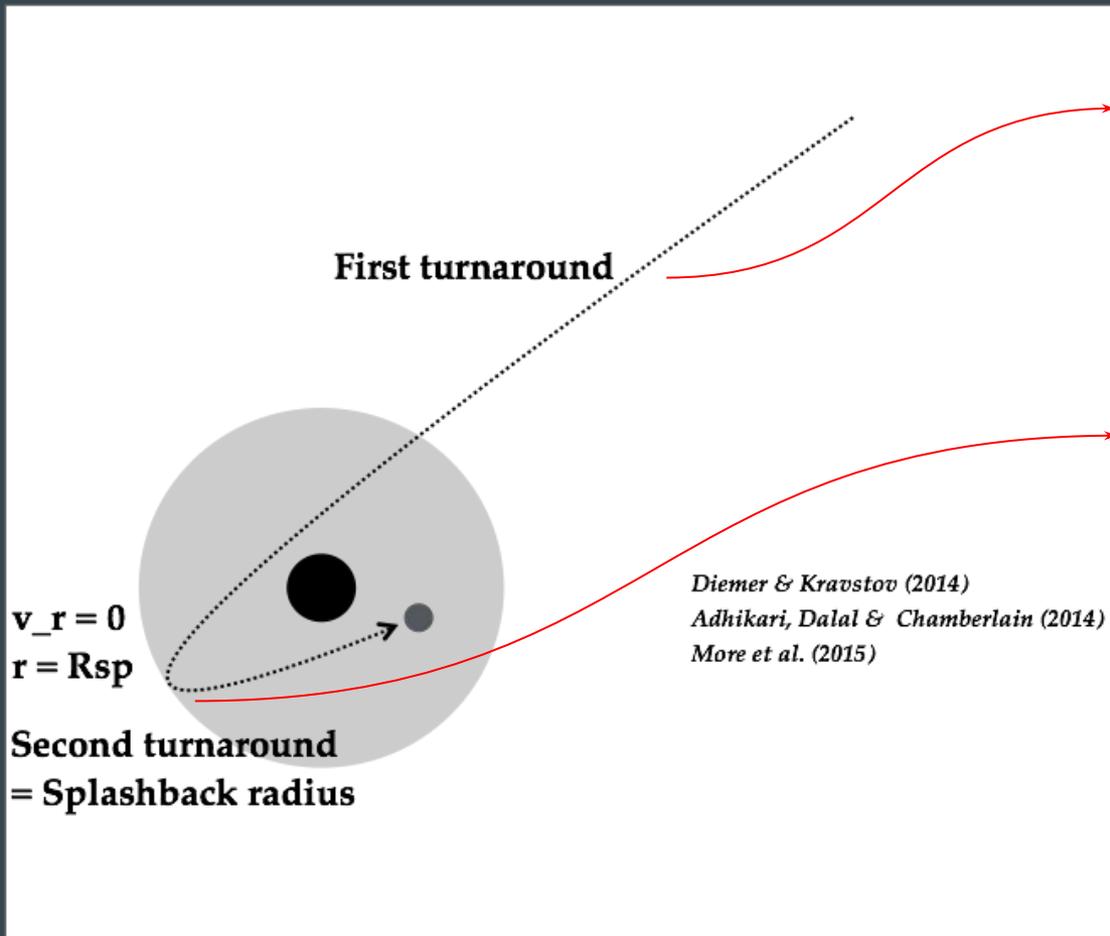


Galaxy clusters live in the **high-mass tail** of the halo mass function  
 $\Rightarrow$  very sensitive to the growth of the structure ( $\Omega_m$  and  $\sigma_8$ )

Thus, it is important to accurately define/measure the mass of the cluster

Preliminary work by Diemer et al. illuminates that **the mass function becomes more universal** against redshift when we use so-called “splashback radius” as the physical boundary of the dark matter halos

# Background

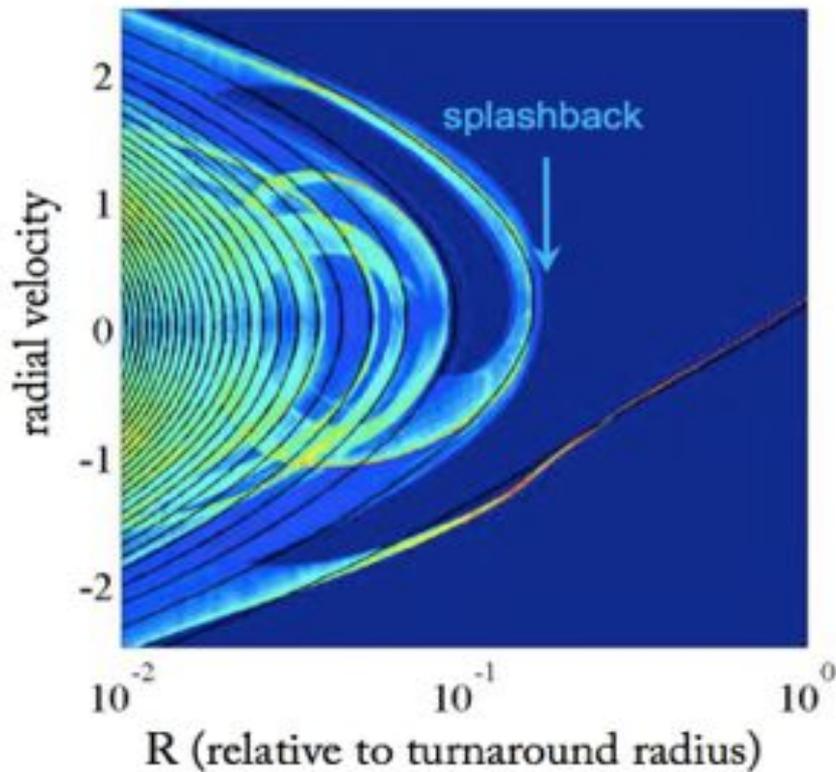


- Galaxies fall into the cluster potential, escaping from the Hubble flow
- They form a sharp “physical” boundary around their first apocenters after the infall, which we call “splashback radius”

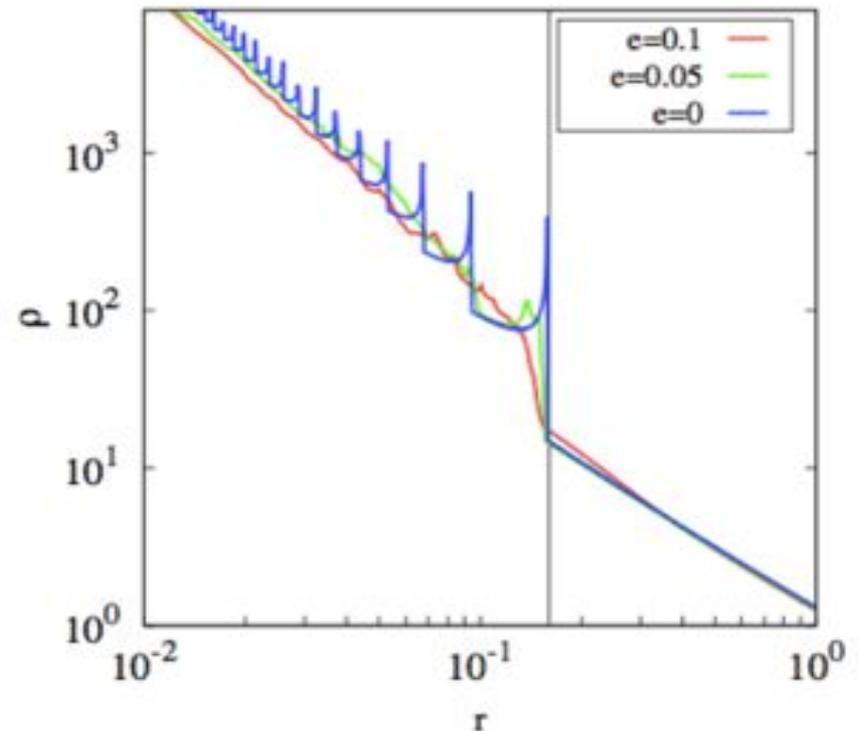
Fig. from Chihway Chang (UChicago)

# Background

- A simple spherical collapse model can predict the existence of the splashback feature (Gunn & Gott 1972, Fillmore & Goldreich 1984, Bertschinger 1985, Adhikari et al. 2014)



Adhikari et al. (2014)



# Background

- Galaxy clusters exhibit a **sharp decline in density profile** around the first orbital apocenters of accreting particles
- Splashback radius,  $r_{sp}$ , represents the location of **the steepest logarithmic slope** and it mostly depends on **accretion rate** of the matter into the clusters, and mass of the clusters

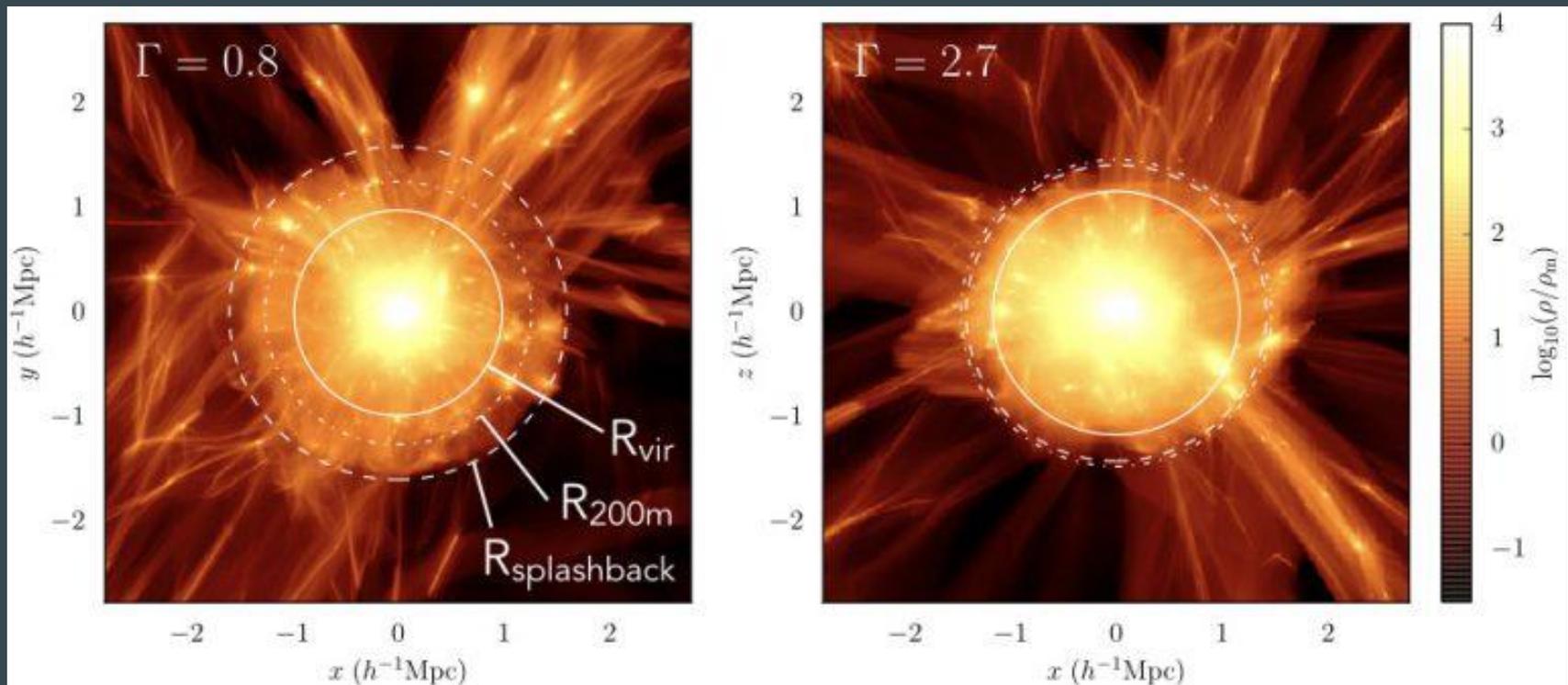
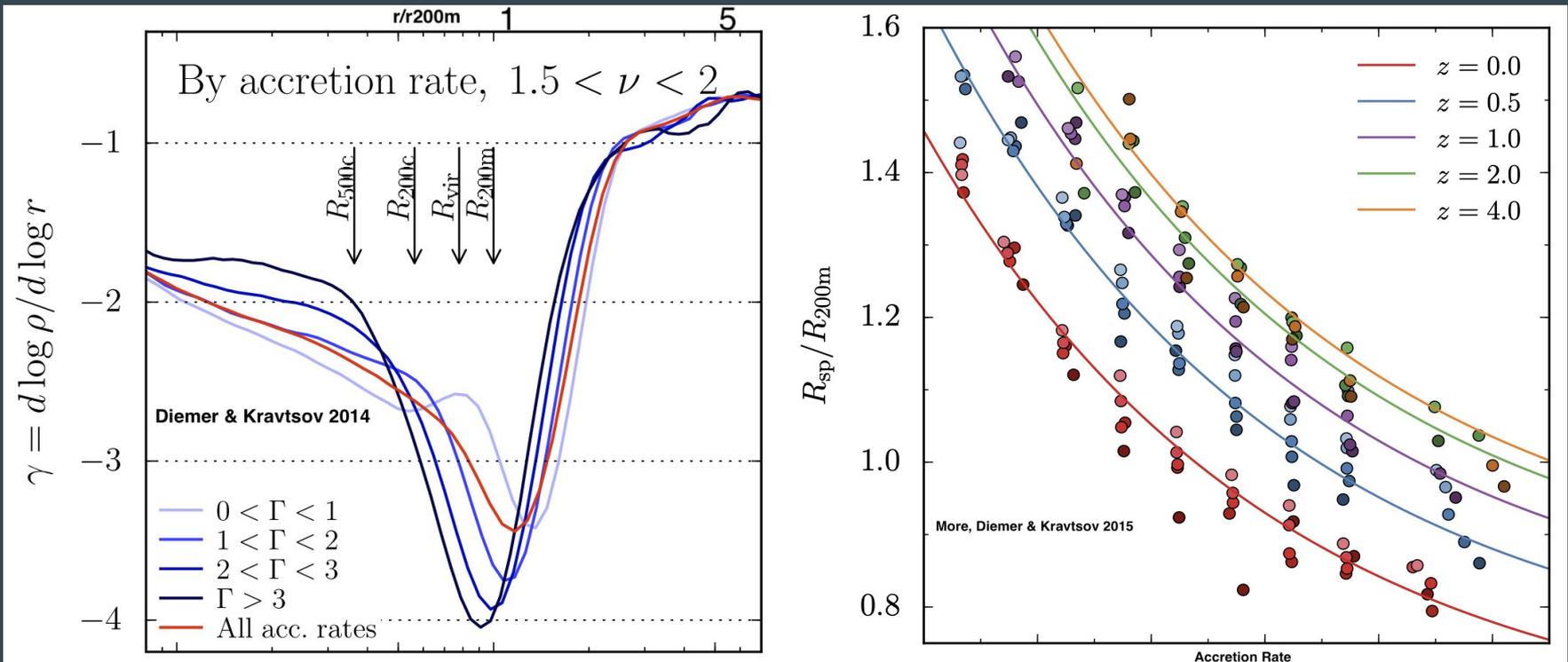


Fig. from More et al. 2015

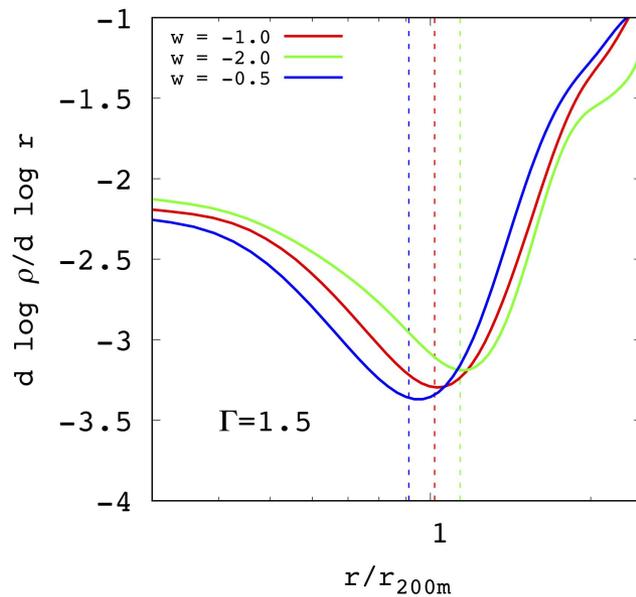
# Background

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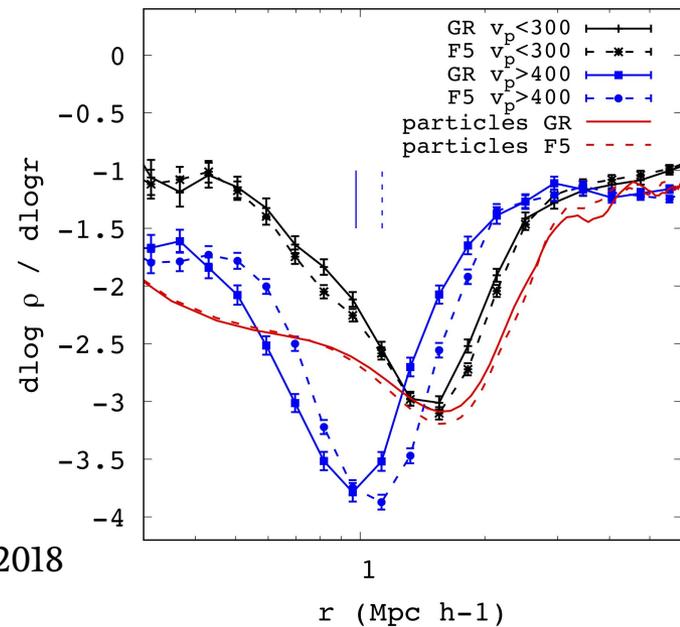


# Background

- Furthermore, splashback feature also depends on the cosmology ( $w$ ), gravity and SIDM etc.



Adhikari et al. 2018



# Previous Studies

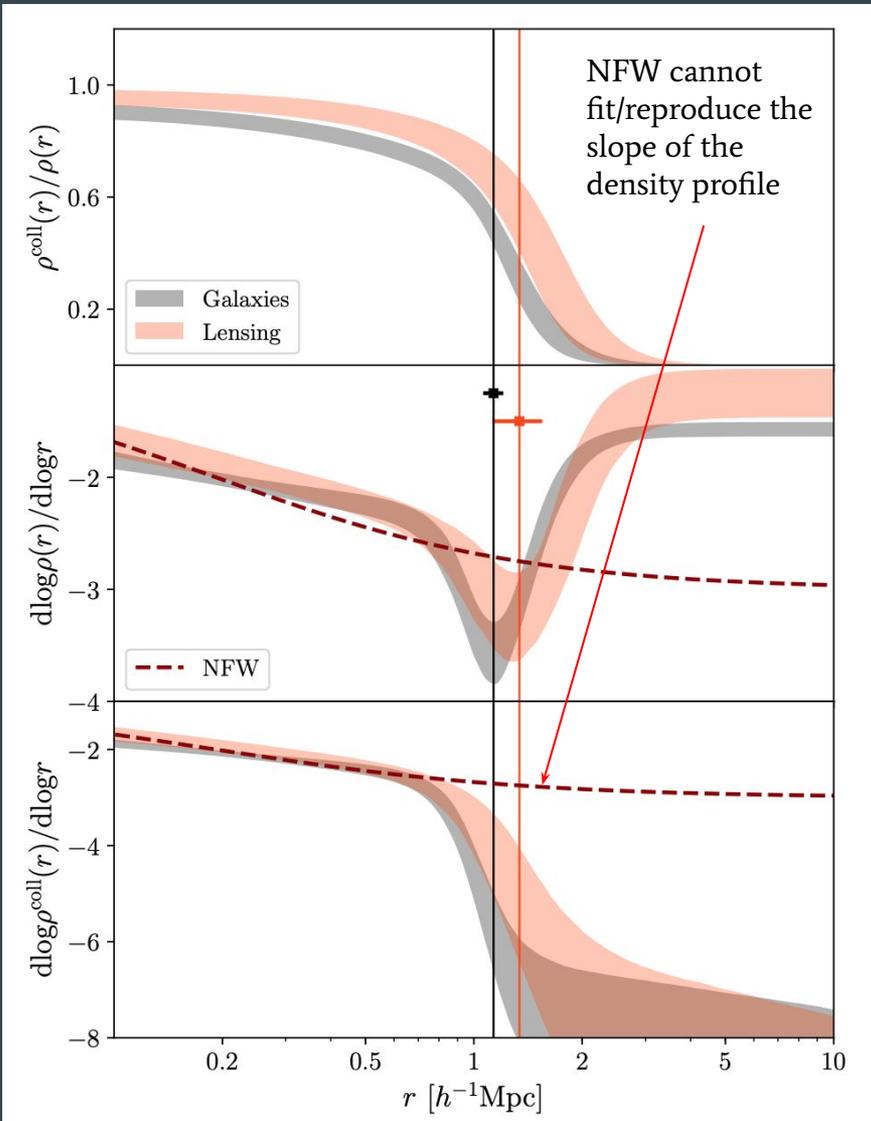
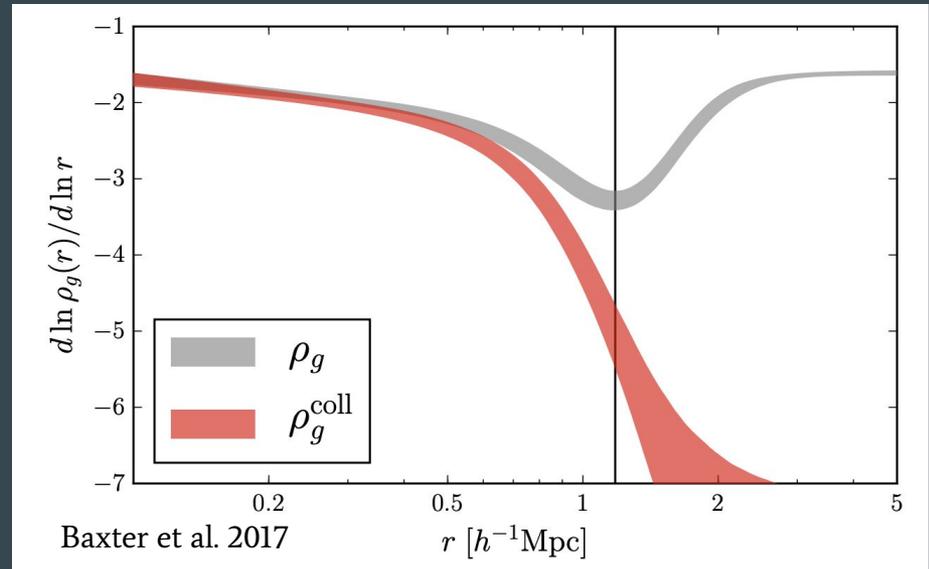
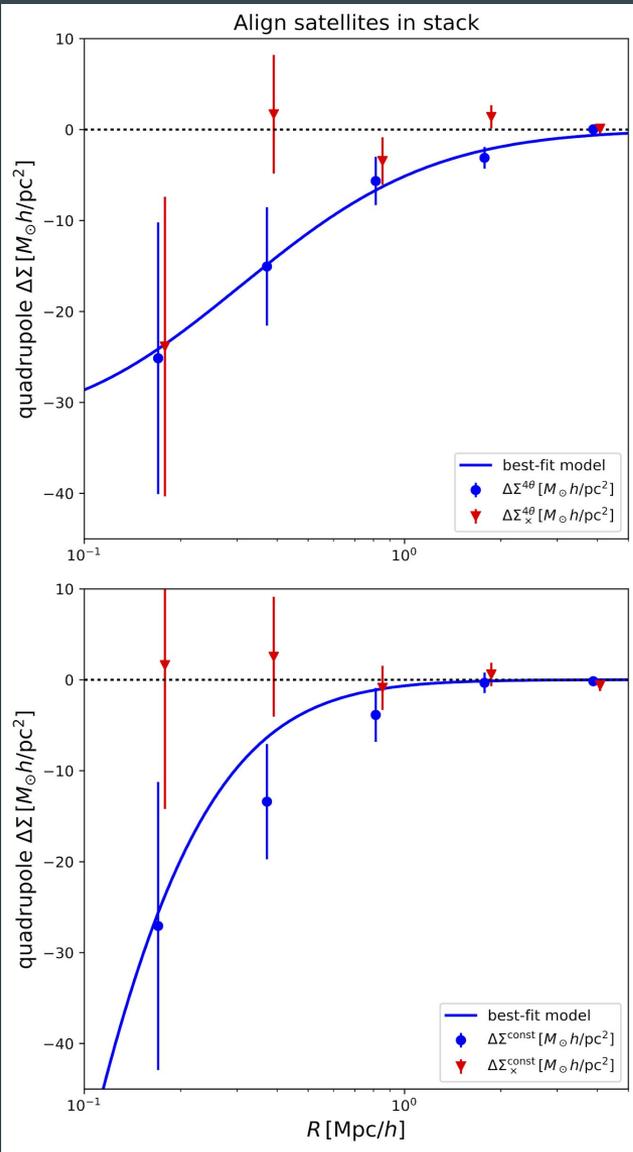


Fig. from Chang et al. 2017



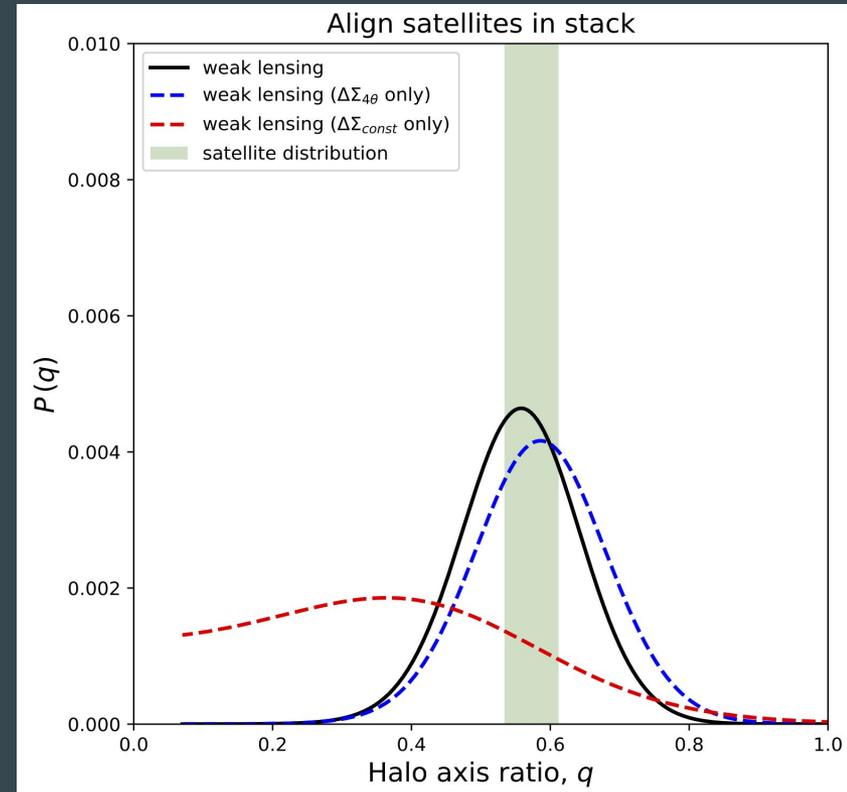
- First detection was reported in More et al. 2016
- Re-analysis with SDSS and DES by Baxter et al. 2017 and Chang et al. 2017, with **optically selected clusters (redMaPPer)**. Detection in lensing also.
- However, the location of  $r_{\text{sp}}$  is **~20% smaller** than the theory

# Ellipticity of Galaxy Clusters (Shin et al. 2018)



← quadrupole weak-lensing signal around stacked SDSS RM clusters

Constrained 2-D axis ratio  $b/a$  from the lensing (black) and from the satellite distribution (green band) ⇒



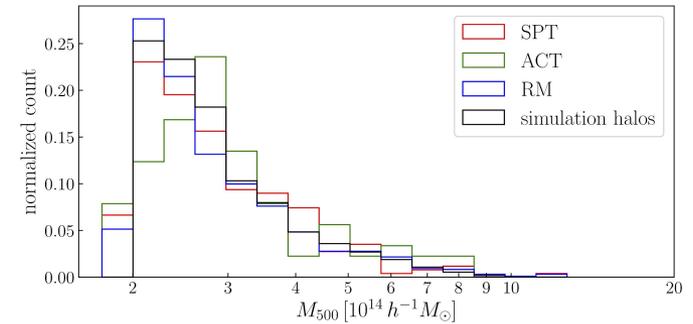
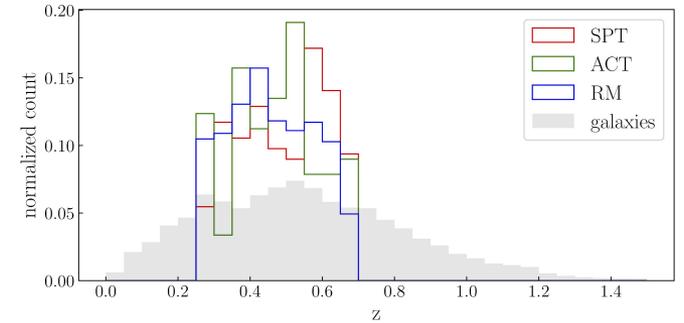
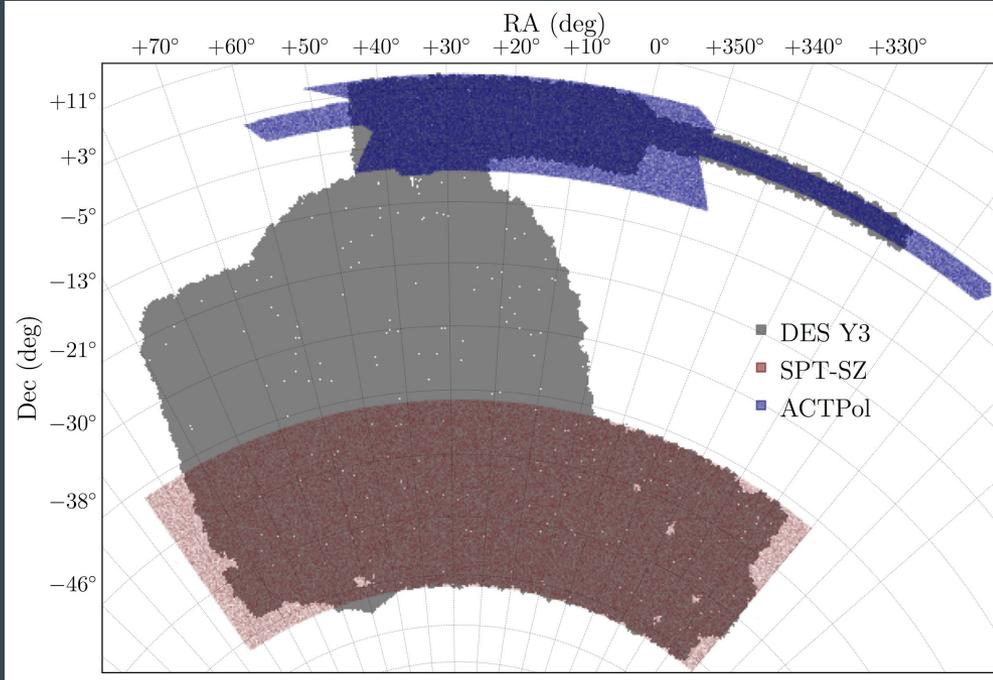
The shape and the boundary of galaxy clusters are typically **anisotropic** ⇒ Splashback feature on the major/minor axes of the underlying halos could tell us about the accretion along the filamentary structure (FUTURE WORK TOPIC)

# Optical vs SZ Cluster Samples

# SZ clusters

- We perform a similar analysis with clusters selected by **Sunyaev-Zel'dovich effect (SZ)** which identifies clumps of hot gas in the clusters (integrated pressure)
  - The SZ observable is completely **independent of all the observables in optical surveys** used to measure the feature (in particular, the galaxy density)
  - The SZ signal is expected to **correlate more tightly with cluster mass than optical richness**, reducing the impact of scatter in the mass-observable relation (easier comparison w.r.t. the simulation)
  - SZ-selection is expected to be **less affected by projection effects** than optical cluster finders ( $y \sim M^{5/3}$ )
  - The SZ-selected cluster samples employed here allow us to extend splashback measurements to the **high-mass, high-redshift** regime
- These SZ clusters are cross-correlated to the DES galaxies

# Data



Redshift Range:  $0.25 < z < 0.7$

Clusters:

~300 SPT clusters:  $\text{SNR} > 4.5$ ,  $\langle M_{500c} \rangle = 3.0e14 M_{\text{sun}}/h$ ,  $\langle z \rangle = 0.49$

~100 ACT clusters:  $\text{SNR}^* > 4.0$ ,  $\langle M_{500c} \rangle = 3.3e14 M_{\text{sun}}/h$ ,  $\langle z \rangle = 0.49$

~1,000 Optical (redMaPPer) clusters:  $\lambda > 58$ ,  $\langle z \rangle = 0.46$ ; mass matched to SPT clusters

Galaxies: DES galaxies with absolute magnitude cut at  $M_i < -19.87$

(apparent magnitude cut  $m_i = 22.5$  at the maximum redshift 0.7)

# Data

- Profiles of subhalos and dark matter particles are drawn from MultiDark Planck 2 (MDPL2) simulation
  - 1 (Gpc/h)<sup>3</sup> box size
  - redshift snapshot at  $z=0.49$
  - mean halo mass is matched to that of our SZ samples with a minimum mass cut
  - the scatter in SZ observable-mass relation does not change the splashback feature significantly

# Halo Model

We model the mass profile following Diemer & Kravtsov (2014)

⇒ truncated Einasto profile (1-halo) in addition to the power-law infalling term (2-halo)

: good up to  $\sim 9R_{\text{vir}}$  (above it, infall regime breaks down)

$$\rho(r) = \rho^{\text{coll}}(r) + \rho^{\text{infall}}(r)$$

$$\rho^{\text{coll}}(r) = \rho^{\text{Ein}}(r)f_{\text{trans}}(r)$$

$$\rho^{\text{Ein}}(r) = \rho_s \exp\left(-\frac{2}{\alpha} \left[\left(\frac{r}{r_s}\right)^\alpha - 1\right]\right)$$

$$f_{\text{trans}}(r) = \left[1 + \left(\frac{r}{r_t}\right)^\beta\right]^{-\gamma/\beta}$$

$$\rho^{\text{infall}}(r) = \rho_0 \left(\frac{r}{r_0}\right)^{-s_e} \text{ Truncation of the 1-halo Einasto profile}$$

- Integrated along the line of sight into a 2D profile

- MCMC fitting with jackknife covariance

- Priors on  $\alpha$ ,  $\beta$  and  $\gamma$  from previous simulation studies

(Gao+ 2008, Diemer&Kravtsov 2014)

- Priors on miscentering from Saro+2015 (SZ: SPT) and Rykoff+2014 (BCG: ACT/RM)

# Correlation Function

The 2D two-point correlation function measures the **excessive probability** of finding two galaxies being separated by a distance of R

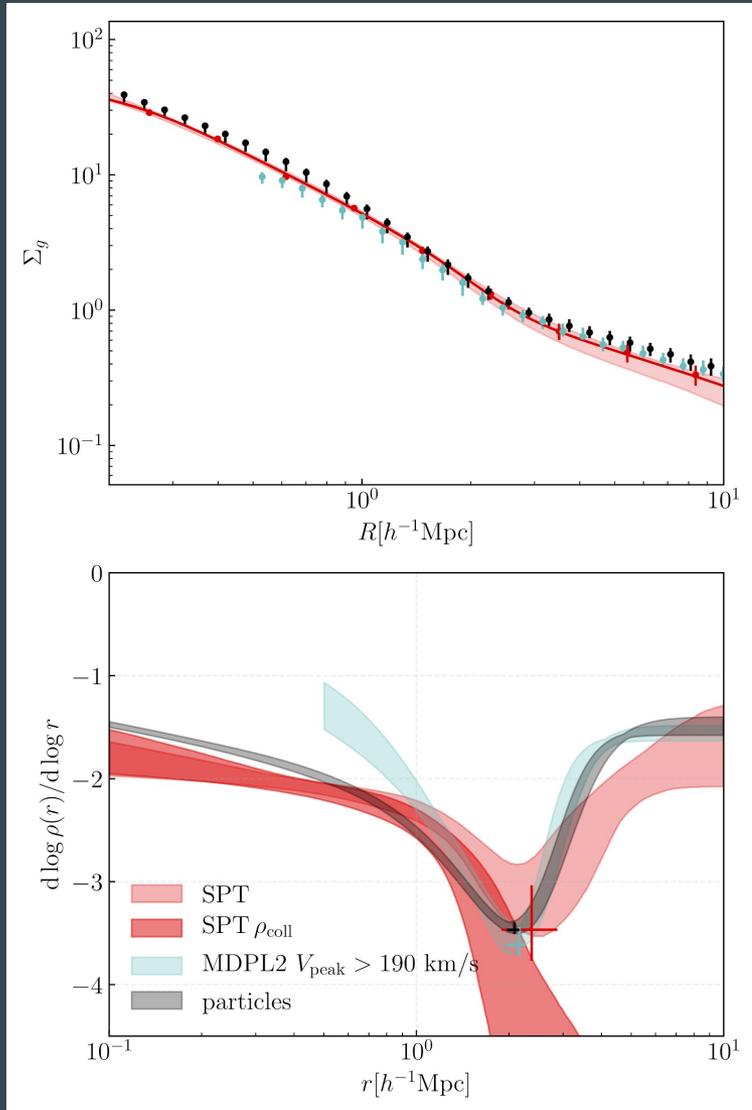
$$dP(R) = n_1 n_2 (1 + \omega(R)) dA_1 dA_2$$

Thus, the mean-subtracted galaxy surface density around the clusters can be expressed as,

$$\Sigma_g(R) = \langle \Sigma_g \rangle \omega(R)$$

When applying the absolute magnitude cut & calculating correlation function, we assume all the galaxies are located at the cluster redshift  
⇒ the correlation function picks up the galaxies that are correlated with the clusters: **avoiding the photo-z uncertainties of the galaxies**

# Result: SPT clusters



$$r_{\text{sp}} = 2.37^{+0.51}_{-0.48} \text{ Mpc}/h$$

$$\text{slope at } r_{\text{sp}} = -3.47^{+0.43}_{-0.30}$$

$$\text{slope of } \rho_{\text{coll}} \text{ at } r_{\text{sp}} = -5.17^{+1.06}_{-0.60}$$

For simulation halos,

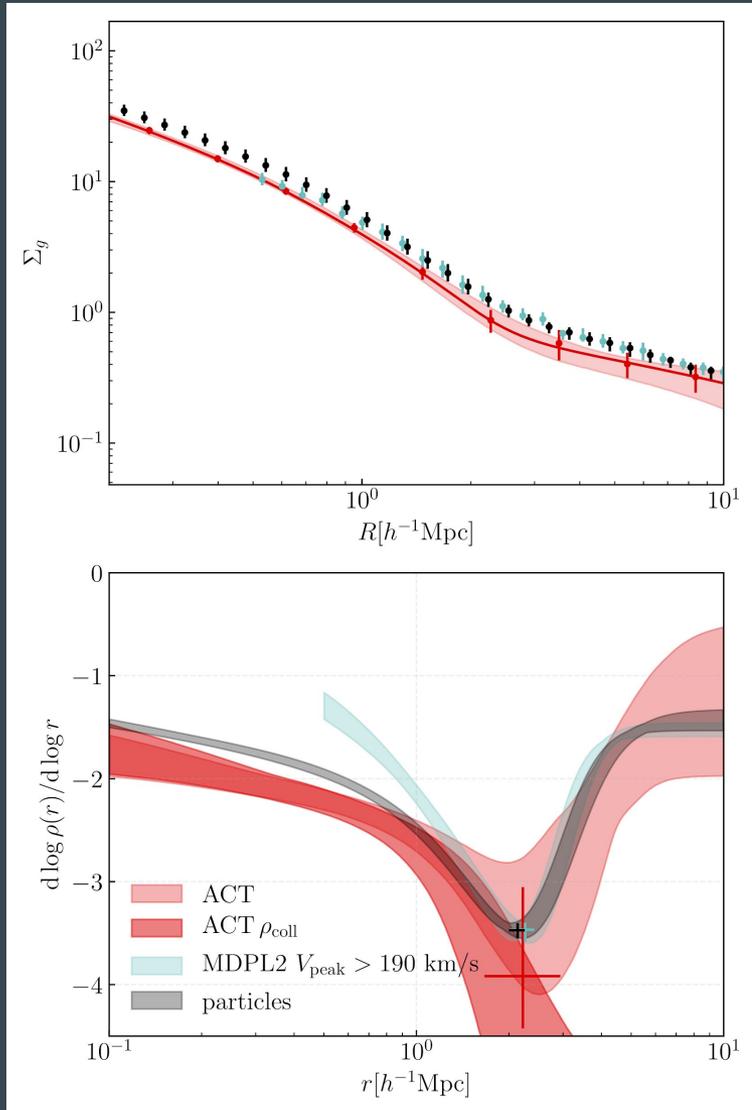
$$r_{\text{sp}} = 2.16^{+0.10}_{-0.20} \text{ Mpc}/h \text{ (subhalos, cyan)}$$

$$2.08^{+0.08}_{-0.11} \text{ Mpc}/h \text{ (particles, black)}$$

⇒ The observed feature agrees with that of simulation within  $1\sigma$

\* The subhalos lose mass due to tidal interactions and pass below the resolution limit in the central regions, resulting in a flattening of the inferred slope

# Result: ACT clusters



$$r_{\text{sp}} = 2.22^{+0.72}_{-0.56} \text{ Mpc/h}$$

$$\text{slope at } r_{\text{sp}} = -3.92^{+0.86}_{-0.51}$$

$$\text{slope of } \rho_{\text{coll}} \text{ at } r_{\text{sp}} = -5.40^{+1.27}_{-0.58}$$

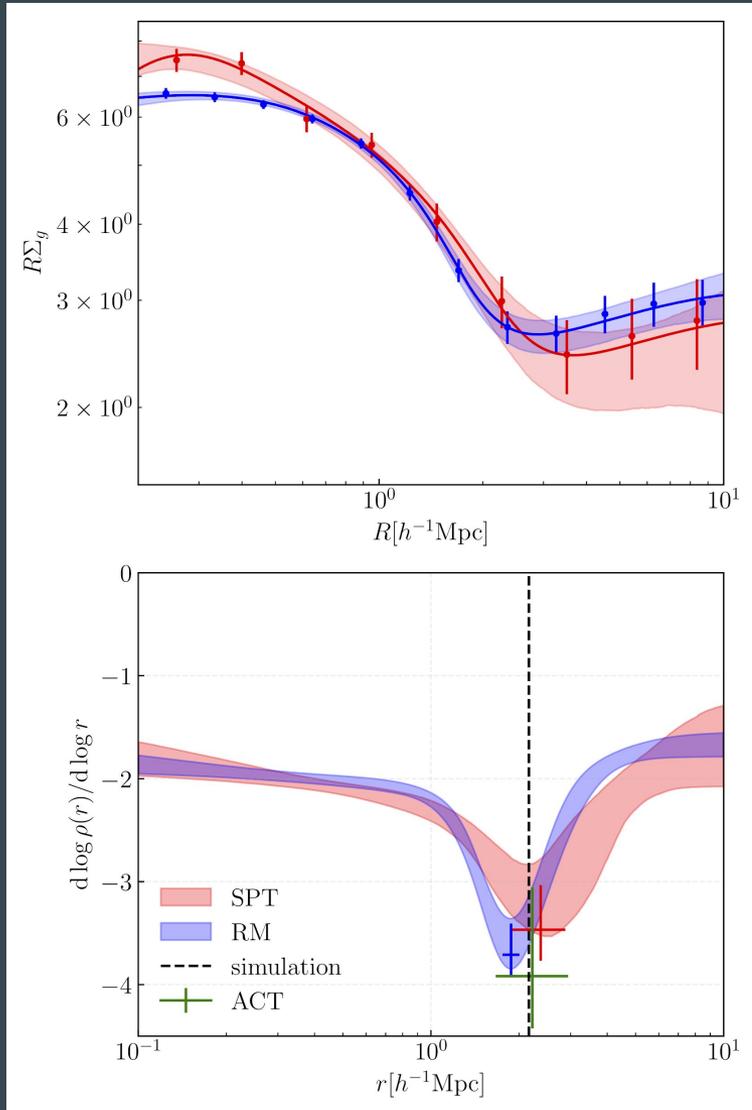
For simulation halos,

$$r_{\text{sp}} = 2.26^{+0.15}_{-0.25} \text{ Mpc/h (subhalos, cyan)}$$

$$2.13^{+0.12}_{-0.14} \text{ Mpc/h (particles, black)}$$

$\Rightarrow$  The observed feature agrees with that of simulation within  $1\sigma$

# Result: RM (vs simulation)



$$r_{\text{sp}} = 1.88^{+0.13}_{-0.12} \text{ Mpc/h (blue)}$$

$$\text{slope at } r_{\text{sp}} = -3.71^{+0.30}_{-0.20}$$

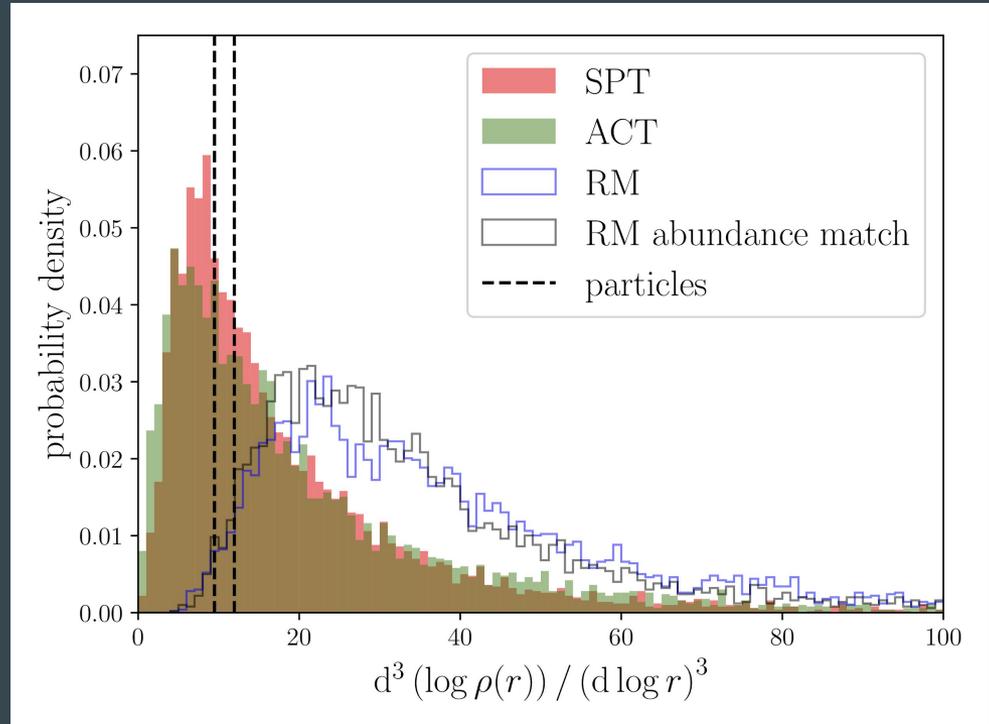
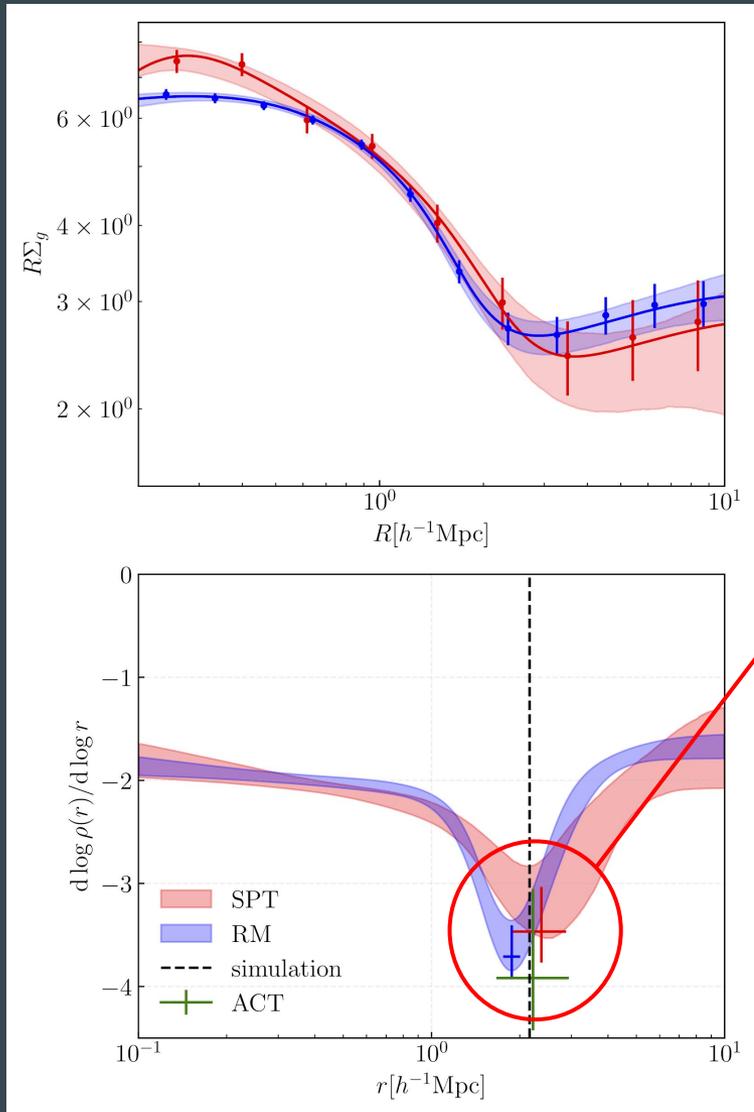
$$\text{slope of } \rho_{\text{coll}} \text{ at } r_{\text{sp}} = -5.52^{+0.88}_{-0.61}$$

For simulation halos,

$$r_{\text{sp}} = 2.16^{+0.10}_{-0.20} \text{ Mpc/h (dashed line)}$$

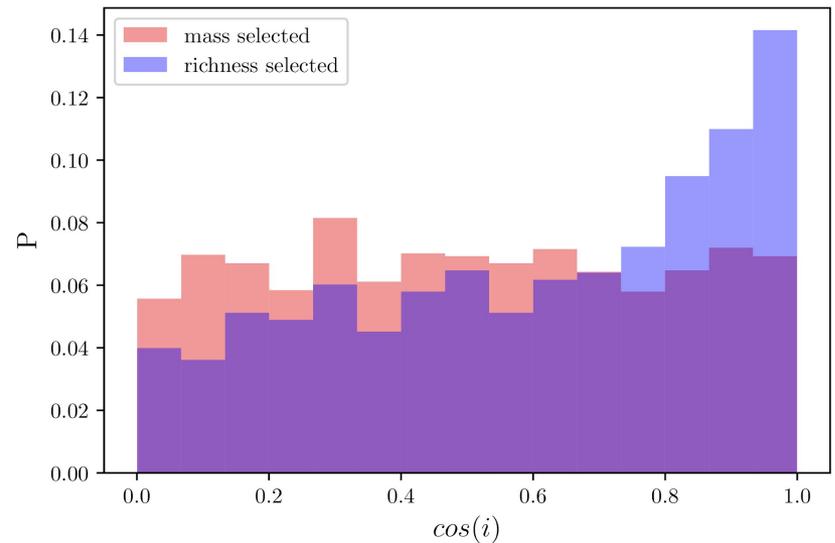
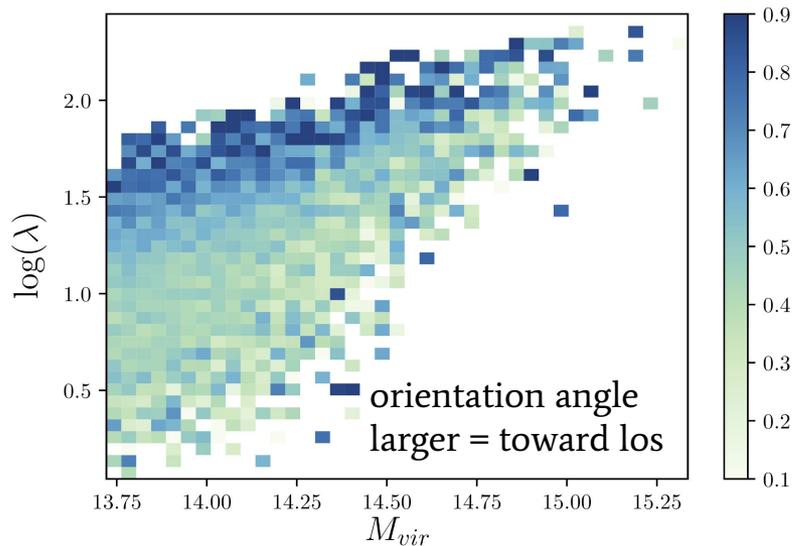
$\Rightarrow$  The observed  $r_{\text{sp}}$  in optically selected RM clusters are  $\sim 2\sigma$  lower than that of simulation (subhalo profile)

# Result: shape of the feature (SZ vs RM vs Sim)



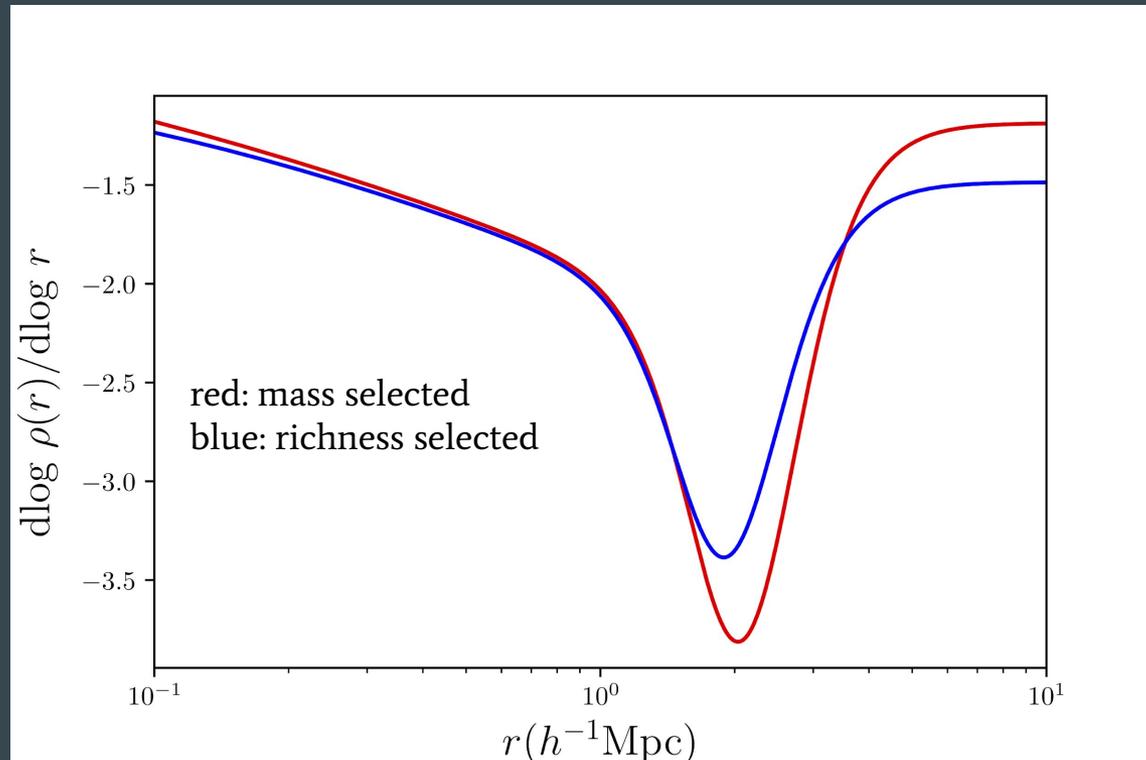
The RM clusters exhibit a **sharper splashback feature** (larger 3<sup>rd</sup> deriv at  $r_{\text{sp}}$ ) than that of SZ clusters and simulation halos

# Systematic test: comparison between mass- and richness-selected sample



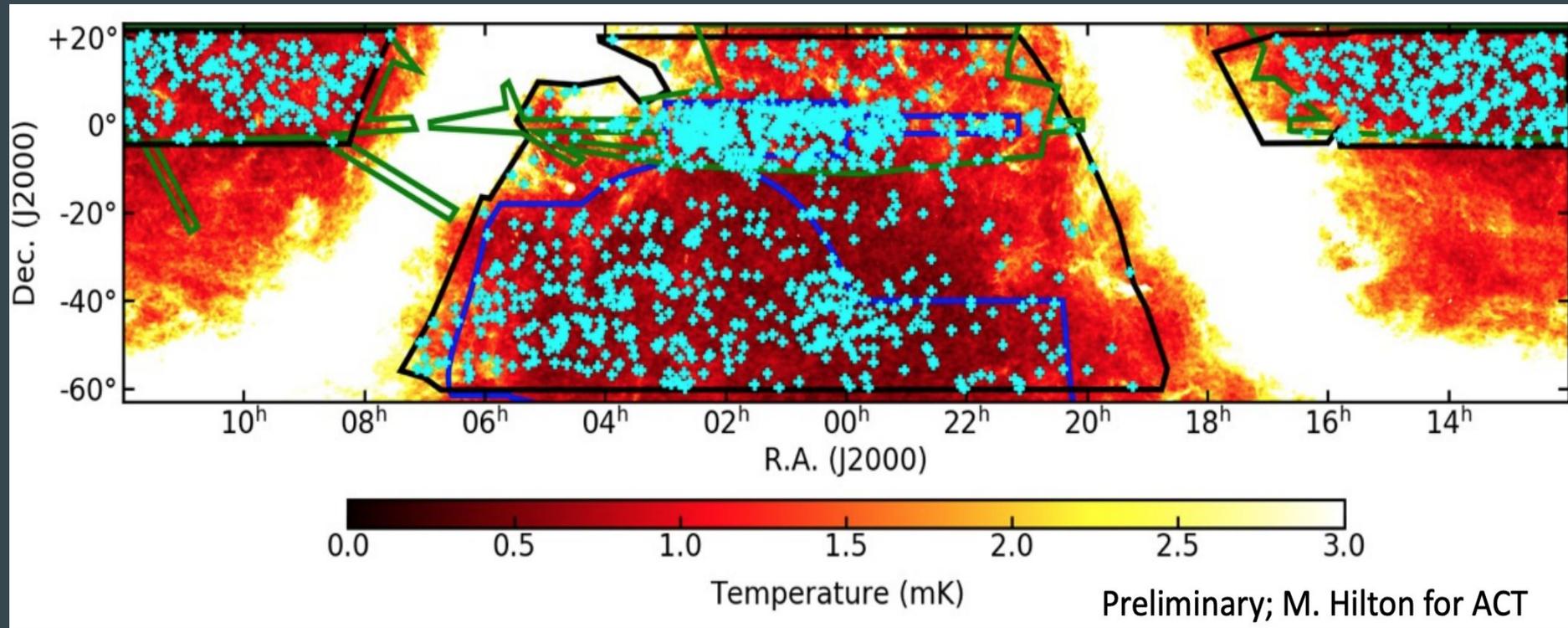
- For a given mass, optical clusters with high richness tend to be more aligned w.r.t. the l.o.s. than those with low richness
- Thus, richness selection ( $>20$ ) results in a biased selection of clusters in terms of their orientation

# Systematic test: comparison between mass- and richness-selected sample



- The difference is  $\sim 6\%$  in the location of the  $r_{\text{sp}}$ , but it is not enough to explain the observed discrepancy between RM and SZ/simulation

# AdvACT clusters

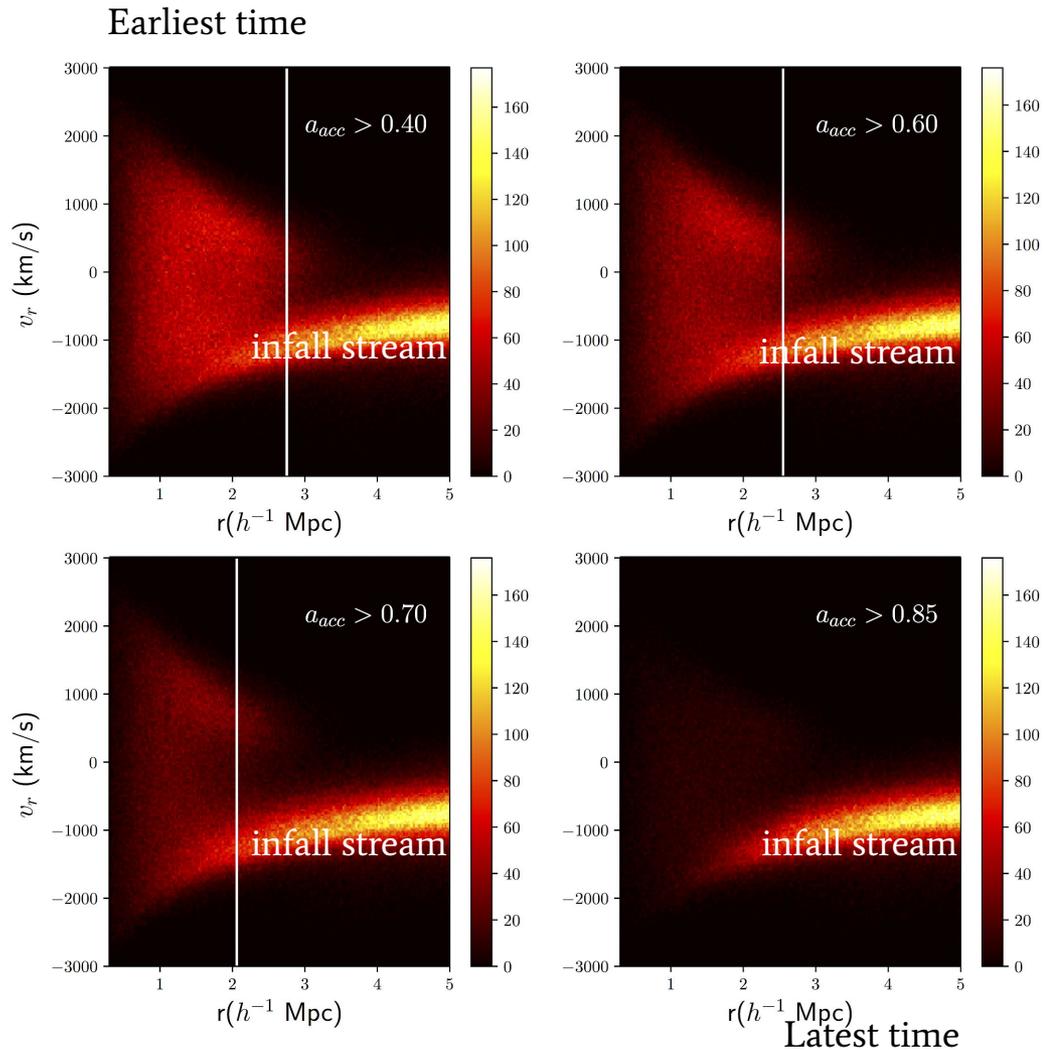


863 clusters (subject to change) in the  
DES footprint having  $\text{SNR} > 4$ , w/  
 $0.15 < z < 0.7$   
 $\langle M_{500c} \rangle = 3.0e14 \text{ Msun}/h$   
 $\langle z \rangle = 0.44$   
ANALYSIS UNDERWAY!

w.r.t. the previous SPT measurements,  
the error bars in the galaxy density  
profile and lensing profile are  
expected to reduce by a factor of  $\sim 2$   
as we will have 3-4x more clusters

# Galaxy Quenching and Splashback

# Infalling particles in phase space

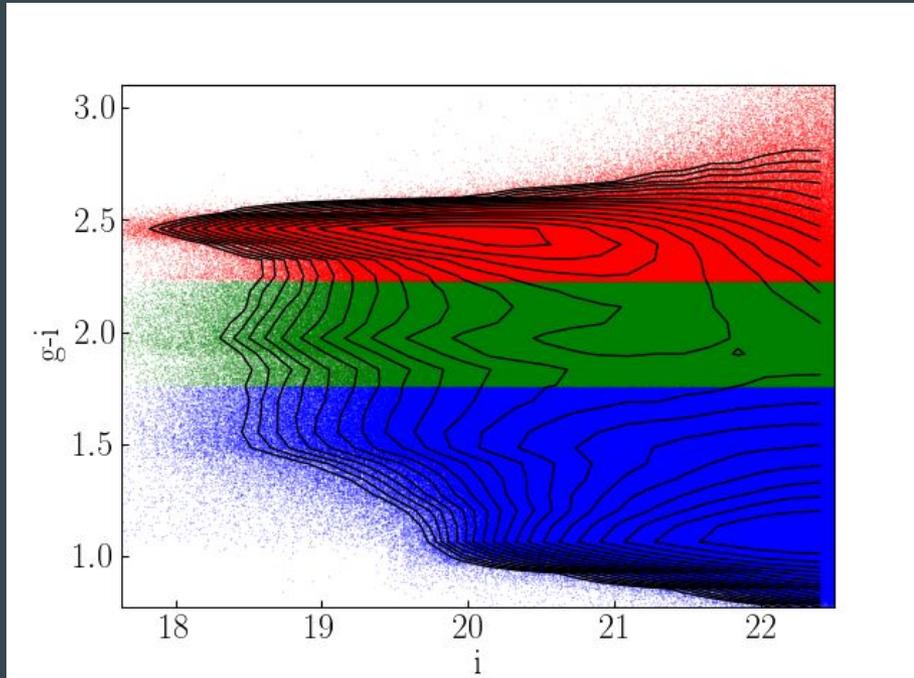


⇐ Subhalos accreted to a cluster at different times in simulation

Galaxies in the **infall stream** do **not** show any splashback feature, while those that have completed at least one crossing show a distinctive splashback feature

⇒ Can we **separate the infall population** from the observational data?

# Split of galaxies in color space



⇐ Galaxies are split in  $g-i$  color in each redshift bin of  $\Delta z=0.025$

: 20% red, 20% green, 60% blue

- The variation in the fraction with the redshift is not significant given our noise level

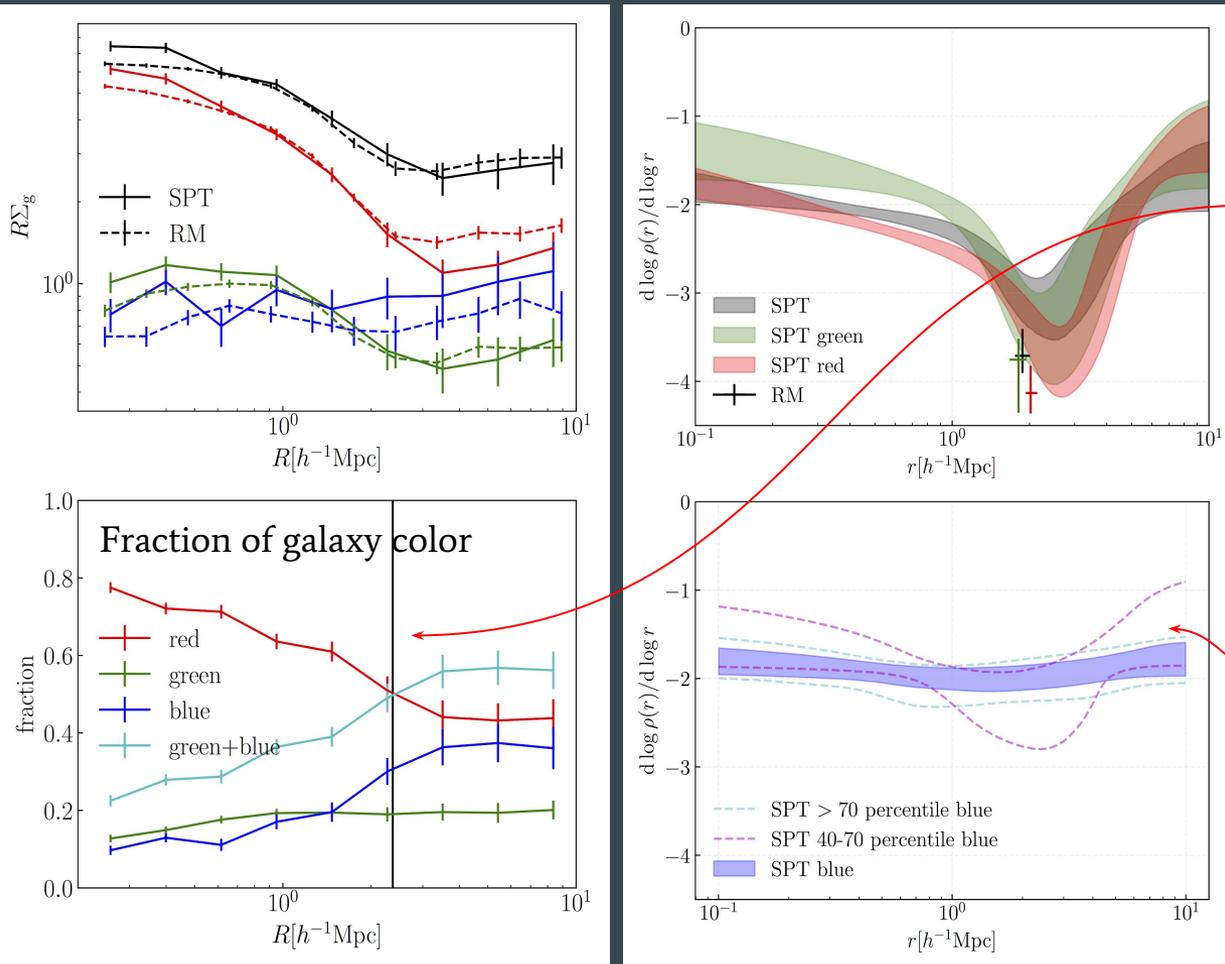
- Blue star-forming galaxies are quenched within clusters, becoming red quiescent galaxies, by various possible processes (Gunn&Gott 1972, Abadi+1999, Larson+1980, Wetzell+2013, von der Linden+2010, Brodwin+2013, Ehlert+2014, Wagner+2015)
- With these color-split galaxies samples, the same analysis has been done with the same SPT cluster as before

# Result: profiles of galaxies with different colors

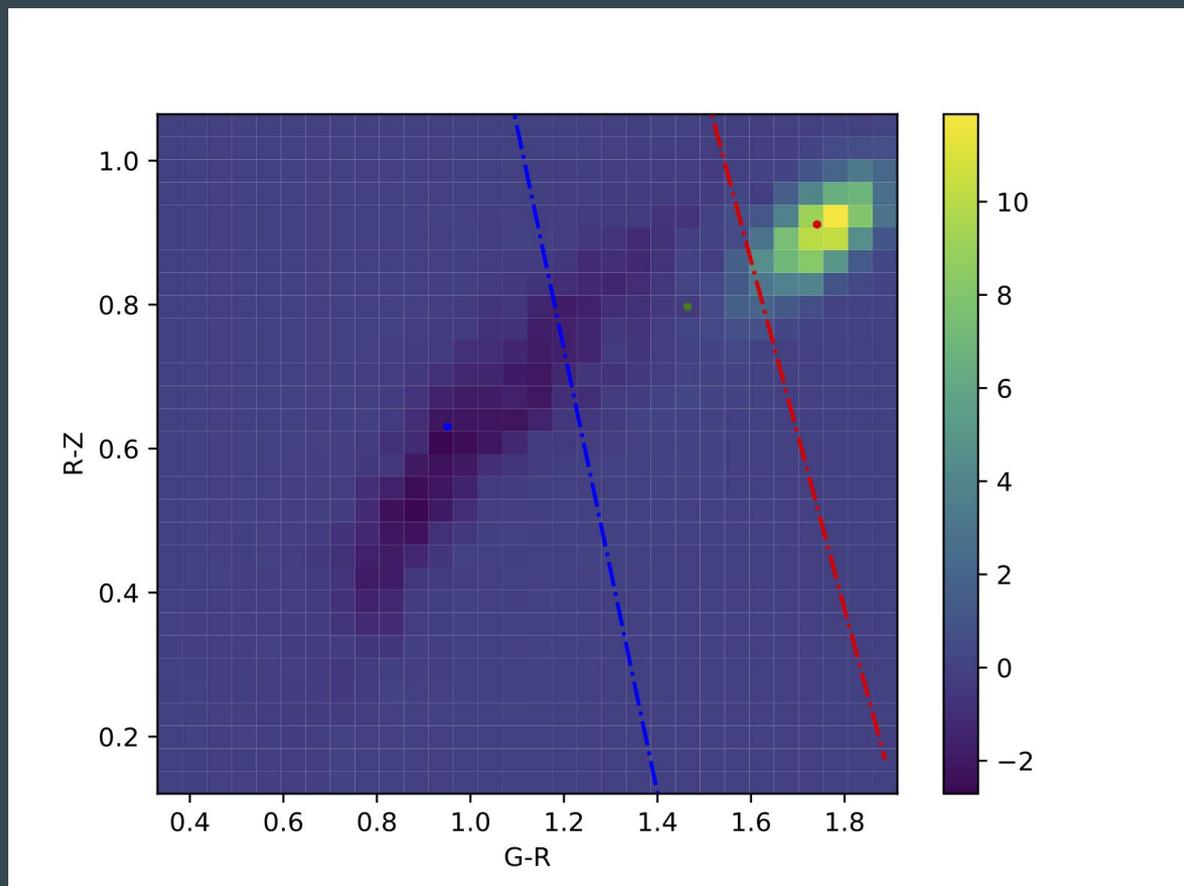
We measure profiles of galaxies split on color.

The upturn of the red fraction around  $r_{\text{sp}}$  = evidence of **quenching of galaxy star formation** inside clusters

Blue galaxies are consistent to a **pure power-law** profile; indicating that they are still on their **first infall passage** (with S. Adhikari)



# New color split scheme (w/ AdvACT SZ clusters)

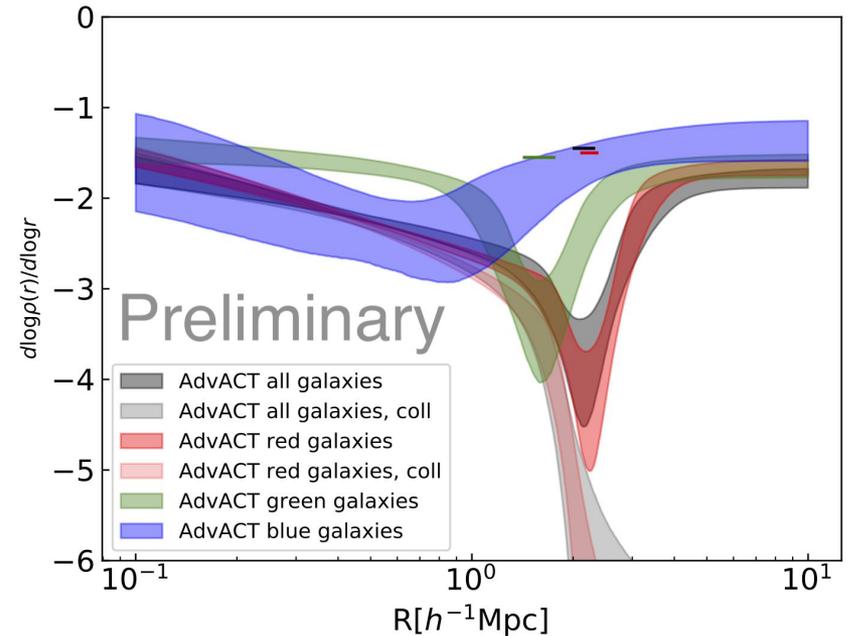
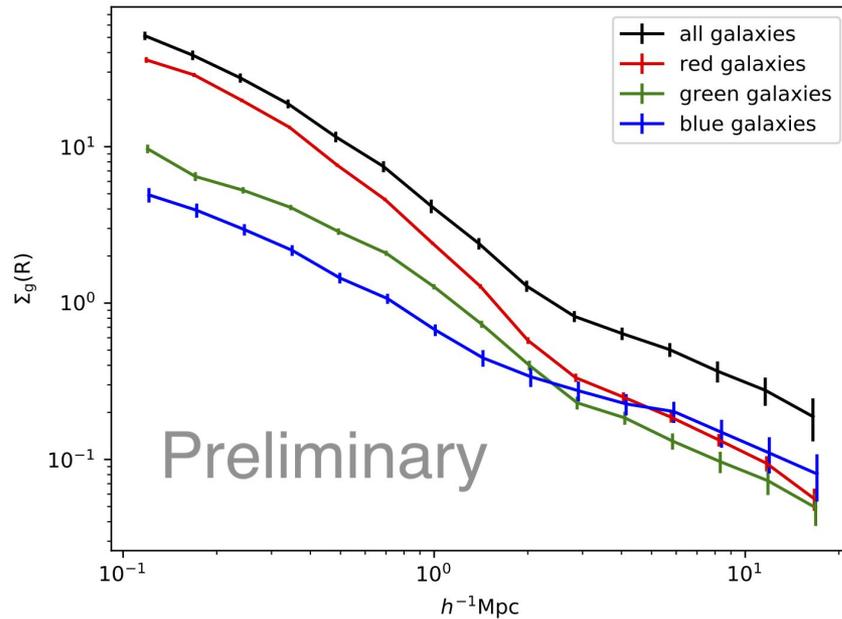


In G-R vs R-Z **color-color space**:

Subtract the density of all galaxies from that  $< 2.5 \text{ Mpc/h}$  from the AdvACT clusters, in each redshift bin of  $dz=0.075$

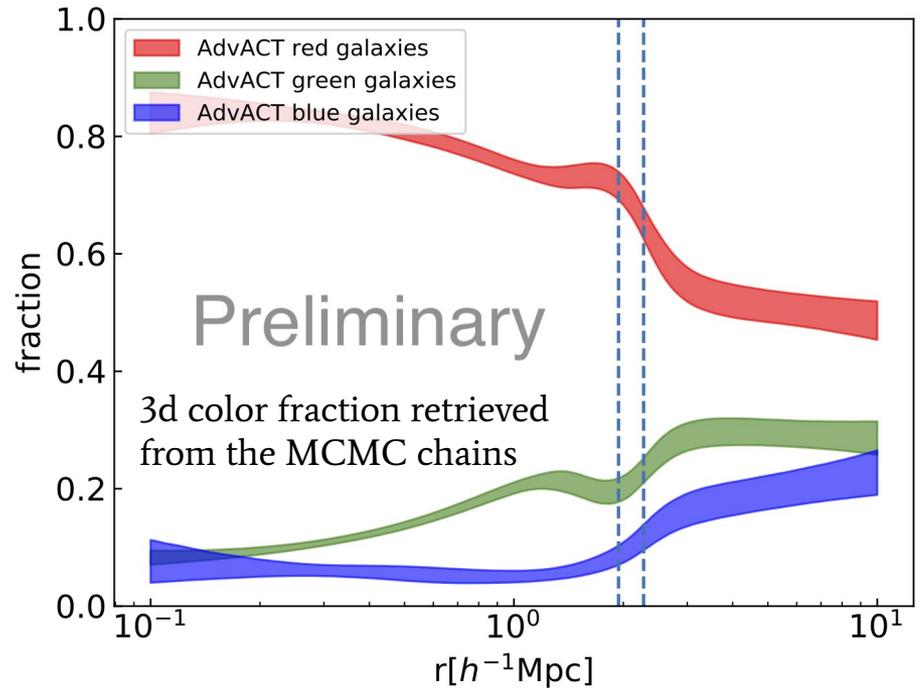
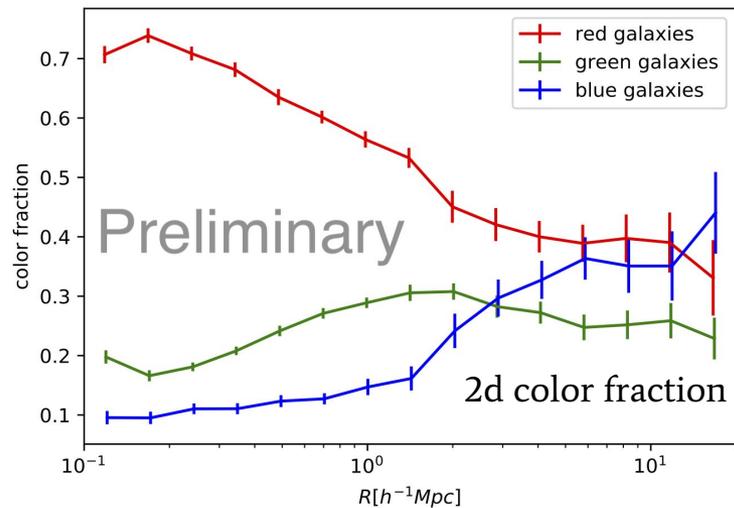
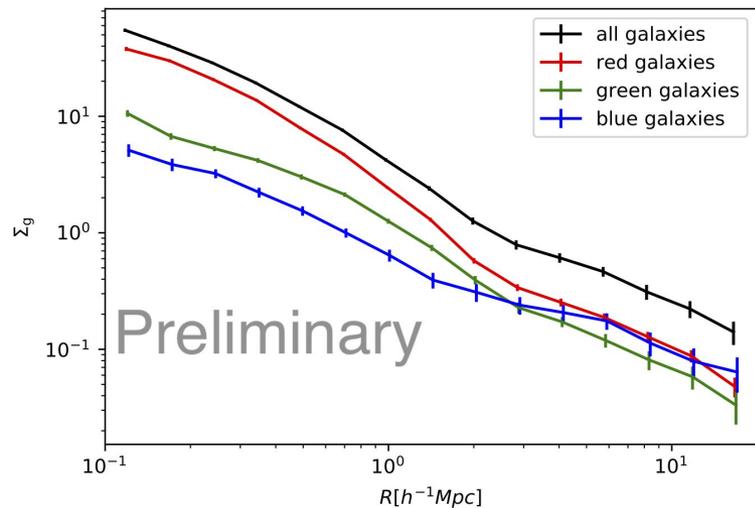
$\Rightarrow$  excess of red galaxies, deficit of blue galaxies as well as the green valley

# (Preliminary) result with AdvACT clusters



Blue galaxy profile is again largely consistent with a power law profile: majority of them are still in their **first infall passage**

# (Preliminary) result with AdvACT clusters



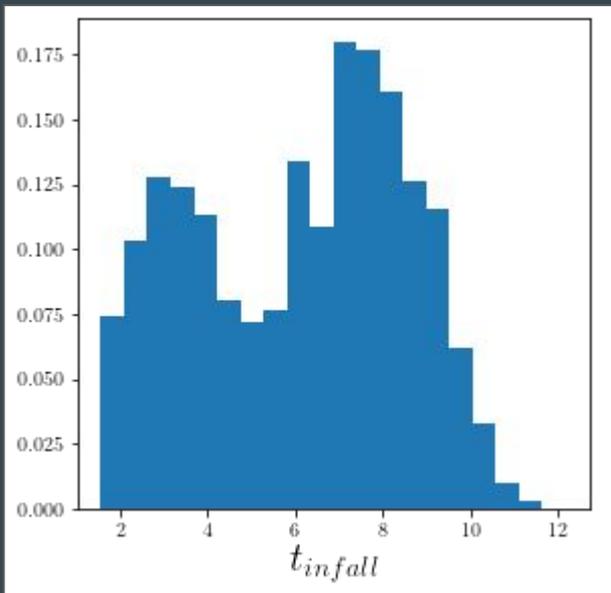
Around the splashback radius, **red fraction** starts to **increase inward**, while blue fraction decreases

⇒ we can use this fraction to constrain the quenching timescale quantitatively, per quenching model

# Constraining SFR quenching timescale

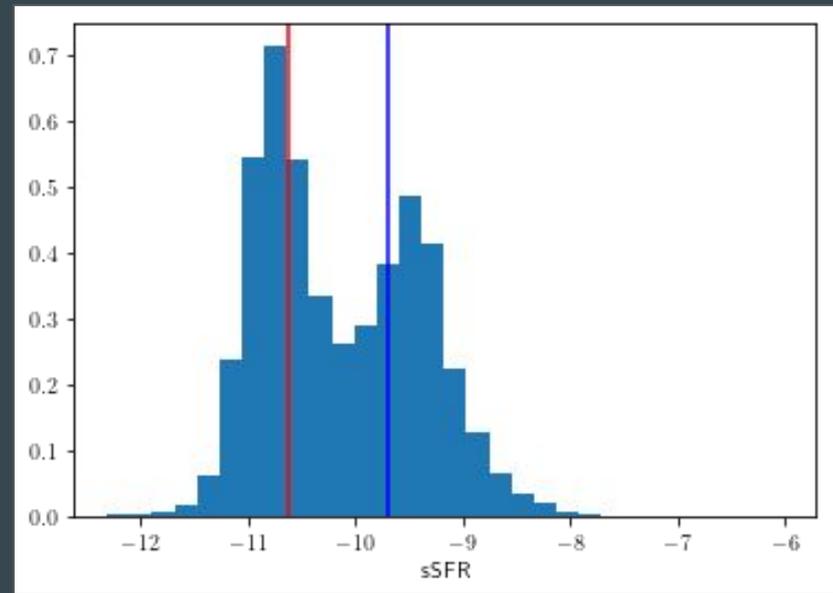
$$\text{SFR}_{\text{sat}}(t) = \begin{cases} \text{SFR}_{\text{cen}}(t) & t < t_{Q, \text{start}} \\ \text{SFR}_{\text{cen}}(t_{Q, \text{start}}) e^{\left\{-\frac{(t-t_{Q, \text{start}})}{\tau_{Q, \text{fade}}}\right\}} & t > t_{Q, \text{start}} \end{cases}$$

quenching starts after  $t_{Q, \text{start}}$  after infall, followed by quenching w/ with timescale of  $\tau_{Q, \text{fade}}$



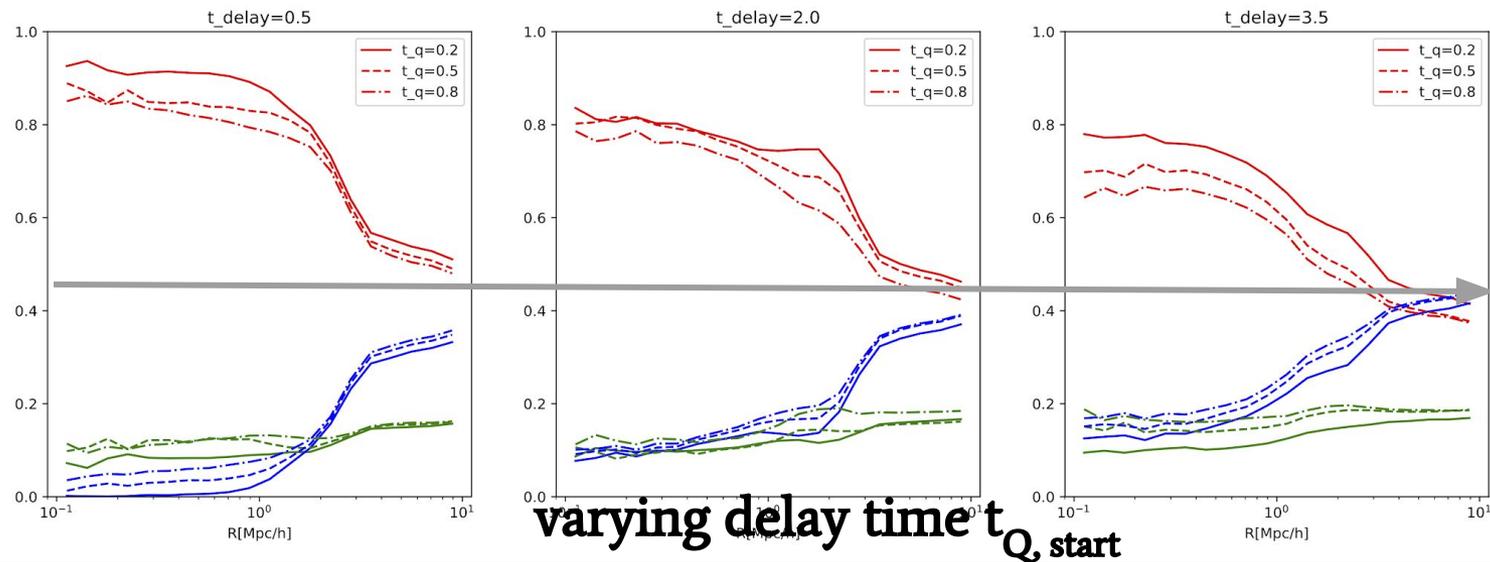
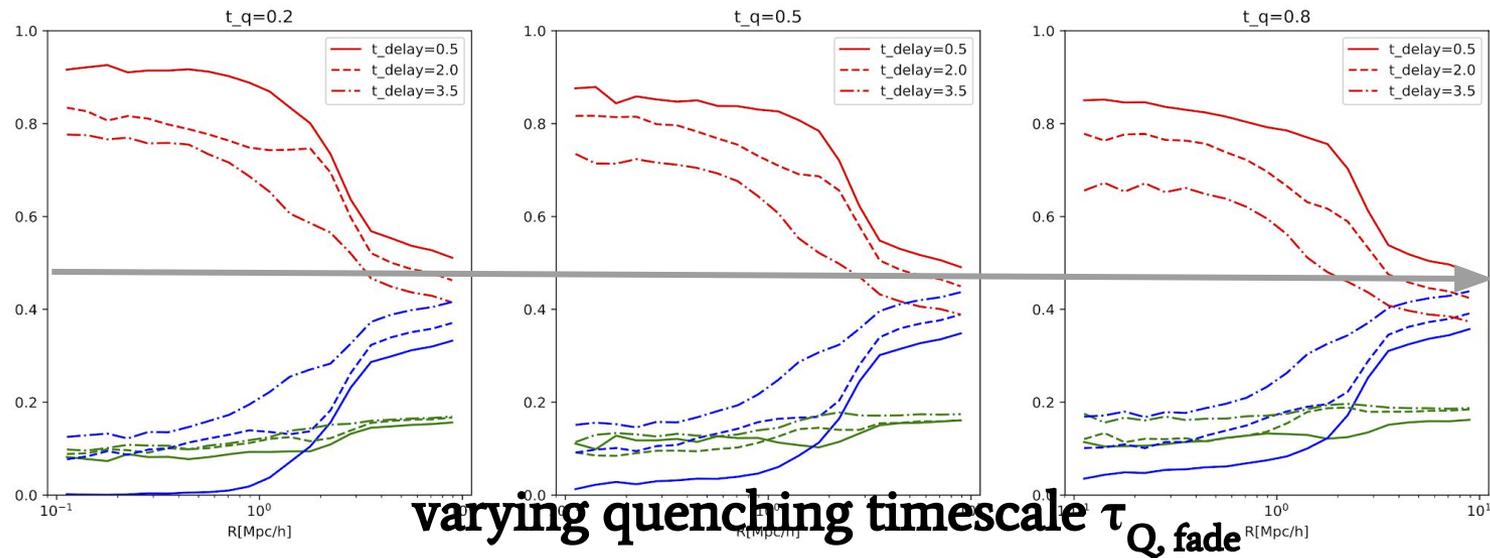
← dist. of infall time in MDPL2 simulation

dist. of sSFR in SDSS spec-z sample ⇒



- Using the sSFR dist. of the field (in the data) and the quenching model, assign each subhalo (in the simulation) a sSFR value ⇒ compare the fraction of color (r/g/b) as a function of radial distance, to the observed value to constrain the quenching params

# Ex) Color fraction in simulation w/ different params



# Summary

- Splashback feature is a plausible **physical** boundary of the halo : it is sensitive to e.g. **accretion rate** of the halo, cosmology (**w**), and **gravity**
- While the observed features in the **optical** cluster samples are located at ~20% **smaller radii** than in simulation halos, the **SZ** cluster samples used in this study show **consistent** splashback features as in the simulation
- Orientation bias and mass calibration does not fully explain the discrepancy in optical clusters (ongoing work)
- The profiles split in galaxy colors suggests that 1) galaxies start to be **quenched at/around  $r_{sp}$**  and 2) **blue** galaxies are mostly still in their **first infall** passage
- With a larger SZ cluster sample w/ AdvACT, we can constrain the parameters in quenching models using the fraction of galaxy colors as a function of radius
- Ongoing and future SZ / X-ray cluster survey (AdvACT, SPT3G, SPTPol, SO, CMB-S4, e-Rosita etc.) will provide additional understandings of physics of galaxy clusters