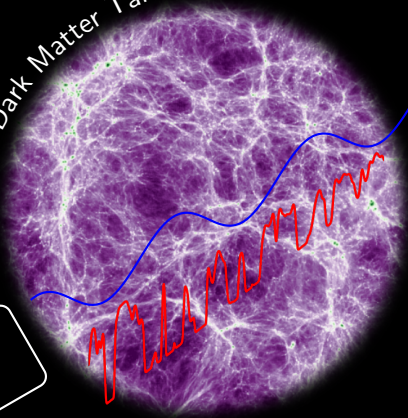


Small scale structure of the IGM:

A Dark Matter Tale



Vid Iršič

© University of Washington

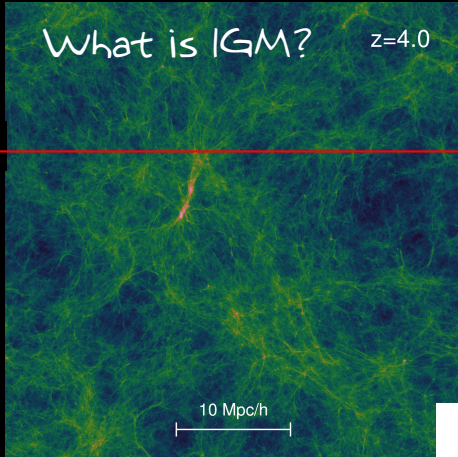


INPA Seminar

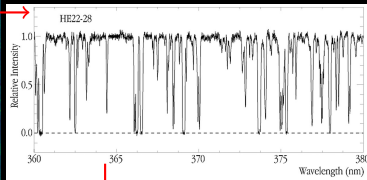
©
Lawrence Berkeley National Laboratory

May 24, 2019

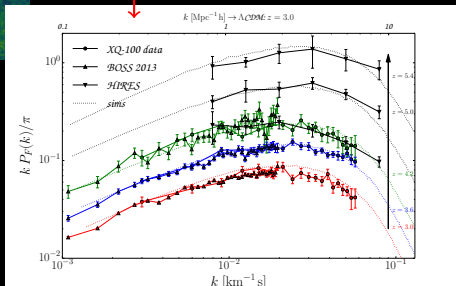
What is IGM? $z=4.0$



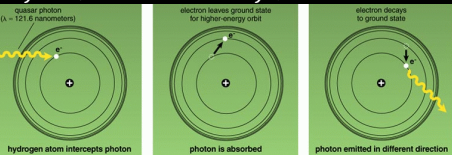
Absorption in Quasar spectra along the line of sight



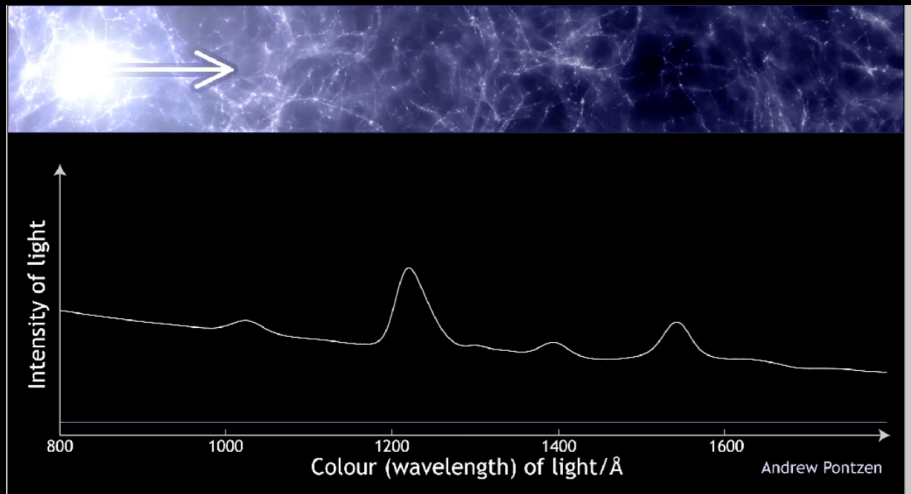
Flux power spectrum



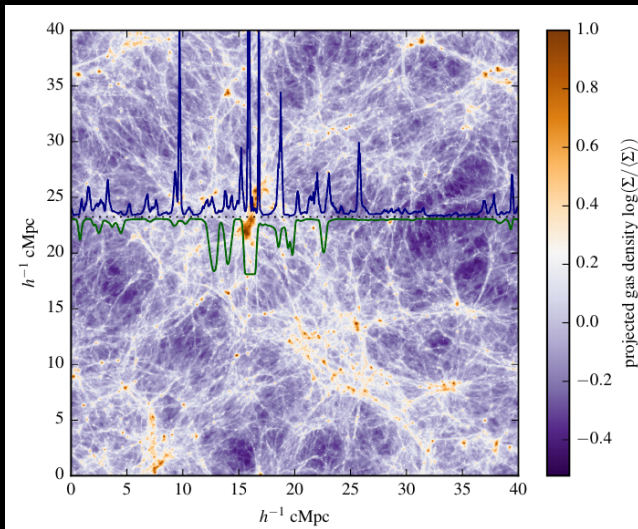
Scattering of the photon on $n = 1 \rightarrow n = 2$ Hydrogen transition (Lyman series)



Observable of the IGM



Why should we care about IGM?



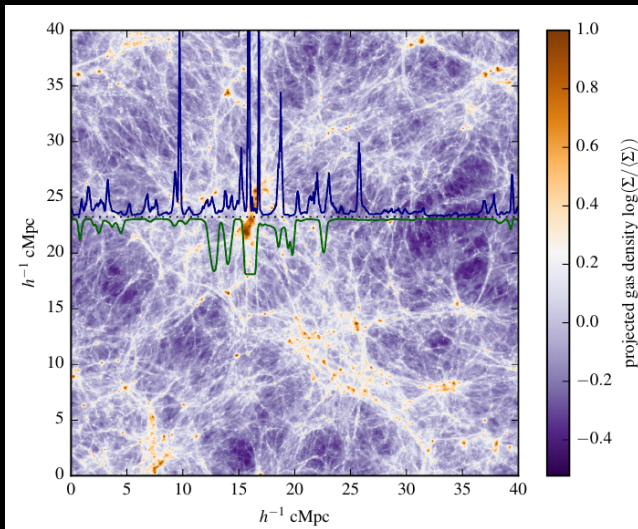
ρ_g - Gas density
 F - Observed flux

Flux is a Biased
tracer of the density:

$$F \sim b \rho_g$$

Sensitive to fluctuations, along the line-of-sight, on scales $\sim 0.1 - 10$ Mpc/h

Why should we care about IGM?



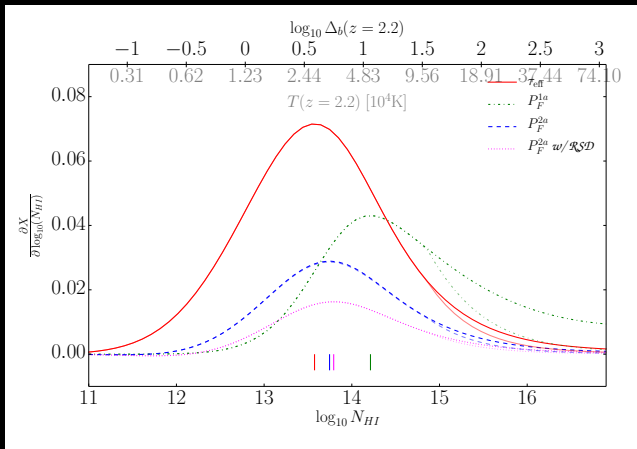
ρ_g - Gas density
 F - Observed flux

Flux is a Biased
tracer of the density:

$$F \sim b \rho_g$$

Sensitive to density fluctuations, along the line-of-sight, on **scales $\sim 0.1 - 10$ Mpc/h**
small scales

Why should we care ABOUT IGM?



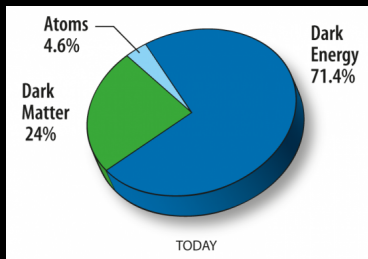
ρ_g - gas density
 F - observed flux

Flux is a biased tracer of the density:

$$F \sim b \rho_g$$

Sensitive to mildly non-linear density fluctuations, along the line-of-sight, on scales $\sim 0.1 - 10 \text{ Mpc}/h$

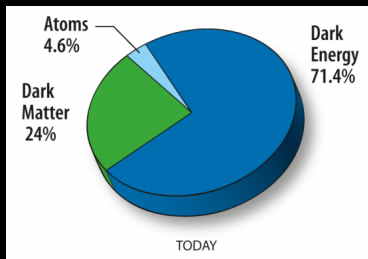
Cold Dark Matter problems (?)



Cold Dark Matter (CDM):

heavy, non-interactive particle(s) → WIMPs

Cold Dark Matter problems (?)



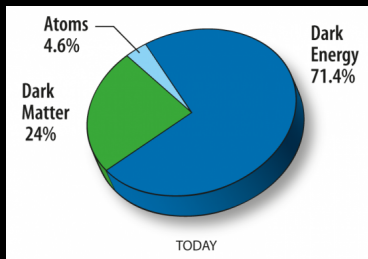
Cold Dark Matter (CDM):

heavy, non-interactive particle(s) → WIMPs

CDM problems of small-scale physics:

- ▶ Missing satellites
- ▶ Core/Cusp problem
- ▶ ...

Cold Dark Matter problems (?)



Cold Dark Matter (CDM):

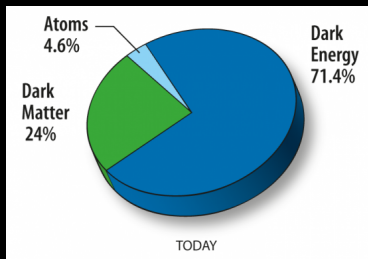
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CDM problems of small-scale physics:

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- ▶ Core/Cusp problem
- ▶ ...

→ Alternative DM models
(Warm DM, Fuzzy DM,
Self-interacting DM, ...)

Cold Dark Matter problems (?)



Baryonic physics

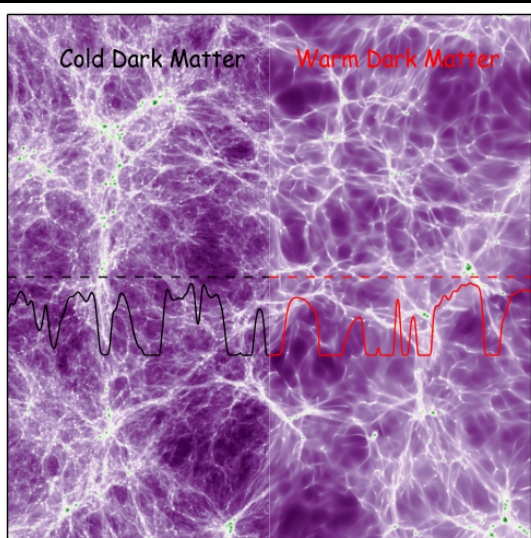
Cold Dark Matter (CDM):
heavy, non-interactive particle(s) \rightarrow WIMPs

CDM problems of small-scale physics:

- ▶ Missing satellites
- ▶ Core/Cusp problem
- ▶ ...

\rightarrow Alternative DM models
(Warm DM, Fuzzy DM,
Self-interacting DM, ...)

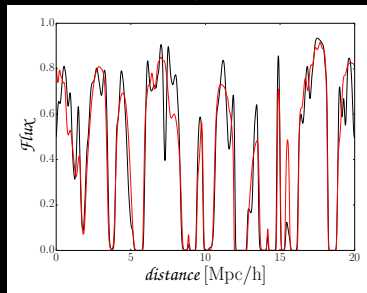
Non-CDM erases small scale structure



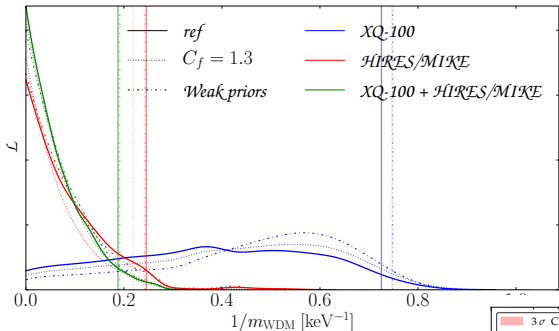
Warm Dark Matter (WDM):
Free-streaming of DM particles
(From the time they decouple
until they become non-relativistic)

Fuzzy Dark Matter (FDM):
de-Broglie wavelength
of ultra-light DM scalar
⇒ erases small scale structure

Typical $\lambda_{\text{FS}} \sim \text{Mpc}/h$



- Typical DM particle mass from local small-scale structure
- $m_{\text{WDM}} \sim 2 - 3 \text{ keV}$ (WDM)
- $m_{\text{FDM}} \sim 1 - 10 \times 10^{-22} \text{ eV}$ (FDM)



'realistic' thermal history:

$$\rightarrow m_{\text{FDM}} > 37 \times 10^{-22} \text{ eV @ } 2\sigma$$

conservative thermal history model:

$$\rightarrow m_{\text{FDM}} > 20 \times 10^{-22} \text{ eV @ } 2\sigma$$

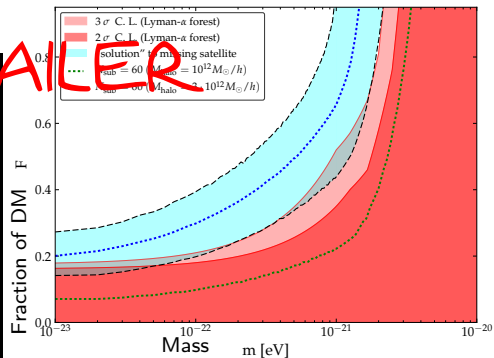
TRAILER

'realistic' thermal history:

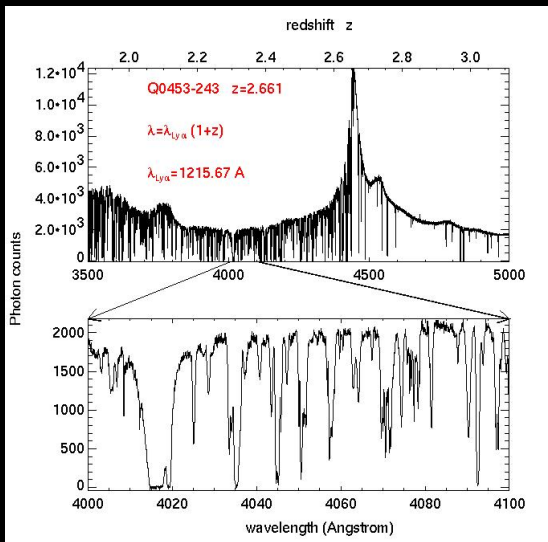
$$\rightarrow m_{\text{WDM}} > 5.3 \text{ keV @ } 2\sigma$$

conservative thermal history model:

$$\rightarrow m_{\text{WDM}} > 3.5 \text{ keV @ } 2\sigma$$



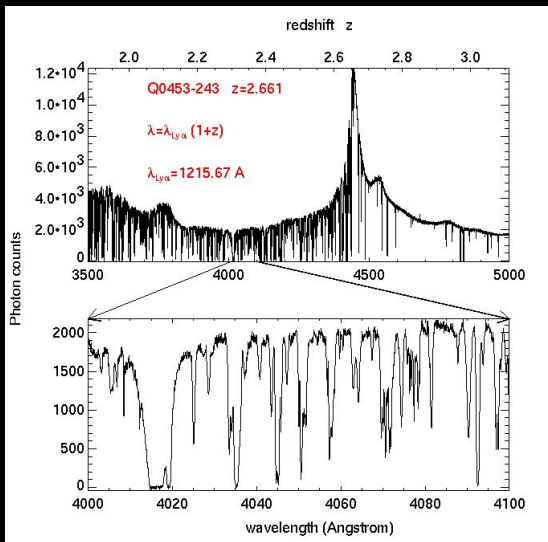
What do we measure?



Measured quantity:

$$f = C \cdot \bar{F} (1 + \delta_F) + n$$

What do we measure?

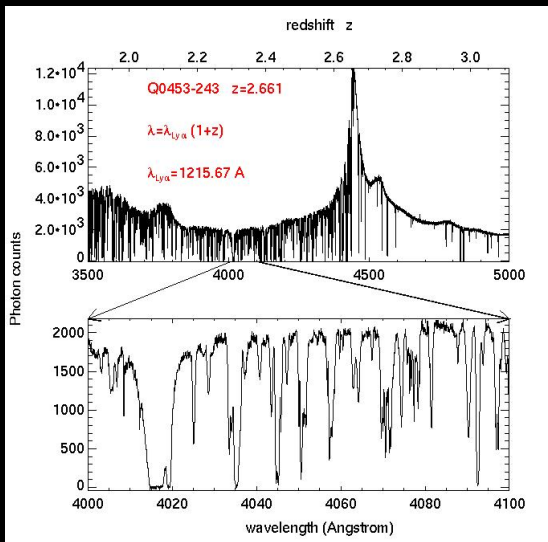


Measured quantity:

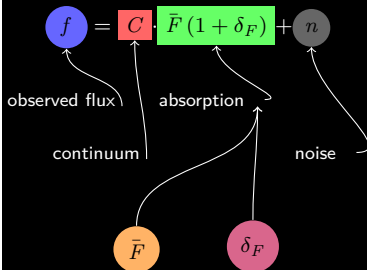
$$f = C \cdot \bar{F} (1 + \delta_F) + n$$

observed flux \leftarrow continuum \leftarrow absorption \leftarrow noise

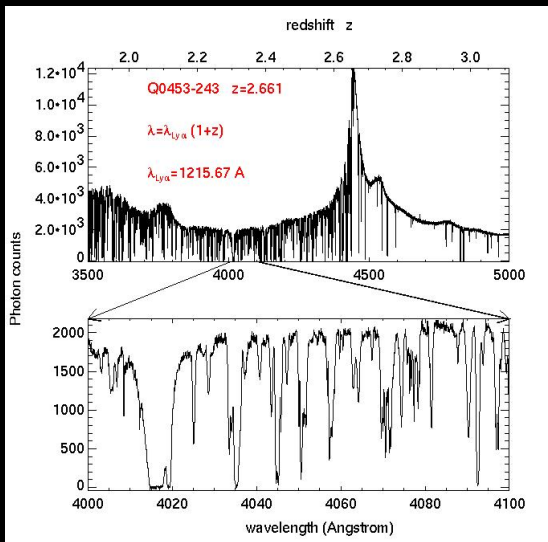
What do we measure?



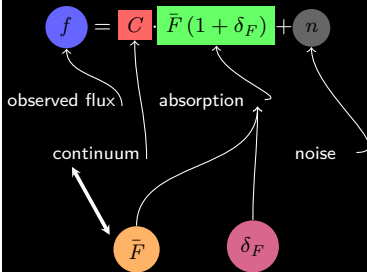
Measured quantity:



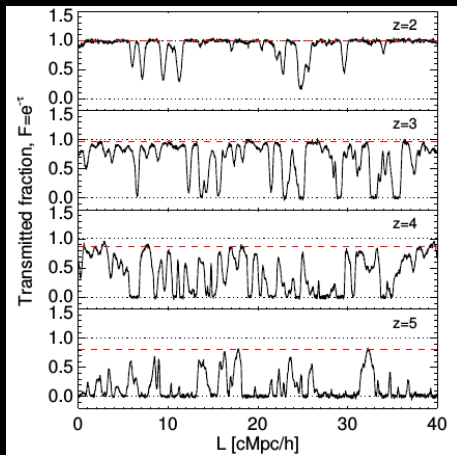
What do we measure?



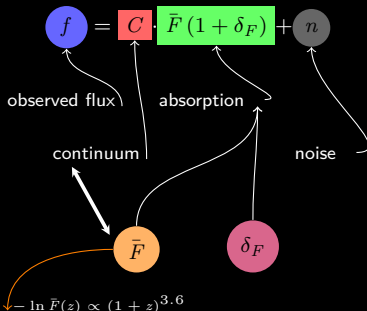
Measured quantity:



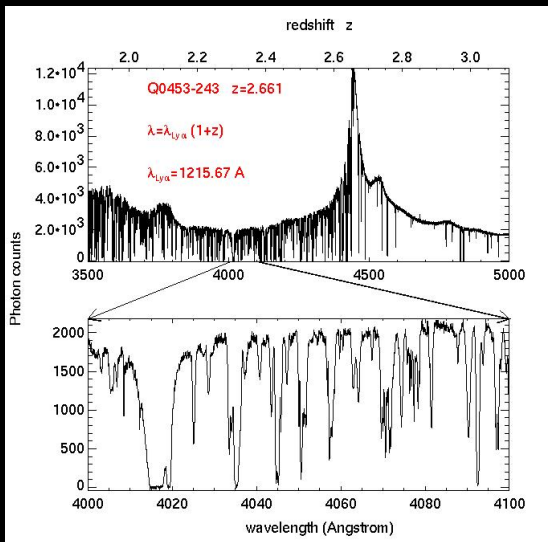
What do we measure?



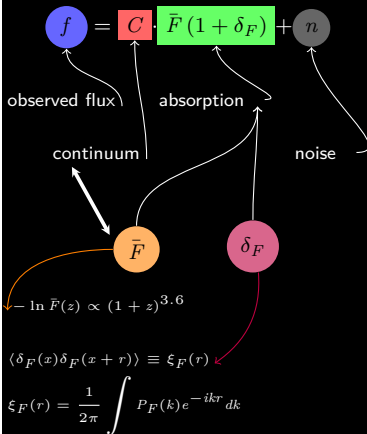
Measured quantity:



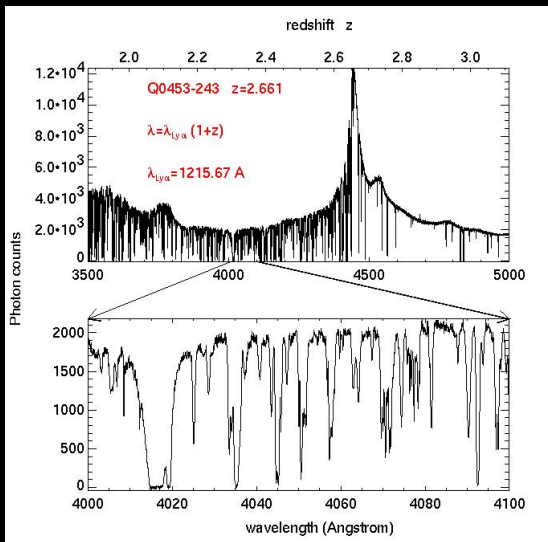
What do we measure?



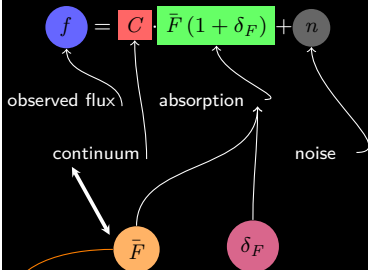
Measured quantity:



What do we measure?



Measured quantity:



$$-\ln \bar{F}(z) \propto (1+z)^{3.6}$$

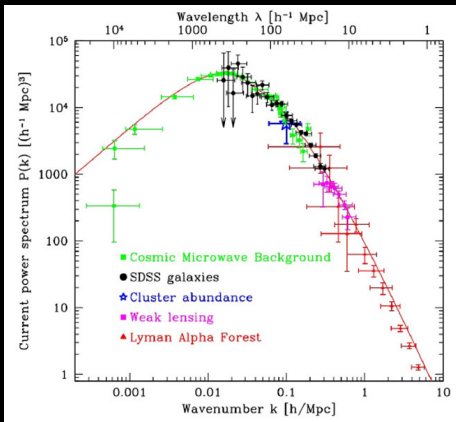
$$\langle \delta_F(x) \delta_F(x+r) \rangle \equiv \xi_F(r)$$

$$\xi_F(r) = \frac{1}{2\pi} \int P_F(k) e^{-ikr} dk$$

1D

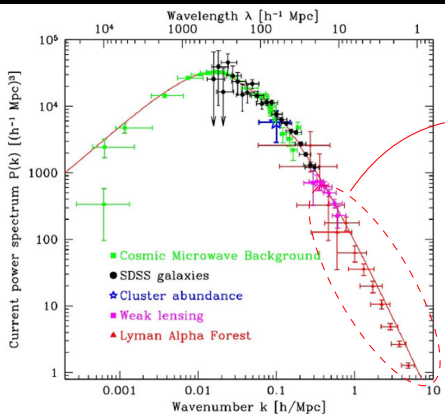
Matter \leftrightarrow Lyman- α

3D: Matter Power spectrum $P_m(k, z)$

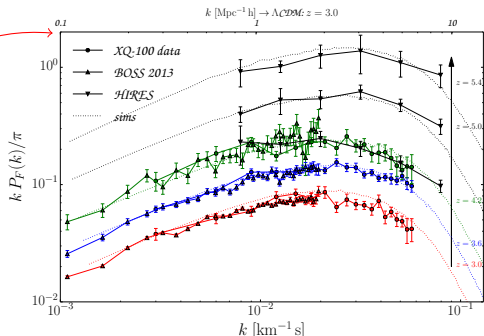


Matter ↔ Lyman- α

3D: Matter Power spectrum $P_m(k, z)$



1D: Flux Power spectrum $P_F(k, z)$



1D: small scales, gas physics, m_{wdm} , ...

$F(\rho_g)$: Complex non-linear relationship

Flux fluctuations in Ly α forest trace matter density fluctuations

$$F = \exp[-\tau(\delta)]$$

- ▶ Intergalactic medium (IGM) is mainly highly ionized hydrogen gas (Gunn & Petterson)
- ▶ UV photo-ionization in equilibrium with recombination
- ▶ Data & simulations suggest the state of IGM: 10^4 K and low densities 10^{-4} cm^{-3}
- ▶ Equation of state of the IGM can be approximately described by $T \propto \rho^\gamma$

Highly nonlinear relation between flux and density

$$F = \exp[-A(1 + \delta)^P]$$

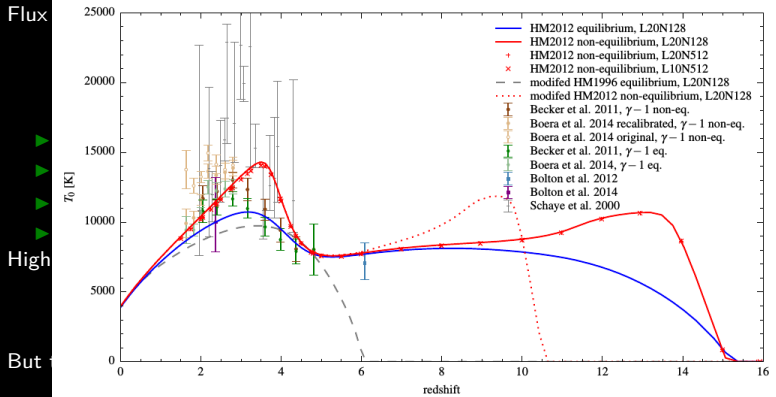
But that is not all... Temperature + peculiar velocity effects:

$$F(v) = \exp \left[-A(\bar{z}; \Omega_i) \int ds \left(1 + \delta_b(s)\right)^2 T(s)^{-0.7} \Gamma_{\gamma, HI}^{-1} V\left(v - s - v_p(s); T(s)\right) \right]$$

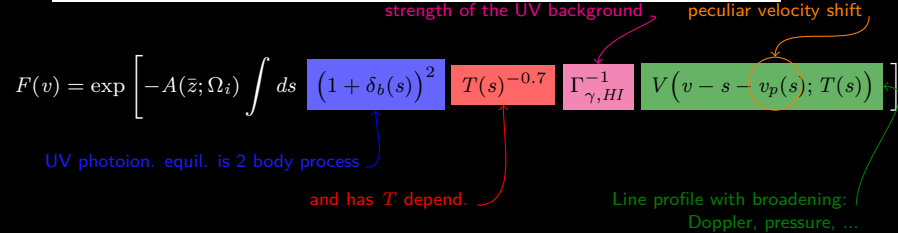
Annotations:

- strength of the UV background (points to $A(\bar{z}; \Omega_i)$)
- peculiar velocity shift (points to $v_p(s)$)
- UV photoion. equil. is 2 body process (points to $\Gamma_{\gamma, HI}^{-1}$)
- and has T depend. (points to $T(s)^{-0.7}$)
- Line profile with broadening: Doppler, pressure, ... (points to $V(v - s - v_p(s); T(s))$)

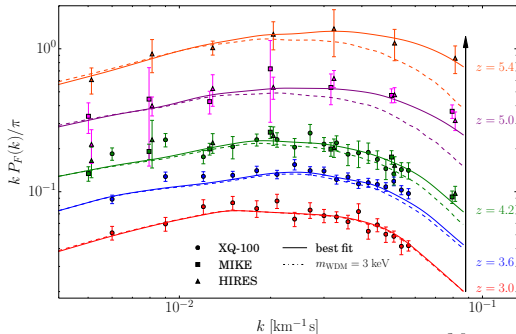
$F(\rho_g)$: Complex non-linear relationship



Petterson)
 cm^{-3}



WDM mass constraints



Low-z data:

XQ-100 ($3 < z < 4.2$)

High-z data:

HIRES/MIKE ($4.2 < z < 5.4$)

Typical DM mass:

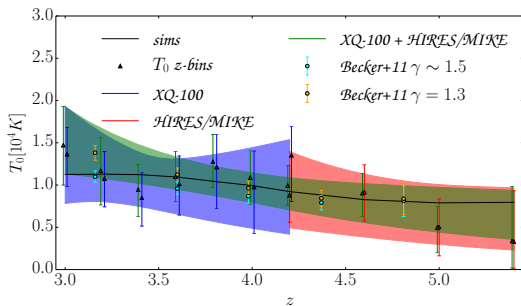
$m_{\text{WDM}} \sim 2 - 3 \text{ keV}$

$T_0(z)$ is power-law

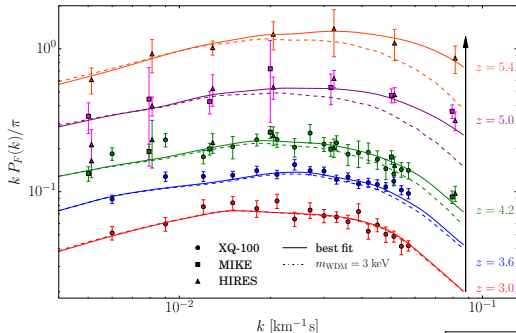
$$\rightarrow m_{\text{WDM}} > 5.3 \text{ keV @ } 2\sigma$$

$T_0(z)$ free + $\frac{\partial T_0}{\partial z}$ bounded

$$\rightarrow m_{\text{WDM}} > 3.5 \text{ keV @ } 2\sigma$$



WDM mass constraints



Low-z data:

XQ-100 ($3 < z < 4.2$)

High-z data:

HIRES/MIKE ($4.2 < z < 5.4$)

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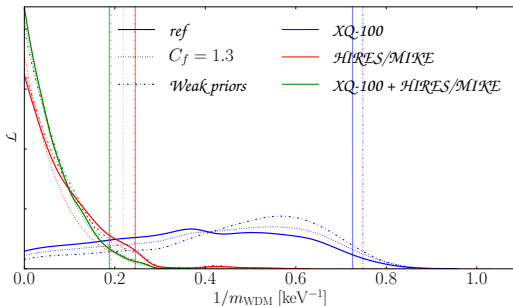
$m_{\text{WDM}} \sim 2 - 3 \text{ keV}$

$T_0(z)$ is power-law

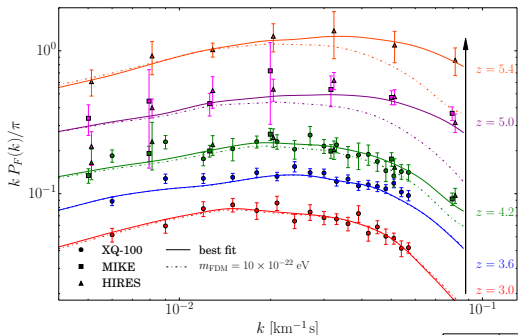
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$T_0(z)$ free + $\frac{\partial T_0}{\partial z}$ bounded

$$\rightarrow m_{\text{WDM}} > 3.5 \text{ keV @ } 2\sigma$$



FDM mass constraints



Results later confirmed by independent groups:

Yeche et al. 2017 (WDM)

Armengaud et al. 2017 (FDM)

Typical DM mass:

$$m_{\text{FDM}} \sim 1 - 10 \times 10^{-22} \text{ eV}$$

$T_0(z)$ is power-law

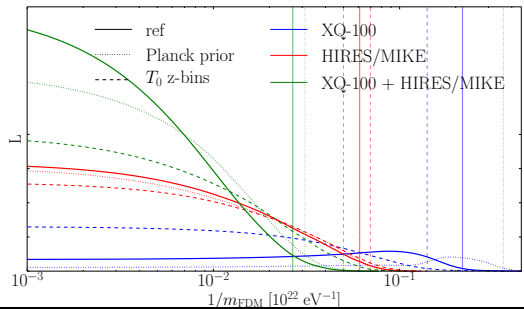
$$\rightarrow m_{\text{FDM}} > 37.5 \times 10^{-22} \text{ eV @ } 2\sigma$$

$T_0(z)$ free + $\frac{\partial T_0}{\partial z}$ bounded

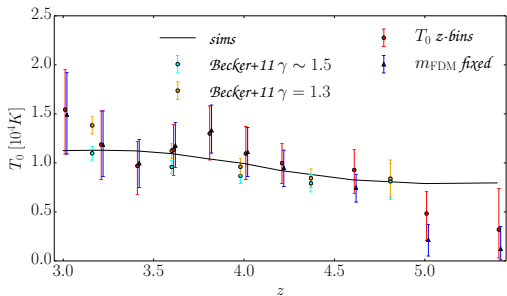
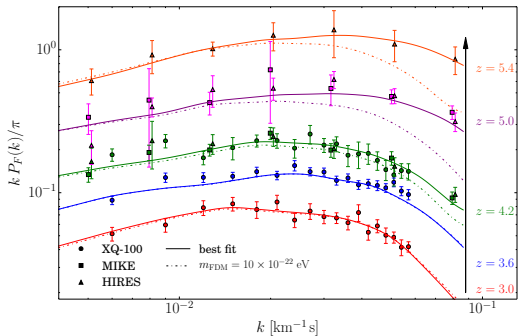
$$\rightarrow m_{\text{FDM}} > 20.4 \times 10^{-22} \text{ eV @ } 2\sigma$$

FDM + Quantum Pressure

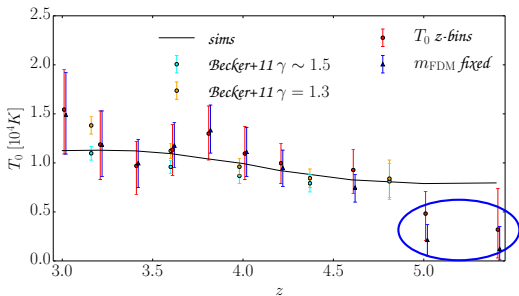
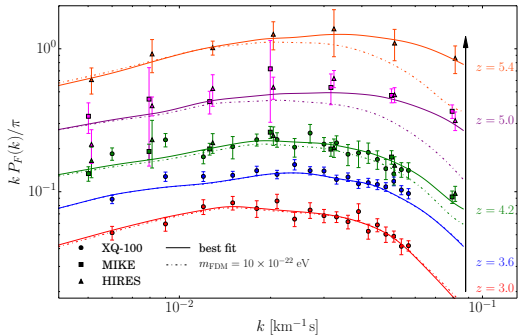
$$\rightarrow m_{\text{FDM}} > 21.1 \times 10^{-22} \text{ eV @ } 2\sigma$$



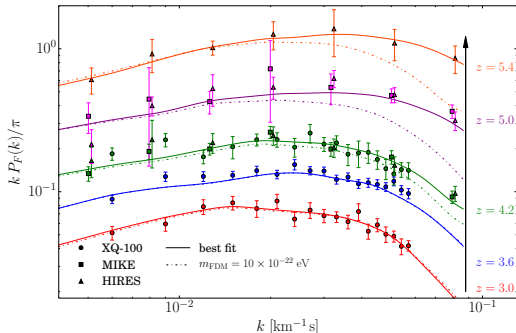
How cold is too cold?



How cold is too cold?



How cold is too cold?



Simple model:

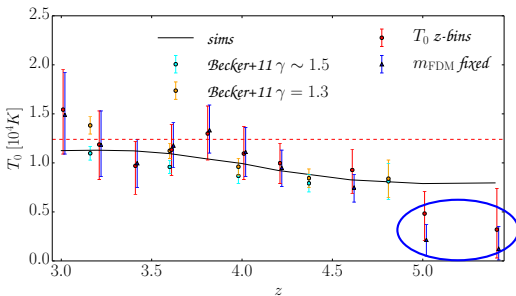
- ▶ instantaneous H reionisation at $z_{\text{rei}} = 9$
- ▶ HI photo-heating, depends on spectral index of UV intensity $\alpha_{bk} = 0$
- ▶ Compton cooling + adiabatic expansion

$$T_0(z = 5.0) = 12,400 \text{ K}$$

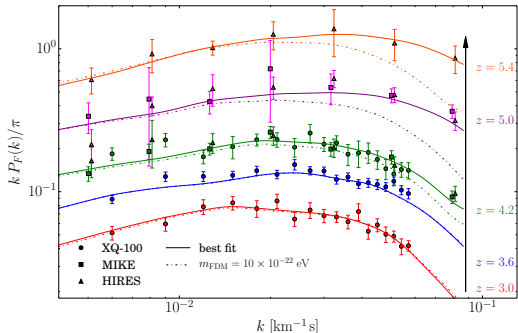
McQuinn & Upton Sanderback (2015),
Upton Sanderback et al. (2016)

Other things assumed:

- ▶ T fluctuations increase above this temperature
- ▶ He I and He II photo-heating only increases the temperature
- ▶ H II, He III recombination cooling decreases temperature by \sim few %
- ▶ Planck Λ CDM Cosmology
- ▶ $T_{\text{rei}} = 10,000 \text{ K}$ (more realistic would be 20,000 K)



How cold is too cold?



Simple model:

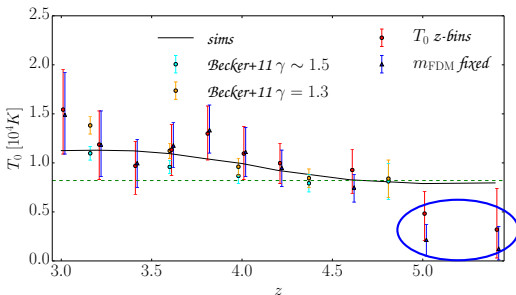
- ▶ instantaneous H reionisation at $z_{\text{rei}} = 9$
- ▶ HI photo-heating, depends on spectral index of UV intensity $\alpha_{bk} = 2$
- ▶ Compton cooling + adiabatic expansion

$$T_0(z = 5.0) = 8,200 \text{ K}$$

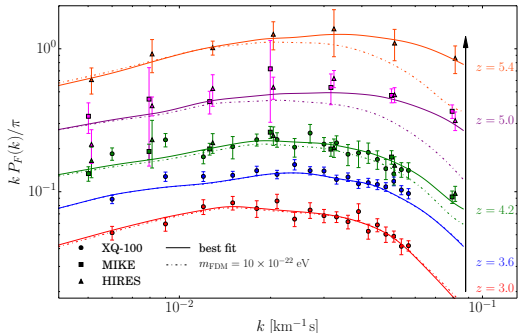
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How cold is too cold?



Simple model:

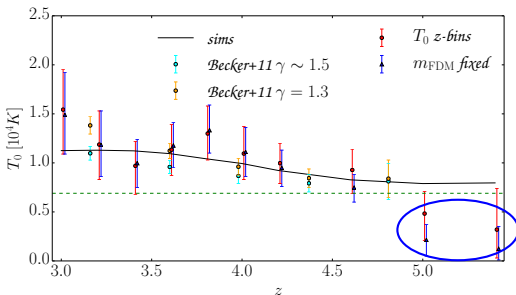
- ▶ instantaneous H reionisation at $z_{\text{rei}} = 15$
- ▶ HI photo-heating, depends on spectral index of UV intensity $\alpha_{bk} = 2$
- ▶ Compton cooling + adiabatic expansion

$$T_0(z = 5.0) = 6,900 \text{ K}$$

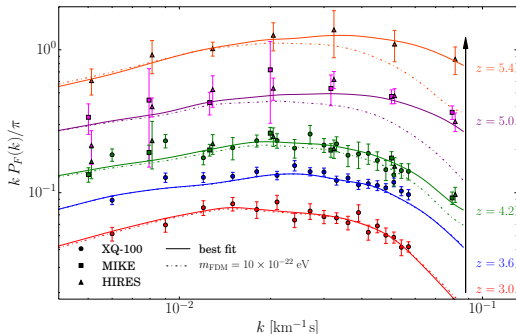
McQuinn & Upton Sanderback (2015),
Upton Sanderback et al. (2016)

Other things assumed:

- ▶ T fluctuations increase above this temperature
- ▶ He I and He II photo-heating only increases the temperature
- ▶ H II, He III recombination cooling decreases temperature by \sim few %
- ▶ Planck Λ CDM Cosmology
- ▶ $T_{\text{rei}} = 10,000 \text{ K}$ (more realistic would be 20,000 K)



How cold is too cold?



Simple model:

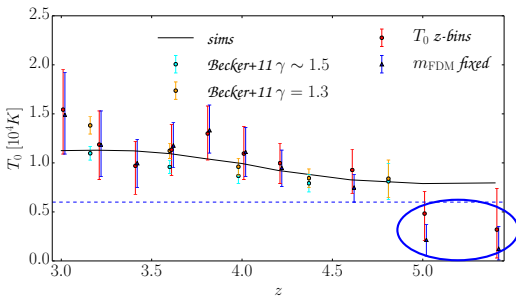
- ▶ instantaneous H reionisation at $z_{\text{rei}} = 15$
- ▶ HI photo-heating, depends on spectral index of UV intensity $\alpha_{bk} = 3$
- ▶ Compton cooling + adiabatic expansion

$$T_0(z = 5.0) = 6,000 \text{ K}$$

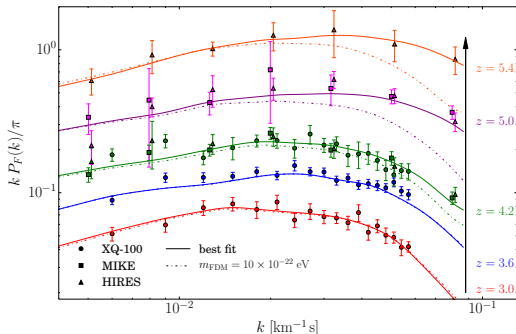
McQuinn & Upton Sanderback (2015),
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Other things assumed:

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How cold is too cold?



Simple model:

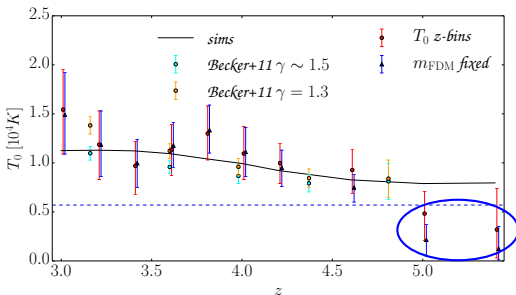
- ▶ instantaneous H reionisation at $z_{\text{rei}} = 20$
- ▶ HI photo-heating, depends on spectral index of UV intensity $\alpha_{\text{bk}} = 3$
- ▶ Compton cooling + adiabatic expansion

$$T_0(z = 5.0) = 5,700 \text{ K}$$

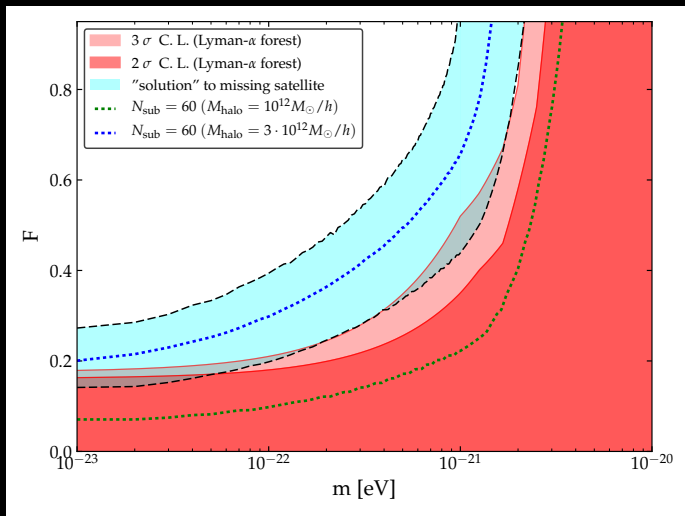
McQuinn & Upton Sanderback (2015),
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- ▶ H II, He III recombination cooling decreases temperature by \sim few %
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Overlapping constraints with different probes

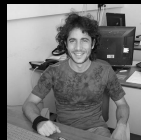
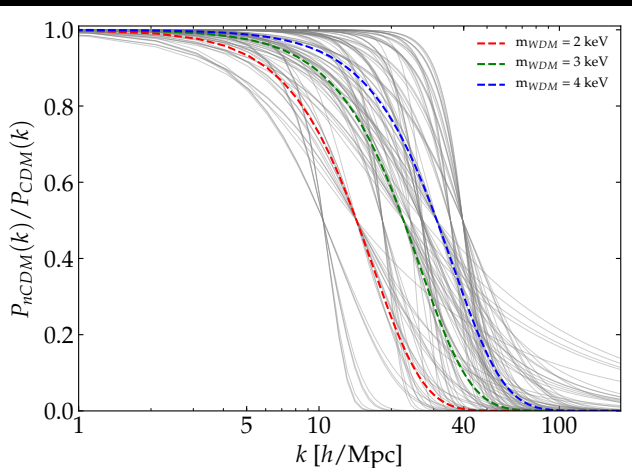


with T. Kobayashi
(SISSA)

General non-CDM models

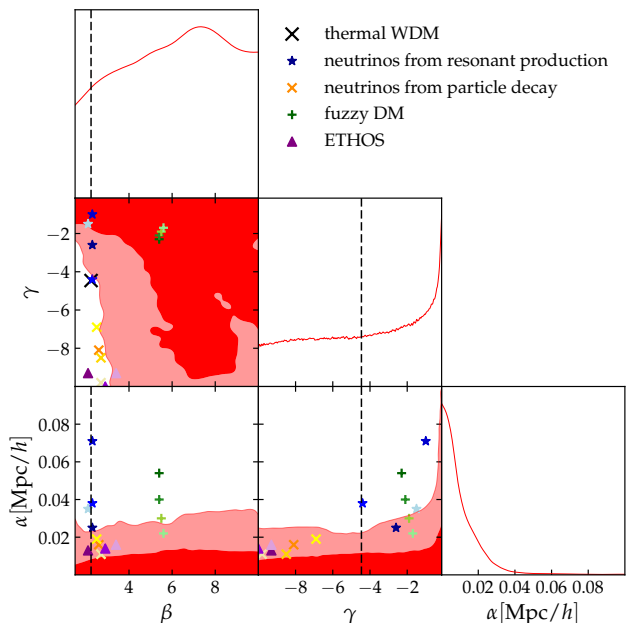
General transfer function for DM: $T(k) = \sqrt{\frac{P_{nCDM}}{P_{CDM}}} = [1 + (\alpha k)^\beta]^\gamma,$

E.g. for thermal WDM: $\beta = 2.24, \gamma = -4.46, \alpha \propto 0.049 \left(\frac{m_{WDM}}{1 \text{ keV}}\right)^{-1.11} h^{-1} \text{ Mpc}$



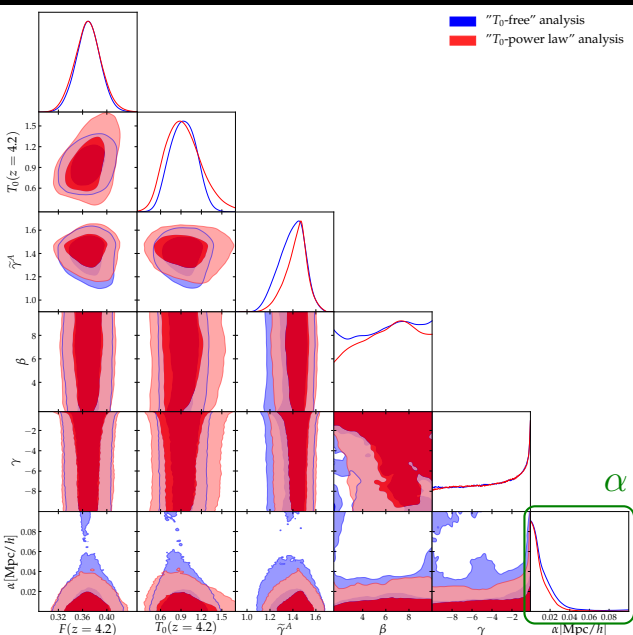
with R. Murgia
(SISSA)

Constraints on the shape of the n CDM $T(k)$



	α [Mpc/h]	β	γ
Neutrinos from resonant production	0.025	2.3	-2.6
	0.071	2.3	-1.0
	0.038	2.3	-4.4
	0.035	2.1	-1.5
Neutrinos from particle decay	0.016	2.6	-8.1
	0.011	2.7	-8.5
	0.019	2.5	-6.9
	0.011	2.7	-9.8
Mixed models	0.16	3.2	-0.4
	0.20	3.7	-0.18
	0.21	3.7	-0.1
	0.21	3.4	-0.053
Fuzzy DM	0.054	5.4	-2.3
	0.040	5.4	-2.1
	0.030	5.5	-1.9
	0.022	5.6	-1.7
ETHOS models	0.0072	1.1	-9.9
	0.013	2.1	-9.3
	0.014	2.9	-10.0
	0.016	3.4	-9.3

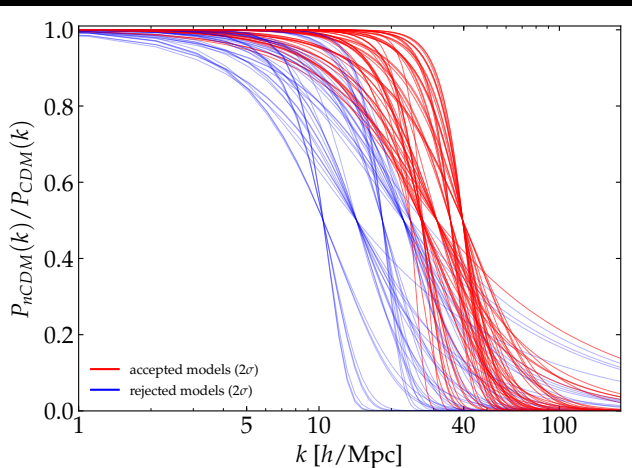
Stable limit on the scale of suppression



$\alpha < 0.03 \text{ Mpc}/h$ (2σ)

Stable limit on the scale of suppression

$$\alpha < 0.03 \text{ Mpc}/h \text{ (} 2\sigma \text{)}$$



Conclusions

Cosmological & Astrophysical Constraints on WDM:

- ▶ Combined data: XQ-100 + HIRES/MIKE (high resolution, high redshift)
- ▶ Large redshift range and probing small scales
- ▶ Good agreement with thermal history evolution (a paper by Becker+11)
- ▶ Constraints on WDM from combined data: $m_{\text{WDM}} > 5.3 \text{ keV}$ at 2σ .
- ▶ Constraints on WDM from combined data: $m_{\text{WDM}} > 3.5 \text{ keV}$ at 2σ (conservative thermal history)
- ▶ The paper: [astro-ph/1702.01764](https://arxiv.org/abs/1702.01764)

Cosmological & Astrophysical Constraints on FDM:

- ▶ First FDM constraints from Lyman- α forest
- ▶ Constraints on FDM from combined data: $m_{\text{FDM}} > 37.5 \times 10^{-22} \text{ eV}$ at 2σ .
- ▶ Constraints on FDM from combined data: $m_{\text{FDM}} > 20.0 \times 10^{-22} \text{ eV}$ at 2σ (conservative thermal history)
- ▶ m_{FDM} value from "local" Universe leads to unphysically small high- z temperature
- ▶ FDM parameter space greatly constrained: it is hard to solve missing satellite problem and satisfy Ly α constraints.
- ▶ The papers: [astro-ph/1703.04683](https://arxiv.org/abs/1703.04683), [astro-ph/1708.00015](https://arxiv.org/abs/1708.00015)