

The Energy Minimisation Principle

Linking protohalo counts and shapes to anisotropic infall

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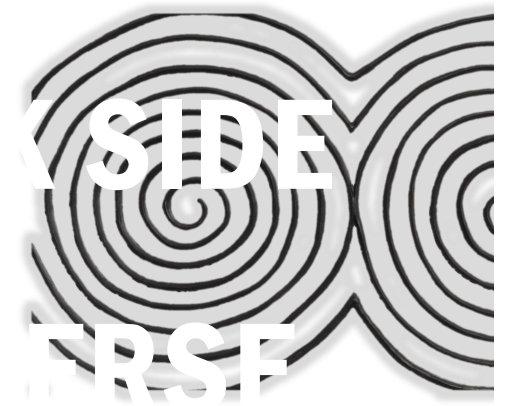


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Organizing DSU2023 in Kigali, Rwanda.
Hope to see you there!

<https://eaifr.ictp.it/events/dsu-2023/>



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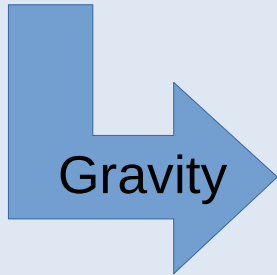
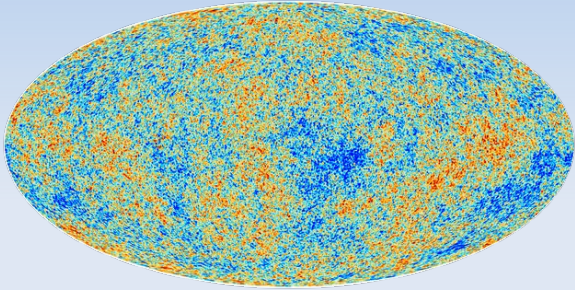
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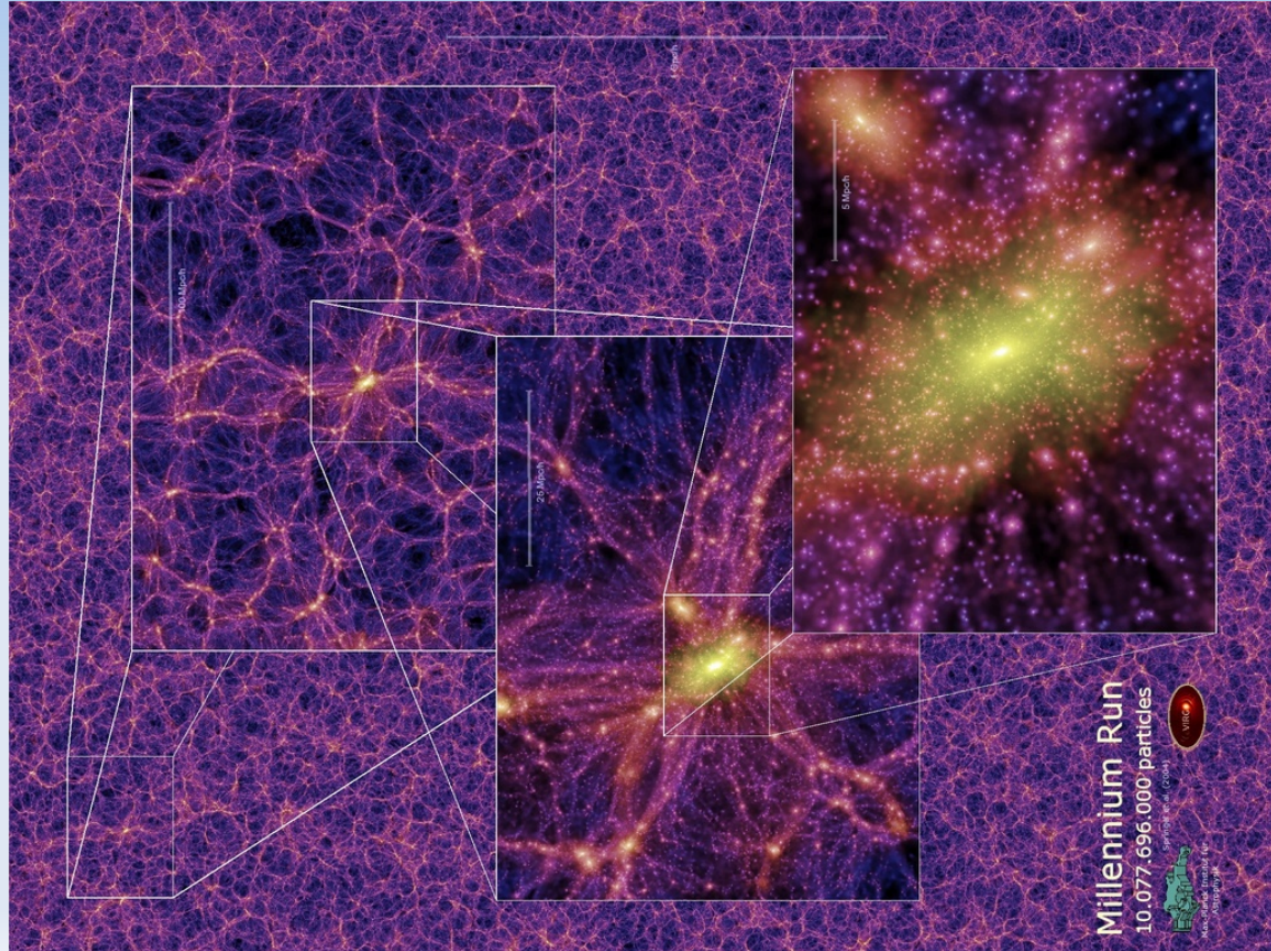
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The non-linear Universe

~Gaussian initial conditions



Filaments, nodes, voids...

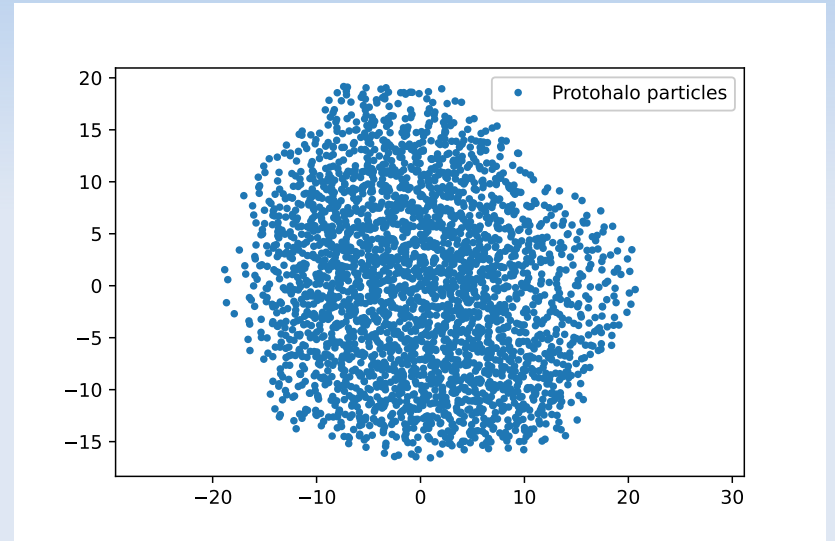
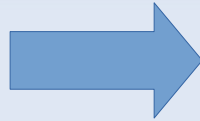
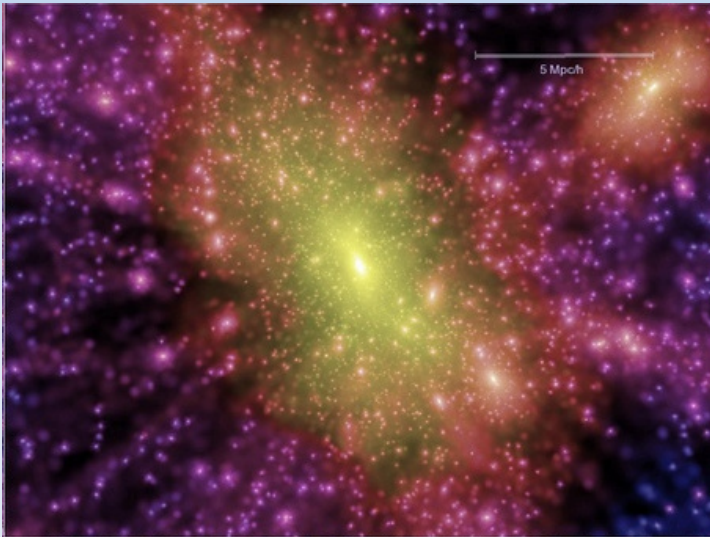


Why analytical models?

- Predict/interpret features of the outcome without running the simulation
- Conceptual grasp of the interplay of the many variables involved
- Useful to quickly scan the parameter space. Simulations are still costly!
- Explore beyond- Λ CDM scenarios (DE, ModGrav, axions...) that are difficult to simulate
- Physically motivated fitting formulae

The rules of the game

- Each halo has a protohalo (Lagrangian region occupied by the halo particles)

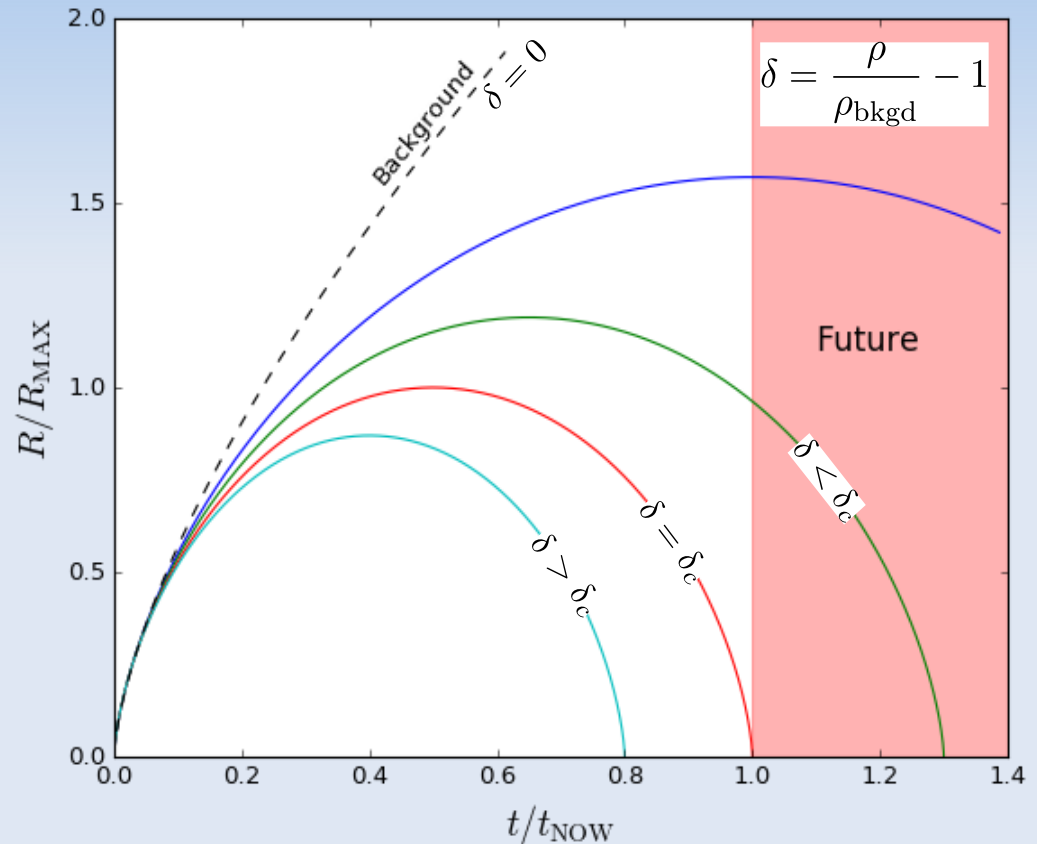


- Protohaloes live in a Gaussian field. Their stats are non-Gaussian, but can be predicted
- Need to identify them!

The critical density

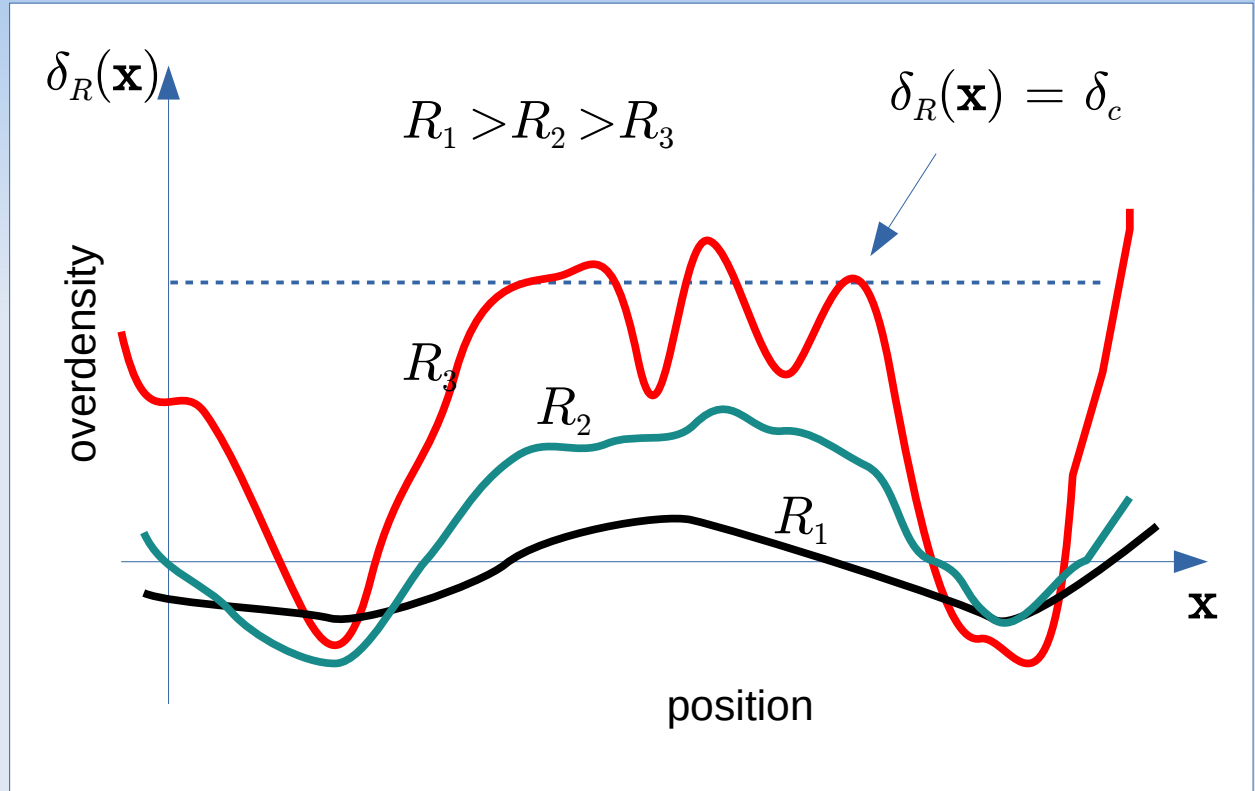
- Protohalos are approximately spherical, follow approx spherical collapse
- Time of collapse depends on **mean overdensity**: $tH_0 = \pi(3/5\delta_R)^{3/2}$
- Maximum (turnaround) radius:
 $r_{max} = 3R/5\delta_R$
- A shell of Lagrangian radius R collapses today if $\delta_R = \delta_c$

$$\delta_c \equiv \frac{3}{10} (18\pi^2)^{1/3} \simeq 1.686$$



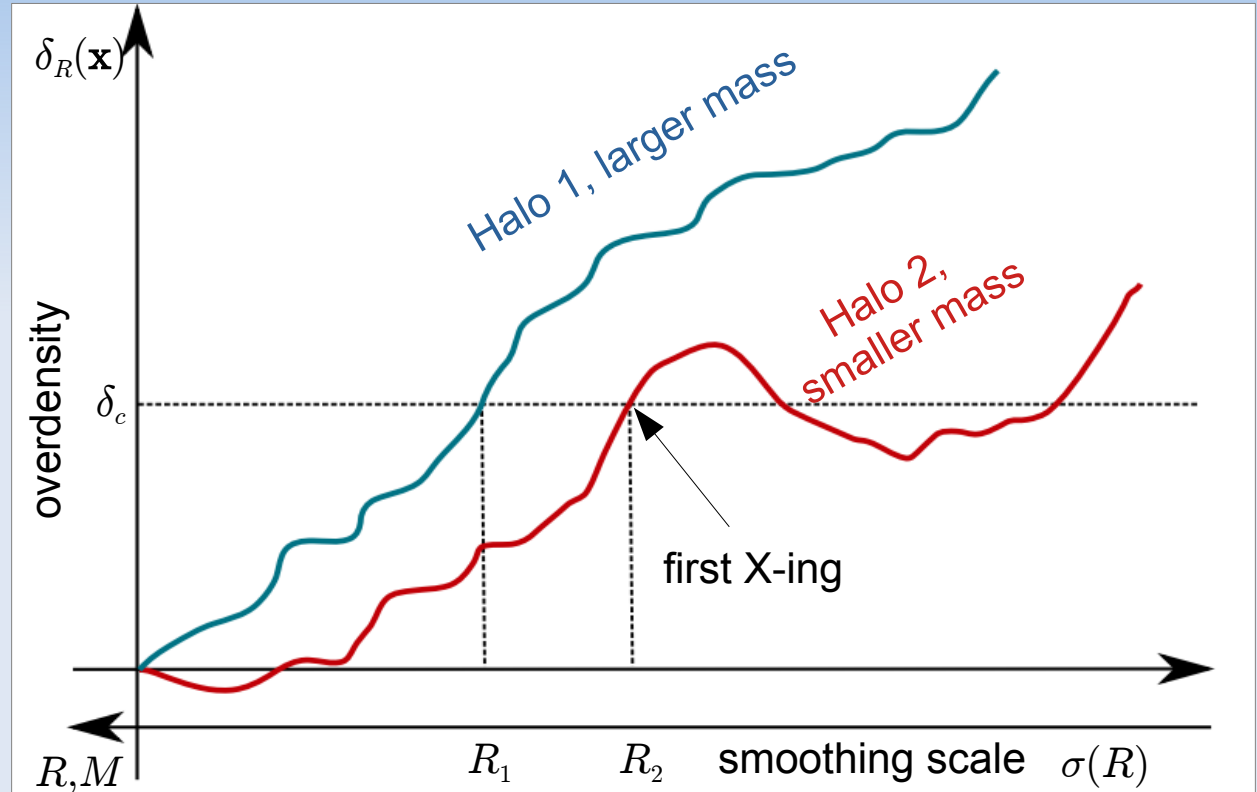
Density peaks

- Protohaloes collapse faster than neighboring regions
- They must be local maxima of the mean (=smoothed) density field
- Which smoothing scale?



Excursion sets

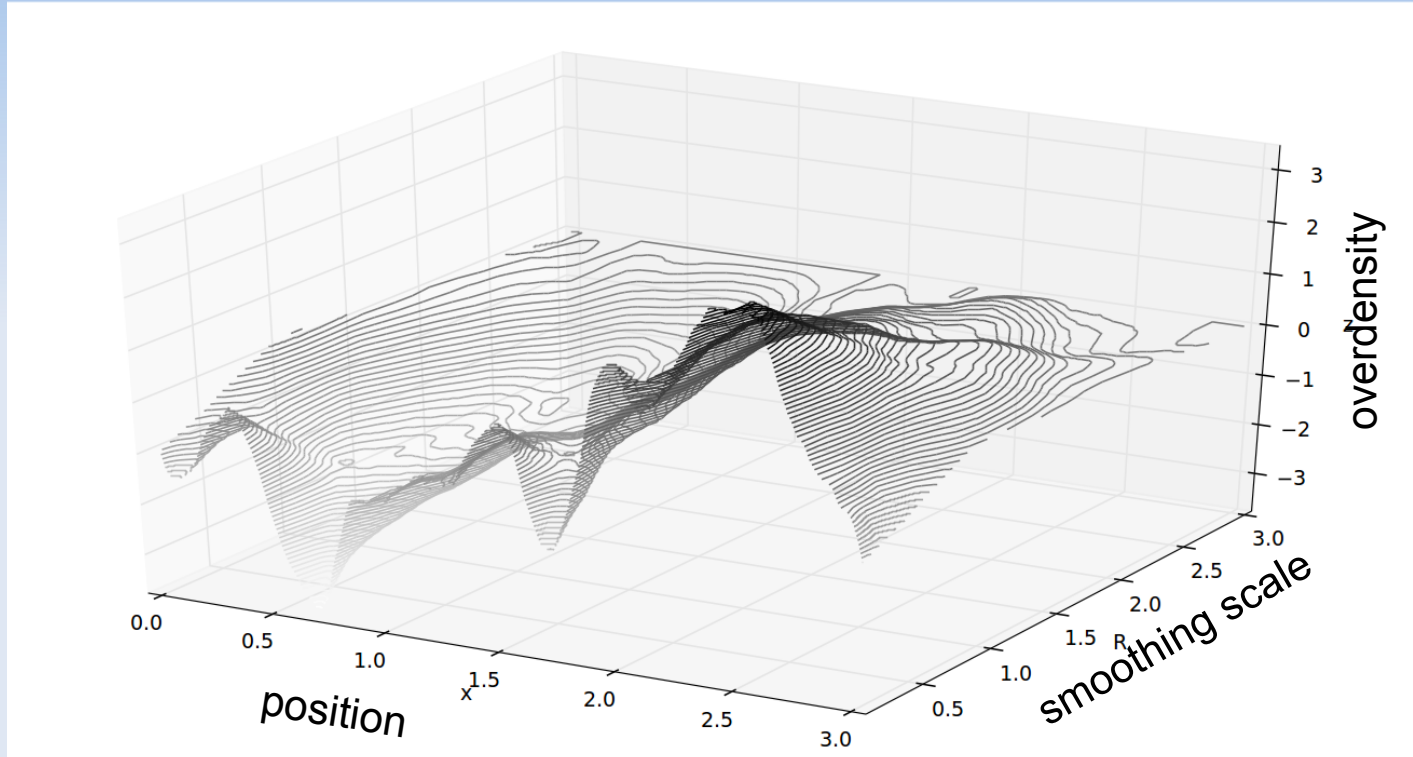
- At fixed \mathbf{x} , the value of $\delta_R(\mathbf{x})$ gives the collapse time of the patch
- Imposing $\delta_R(\mathbf{x}) = \delta_c$ fixes the smoothing scale R
- Look for the largest such R (“first crossing”)



Press & Schechter 74,
Bond et al 91,

Excursion sets + peaks

- Combine the two ideas
- After smoothing, 4-dimensional landscape in x and R
- Look for peaks that first cross the critical height
- Peak constraint fixes position x
- Threshold fixes the smoothing scale R



The standard game

- Protohalos are peaks of the (smoothed) initial density field
- Smooth the initial overdensity field $\delta_R(\mathbf{x})$ on ALL possible scales R
- **Peak constraint:** $\nabla\delta_R = 0 \rightarrow$ fixes the position \mathbf{x}
- Peaks of different height collapse at different times
- **Critical peak height:** $\delta_R = \delta_c \rightarrow$ fixes the smoothing scale R
- Mass is conserved. The mass of the final halo is $M = 4\pi\bar{\rho}R^3/3$

The good, the bad and the ugly

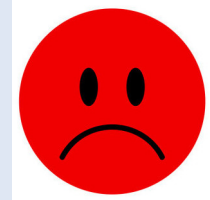
- The good. Simple and intuitive.
An ad-hoc stochastic threshold gives good (~5-10%) mass function AND bias (success for bias is non-trivial)

Paranjape, Sheth & Desjacques (2013); Castorina et al. (2016)

- The bad. Scatter in δ_c agrees with N-body only up to 30%
Only 75% of small-mass protohalos are peaks

e.g. Ludlow, Borszykowski, Porciani (2014)

- The ugly. Top-Hat filter gives divergences in Λ CDM
No connection with perturbation theory
No information on the shape
Why is δ_c stochastic?



1. Excursion set peaks in energy

Or: getting the position right

(based on 1907.09147 with Ravi Sheth)

Matter vs energy peaks

- Geometrical radius: $R^3 = 3V/4\pi$

- $\frac{\dot{R}^2}{2} - \frac{GM}{R} - \frac{\Lambda}{6}R^2 \simeq -\frac{5}{3}\frac{GM}{R_{\text{in}}}\delta_{R,\text{in}}$

- Mass: $M = \frac{4\pi}{3}\bar{\rho}(1 + \delta_R)R^3$

vs

- Inertial radius: $\frac{R_I^2}{5} \equiv \int_V \frac{d^3r}{3M}\rho |\mathbf{r} - \mathbf{r}_{\text{cm}}|^2$

- $\frac{\dot{R}_I^2}{2} - \frac{GM_I}{R} - \frac{\Lambda}{6}R_I^2 \simeq -\frac{5}{3}\frac{GM_I}{R_{I,\text{in}}}\epsilon_{\text{in}}$

- Inertial mass: $M_I = \frac{4\pi}{3}\bar{\rho}(1 + \epsilon_R)R_I^3$

- Governed by **matter** overdensity

$$\delta_R \equiv \frac{1}{V} \int_V d^3r \delta(\mathbf{r})$$

- Governed by **energy** overdensity

$$\epsilon_R \equiv 5 \int_V \frac{d^3r}{MR_I^2}\rho(\mathbf{r}) \mathbf{r} \cdot (\nabla\phi - \nabla\phi_{\text{cm}})$$

Matter vs energy peaks

- Characteristic time $\sim (1/\delta_R)^{3/2}$
- Halos of mass M are peaks of $\delta_R(\mathbf{x})$
- In Fourier space:

$$\delta_R = \int \frac{d^3k}{(2\pi)^3} \delta(\mathbf{k}) \frac{3j_1(kR)}{kR}$$

VS

- Characteristic time $\sim (1/\epsilon_R)^{3/2}$
- Halos of mass M are peaks of $\epsilon(\mathbf{x})$
- In Fourier space:

$$\epsilon_R = \int \frac{d^3k}{(2\pi)^3} \delta(\mathbf{k}) \frac{15j_2(kR)}{(kR)^2}$$

(extra power of $1/k$)

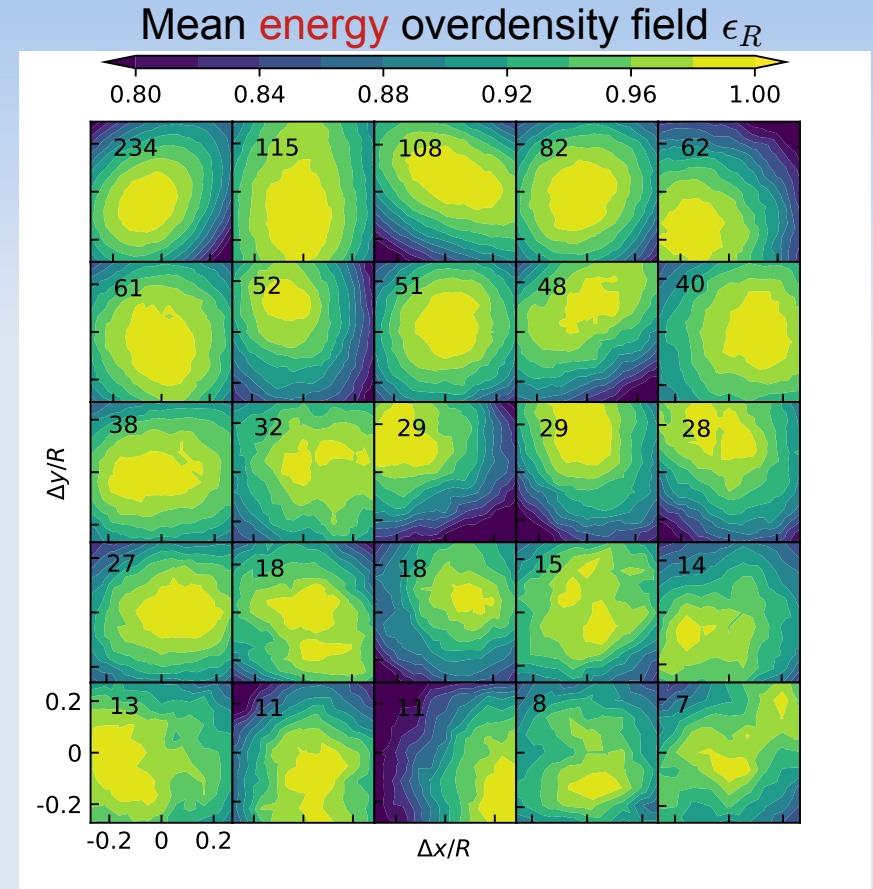
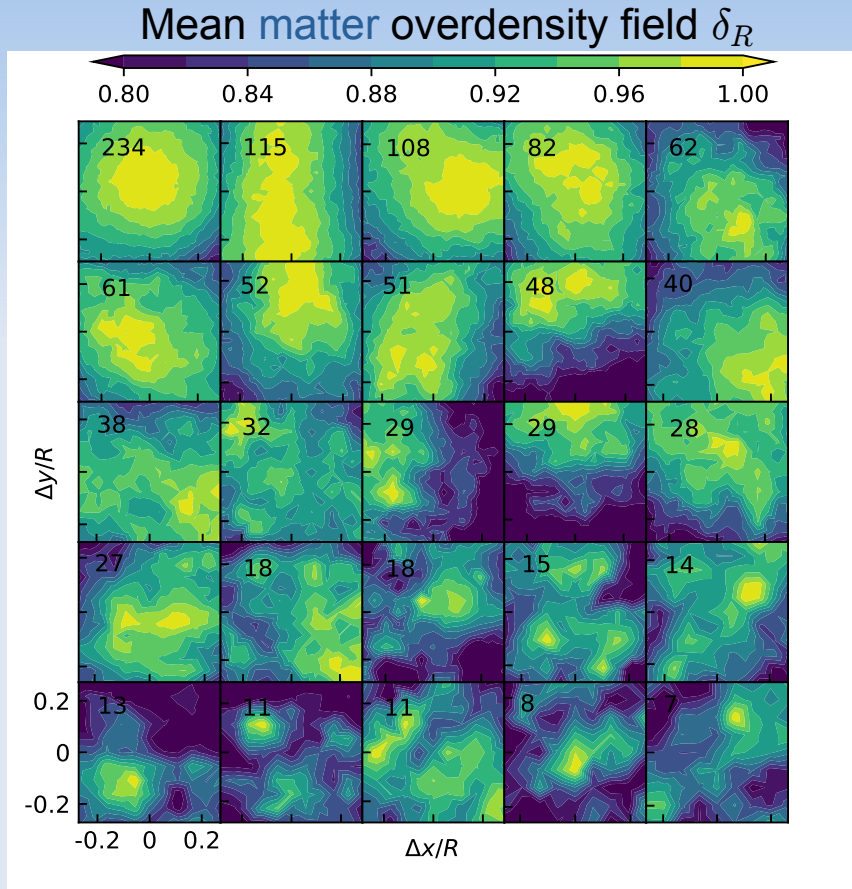
What is the advantage?

- R is very sensitive to the halo boundary
- No dynamical meaning in $\nabla\delta_R = 0$
- More small-scale power.
 $\langle(\nabla^2\delta_R)^2\rangle$ diverges in Λ CDM.
- Usually resort to Gaussian filter.
Blurred physical interpretation

vs

- R_I is density weighted, less sensitive to halo boundary at late times
- $\nabla\epsilon_R \sim$ dipole moment.
 $\nabla\epsilon_R = 0$ implies radial infall
- Less small-scale power.
 $\langle(\nabla^2\epsilon_R)^2\rangle$ remains finite.
- No need to “tweak” the filter.
Clearly rooted in the EoM

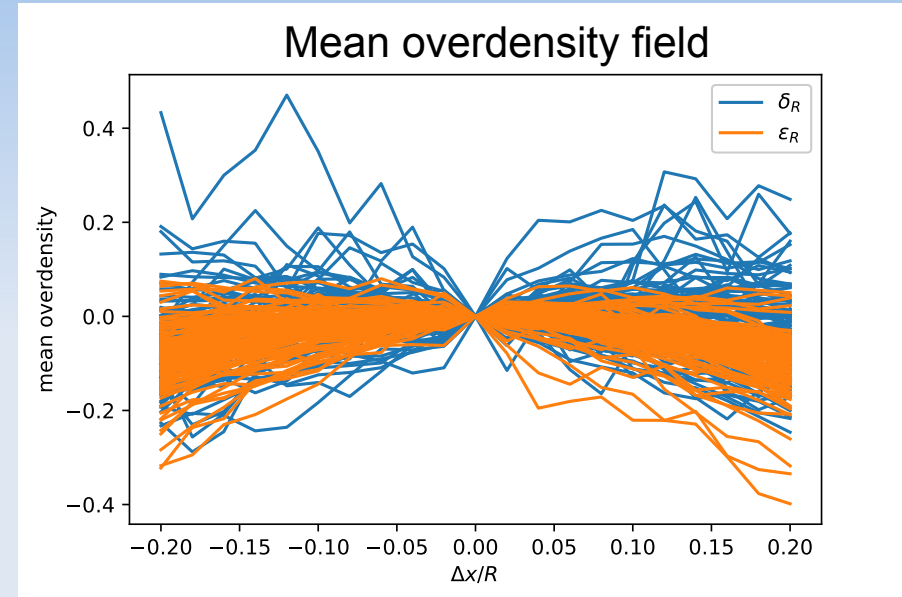
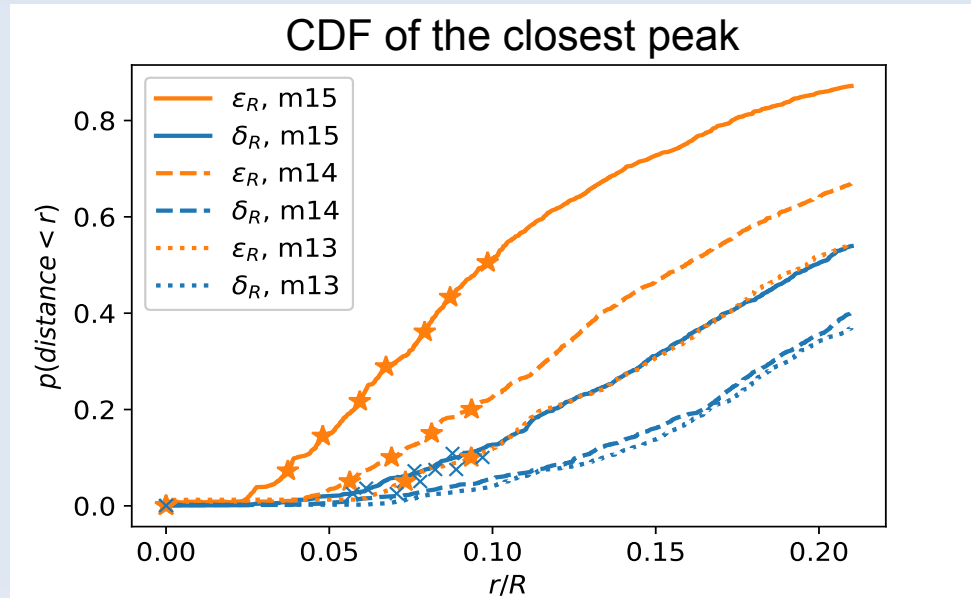
Testing the energy peak ansatz



- Energy peaks are a better proxy for protohalo centers!

Testing the energy peak ansatz

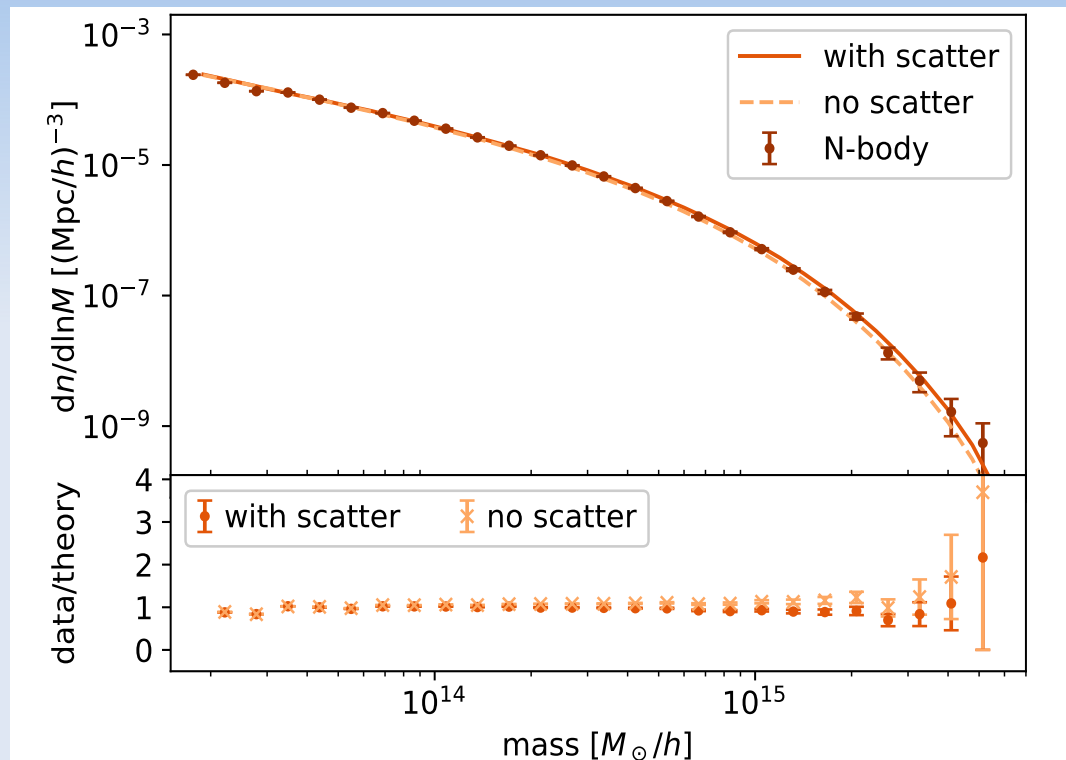
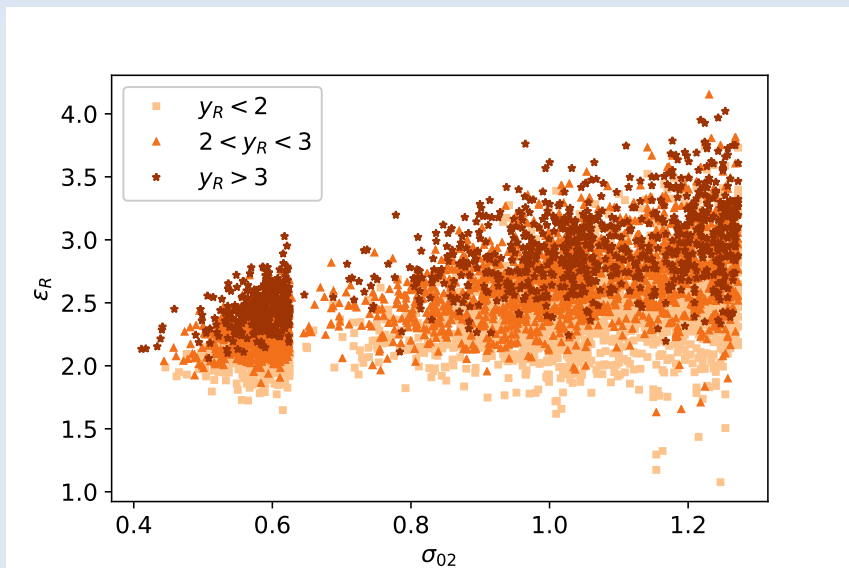
- Regions around protohalo centers are more likely to be energy peaks than matter density ones:



- Protohaloes are more likely to be close to energy peaks

Halo mass function

- Predicted, fitting the scatter of ϵ_R



- Scatter can describe assembly bias.

Other filters?

- There are other popular filter choices (Gaussian, sharp-k...)
- They have interesting mathematical properties that can make calculations easier
- However, they are not obviously connected to physical quantities
- They don't have a clear dynamical meaning (to me...)

2. The Minimum Energy Principle

Or: getting the shape right

(based on 2303.02142 with Ravi Sheth)

Shape of maximal ϵ

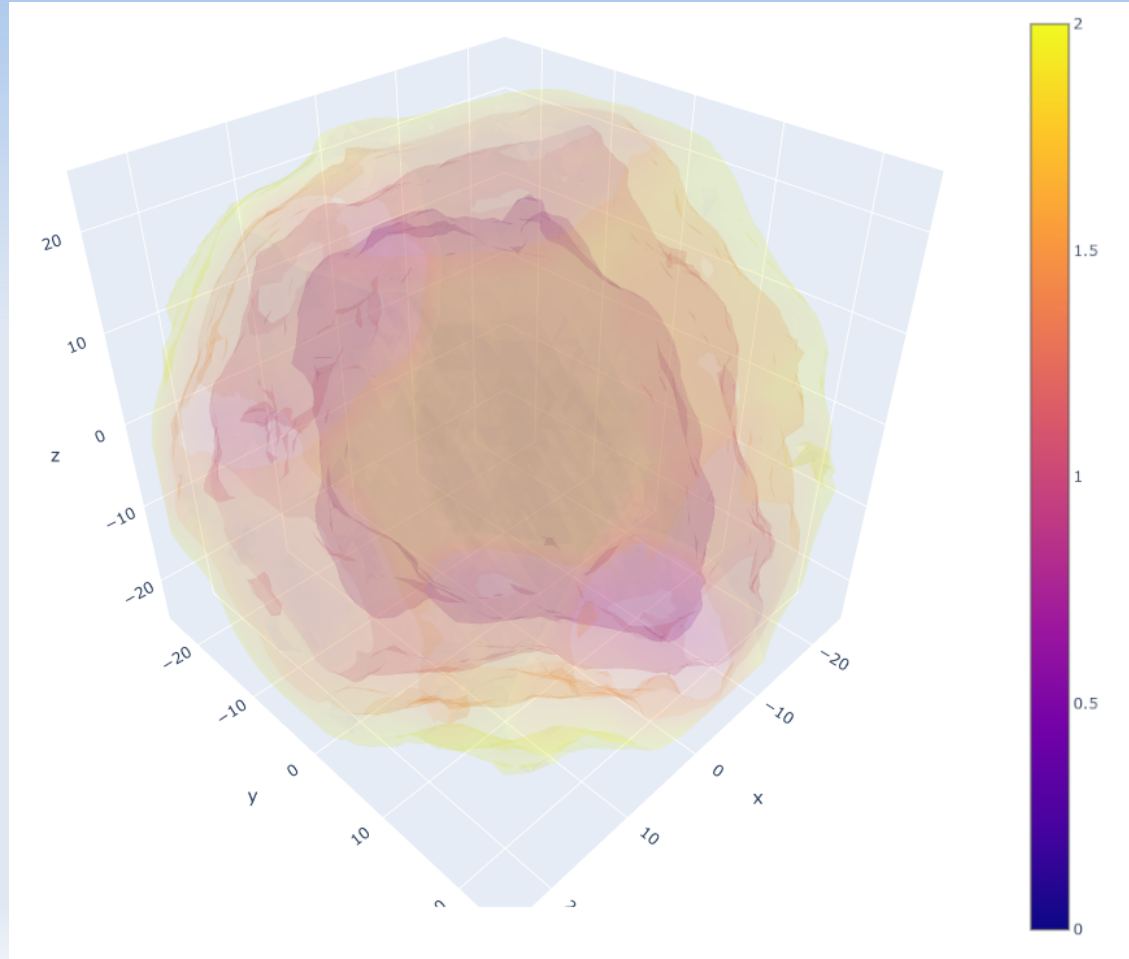
- Once a spherical peak is found, one can further increase ϵ (decrease E) by deforming the sphere at fixed volume.
- The inertial radius R_I of the deformed region collapses even faster
- The boundary of the region of maximal ϵ (minimal E) must be an isocontour of

$$\mathcal{V}(\mathbf{r}) \equiv (\mathbf{r} - \mathbf{r}_{\text{cm}}) \cdot \left[\nabla\phi - \nabla\phi_{\text{cm}} - \frac{\epsilon}{3}(\mathbf{r} - \mathbf{r}_{\text{cm}}) \right]$$

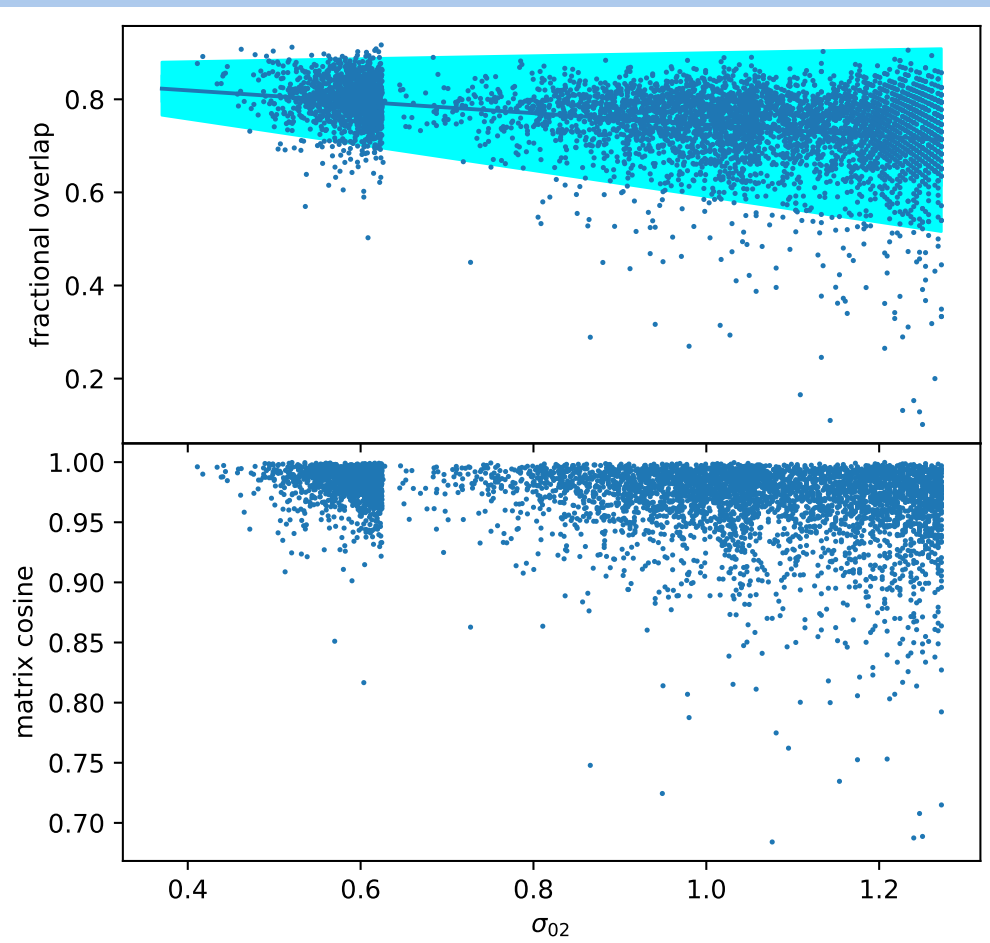
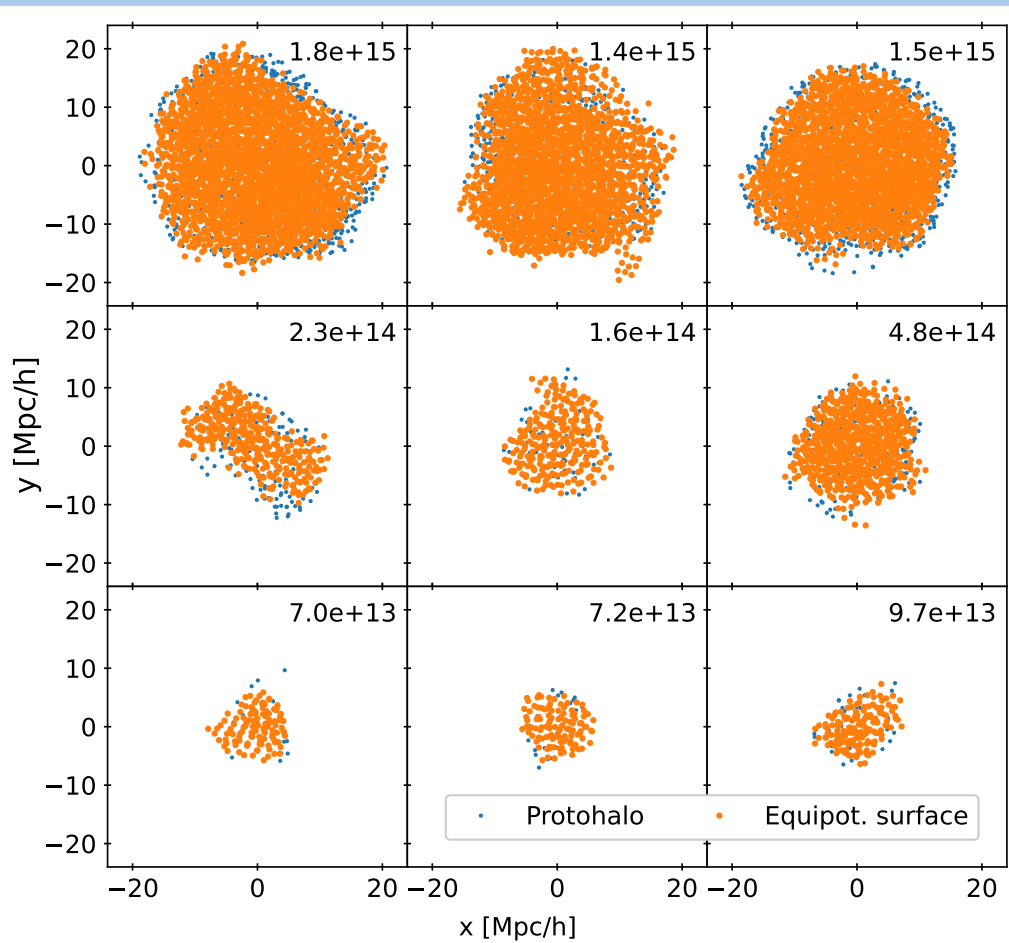
- Proxy for protohalo shape and boundary!
- Longest axis in the direction of maximum compression (orthogonal to the filament)
- Can predict initial torques

Equipotential surfaces

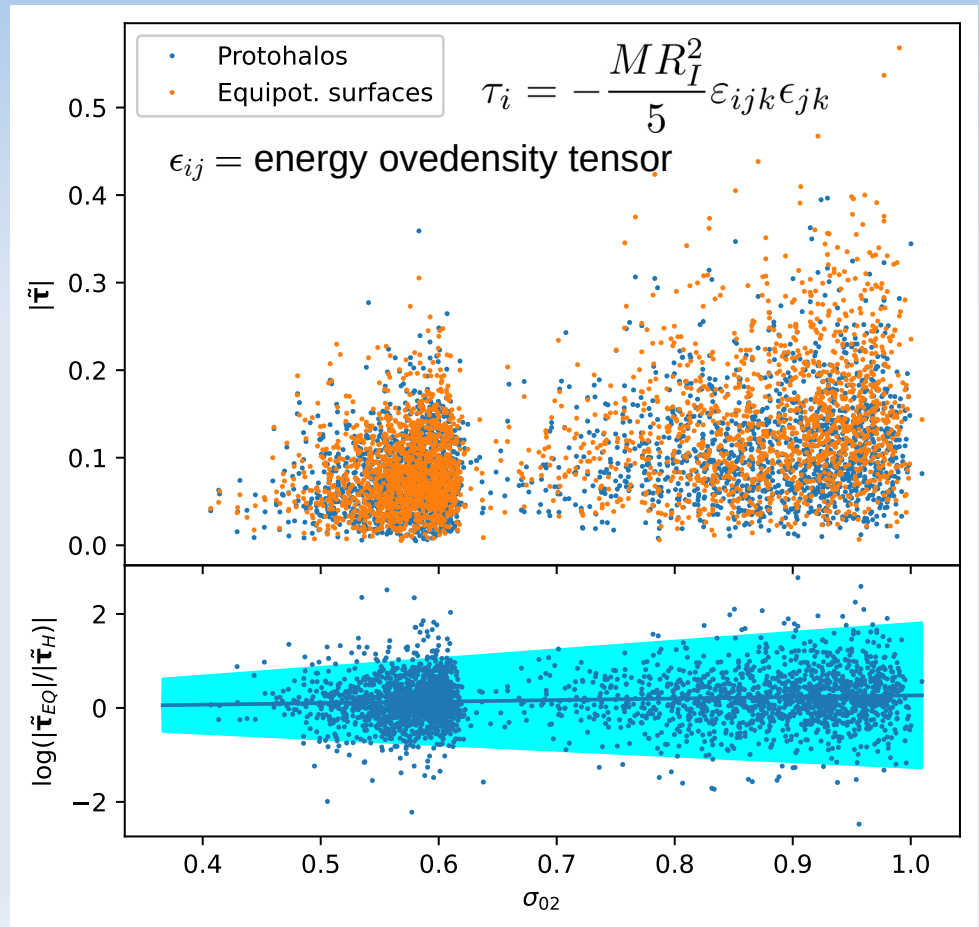
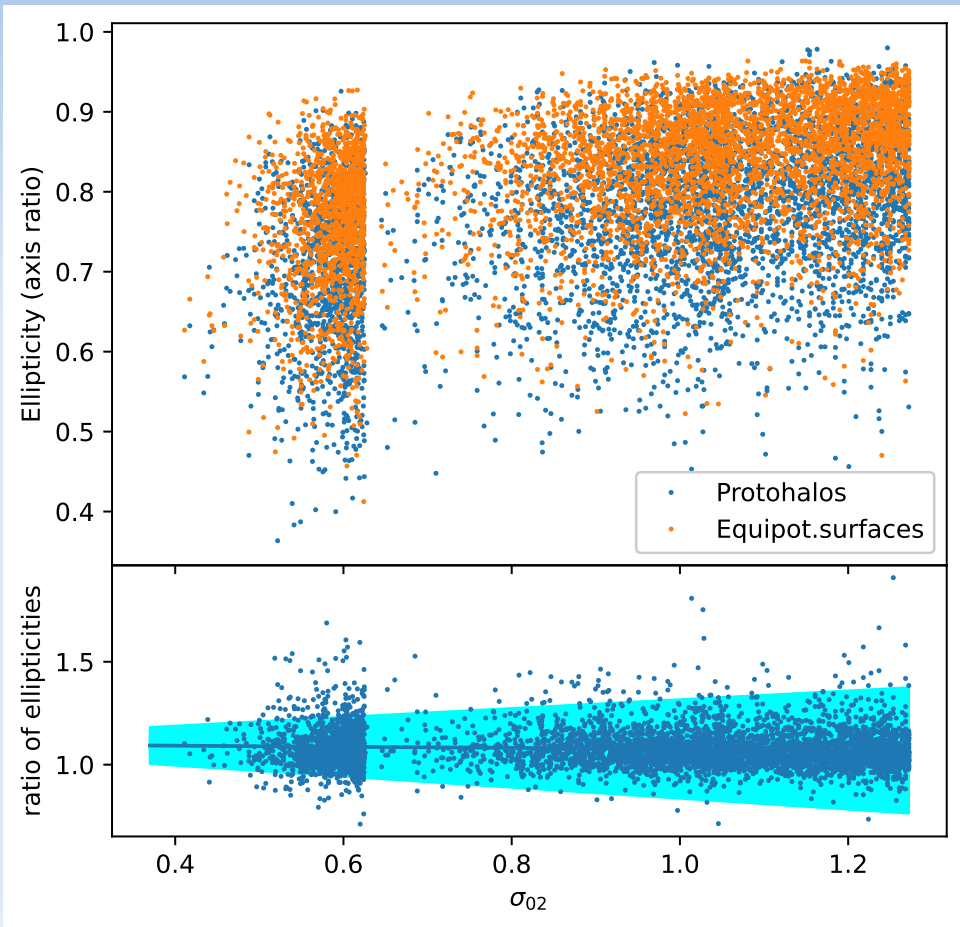
- Nested equipotential surfaces with different overdensity ϵ and volume V describe the mass accretion history
- Excursion sets of peaks of arbitrary shape!



Protohaloes vs equipotential surfaces

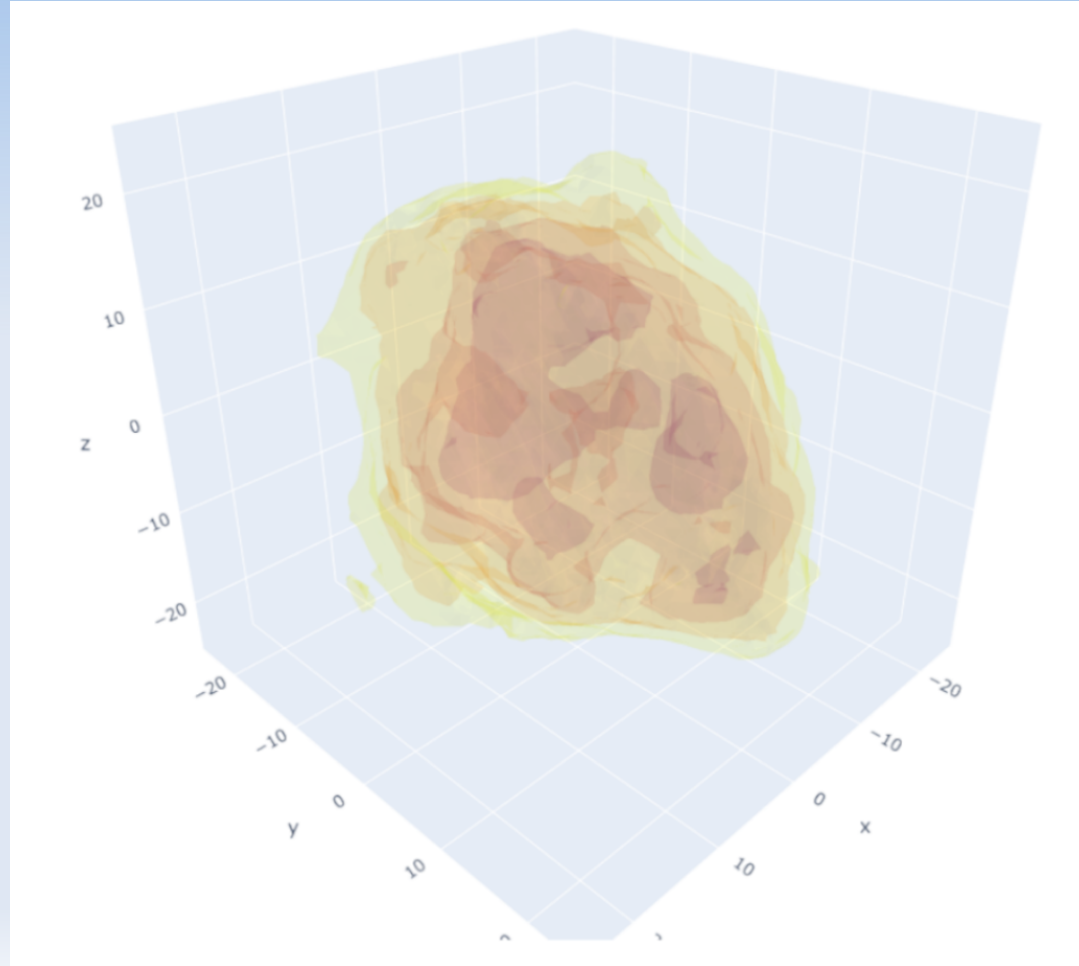


Ellipticities and torques



Mergers

- Zooming in, the surfaces of constant infall potential \mathcal{V} may fragment
- Natural prediction of merger events!



Conclusions

- Protohaloes are peaks of the initial **energy** overdensity field. Not densest but **most energetically bound** initial regions, having fastest collapse times.
- Peaks in ϵ_R are **convergent matter flows**. Initial evolution matches perturbation theory. Final high mean density results **dynamically**, not put in “by hand”.
- Using ϵ_R instead of δ_R simply means changing the filter (to a more convergent one)
- Energy density peaks are better behaved, and **better proxies for protohalo centers**
- **Protohalo shapes** are very well described by **equipotential surfaces**
- Excellent prediction of ellipticities, shear-shape alignments and torques

Open questions and outlook

- Can we predict critical value ϵ_c ? Must model virialization (in progress)
- Relation with halo finder? Ellipsoidal? FOF? Energy-based?
- Angular momentum? (in progress)
- How to improve even more? Account for non-conservation of energy?
- Final shear/shape alignments?
- (Assembly) bias? Voids? Skeleton/cosmic web?
- Primordial BHs?
- ...

Thank you!!