

# Large Scale Structure Beyond the Two-Point Function

### **Oliver Philcox (Princeton / IAS)**

Physics Division Seminar, LBNL 11/23/21

- 1. What is the nature of dark matter and dark energy?
- 2. How fast is the Universe expanding?
- 3. What happened during inflation?
- 4. Is General Relativity correct?
- 5. What is the **neutrino** mass?

Galaxy surveys can shed light on these...

### THE QUANTUM UNIVERSE

Inflation produces quantum fluctuations in the primordial matter distribution

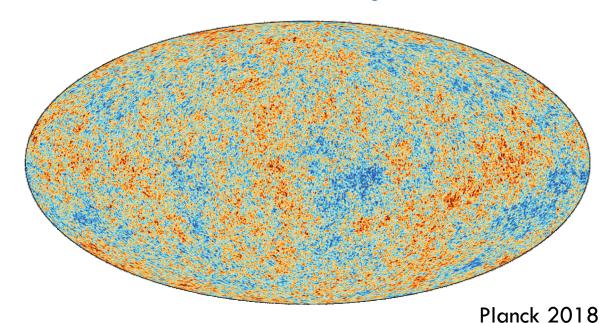
$$\delta(\mathbf{x}) = \frac{\rho(\mathbf{x}) - \langle \rho \rangle}{\langle \rho \rangle}$$

These are (mostly) adiabatic and Gaussian

 $\delta(\mathbf{k}) \sim \mathcal{N}\left(0, P_L(\mathbf{k})\right)$ 

> All information is in the **power spectrum** 

Quantum Fluctuations in the Cosmic Microwave Background



### THE QUANTUM UNIVERSE

Inflation produces quantum fluctuations in the primordial matter distribution

 $\delta(\mathbf{x}) = \frac{\rho(\mathbf{x}) - \langle \rho \rangle}{\langle \rho \rangle}$ 

These are (mostly) adiabatic and Gaussian

 $\delta(\mathbf{k}) \sim \mathcal{N}\left(0, P_L(\mathbf{k})\right)$ 

All information is in the power spectrum or two-point correlation function (2PCF) Fourier Space  $\langle \delta(\mathbf{k}) \delta(\mathbf{k}') \rangle = (2\pi)^3 \delta_{\mathrm{D}} (\mathbf{k} + \mathbf{k}') P_L(\mathbf{k})$  $\langle \delta(\mathbf{k}) \delta(\mathbf{k}') \delta(\mathbf{k}'') \rangle = 0$ 

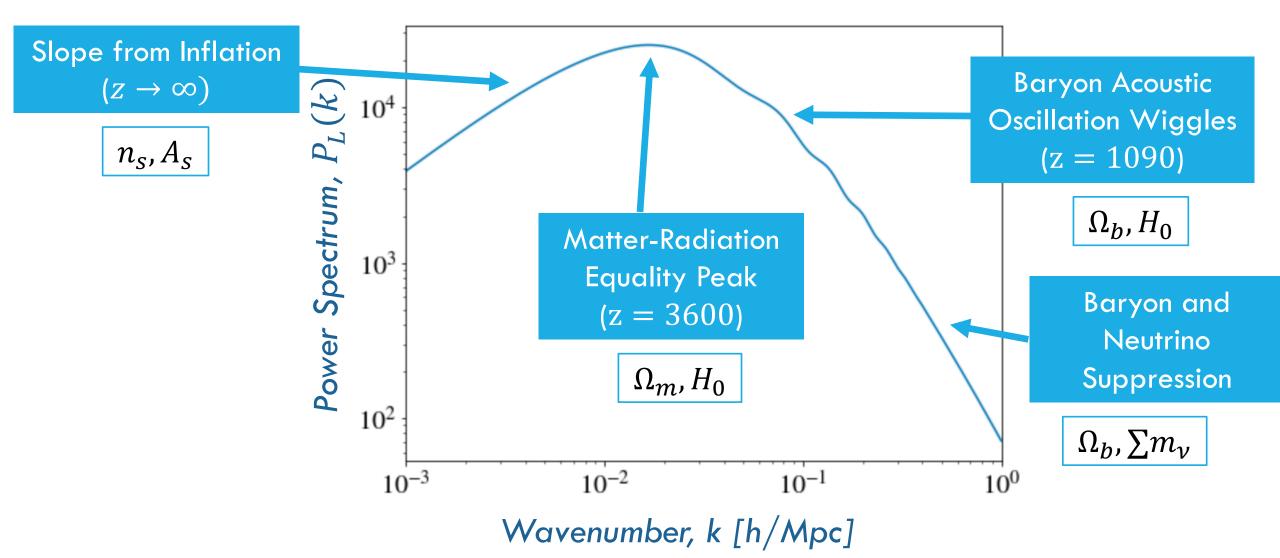
**Configuration Space** 

 $\langle \delta(\mathbf{r}_1) \delta(\mathbf{r}_2) \rangle = \xi_L(\mathbf{r}_1 - \mathbf{r}_2)$ 

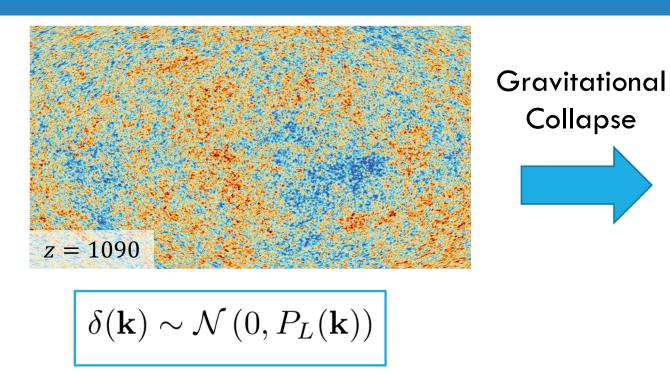
 $\langle \delta(\mathbf{r}_1) \delta(\mathbf{r}_2) \delta(\mathbf{r}_3) \rangle = 0$ 

### LINEAR POWER SPECTRUM

### $\langle \delta(\mathbf{k})\delta(\mathbf{k}')\rangle = (2\pi)^3 \delta_{\rm D} \left(\mathbf{k} + \mathbf{k}'\right) P_L(\mathbf{k})$



## THE LATE UNIVERSE IS NOT GAUSSIAN



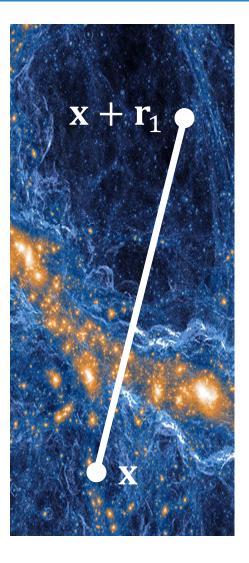
- > All information contained in the power spectrum  $[P_L(\mathbf{k})]$
- **No** higher order statistics needed!

z = 0

$$\delta(\mathbf{k}) \not\sim \mathcal{N}\left(0, P_L(\mathbf{k})\right)$$

- Not all information contained in the power spectrum
- Higher-order statistics needed!

### NON-GAUSSIAN DENSITY $\Rightarrow$ NON-GAUSSIAN STATISTICS



### Gaussian

- 1. Power Spectrum:
- $P(\mathbf{k}_1) = \langle \delta(\mathbf{k}_1) \delta(\mathbf{k}_2) \rangle'$
- 2. 2-Point Correlation Function:

$$\xi(\mathbf{r}_1) \equiv \langle \delta(\mathbf{x}) \delta(\mathbf{x} + \mathbf{r}_1) \rangle$$

### Non-Gaussian

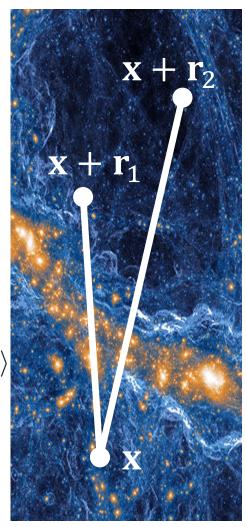
1. Bispectrum:

$$B(\mathbf{k}_1,\mathbf{k}_2) = \langle \delta(\mathbf{k}_1)\delta(\mathbf{k}_2)\delta(\mathbf{k}_3) \rangle'$$

2. 3-Point Correlation Function:

$$\zeta_3(\mathbf{r}_1, \mathbf{r}_2) \equiv \langle \delta(\mathbf{x}) \delta(\mathbf{x} + \mathbf{r}_1) \delta(\mathbf{x} + \mathbf{r}_2) \rangle$$

And beyond...



### WHAT MAKES UP THE BISPECTRUM?

$$B_g(\mathbf{k}_1, \mathbf{k}_2) = \left[ 2b_1^3 F_2(\mathbf{k}_1, \mathbf{k}_2) + b_2 b_1^2 + 2b_{s^2} b_1^2 \left( \hat{\mathbf{k}}_1 \cdot \hat{\mathbf{k}}_2 - 1/3 \right) \right] P_L(k_1) P_L(k_2) + 2 \text{ perms.}$$

The galaxy bispectrum depends on galaxy formation physics, gravity, and early-Universe cosmology.\*

 $\triangleright$  To obtain **all** the large-scale information in the initial conditions, we need: $^{*}$ 

- Power Spectra
- Bispectra
- Trispectra

 $\sim P_L(k)$   $\sim P_L^2(k)$  $\sim P_L^3(k)$ 

• etc.

8 \*ignoring higher-order perturbative effects, redshift-space distortions, renormalization, etc.

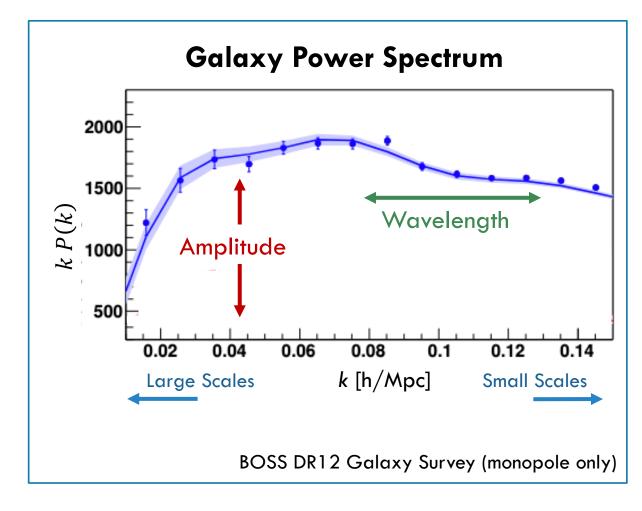
e.g. Ivanov, Philcox+21

Analyze the galaxy power spectrum using a scaling analysis

This measures:

- Overall **amplitude** ( = primordial amplitude)
- **Wiggle** positions ( = BAO feature)

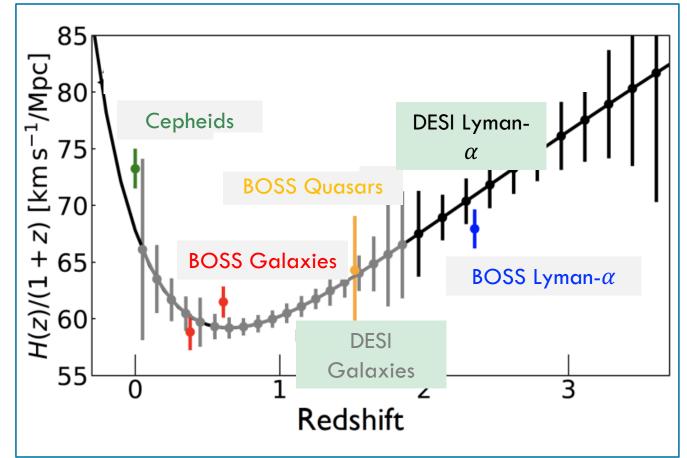
Robust way to constrain growth rate and expansion history H(z)



Analyze the galaxy power spectrum using a scaling analysis

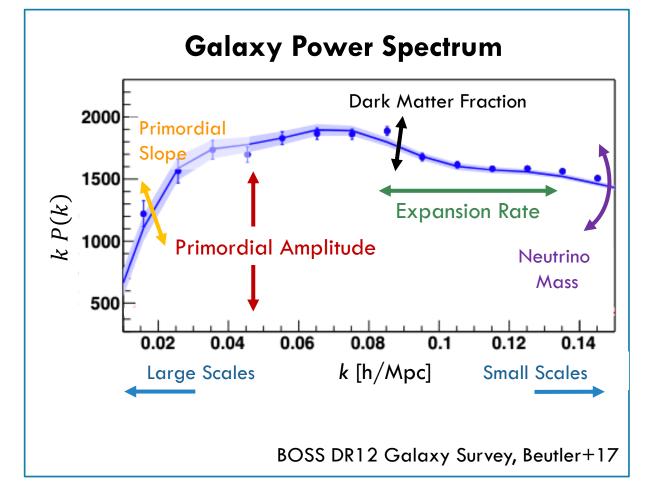
Measure wiggle positions ( = BAO feature) and overall amplitude

Robust way to constrain growth rate and expansion history H(z)



 $\triangleright$  This is not all the available information!

Measure parameters directly from the full shape of the galaxy power spectrum

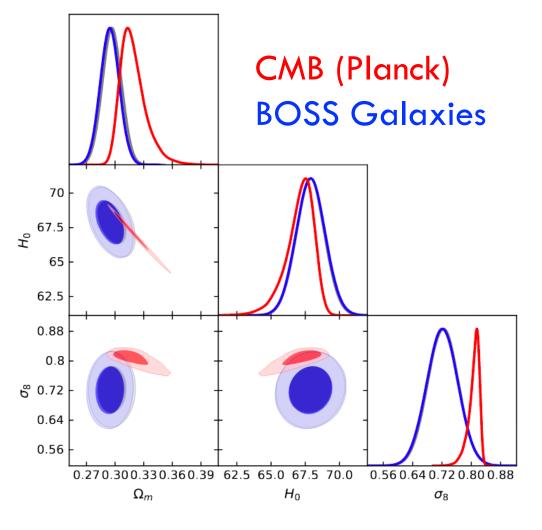


 $\triangleright$  This is not all the available information!

Measure parameters directly from the full shape of the galaxy power spectrum

Constrain parameters in new ways e.g. expansion rate from equality scale

Can we go **beyond** the power spectrum?



12 e.g. Ivanov+19,20, d'Amico+19,20, **Philcox**+20ab, Chen+21, Kobayashi+21

### WHY USE HIGHER-ORDER STATISTICS?

**Sharpen** parameter constraints!

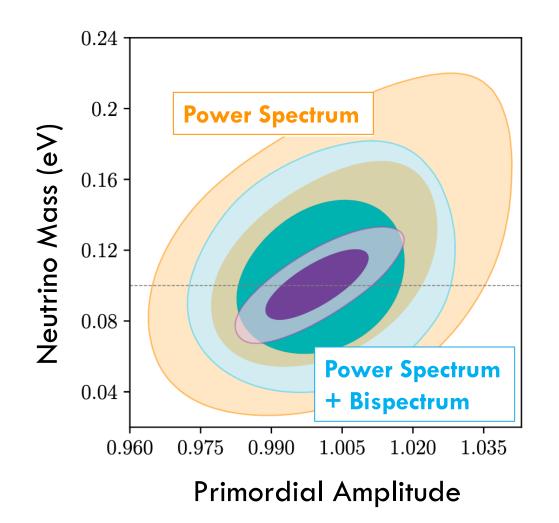
**Break** parameter **degeneracies**!

[e.g.  $P_g \sim b_1^2 \sigma_8^2$ ,  $B_g \sim b_1^3 \sigma_8^4$ ]

#### **Euclid Clustering Forecast**

 $\triangleright$  Bispectrum improves constraints by pprox 40%

 $\triangleright 1\sigma$  constraint of  $\sigma_{M_{\nu}} = 13$  meV [including Planck]



### **NON-GAUSSIAN INFLATION**

Are the primordial perturbations Gaussian?

### **Standard Model of Inflation:**

Scalar field  $\phi$  rolling down a potential  $V(\phi)$  $S = \int d^4x \sqrt{-g} \left[ \frac{1}{2}R - \frac{1}{2}\partial^{\mu}\phi\partial_{\mu}\phi - V(\phi) \right]$ Gravity Kinetic Energy Potential  $B_{\zeta}(\mathbf{k}_1, \mathbf{k}_2) \approx \frac{6}{5} f_{\mathrm{NL}} P_{\zeta}(k_1) P_{\zeta}(k_2) + 2 \text{ perms.}$  **Simplest Inflationary Model:** 

$$f_{\rm NL} \sim (1 - n_s) \ll 1$$

Non-standard inflation can beat this:
Multifield Inflation [Local Bispectrum]
New Kinetic Terms [Equilateral Bispectrum]
New Vacuum States [Folded Bispectrum]

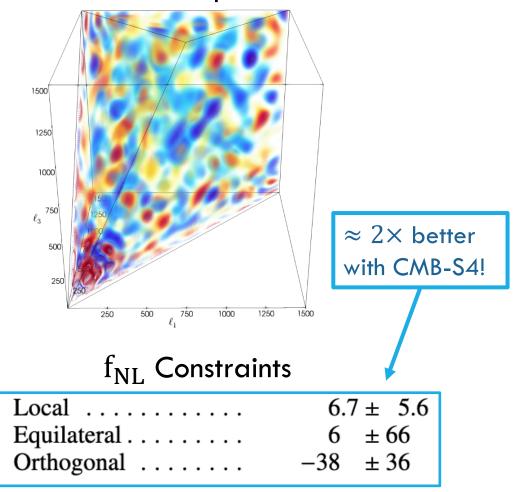
#### Planck 2018 IX

### **NON-GAUSSIAN INFLATION**

#### How do we measure this?

1. CMB Bispectrum

#### Planck TTT Bispectrum



 $B_{\zeta}(\mathbf{k}_1, \mathbf{k}_2) \approx \frac{6}{5} f_{\rm NL} P_{\zeta}(k_1) P_{\zeta}(k_2) + 2 \text{ perms.}$ 

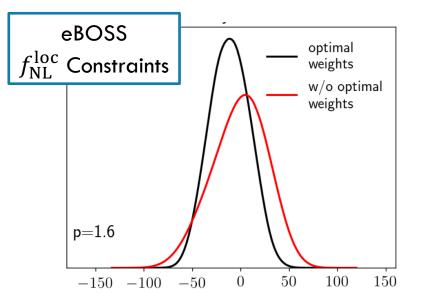
## **NON-GAUSSIAN INFLATION**

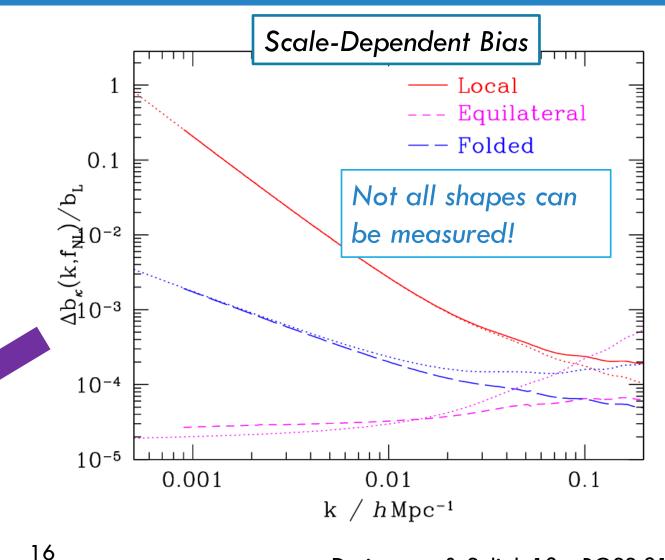
 $B_{\zeta}(\mathbf{k}_1, \mathbf{k}_2) \approx \frac{6}{5} f_{\rm NL} P_{\zeta}(k_1) P_{\zeta}(k_2) + 2 \text{ perms.}$ 

How do we measure this?

1. CMB Bispectrum

#### 2. Galaxy Power Spectrum

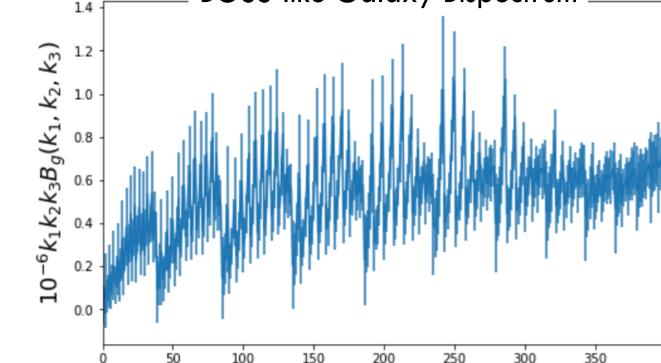




Desjacques & Seljak 10, eBOSS 21

### **BOSS-like Galaxy Bispectrum**

 $B_{\zeta}(\mathbf{k}_1, \mathbf{k}_2) \approx \frac{6}{5} f_{\mathrm{NL}} P_{\zeta}(k_1) P_{\zeta}(k_2) + 2 \text{ perms.}$ 



Triangles

### **NON-GAUSSIAN INFLATION**

How do we measure this?

- CMB Bispectrum
- 2. **Galaxy Power Spectrum**

**Galaxy Bispectrum** 3.

### **CHERN-SIMONS INTERACTIONS VIOLATE PARITY**

$$S = \int d^4x \sqrt{-g} \left[ \frac{1}{2}R - \frac{1}{2}\partial^\mu \phi \partial_\mu \phi - V(\phi) - \frac{1}{4}f(\phi)F_{\mu\nu}F^{\mu\nu} + \frac{\gamma}{4}f(\phi)F_{\mu\nu}\tilde{F}^{\mu\nu} \right]$$

▷ Add a **gauge field**  $A_{\mu}$  to the inflationary action, via  $F_{\mu\nu} = \partial_{\nu}A_{\mu} - \partial_{\mu}A_{\nu}$ 

 $\triangleright$  This can include a Chern-Simons coupling to the (pseudo-)scalar  $\phi$  [motivated by baryogenesis]

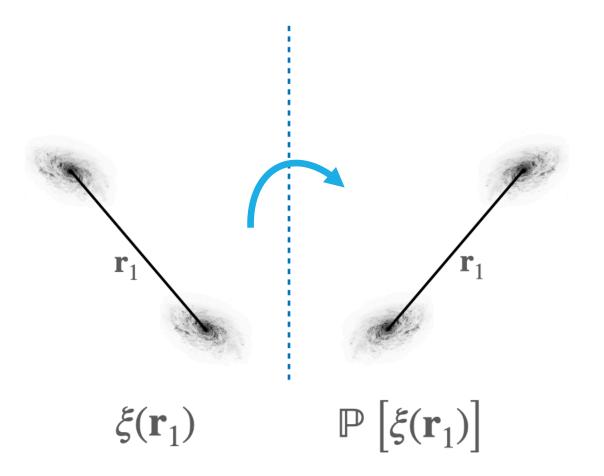
 $ightarrow f(\phi)F_{\mu\nu}\tilde{F}^{\mu\nu}$  violates **parity symmetry**  $\Rightarrow$  parity-violating correlators!

Where should we look for these signatures?

### THE 2PCF + 3PCF ARE NOT SENSITIVE TO PARITY

#### 2-Point Correlation Function (2PCF):

Parity Inversion = Rotation



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Lue+98, Jeong+12, Shiraishi 16, Cahn+21, **Philcox** (in prep.)

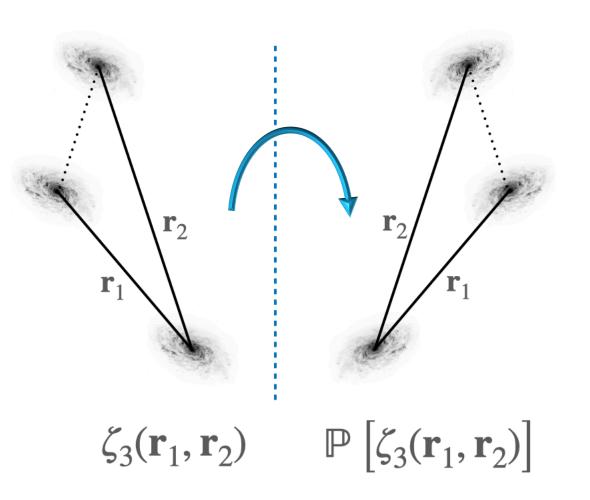
### THE 2PCF + 3PCF ARE NOT SENSITIVE TO PARITY

**2-Point Correlation Function (2PCF):** 

Parity Inversion = Rotation

#### **3-Point Correlation Function (3PCF):**

Parity Inversion = Rotation



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Lue+98, Jeong+12, Shiraishi 16, Cahn+21, Philcox (in prep.)

## THE 2PCF + 3PCF ARE NOT SENSITIVE TO PARITY

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#### 2-Point Correlation Function (2PCF):

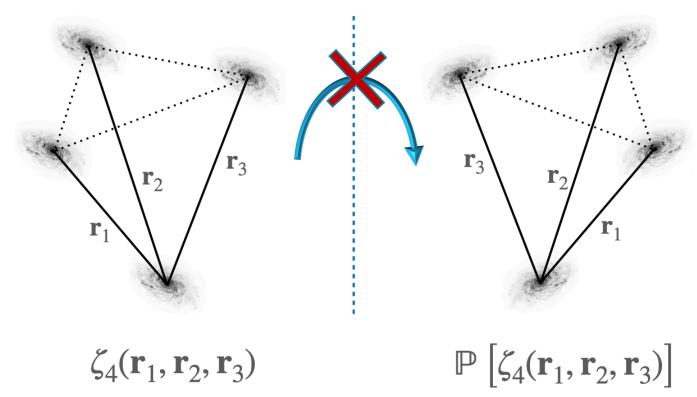
Parity Inversion = Rotation

#### **3-Point Correlation Function (3PCF):**

Parity Inversion = Rotation

#### **4-Point Correlation Function (4PCF):**

Parity Inversion  $\neq$  Rotation



Lue+98, Jeong+12, Shiraishi 16, Cahn+21, **Philcox** (in prep.)

### WHY USE HIGHER-ORDER STATISTICS?

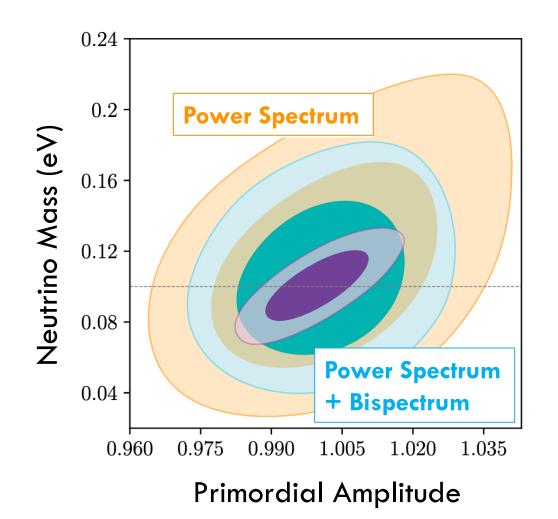
**Sharpen** parameter constraints!

**Break** parameter **degeneracies**!

Test non-standard physics models!

#### Why Use Large Scale Structure?

- Signal-to-Noise is **cubic** in number of modes unlike CMB
- New physics constraints **don't** dilute with redshift



### HOW TO MEASURE A BISPECTRUM

$$\hat{B}_{g}(k_{1},k_{2},k_{3}) = \int_{\mathbf{k}_{1},\mathbf{k}_{2},\mathbf{k}_{3}\in\text{bins}} \delta_{g}(\mathbf{k}_{1})\delta_{g}(\mathbf{k}_{2})\delta_{g}(\mathbf{k}_{3})(2\pi)^{3}\delta_{\text{D}}\left(\mathbf{k}_{1}+\mathbf{k}_{2}+\mathbf{k}_{3}\right)$$

**Problem:** We don't measure the density field directly.

$$\delta_g(\mathbf{r}) \to W(\mathbf{r})\delta_g(\mathbf{r}) \qquad \delta_g(\mathbf{k}) \to \int \frac{d\mathbf{p}}{(2\pi)^3} W(\mathbf{k} - \mathbf{p})\delta_g(\mathbf{p})$$
  
Window Function

The measured bispectrum is a triple convolution

$$B_g(\mathbf{k}_1, \mathbf{k}_2) \to \int_{\mathbf{p}_1 \mathbf{p}_2} W(\mathbf{k}_1 - \mathbf{p}_1) W(\mathbf{k}_2 - \mathbf{p}_2) W(\mathbf{p}_1 + \mathbf{p}_2 - \mathbf{k}_1 - \mathbf{k}_2) B_g(\mathbf{p}_1, \mathbf{p}_2)$$

Solution: Convolve the theory model too

#### Survey Window Function



### **CONVOLUTION IS EXPENSIVE**

$$B_g^{\min}(\mathbf{k}_1, \mathbf{k}_2) = \int_{\mathbf{p}_1 \mathbf{p}_2} W(\mathbf{k}_1 - \mathbf{p}_1) W(\mathbf{k}_2 - \mathbf{p}_2) W(\mathbf{p}_1 + \mathbf{p}_2 - \mathbf{k}_1 - \mathbf{k}_2) B_g(\mathbf{p}_1, \mathbf{p}_2)$$

> Window convolution is too costly to do repeatedly!

Common approximation: apply the window **only** to the power spectrum

 $B_g(\mathbf{k}_1, \mathbf{k}_2) \supset P_L(k_1) P_L(k_2)$ 

But:

- This gives systematic errors on large scales
- Spectra cannot be used to search for new physics!

### **BISPECTRA WITHOUT WINDOWS**

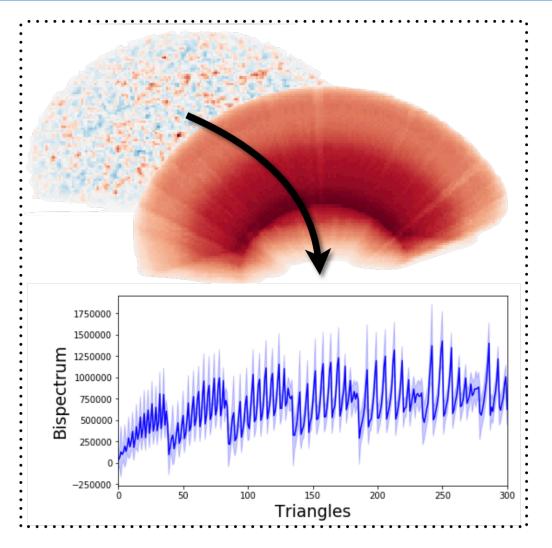
## Alternatively: estimate the unwindowed bispectrum directly

 $B_g^{\min}(\mathbf{k}_1, \mathbf{k}_2) = \int_{\mathbf{p}_1 \mathbf{p}_2} W(\mathbf{k}_1 - \mathbf{p}_1) W(\mathbf{k}_2 - \mathbf{p}_2) W(\mathbf{p}_1 + \mathbf{p}_2 - \mathbf{k}_1 - \mathbf{k}_2) B_g(\mathbf{p}_1, \mathbf{p}_2)$ 

Derive a maximum-likelihood estimator for the true bispectrum

Effectively **deconvolves** the window

$$\nabla_{B_g} L[\text{data}|B_g] = 0 \quad \Rightarrow \quad \widehat{B}_g = \cdots$$

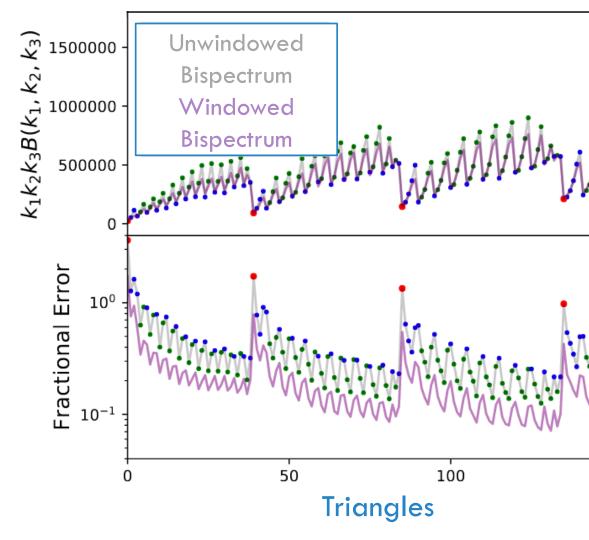


## **BISPECTRA WITHOUT WINDOWS**

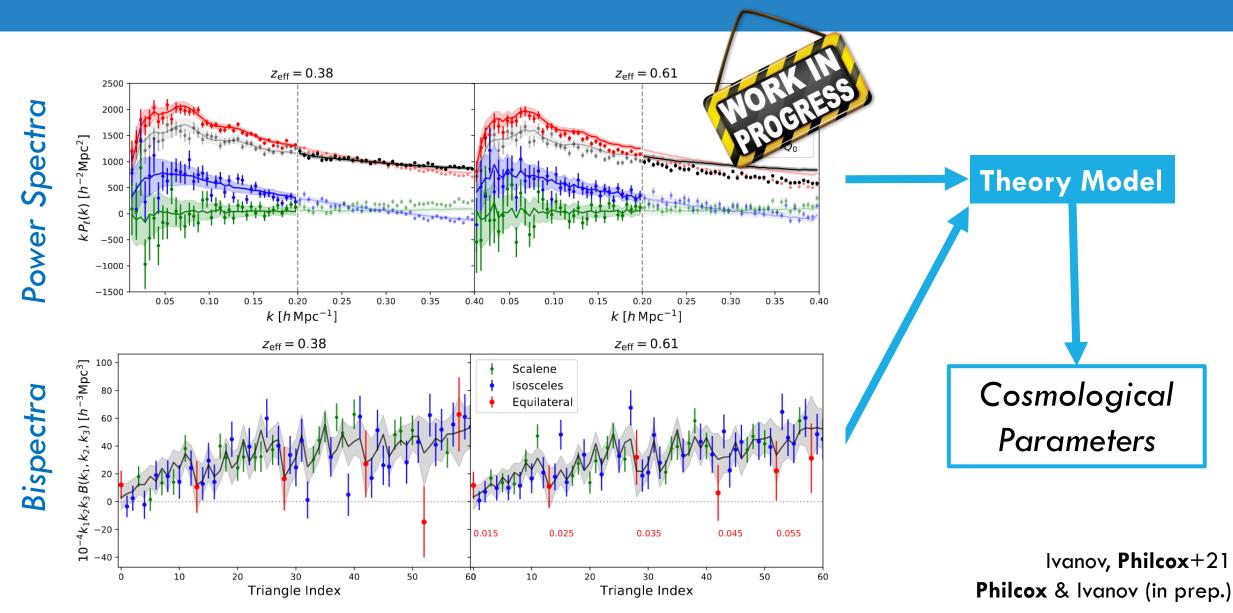
#### Properties of the **cubic estimator**:

- 1. Unbiased
- 2. Minimum variance [as  $B(k_1, k_2, k_3) \rightarrow 0$ ]
- 3. Window-free [effectively a deconvolution]

Requires various tricks for dealing with high-dimensional data [e.g. conjugate gradient descent, Monte Carlo estimation etc.]



### **BOSS WITHOUT WINDOWS**



### WHAT WILL WE MEASURE?

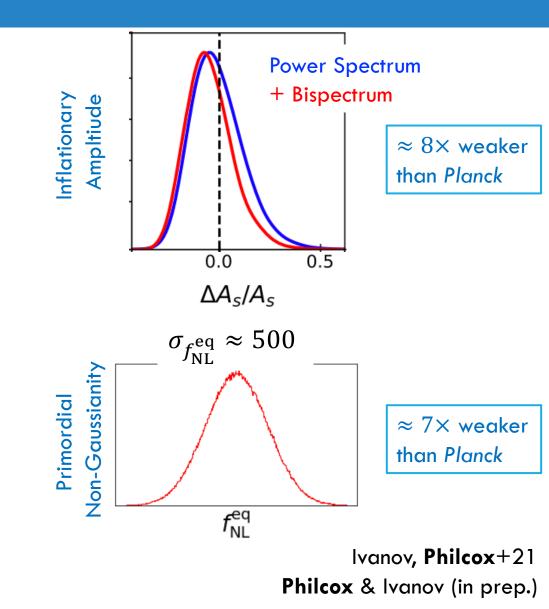
Tighter constraints on cosmological and galaxy formation parameters

 $ightarrow \sigma_8$  improves by 10%

▷ Tidal bias improves by 50%

Bounds on all flavors of Primordial Non-Gaussianity

First equilateral-type measurement from LSS



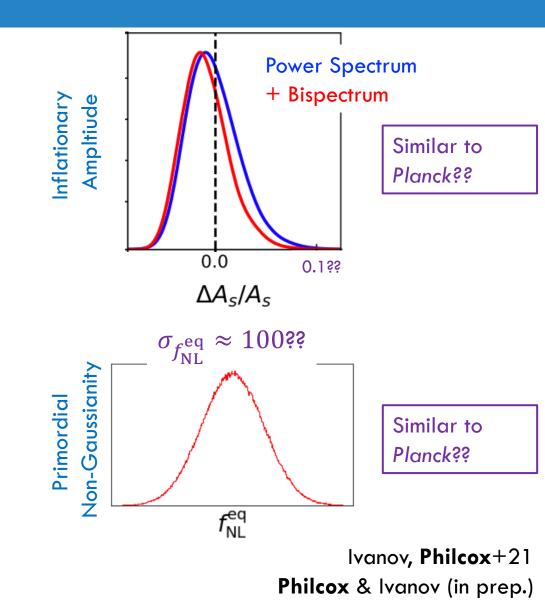
### WHAT'S NEXT FOR BISPECTRA?

- Improve bispectrum modeling
- ▷ Higher-order perturbation theory
- > Add **redshift-space** information
- Better treatment of fingers-of-God
- > Apply to **DESI** data
  - Pipelines already available and tested

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 $\triangleright$  Expect  $O(5) \times$  stronger constraints

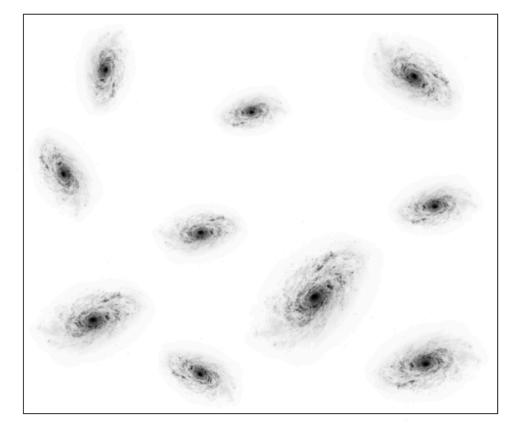




### CORRELATION FUNCTIONS = GALAXY COUNTS

$$\hat{\zeta}_{g,4}(r_1, r_2, r_3, \hat{\mathbf{r}}_1 \cdot \hat{\mathbf{r}}_2, \hat{\mathbf{r}}_1 \cdot \hat{\mathbf{r}}_3, \hat{\mathbf{r}}_2 \cdot \hat{\mathbf{r}}_3) = \int \frac{d\mathbf{x}}{V} \int_{\mathbf{r}_1, \mathbf{r}_2, \mathbf{r}_3 \in \text{bins}} \delta_g(\mathbf{x}) \delta_g(\mathbf{x} + \mathbf{r}_1) \delta_g(\mathbf{x} + \mathbf{r}_2) \delta_g(\mathbf{x} + \mathbf{r}_3)$$

Measuring the 4PCF involves counting **quadruplets** of galaxies



### CORRELATION FUNCTIONS = GALAXY COUNTS

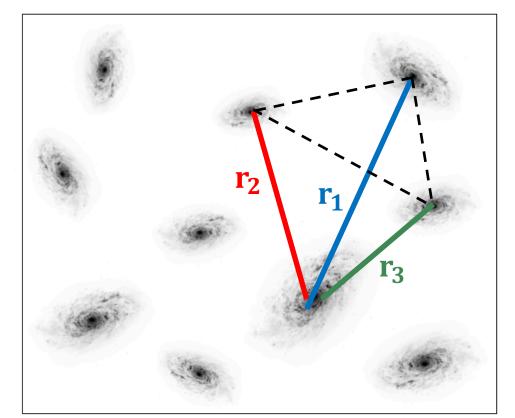
$$\hat{\zeta}_{g,4}(r_1, r_2, r_3, \hat{\mathbf{r}}_1 \cdot \hat{\mathbf{r}}_2, \hat{\mathbf{r}}_1 \cdot \hat{\mathbf{r}}_3, \hat{\mathbf{r}}_2 \cdot \hat{\mathbf{r}}_3) = \int \frac{d\mathbf{x}}{V} \int_{\mathbf{r}_1, \mathbf{r}_2, \mathbf{r}_3 \in \text{bins}} \delta_g(\mathbf{x}) \delta_g(\mathbf{x} + \mathbf{r}_1) \delta_g(\mathbf{x} + \mathbf{r}_2) \delta_g(\mathbf{x} + \mathbf{r}_3)$$

Measuring the 4PCF involves counting **quadruplets** of galaxies

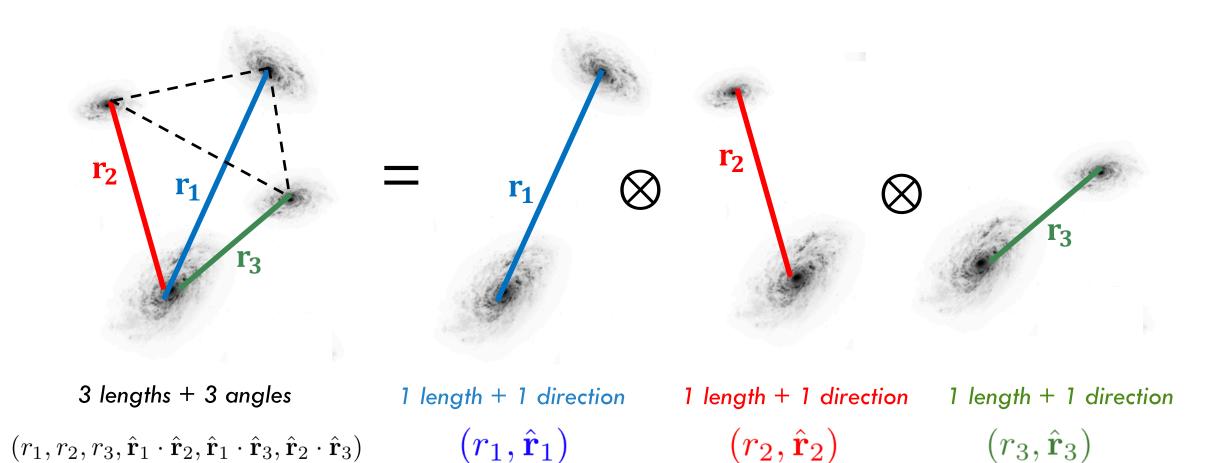
Total number of quadruplets:

 $O(N_g^4)$ 

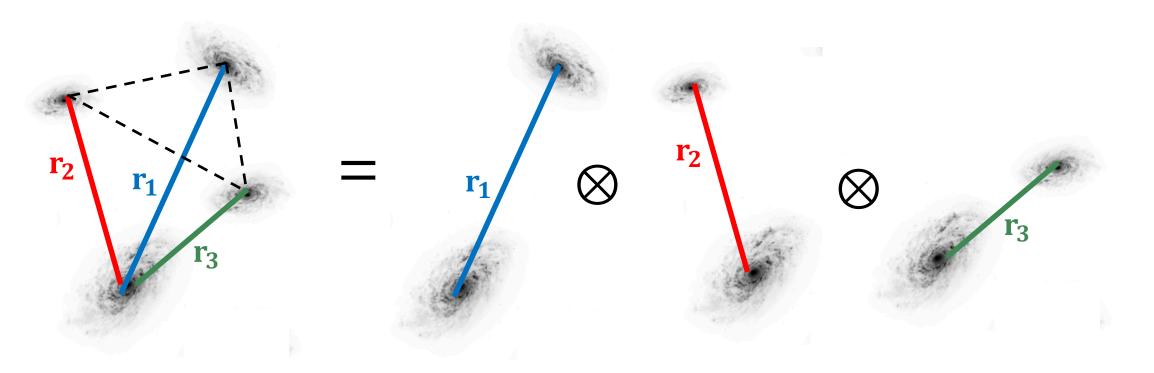
This is too many to count...



### ONE TETRAHEDRON = THREE VECTORS



### ONE TETRAHEDRON = THREE VECTORS



3 lengths + 3 multipoles $(r_1, r_2, r_3, \ell_1, \ell_2, \ell_3)$ 

1 length + 2 multipoles1 length + 2 multipoles1 length + 2 multipoles $(r_1, \ell_1, m_1)$  $(r_2, \ell_2, m_2)$  $(r_3, \ell_3, m_3)$ 

### ANGULAR MOMENTUM BASIS

Expand 4PCF in basis of **isotropic functions** 

$$\zeta_4(\mathbf{r}_1, \mathbf{r}_2, \mathbf{r}_3) = \sum_{\ell_1 \ell_2 \ell_3} \zeta_{\ell_1 \ell_2 \ell_3}(r_1, r_2, r_3) \mathcal{P}_{\ell_1 \ell_2 \ell_3}(\hat{\mathbf{r}}_1, \hat{\mathbf{r}}_2, \hat{\mathbf{r}}_3)$$

$$\uparrow$$
Coefficients

Basis formed from **angular momentum addition** in 3D

$$\mathcal{P}_{\ell_1\ell_2\ell_3}(\hat{\mathbf{r}}_1, \hat{\mathbf{r}}_2, \hat{\mathbf{r}}_3) = \sum_{m_1m_2m_3} \begin{pmatrix} \ell_1 & \ell_2 & \ell_3 \\ m_1 & m_2 & m_3 \end{pmatrix} Y^*_{\ell_1m_1}(\hat{\mathbf{r}}_1) Y^*_{\ell_2m_2}(\hat{\mathbf{r}}_2) Y^*_{\ell_3m_3}(\hat{\mathbf{r}}_3)$$

This is separable in  $\widehat{r}_1,\,\widehat{r}_2,\,\widehat{r}_3$ 

 $\mathbf{r}_1$ 

 $\otimes$ 

### A SEPARABLE BASIS $\Rightarrow$ A QUADRATIC ESTIMATOR

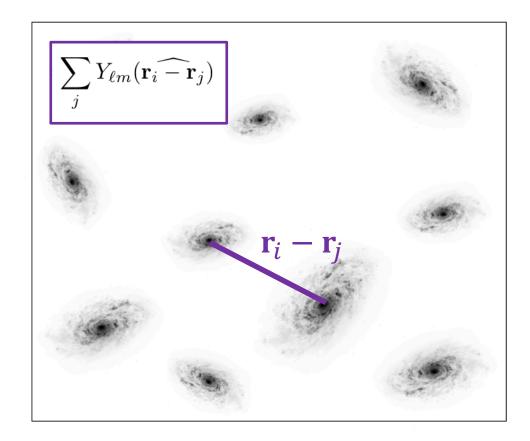
$$\hat{\zeta}_{\ell_1\ell_2\ell_3}(r_1, r_2, r_3) = \sum_{m_1m_2m_3} \begin{pmatrix} \ell_1 & \ell_2 & \ell_3 \\ m_1 & m_2 & m_3 \end{pmatrix} \int d\mathbf{x} \, \delta_g(\mathbf{x}) \left[ \int_{\mathbf{r}_1} \delta_g(\mathbf{x} + \mathbf{r}_1) Y_{\ell_1m_1}(\hat{\mathbf{r}}_1) \right] \left[ \int_{\mathbf{r}_2} \delta_g(\mathbf{x} + \mathbf{r}_2) Y_{\ell_2m_2}(\hat{\mathbf{r}}_2) \right] \left[ \int_{\mathbf{r}_3} \delta_g(\mathbf{x} + \mathbf{r}_3) Y_{\ell_3m_3}(\hat{\mathbf{r}}_3) \right] d\mathbf{x} \, \delta_g(\mathbf{x} + \mathbf{r}_3) \left[ \int_{\mathbf{r}_3} \delta_g(\mathbf{x} + \mathbf{r}_3) Y_{\ell_3m_3}(\hat{\mathbf{r}}_3) \right] d\mathbf{x} \, \delta_g(\mathbf{x} + \mathbf{r}_3) \left[ \int_{\mathbf{r}_3} \delta_g(\mathbf{x} + \mathbf{r}_3) Y_{\ell_3m_3}(\hat{\mathbf{r}}_3) \right] d\mathbf{x} \, \delta_g(\mathbf{x} + \mathbf{r}_3) \left[ \int_{\mathbf{r}_3} \delta_g(\mathbf{x} + \mathbf{r}_3) Y_{\ell_3m_3}(\hat{\mathbf{r}}_3) \right] d\mathbf{x} \, \delta_g(\mathbf{x} + \mathbf{r}_3) \left[ \int_{\mathbf{r}_3} \delta_g(\mathbf{x} + \mathbf{r}_3) Y_{\ell_3m_3}(\hat{\mathbf{r}}_3) \right] d\mathbf{x} \, \delta_g(\mathbf{x} + \mathbf{r}_3) \left[ \int_{\mathbf{r}_3} \delta_g(\mathbf{x} + \mathbf{r}_3) Y_{\ell_3m_3}(\hat{\mathbf{r}}_3) \right] d\mathbf{x} \, \delta_g(\mathbf{x} + \mathbf{r}_3) \left[ \int_{\mathbf{r}_3} \delta_g(\mathbf{x} + \mathbf{r}_3) Y_{\ell_3m_3}(\hat{\mathbf{r}}_3) \right] d\mathbf{x} \, \delta_g(\mathbf{x} + \mathbf{r}_3) \left[ \int_{\mathbf{r}_3} \delta_g(\mathbf{x} + \mathbf{r}_3) Y_{\ell_3m_3}(\hat{\mathbf{r}}_3) \right] d\mathbf{x} \, \delta_g(\mathbf{x} + \mathbf{r}_3) \left[ \int_{\mathbf{r}_3} \delta_g(\mathbf{x} + \mathbf{r}_3) Y_{\ell_3m_3}(\hat{\mathbf{r}}_3) \right] d\mathbf{x} \, \delta_g(\mathbf{x} + \mathbf{r}_3) \left[ \int_{\mathbf{r}_3} \delta_g(\mathbf{x} + \mathbf{r}_3) Y_{\ell_3m_3}(\hat{\mathbf{r}}_3) \right] d\mathbf{x} \, \delta_g(\mathbf{x} + \mathbf{r}_3) \left[ \int_{\mathbf{r}_3} \delta_g(\mathbf{x} + \mathbf{r}_3) Y_{\ell_3m_3}(\hat{\mathbf{r}}_3) \right] d\mathbf{x} \, \delta_g(\mathbf{x} + \mathbf{r}_3) \left[ \int_{\mathbf{r}_3} \delta_g(\mathbf{x} + \mathbf{r}_3) Y_{\ell_3m_3}(\hat{\mathbf{r}}_3) \right] d\mathbf{x} \, \delta_g(\mathbf{x} + \mathbf{r}_3) \left[ \int_{\mathbf{r}_3} \delta_g(\mathbf{x} + \mathbf{r}_3) Y_{\ell_3m_3}(\hat{\mathbf{r}}_3) \right] d\mathbf{x} \, \delta_g(\mathbf{x} + \mathbf{r}_3) \left[ \int_{\mathbf{r}_3} \delta_g(\mathbf{x} + \mathbf{r}_3) Y_{\ell_3m_3}(\hat{\mathbf{r}}_3) \right] d\mathbf{x} \, \delta_g(\mathbf{x} + \mathbf{r}_3) \left[ \int_{\mathbf{r}_3} \delta_g(\mathbf{x} + \mathbf{r}_3) Y_{\ell_3m_3}(\hat{\mathbf{r}}_3) \right] d\mathbf{x} \, \delta_g(\mathbf{x} + \mathbf{r}_3) \left[ \int_{\mathbf{r}_3} \delta_g(\mathbf{x} + \mathbf{r}_3) Y_{\ell_3m_3}(\hat{\mathbf{r}}_3) \right] d\mathbf{x} \, \delta_g(\mathbf{x} + \mathbf{r}_3) \left[ \int_{\mathbf{r}_3} \delta_g(\mathbf{x} + \mathbf{r}_3) Y_{\ell_3m_3}(\hat{\mathbf{r}}_3) \right] d\mathbf{x} \, \delta_g(\mathbf{x} + \mathbf{r}_3) \left[ \int_{\mathbf{r}_3} \delta_g(\mathbf{x} + \mathbf{r}_3) Y_{\ell_3m_3}(\hat{\mathbf{r}}_3) \right] d\mathbf{x} \, \delta_g(\mathbf{x} + \mathbf{r}_3) \left[ \int_{\mathbf{r}_3} \delta_g(\mathbf{x} + \mathbf{r}_3) Y_{\ell_3m_3}(\hat{\mathbf{r}}_3) \right] d\mathbf{x} \, \delta_g(\mathbf{x} + \mathbf{r}_3) \left[ \int_{\mathbf{r}_3} \delta_g(\mathbf{x} + \mathbf{r}_3) Y_{\ell_3m_3}(\hat{\mathbf{r}}_3) \right] d\mathbf{x} \, \delta_g(\mathbf{x} + \mathbf{r}_3) \left[ \int_{\mathbf{r}_3} \delta_g(\mathbf{x} + \mathbf{r}_3) Y_{\ell_3m_3}(\hat{\mathbf{r}}_3) \right] d\mathbf{x} \, \delta_g(\mathbf{x} + \mathbf{r}_3) \left[ \int_{\mathbf{r}_3} \delta_g(\mathbf{x} + \mathbf{r}_3) Y_{\ell$$

The estimator **factorizes** into **independent** pieces

To compute the 4PCF: count pairs of galaxies

Total number of pairs:  $\mathcal{O}ig(N_{\mathrm{g}}^2ig)$ 

This can be computed!



## **ENCORE**: ULTRA-FAST N-POINT FUNCTIONS

#### Public C++/CUDA code

Corrects for **survey geometry** 

Requires ~ 10 CPU-hours to compute 4PCF of current data

Extends to curved geometries, higher-dimensions, 5PCF, 6PCF and beyond...

### oliverphilcox/ encore



encore: Efficient isotropic 2-, 3-, 4-, 5- and 6-point correlation functions in C++ and CUDA

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Contributors	Issues	Stars	Fork	

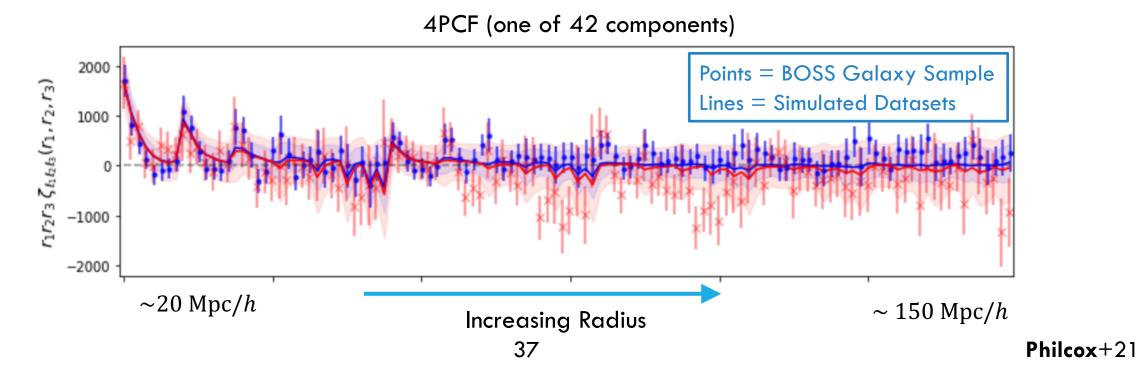
#### oliverphilcox/encore

encore: Efficient isotropic 2-, 3-, 4-, 5- and 6-point correlation functions in C++ and CUDA - oliverphilcox/encore  $\oslash$  github.com

### **MEASURING THE 4-POINT FUNCTION**

Compute the (connected) 4PCF from  $\sim 10^6$  BOSS galaxies

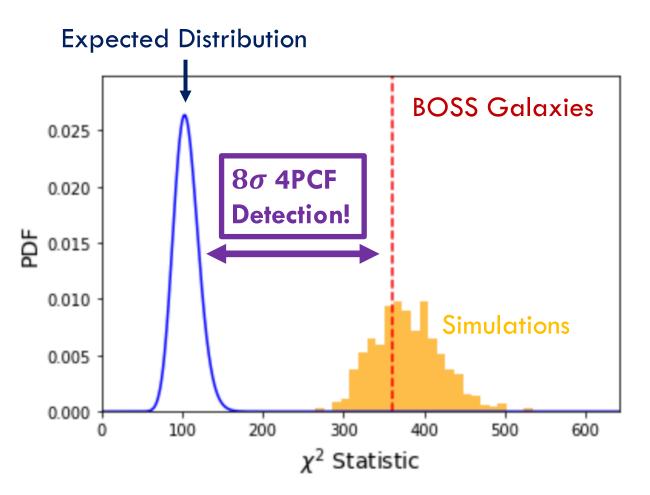
#### Do we detect a signal?



## CAN WE DETECT THE GRAVITATIONAL 4PCF?

> Perform a  $\chi^2$ -test to search for a gravitational 4PCF

- $\triangleright$  Null Hypothesis: **4PCF = 0**.
- **Strong** detection of non-Gaussianity!



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Philcox+21

### WHAT'S NEXT FOR THE 4-POINT FUNCTION?

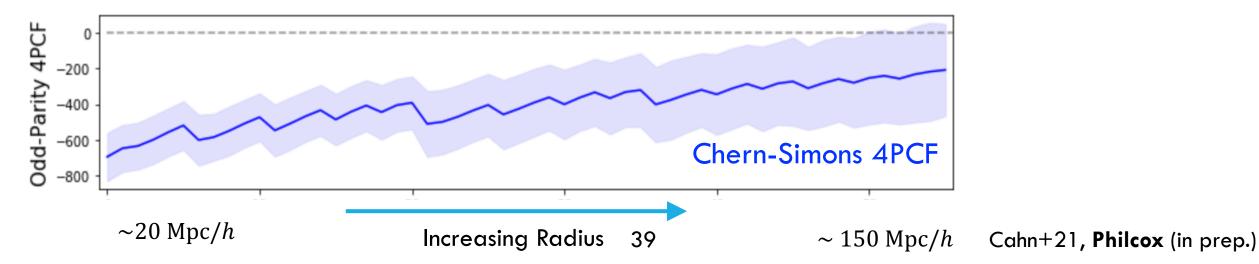
Create a theory model and quantify information content:

 $\triangleright$  Allows **ACDM information** to be extracted

Search for parity-violating physics in the BOSS 4PCF

$$\zeta_4(\mathbf{r}_1,\mathbf{r}_2,\mathbf{r}_3) - \mathbb{P}[\zeta_4(\mathbf{r}_1,\mathbf{r}_2,\mathbf{r}_3)]$$

 $\triangleright$  Apply to **DESI** data [2× higher precision] and combine with the **CMB** 



**crXiv** 2012.09389 2105.08722 2106.10278 2107.06287 2108.01670 2110.10161

### **Contact** ep2@cantab.ac @oliver\_philcox

# CONCLUSIONS

#### o Non-Gaussian statistics:

- 1. Sharpen cosmological constraints
- 2. Probe **non-standard** physics in the early Universe

• Fast and accurate estimators now available

Extract more information from LSS surveys
 without additional cost