



The structure (and substructures) of dark matter haloes

Ian G. McCarthy

Astrophysics Research Institute, Liverpool JMU

Collaborators:

Shaun Brown, Andreea Font, Sam Stafford (LJMU),
Benedikt Diemer (Maryland)

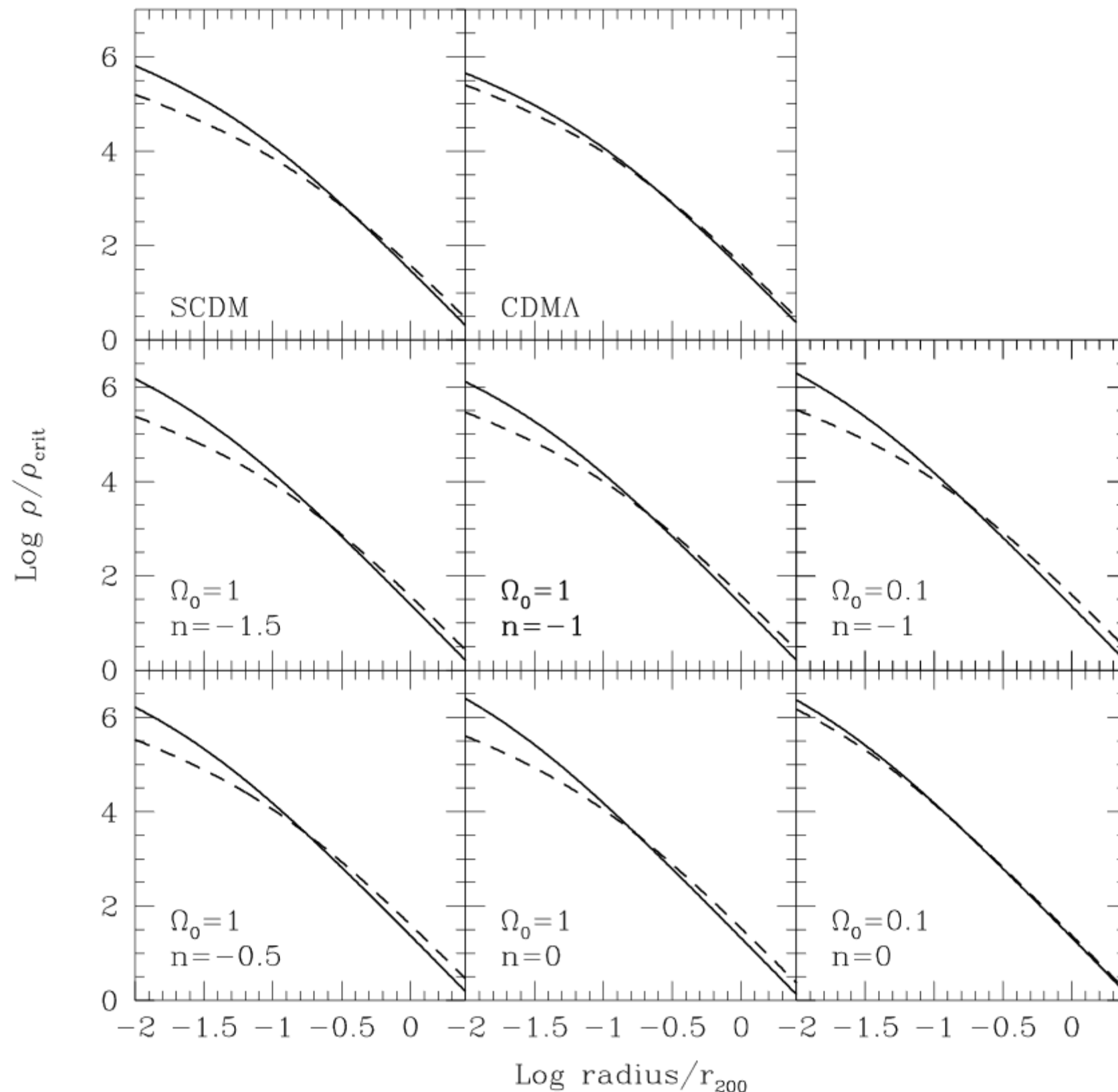
Talk objectives

- 1.) Does the overall structure (radial profiles) of dark matter haloes depend on the initial conditions?
- 2.) How can we observationally differentiate between a universe where low-mass subhaloes are completely dark vs. one where they are not present at all?

Part I: Structure of dark matter haloes

(Brown, IGM+2020, MNRAS, 495, 4994)

The self-similarity of dark matter haloes



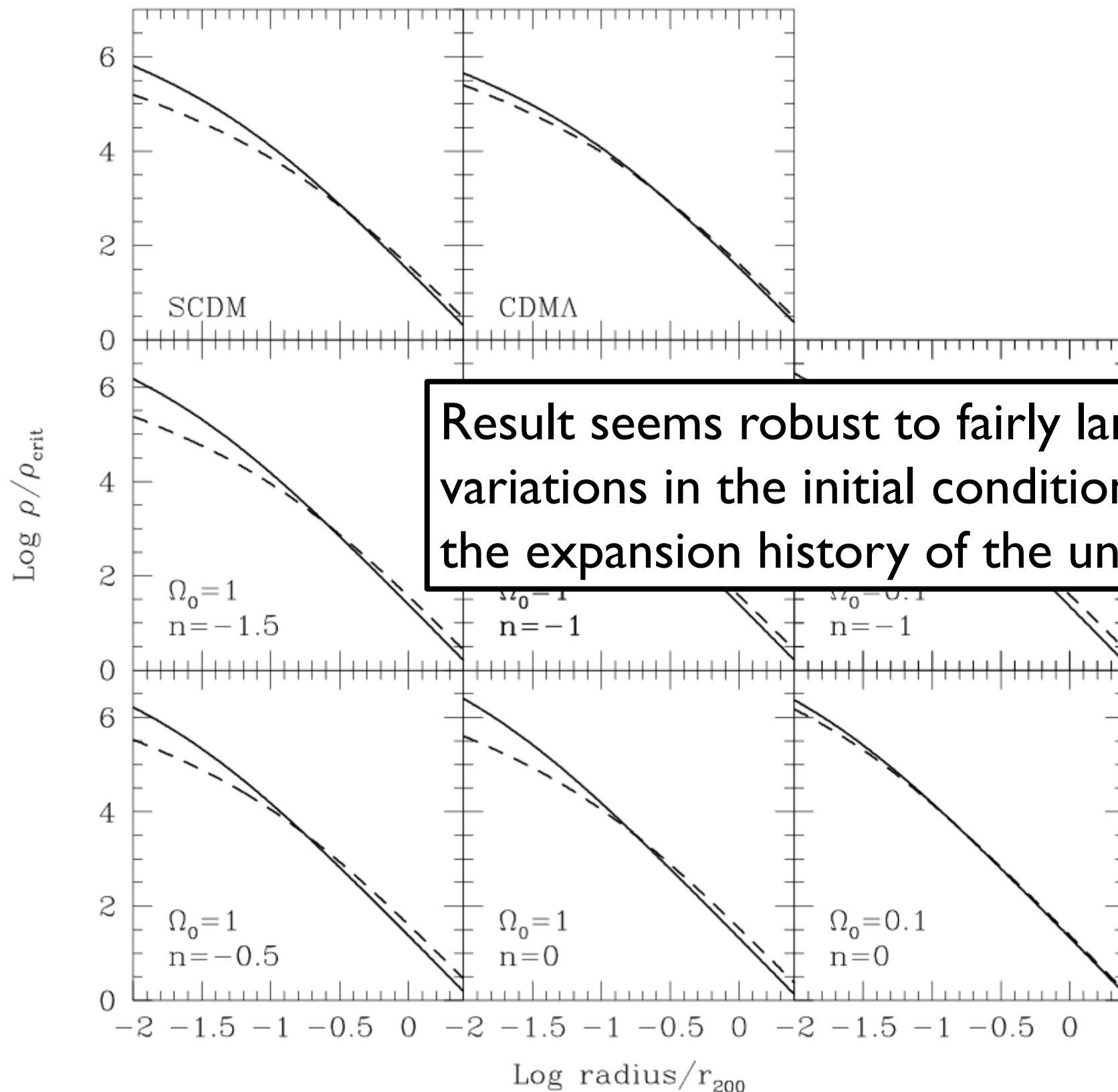
When radial coordinate is normalised by virial radius, profiles are approximately self-similar.

Underpredicts the density in the inner regions, so concentration (or scale radius) must vary with halo mass.

Linked to halo formation time.

$$\frac{\rho(r)}{\rho_{\text{crit}}} = \frac{\delta_c}{(r/r_s)(1 + r/r_s)^2}$$

The self-similarity of dark matter haloes



Result seems robust to fairly large variations in the initial conditions and the expansion history of the universe

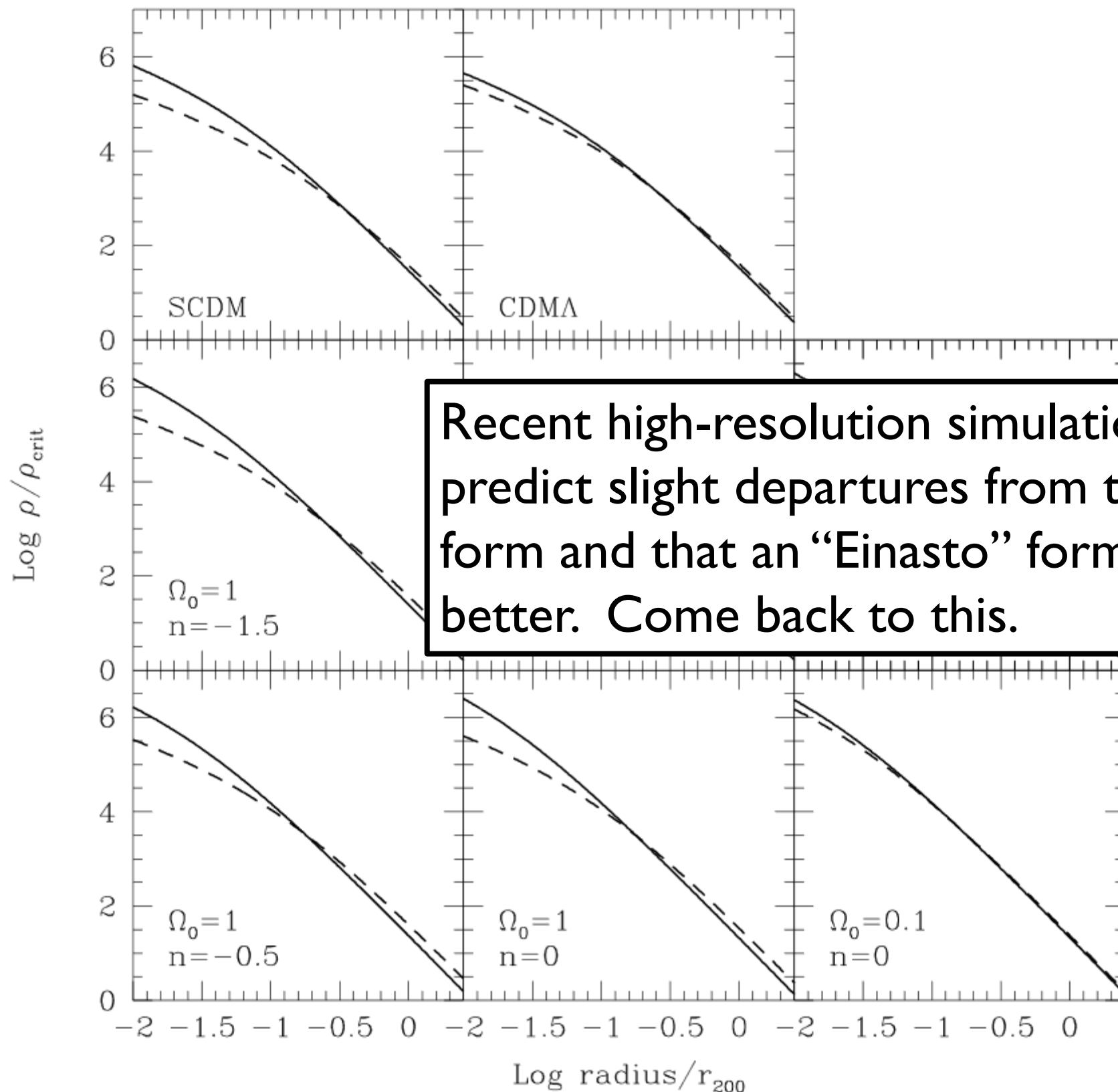
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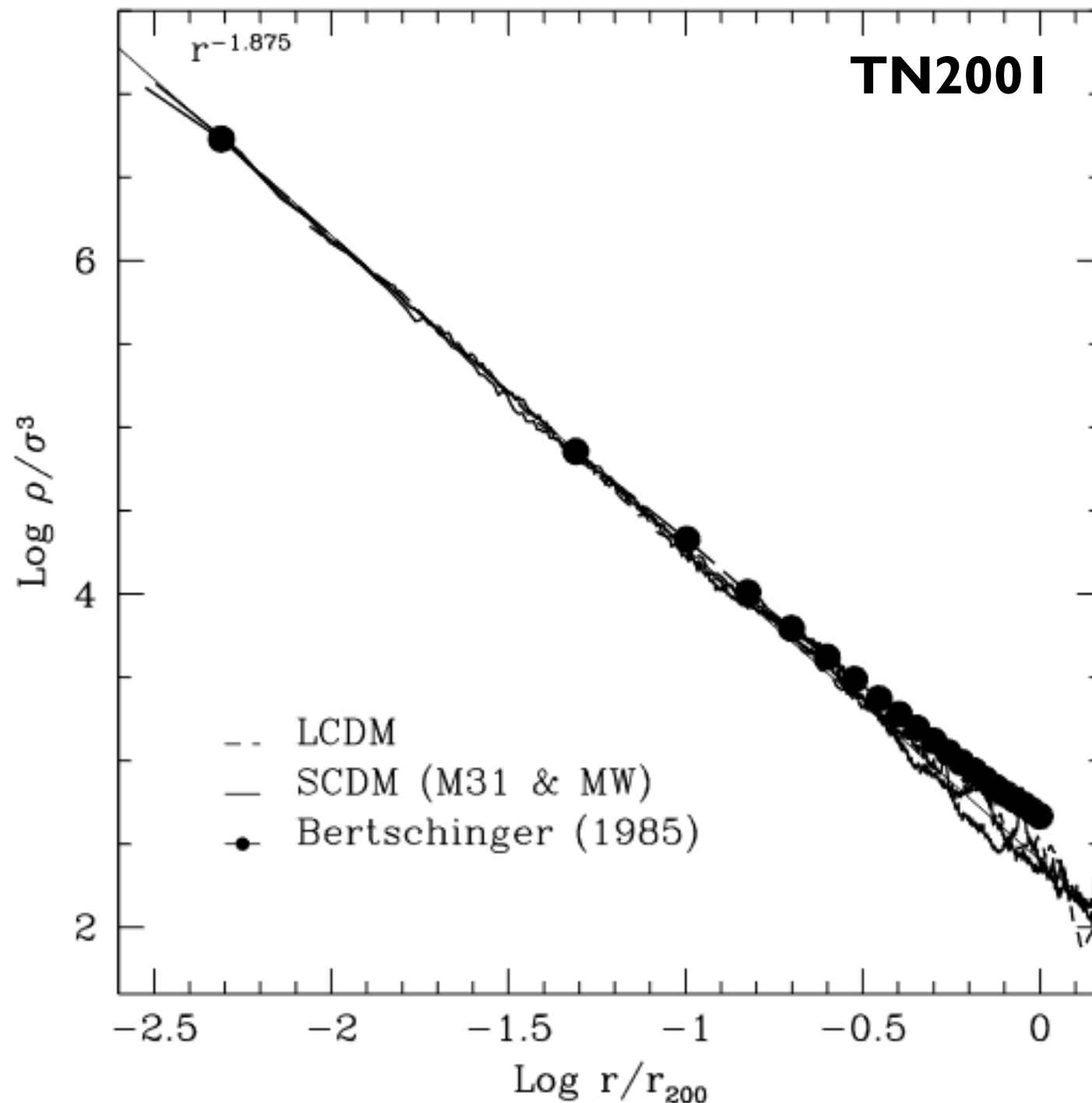
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The “phase-space” density of DM haloes



Curious result:

Taylor & Navarro (2001) showed that a quantity called the “phase space” density is approx. a pure power law, independent of mass.

Is this more fundamental?
What sets the slope?

Why does Bertschinger (appear to) work?

Aside: “phase space” density or entropy?

Phase space density defined as:

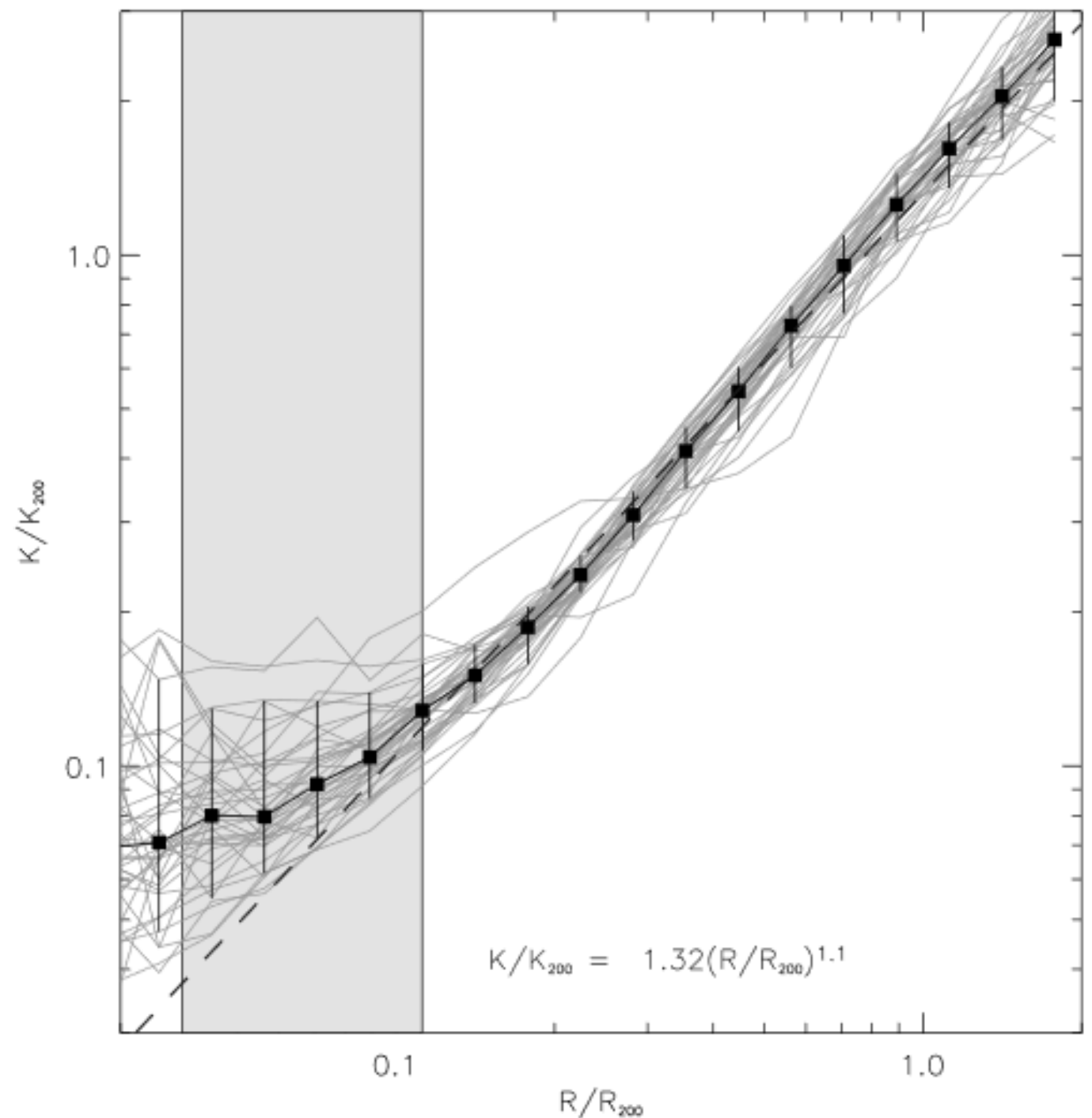
$$\frac{\rho}{\sigma^3}$$

Writing $\sigma^2 \rightarrow kT$ and taking to $-2/3$ power, one gets:

$$\frac{kT}{\rho^{2/3}}$$

Sometimes called the “entropy”.

Interestingly, the intracluster medium has the same entropy profile as the dark matter.



Voit+2005

Aside: “phase space” density or entropy?

Phase space density defined as:

$$\frac{\rho}{\sigma^3}$$

Writing $\sigma^2 \rightarrow kT$ and power, one gets:

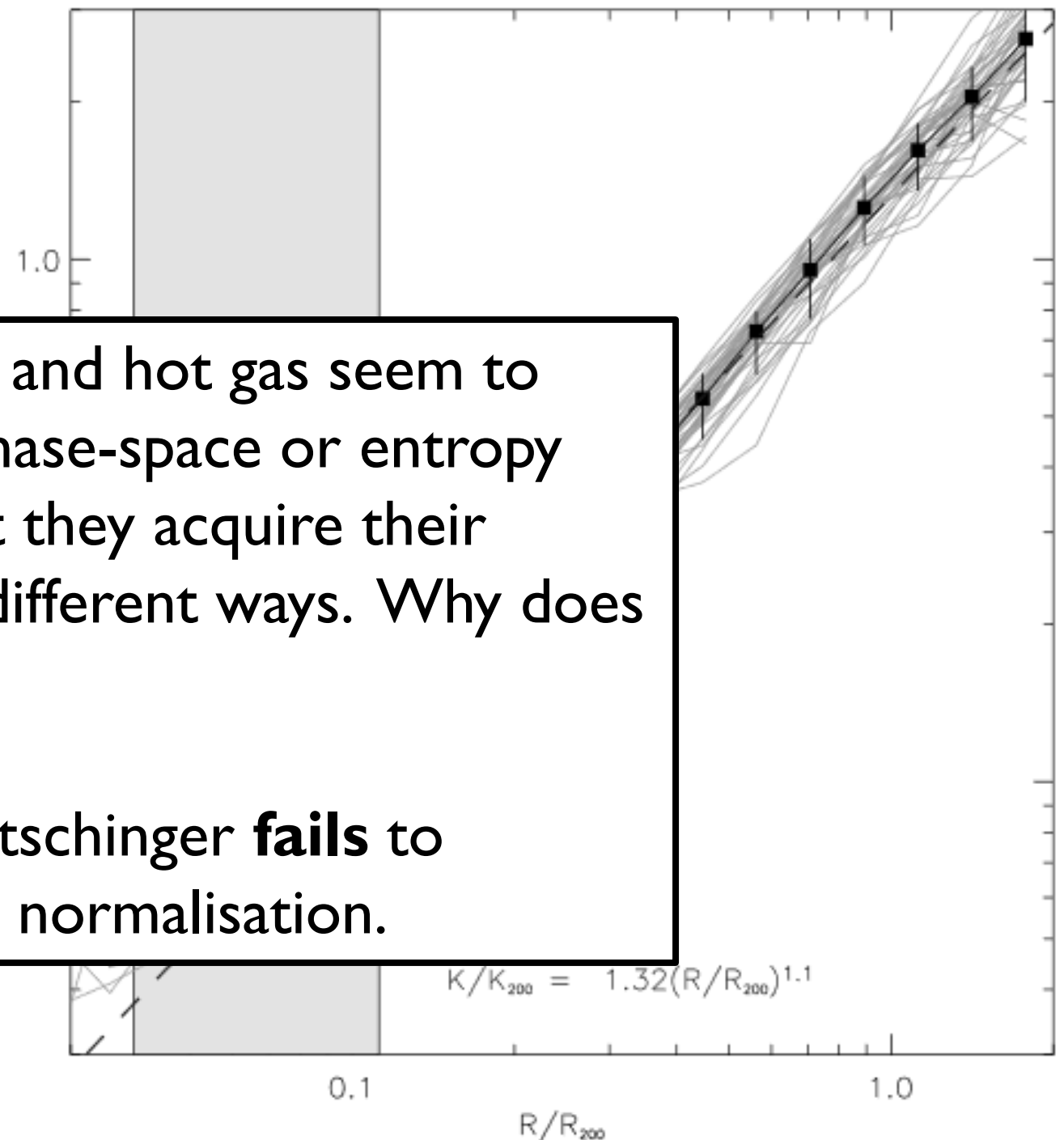
$$\frac{kT}{\rho^{2/3}}$$

Sometimes called the entropy.

Interestingly, the intracluster medium has the same entropy profile as the dark matter.

The dark matter and hot gas seem to have the same phase-space or entropy distributions. Yet they acquire their entropy in very different ways. Why does this work?

Voit showed Bertschinger **fails** to correctly predict normalisation.



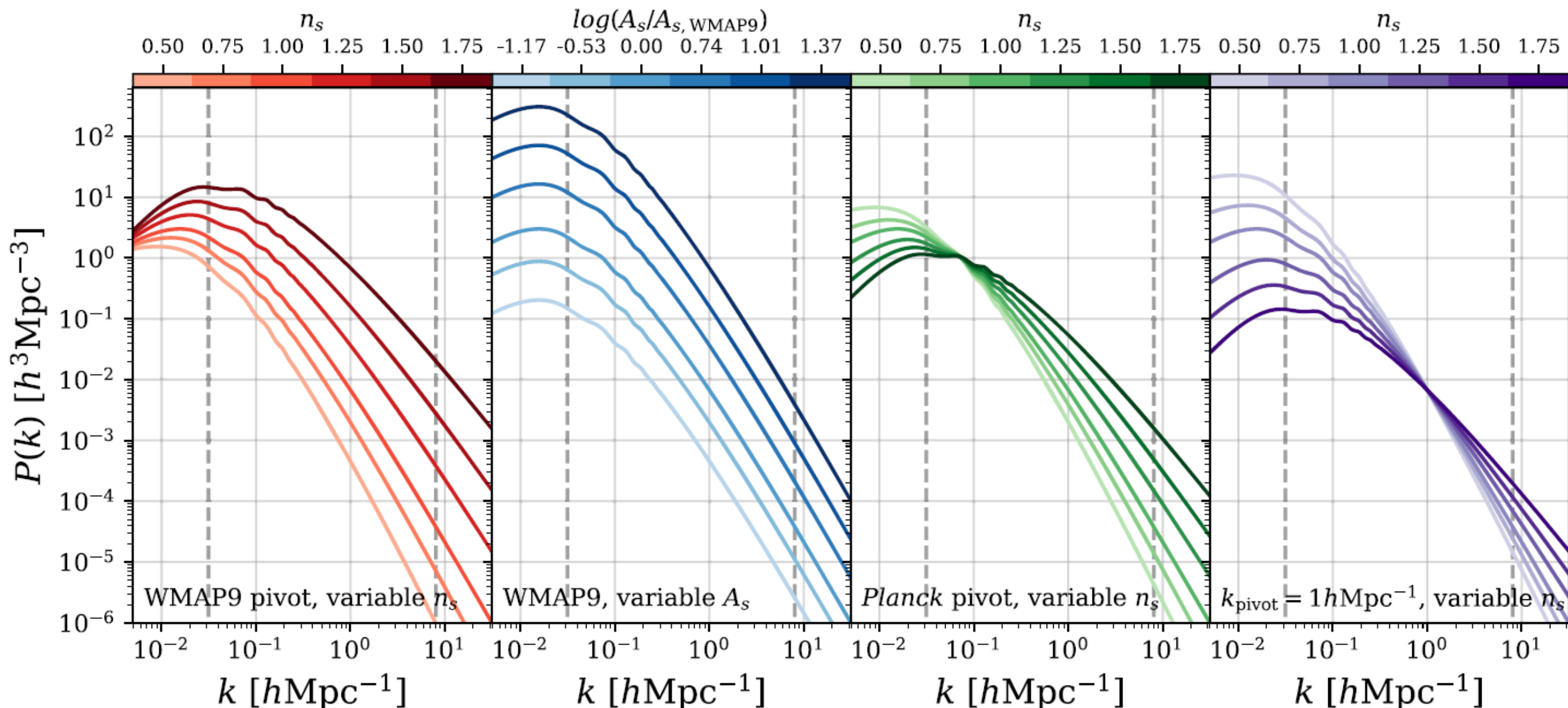
Voit+2005

Universal or not?

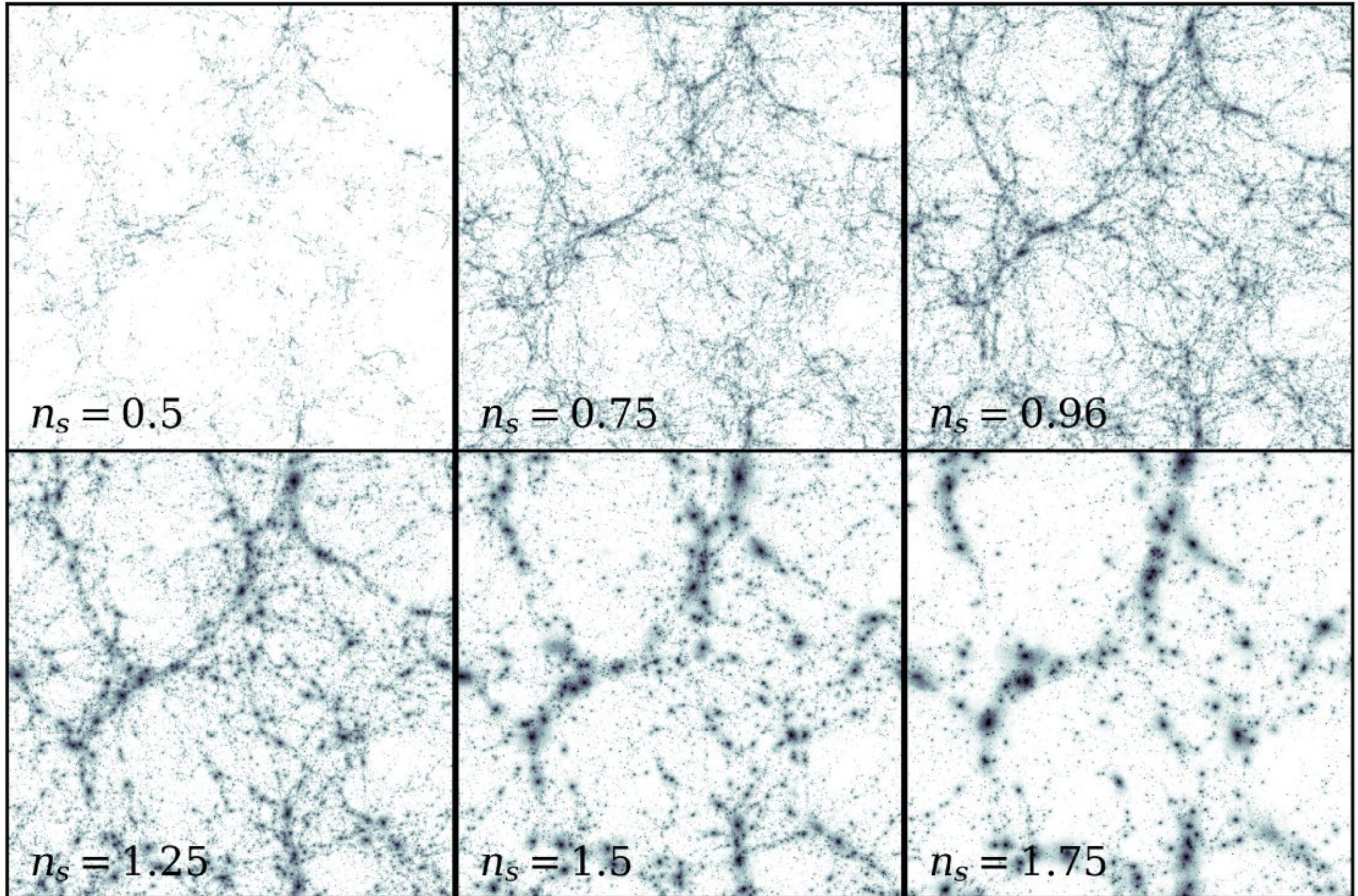
A numerical experiment:

Revisit the sensitivity to the initial conditions. **Brown+2020, MNRAS, 495, 4994**

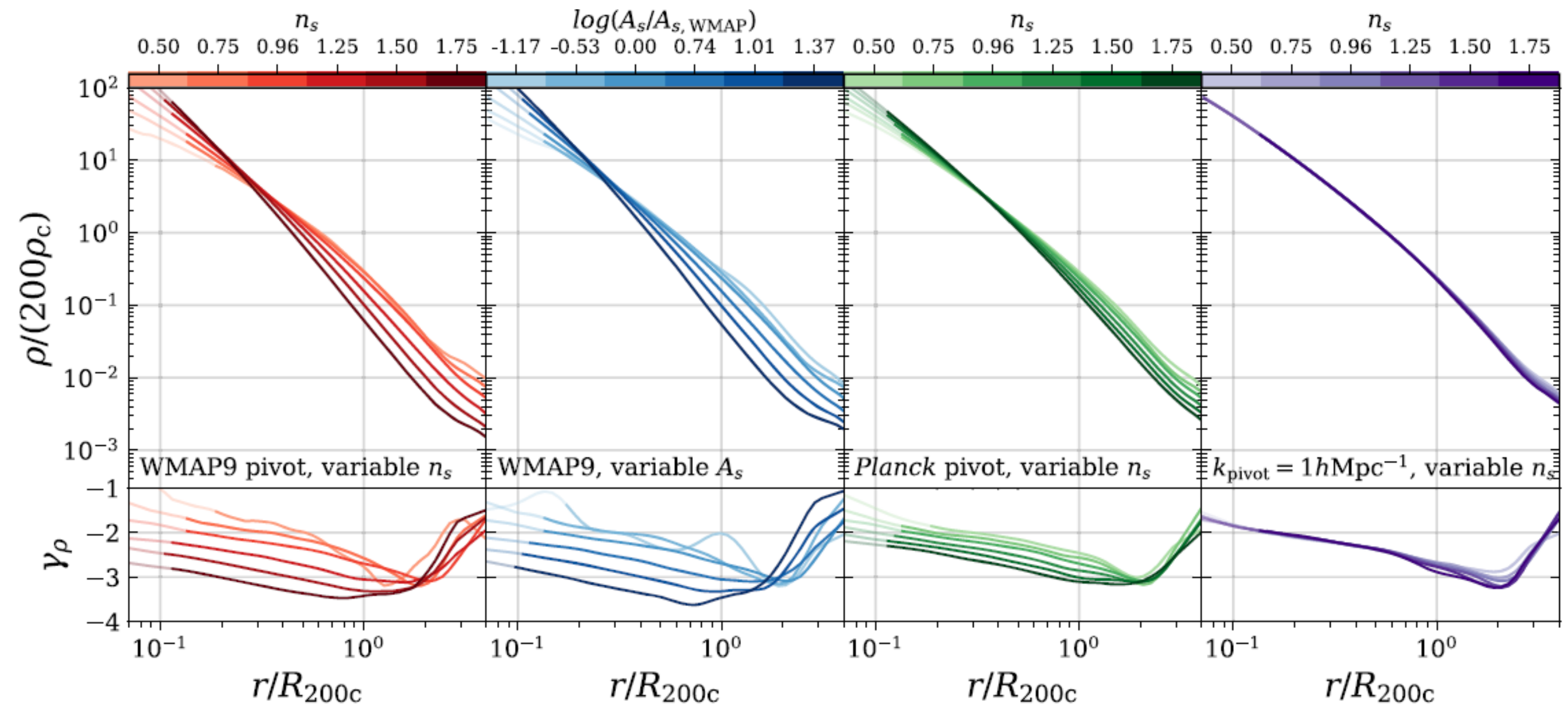
Noteworthy that previous studies did **not** explore a wide range of power spectrum normalisations (i.e., σ_8 was always ~ 1 ; “CMB-normalised”).



Resulting structure (visual guide)



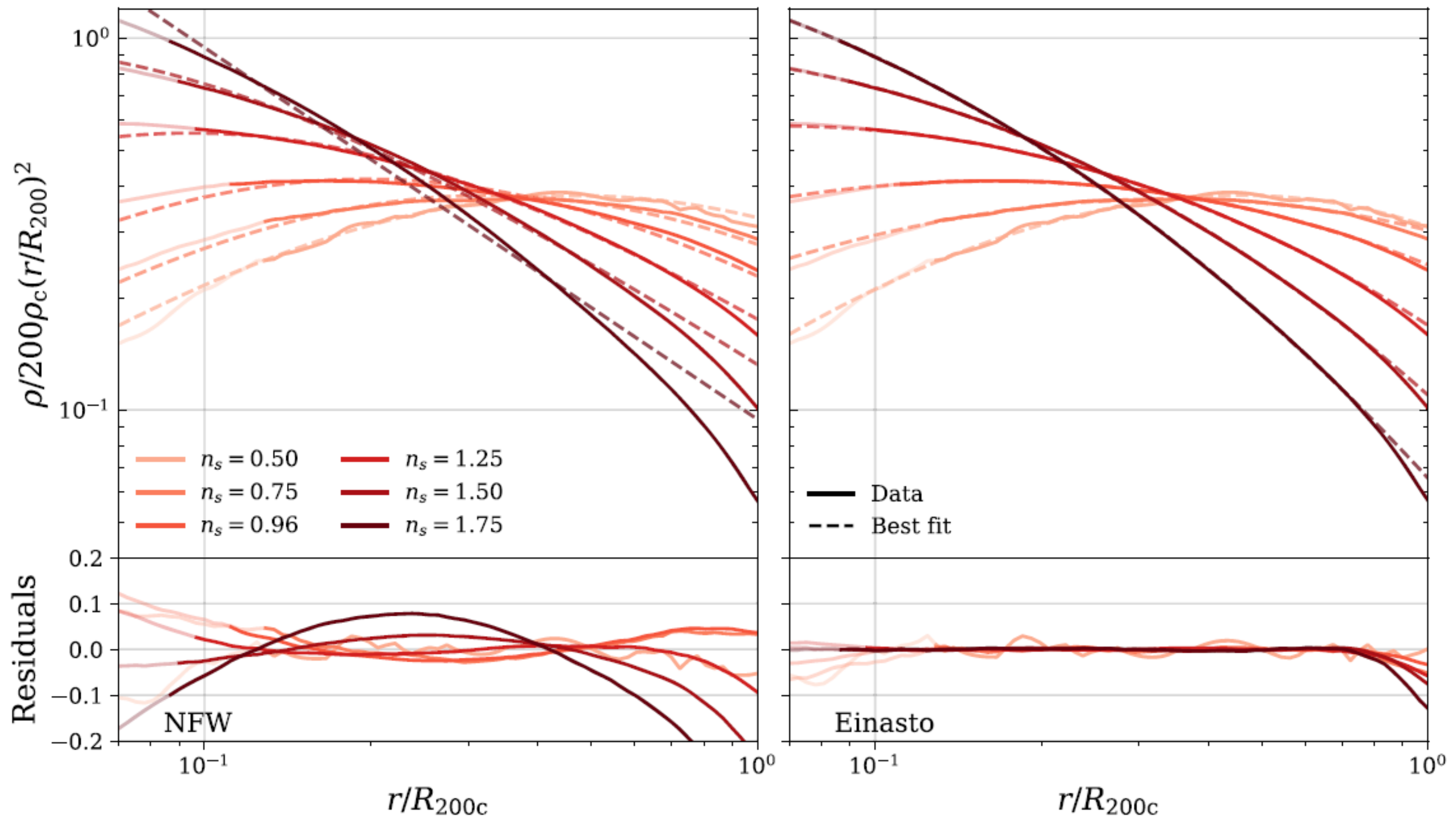
Density profiles



For haloes with $M_{200c} \sim 10^{13} M_\odot$

Brown, IGM+2020

Testing NFW and Einasto



For haloes with $M_{200c} \sim 10^{13} M_{\odot}$.

Brown, IGM+2020

Predicting the trends (in hindsight)

With a high-normalisation primordial power spectrum, the simulations show logarithmic slopes steeper than -3 well within R_{200c} . Could this have been anticipated? Potentially, yes.

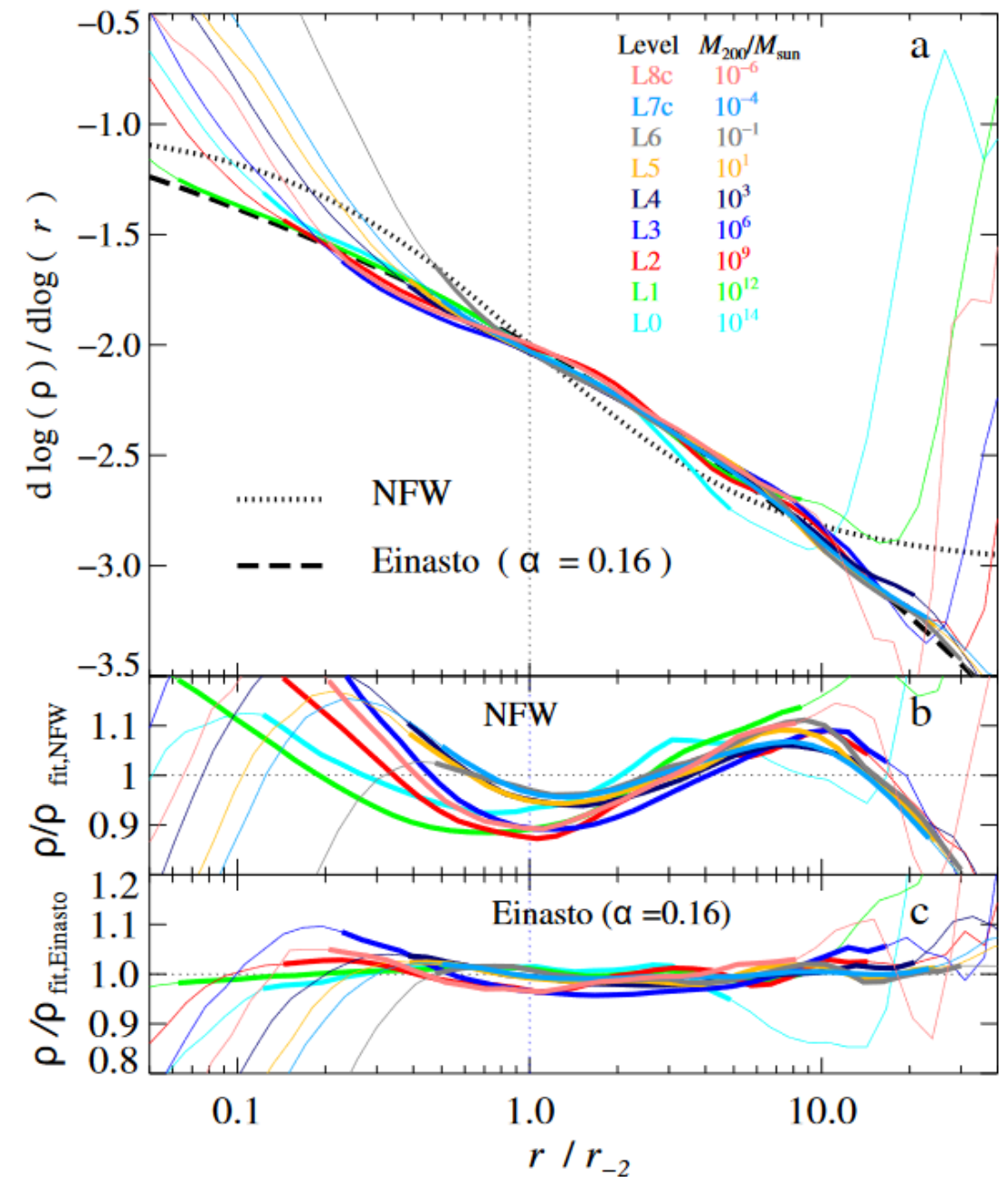
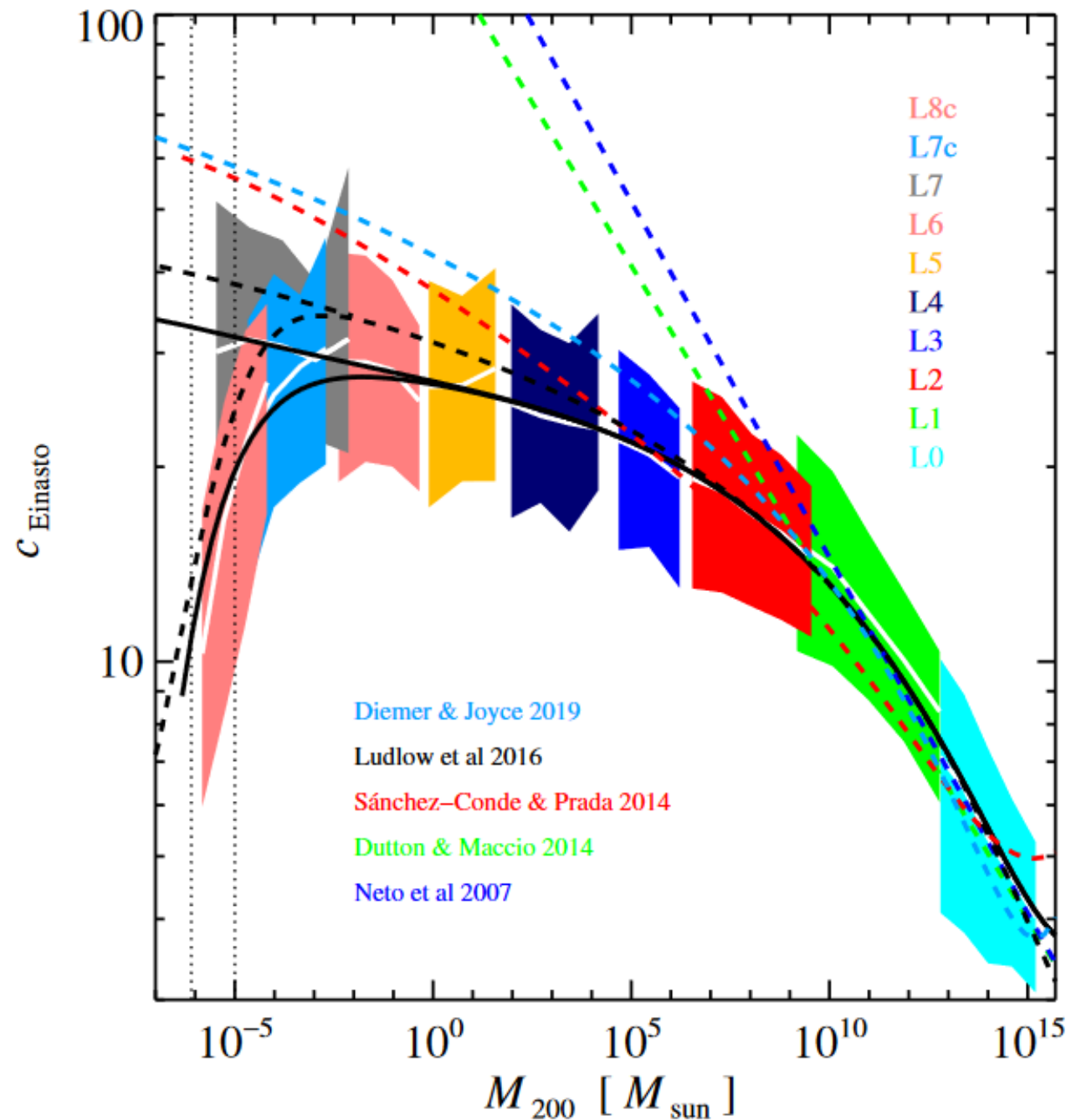
Logarithmic slope of Einasto profile is:

$$\gamma(r) = -2 \left(\frac{cr}{R_{200c}} \right)^\alpha$$

Adopting $\alpha = 0.16$, one predicts $\gamma(r) < -3$ at $r = R_{200c}$ for $c > 12.6$, Einasto is correct. Simulations verify this (i.e., Einasto works well).

Does this ever happen for CMB-normalised cosmologies? Yes!

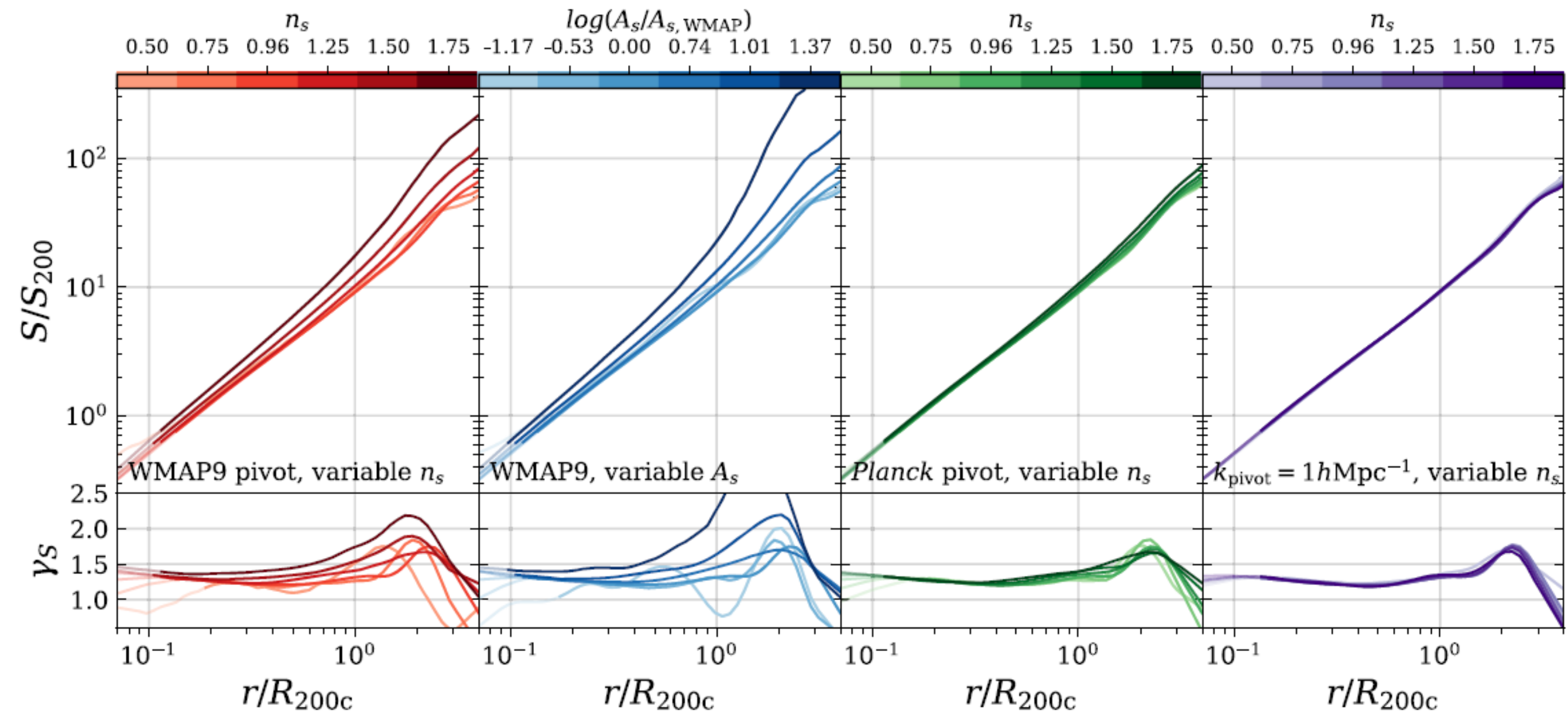
Predicting the trends (in hindsight)



Low-mass haloes in CMB-normalised cosmologies are sufficiently concentrated to show strong departure from NFW.

Wang, Bose+2020

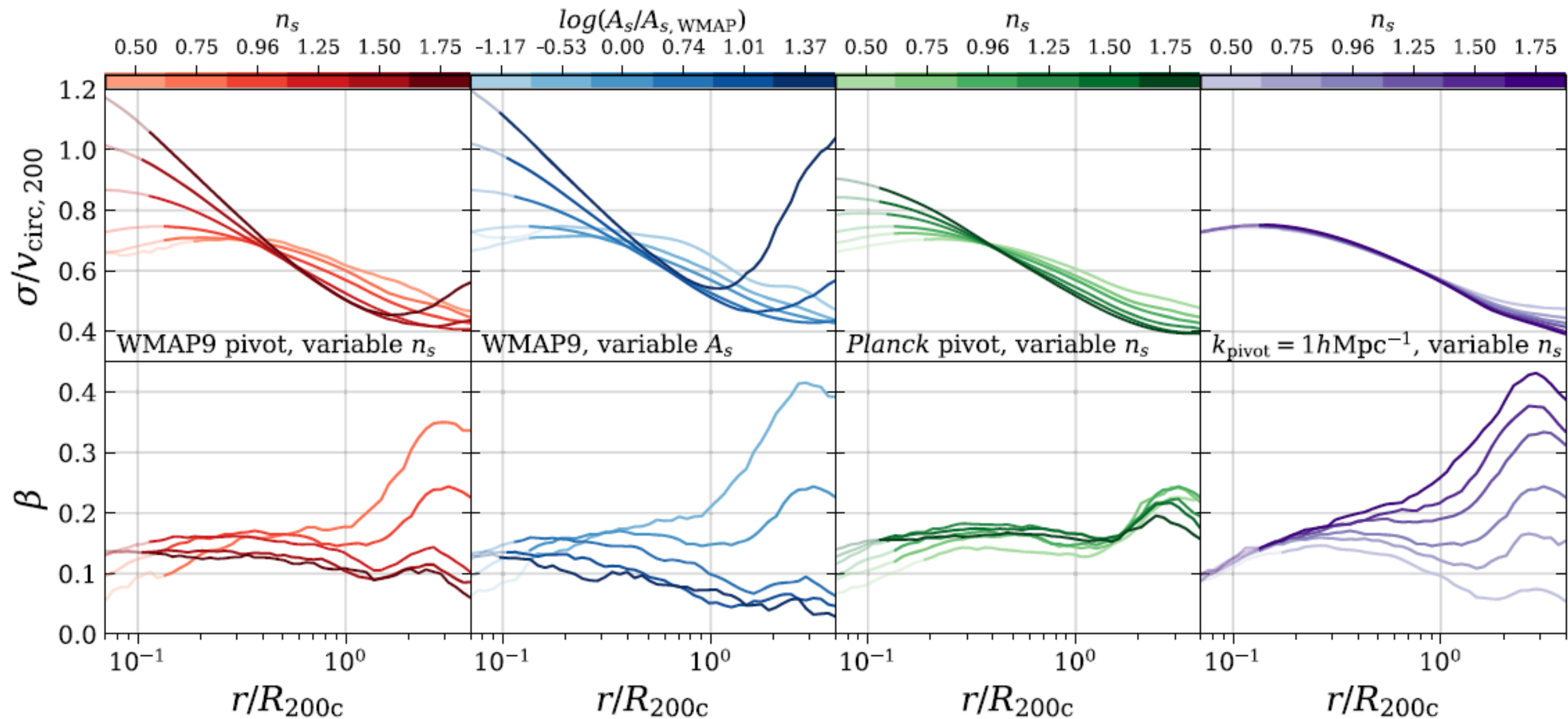
Entropy profiles



For haloes with $M_{200c} \sim 10^{13} M_{\odot}$

Brown, IGM+2020

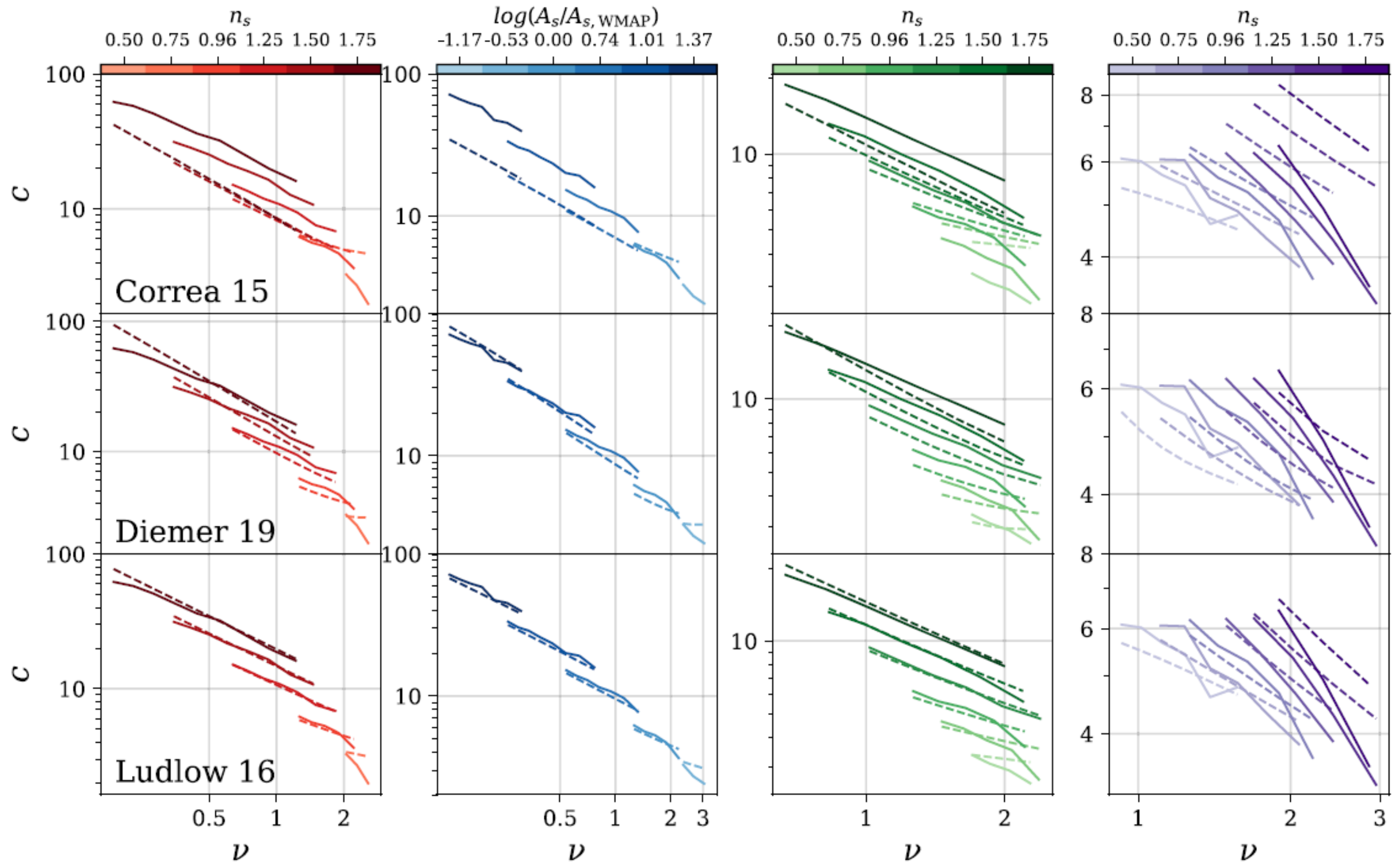
Velocity dispersion and velocity anisotropy



For haloes with $M_{200c} \sim 10^{13} M_{\odot}$

Brown, IGM+2020

Concentration – peak height relations



Conclusions (part I)

- The NFW form has been around for almost 25 years, appearing to be robust to changes in the expansion history and initial conditions of the Universe.
- However, previous numerical experiments were mainly limited to “CMB-normalised” cosmologies with $\sigma_8 \sim 1$, which we find to be key.
- We have shown that **the structure and kinematics of dark matter haloes depend strongly on the amplitude of the primordial power spectrum** at a characteristic wavenumber.
- Increasing (decreasing) the amplitude steepens (shallows) the density and entropy profiles. In this sense, there is no universal profile.
- The shape of the power spectrum at the characteristic wavenumber plays a secondary (but still relevant) role.

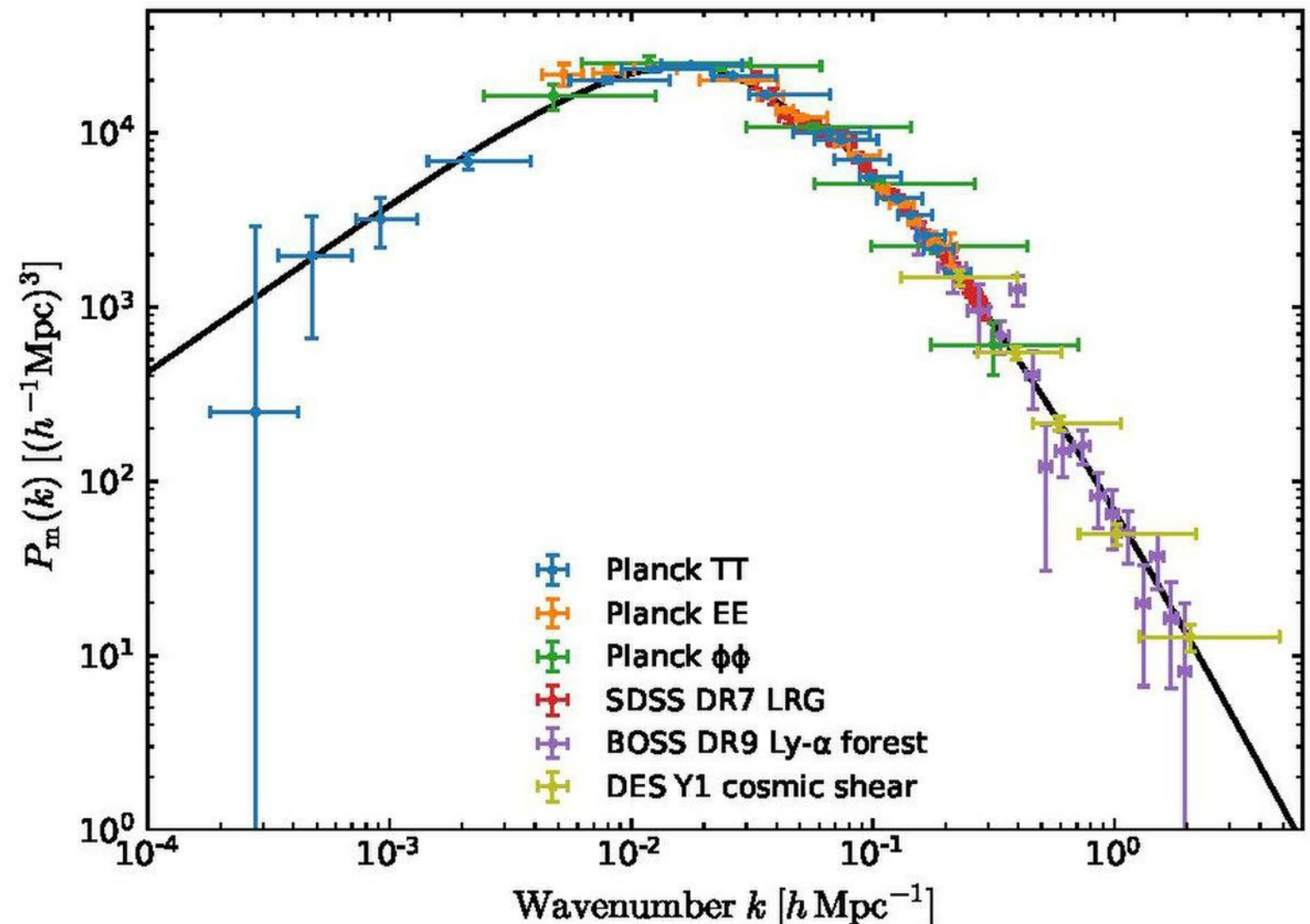
Part II: Substructures of dark matter haloes

(IGM+2020, MNRAS, 499, 3255)

The Λ CDM model

The standard model of cosmology posits:

- GR is the correct theory of gravity
- Universe is isotropic, homogeneous, and spatially flat.
- Constituent energy components are: radiation, baryons, dark matter, dark energy
- Dark matter is assumed to interact only via gravity and is “cold”



In general, has had **remarkable success on “large” scales.**

Great, but what does that mean?

Predictions on small scales

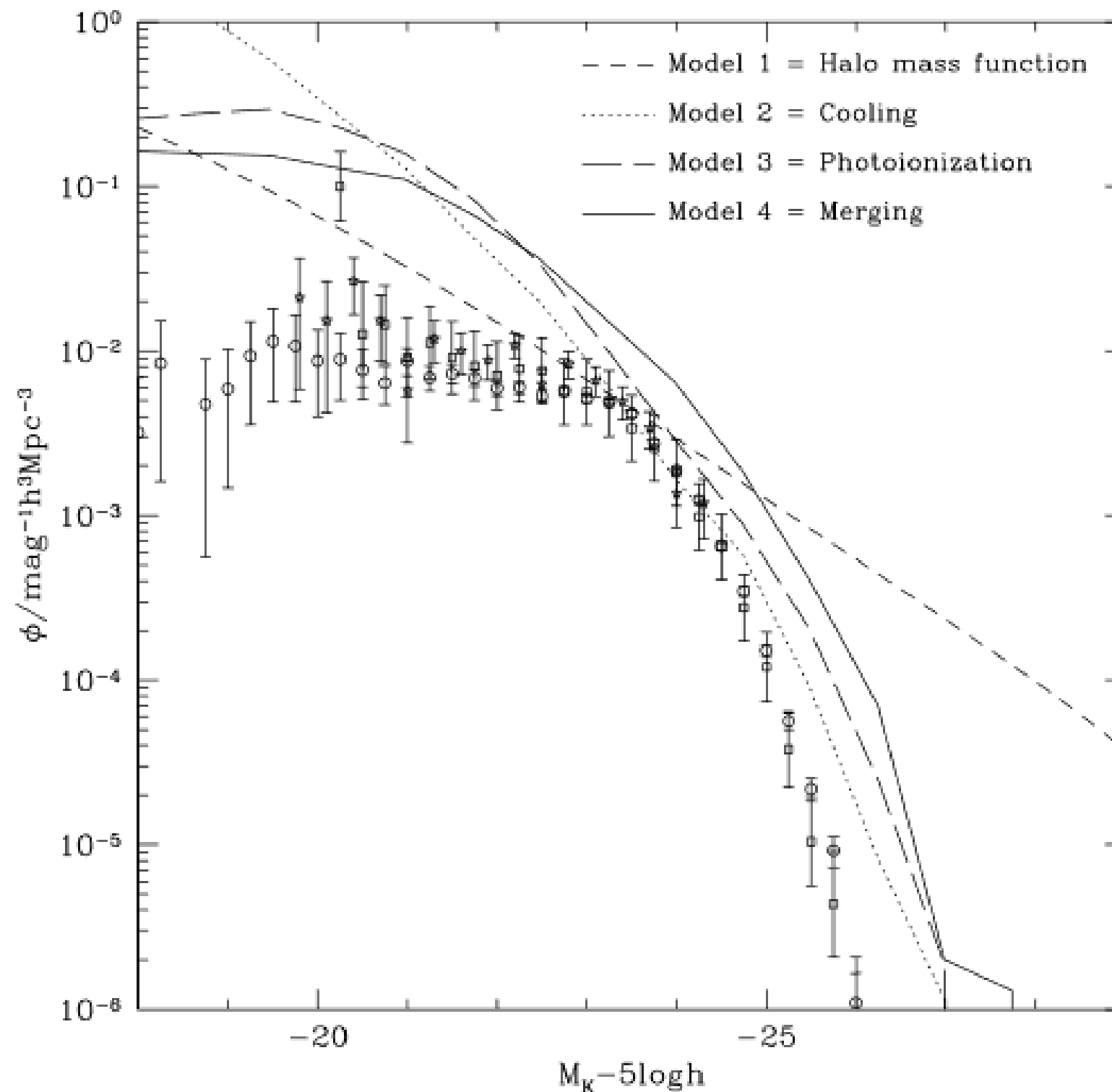
Aquarius “zoom” simulation



Springel+2008

A typical Milky Way-mass halo

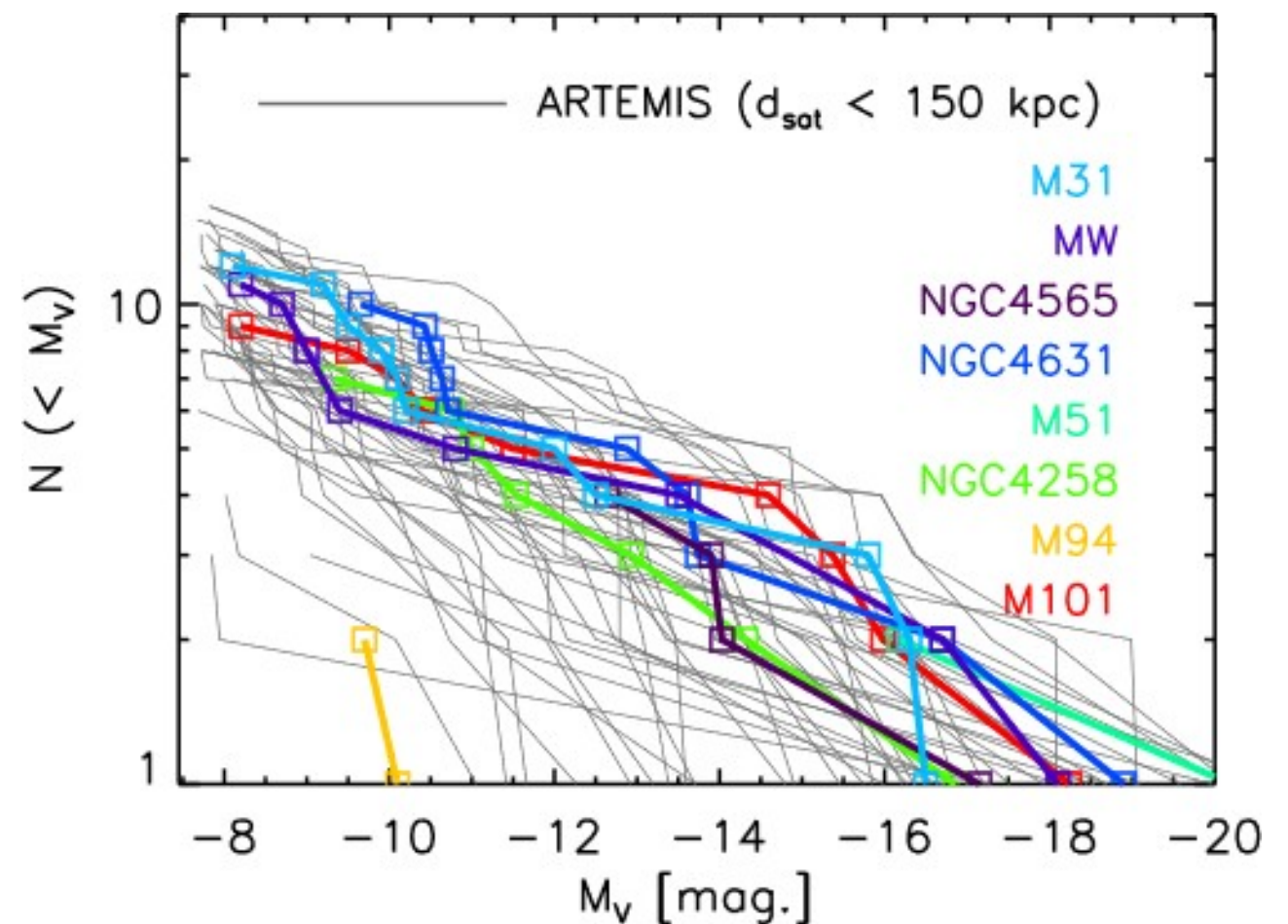
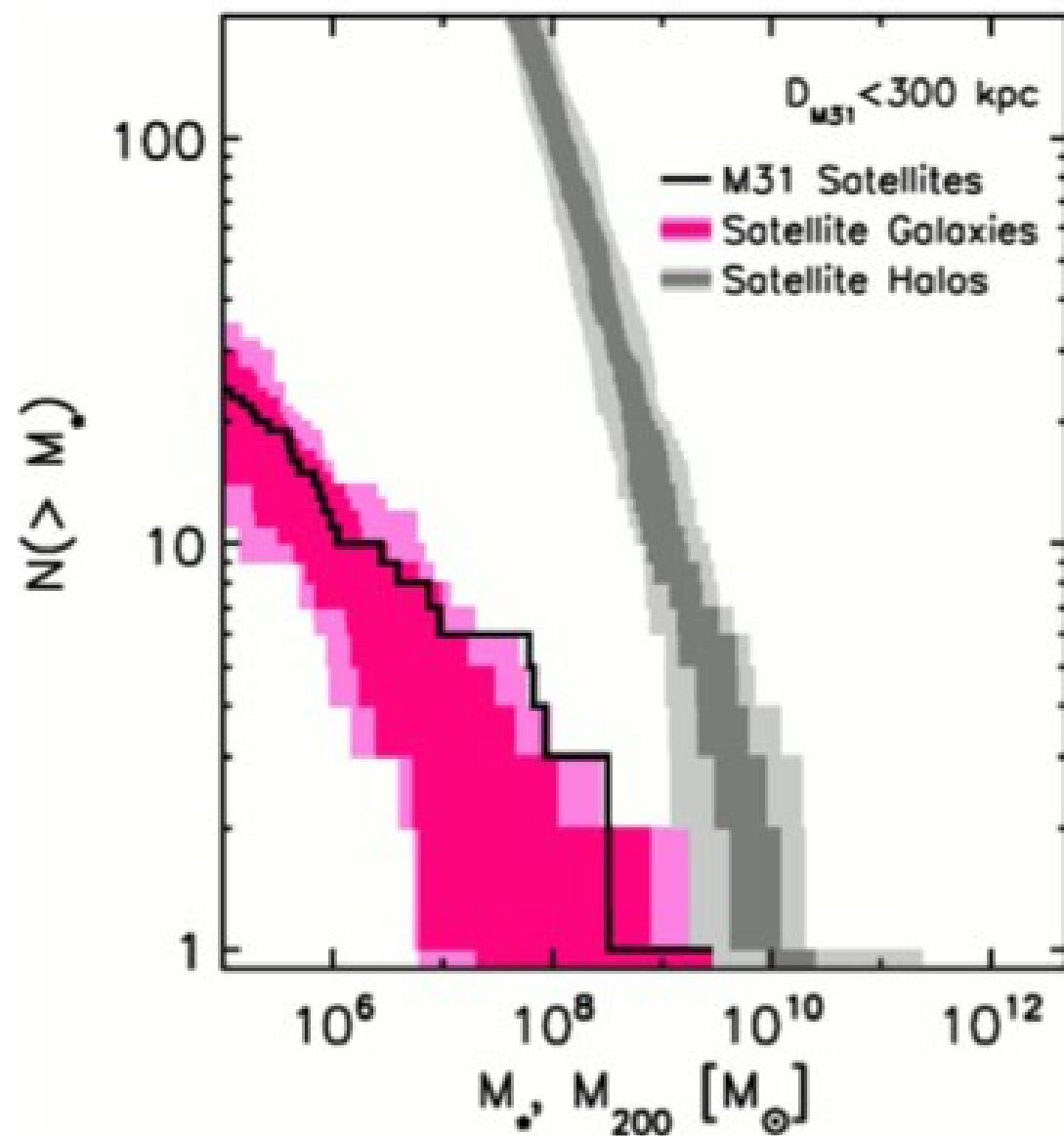
Impact of galaxy formation



Photoionisation and stellar feedback combine to prevent significant star formation in low-mass subhaloes.

Photoionisation is particularly effective for haloes with circular velocities $< \sim 10$ km/s

Impact of galaxy formation

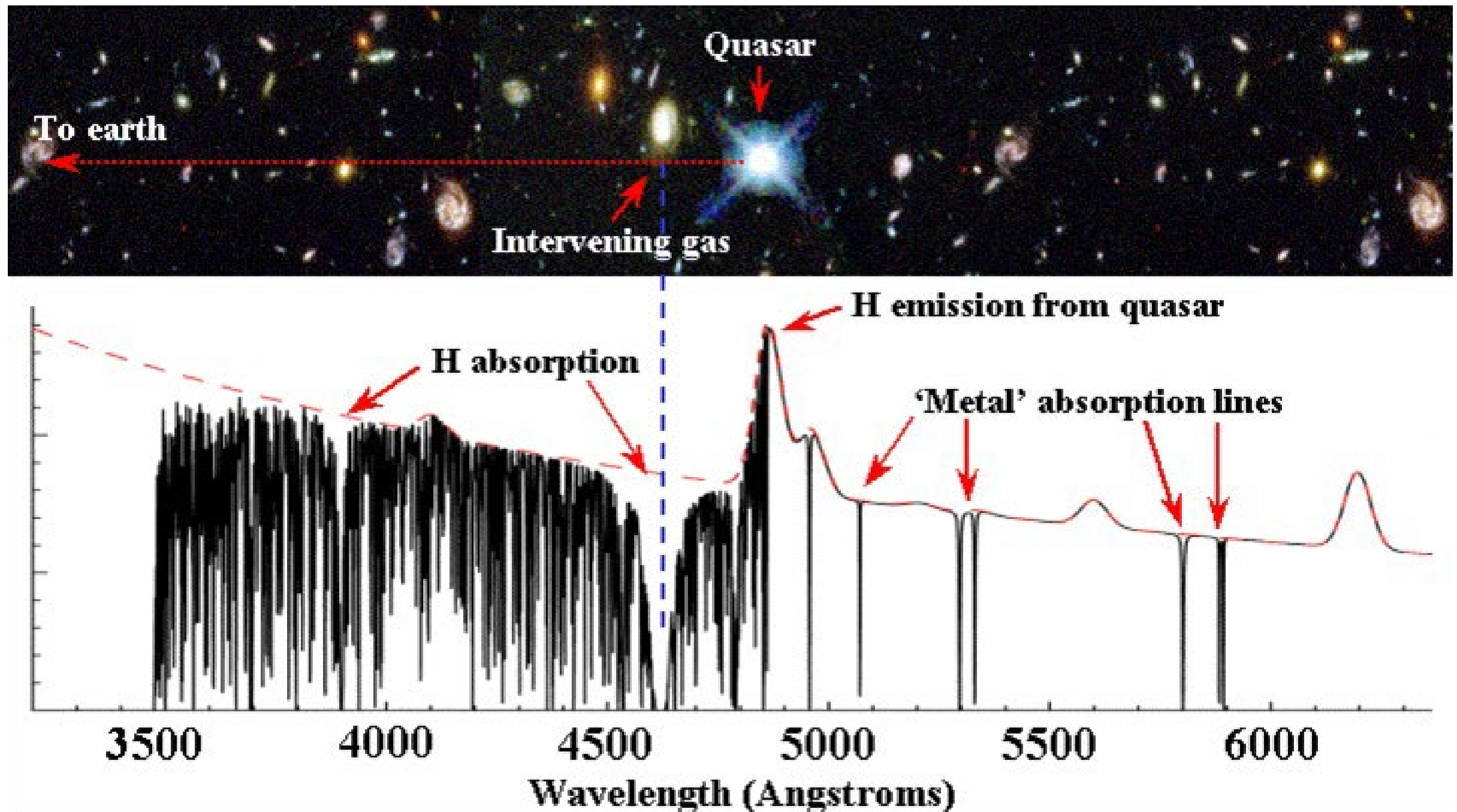


APOSTLE simulations (Sawala+2016)

ARTEMIS simulations (Font+2020)

**But the basic prediction is unchanged...
there should be gazillions of subhaloes**

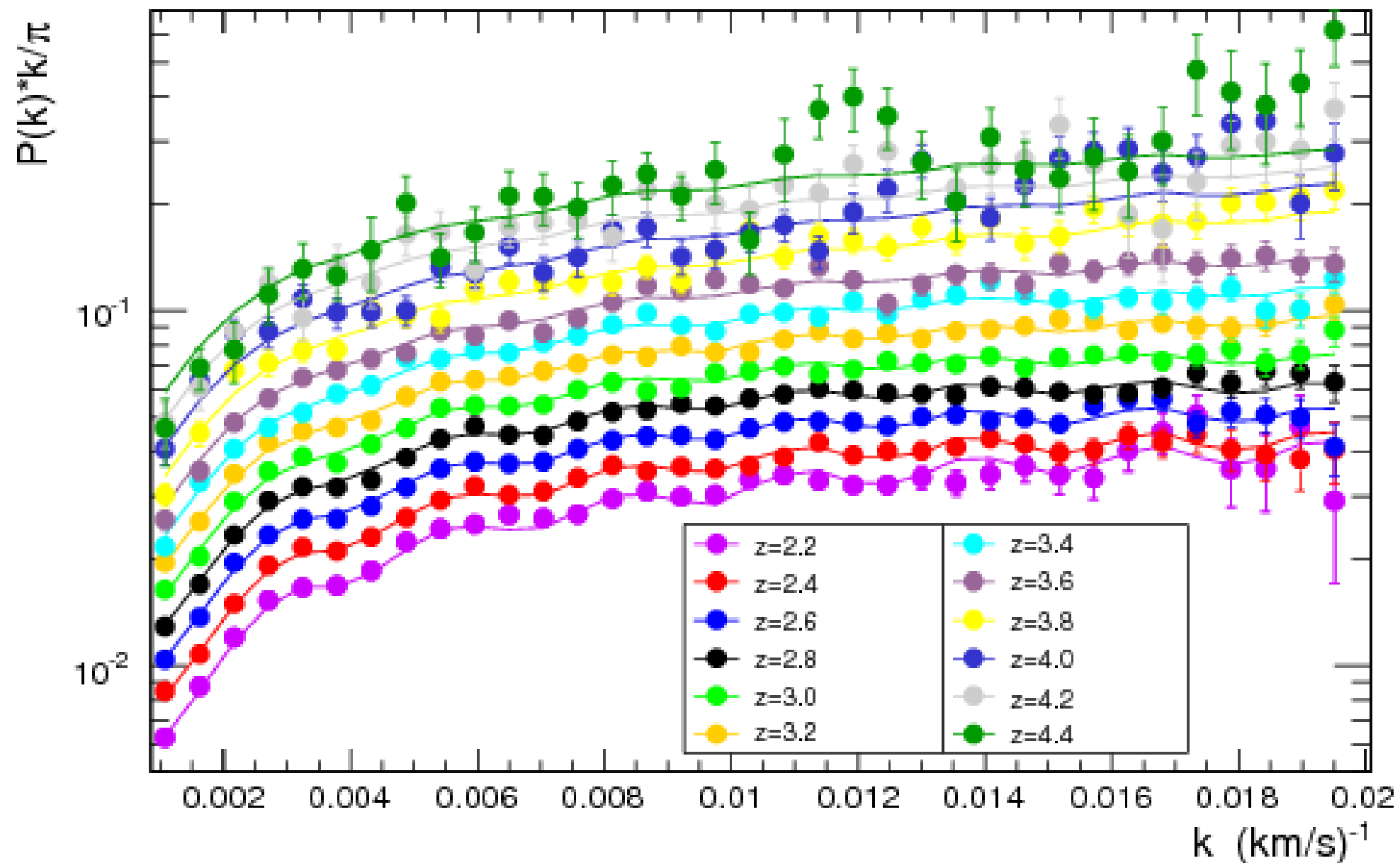
Current observational tests of dark subhaloes



Lyman-alpha forest

Current observational tests of dark subhaloes

(e.g., Viel+2005, Palanque-Delabrouille+2013)



Lyman-alpha forest flux power spectrum → matter power spectrum.
Constraint on nature of dark matter. Complex and limited.

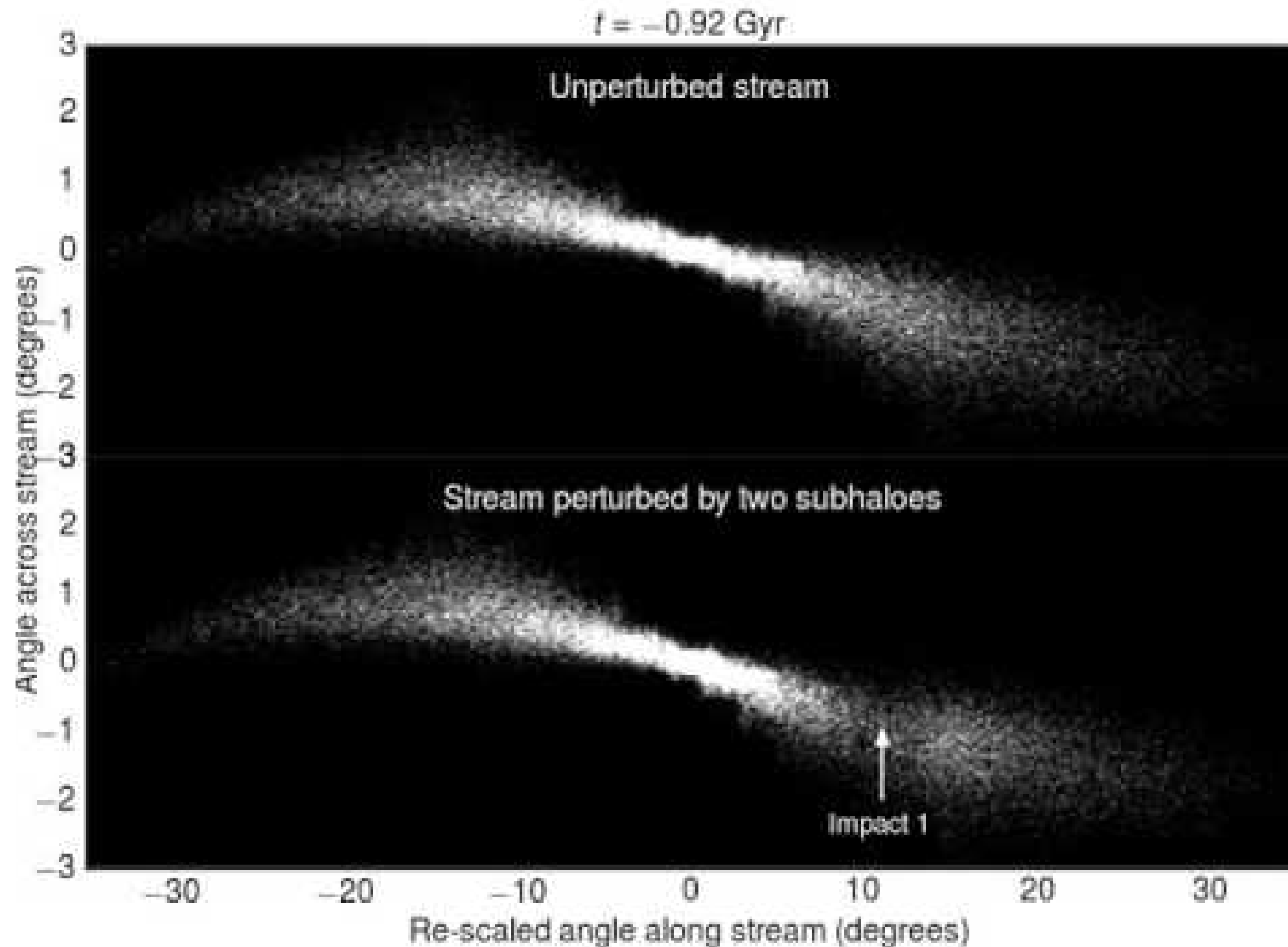
Current observational tests of dark subhaloes



(e.g., Koopmans 2005)

Perturbations of strong lensing arcs

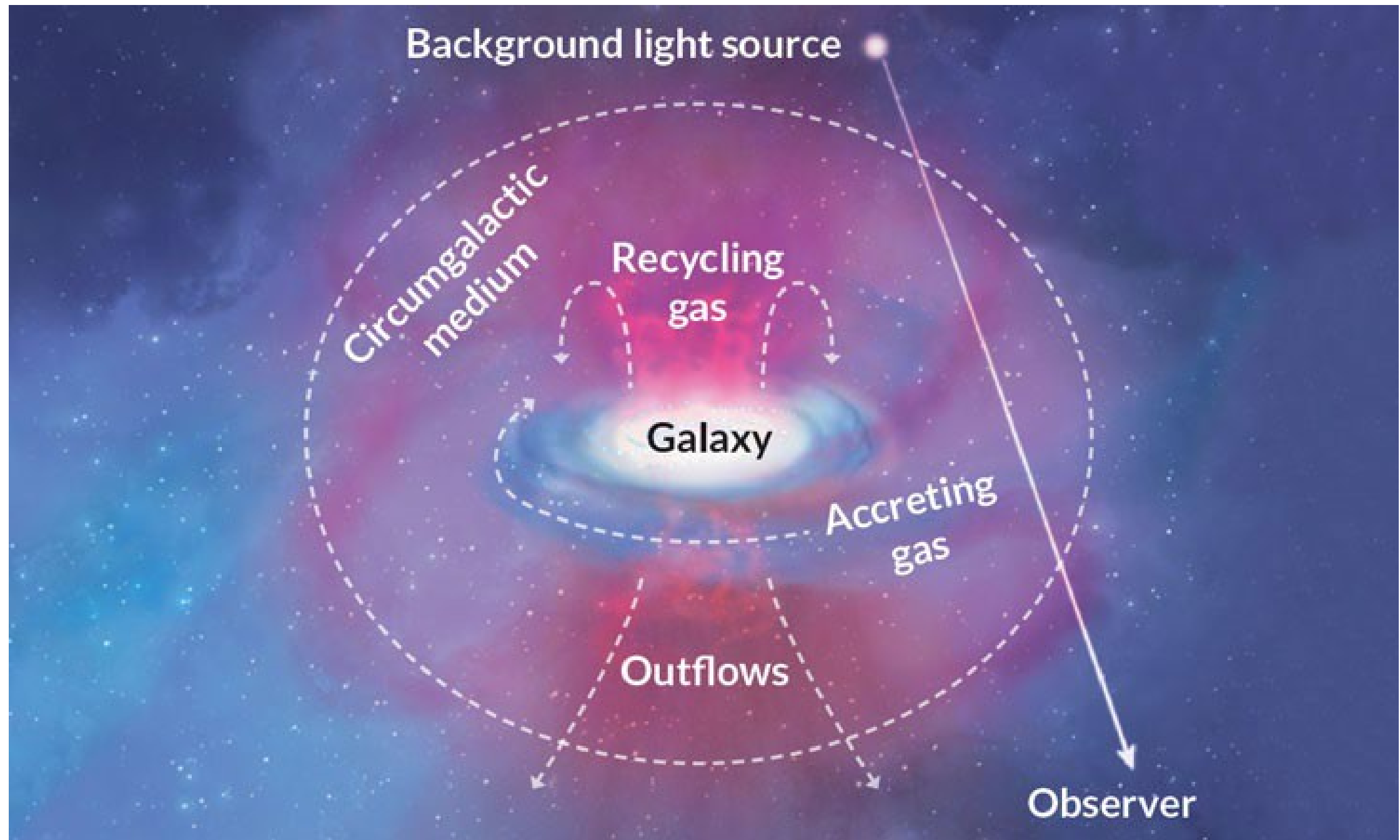
Current observational tests of dark subhaloes



(e.g., Erkal & Belokurov 2015)

Perturbations of stellar streams

Proposal: fluctuations in the CGM



Advantages: most baryons are in hot phase. Volume-filling.

Disadvantages: not easy to observe and CGM is dynamic.

CGM fluctuations: the physics

Gas that is flowing over the top of a subhalo will try to satisfy HSE:

$$\frac{dP_{\text{cgm}}(r)}{dr} = - \frac{GM_{\text{grav}}(< r) \rho_{\text{cgm}}(r)}{r^2}$$

So long as the flow is not supersonic, HSE will be about right. Also, **compression will be adiabatic**. Under these conditions, one can derive the density enhancement profile centred on a subhalo:

$$\frac{\rho_{\text{cgm}}(r)}{\rho_0} = \left[1 + \frac{4}{5} \frac{T_{200}}{T_0} \ln \left(\frac{R}{r} \right) \right]^{3/2} \quad \text{where} \quad \frac{k_B T_{200}}{\mu m_H} = \frac{GM_{200}}{2r_{200}}$$

Depends only on the ratio of the subhalo gravitational temperature (T_{200}) to the CGM temperature (T_0), not on the density of the CGM itself.

Fall off with radius depends on dark matter profile (assumed isothermal here).

Idealised hydrodynamical simulations



Run with new **SWIFT** code (Schaller et al.)

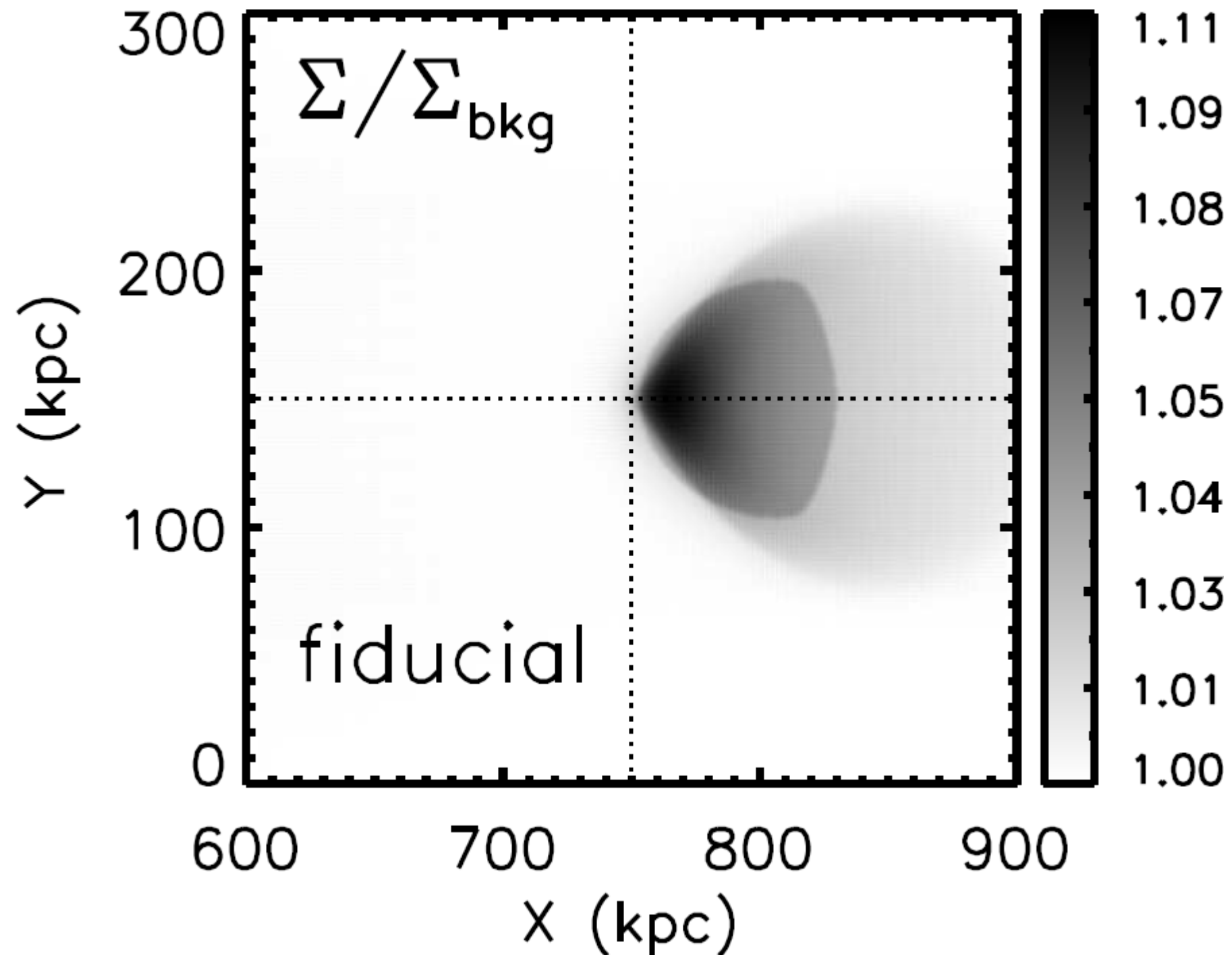
Fiducial setup:

- $M_{\text{sub}} = 10^{10} M_{\text{sun}}$
- $T_{\text{cgm}} = 10^6 \text{ K}$
- $V_{\text{sub}} = 200 \text{ km/s}$
- $n_{\text{cgm}} = 10^{-4} \text{ g/cm}^3$ (doesn't matter)
- Very high resolution!

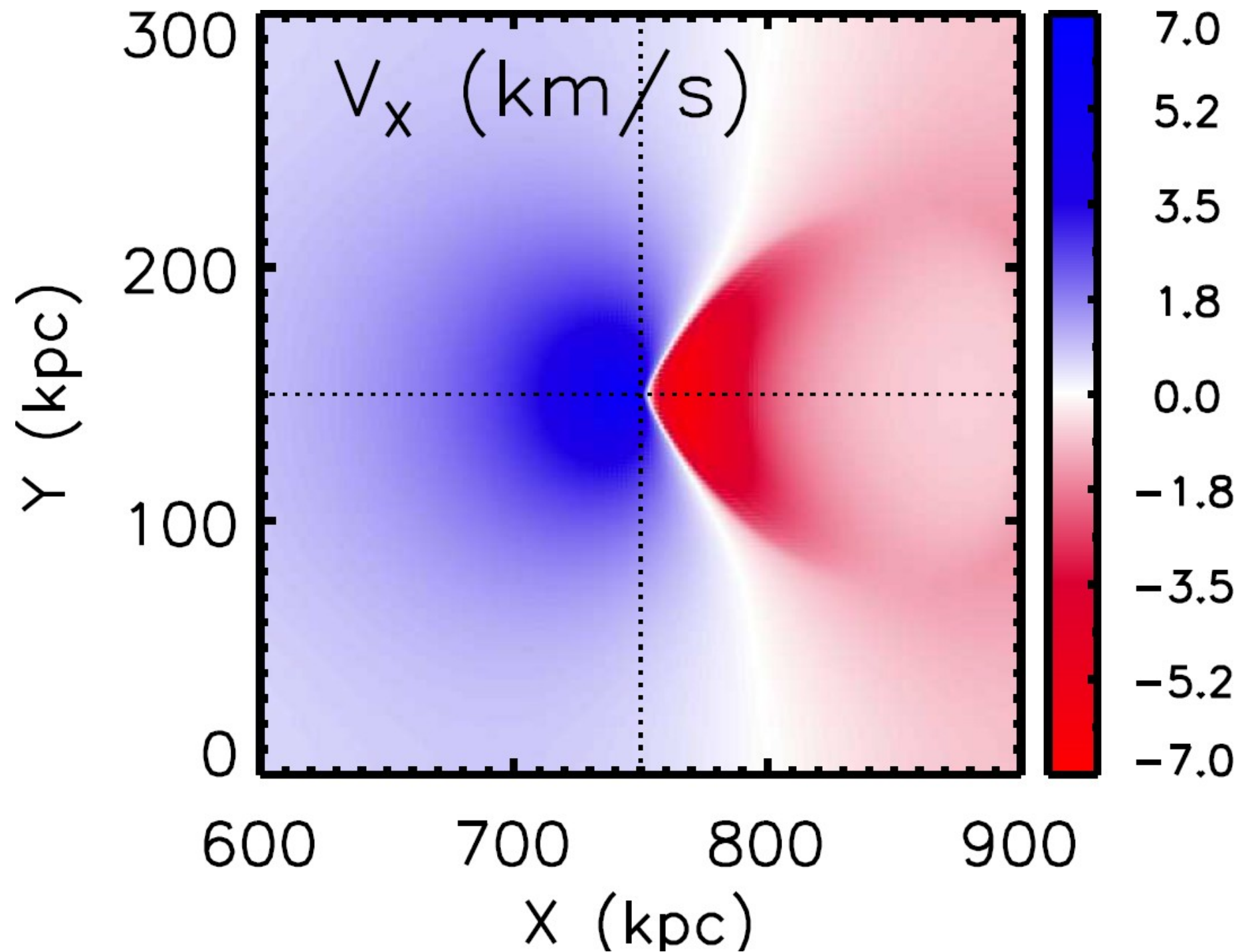
Run on only 32 cores for about 5 hours!

Vary subhalo mass, CGM temperature, relative velocity.

Idealised hydrodynamical simulations



Idealised hydrodynamical simulations

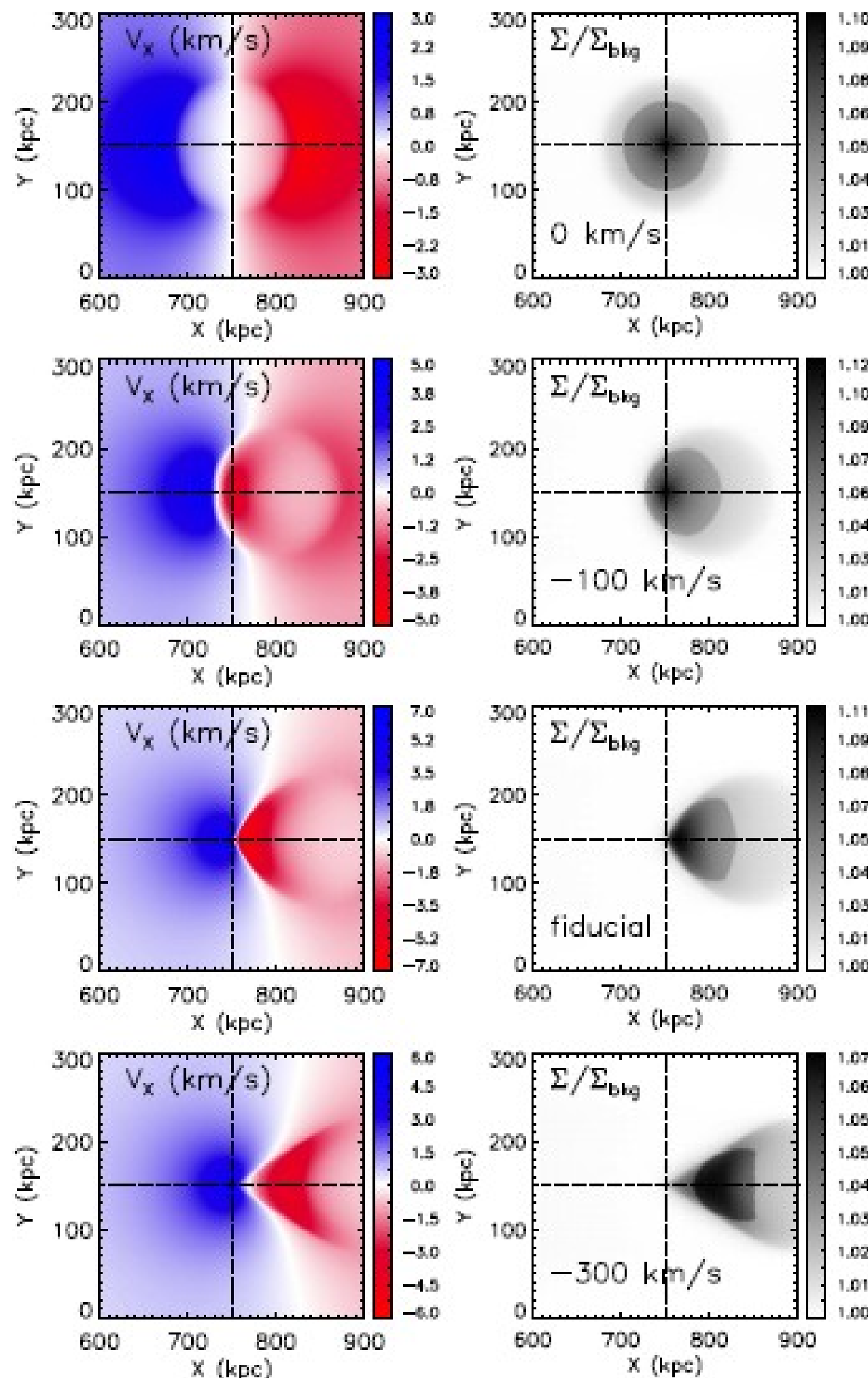


Varying the relative velocity

Faster the motion the larger the offset is between the subhalo and the CGM enhancement.

Morphology changes.

Amplitude of the effect is only mildly altered.

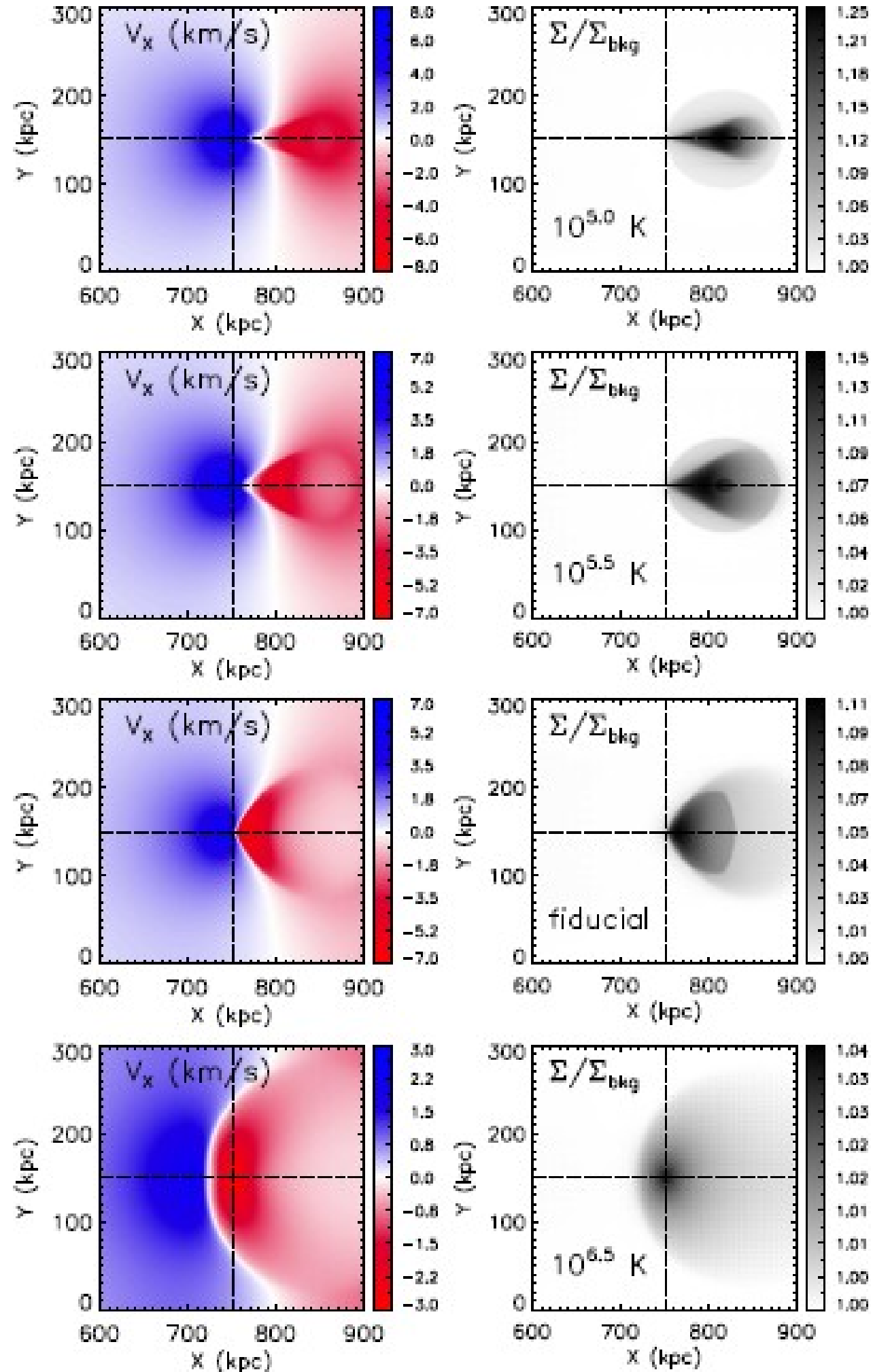


Varying the CGM temperature

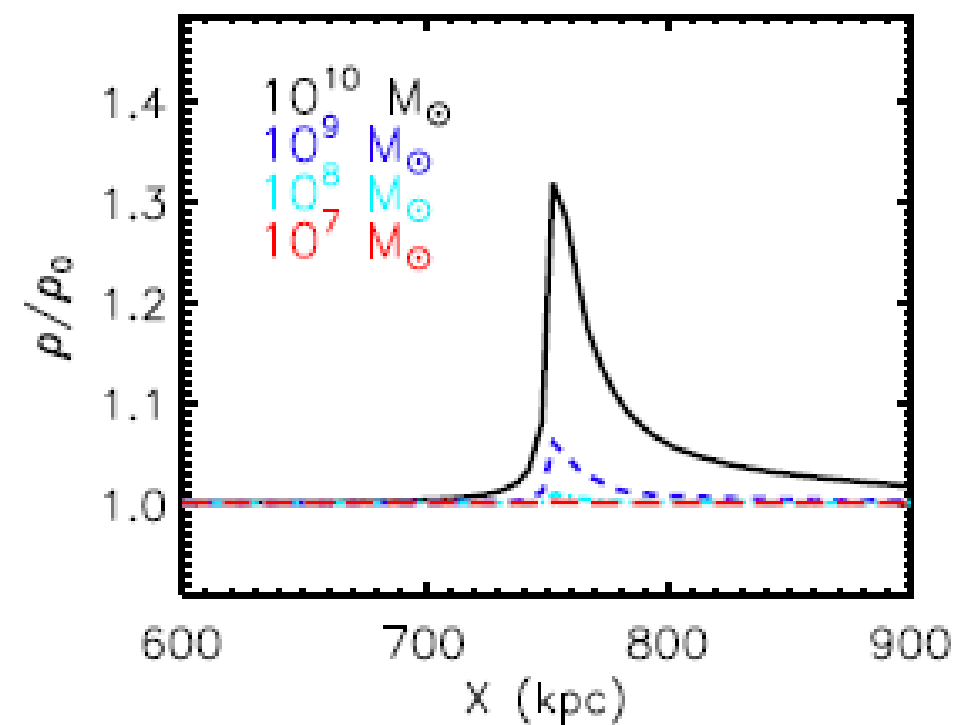
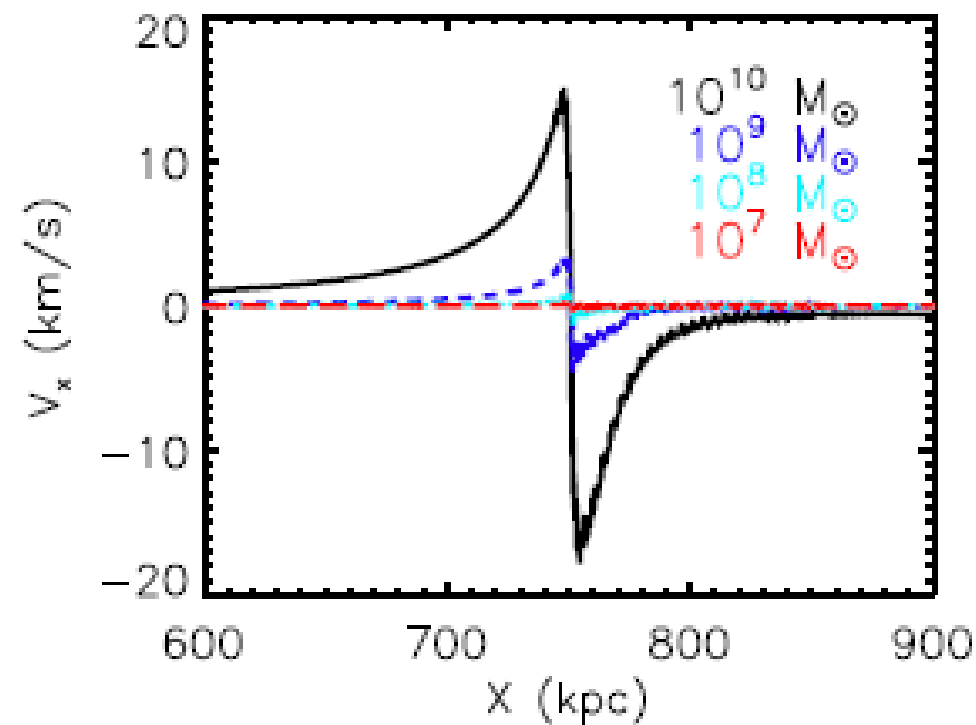
Cooler the CGM the larger the offset is between the subhalo and the CGM enhancement.

Morphology changes.

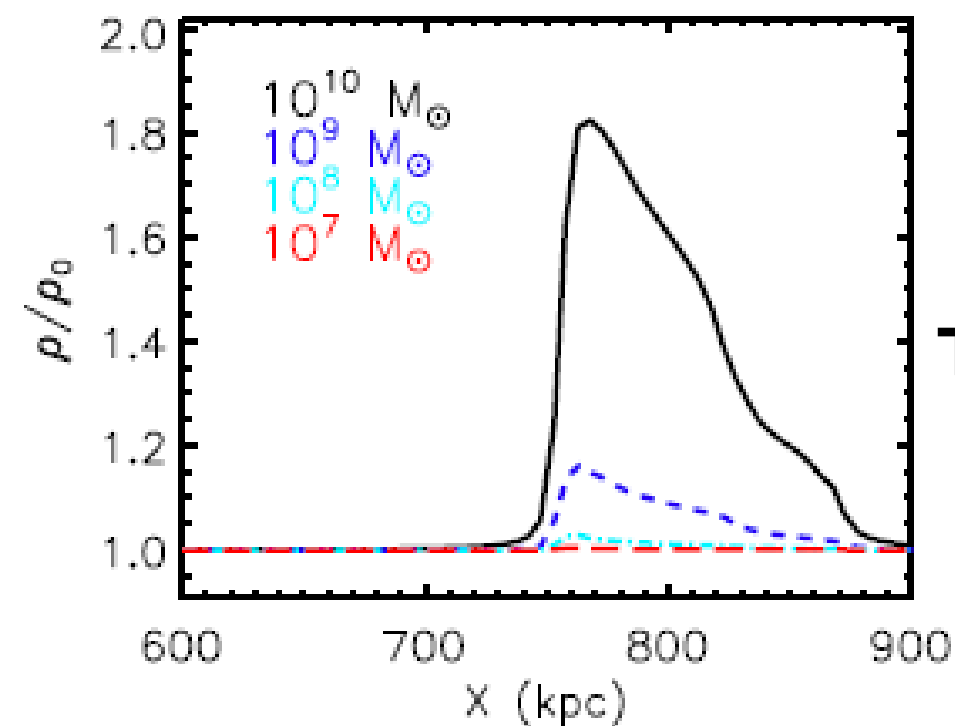
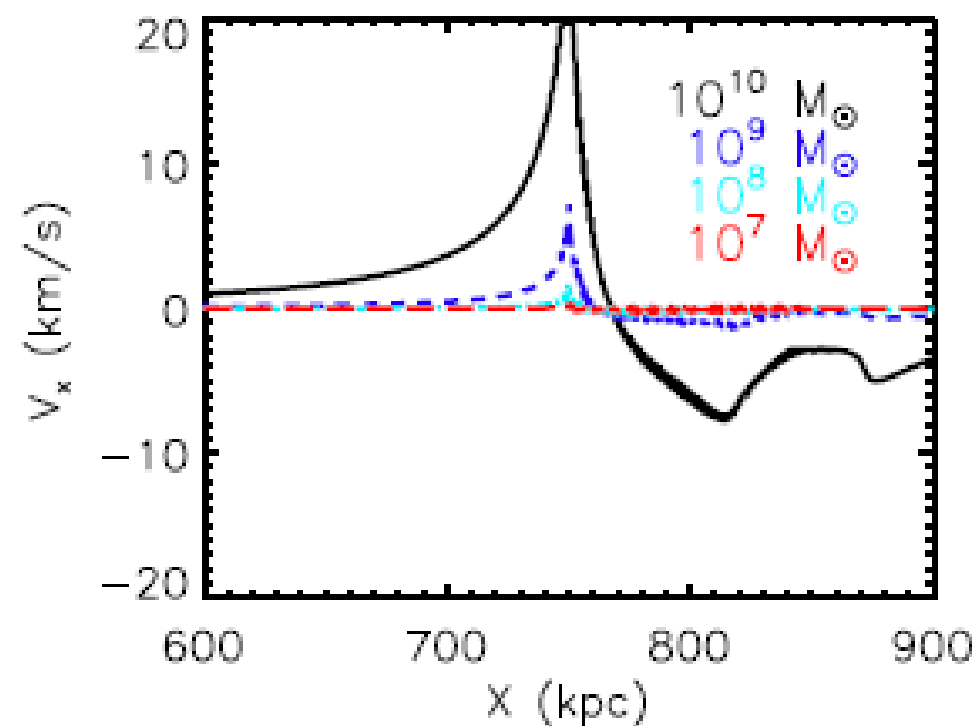
Amplitude of the effect is strongly altered (as expected based on analytics). Cooler CGM is more easily compressed.



Idealised hydrodynamical simulations



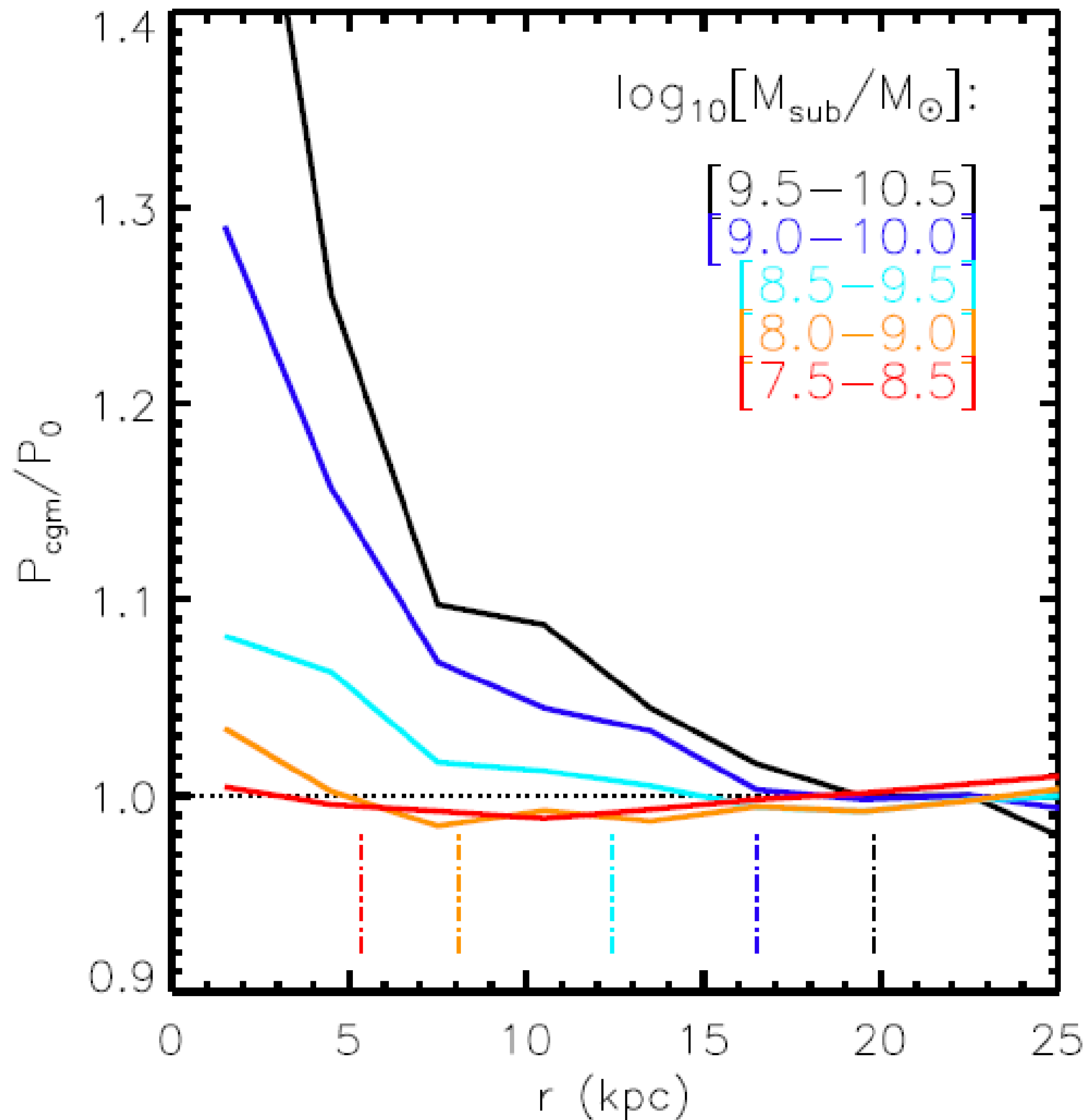
$T_{\text{cgm}} = 10^6 \text{ K}$



$T_{\text{cgm}} = 10^5 \text{ K}$

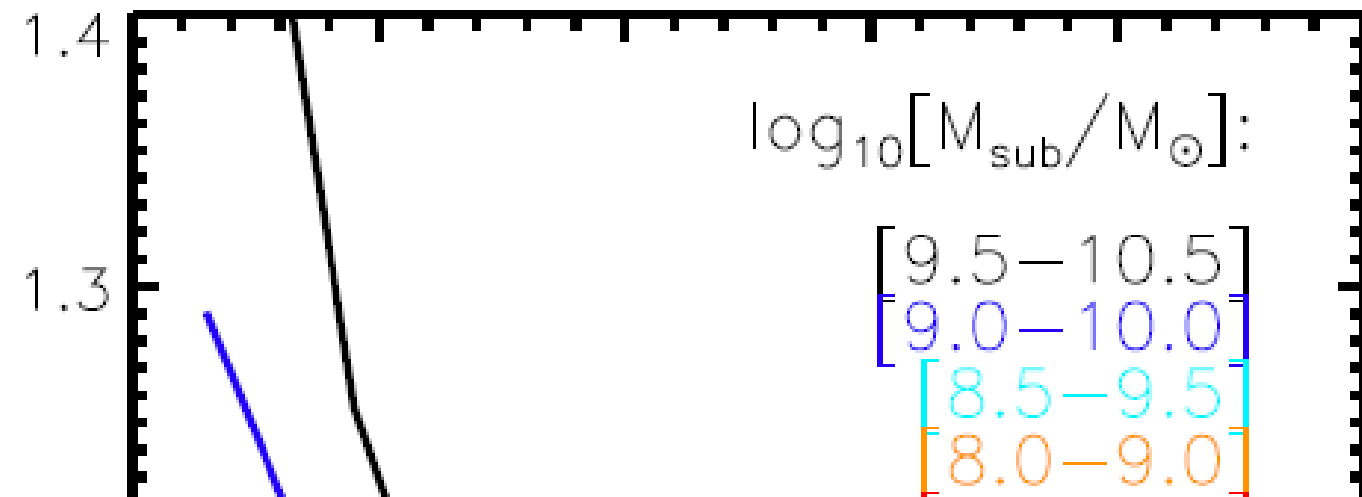
Full cosmological simulations: ARTEMIS

CGM pressure profiles centred on dark subhaloes of ARTEMIS host haloes

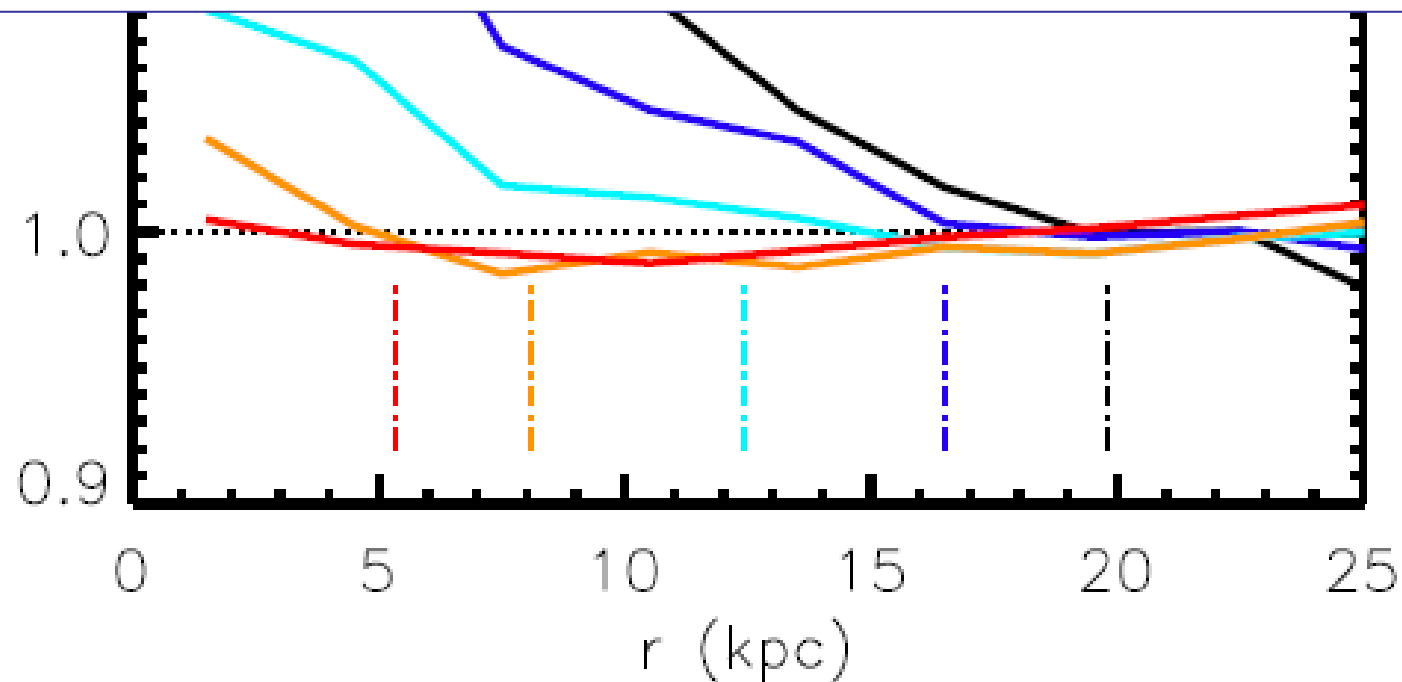


Full cosmological simulations: ARTEMIS

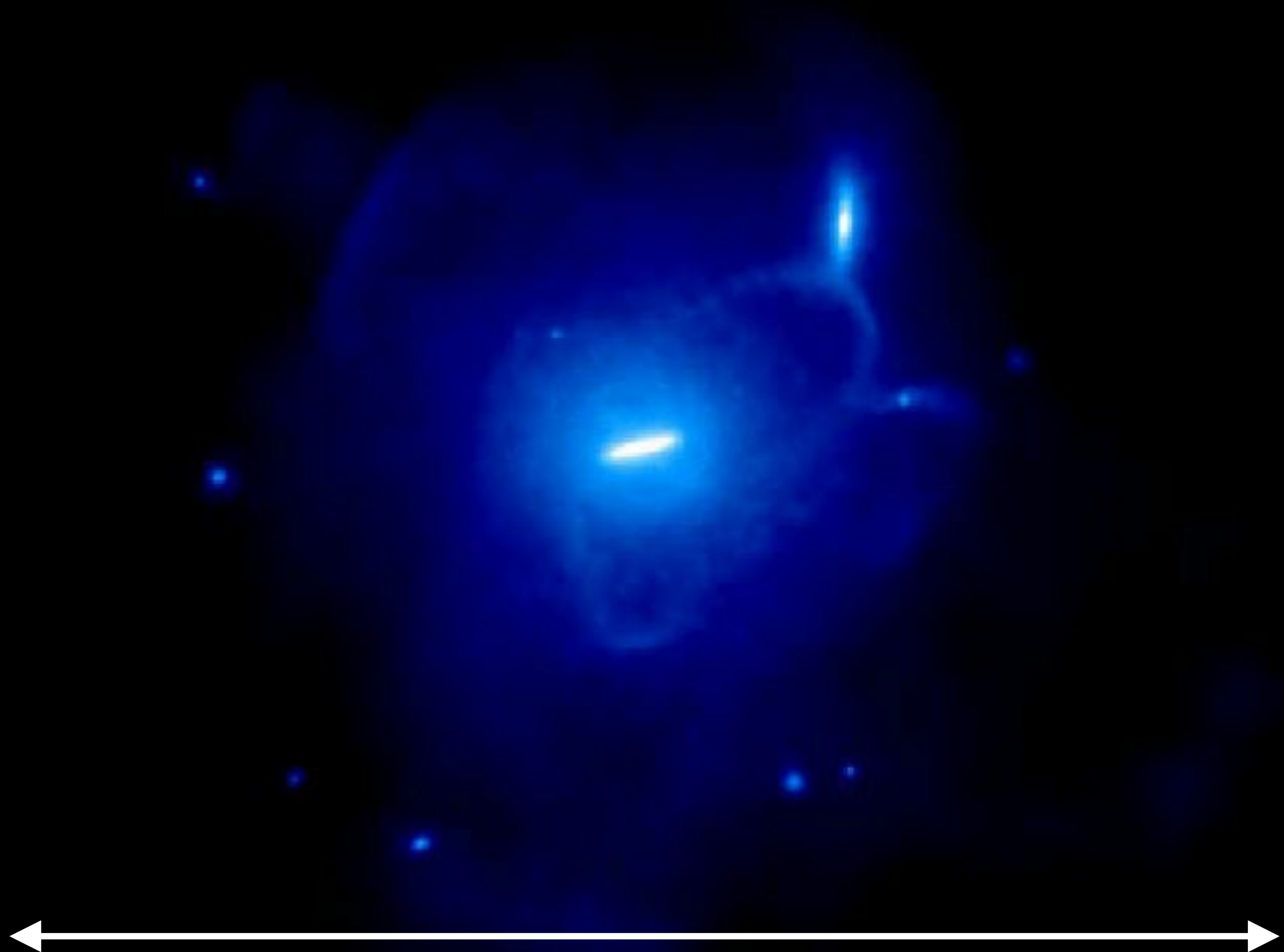
CGM pressure profiles centred on dark subhaloes of ARTEMIS host haloes



Great it exists, but how do we observe this in real data?



Stellar mass image of ARTEMIS halo



400 kpc

IGM, in prep

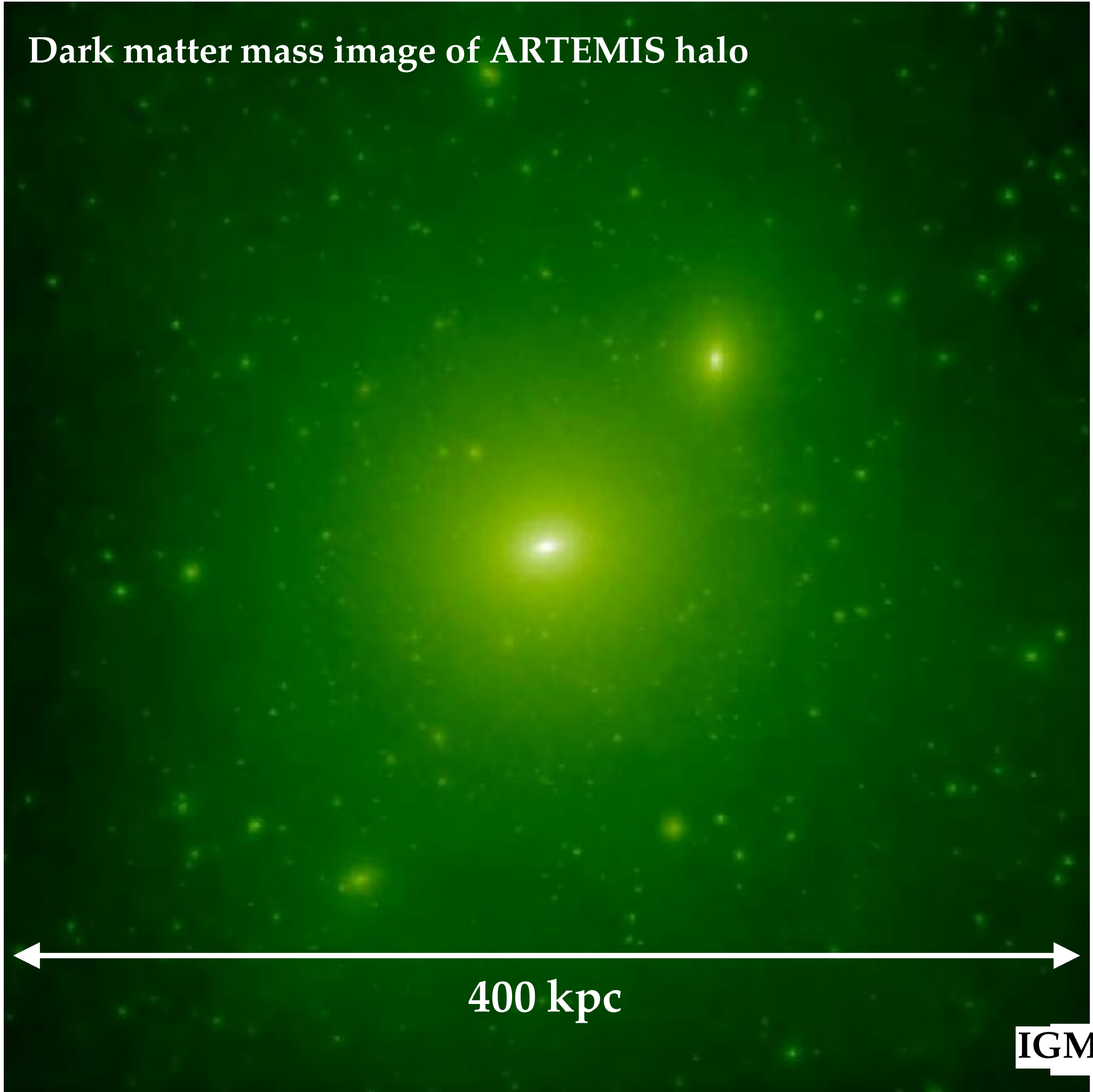
CGM mass image of ARTEMIS halo



400 kpc

IGM, in prep

Dark matter mass image of ARTEMIS halo



400 kpc

IGM, in prep

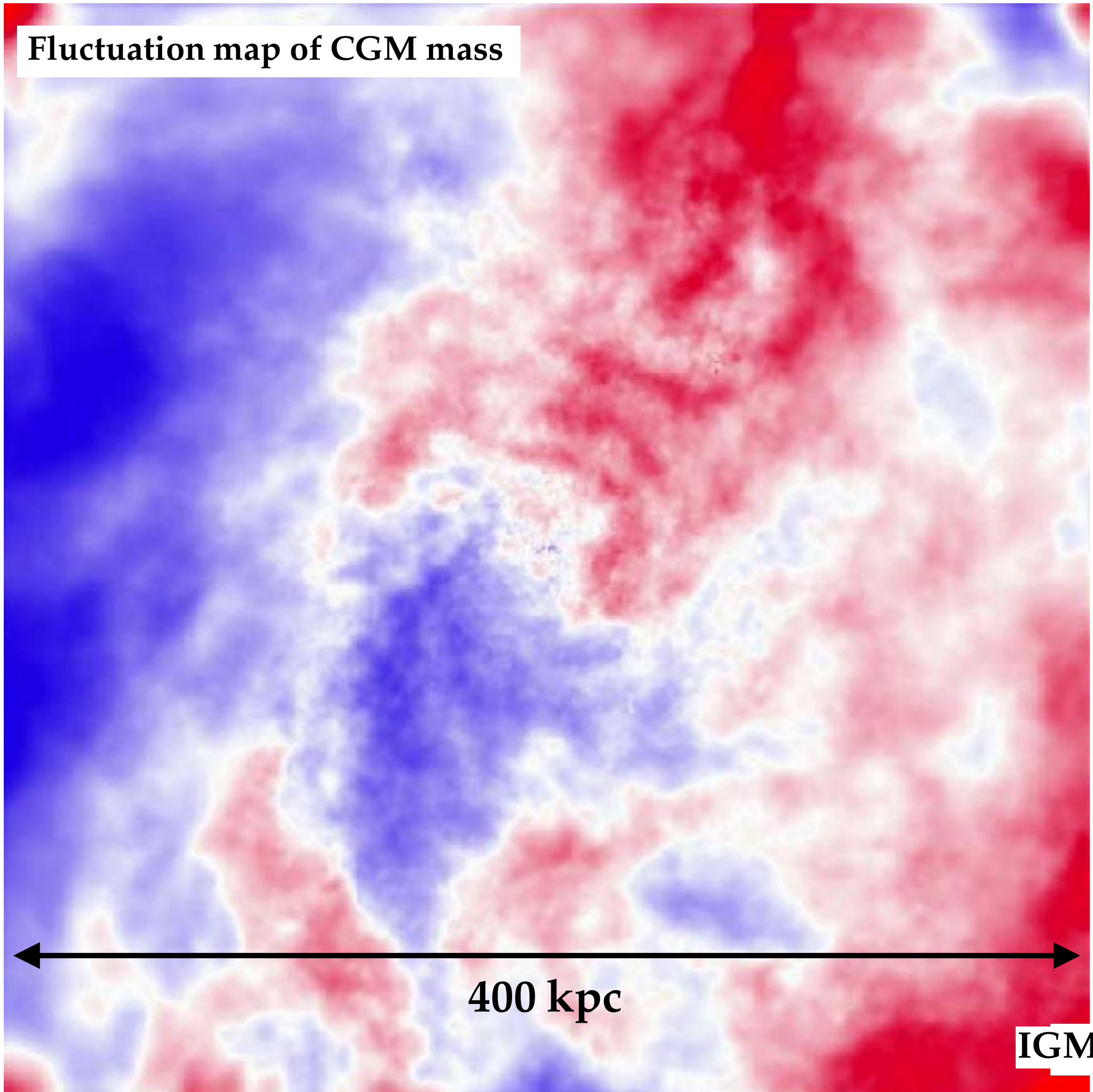
CGM mass image of ARTEMIS halo



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IGM, in prep

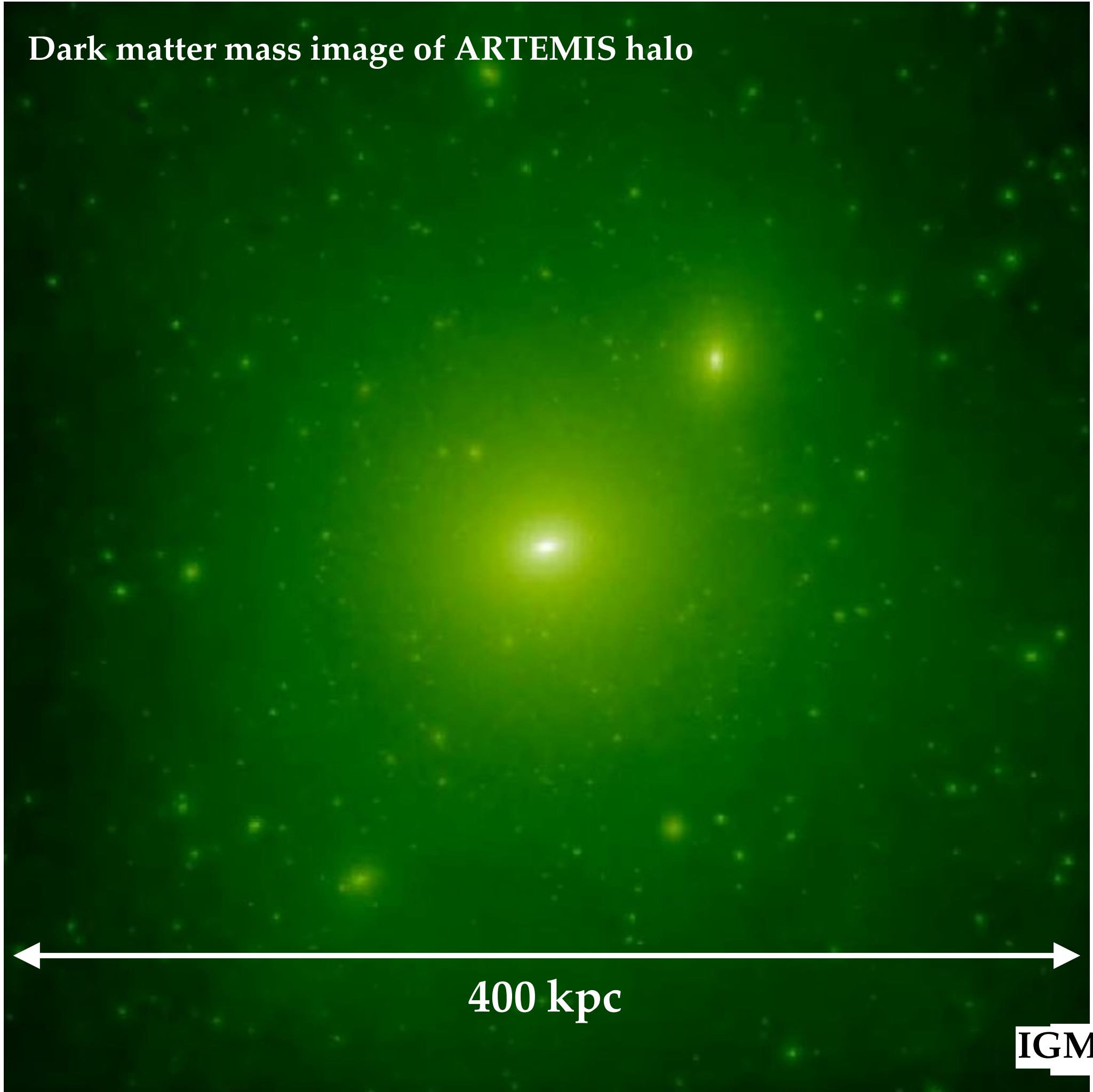
Fluctuation map of CGM mass



400 kpc

IGM, in prep

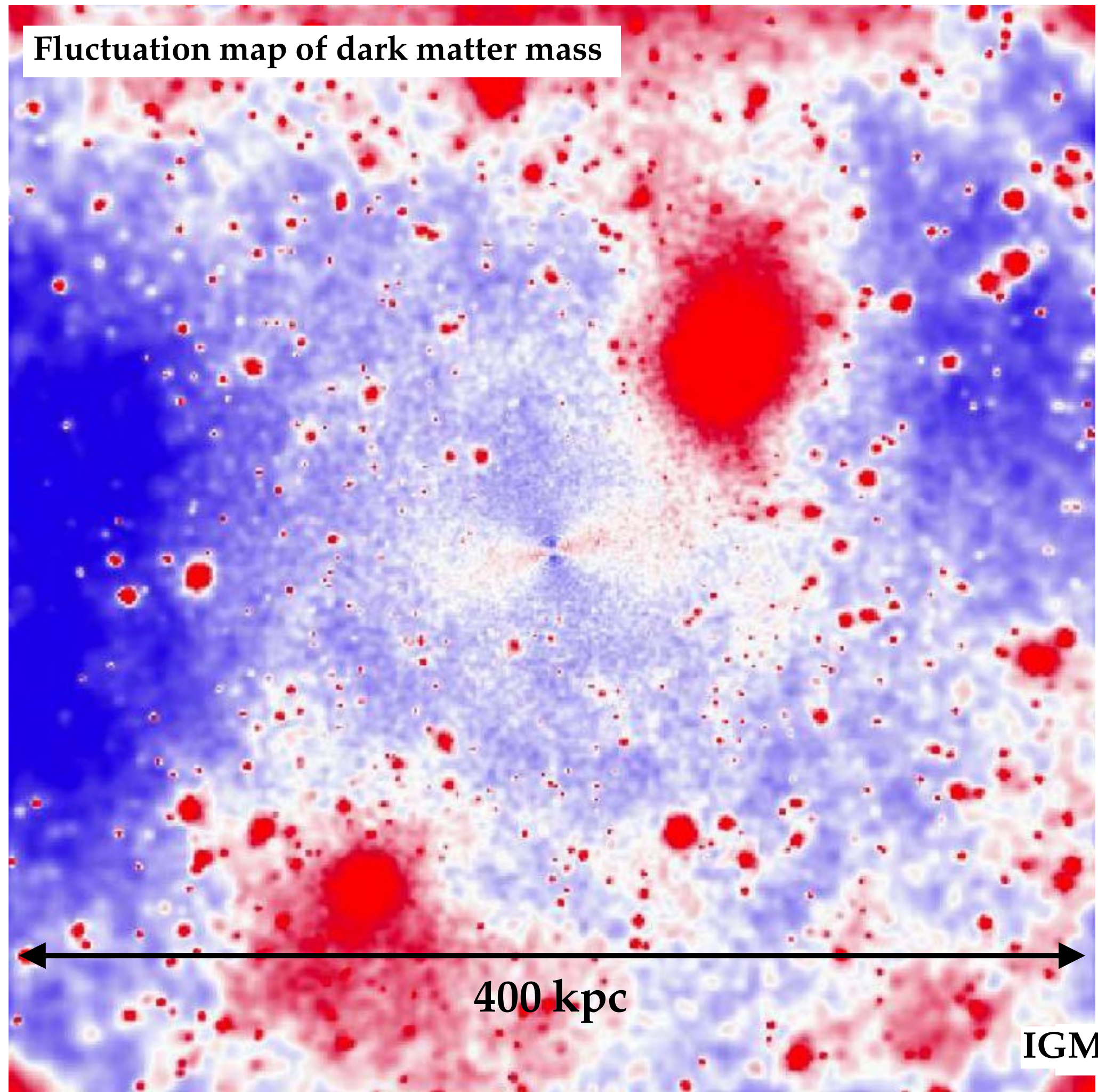
Dark matter mass image of ARTEMIS halo



400 kpc

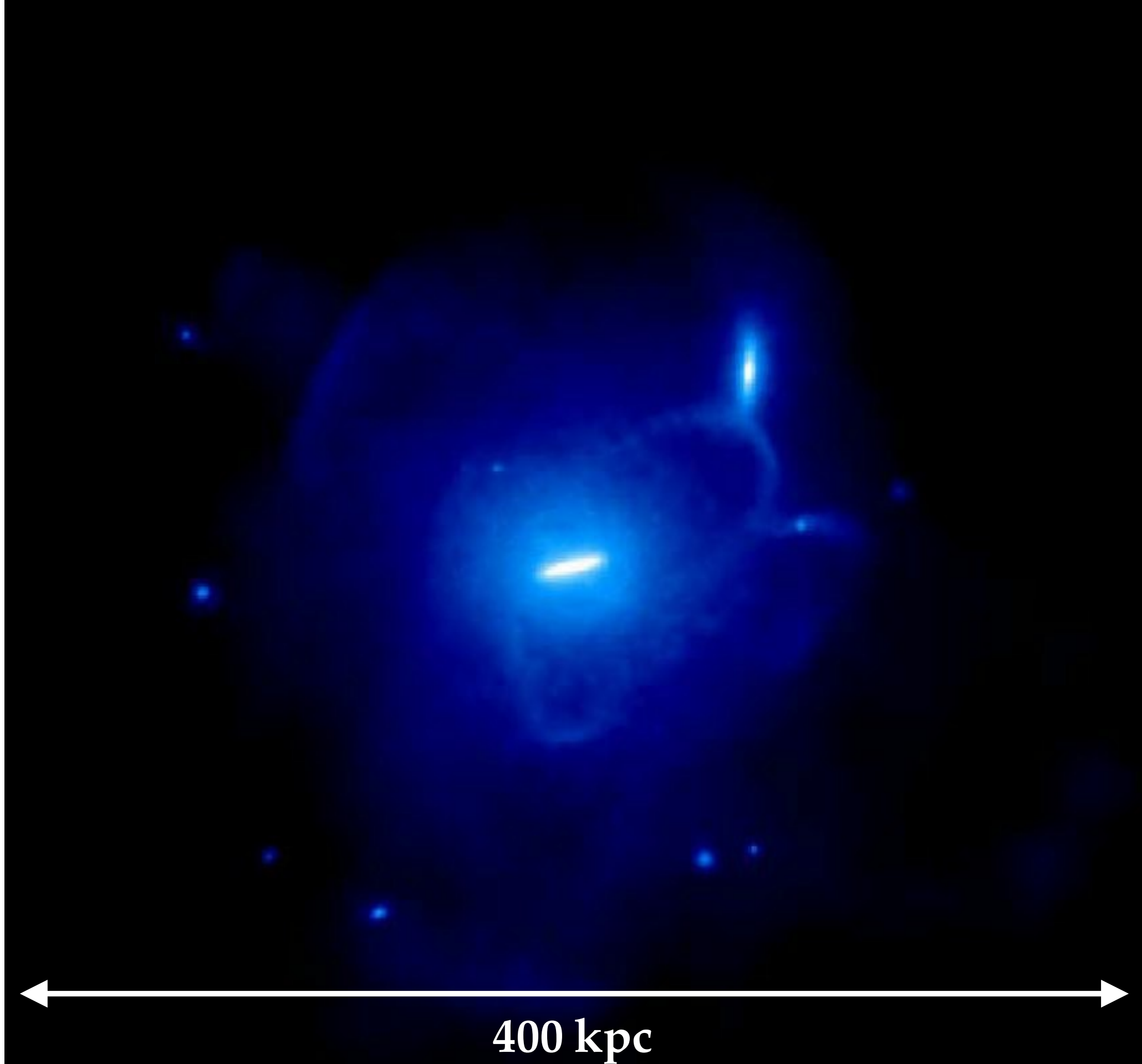
IGM, in prep

Fluctuation map of dark matter mass



400 kpc

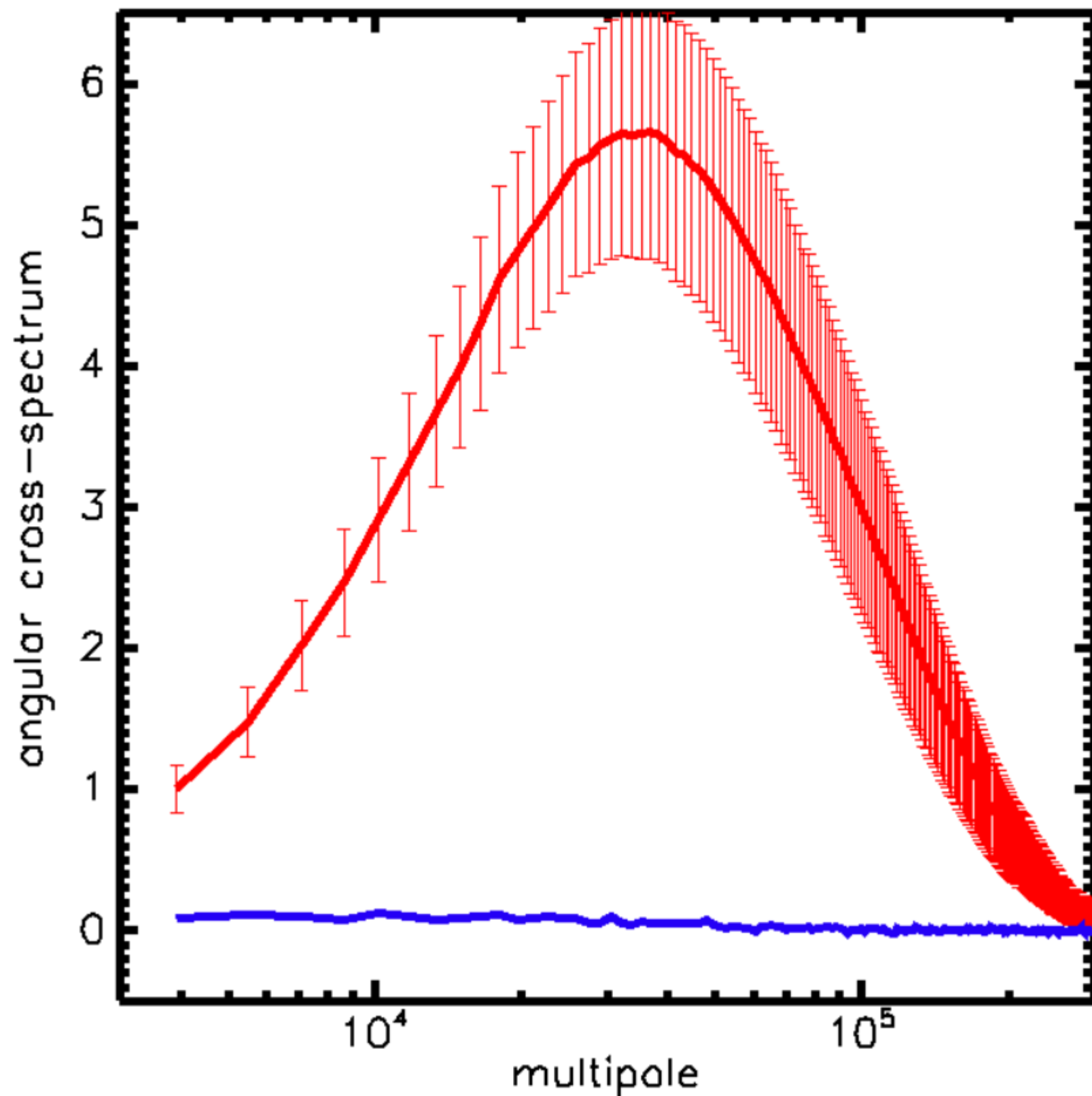
IGM, in prep



400 kpc

IGM, in prep

CGM-dark matter angular cross-spectrum



IGM, in prep

Conclusions (part 2)

- If Λ CDM is correct, there should be countless dark matter subhaloes around galaxies like the Milky Way.
- Galaxy formation (feedback) provides a plausible explanation for why most of them are dark. **But how can we distinguish between being present and dark or just not being present?**
- Recent searches for dark subhaloes have relied on stellar mass/cool gas. Requires a bit of fortune.
- **New proposal:** CGM has most of the mass and is volume-filling.
- Effect is indeed present at few percent up to tens of percent (in density).
- Observations that probe the CGM (X-ray, tSZ, absorption/emission line studies, and radio dispersion) can potentially measure the fluctuation power spectrum.