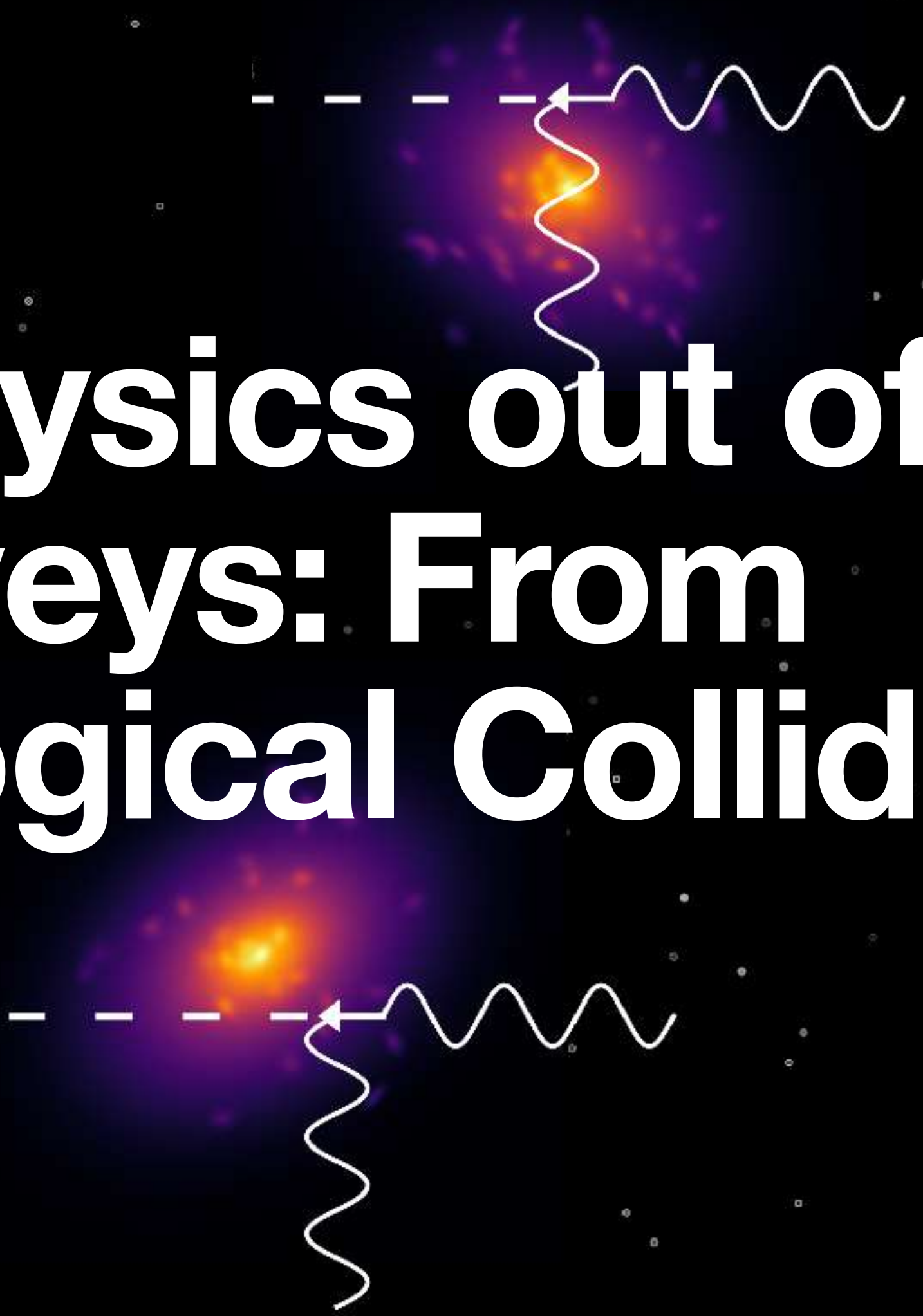
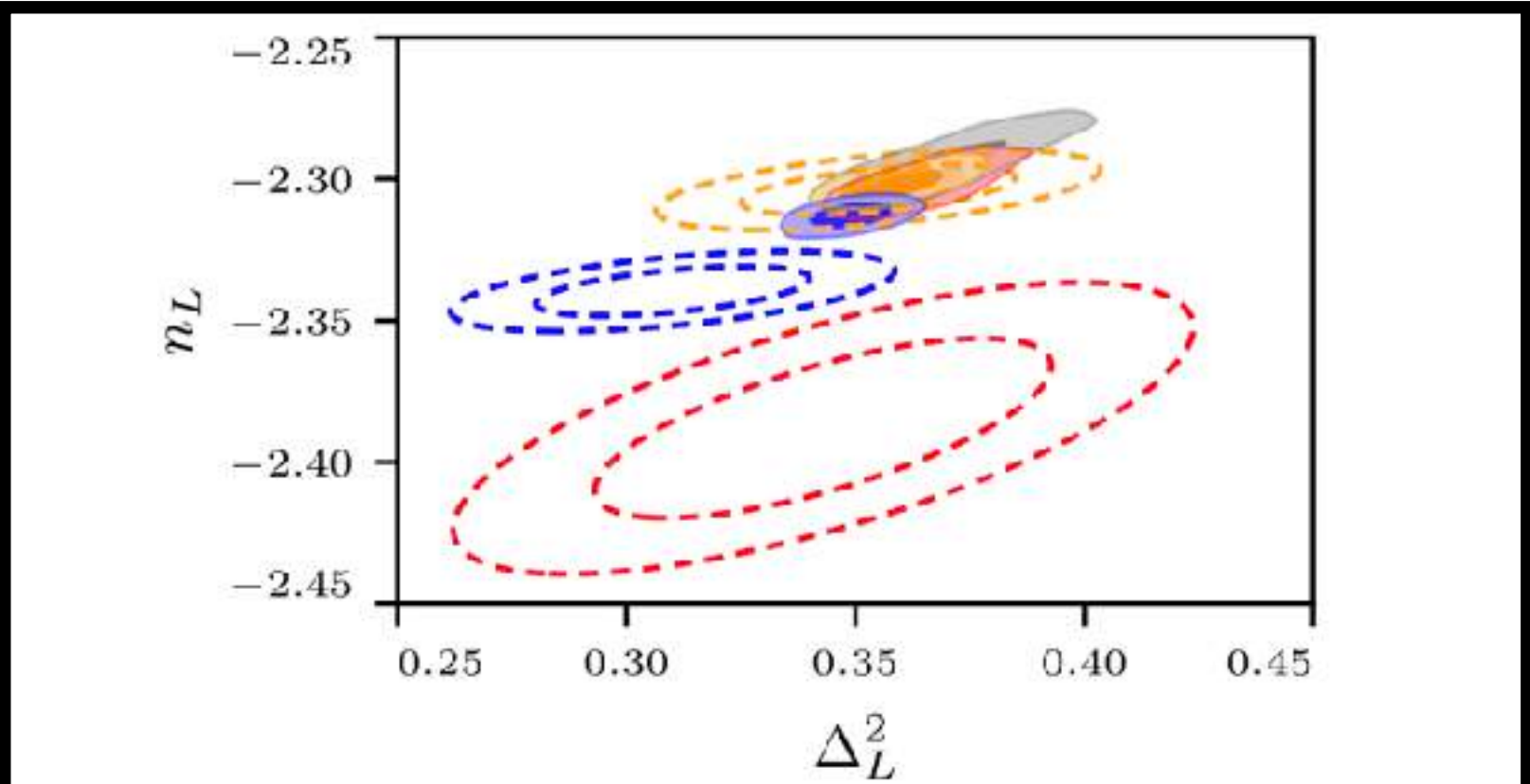


Squeezing New Physics out of Cosmological Surveys: From Axions to Cosmological Colliders

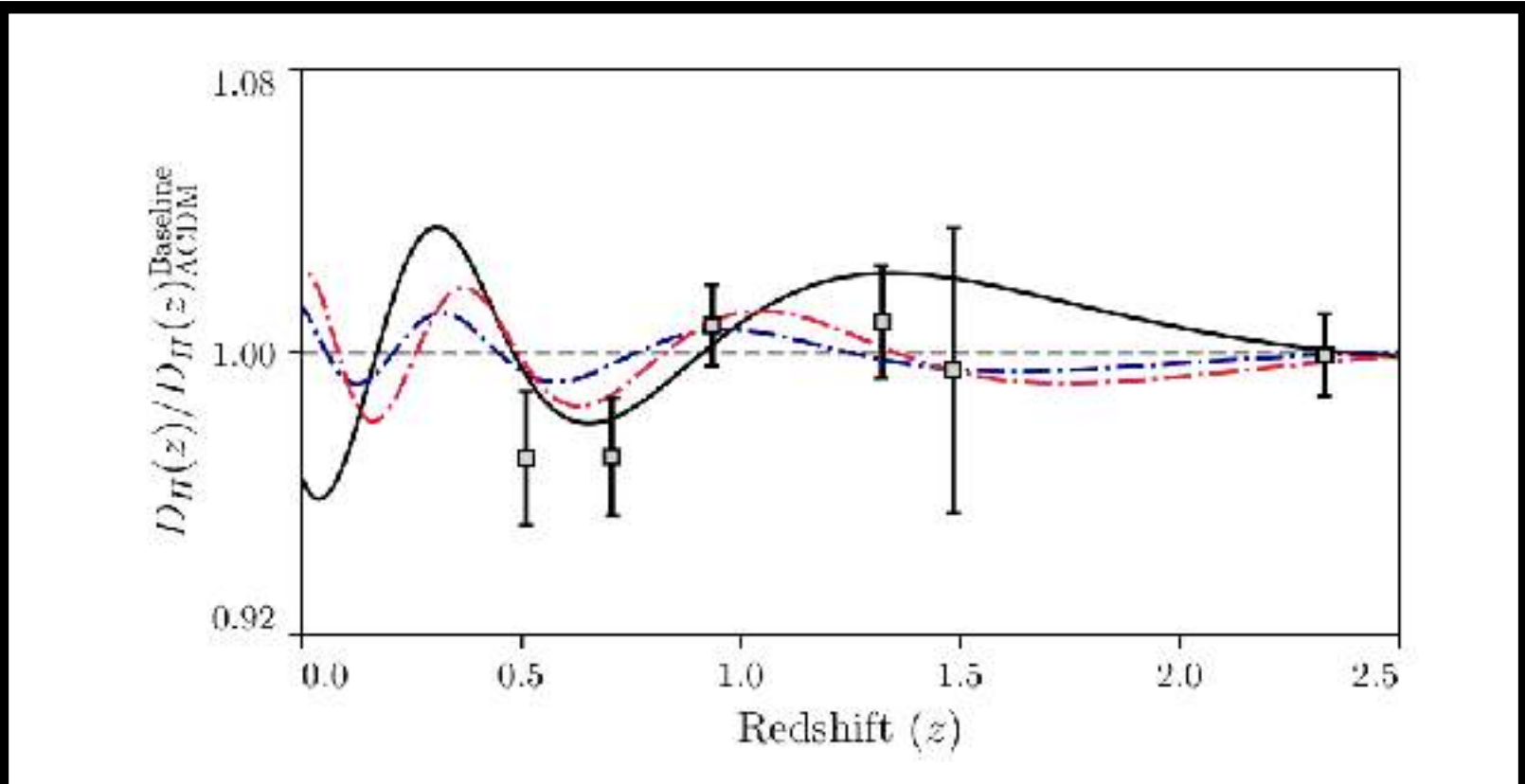


Sam Goldstein

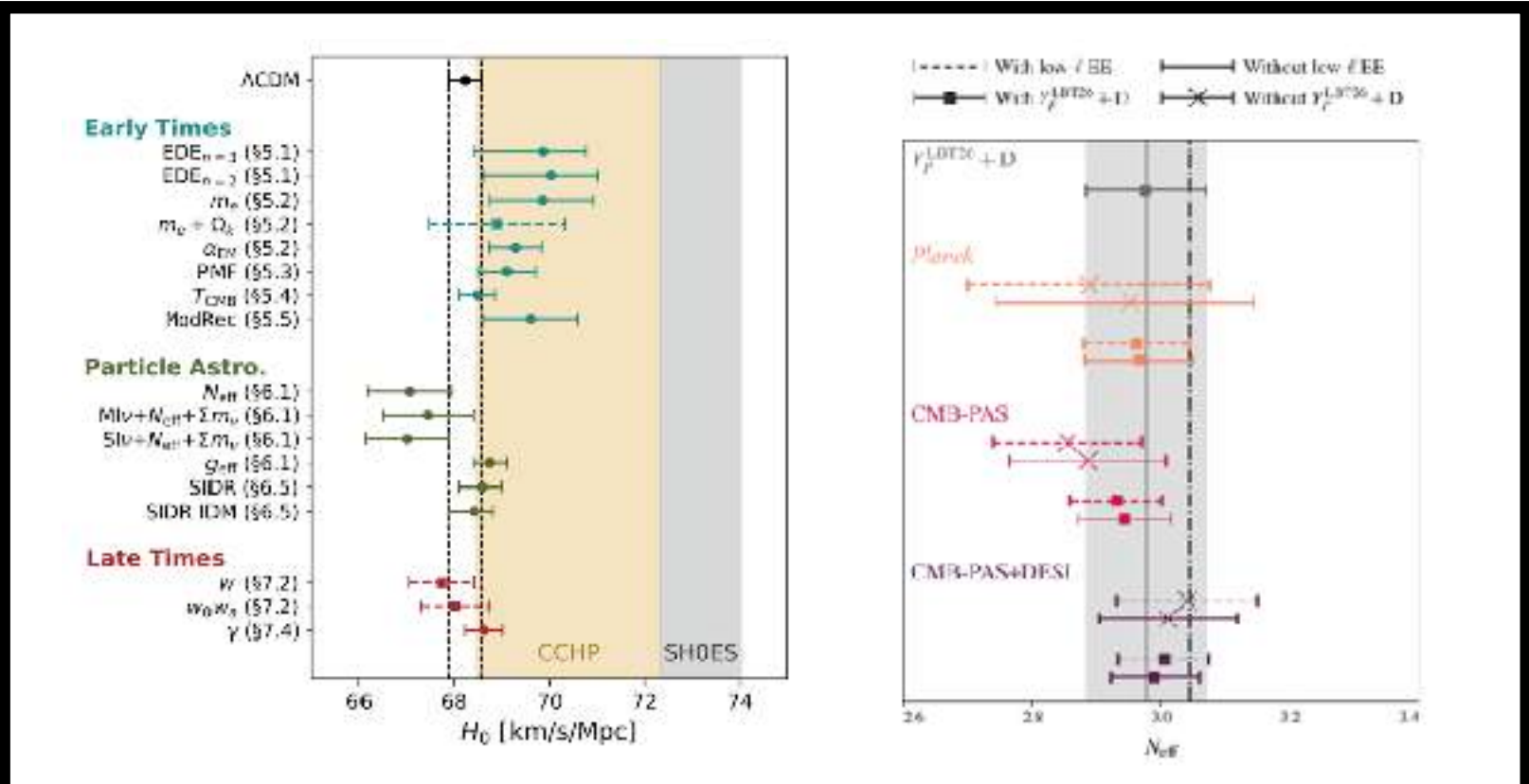
Things I won't talk about today, but I am around for the next 6 weeks!



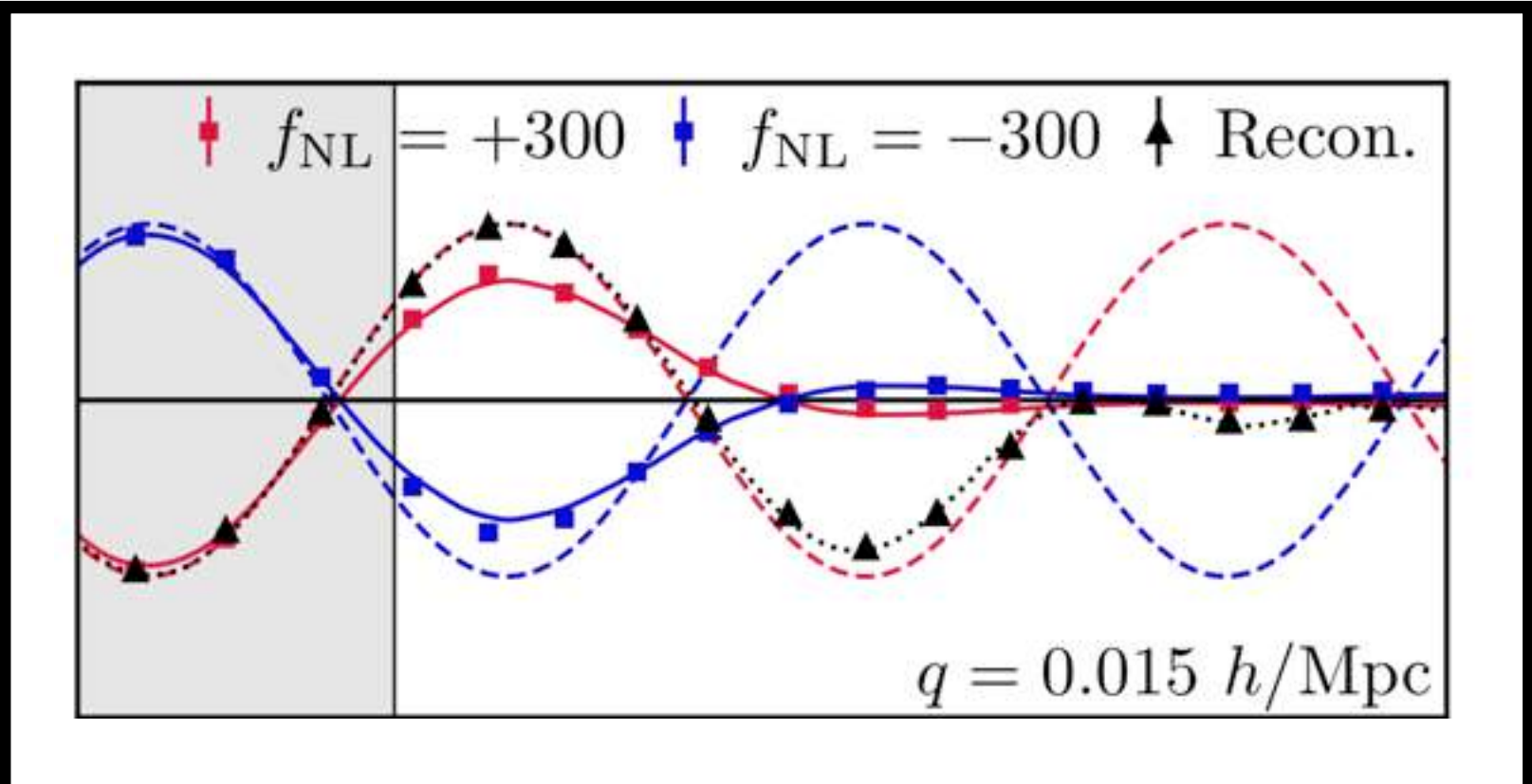
Early dark energy and Ly- α forest
SG++ 2303.00746 (PRL), Editors' Suggestion



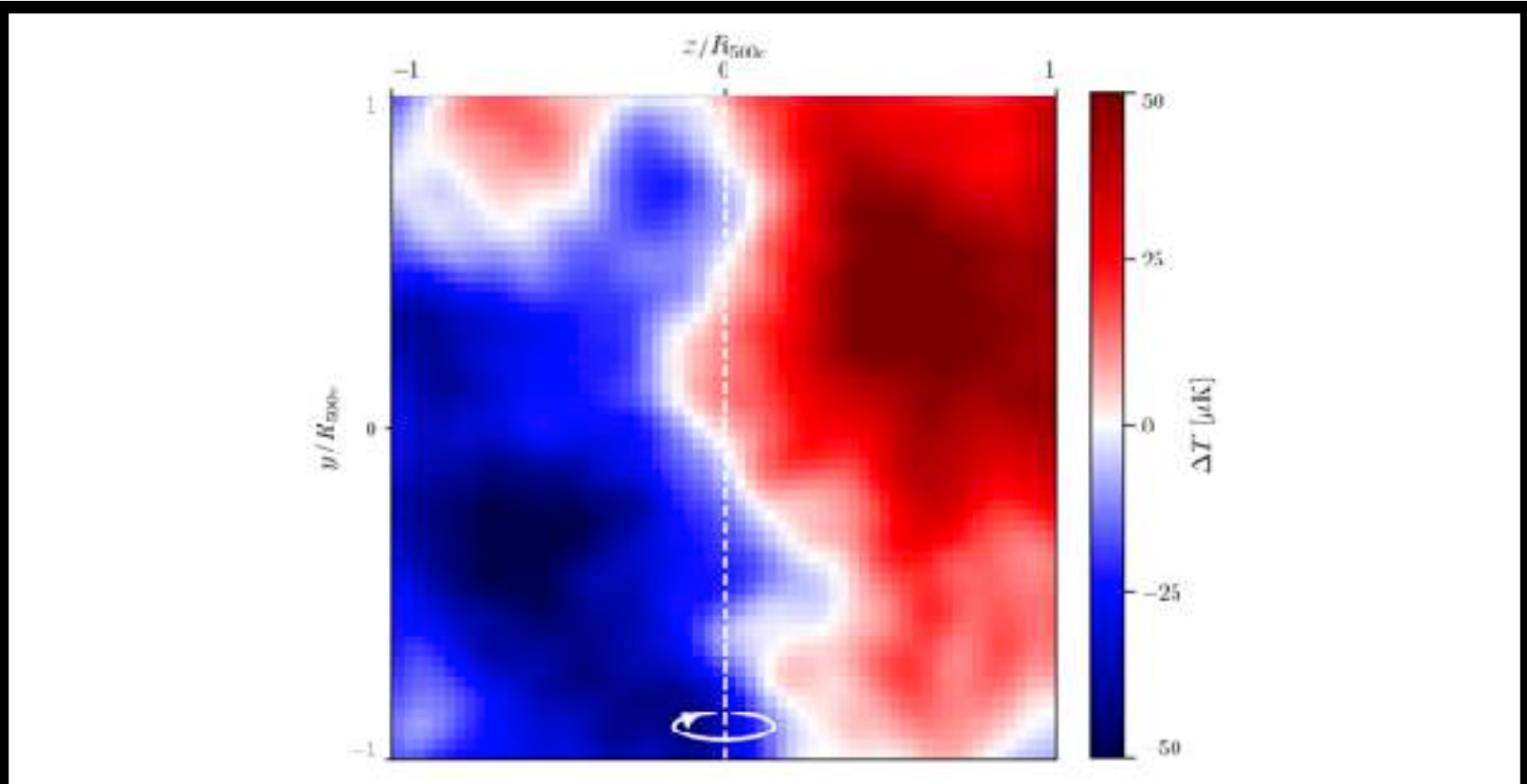
(Over?)fitting wiggles in DESI BAO
SG++ 2507.16970 (PRD)



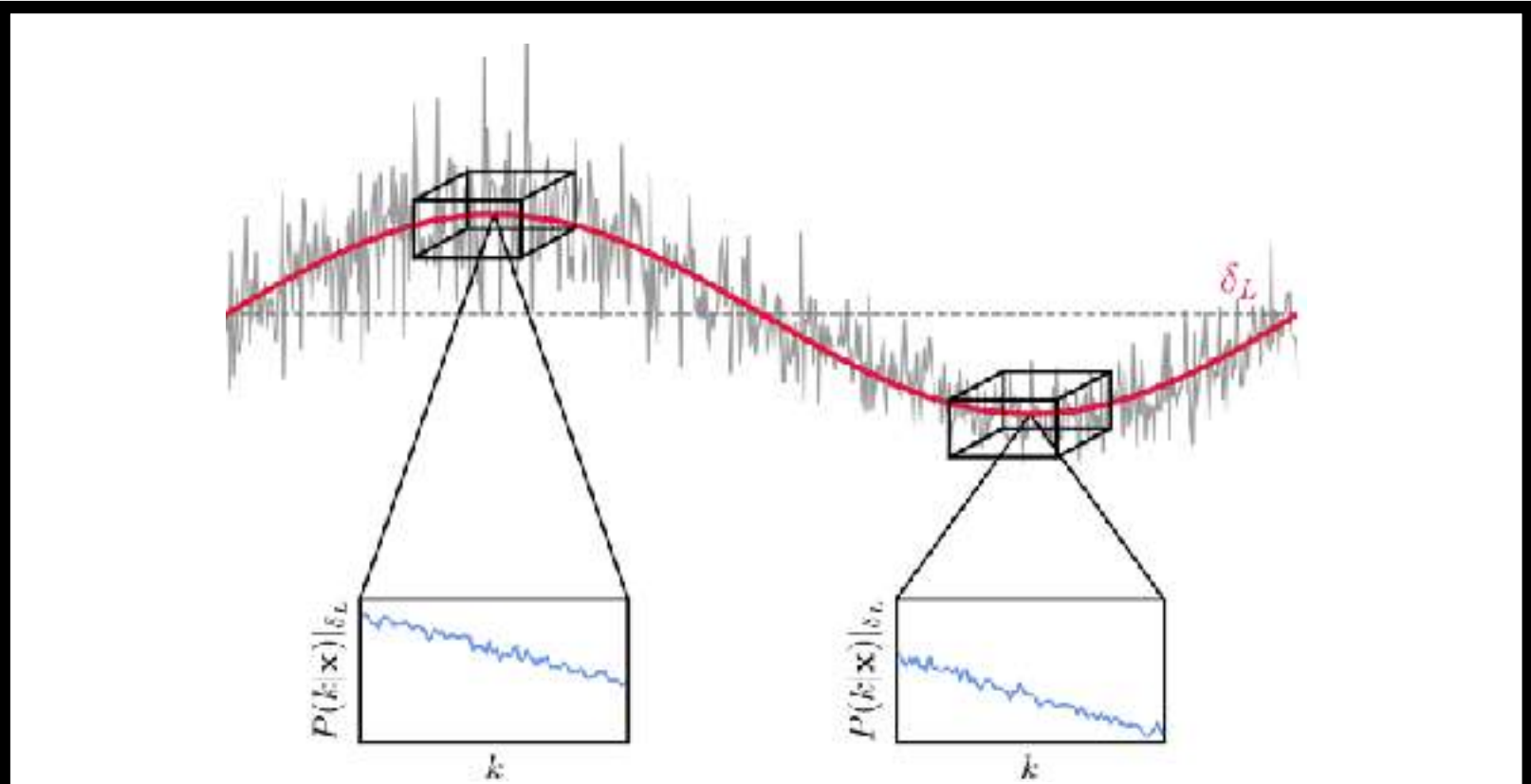
Cosmology from the primary CMB
ACT DR6 2503.14454 (JCAP) SG++2603.13226 (PRD)



Simulating the cosmological collider
SG++2407.08731 (PRD) SG++2505.13443 (PRD)



Rotational kSZ effect
SG++ (PRL in press), Editors' Suggestion

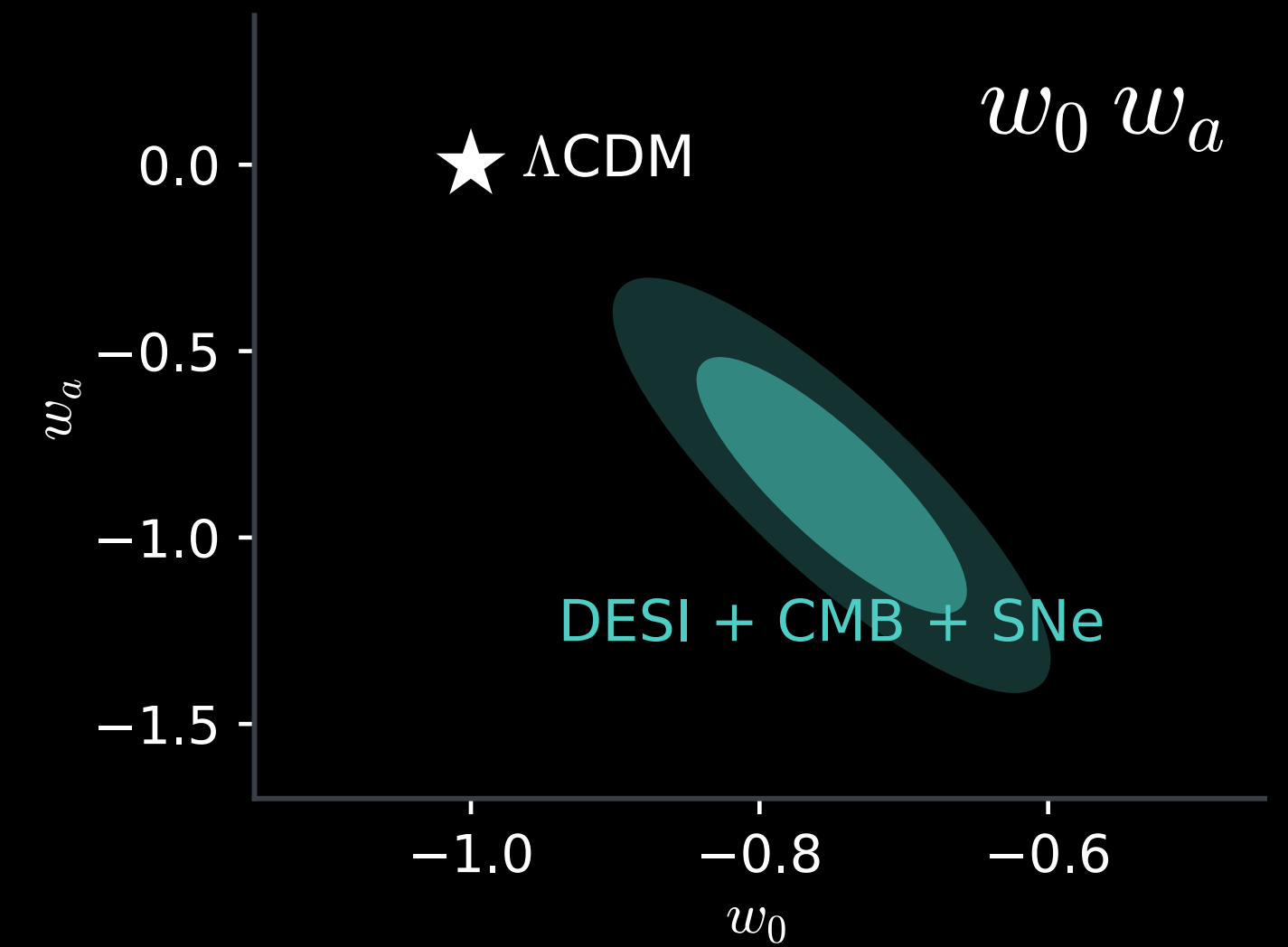
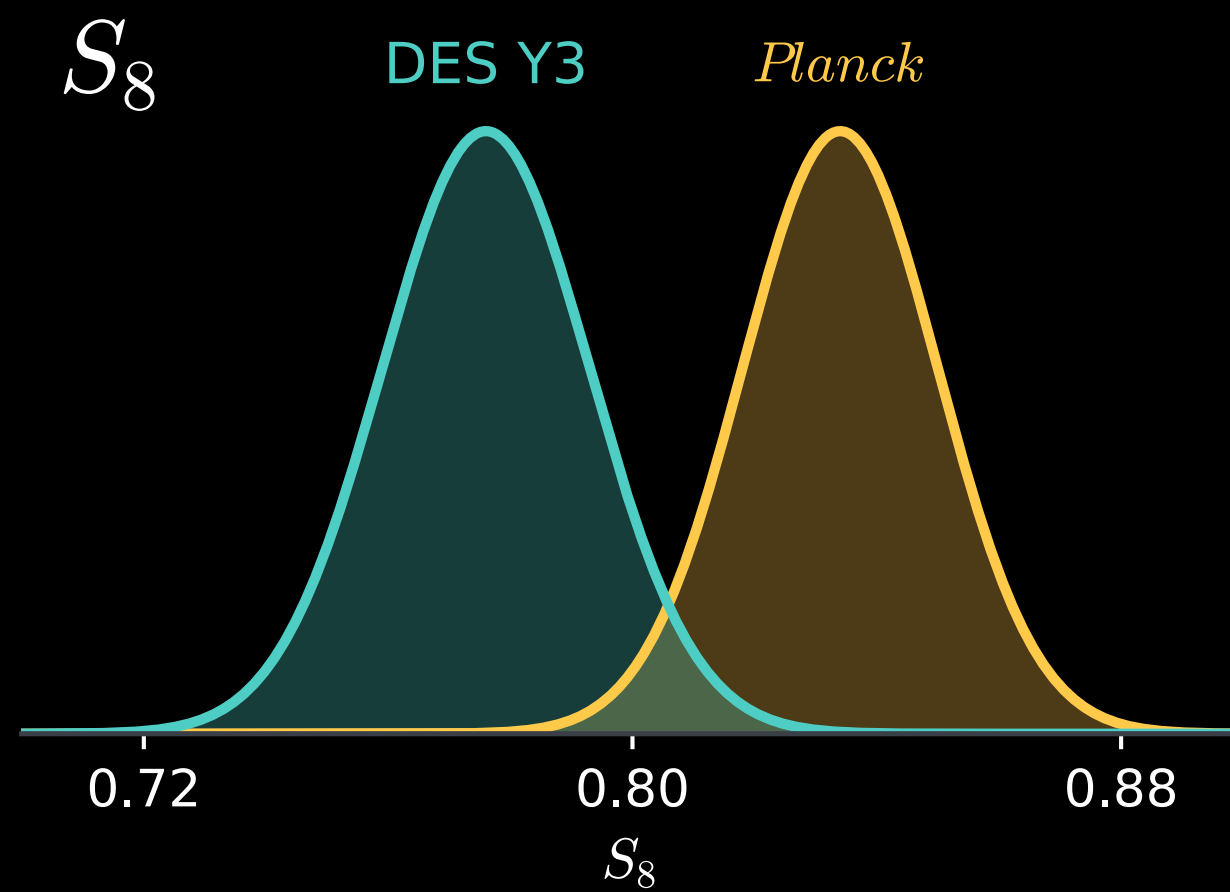
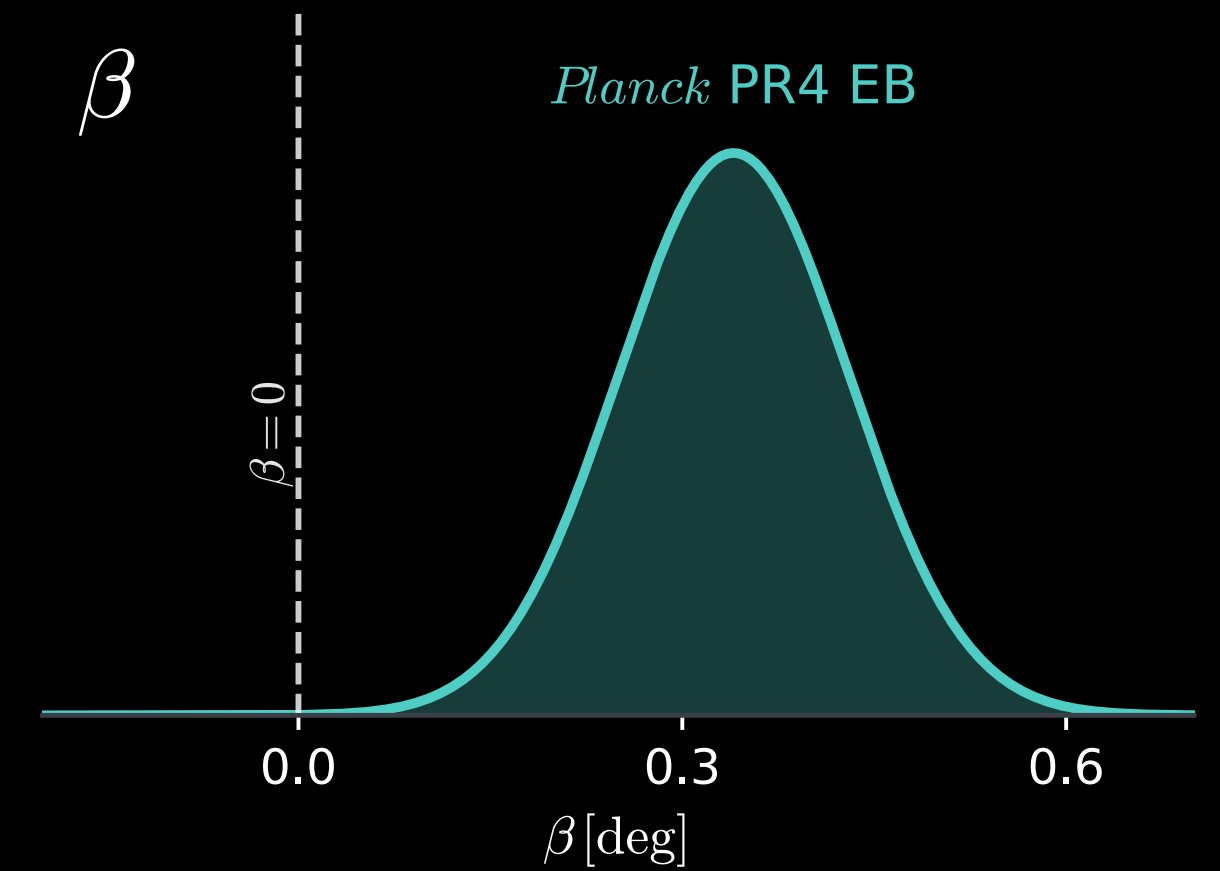
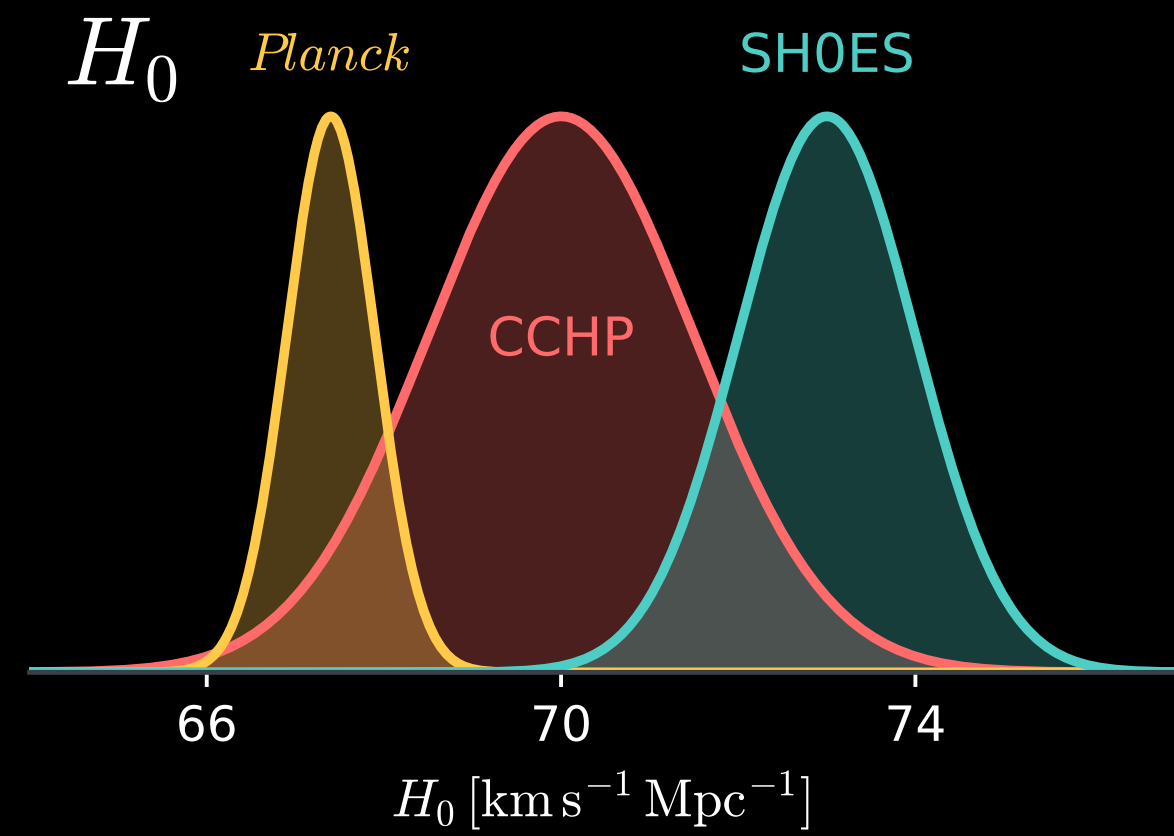


Power spectrum covariance estimation
SG++ 2604.18581 (PRD under review)

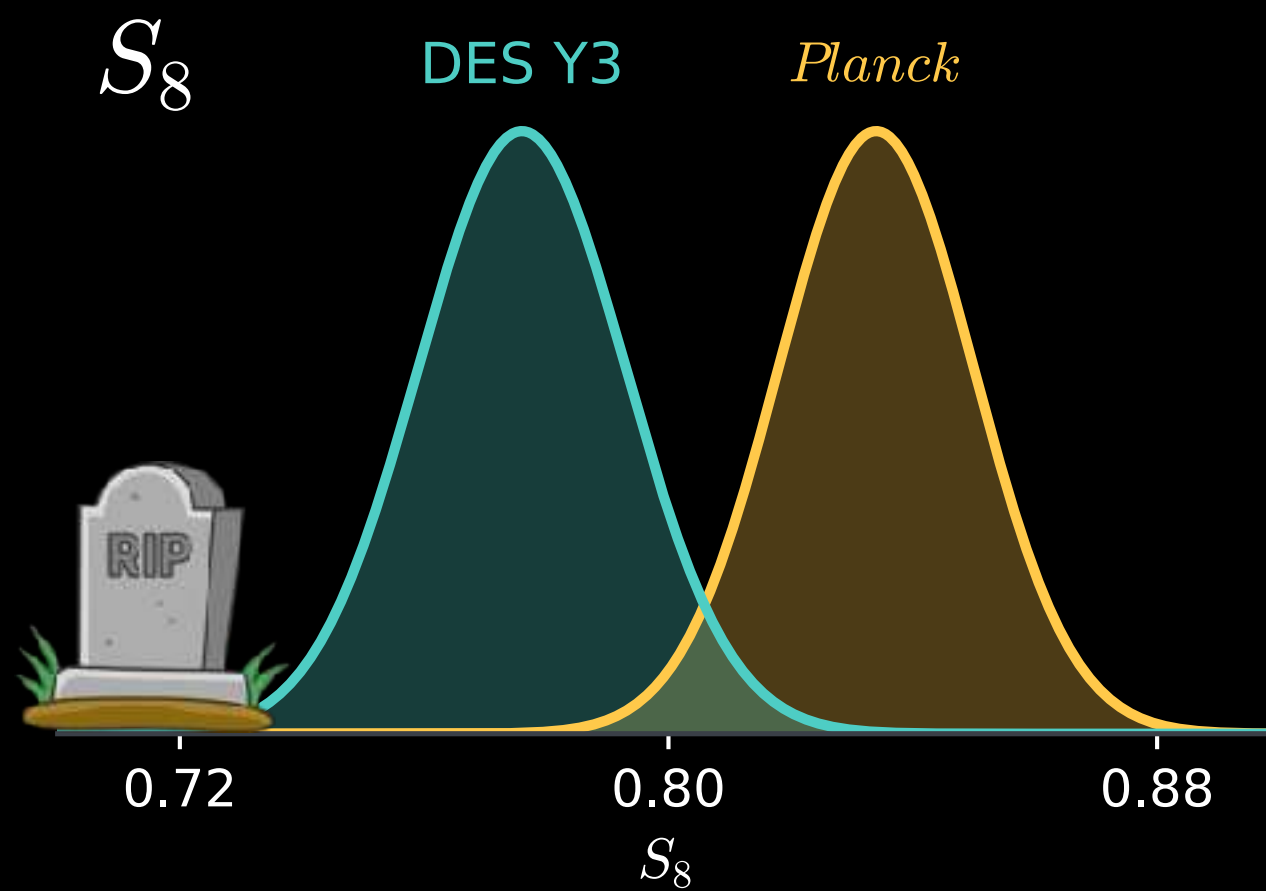
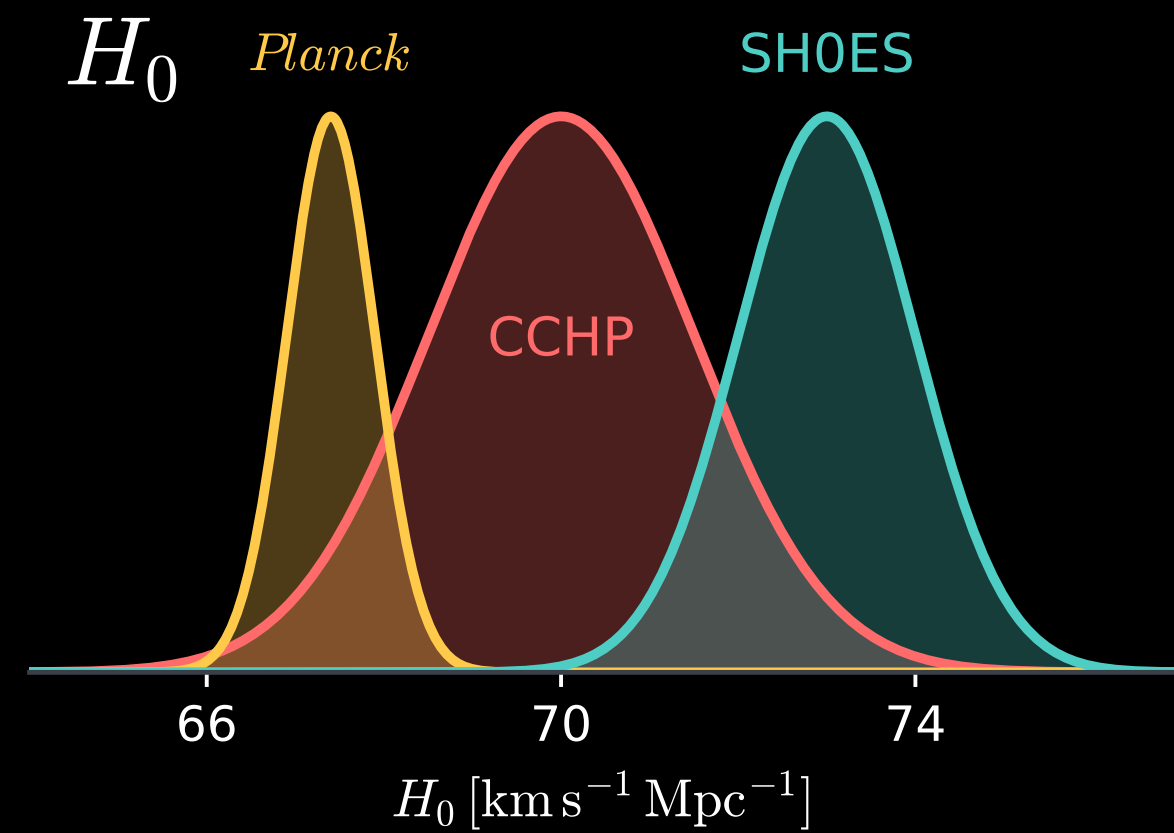
Why search for new physics?



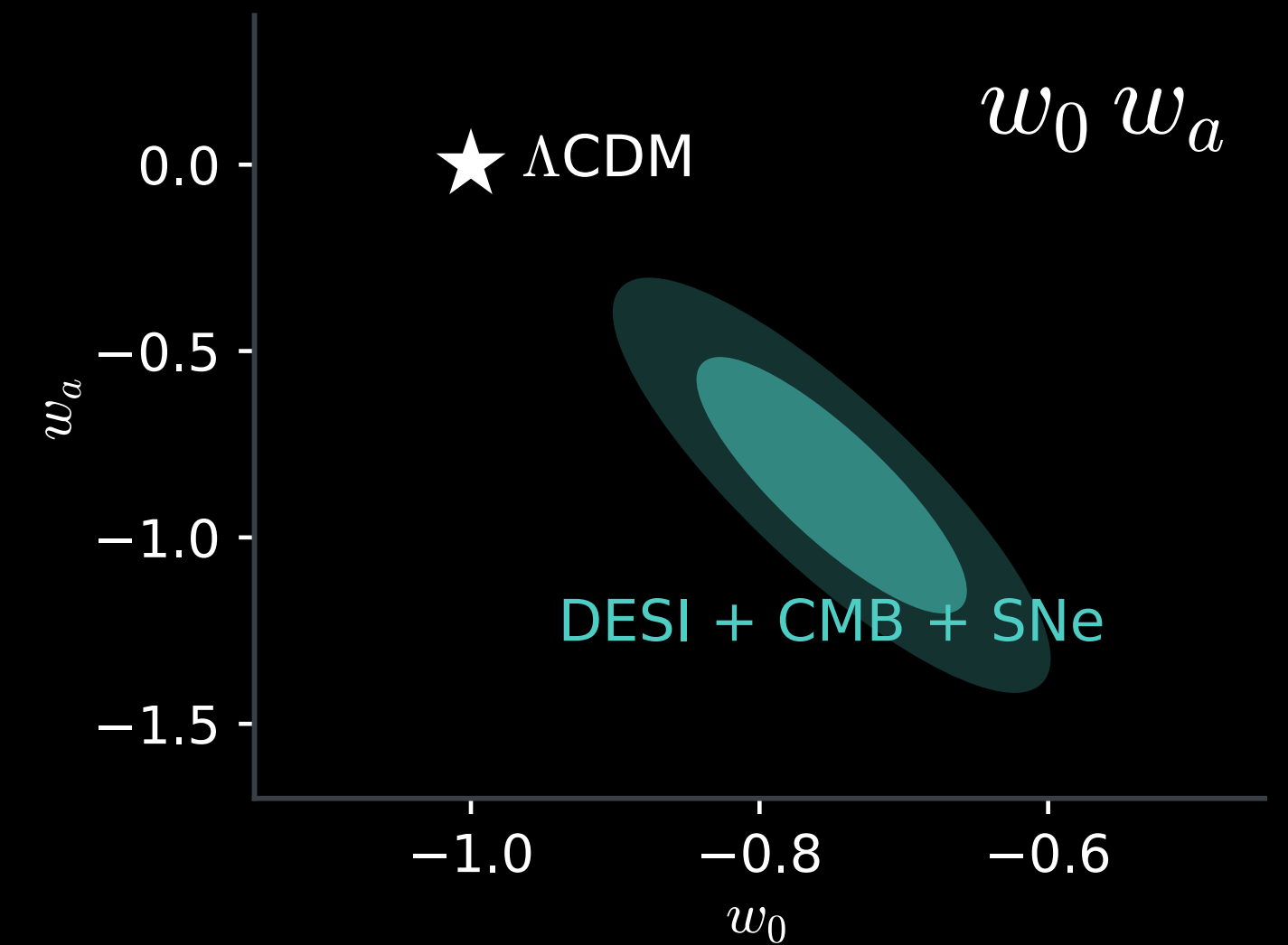
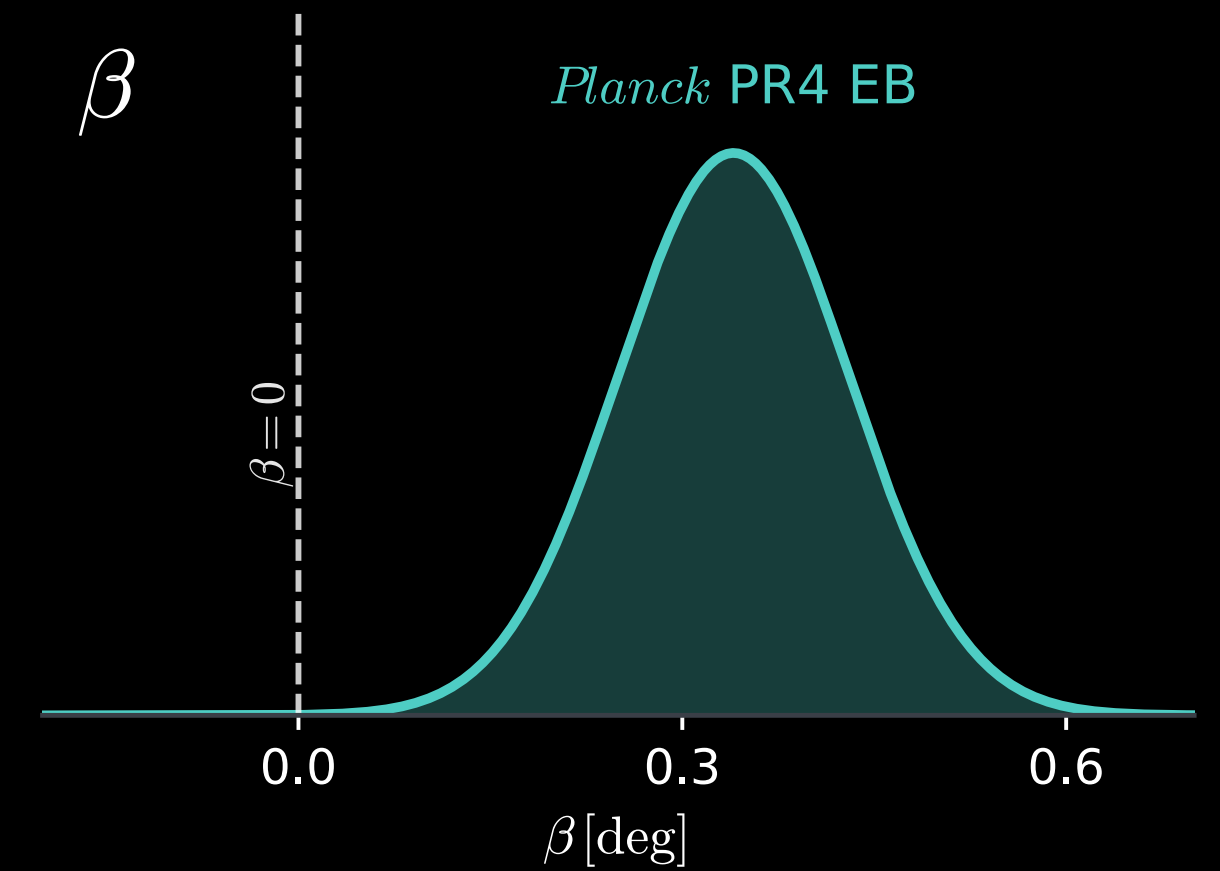
Why search for new physics?



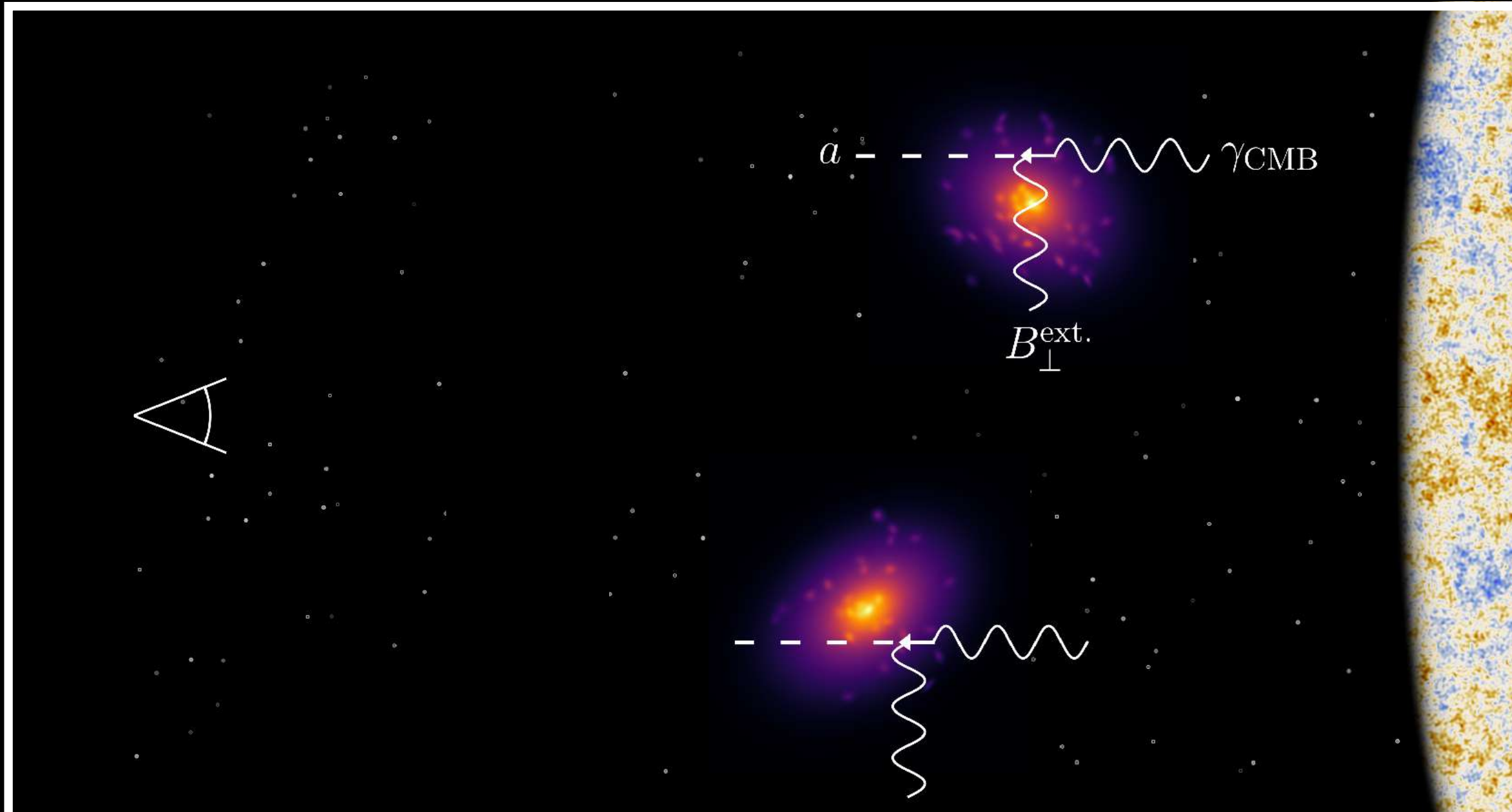
Why search for new physics?



hints of new physics
come and go
but the friends ~~we~~ *pipelines*
we made along the way
are forever

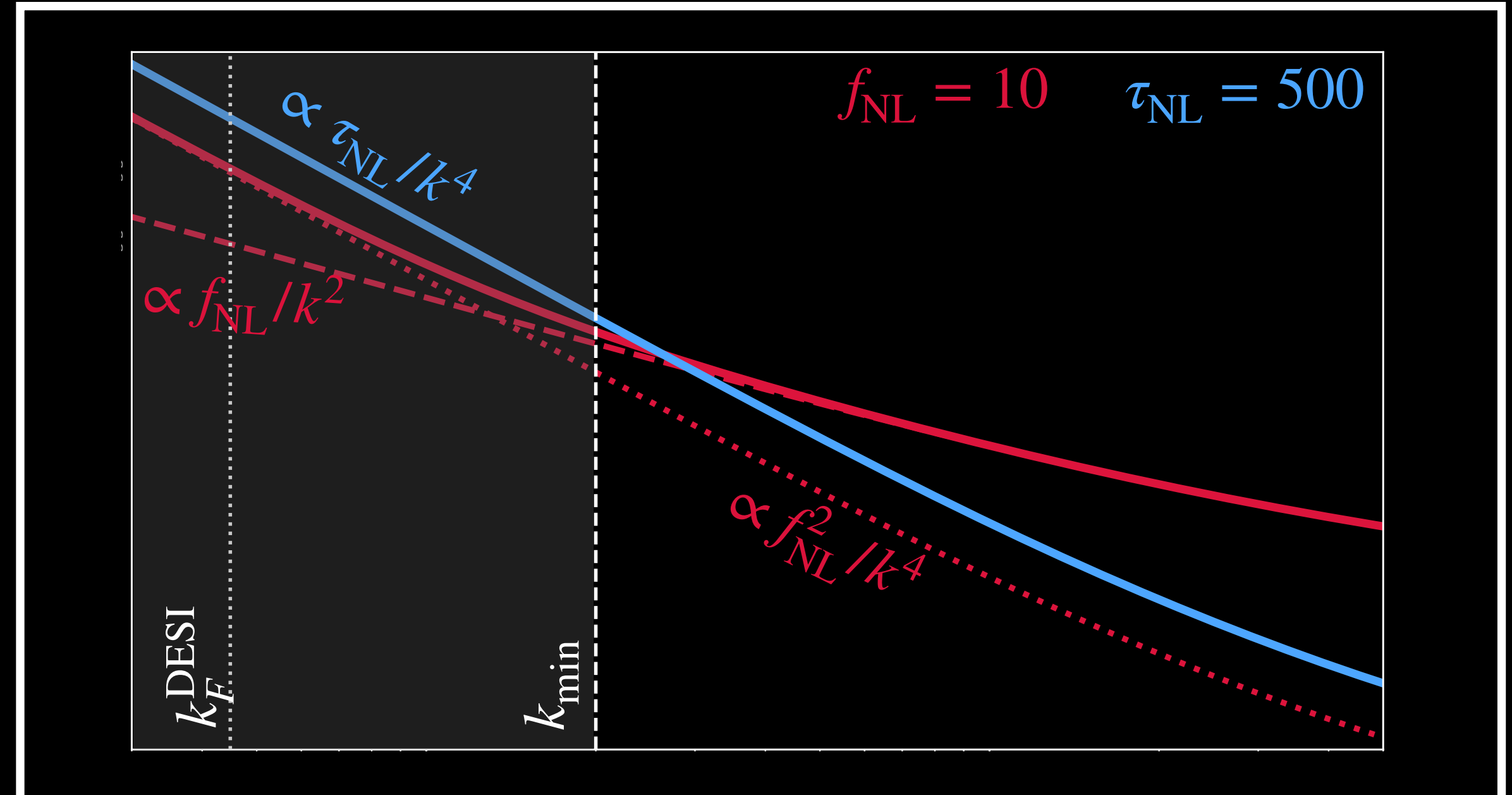


Outline



Part 1: Searching for axions with CMB $\times$ LSS

- *Background*: axion-induced screening of the CMB
- *Challenge*: foregrounds and analyzing high-resolution CMB data
- *Results*: axion constraints and a first look at Simons Observatory LAT data



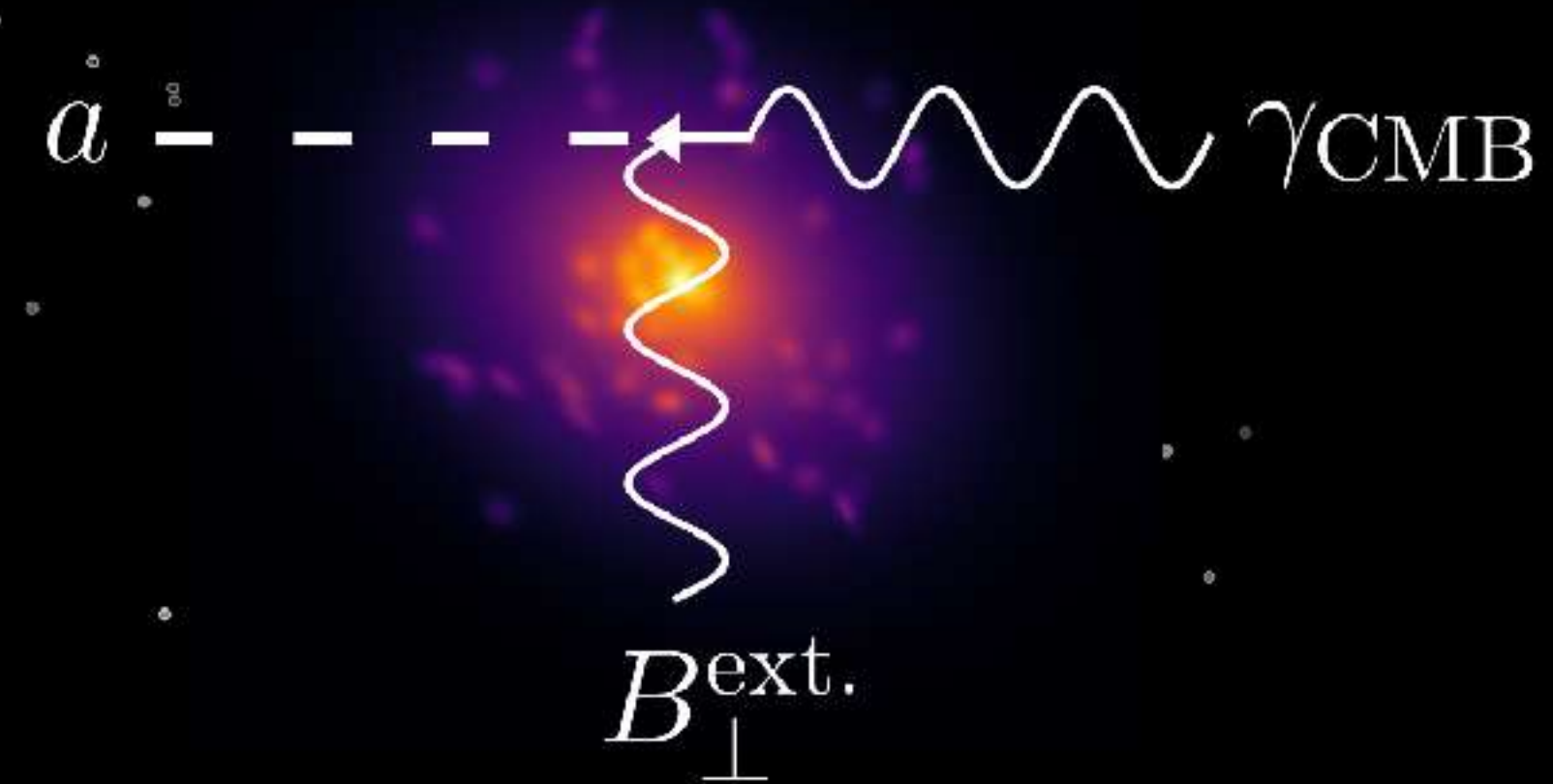
Part 2: Testing inflation with galaxy surveys

- *Background*: galaxy bias as a probe of inflation
- *Challenge*: galaxy clustering on ultra-large scales
- *Results*: first LSS constraint on τ_{NL} with DESI, plus a pipeline for its largest scales

Axion-induced screening of the CMB

Central questions:

1. *Can CMB photons convert into axions as they travel to us?*
2. *If so, how do we search for this?*



References

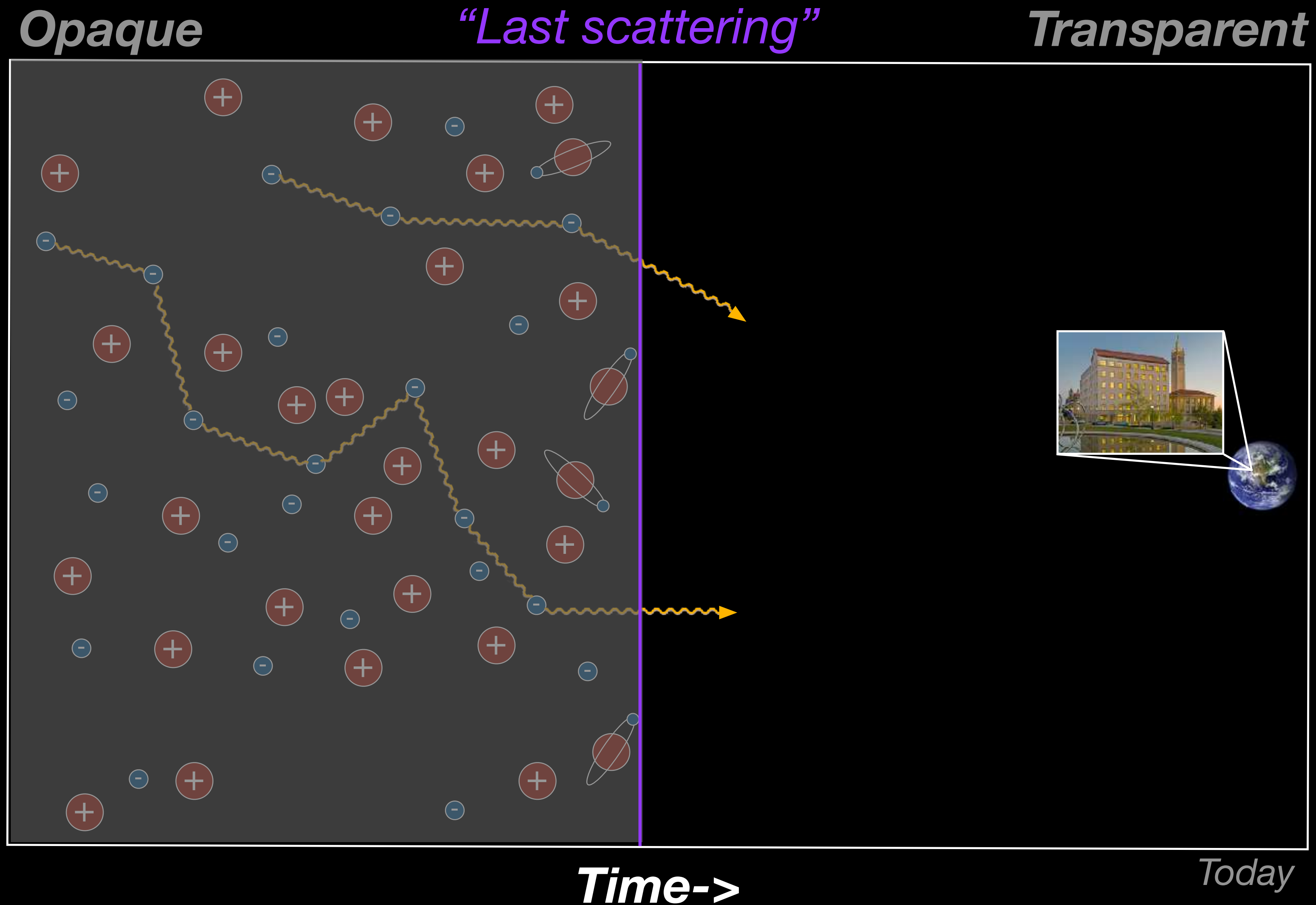
SG, McCarthy, Mondino, Hill, Huang, Johnson ([*PRL*, 2409.10514](#))

SG++ (Atacama Cosmology Telescope) in prep.

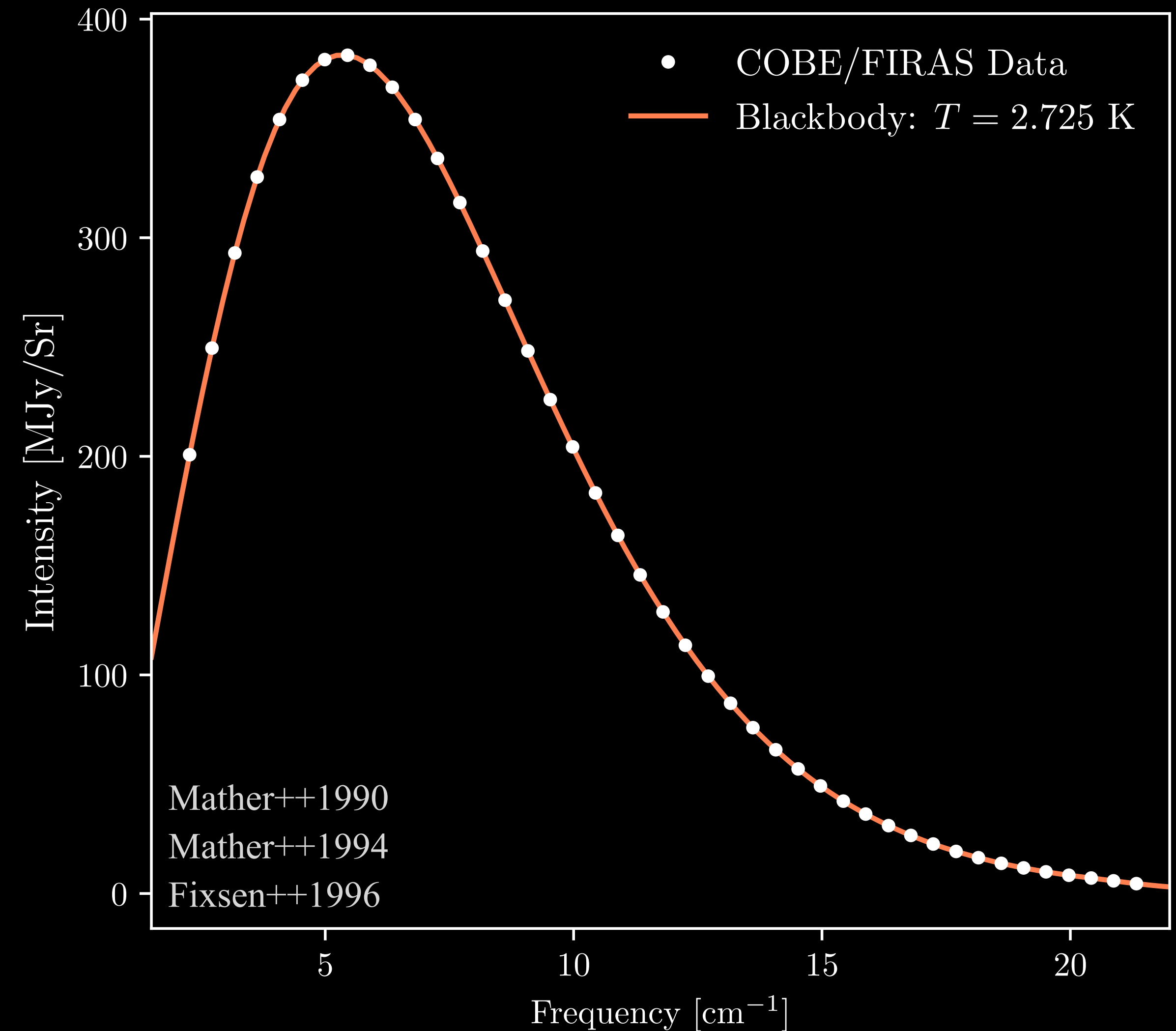
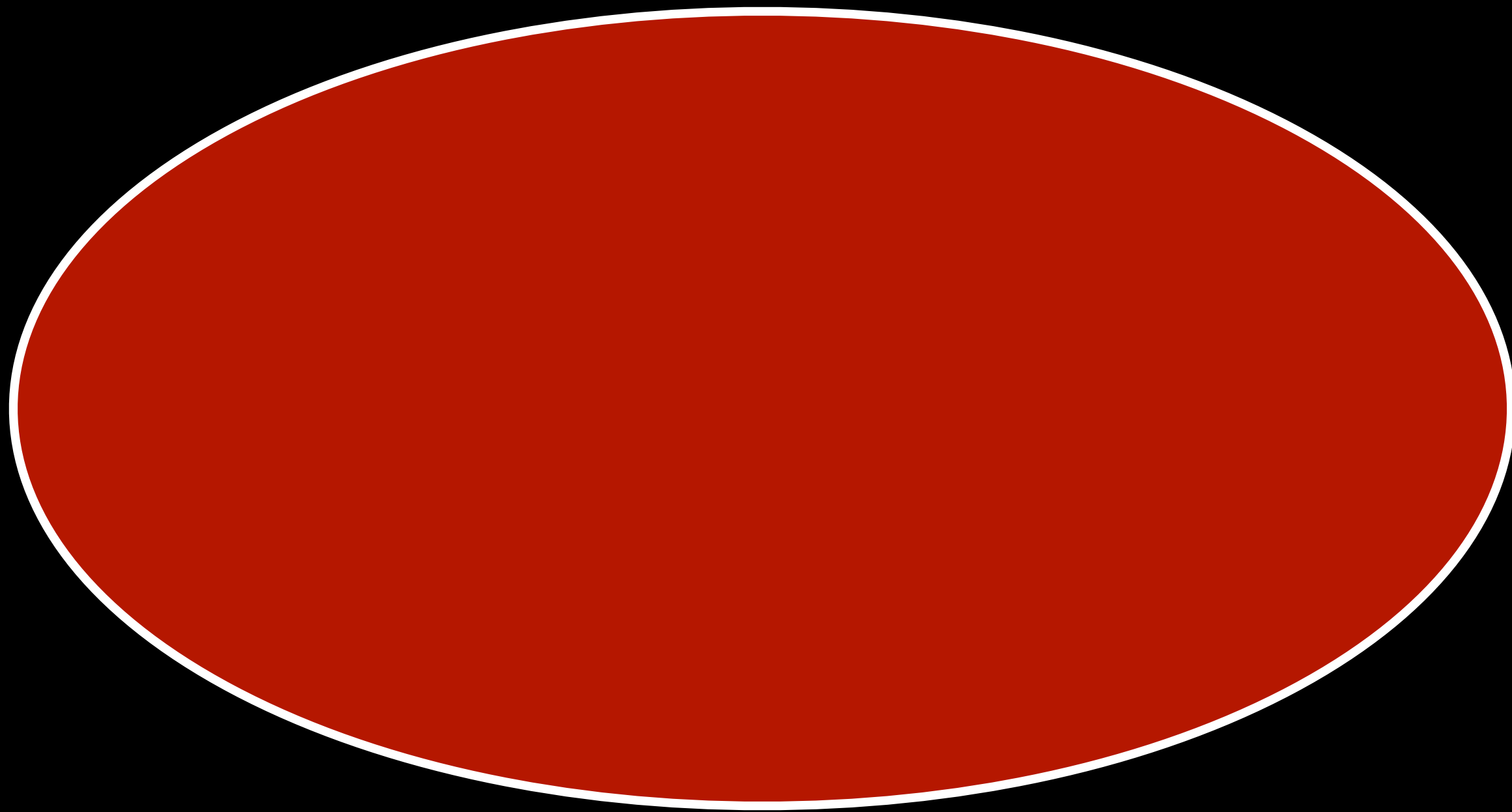
SG++ (Simons Observatory) in prep.

with J. Colin Hill, Junwu Huang, Matthew Johnson, Fiona McCarthy, Cristina Mondino

Cosmic microwave background background (CMB²)

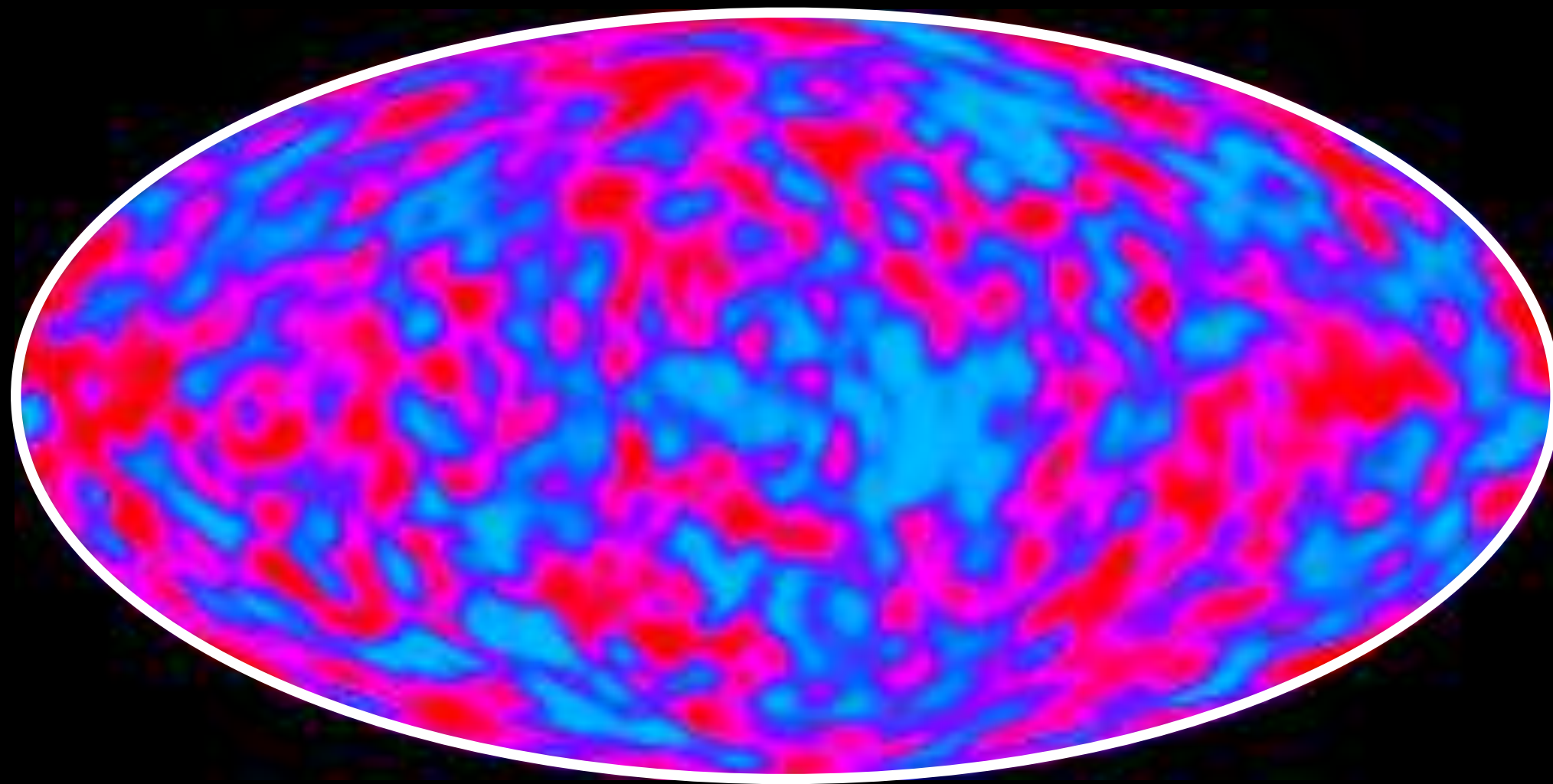


Cosmic microwave background background (CMB²)



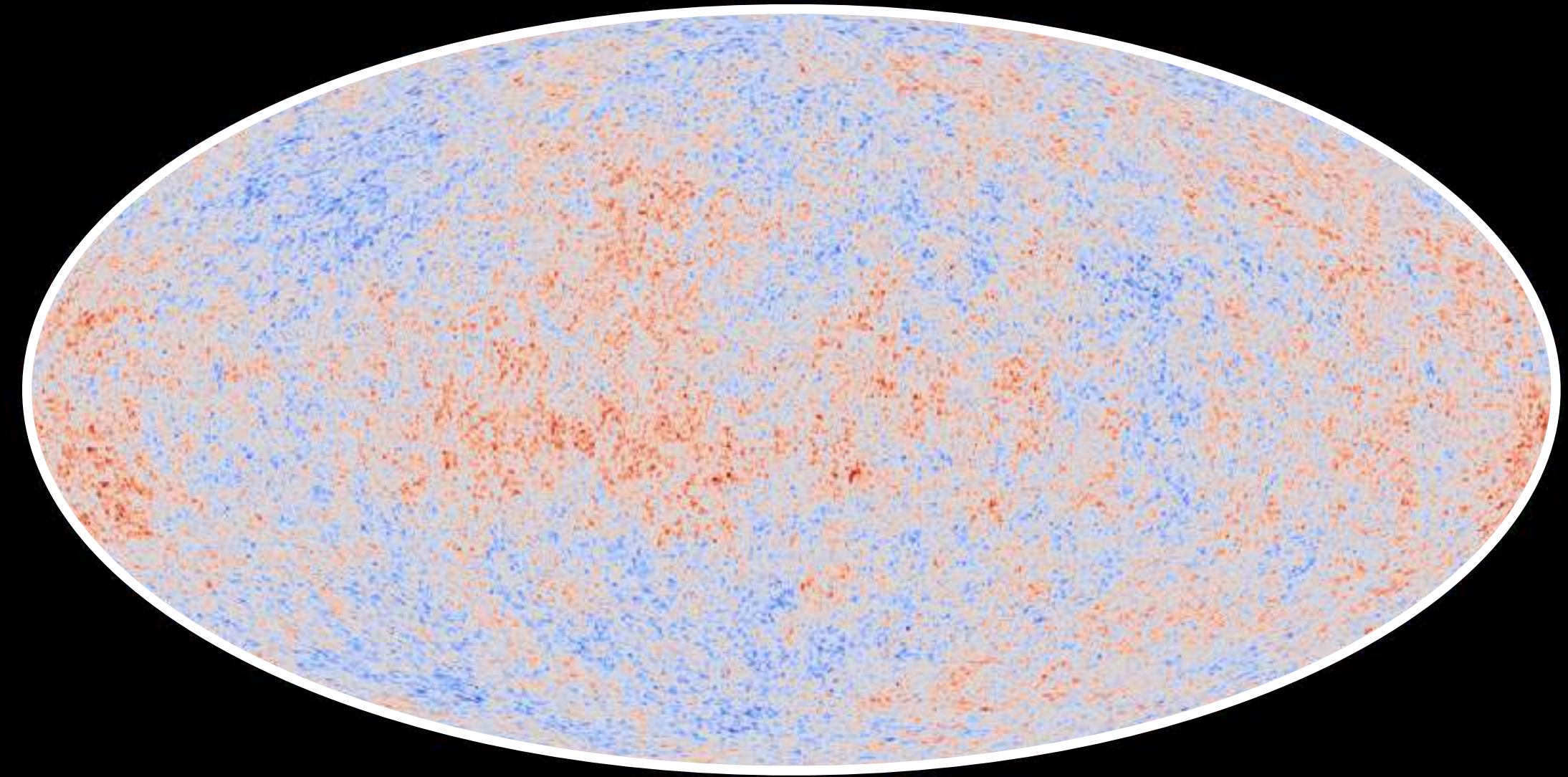
Sky-averaged (monopole) CMB is a *near*-perfect **blackbody**

CMB primary anisotropies



COBE

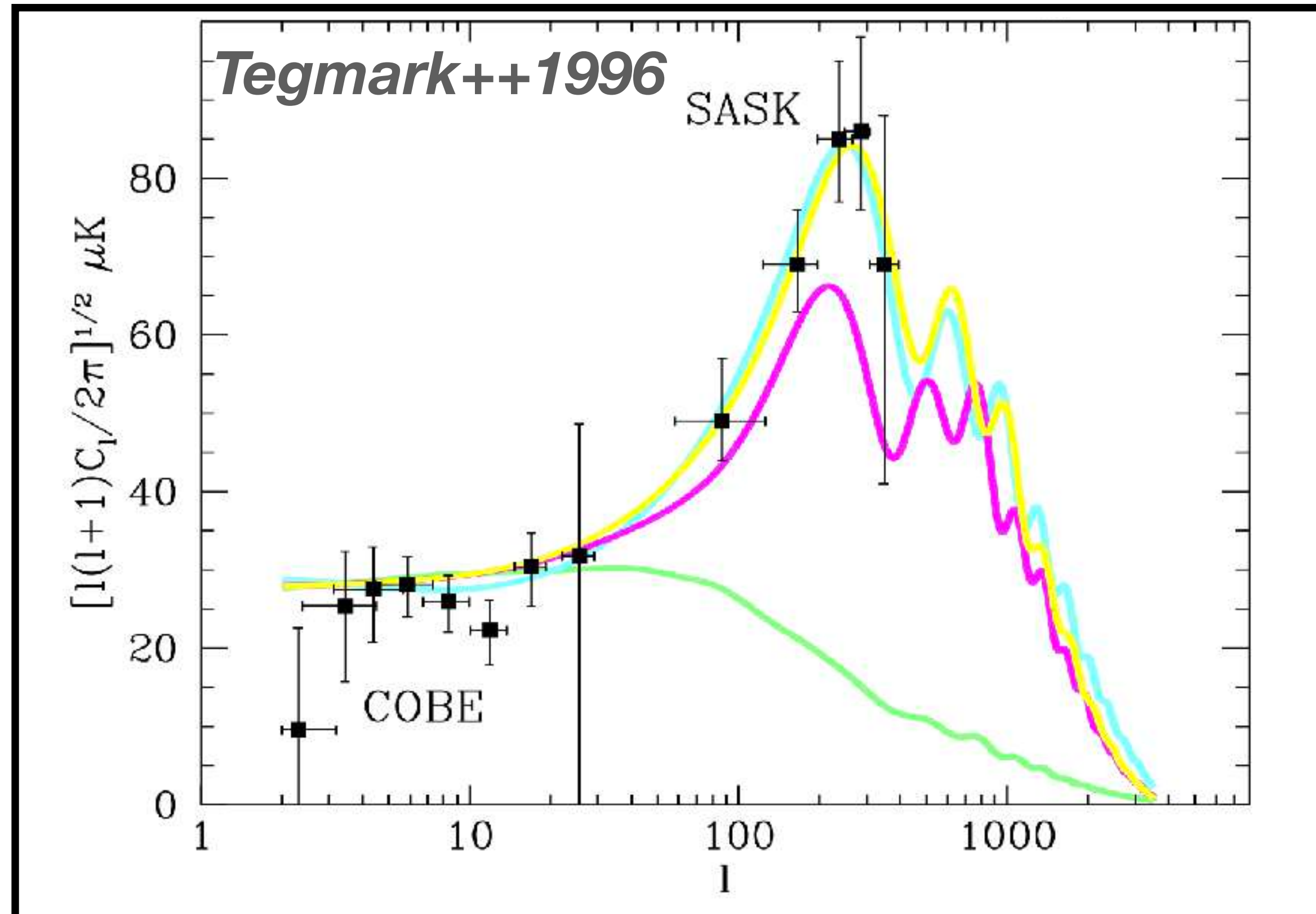
Smoot++1992, Bennett++1996



Planck

Planck Collaboration 2018, IV

Evolution of CMB power spectrum measurements

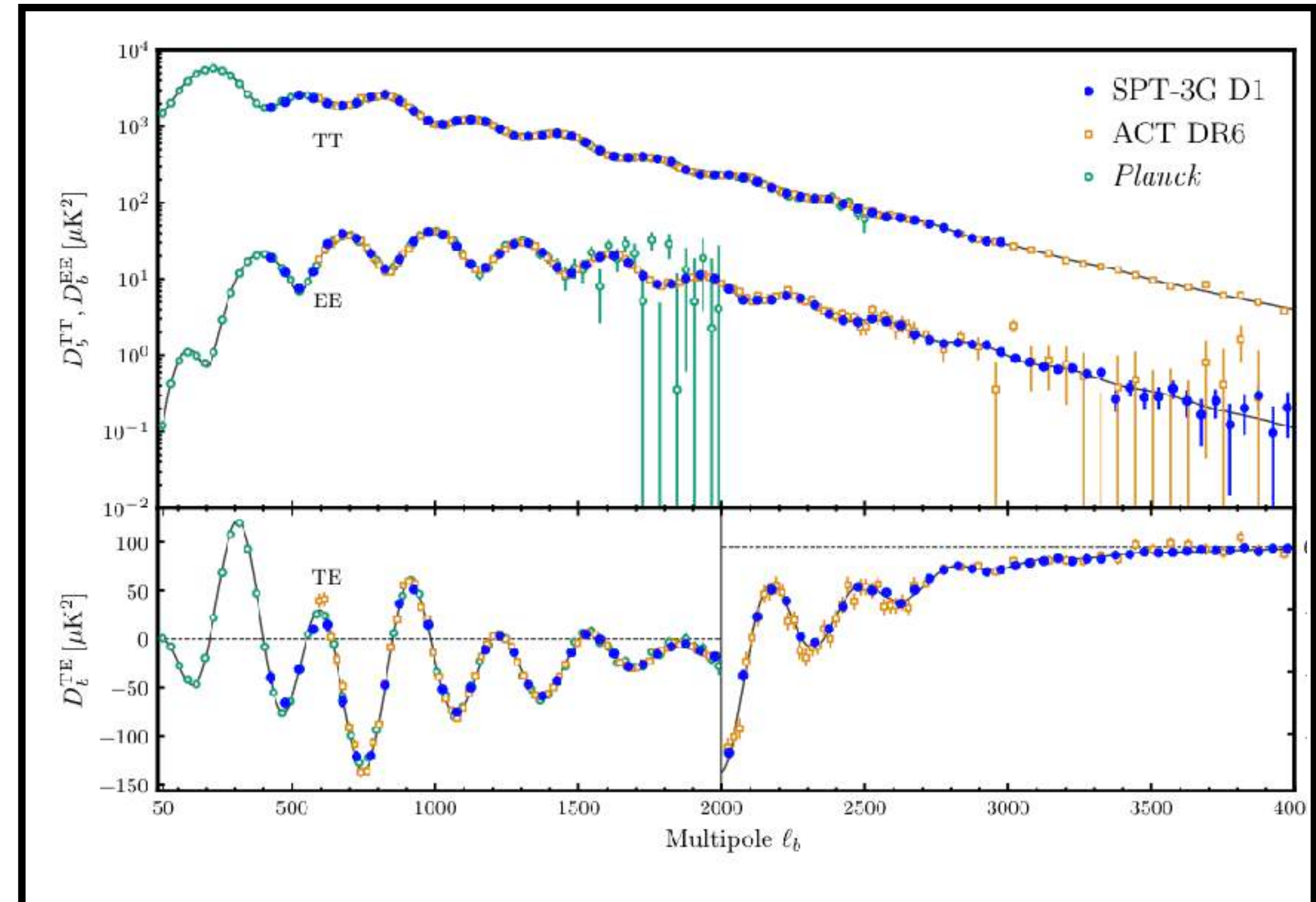


COBE

Smoot++1992

SASK

Netterfield++1996



ACT DR6:

Naess++2025

Louis++2025

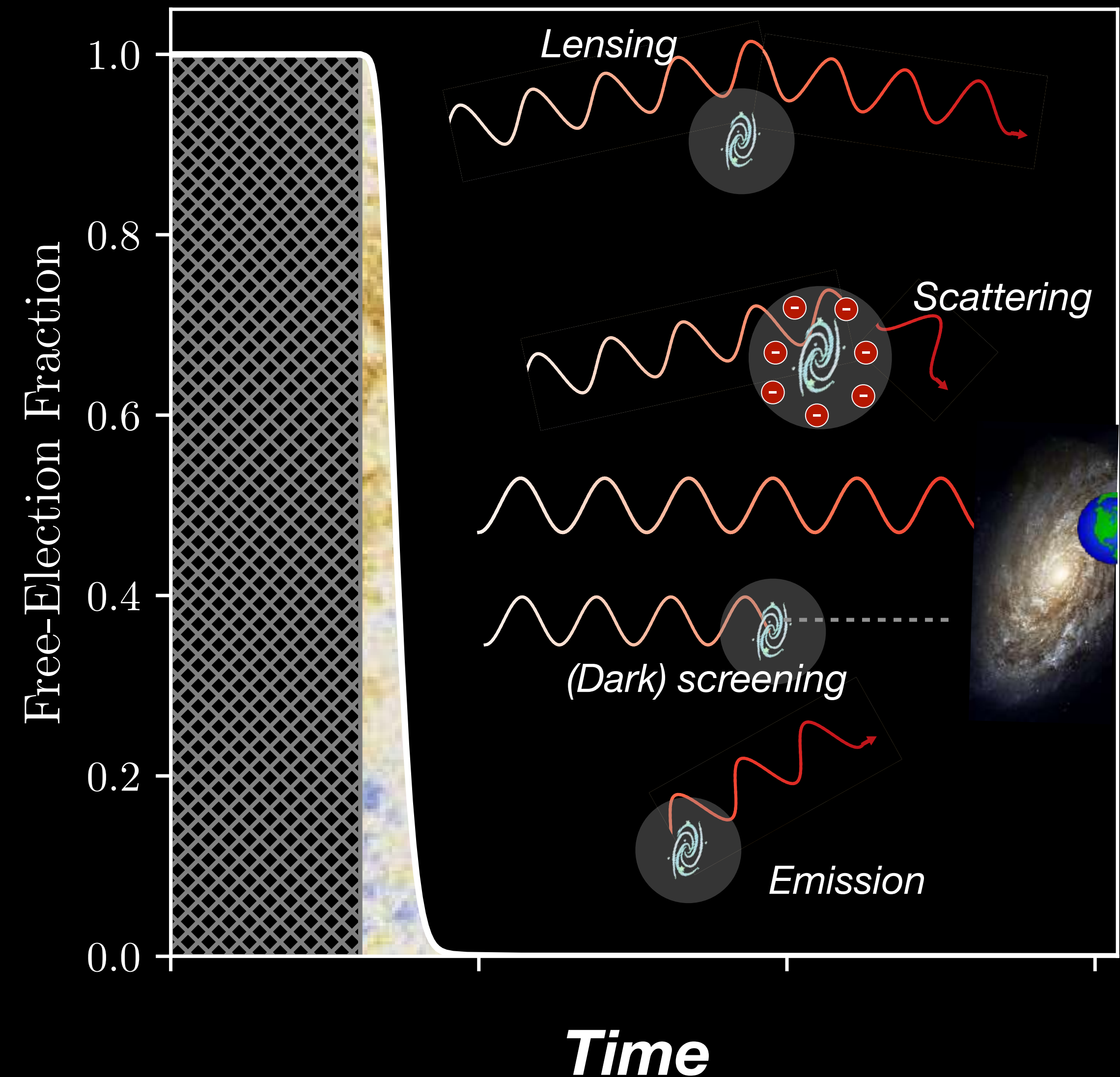
Calabrese++2025

SPT 3G:

Camhuis++2025

CMB Secondary Anisotropies

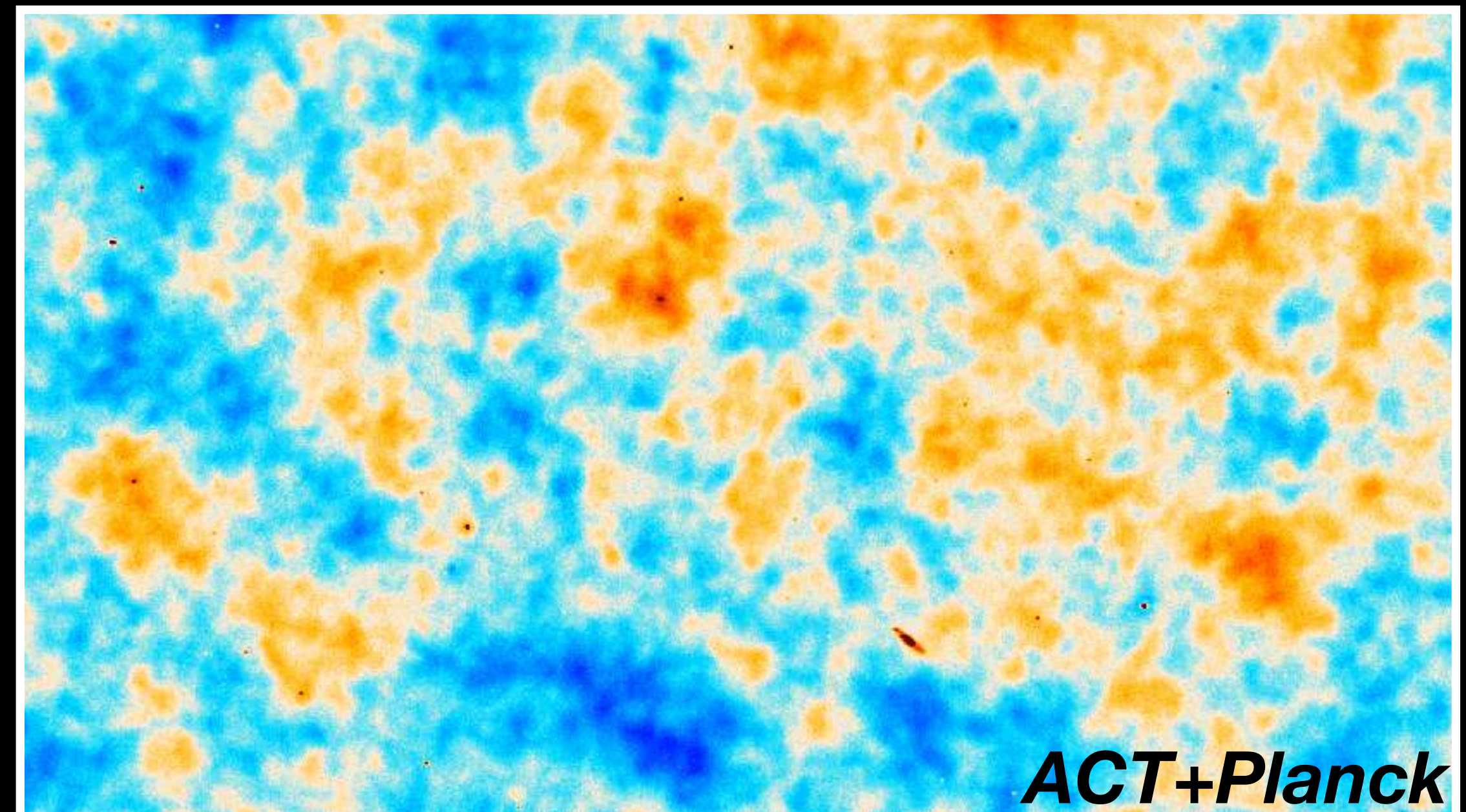
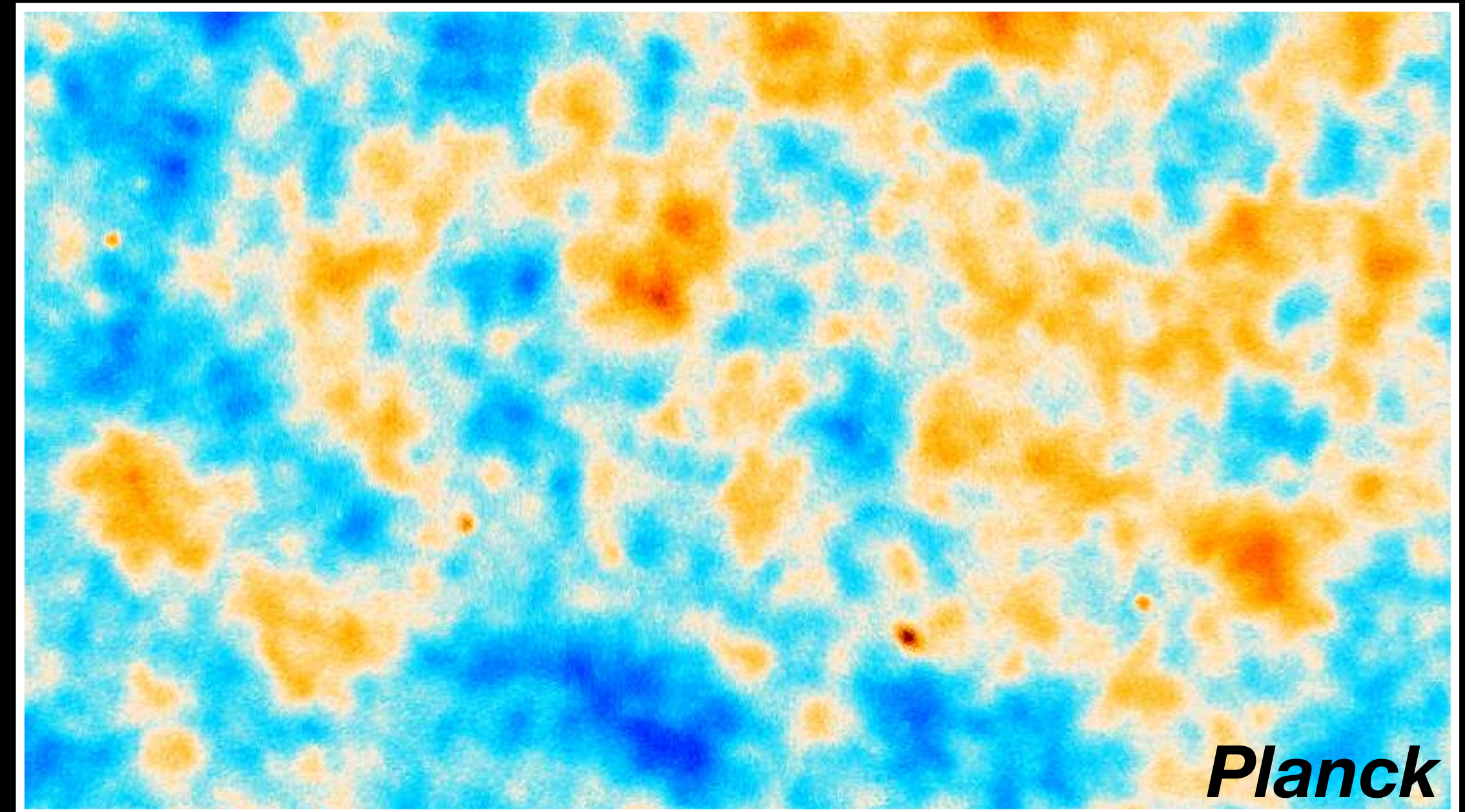
- *Secondary anisotropies*: modifications as photons propagate after last scattering
- **Powerful**: contain valuable information
 - **Cosmology**: new large-scale structure (LSS) tracers
 - **Astrophysics**: unique sensitivity to baryons
- **Challenging**: difficult to isolate/separate



CMB Secondary Anisotropies

- *Secondary anisotropies*: modifications as photons propagate after last scattering
- **Powerful**: contain valuable information
 - **Cosmology**: new large-scale structure (LSS) tracers
 - **Astrophysics**: unique sensitivity to baryons
- **Challenging**: difficult to isolate/separate

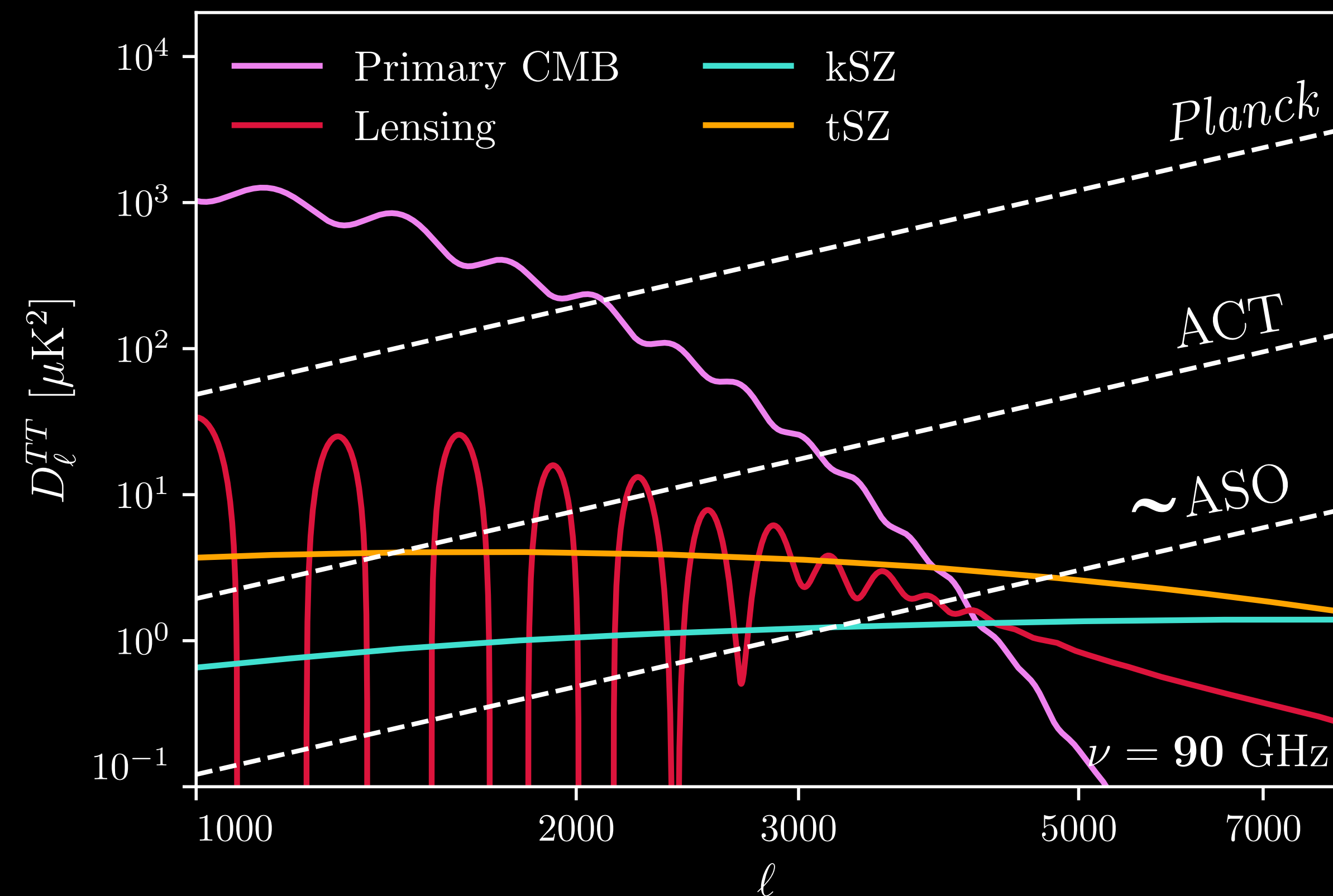
We are entering an exciting era for CMB secondaries with ACT, SPT, SO!



CMB Secondary Anisotropies

- *Secondary anisotropies*: modifications as photons propagate after last scattering
- **Powerful**: contain valuable information
 - **Cosmology**: new large-scale structure (LSS) tracers
 - **Astrophysics**: unique sensitivity to baryons
- **Challenging**: difficult to isolate/separate

We are entering an exciting era for CMB secondaries with ACT, SPT, SO!



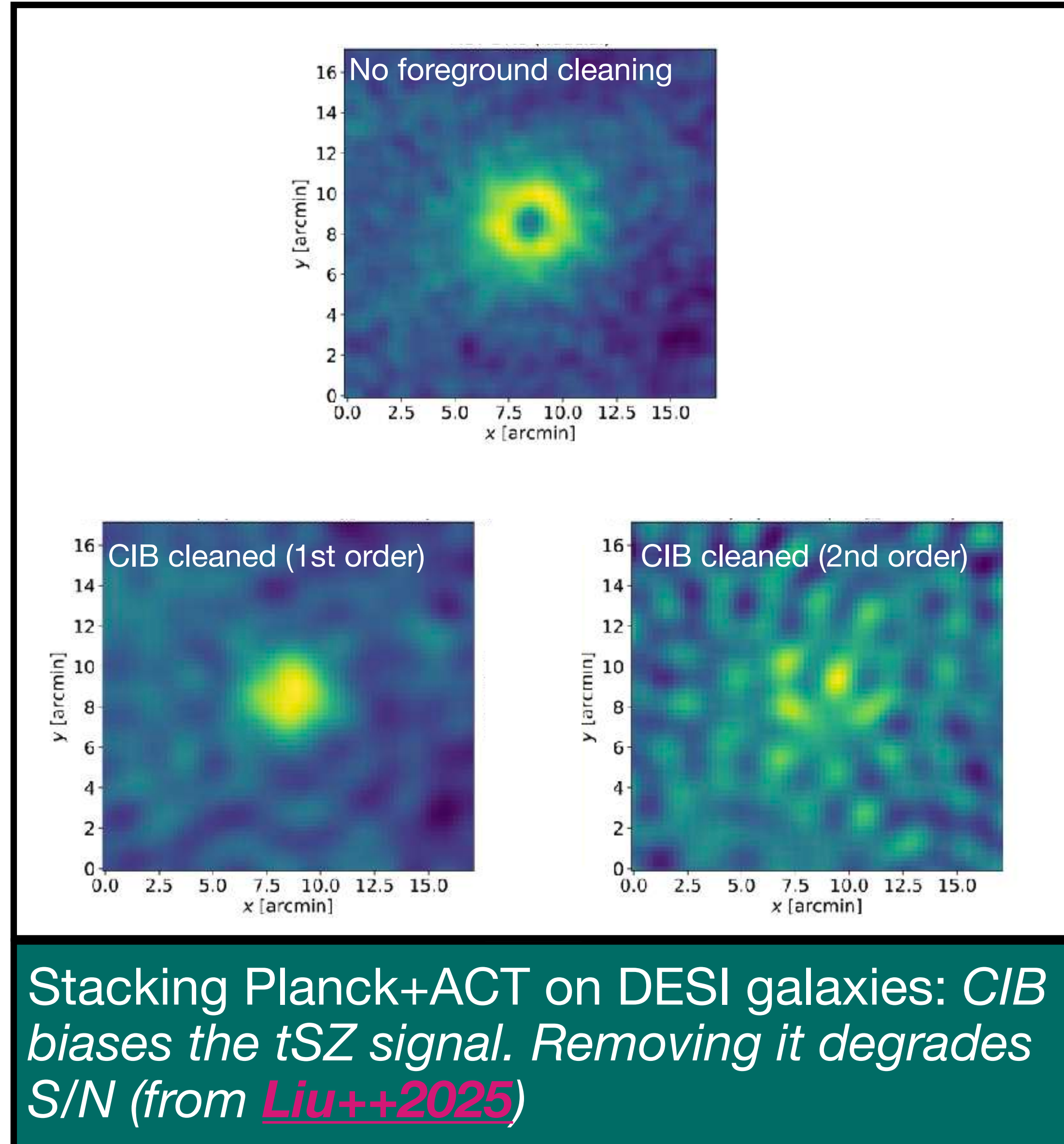
Qualitative illustration of near-term sensitivity to CMB secondaries.

A familiar example: the thermal Sunyaev–Zel’dovich (tSZ) effect

- **tSZ:** CMB photons inverse-Compton scatter off hot electrons

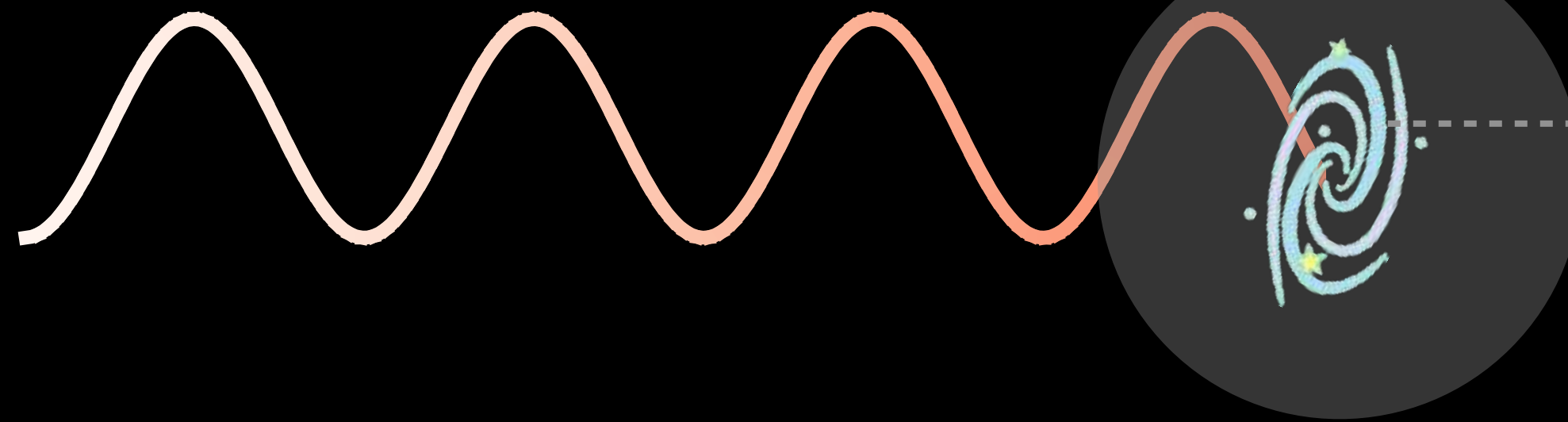
$$y(\hat{\mathbf{n}}) = \frac{\sigma_T}{m_e c^2} \int dl P_e(l)$$

- **In theory:** powerful probe
 - *Cosmology:* cluster abundance/ σ_8
 - *Baryons:* gas thermodynamics/feedback
- **In practice:** challenging
 - *Measurements:* foregrounds, esp. cosmic infrared background (CIB)
 - *Modeling:* uncertain galaxy-halo connection

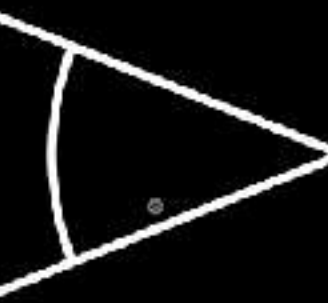


A new secondary? Patchy dark screening

CMB photon

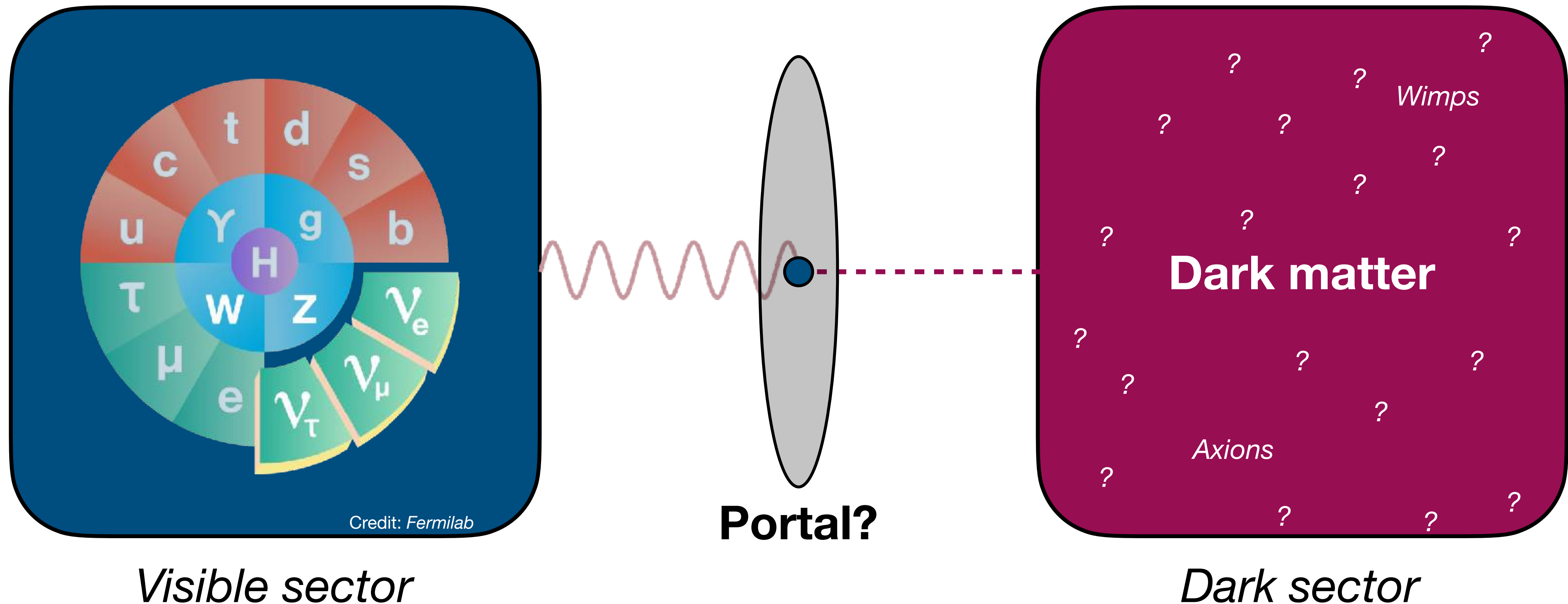


Dark sector (BSM)
particle



The dark sector

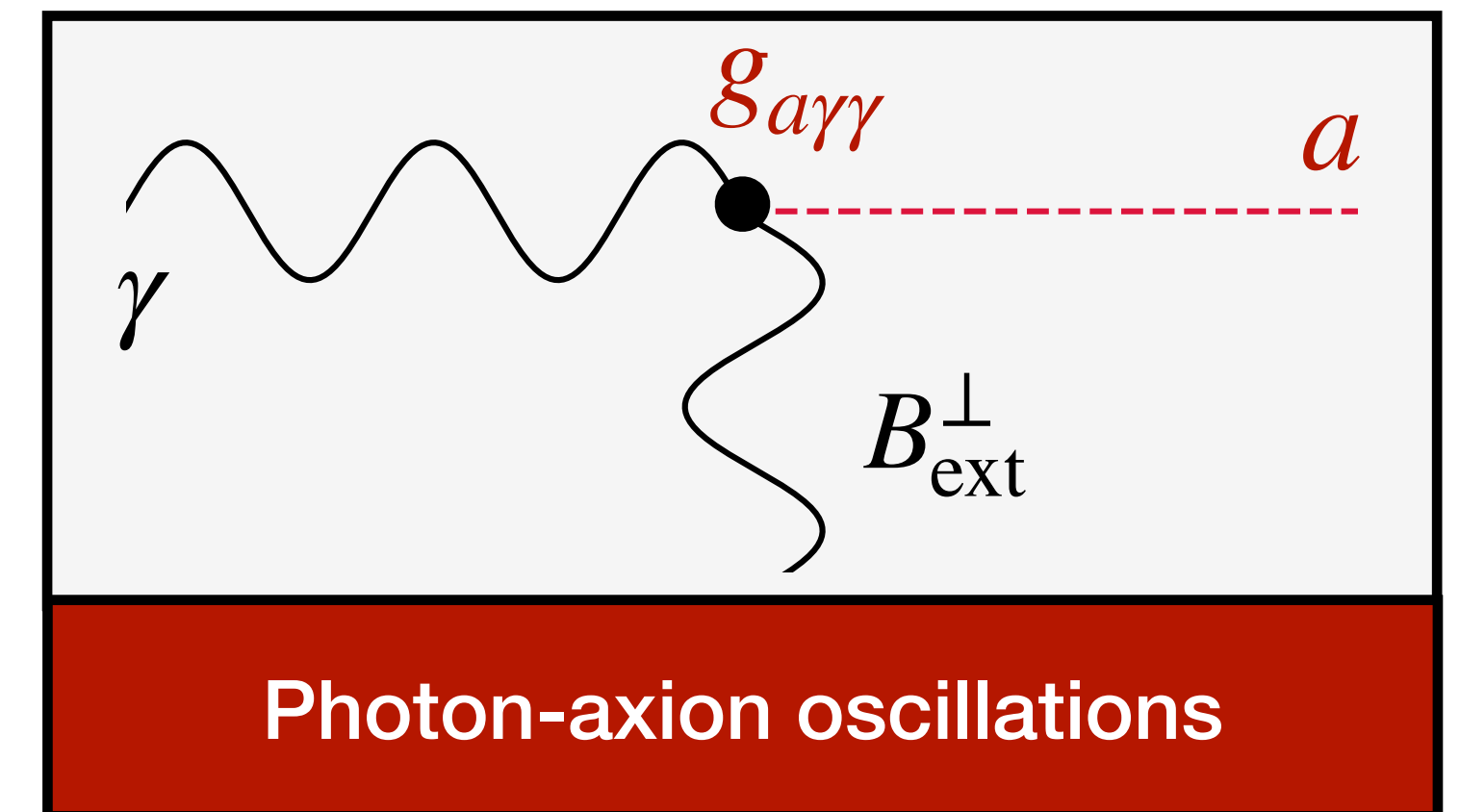
- **Dark matter** and dark energy provide significant motivation to search for beyond standard model (BSM) fields



The dark sector

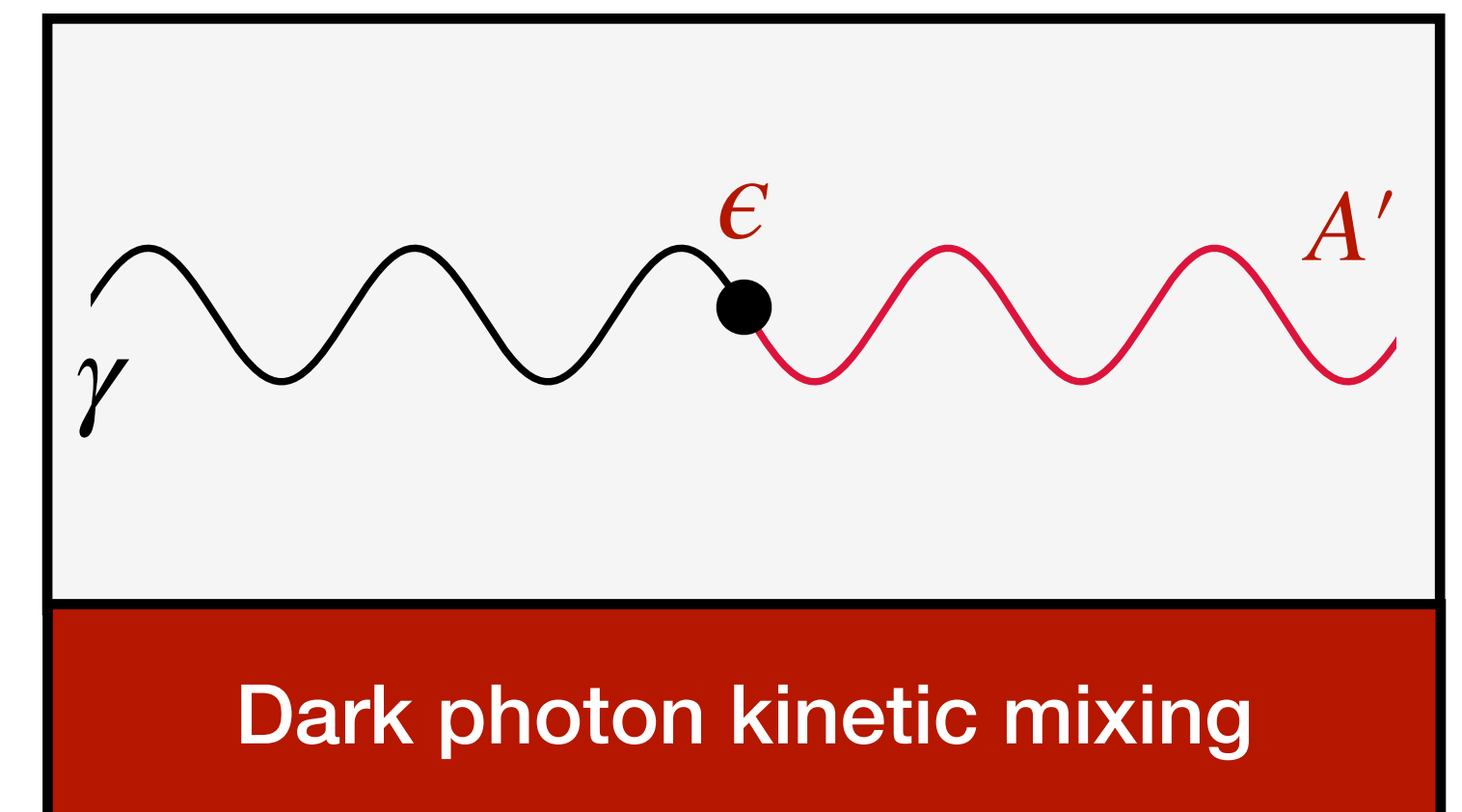
- What can the CMB photons interact with?
- (Pseudo)-scalar portal: axion(-like particles)*

$$\mathcal{L} \supset -\frac{1}{2}\partial_\mu a \partial^\mu a - \frac{1}{2}m_a^2 a^2 - \frac{1}{4}g_{a\gamma\gamma} a F_{\mu\nu} \tilde{F}^{\mu\nu} \longrightarrow$$



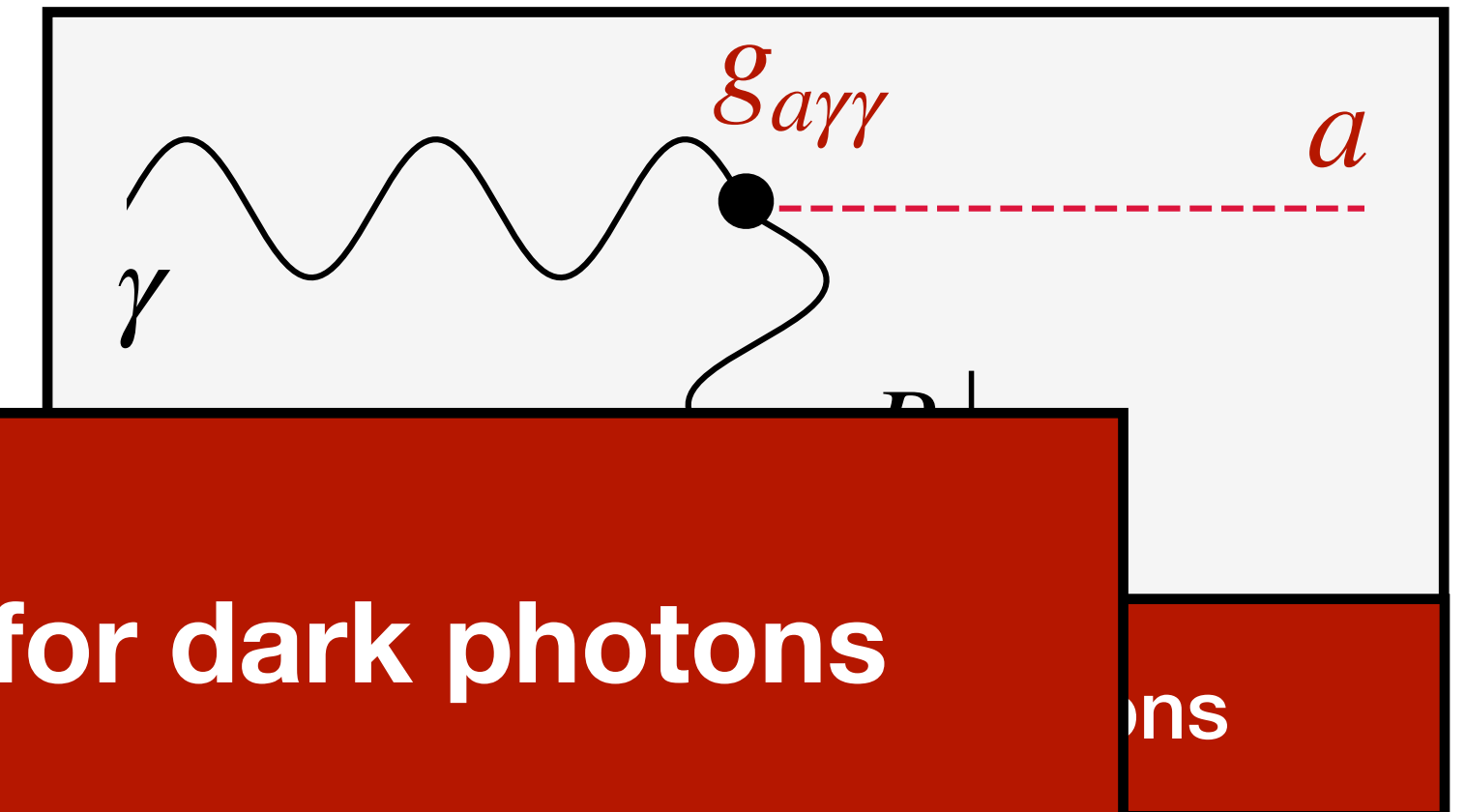
- Vector portal: dark photons*

$$\mathcal{L} \supset -\frac{1}{4}F'_{\mu\nu} F'^{\mu\nu} - \frac{1}{2}m_{\text{DP}}^2 A'_\mu A'^\mu - \frac{\epsilon}{2} F_{\mu\nu} F'^{\mu\nu} \longrightarrow$$



The dark sector

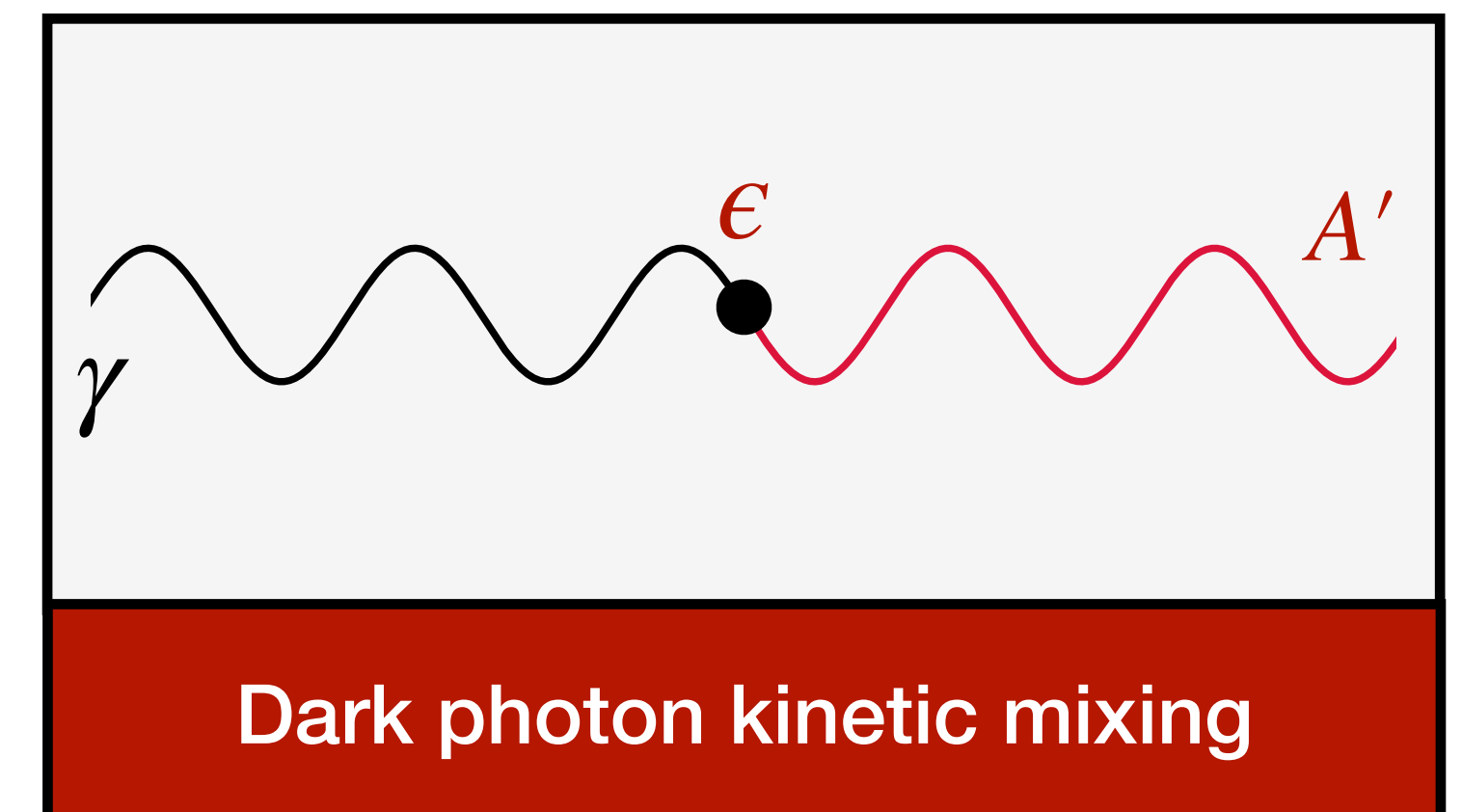
- What can the CMB photons interact with?
- *(Pseudo)-scalar portal: axion(-like particles)*



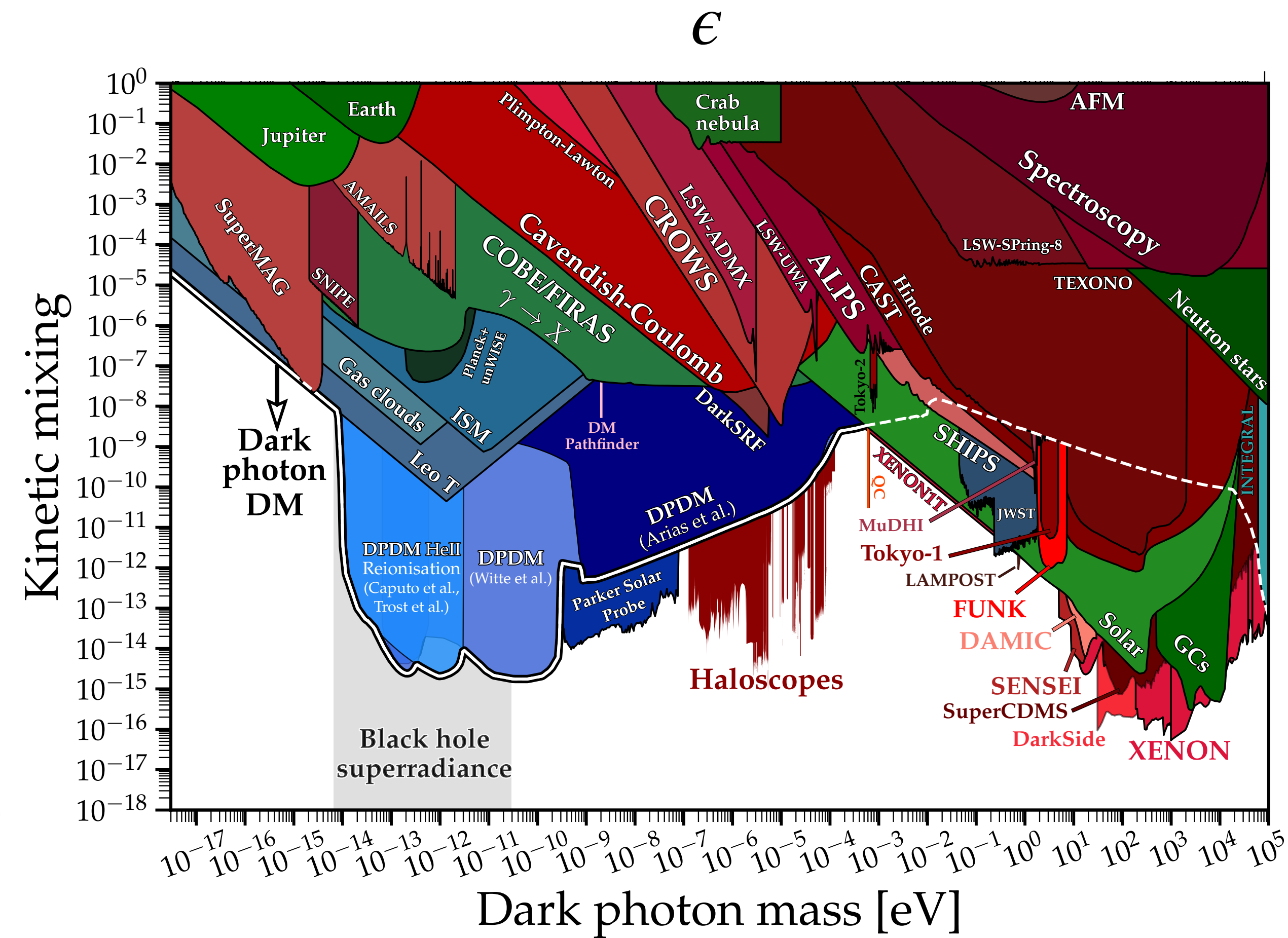
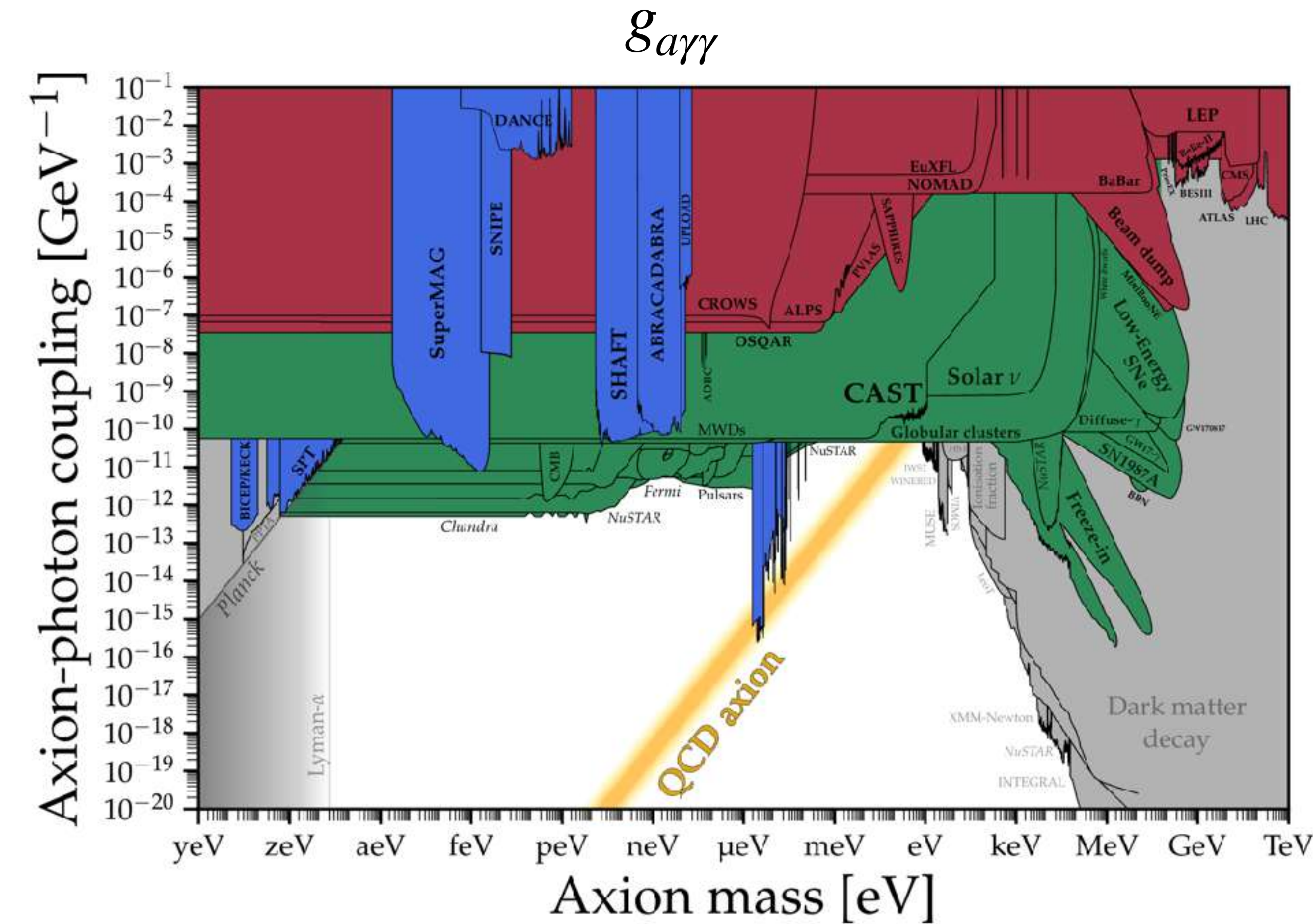
I will mostly focus on axions, but story is similar for dark photons

- *Vector portal: dark photons*

$$\mathcal{L} \supset -\frac{1}{4}F'_{\mu\nu}F'^{\mu\nu} - \frac{1}{2}m_{\text{DP}}^2 A'_\mu A'^\mu - \frac{\epsilon}{2}F_{\mu\nu}F'^{\mu\nu} \longrightarrow$$



Substantial effort devoted to axion/dark photon searches



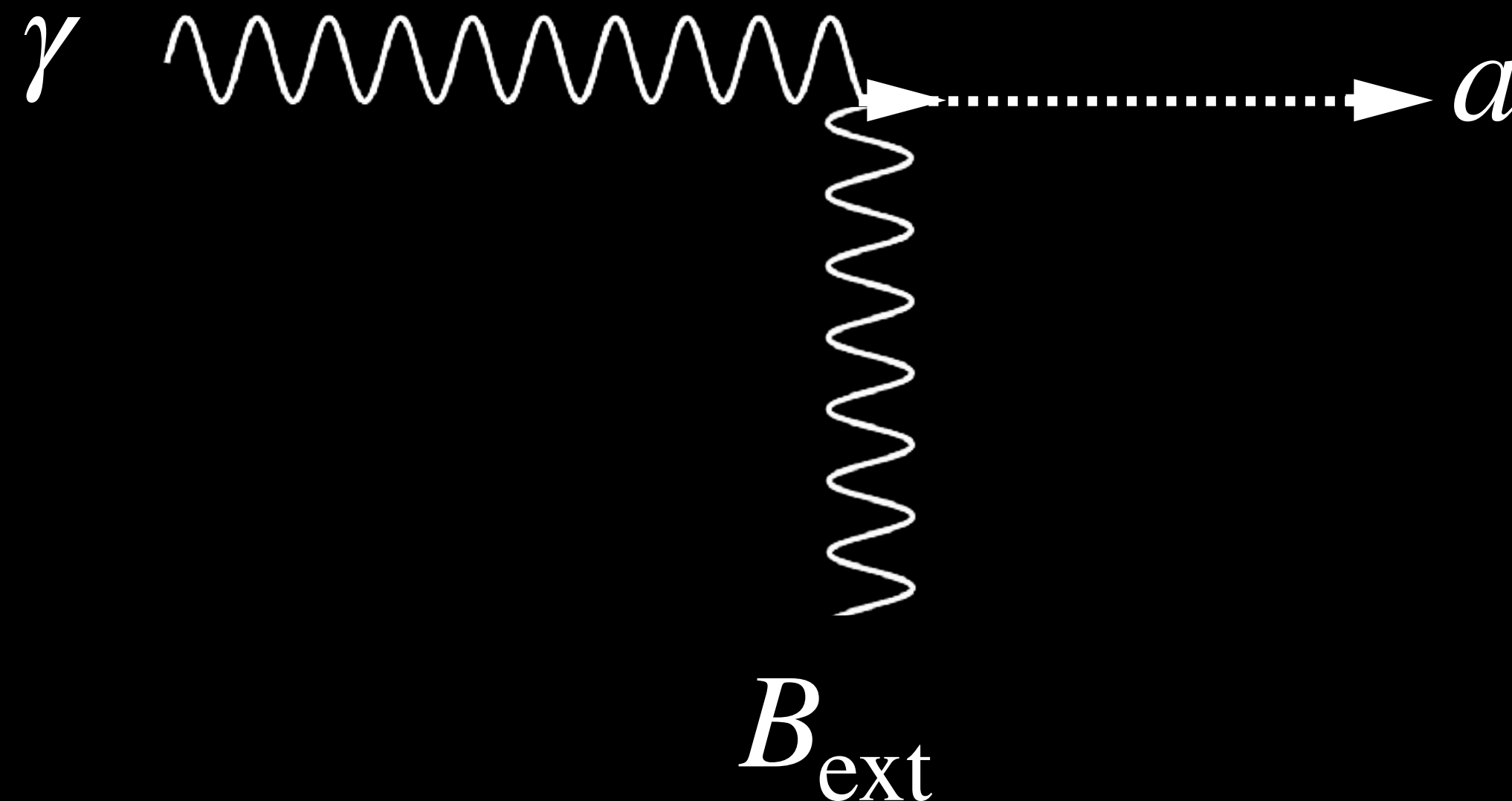
Existing constraints span wide range of masses (<https://cajohare.github.io/AxionLimits/>)

How to find an axion?

- Axions **modify Maxwell's equations**

$$\mathcal{L} \supset g_{a\gamma\gamma} a \mathbf{E} \cdot \mathbf{B} \implies \begin{cases} \nabla \cdot \mathbf{E} = \rho_0 + g_{a\gamma\gamma} (\nabla a) \cdot \mathbf{B} \\ \nabla \times \mathbf{B} = \partial_t \mathbf{E} + \mathbf{J} + g_{a\gamma\gamma} (\mathbf{B} \partial_t a - \mathbf{E} \times \nabla a) \\ \nabla \cdot \mathbf{B} = 0 \\ \nabla \times \mathbf{E} = -\partial_t \mathbf{B} \end{cases}$$

- Photons can convert to axions in a magnetic field



How to find an axion?

- Axions **modify Maxwell's equations**

$$\mathcal{L} \supset g_{a\gamma\gamma} a \mathbf{E} \cdot \mathbf{B} \implies \begin{cases} \nabla \cdot \mathbf{E} = \rho_0 + g_{a\gamma\gamma} (\nabla a) \cdot \mathbf{B} \\ \nabla \times \mathbf{B} = \partial_t \mathbf{E} + \mathbf{J} + g_{a\gamma\gamma} (\mathbf{B} \partial_t a - \mathbf{E} \times \nabla a) \\ \nabla \cdot \mathbf{B} = 0 \\ \nabla \times \mathbf{E} = -\partial_t \mathbf{B} \end{cases}$$

- Photons can convert to axions in a magnetic field

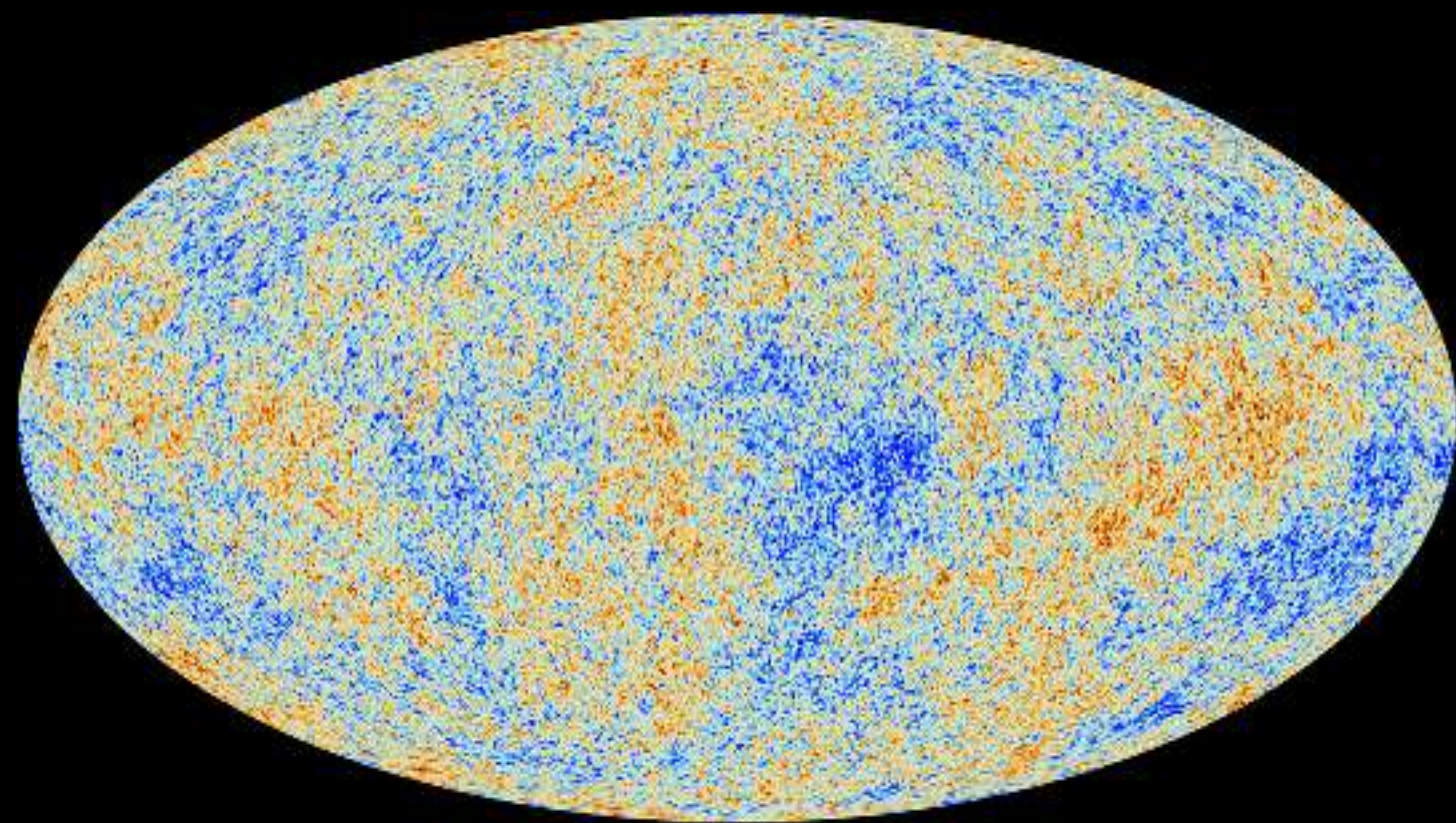


How to find an axion?

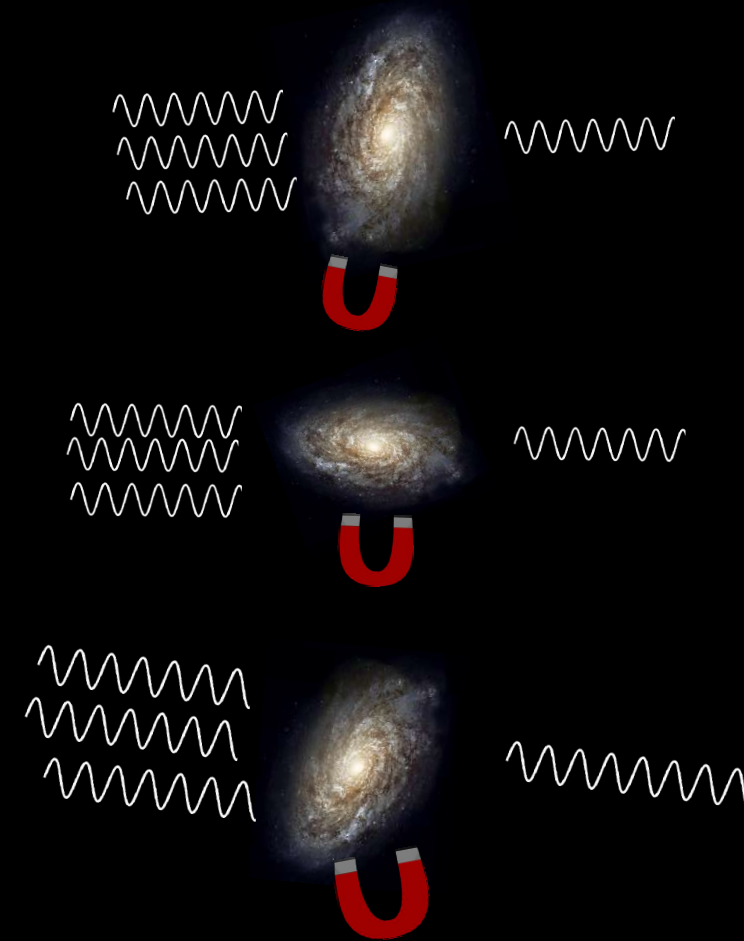
- Axions **modify Maxwell's equations**

$$\mathcal{L} \supset g_{a\gamma\gamma} a \mathbf{E} \cdot \mathbf{B} \implies \begin{cases} \nabla \cdot \mathbf{E} = \rho_0 + g_{a\gamma\gamma} (\nabla a) \cdot \mathbf{B} \\ \nabla \times \mathbf{B} = \partial_t \mathbf{E} + \mathbf{J} + g_{a\gamma\gamma} (\mathbf{B} \partial_t a - \mathbf{E} \times \nabla a) \\ \nabla \cdot \mathbf{B} = 0 \\ \nabla \times \mathbf{E} = -\partial_t \mathbf{B} \end{cases}$$

- Photons can convert to axions in a magnetic field

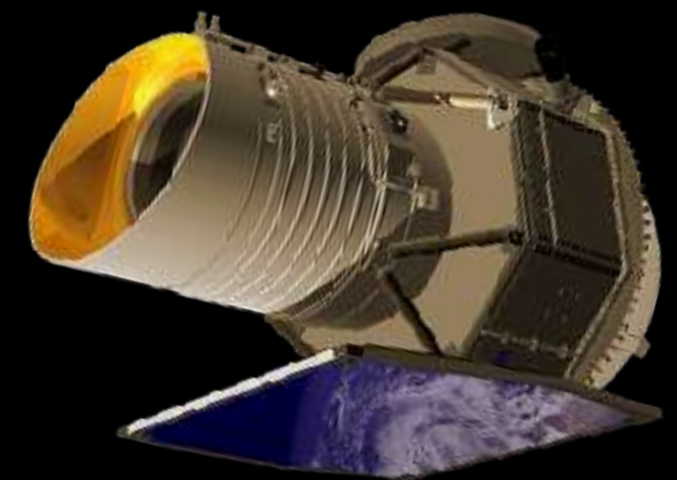
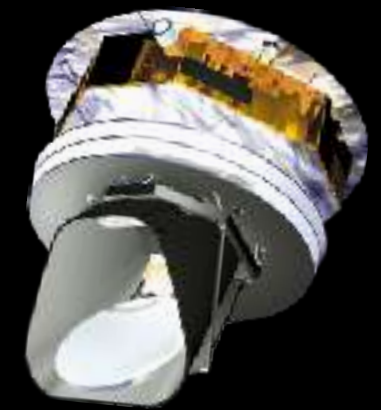


Photon source=CMB

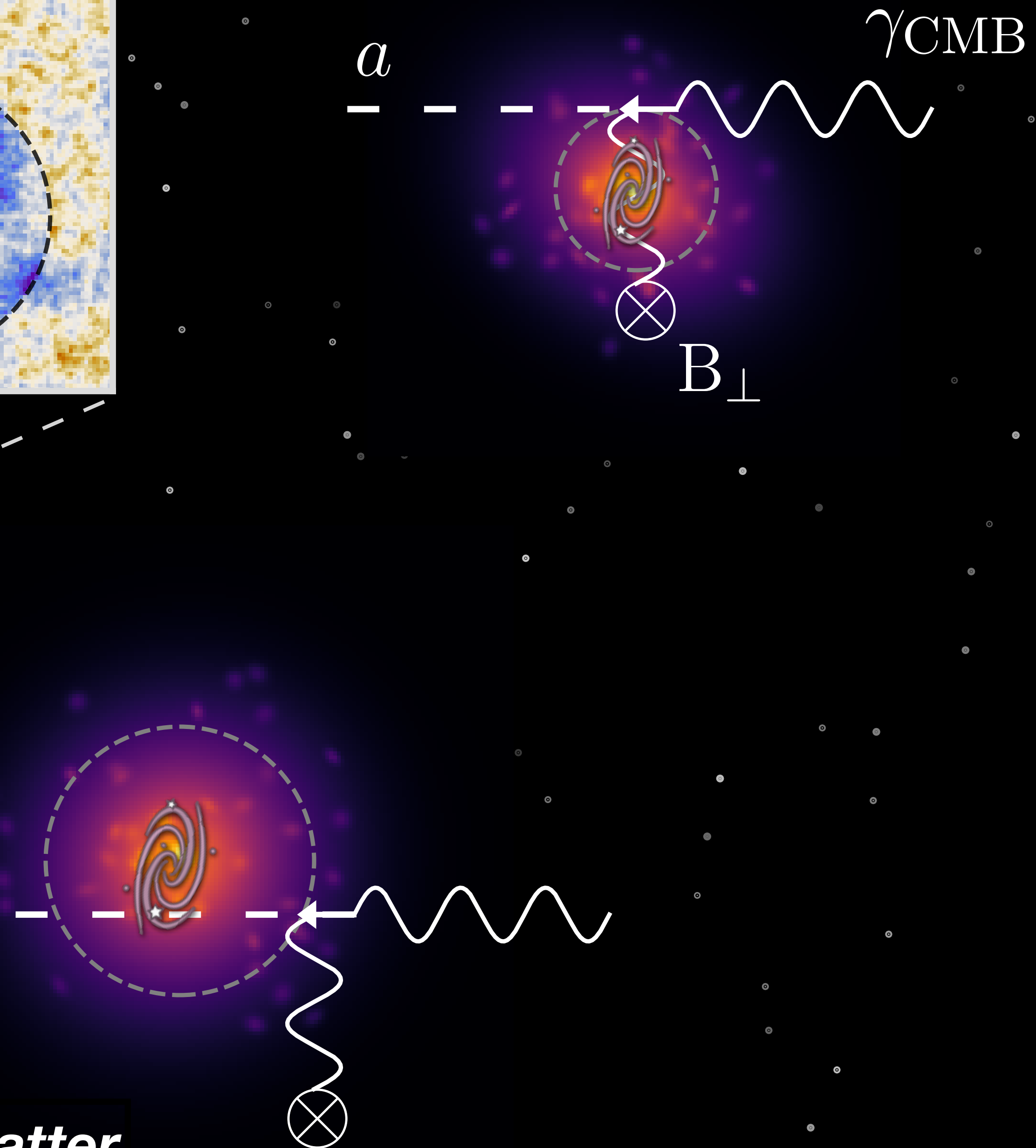
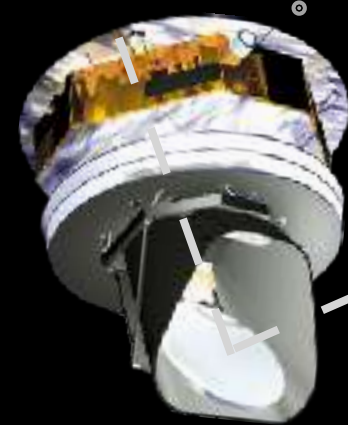
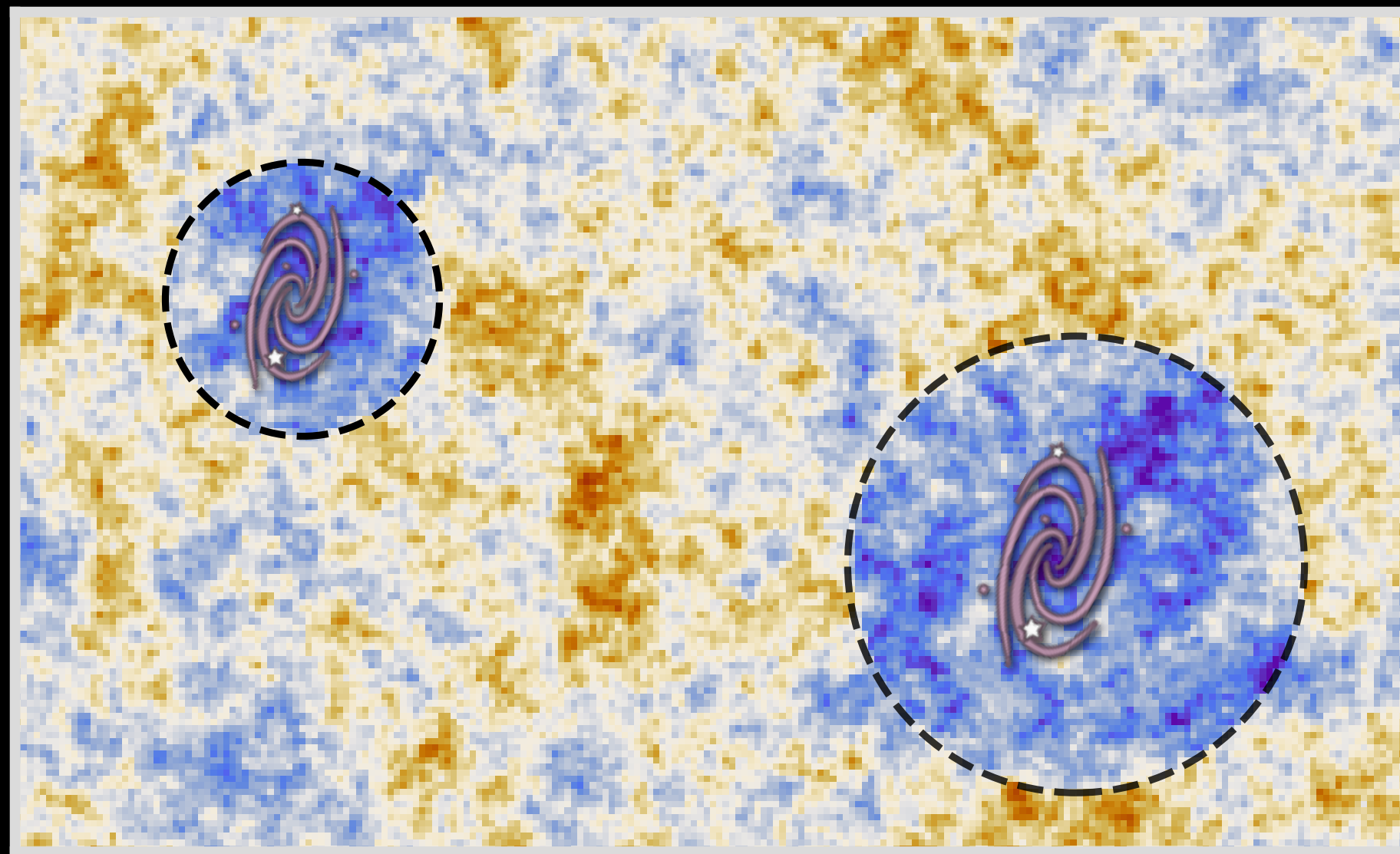


B-field=30 million galaxies

Planck

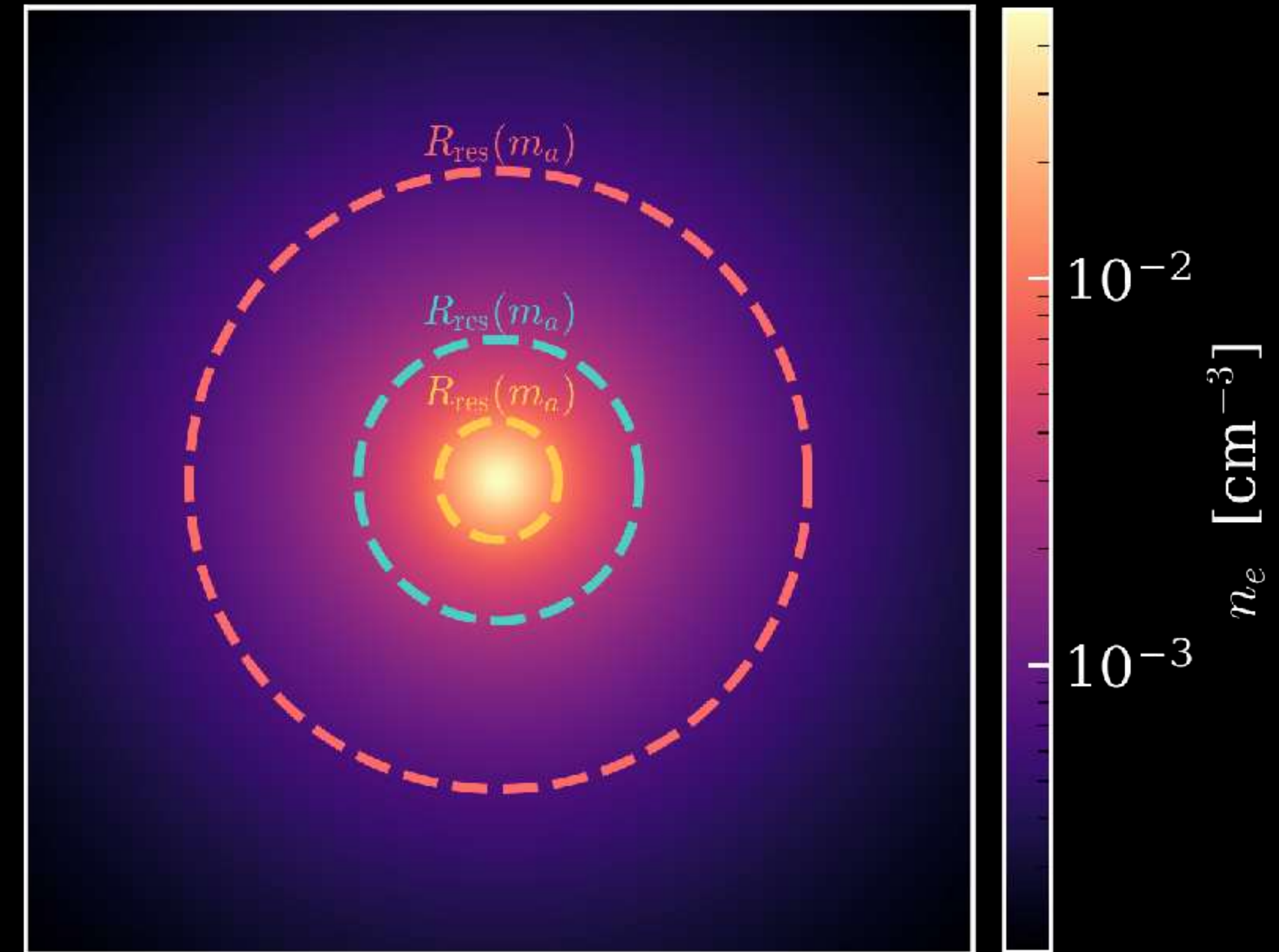
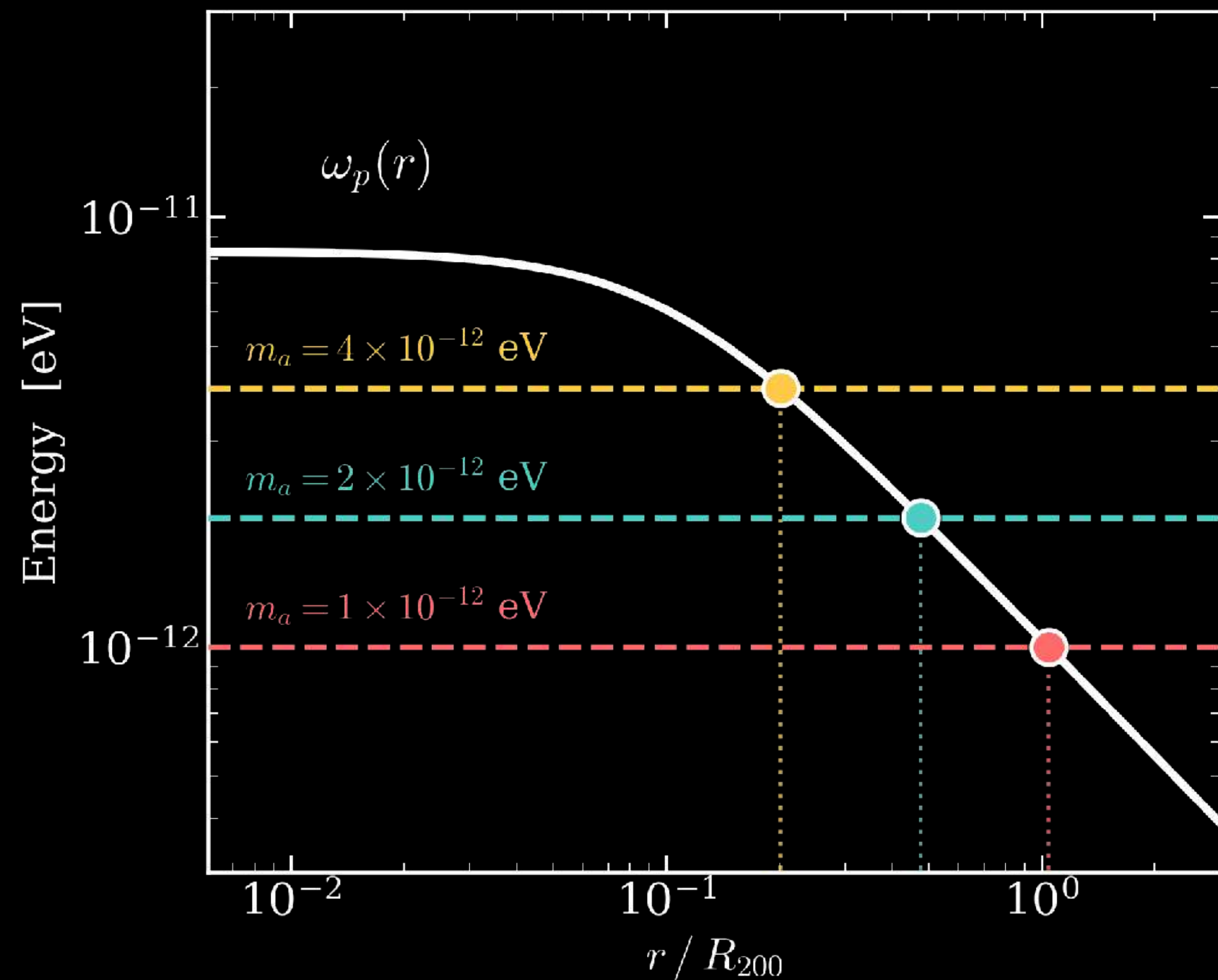


WISE



Note: we do not assume axions are dark matter

Resonant photon-axion conversion in halos

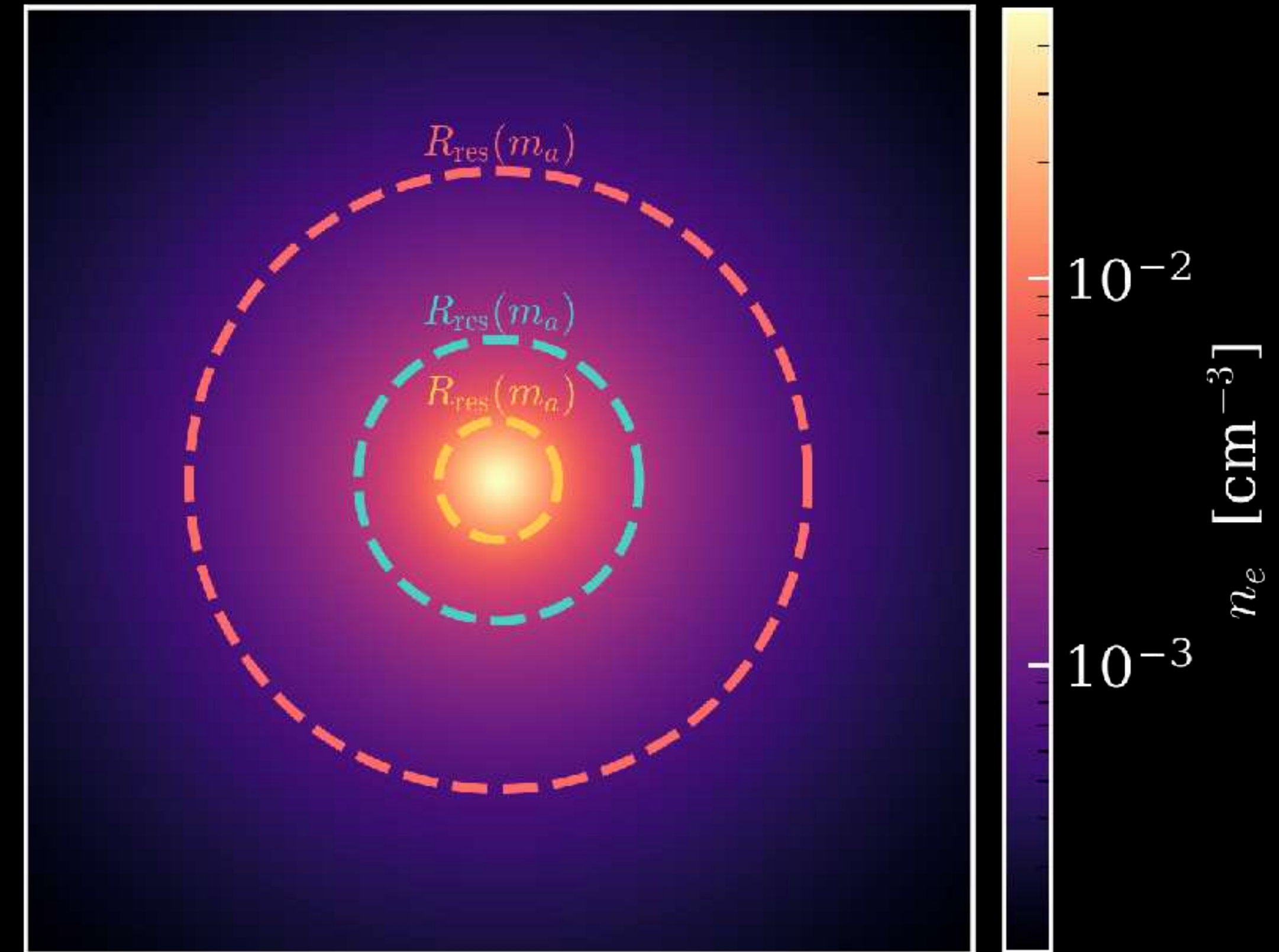
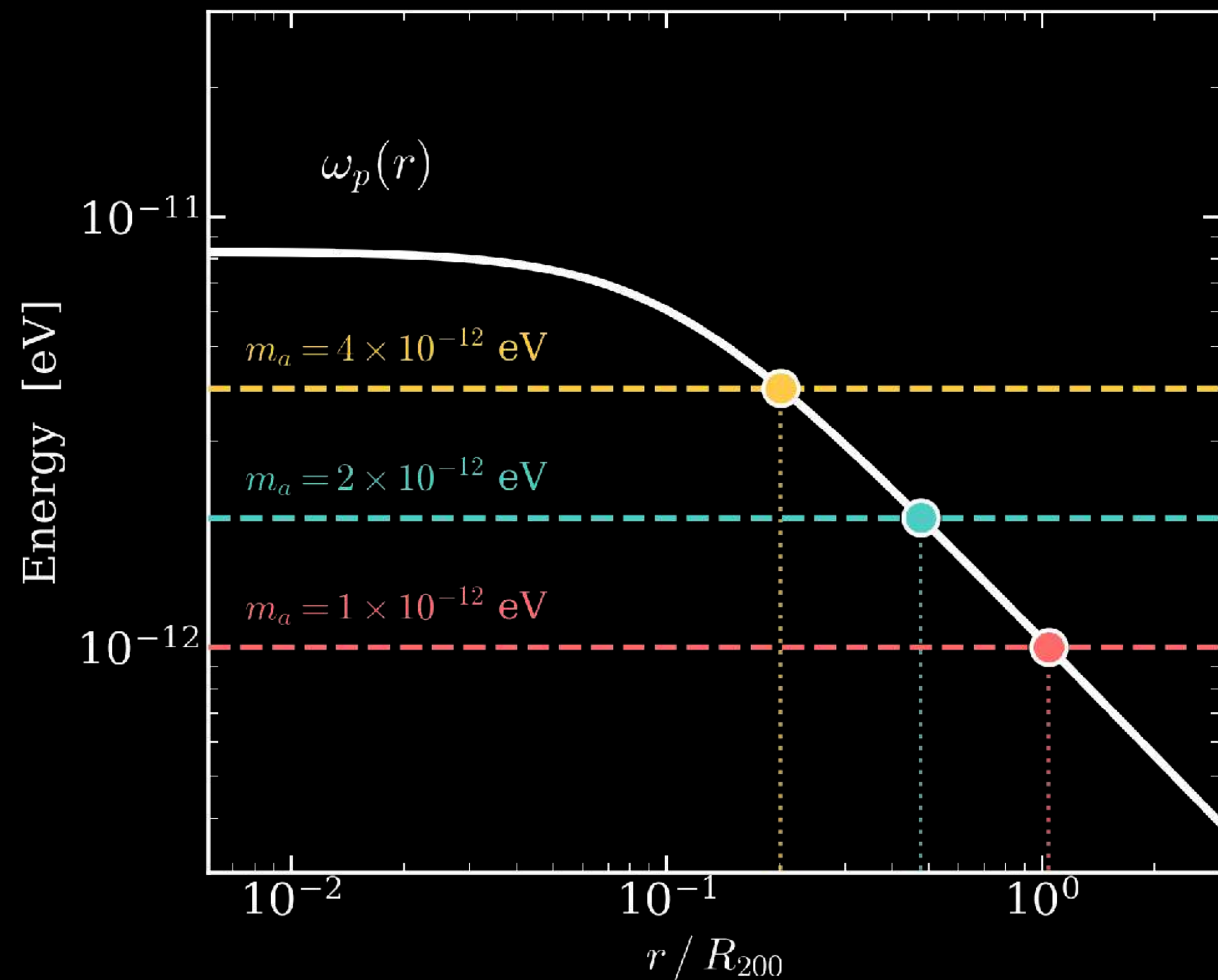


$$P_{\gamma \rightarrow a}^{\text{res}}(R_{\text{res}}, \nu) \propto \frac{\nu g_{a\gamma\gamma}^2 |\mathbf{B}_{\perp}^{\text{res}}|^2}{m_a^2} \left| \frac{d \ln n_e}{dr} \right|_{R_{\text{res}}}^{-1}$$

Conversion prob. given by Landau-Zener

Parke++1986 , Mirizzi++2009,
Tashiro++2013, Pîrvu++2023,
Mondino++2024

Resonant photon-axion conversion in halos



$$P_{\gamma \rightarrow a}^{\text{res}}(R_{\text{res}}, \nu) \propto \frac{\nu g_{a\gamma\gamma}^2 |\mathbf{B}_{\perp}^{\text{res}}|^2}{m_a^2} \left| \frac{d \ln n_e}{dr} \right|_{R_{\text{res}}}^{-1}$$

Conversion prob. given by Landau-Zener

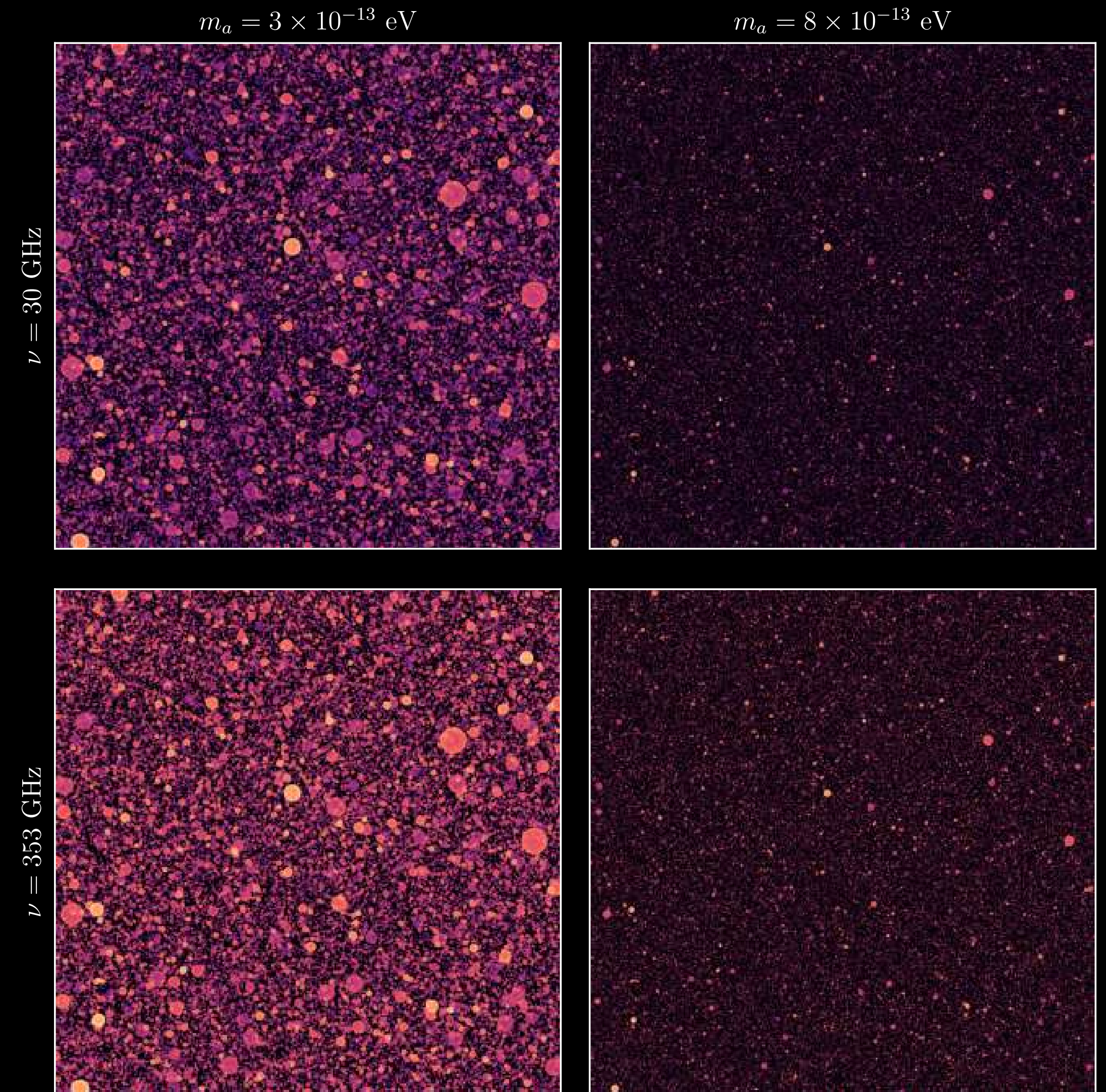
Parke++1986 , Mirizzi++2009,
Tashiro++2013, Pîrvu++2023,
Mondino++2024

The dark-screening optical depth

- CMB is “screened” by dark sector conversions -> described by optical depth

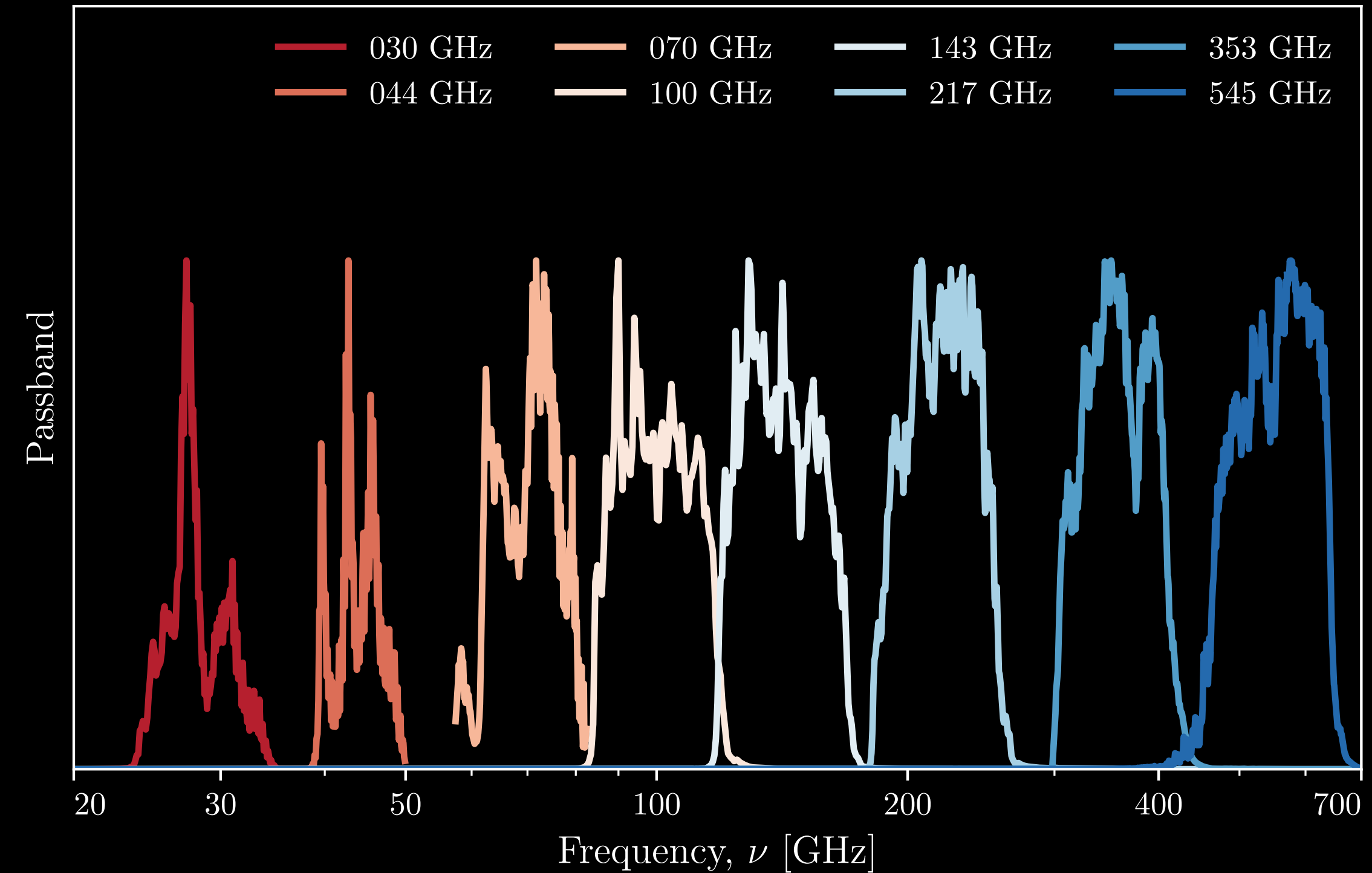
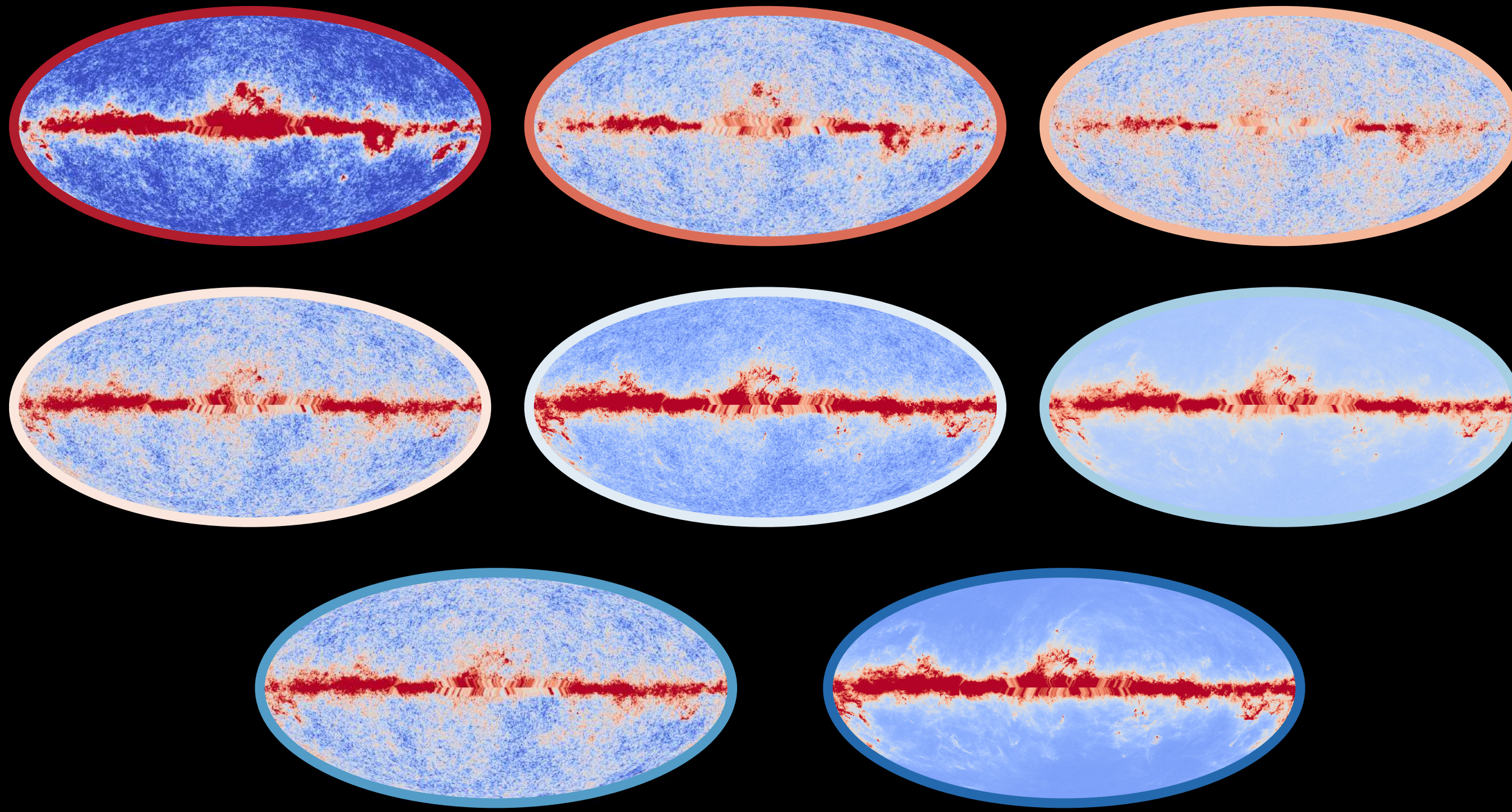
$$\frac{\Delta T}{T_{\text{CMB}}}(\hat{\mathbf{n}}, \nu) \equiv -\tau_a(\hat{\mathbf{n}}, \nu) \propto \sum_{\text{halos along } \hat{\mathbf{n}}} P_{\gamma \rightarrow a}^{\text{res}}(\hat{\mathbf{n}}, \nu)$$

- Two keys for isolating this signal
 - Spectral info:** screening has non-thermal frequency dependence
 - Spatial info:** screening traces large-scale structure



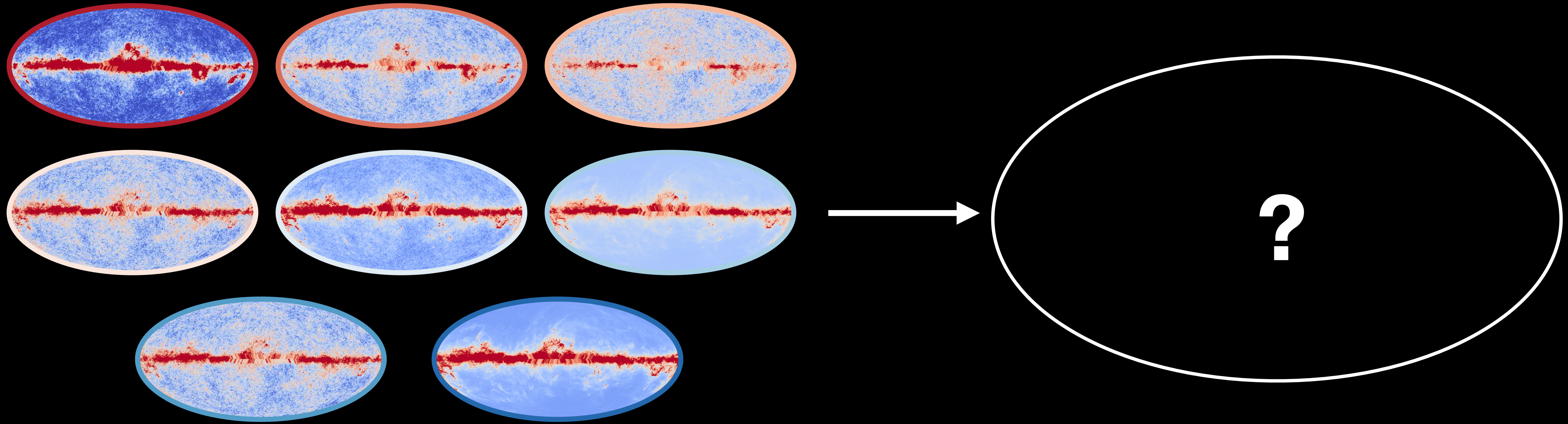
Qualitative illustration of dark screening

Isolating CMB secondaries with multi-frequency data

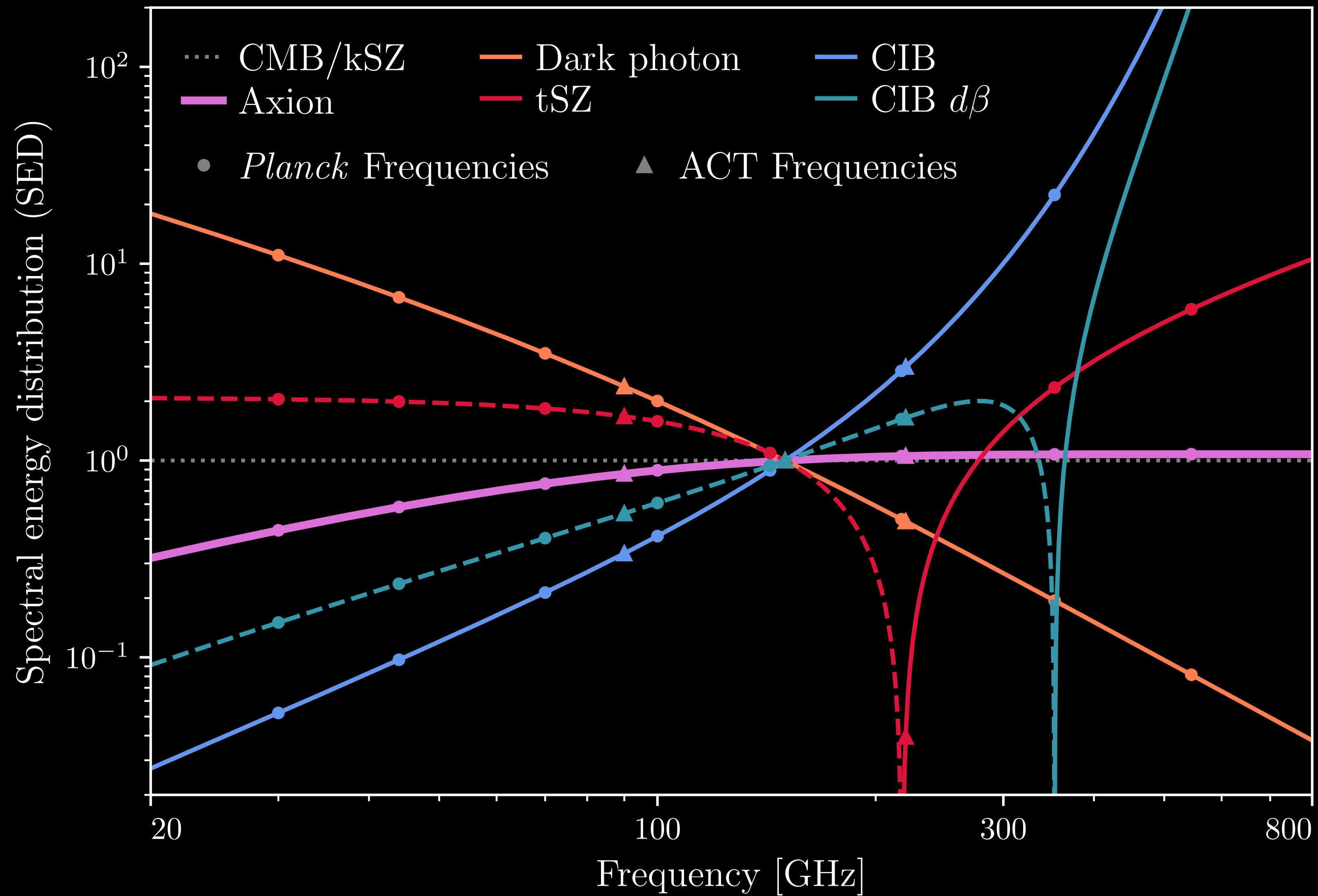


Can we combine multi-frequency “CMB” maps into a single map containing the CMB secondary of interest (axions) and suppressing other contributions?

Isolating CMB secondaries with multi-frequency data

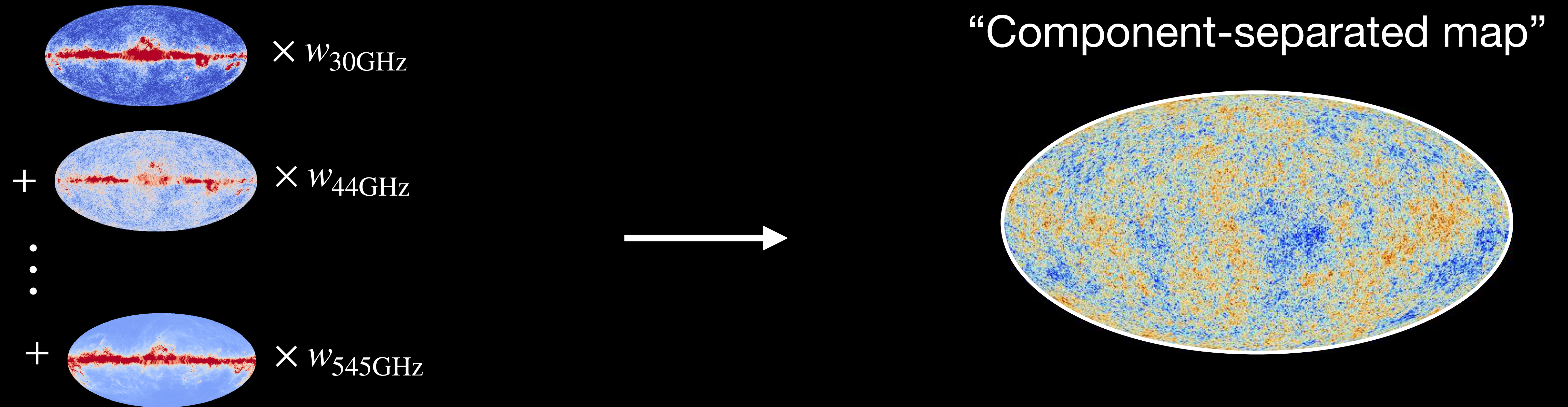


Can we combine multi-frequency “CMB” maps into a single map containing the CMB secondary of interest (axions) and suppressing other contributions?



Yes, because the axion signal has a characteristic frequency dependence!

Internal linear combination (ILC) approach

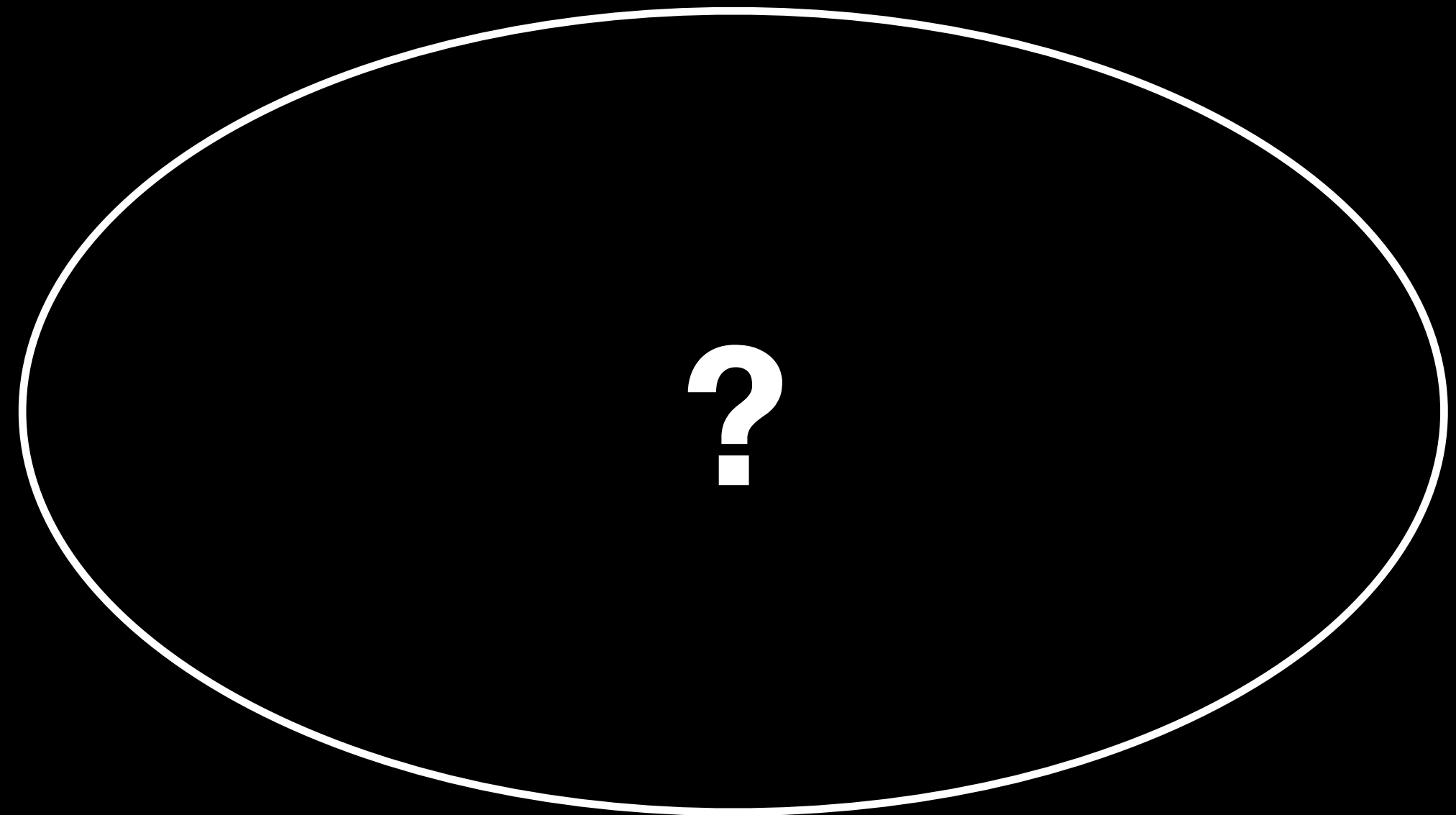
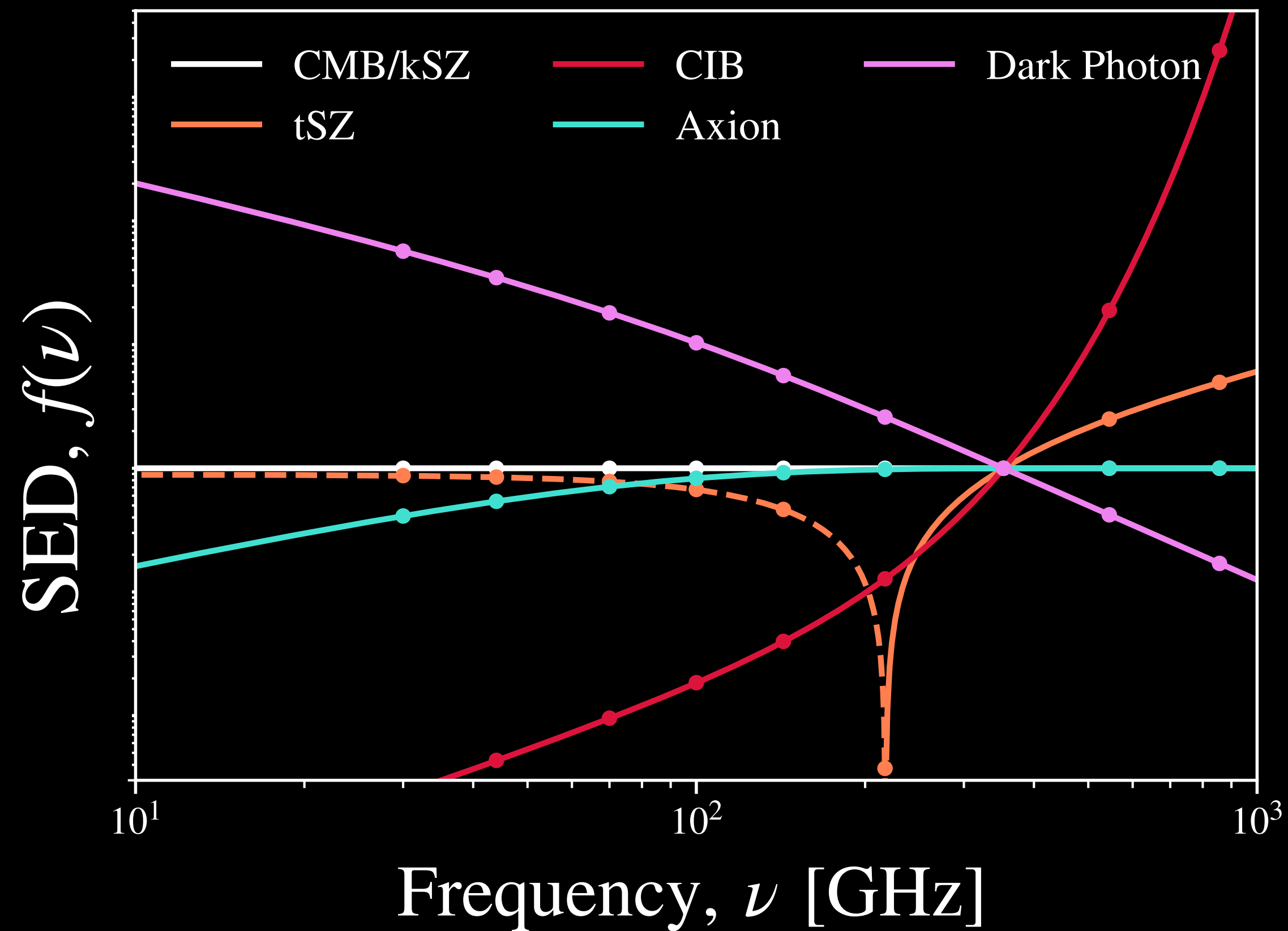


- *Minimum-variance linear combination, subject to unit response to the SED of interest*

$$\text{Minimize } w^T C w \text{ s.t. } w^T a = 1 \implies w = C^{-1} a / (a^T C^{-1} a) \quad \text{\textit{Bennett++03}}$$

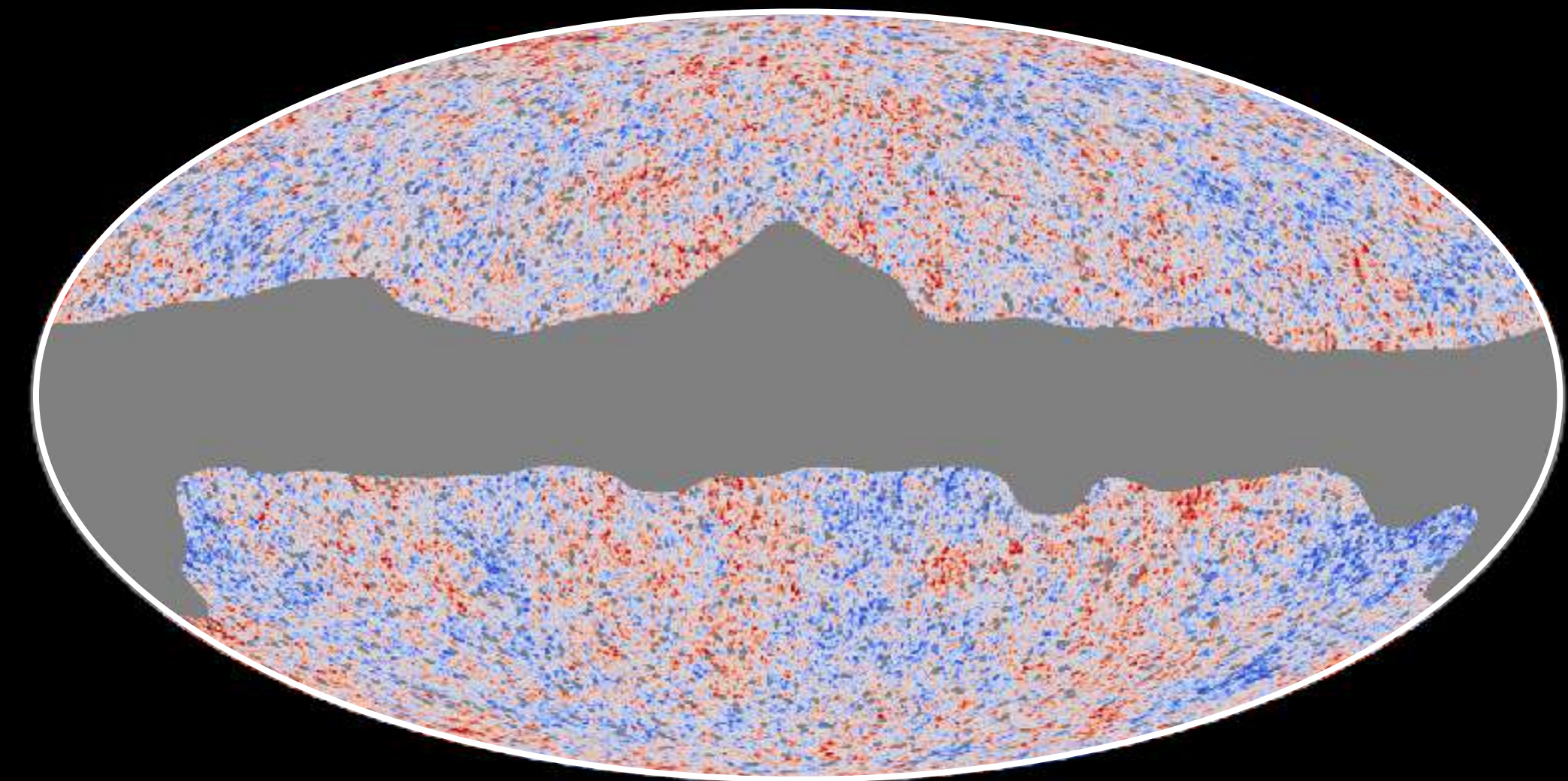
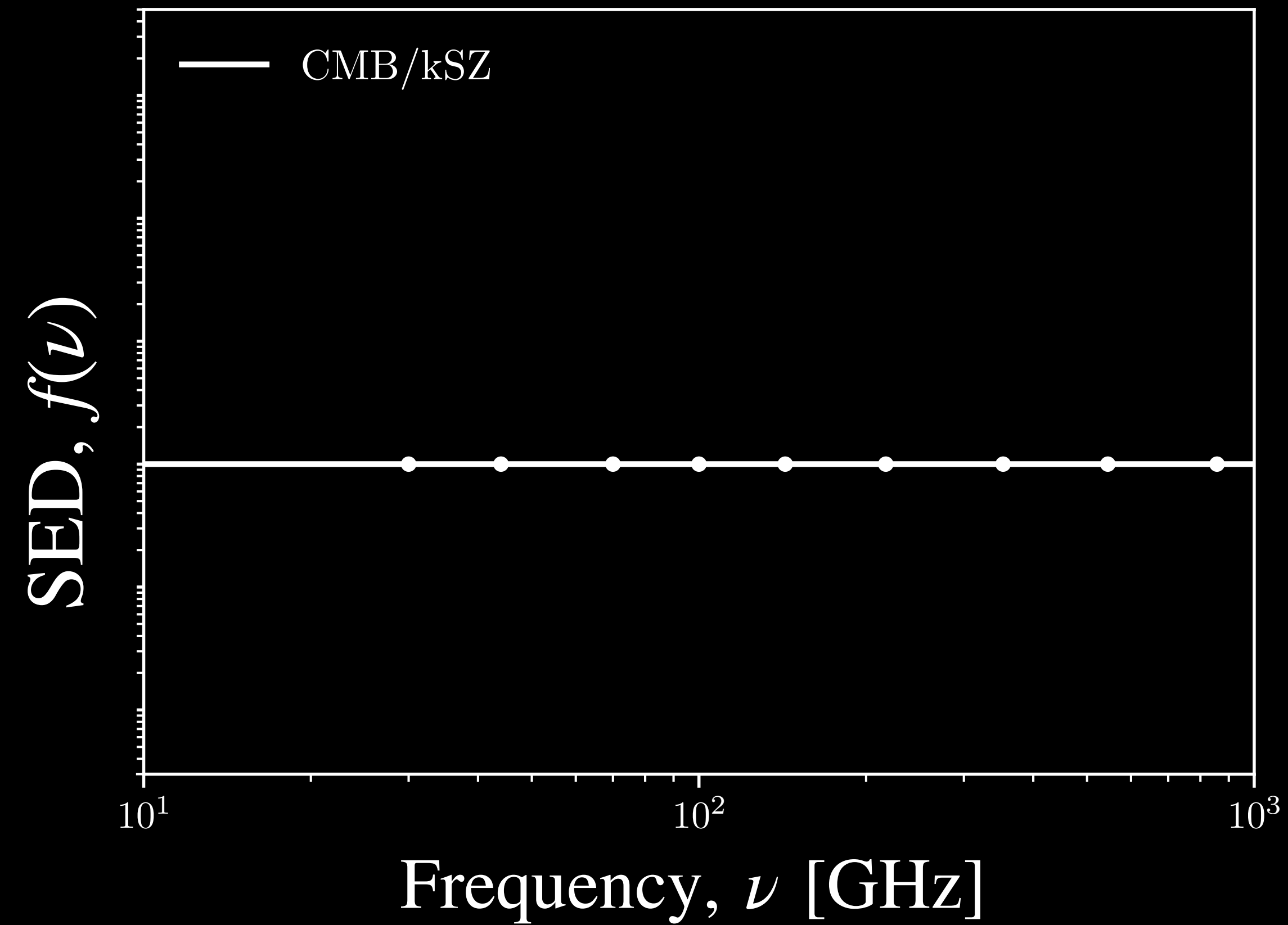
- **Blind:** need only SED you want, not a foreground model
- **Internal:** covariance estimated from data
- **Needlets:** weights vary with position and multipole (Delabrouille++08)

Some example ILC maps



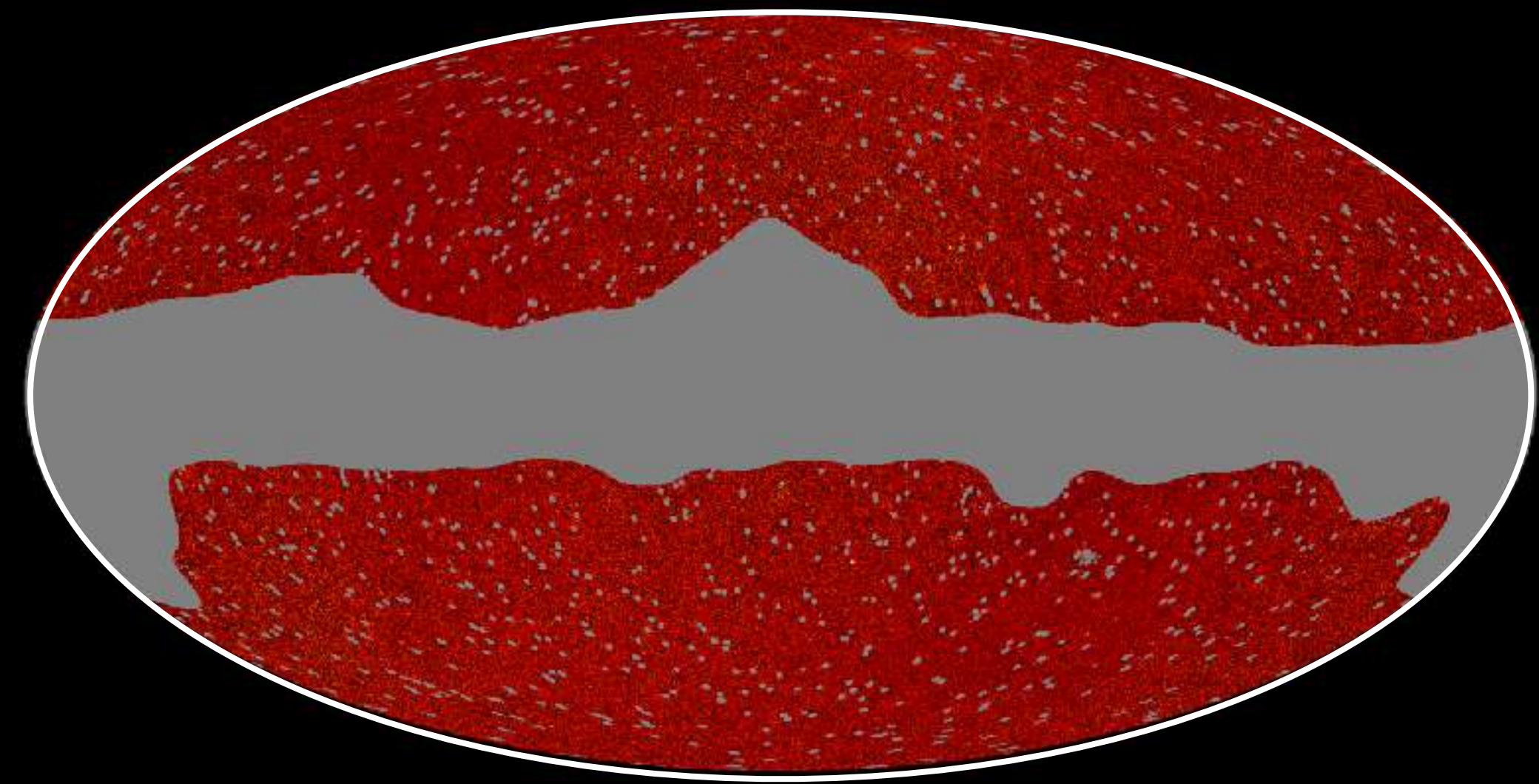
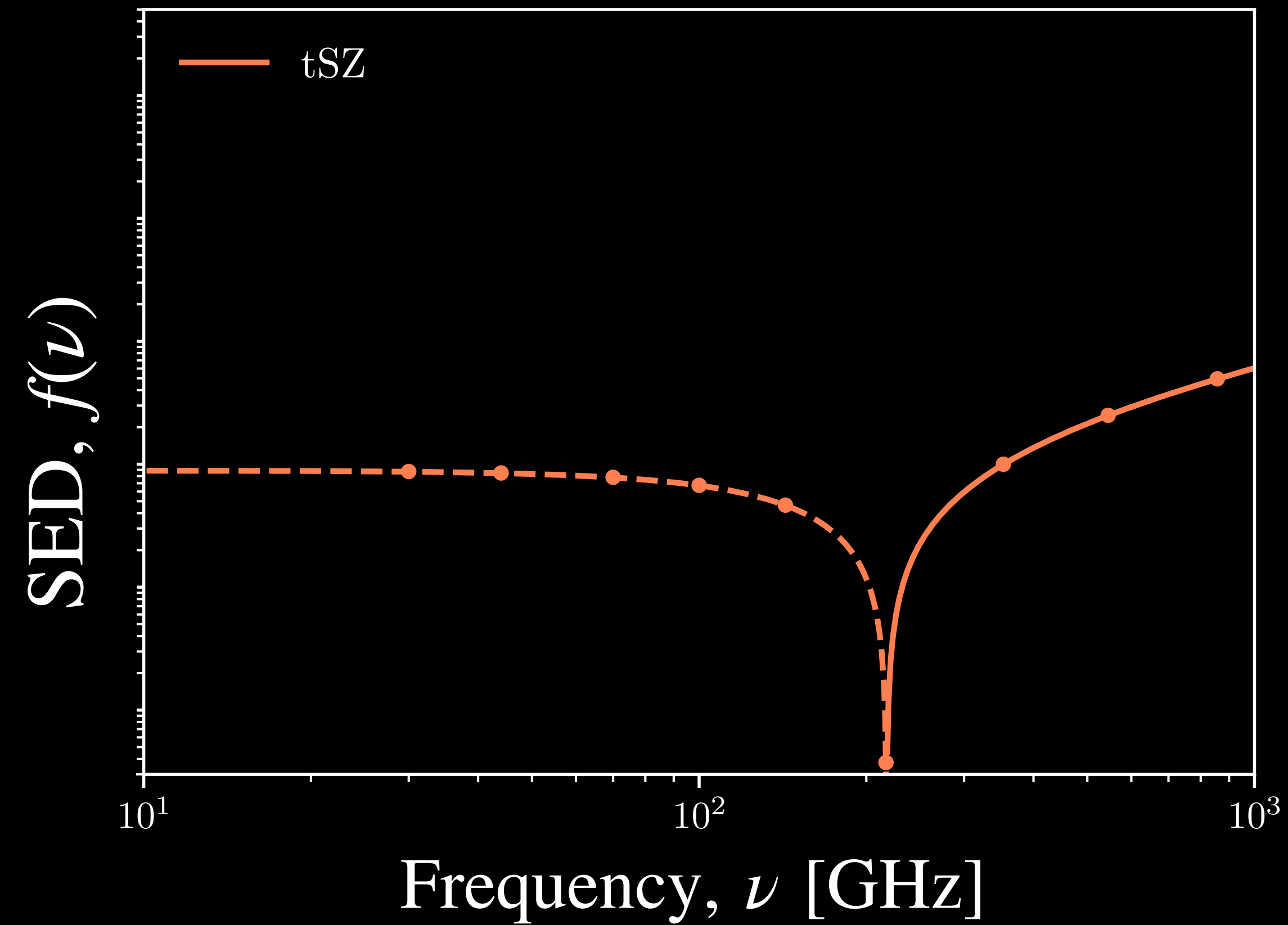
*For the *Planck* analysis, we use *pyilc* implementation ([McCarthy & Hill, 2023](#))

Some example ILC maps



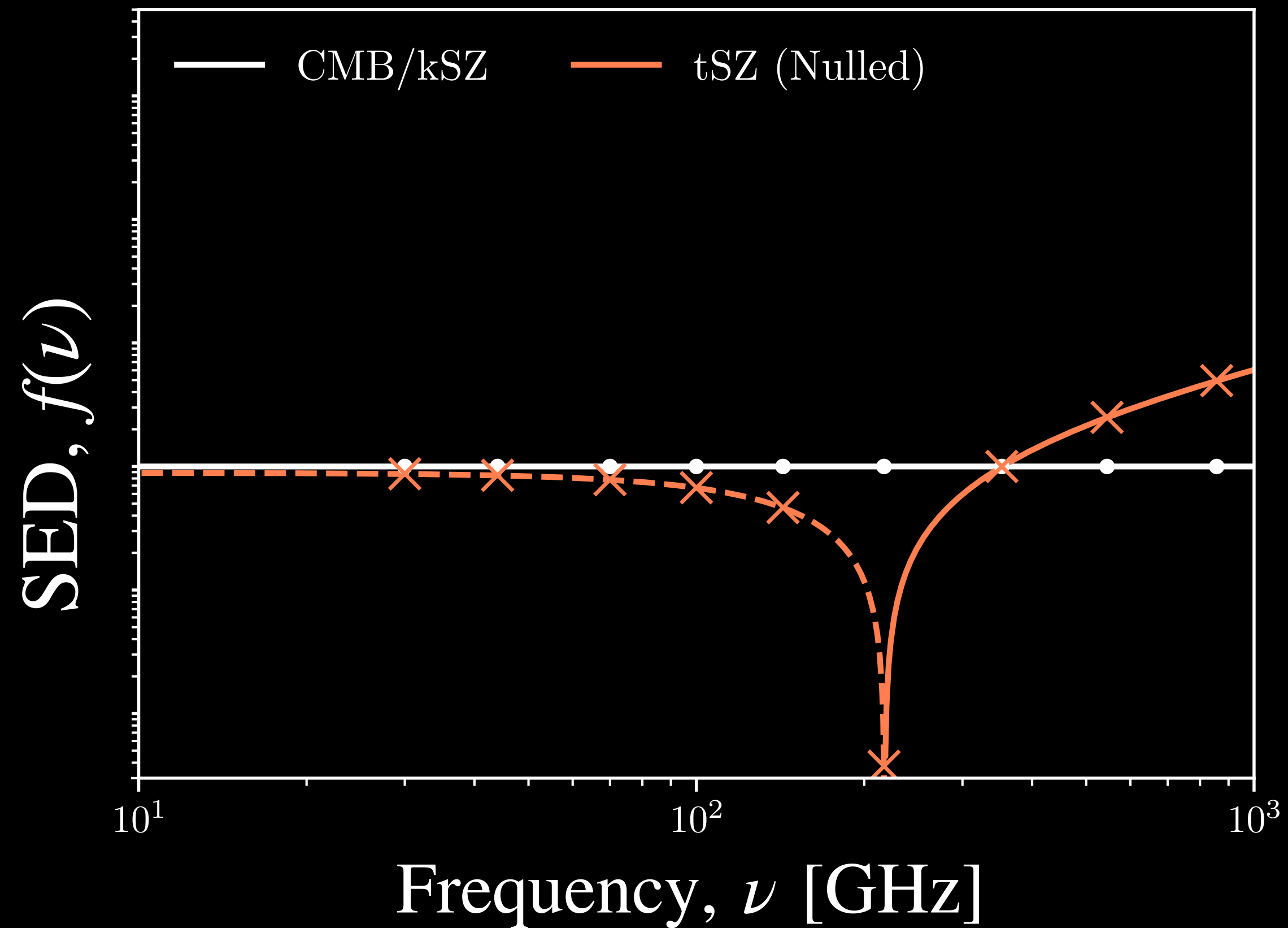
Blackbody T: CMB+kSZ

Some example ILC maps

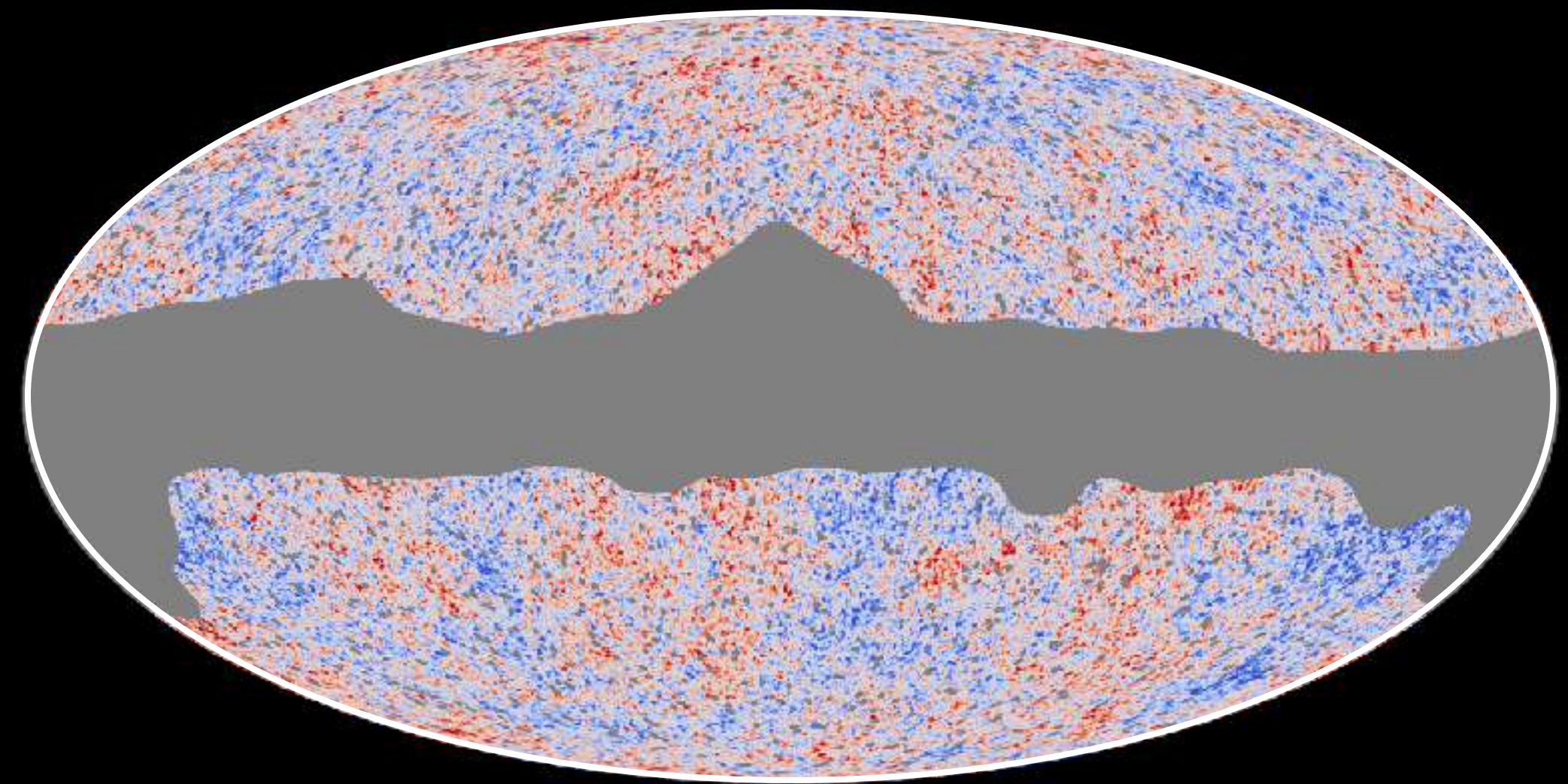


Compton-y (tSZ)

Some example ILC maps



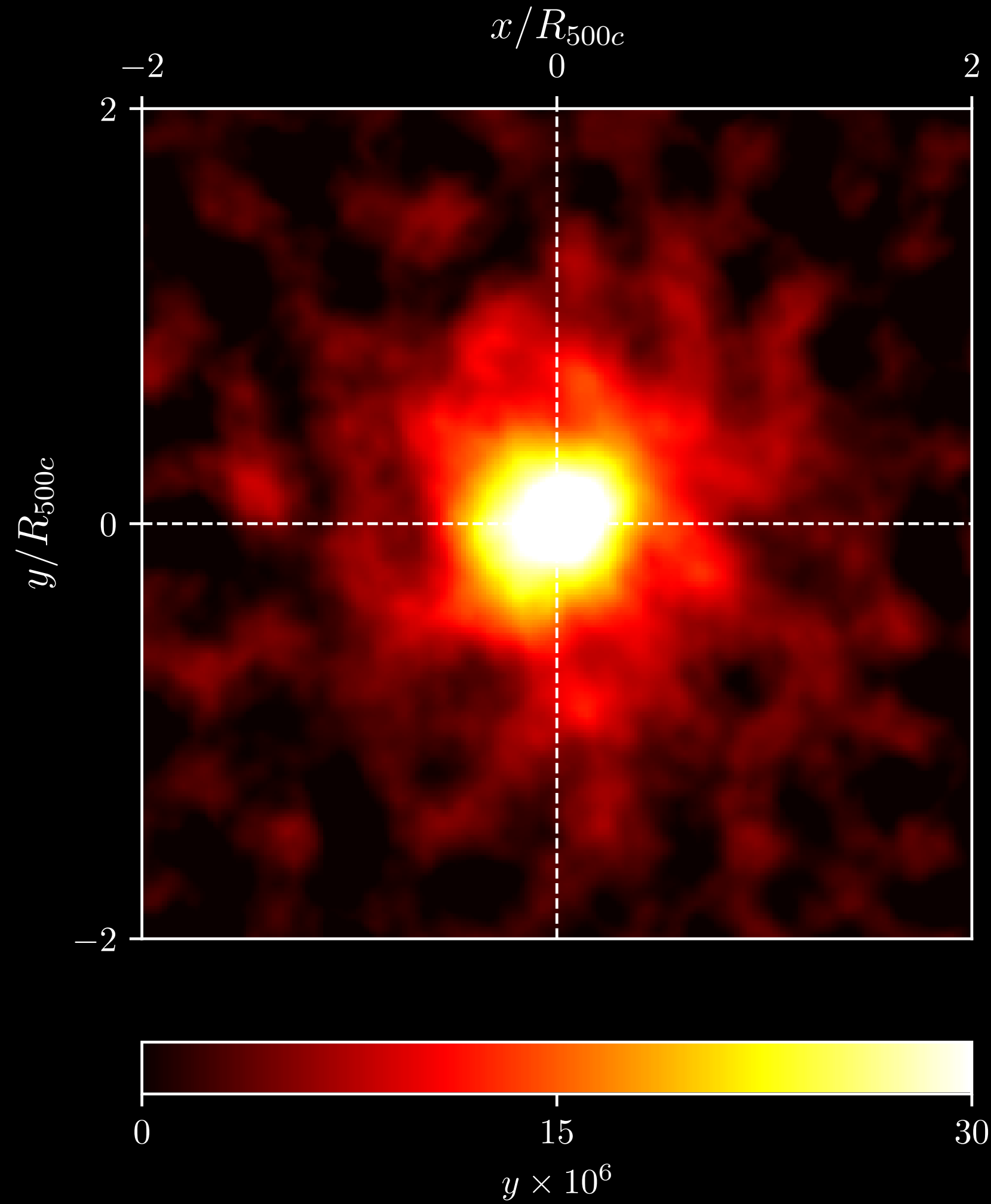
Constrained ILC ([*Remazeilles++10*](#))



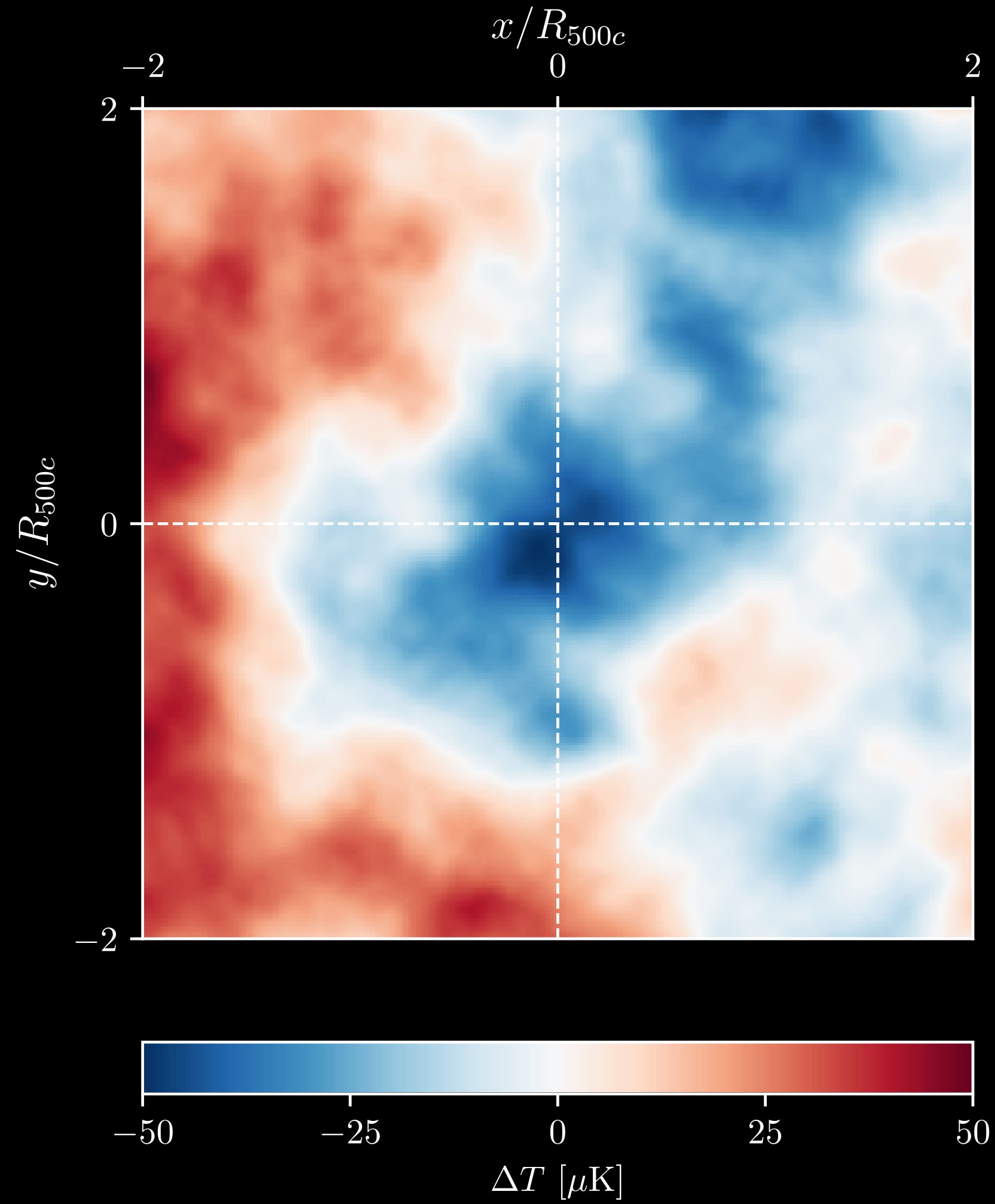
Blackbody T with tSZ deprojected

Stacking NILC maps on 60 X-ray selected clusters

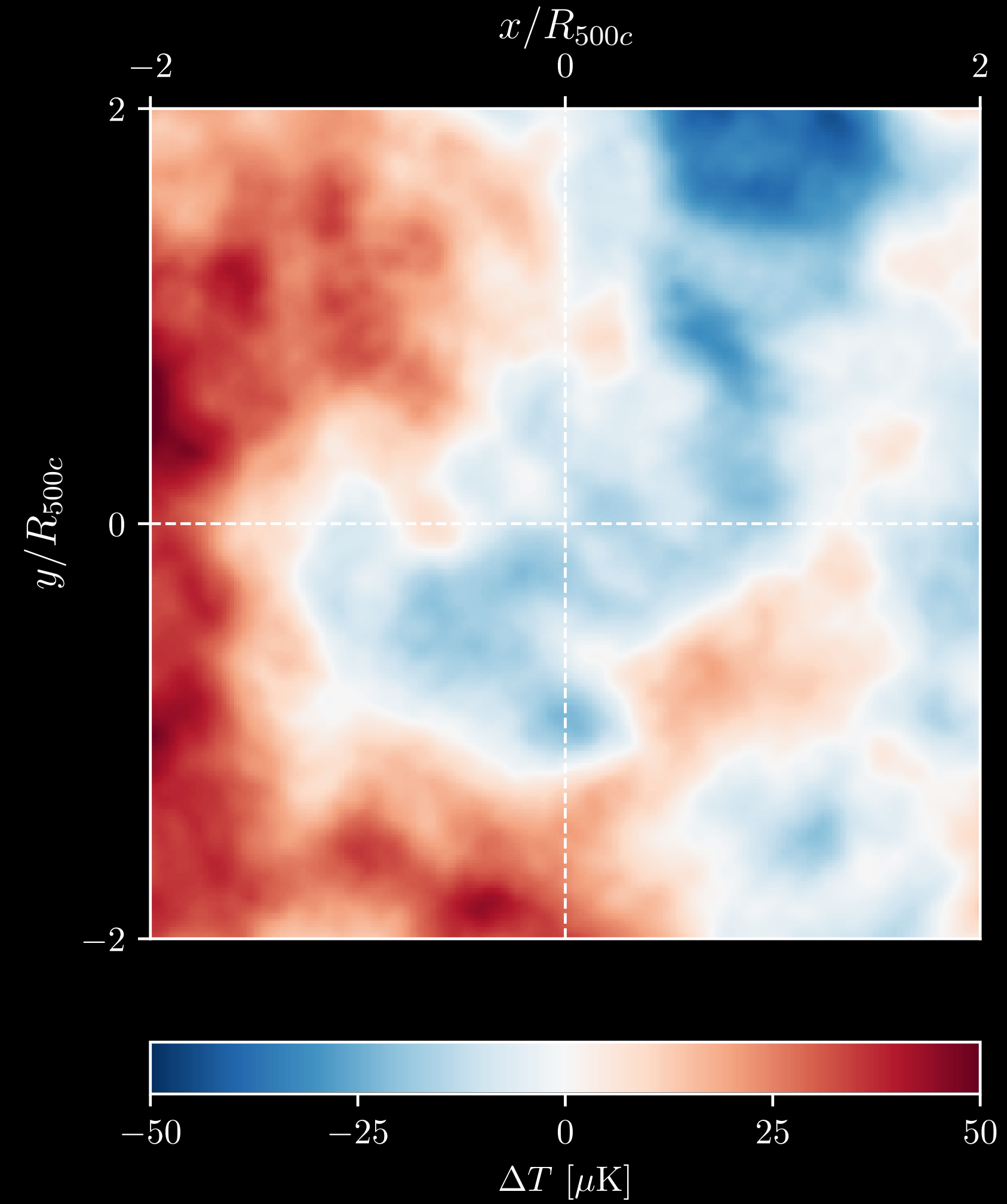
Compton- y NILC Map



CMB NILC Map



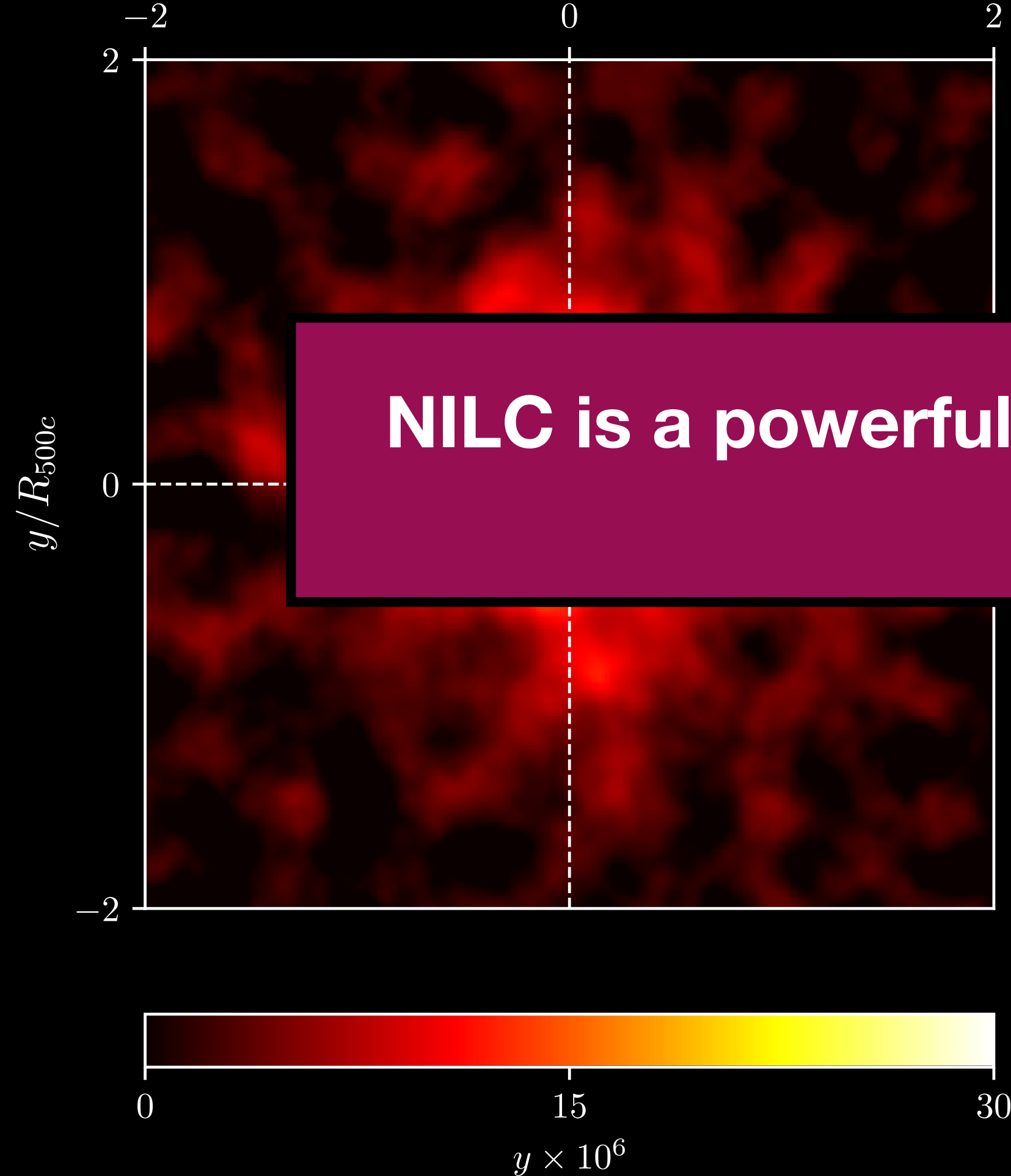
tSZ-Deprojected CMB NILC Map



Stacking NILC maps on 60 X-ray selected clusters

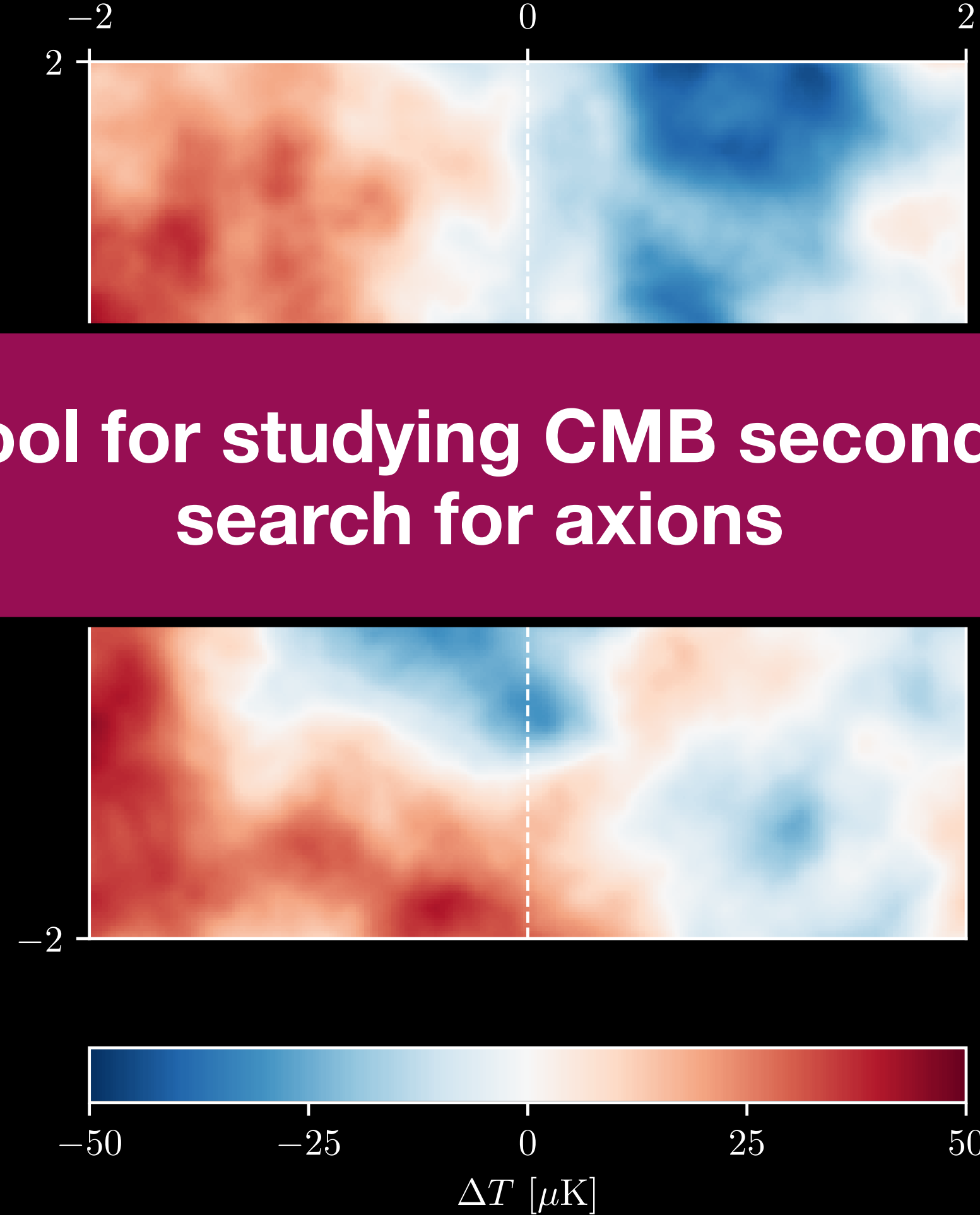
Compton- y NILC Map

x/R_{500c}



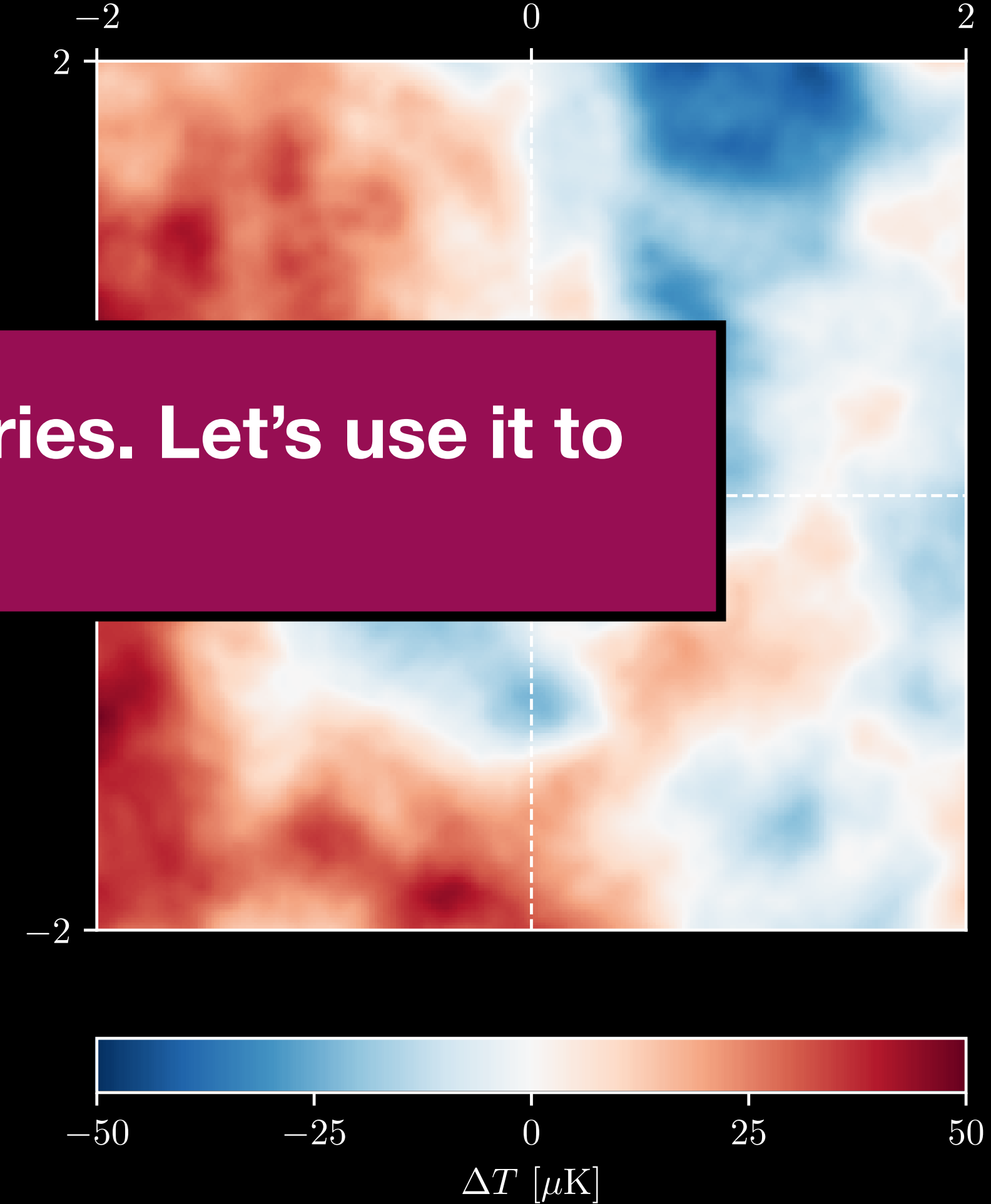
CMB NILC Map

x/R_{500c}



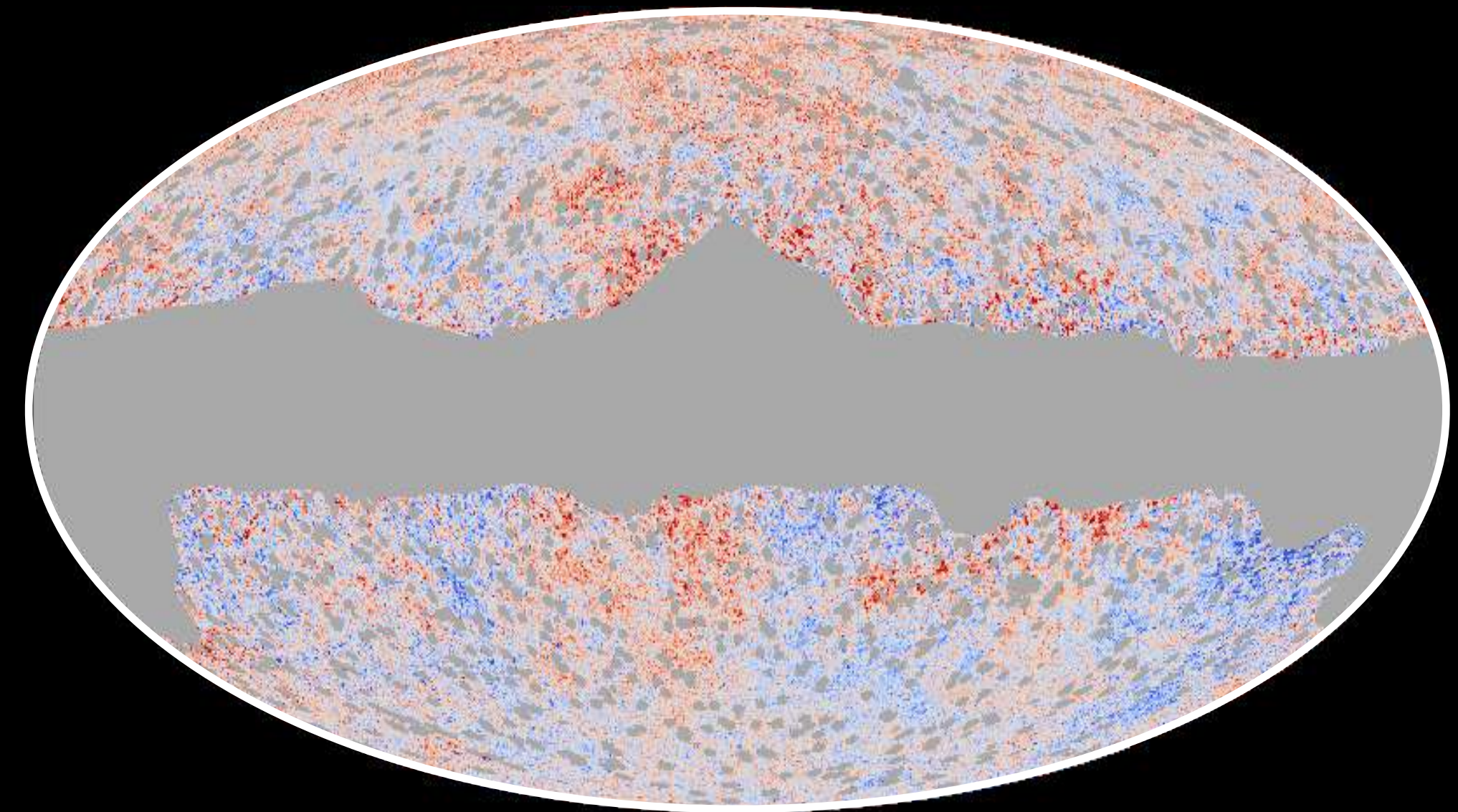
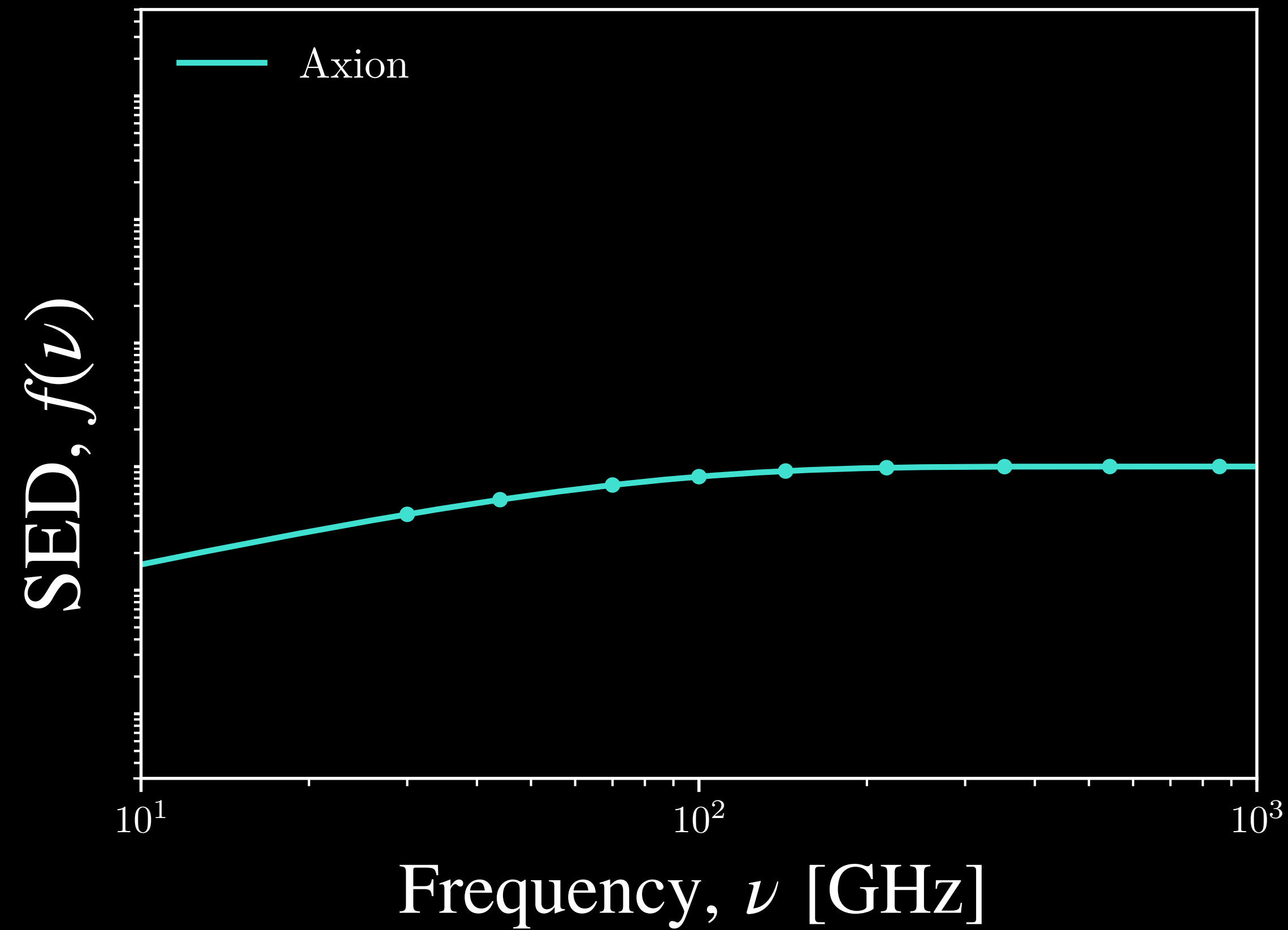
tSZ-Deprojected CMB NILC Map

x/R_{500c}



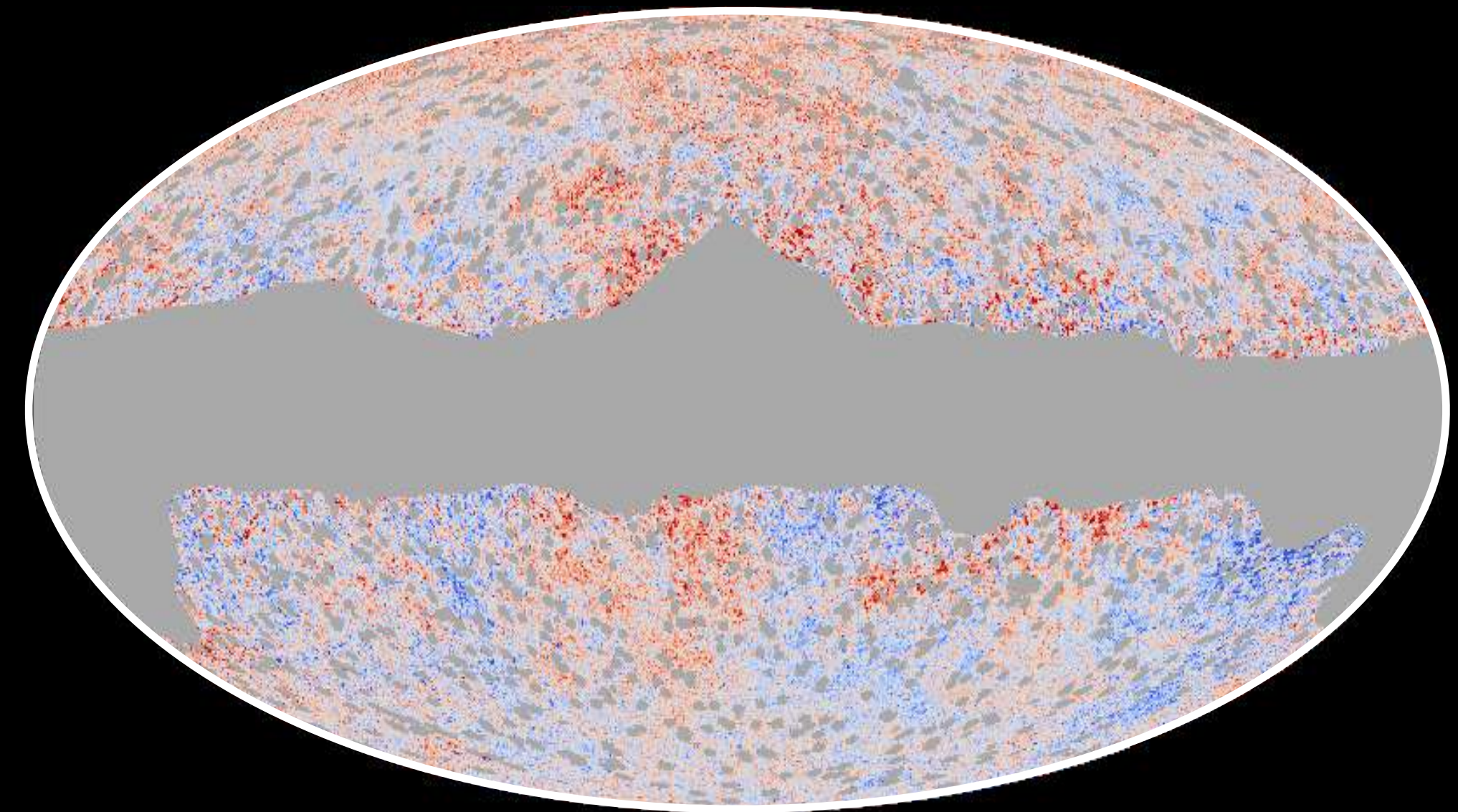
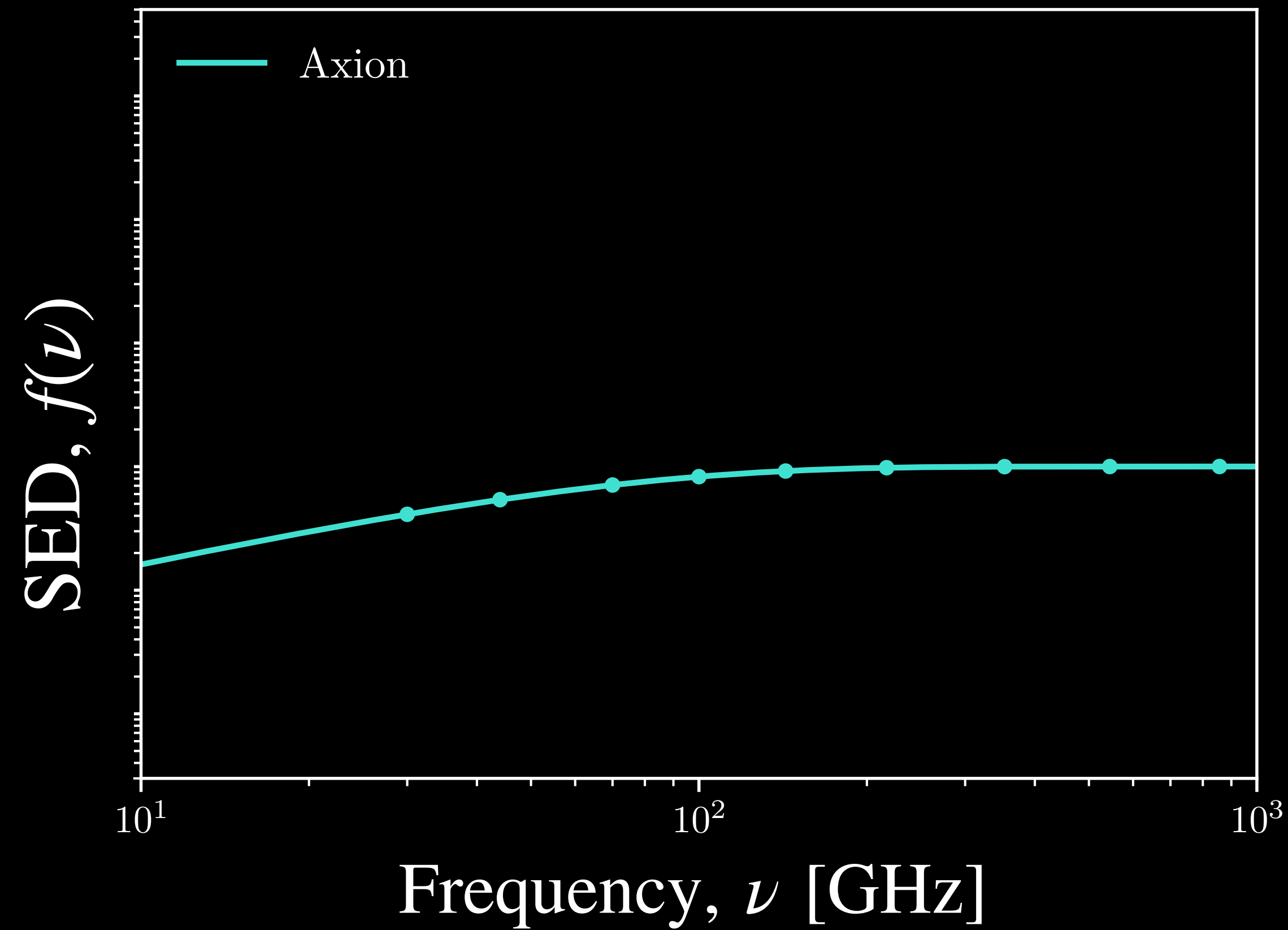
NILC is a powerful tool for studying CMB secondaries. Let's use it to search for axions

Axion-induced screening map



Map = axion signal + “noise”

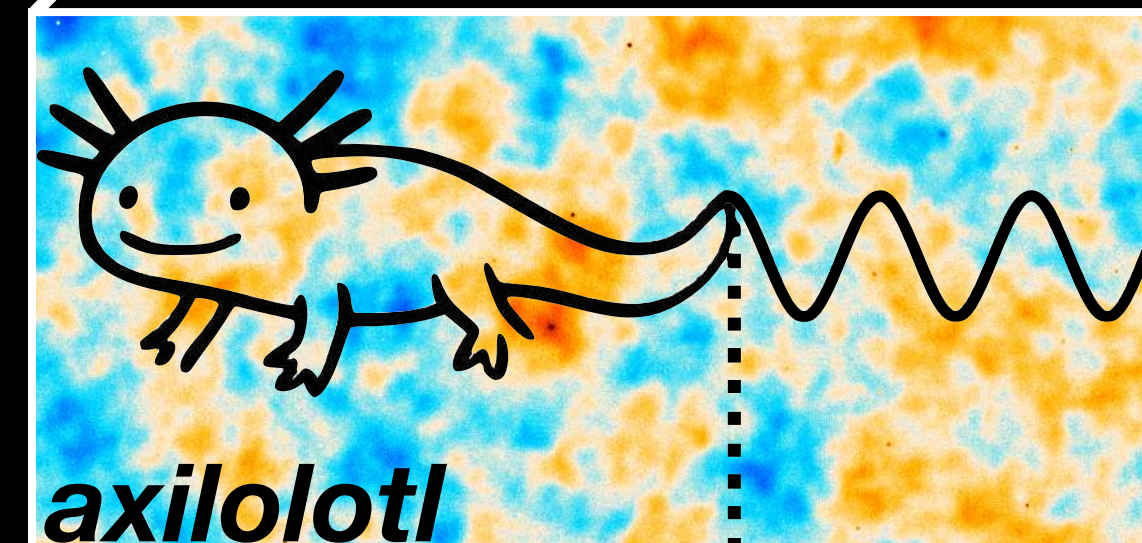
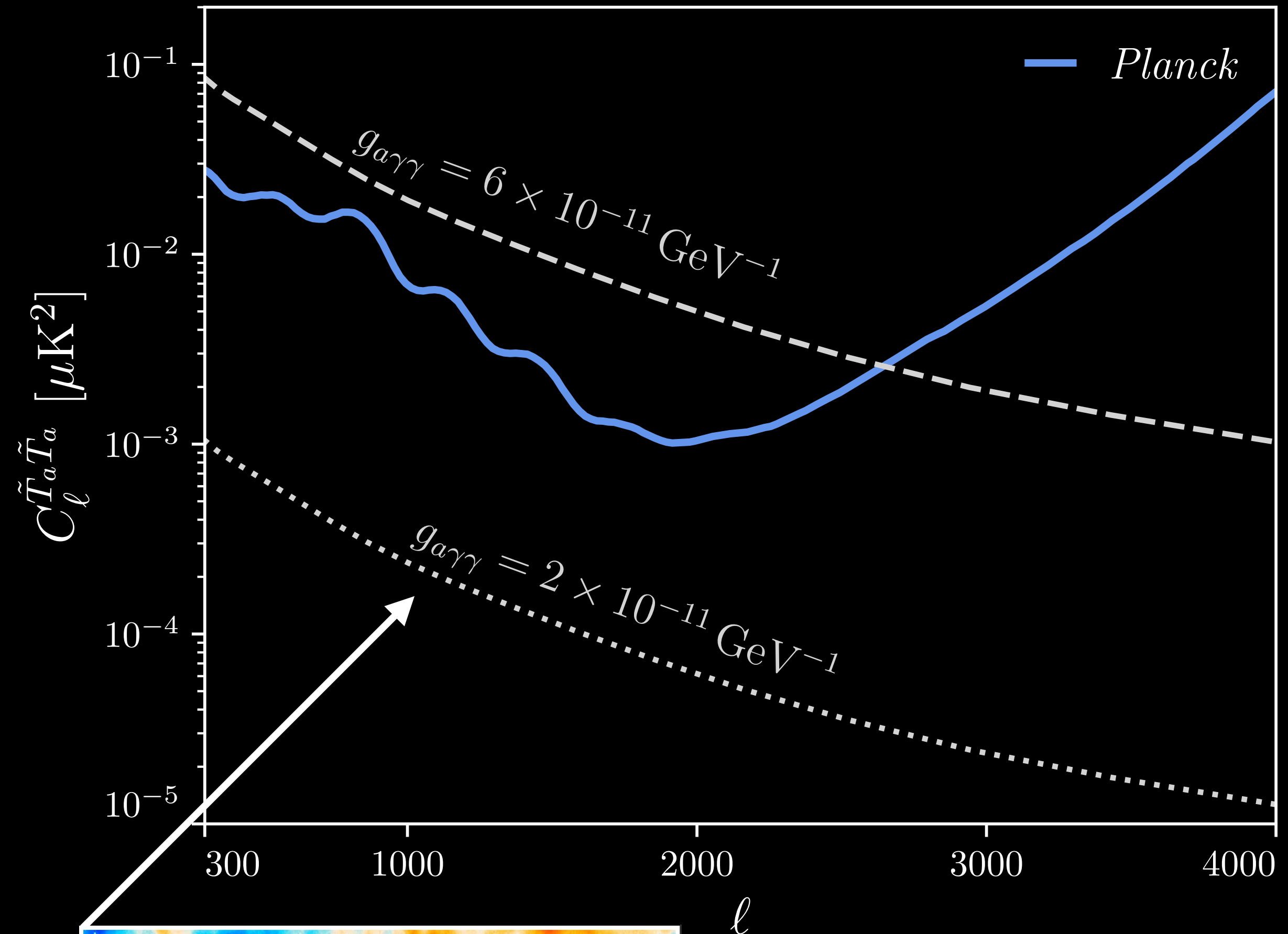
Axion-induced screening map



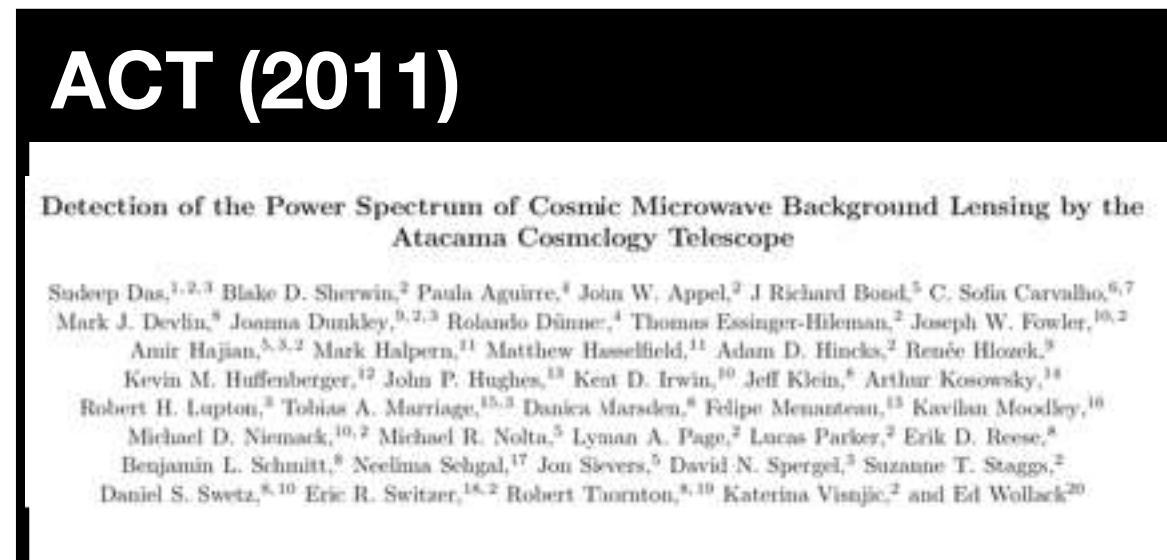
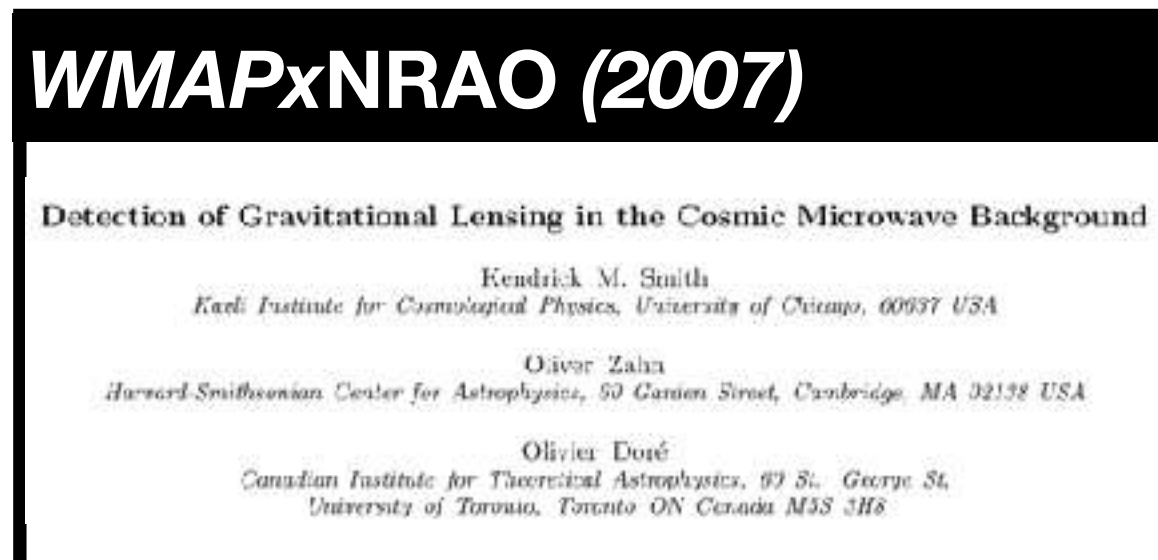
Cool... but what can we do with it?

Option a). Auto-Power Spectrum

- **Extremely** difficult to interpret
 - NILC map = signal+complicated mix of foregrounds and instrumental noise
 - Difficult to make a “detection”
- **Sub-optimal**
 - Signal is correlated with LSS and we have great LSS data
 - Historic example: detection of CMB lensing



Halo model code for dark screening



Intermission: spot the snow leopard (auto-correlation edition)

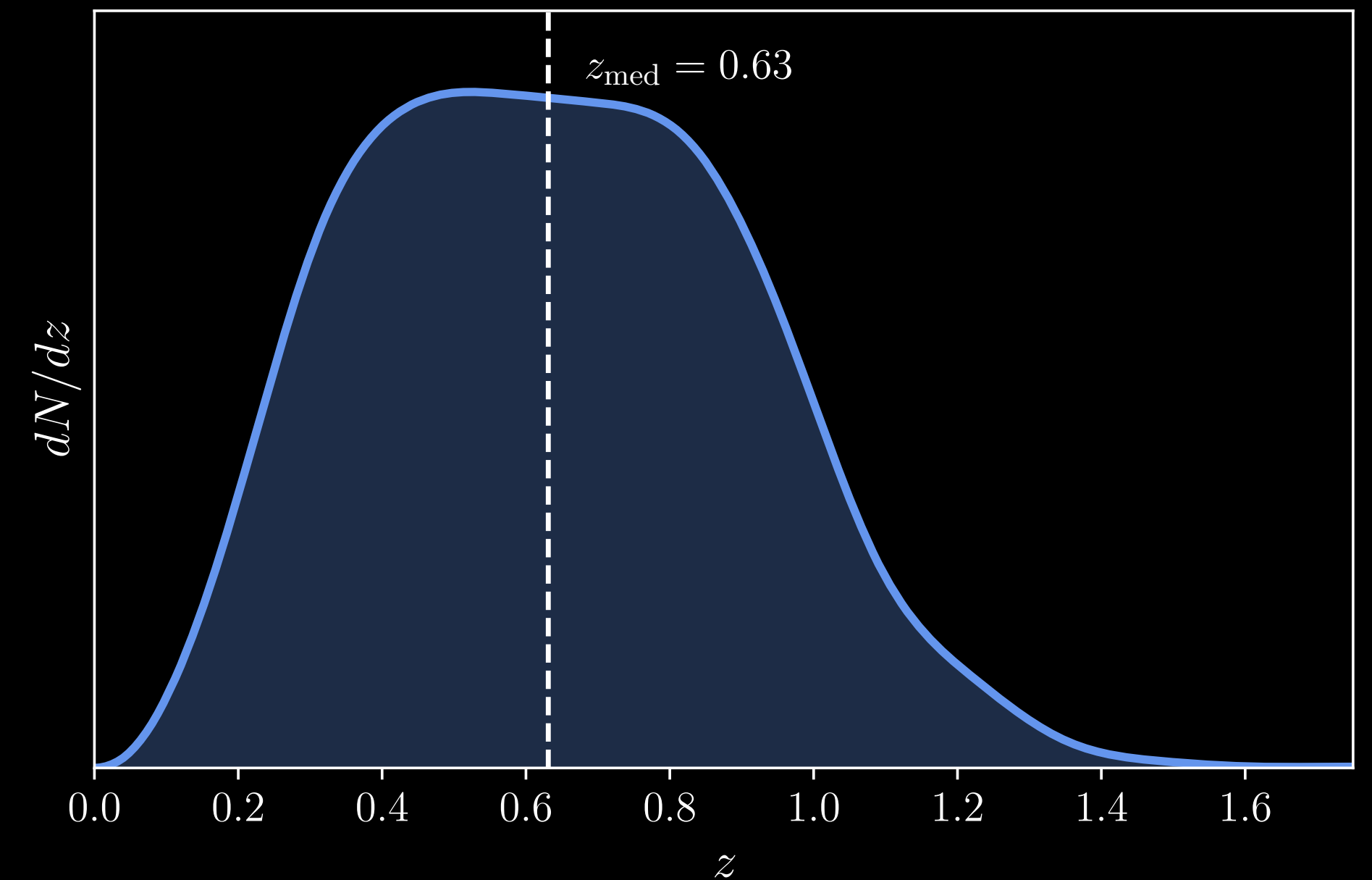


Intermission: spot the snow leopard (cross-correlation edition)



What should we correlate with?

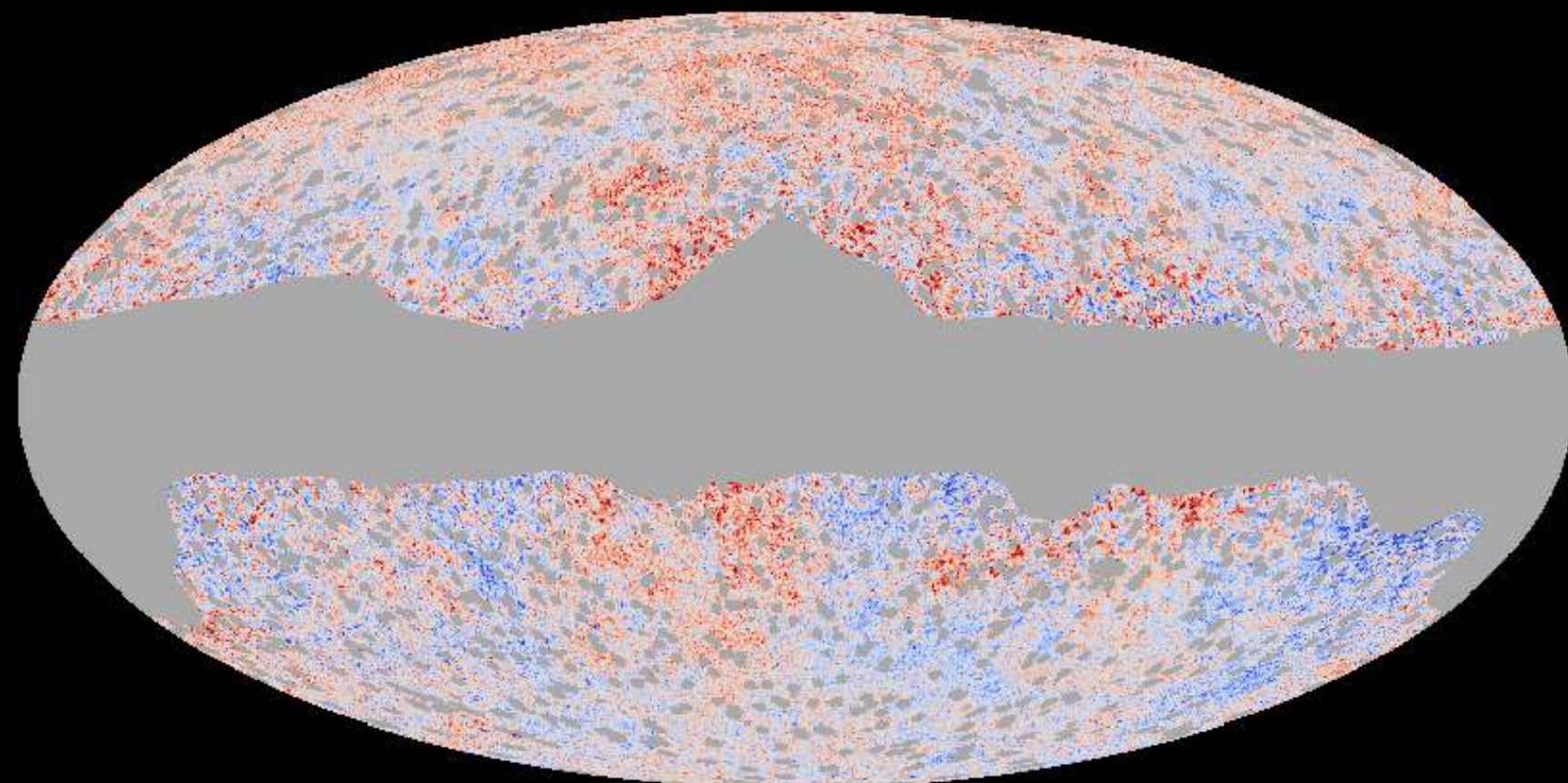
- *unWISE* catalog ([Schlafly++19](#))
 - Based on WISE and NEOWISE
 - Full sky with >500 million galaxies
- Great for cross-correlations
 - CMB lensing/ S_8 ([Krolewski++19](#), [Farren++23](#))
 - Characterized HOD ([Kusiak++22](#))



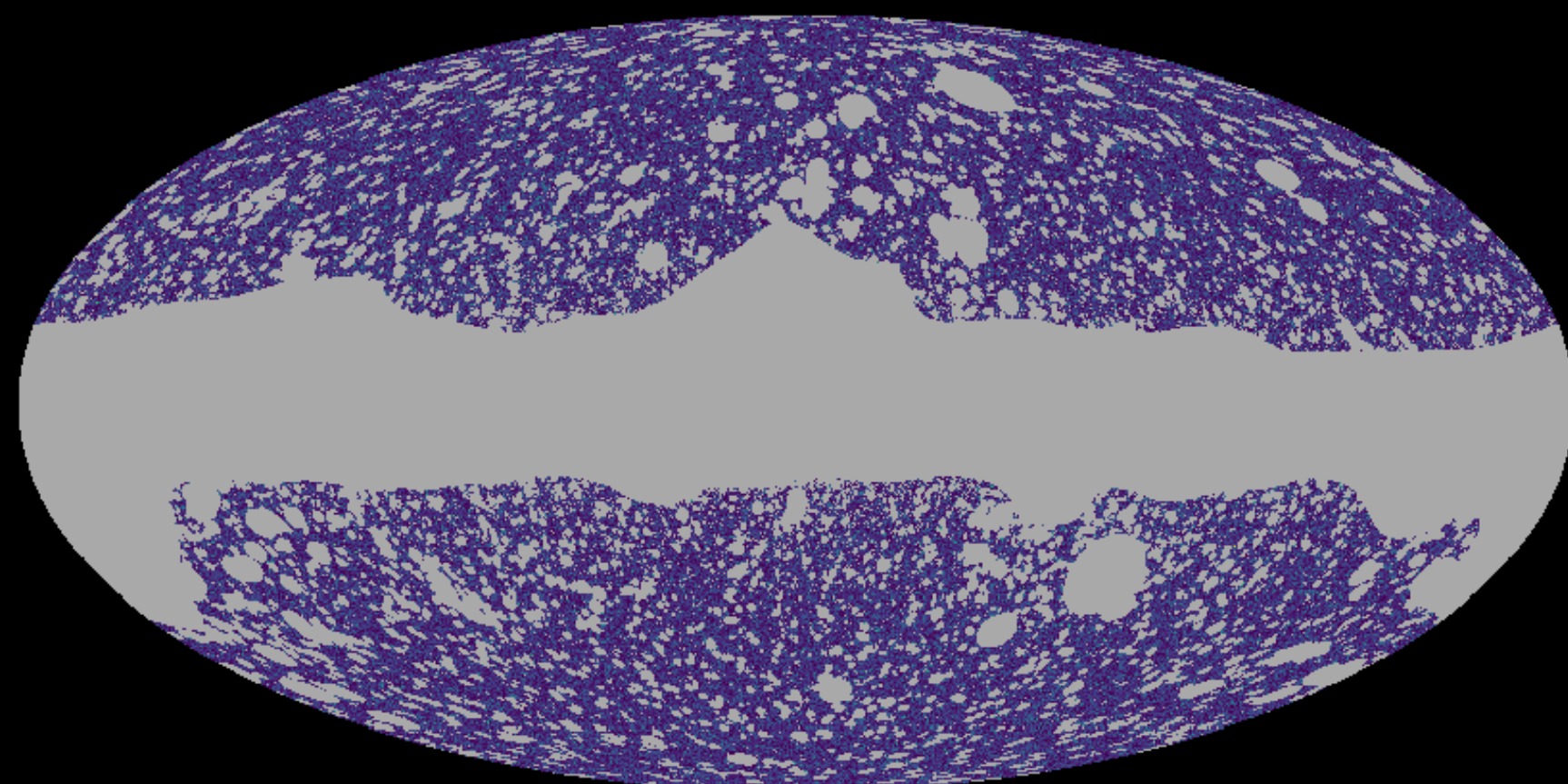
We use the “Blue” subsample

Cross-correlating with unWISE galaxies

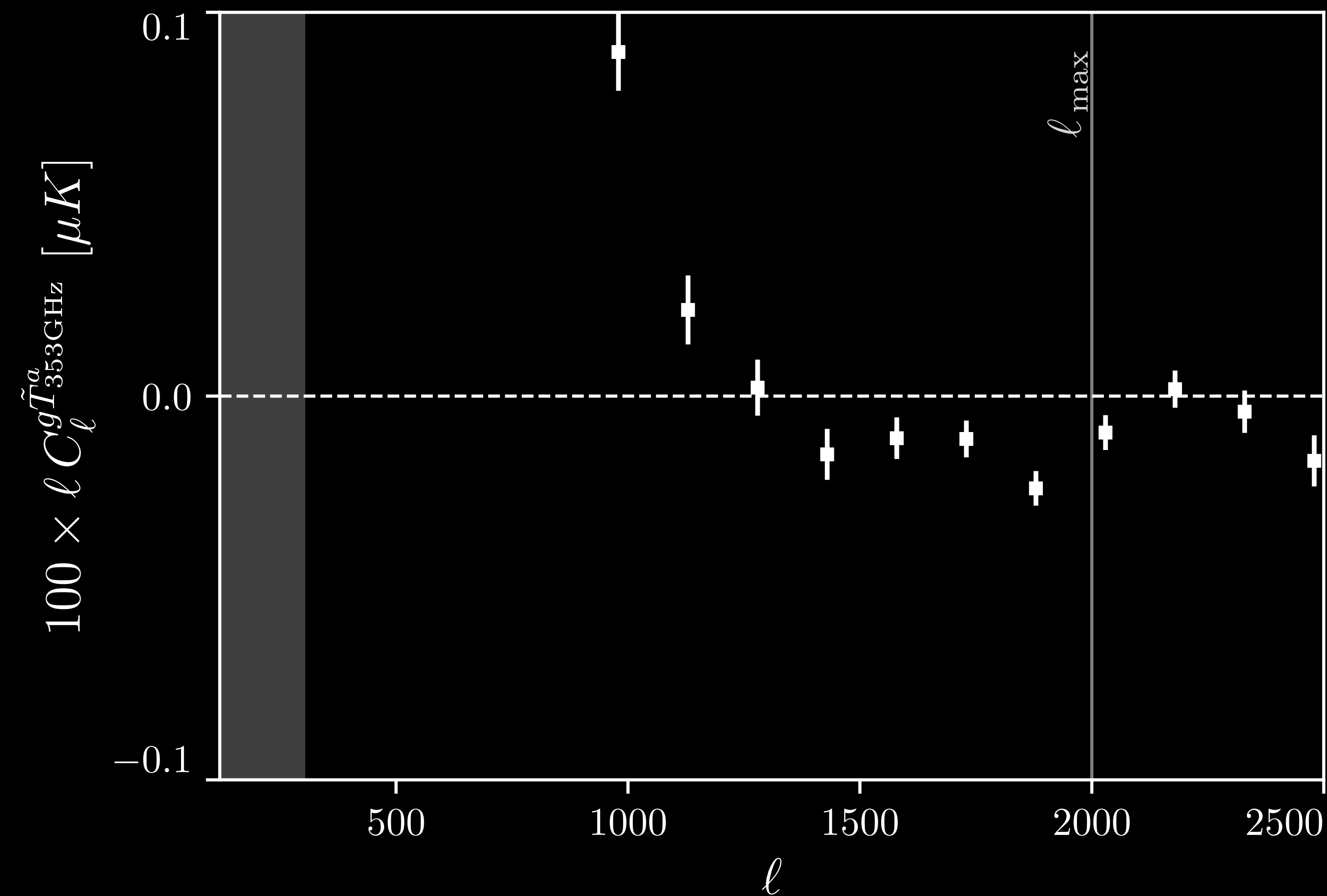
Axion Map



×

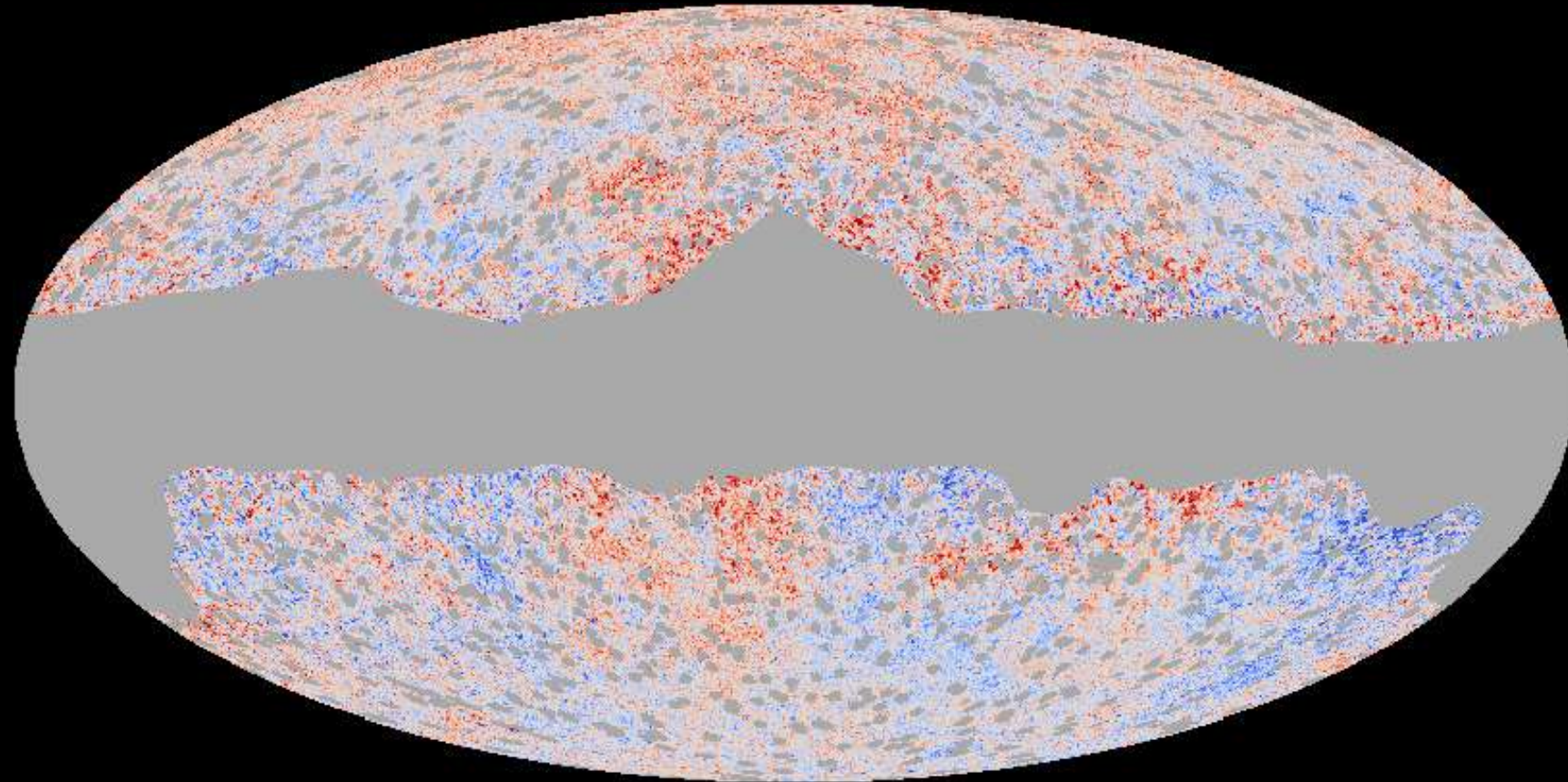


unWISE "Blue" Galaxies

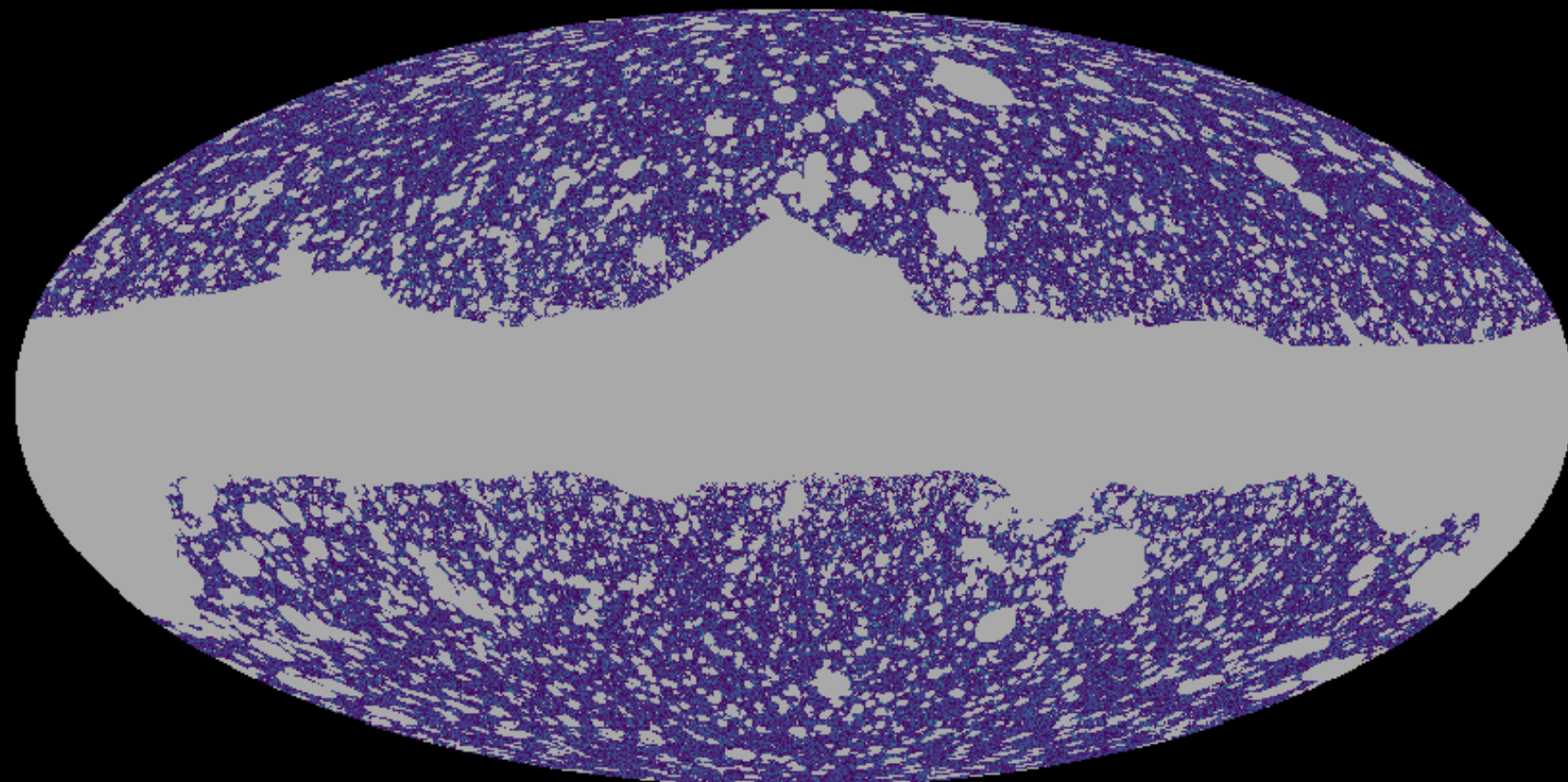


Cross-correlating with unWISE galaxies

Axion Map

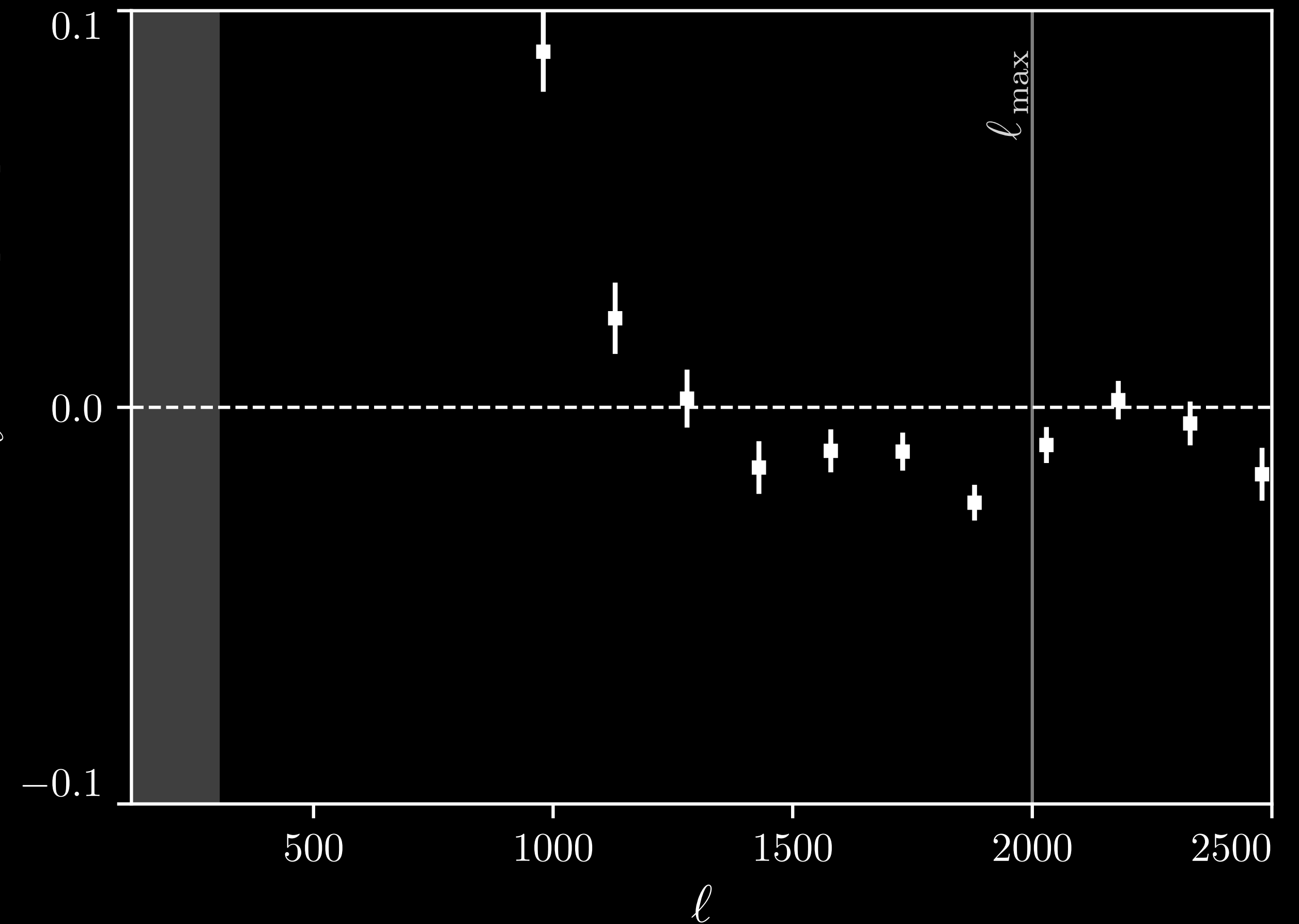


×



unWISE “Blue” Galaxies

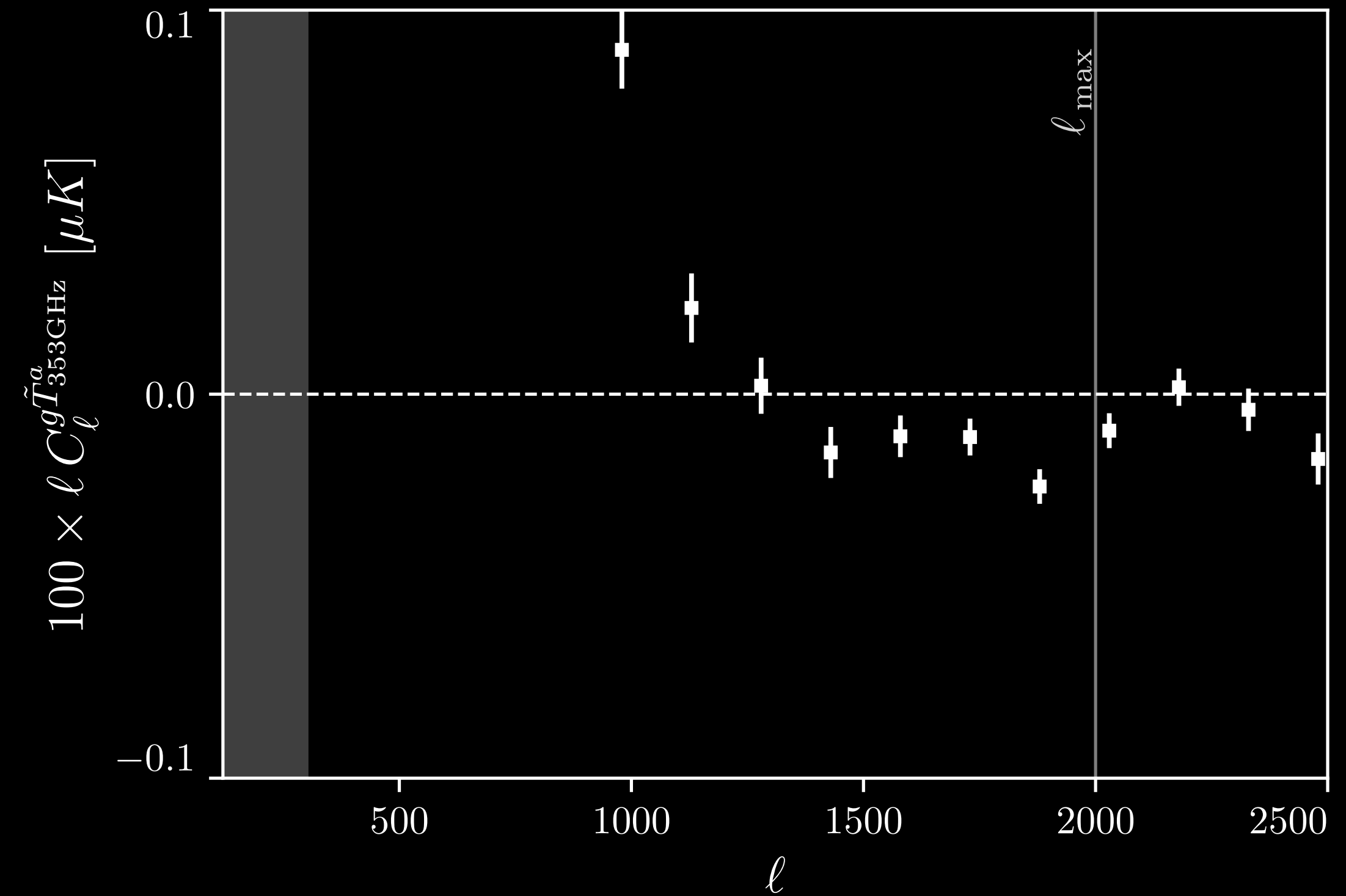
$$100 \times \ell C_\ell^g \tilde{T}_{353\text{GHz}}^a [\mu K]$$



Non-zero -> did we detect axions?

Foreground contamination

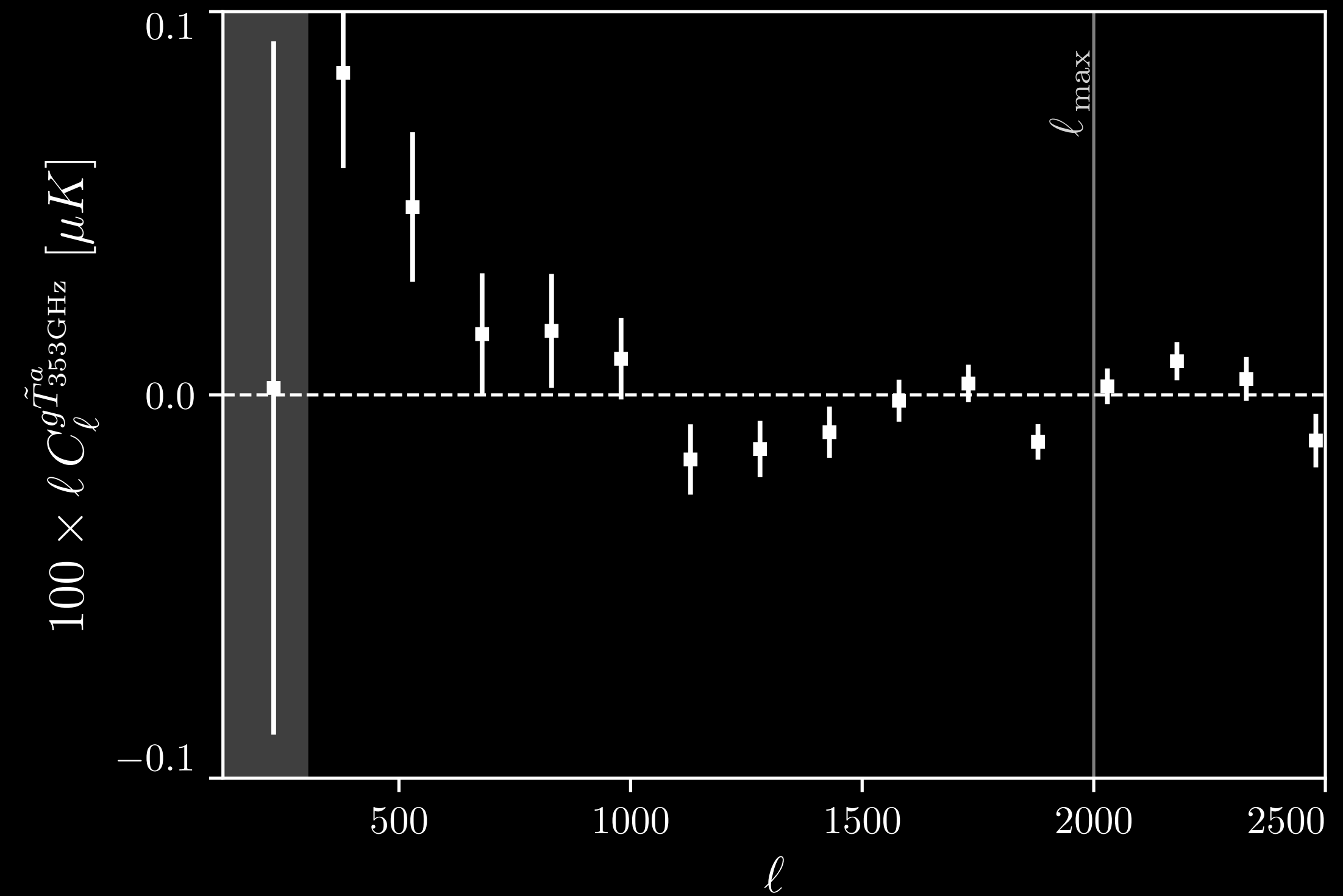
- Measurement is sensitive to any residual foregrounds that correlate with unWISE galaxies



Without foreground deprojection.

Foreground contamination

- Measurement is sensitive to any residual foregrounds that correlate with unWISE galaxies
 1. **tSZ**: significant correlation, but SED is known

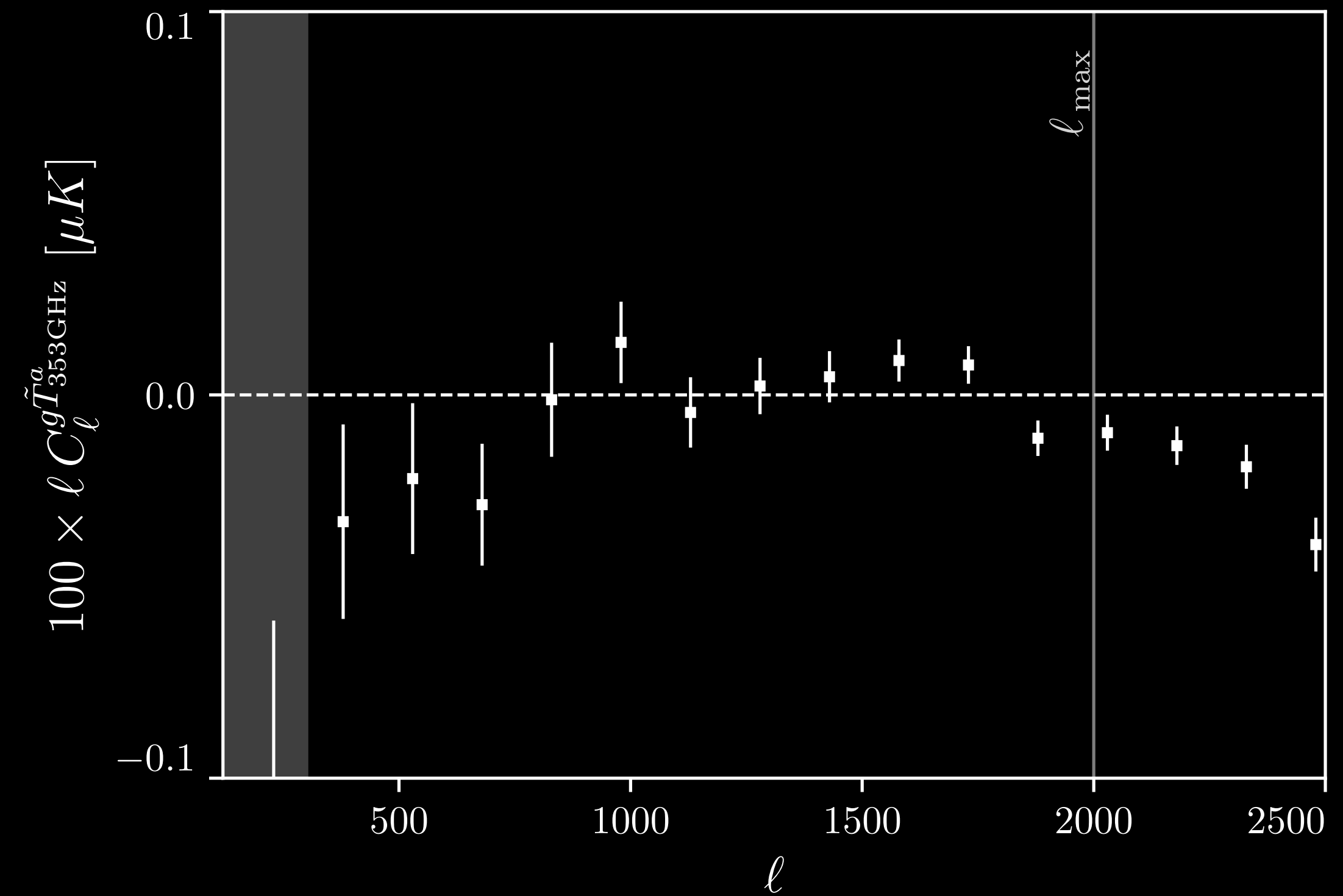


Deproject tSZ

Foreground contamination

- Measurement is sensitive to any residual foregrounds that correlate with unWISE galaxies
 1. **tSZ**: significant correlation, but SED is known
 2. **CIB**: significant correlation ([Yan++2023](#)), uncertainty in SED
- Model as modified blackbody (MBB)

$$\Theta_{\text{CIB}}(\nu | T_{\text{CIB}}^{\text{eff}}, \beta_{\text{CIB}}^{\text{eff}}) = \nu^{\beta_{\text{CIB}}^{\text{eff}}} B(\nu | T_{\text{CIB}}^{\text{eff}})$$



Deproject tSZ+CIB MBB

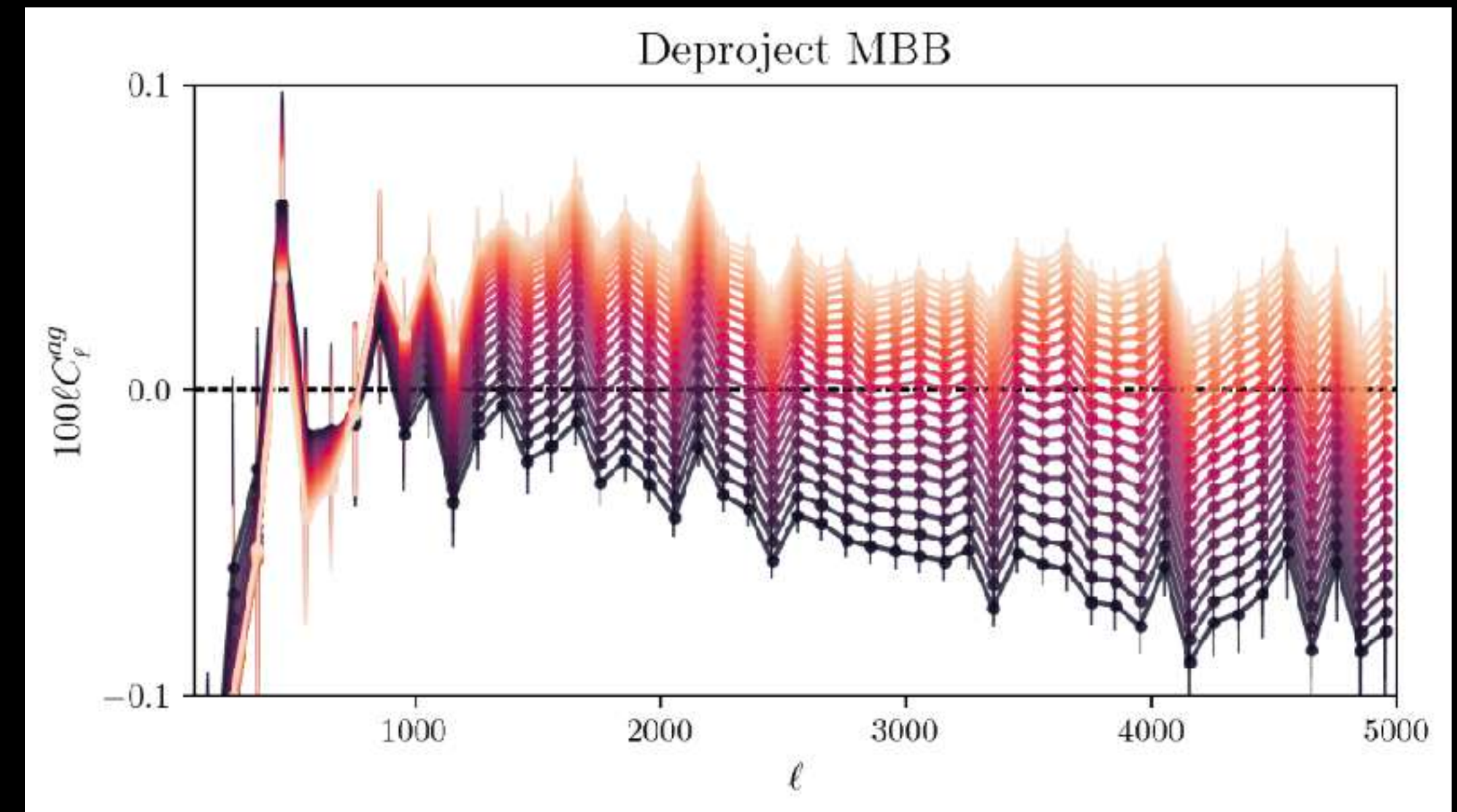
$$(T_{\text{CIB}}^{\text{eff}}, \beta_{\text{CIB}}^{\text{eff}}) = (20.0 \text{ K}, 1.6)$$

Foreground contamination

- Measurement is sensitive to any residual foregrounds that correlate with unWISE galaxies
 1. **tSZ**: significant correlation, but SED is known
 2. **CIB**: significant correlation ([Yan++2023](#)), uncertainty in SED
- Model as modified blackbody (MBB)

$$\Theta_{\text{CIB}}(\nu | T_{\text{CIB}}^{\text{eff}}, \beta_{\text{CIB}}^{\text{eff}}) = \nu^{\beta_{\text{CIB}}^{\text{eff}}} B(\nu | T_{\text{CIB}}^{\text{eff}})$$

Very sensitive to CIB SED parameters!



Deproject tSZ+CIB MBB

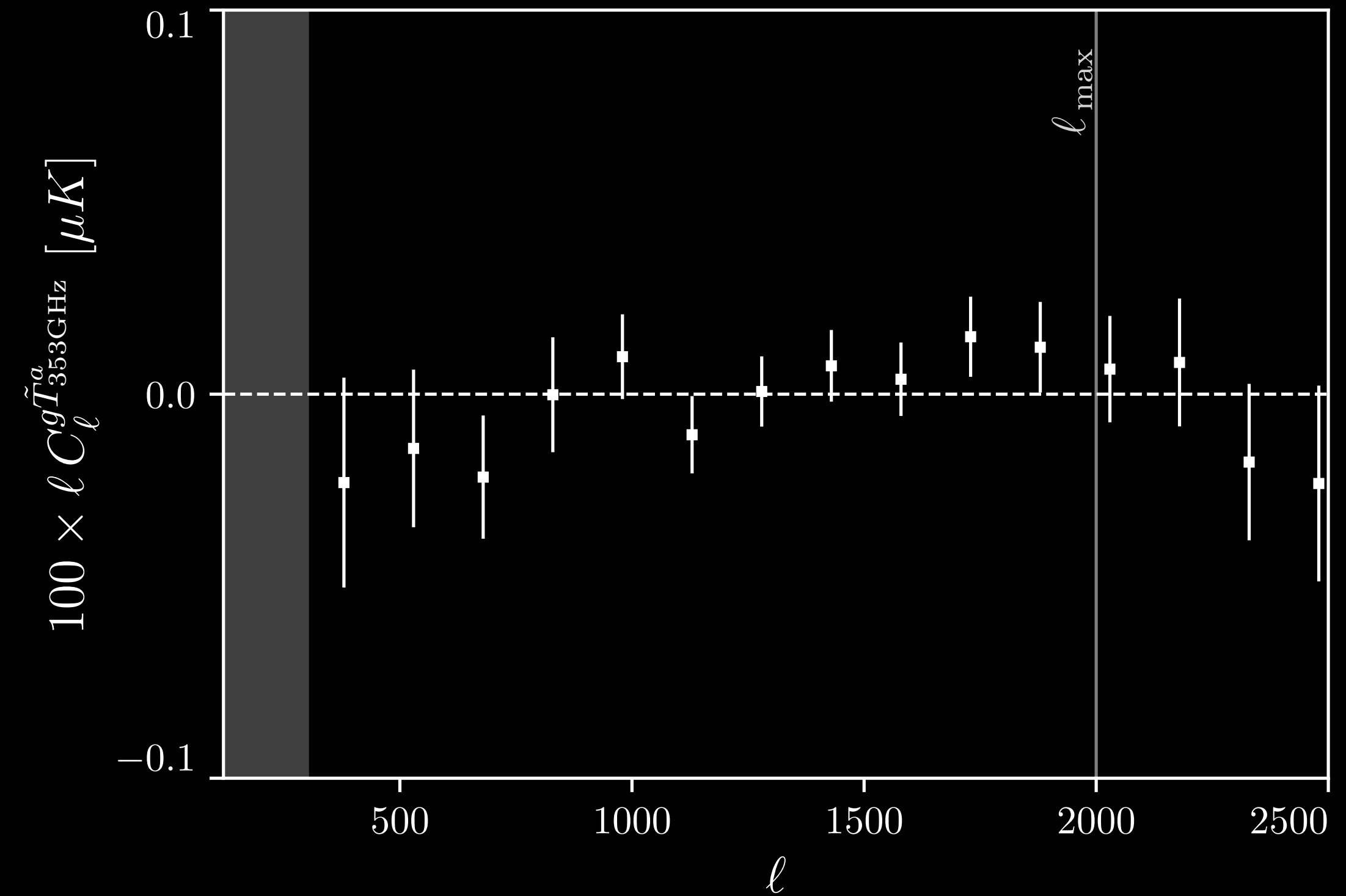
Foreground contamination

- Measurement is sensitive to any residual foregrounds that correlate with unWISE galaxies

1. **tSZ:** significant correlation, but SED is known
 2. **CIB:** significant correlation ([Yan++2023](#)), uncertainty in SED
- Model as modified blackbody (MBB)

$$\Theta_{\text{CIB}}(\nu | T_{\text{CIB}}^{\text{eff}}, \beta_{\text{CIB}}^{\text{eff}}) = \nu^{\beta_{\text{CIB}}^{\text{eff}}} B(\nu | T_{\text{CIB}}^{\text{eff}})$$

- Deproject first moment of MBB ([Chluba++2017](#))



Deproject tSZ+CIB MBB+ $d\beta_{\text{CIB}}$

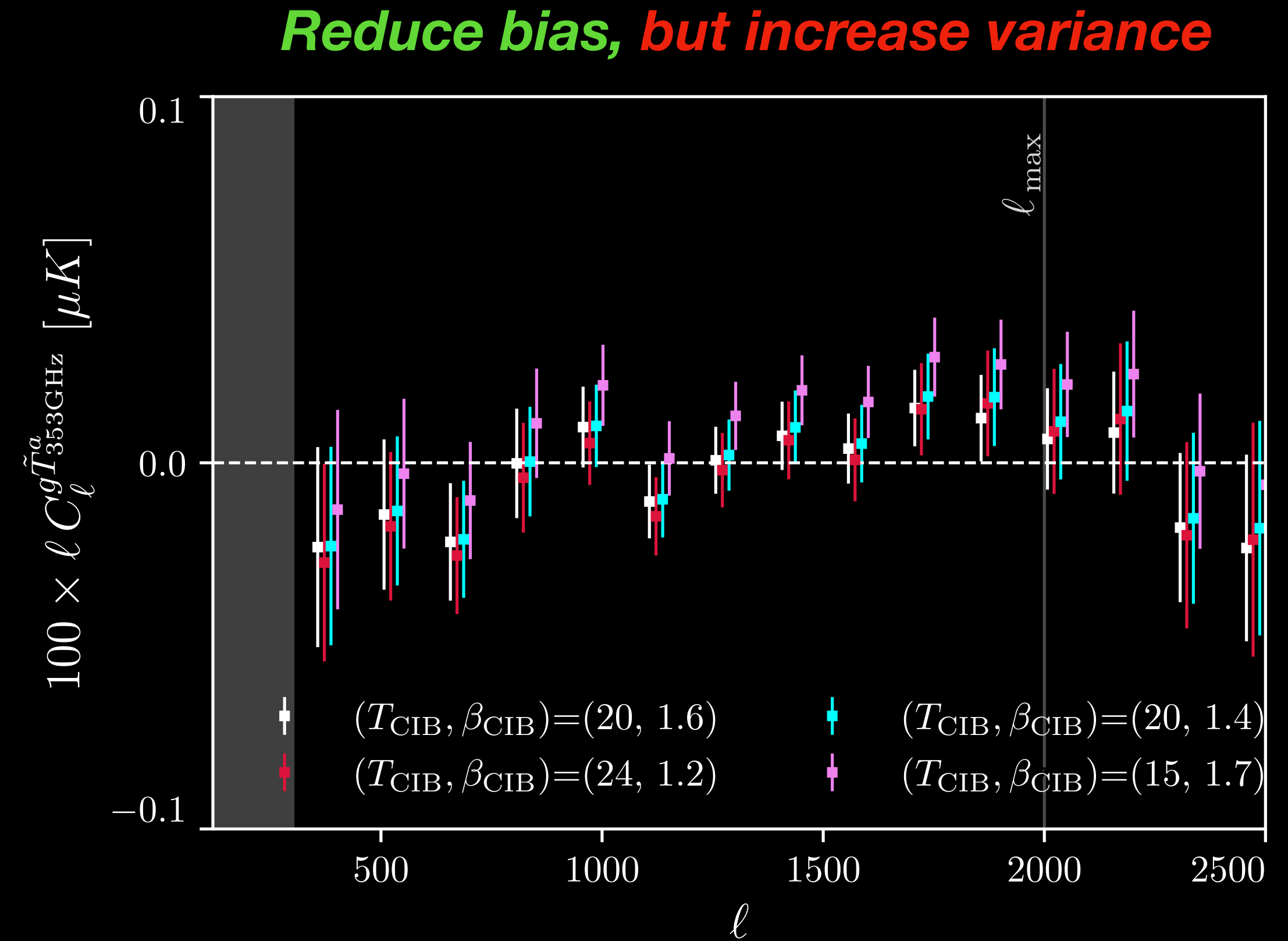
Foreground contamination

- Measurement is sensitive to any residual foregrounds that correlate with unWISE galaxies

- tSZ:** significant correlation, but SED is known
 - CIB:** significant correlation ([Yan++2023](#)), uncertainty in SED
- Model as modified blackbody (MBB)

$$\Theta_{\text{CIB}}(\nu | T_{\text{CIB}}^{\text{eff}}, \beta_{\text{CIB}}^{\text{eff}}) = \nu^{\beta_{\text{CIB}}^{\text{eff}}} B(\nu | T_{\text{CIB}}^{\text{eff}})$$

- Deproject first moment of MBB ([Chluba++2017](#))



Deproject tSZ+CIB MBB+ $d\beta_{\text{CIB}}$

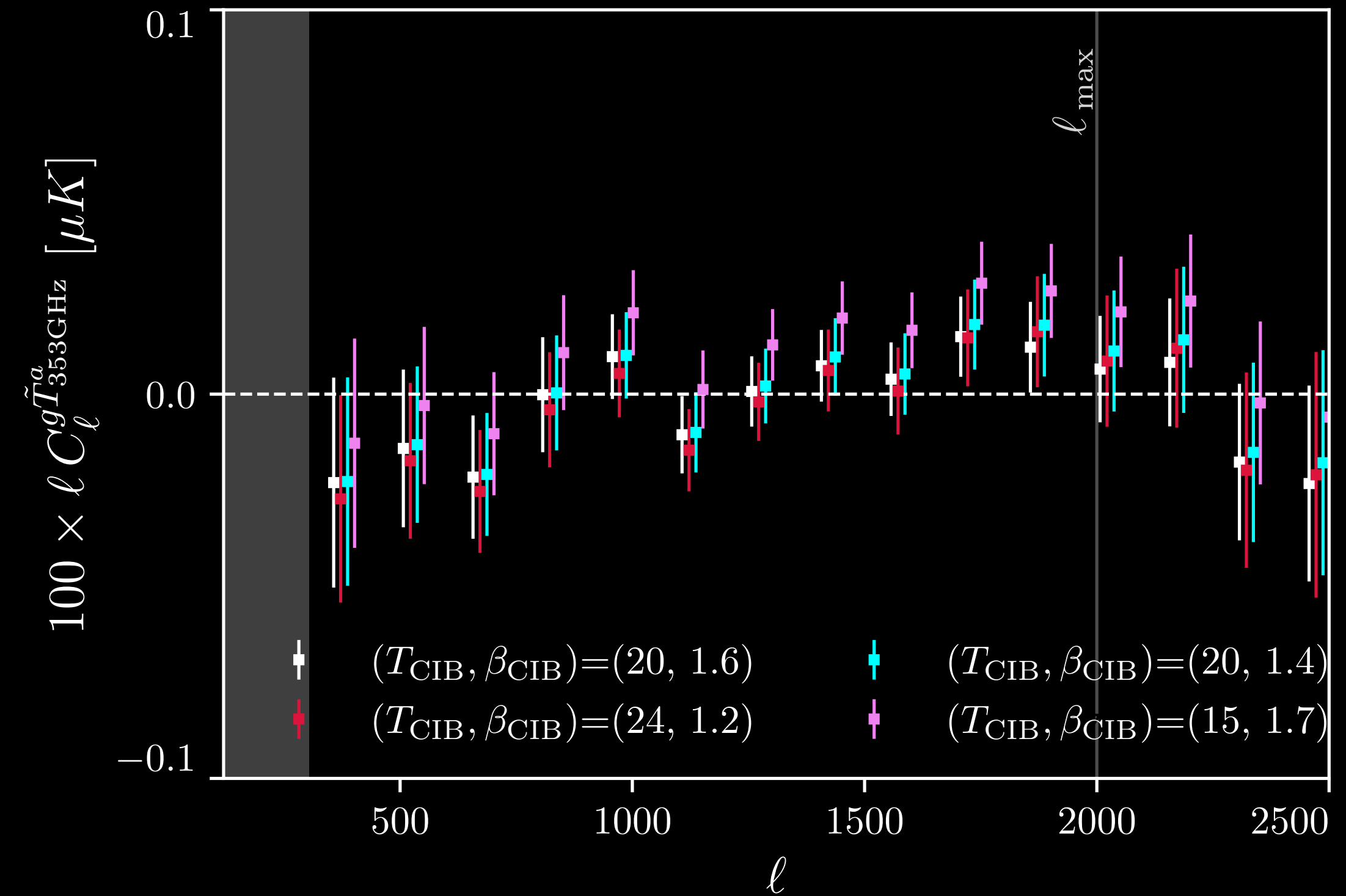
Foreground contamination

- Measurement is sensitive to any residual foregrounds that correlate with unWISE galaxies

- tSZ:** significant correlation, but SED is known
- CIB:** significant correlation ([Yan++2023](#)), uncertainty in SED
 - Model as modified blackbody (MBB)

$$\Theta_{\text{CIB}}(\nu | T_{\text{CIB}}^{\text{eff}}, \beta_{\text{CIB}}^{\text{eff}}) = \nu^{\beta_{\text{CIB}}^{\text{eff}}} B(\nu | T_{\text{CIB}}^{\text{eff}})$$

- Deproject first moment of MBB ([Chluba++2017](#))



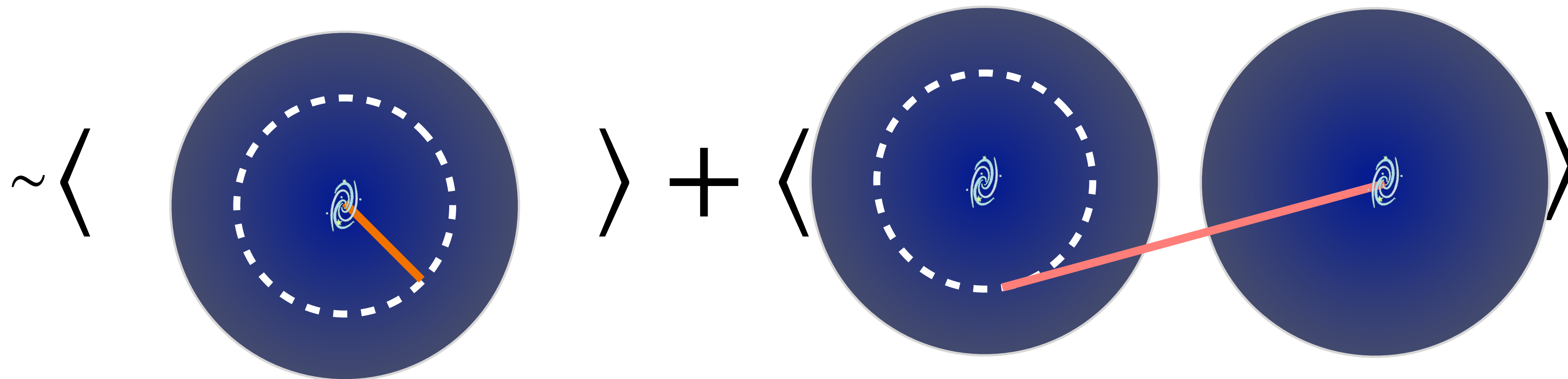
We find no evidence for axion-induced spectral distortions in *Planck*xunWISE!

Deproject tSZ+CIB MBB+ $d\beta_{\text{CIB}}$

Converting non-detection into constraints

- Theoretical predictions are computed using halo model

$$\langle \tilde{T}_a g \rangle = C_{\ell}^{\tilde{T}_a g; 1 \text{ halo}} + C_{\ell}^{\tilde{T}_a g; 2 \text{ halo}}$$



$$\int dz \frac{\chi(z)^2}{H(z)} \int dM n(z, M) \tilde{T}_{\ell}^a(z, M) u_{\ell}^g(z, M)$$

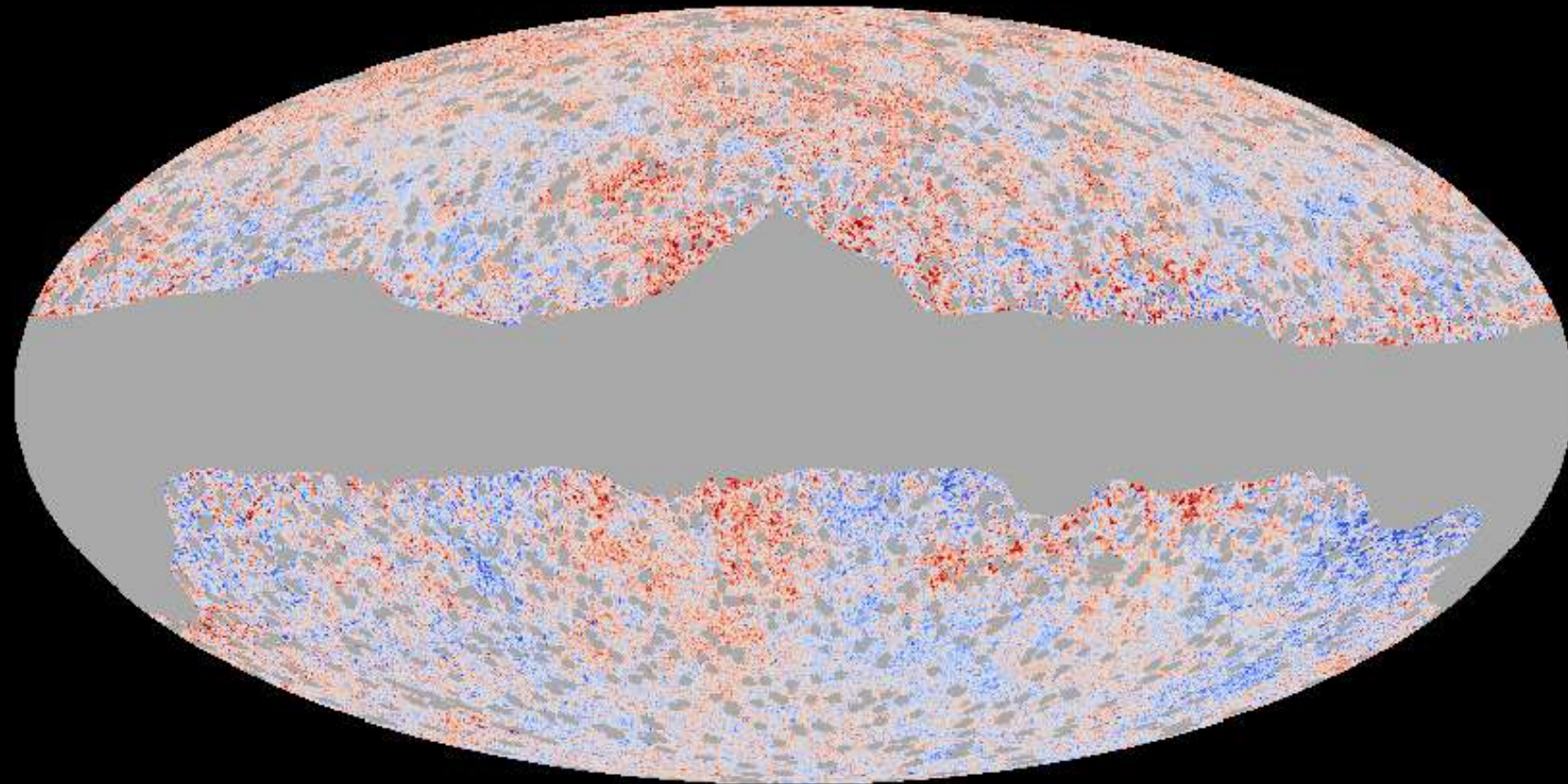
$$\int dz \frac{\chi(z)^2}{H(z)} \left[\prod_{i=1,2} \int dM_i n(z, M_i) b(z, M_i) \right] u_{\ell}^g(z, M_1) \tilde{T}_{\ell}^a(z, M_2) P_{mm}^{\text{lin}} \left(\frac{\ell+1/2}{\chi(z)}, z \right),$$

Additional ingredients

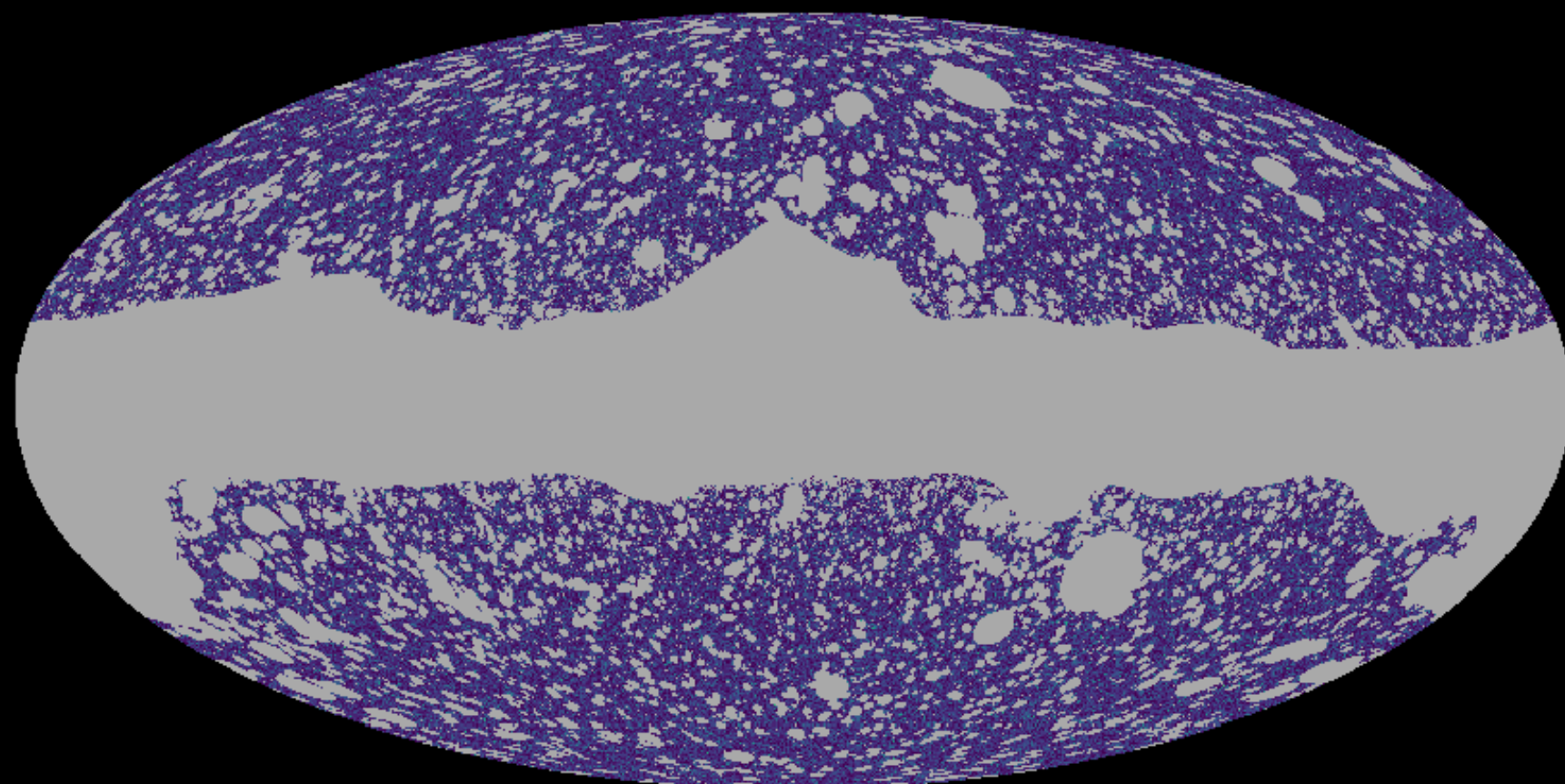
- Halo occupation distribution (HOD) from Kusiak++22
- Electron number density profile from Battaglia++ 2011
- Magnetic field profiles from IllustrisTNG

From measurements to constraints

Axion Map

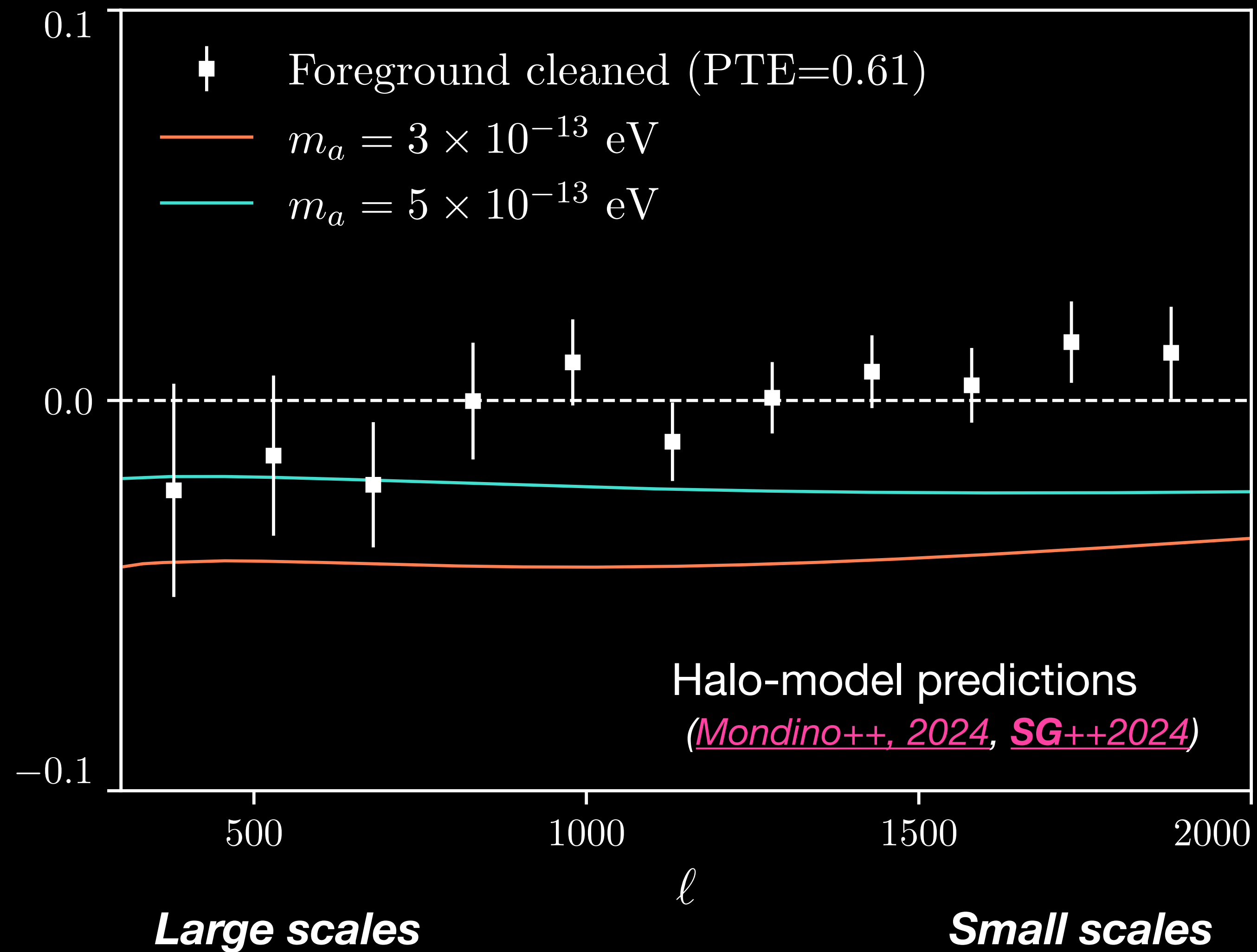


×



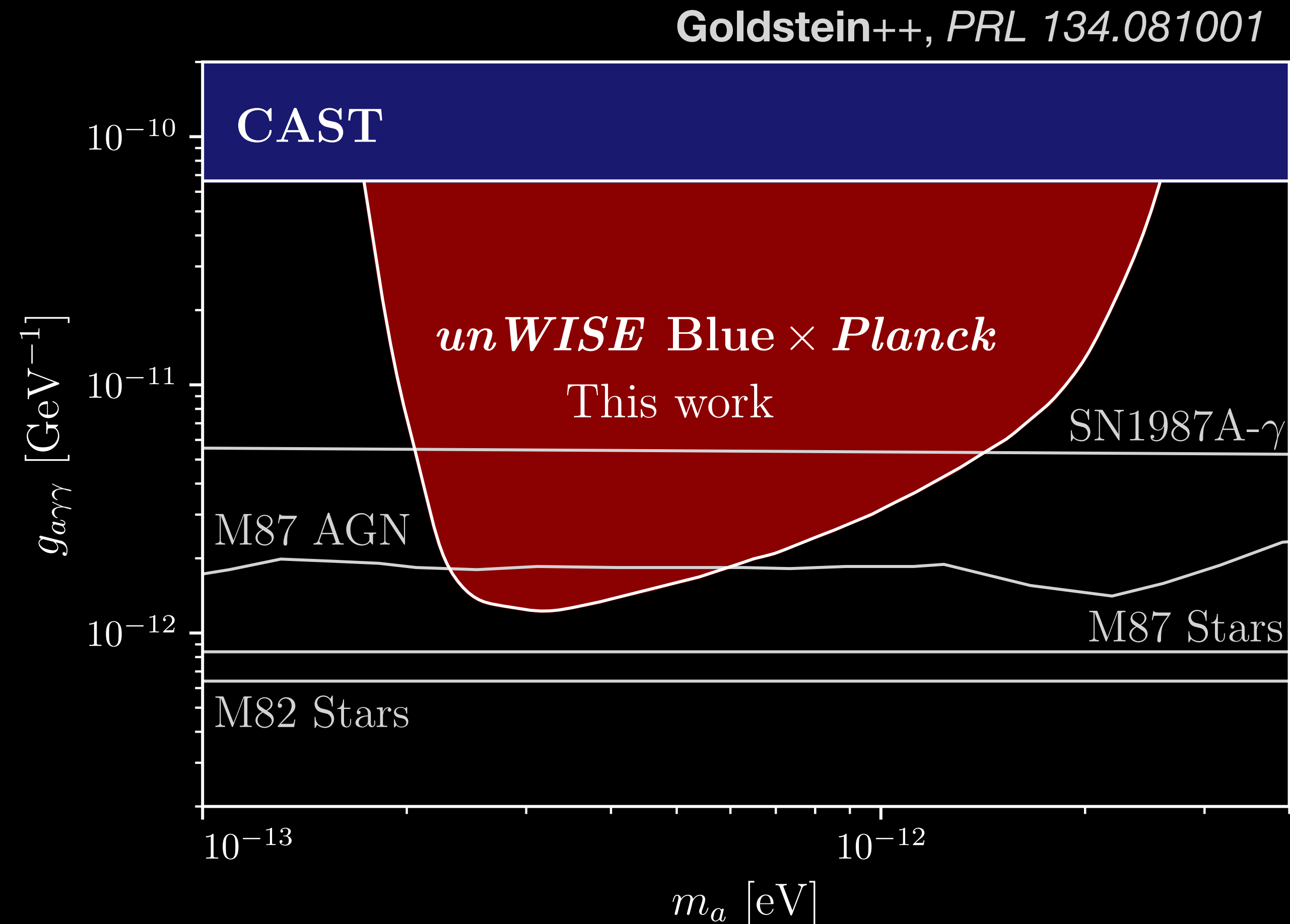
unWISE Galaxies

$$100 \times \ell C_{\ell}^{g\tilde{T}_a^a} [\mu K]$$



Constraints on axion-photon coupling

- We obtain some of the tightest constraints on $g_{a\gamma\gamma}$ for $m_a \approx 5 \times 10^{-13}$ eV axions
- Highly complementary to existing bounds
- Existing bounds: **single astrophysical object**
- Our bounds: ensemble average over **30 million objects**

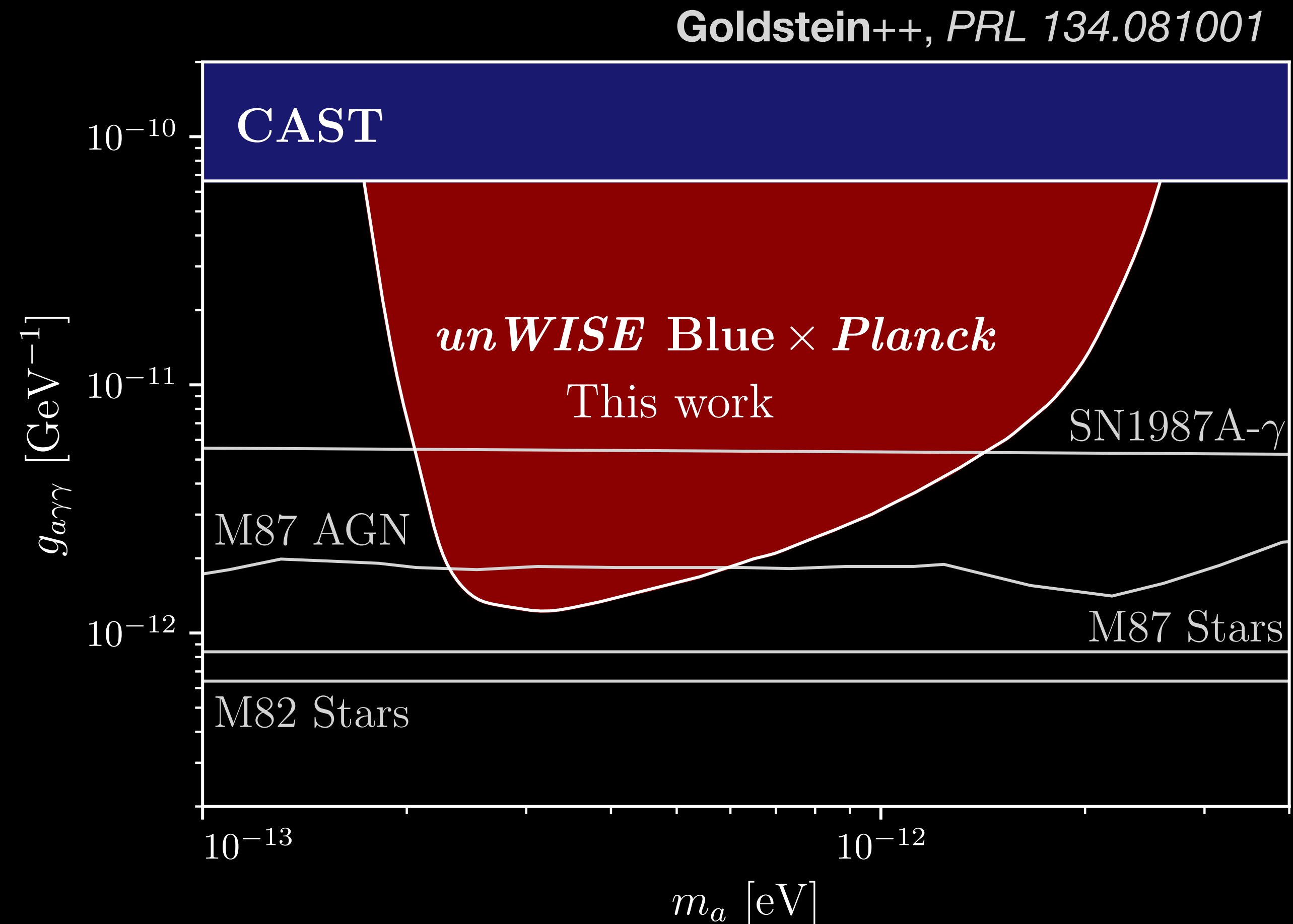


95% confidence interval on $g_{a\gamma\gamma}$
(<https://cajohare.github.io/AxionLimits/>)

Constraints on axion-photon coupling

- We obtain some of the tightest constraints on $g_{a\gamma\gamma}$ for $m_a \approx 5 \times 10^{-13}$ eV axions
- Highly complementary to existing bounds
- Existing bounds: **single astrophysical object**
- Our bounds: ensemble average over **30 million objects**

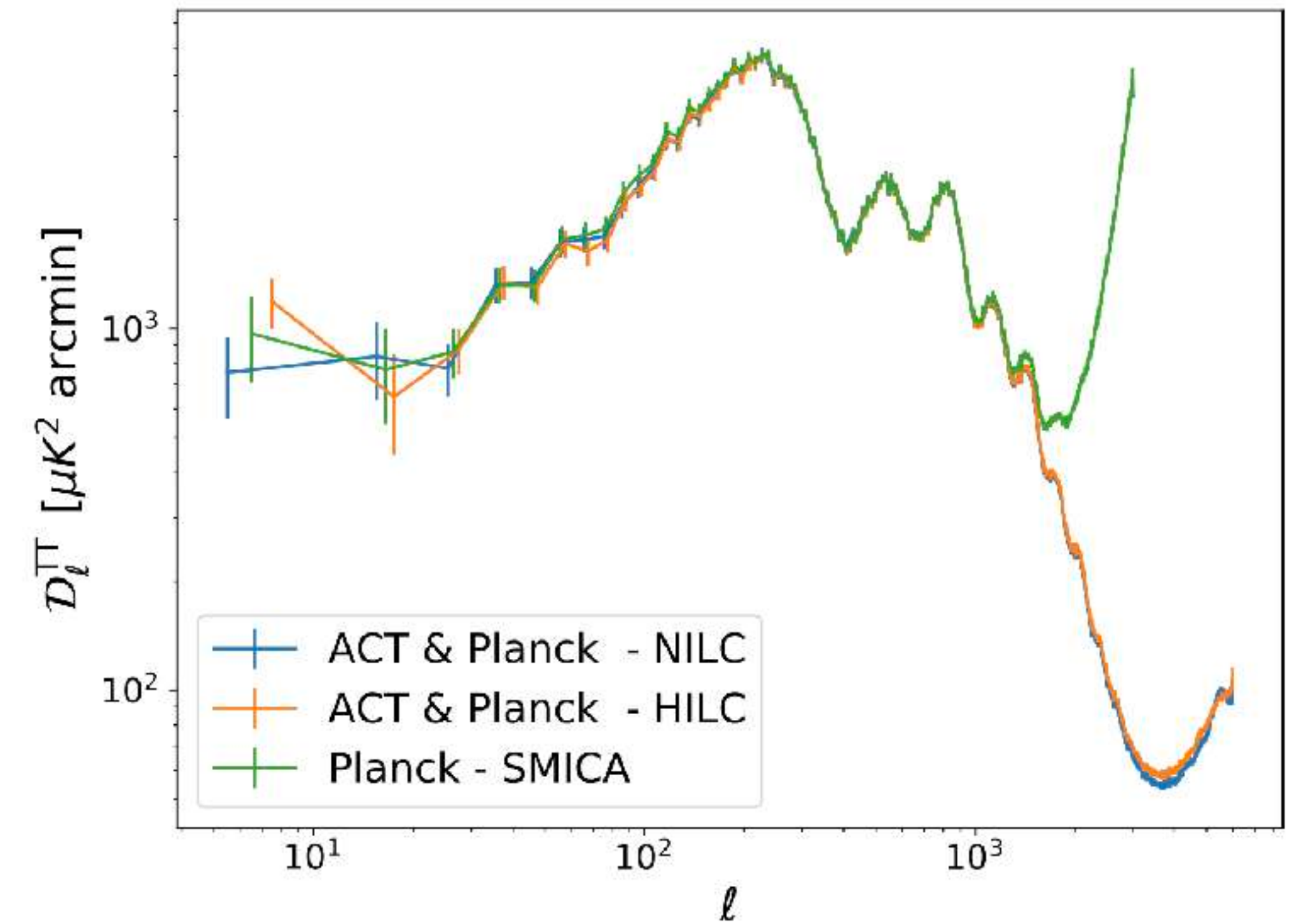
**We can do better with new ACT
DR6 data!**



95% confidence interval on $g_{a\gamma\gamma}$
(<https://cajohare.github.io/AxionLimits/>)

Adding in ACT data

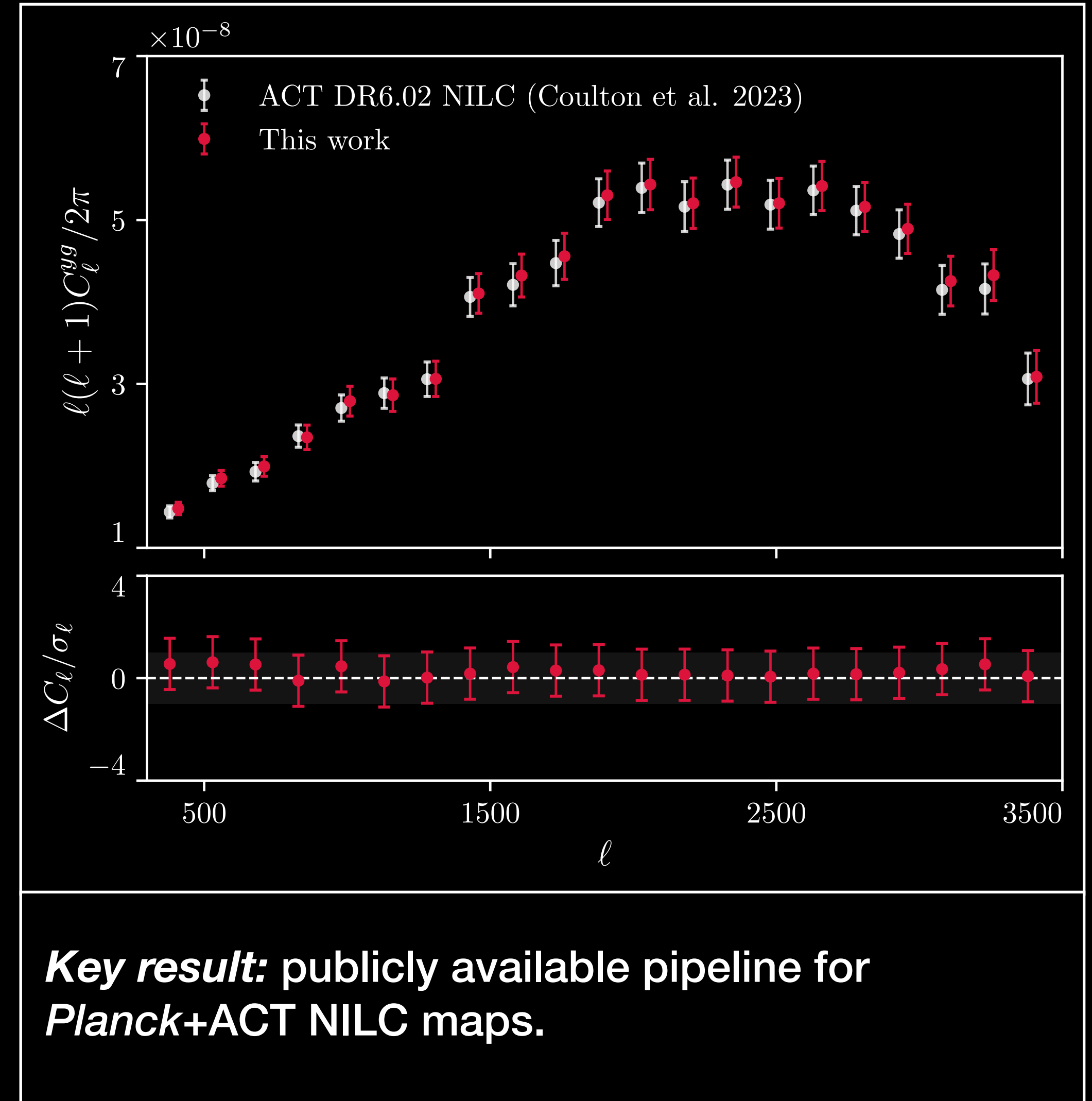
- Screening happens inside halos -> additional info at small scales
- ACT has high-resolution 90, 150, and 220 GHz maps
- **Challenges:**
 - More maps + higher resolution -> need a faster NILC pipeline
 - More sensitive data requires better foreground model and improved treatment of systematic uncertainties



Example of statistical gains in NILC maps after adding ACT data ([Coulton++23](#))

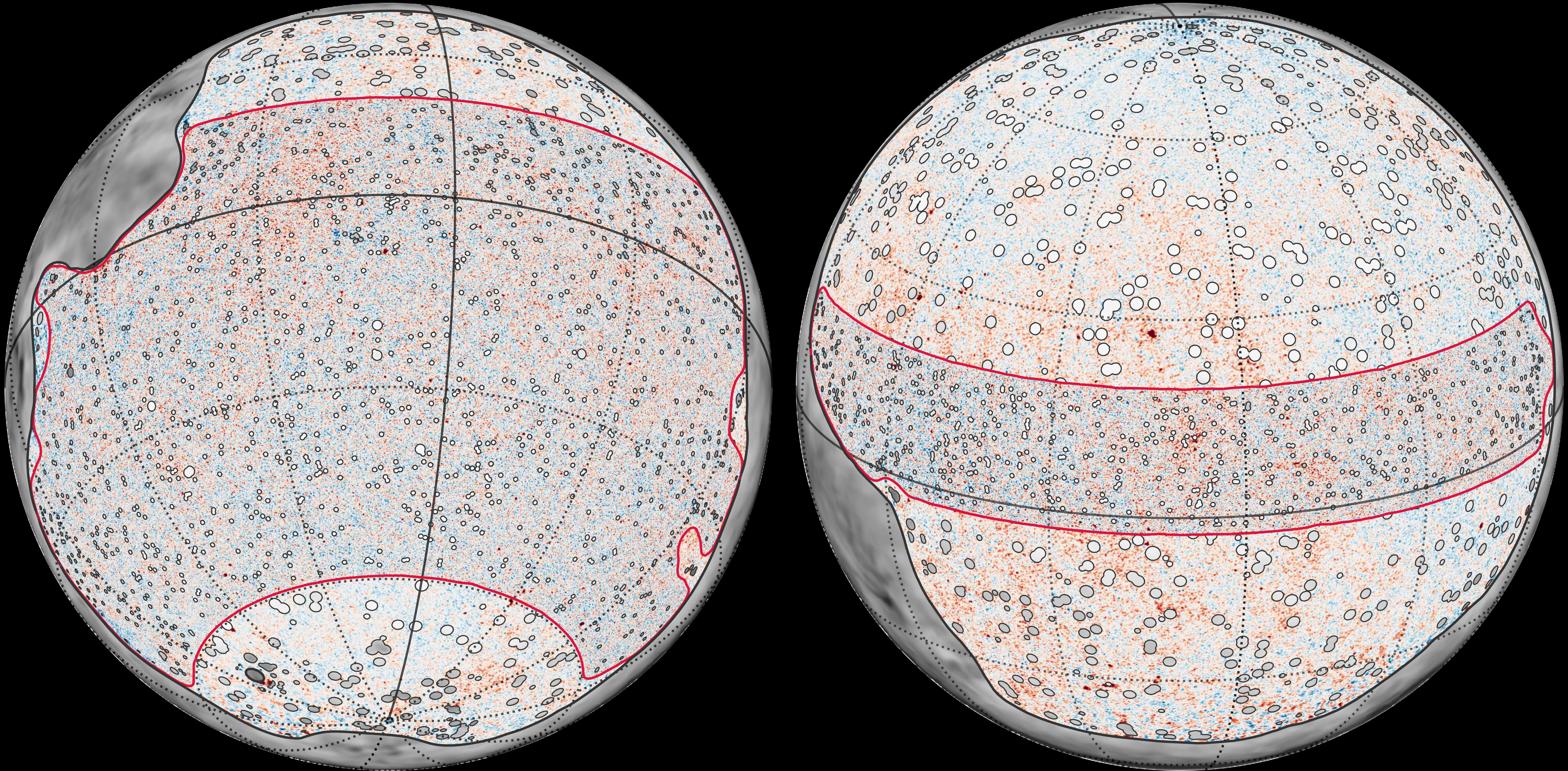
A new NILC pipeline: Coberus

- **Coberus*** dask-based NILC code developed for SO)
 - Combine ACT+Planck up to $\ell_{\text{max}} \sim 10,000$ in ~ 5 minutes on a single node
 - Full sky component separation
- Can now marginalize over calibration, passband, and beam uncertainties

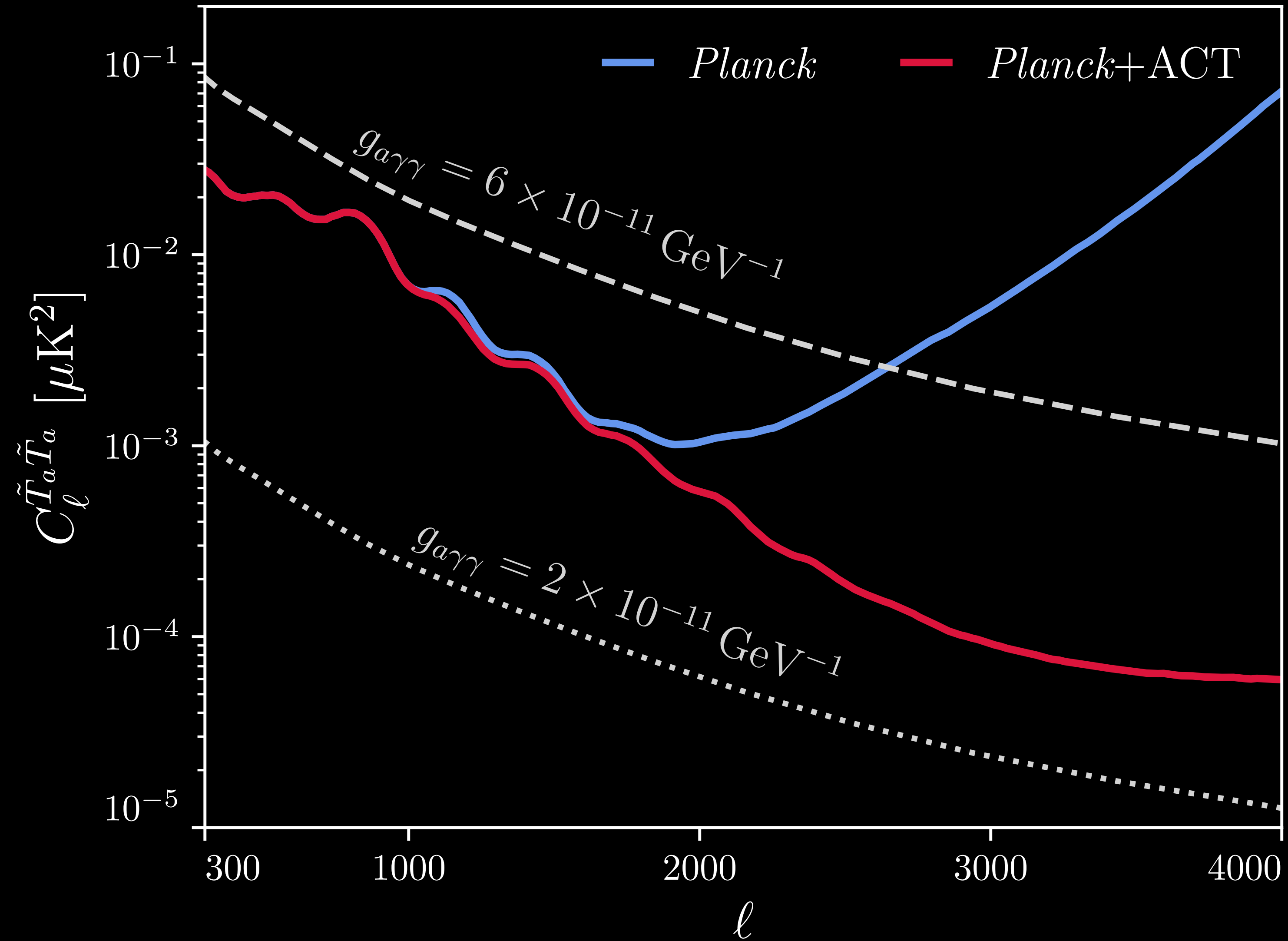


*<https://github.com/simonsobs/coberus/> (with Josh Borrow, Joshua Kim, and Mathew Madhavacheril)

Planck+ACT Axion-induced screening NILC map: no deprojections

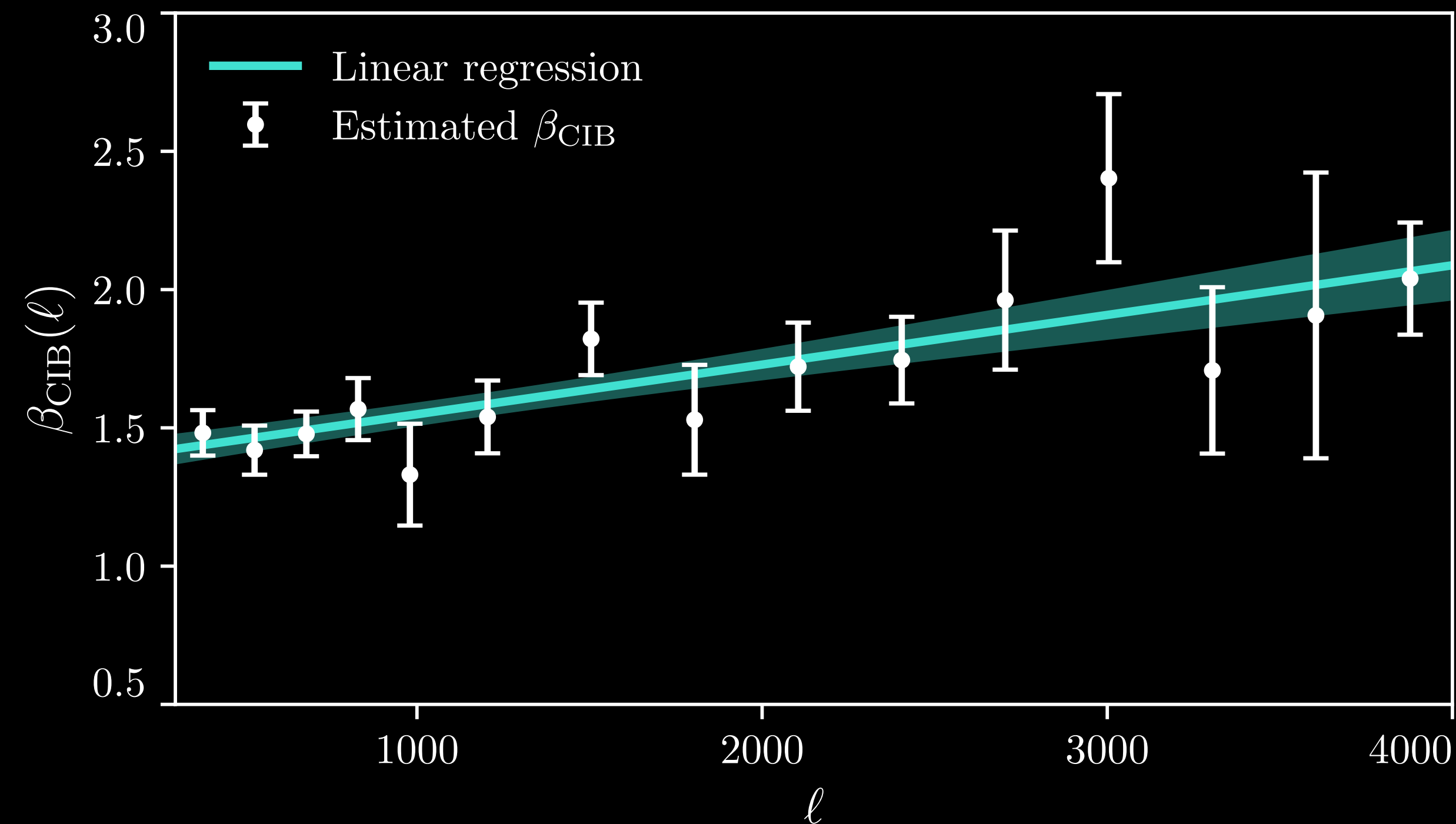
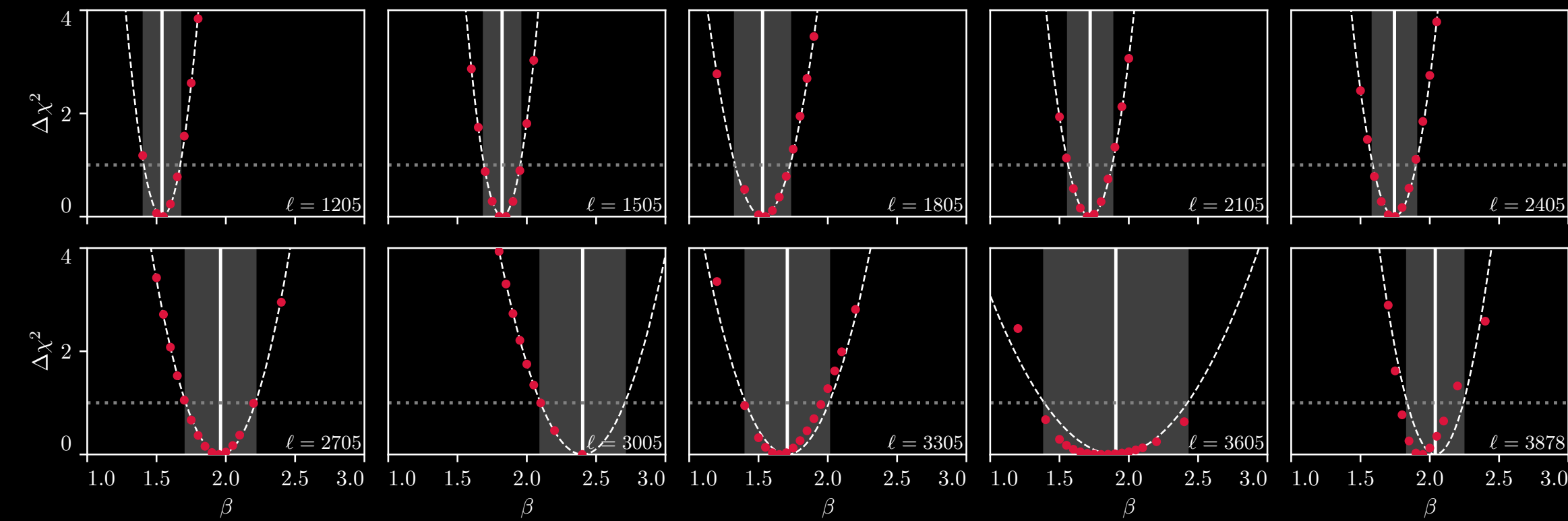


Improvement 1: Adding in ACT data

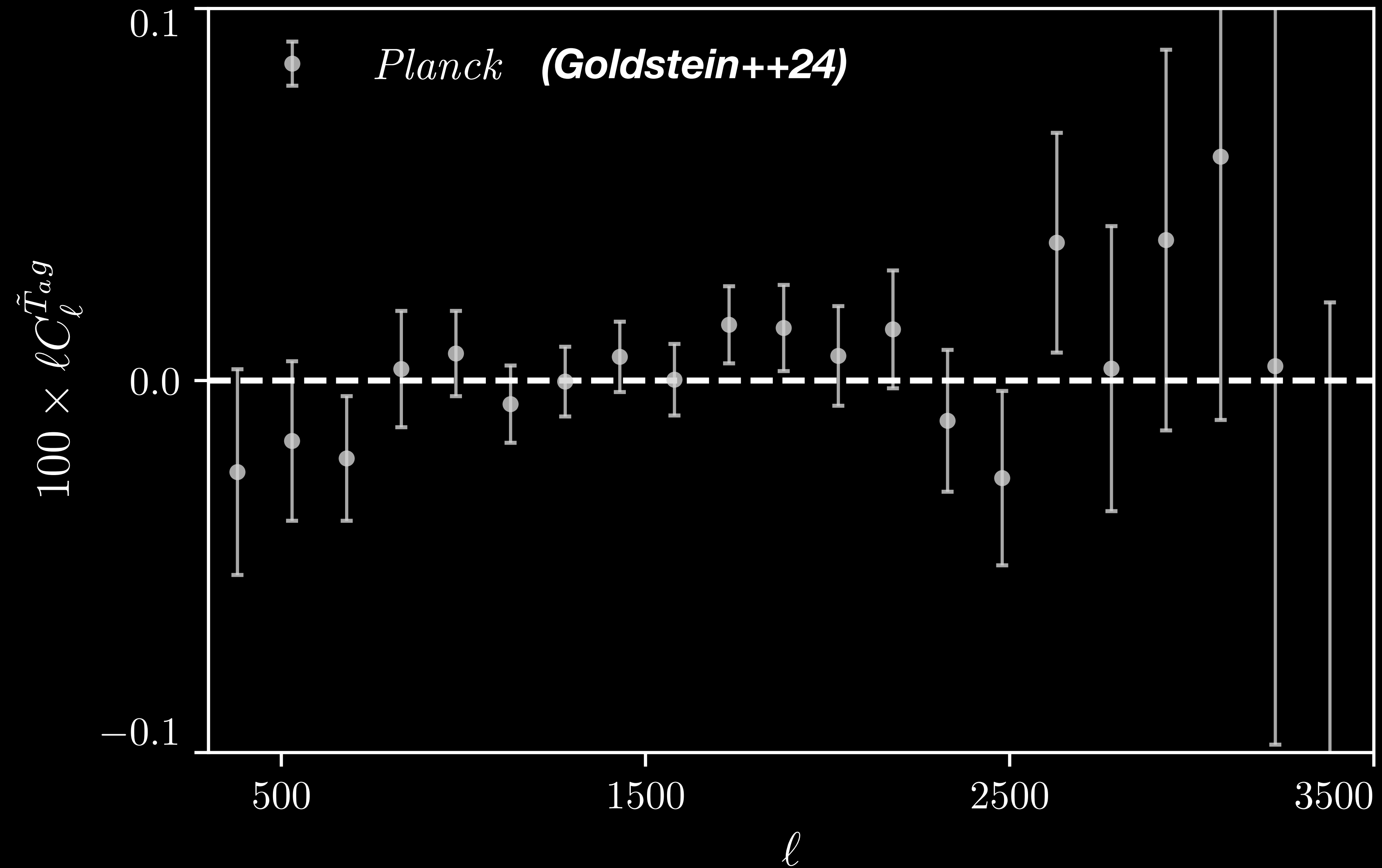


Improvement 2: Better treatment of the CIB SED

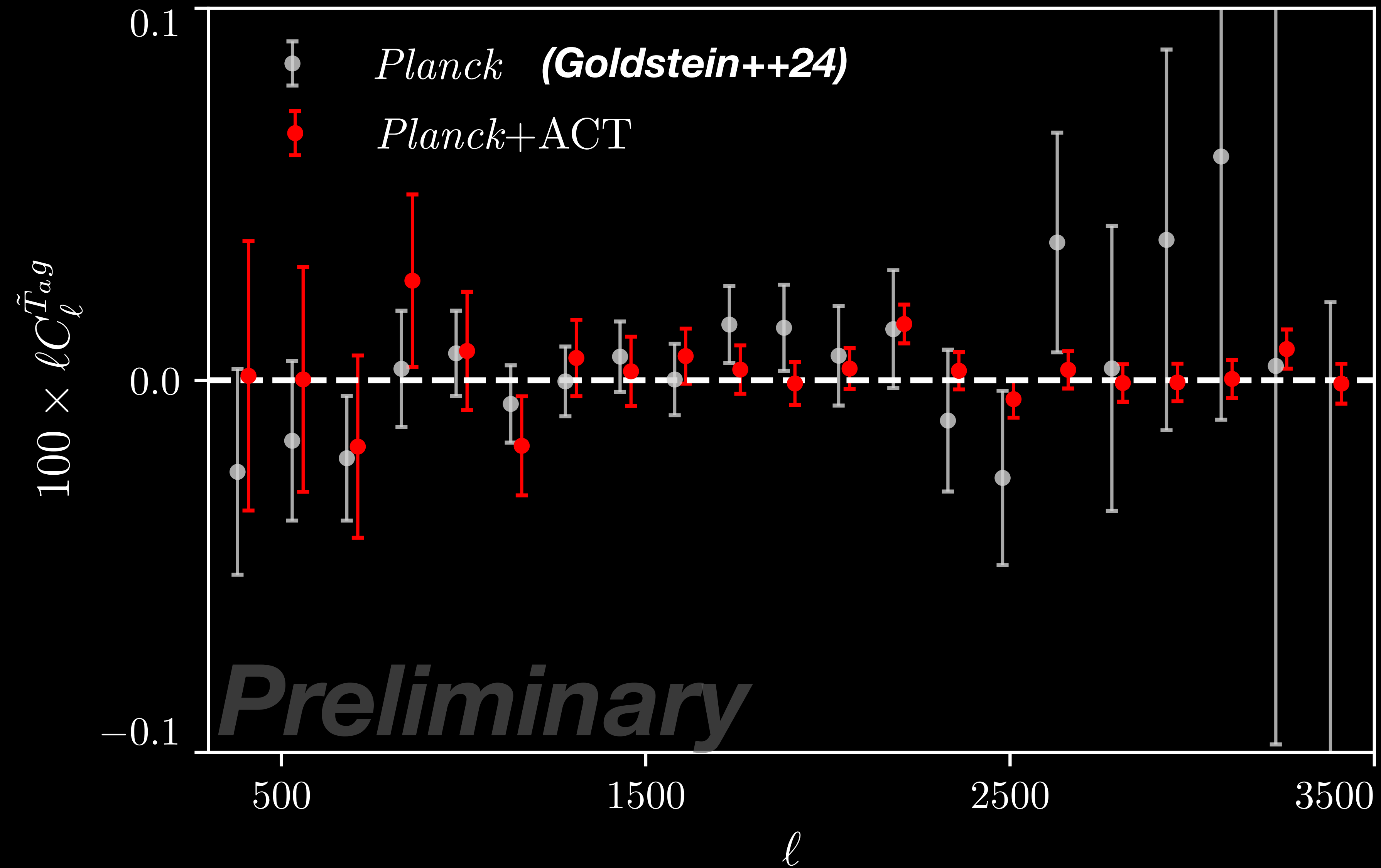
- First application of a new technique for mitigating CIB proposed in [*Surrao++25*](#)
- Data-driven approach to estimate β_{CIB} for a particular LSS tracer
- **Advantage:** no moment deprojection -> lower noise in NILC map
- **Challenge:** requires running 100-1000 NILCs without reusing covariances
 - Infeasible without Coberus



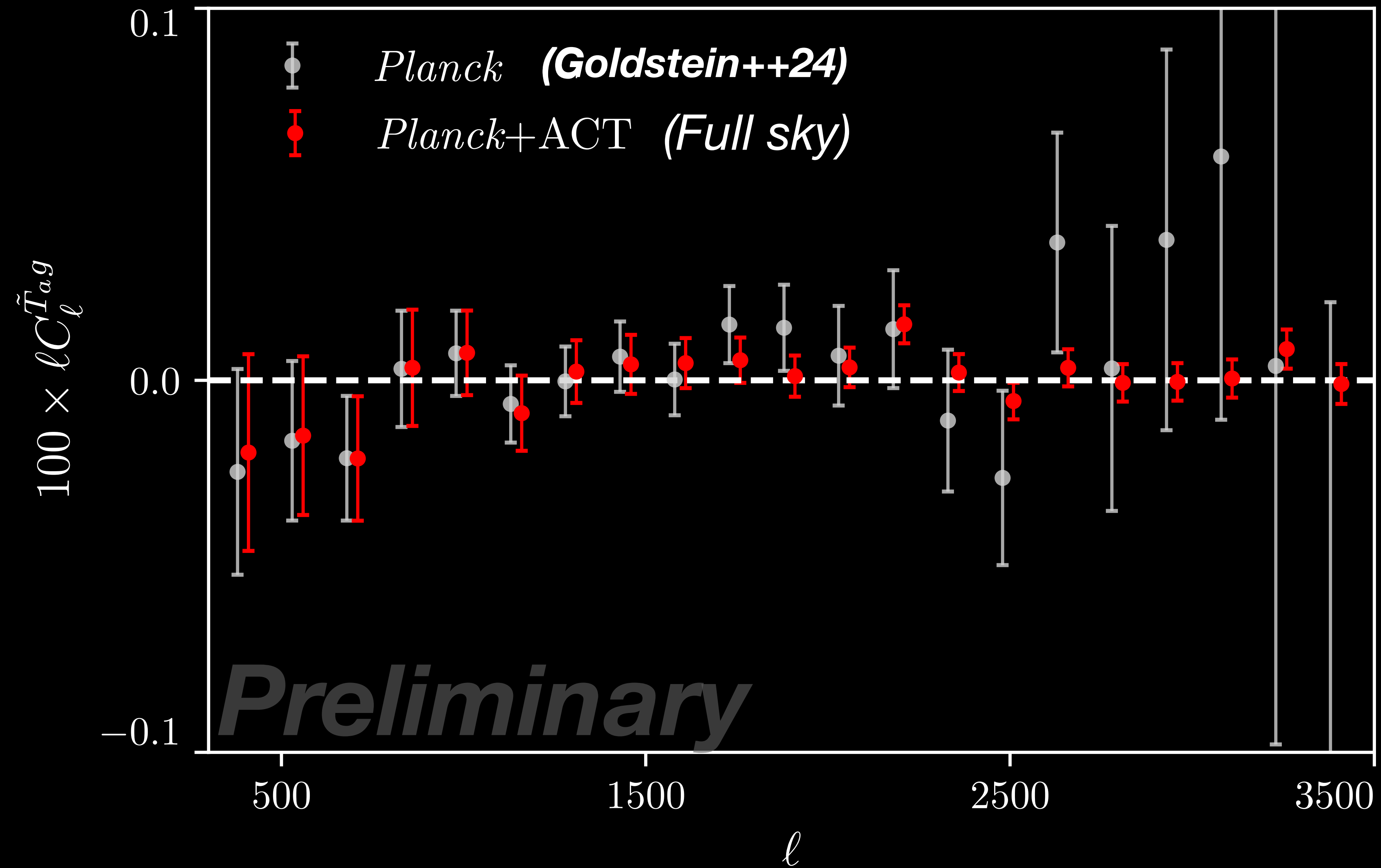
New axion-induced screening measurement



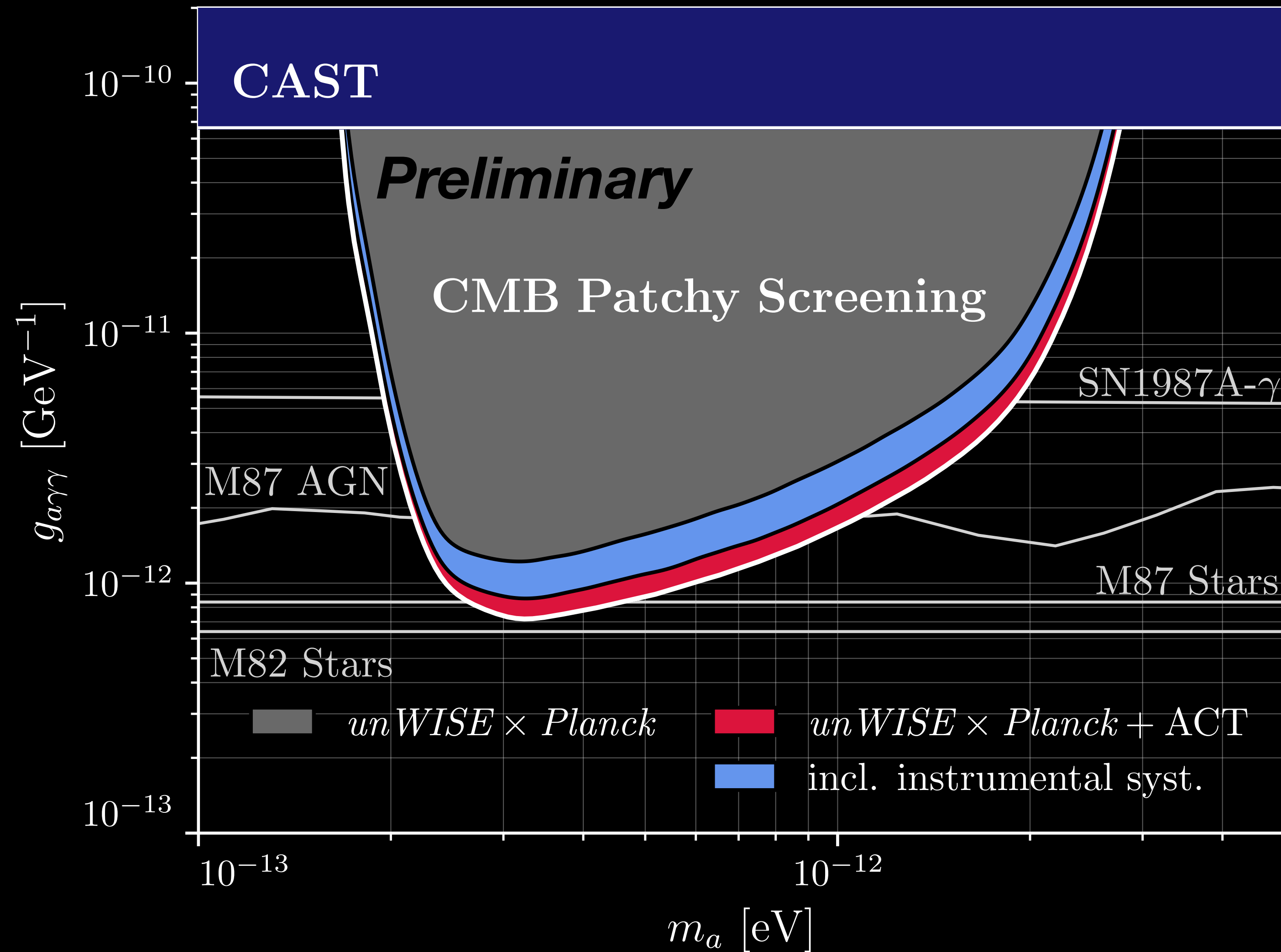
New axion-induced screening measurement



New axion-induced screening measurement

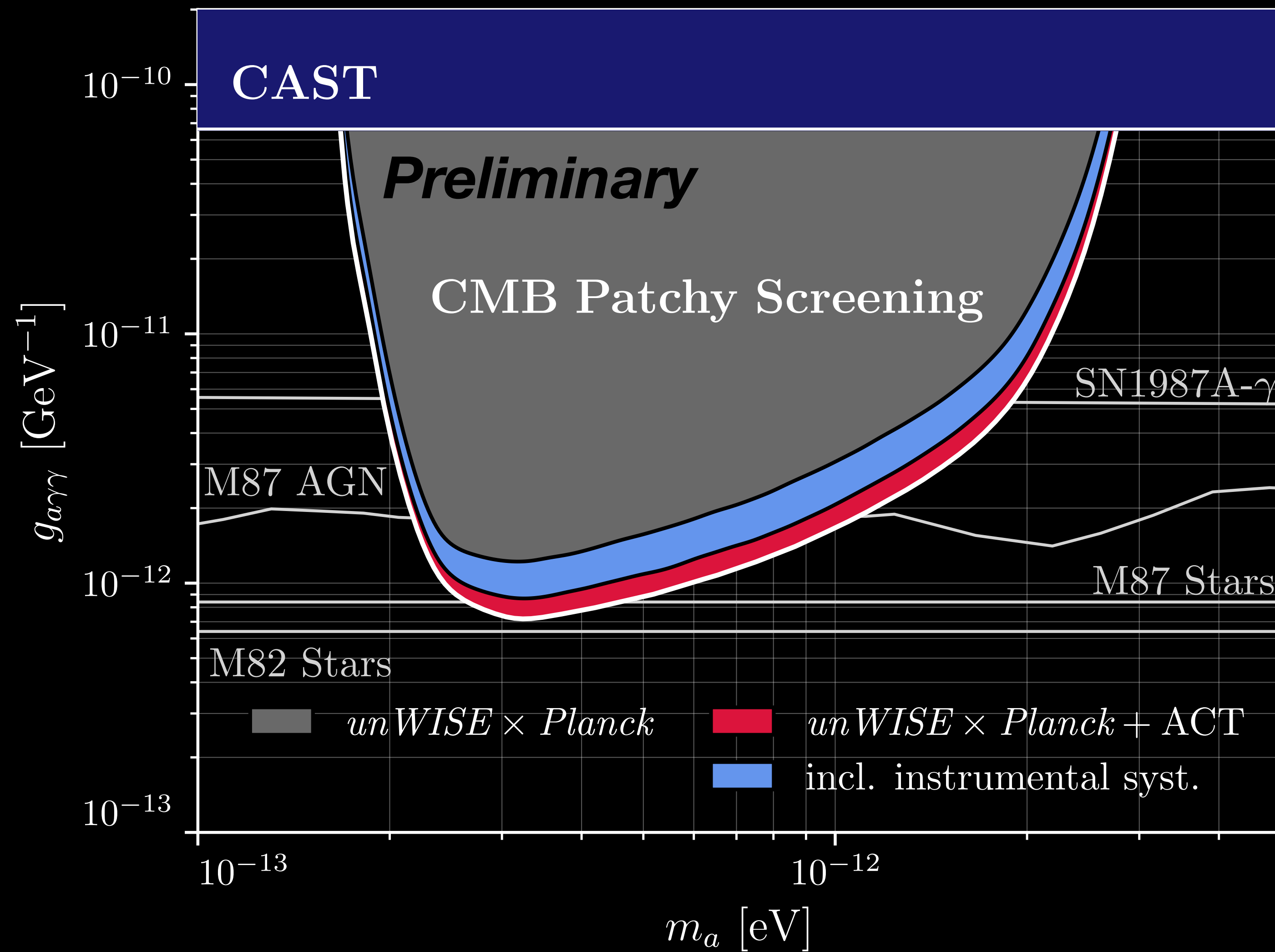


New axion-induced screening constraints



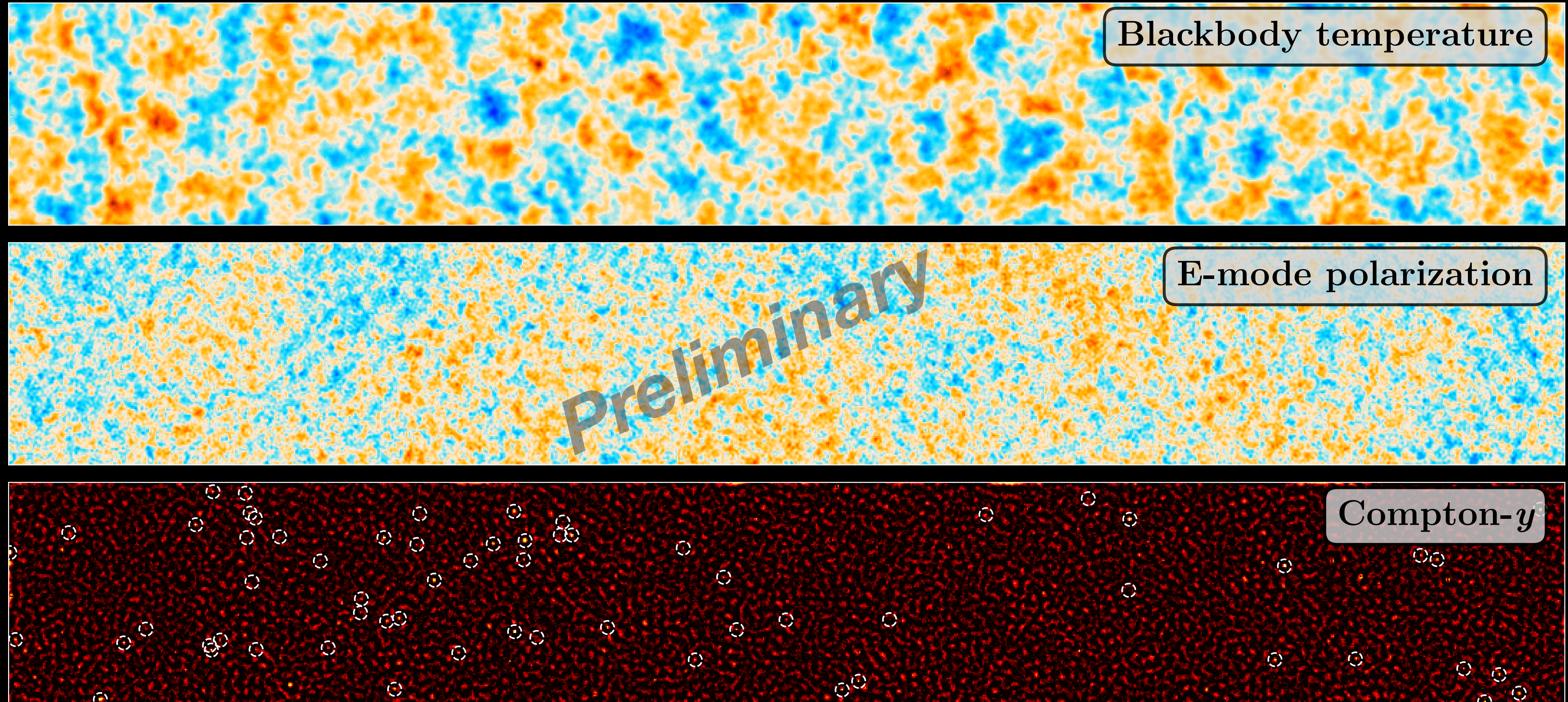
- New data+new techniques -> **nearly a factor of two reduction in stat. error**
- New pipeline enables estimate of **syst. error** from instrument (requires $\mathcal{O}(1000)$ NILC maps)

New axion-induced screening constraints



What's next for dark screening?

Future outlook: improvements from the CMB side

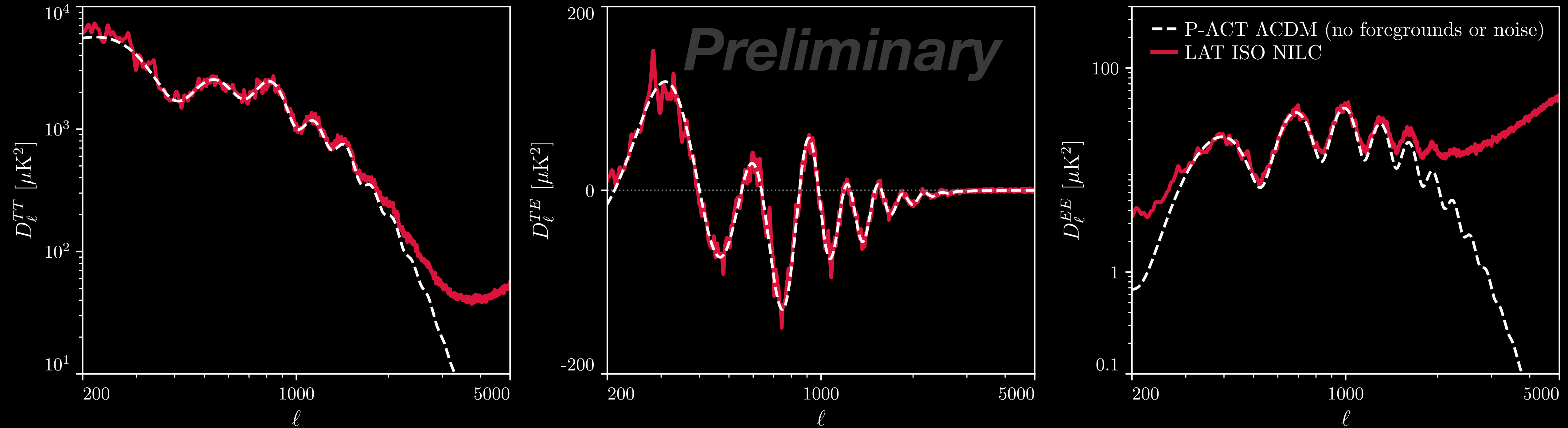


Have run NILC pipeline on early SO Large Aperture Telescope (LAT) data

Credit: Simons Observatory Collaboration

Goldstein++ (in prep.)

Future outlook: improvements from the CMB side

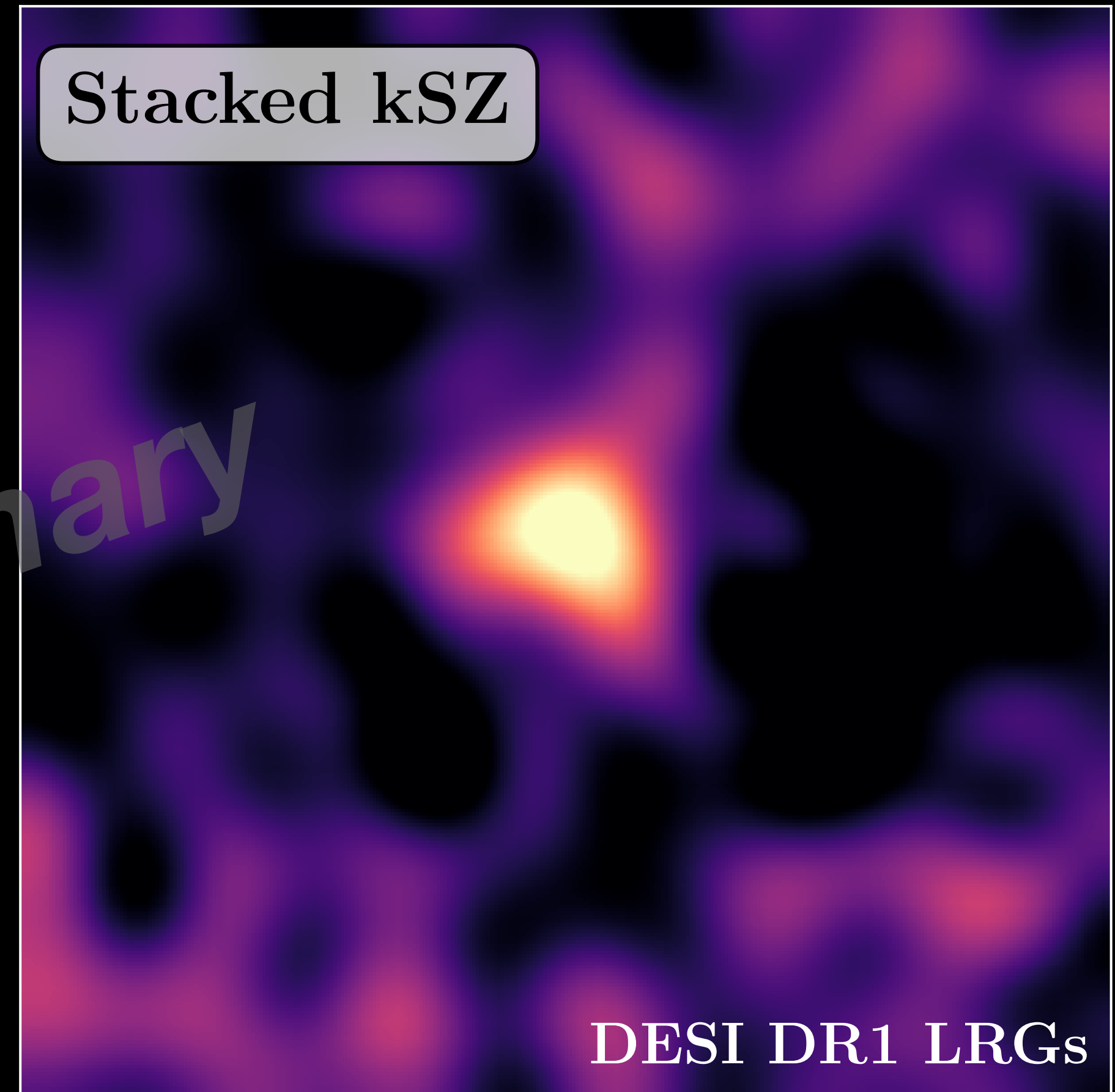
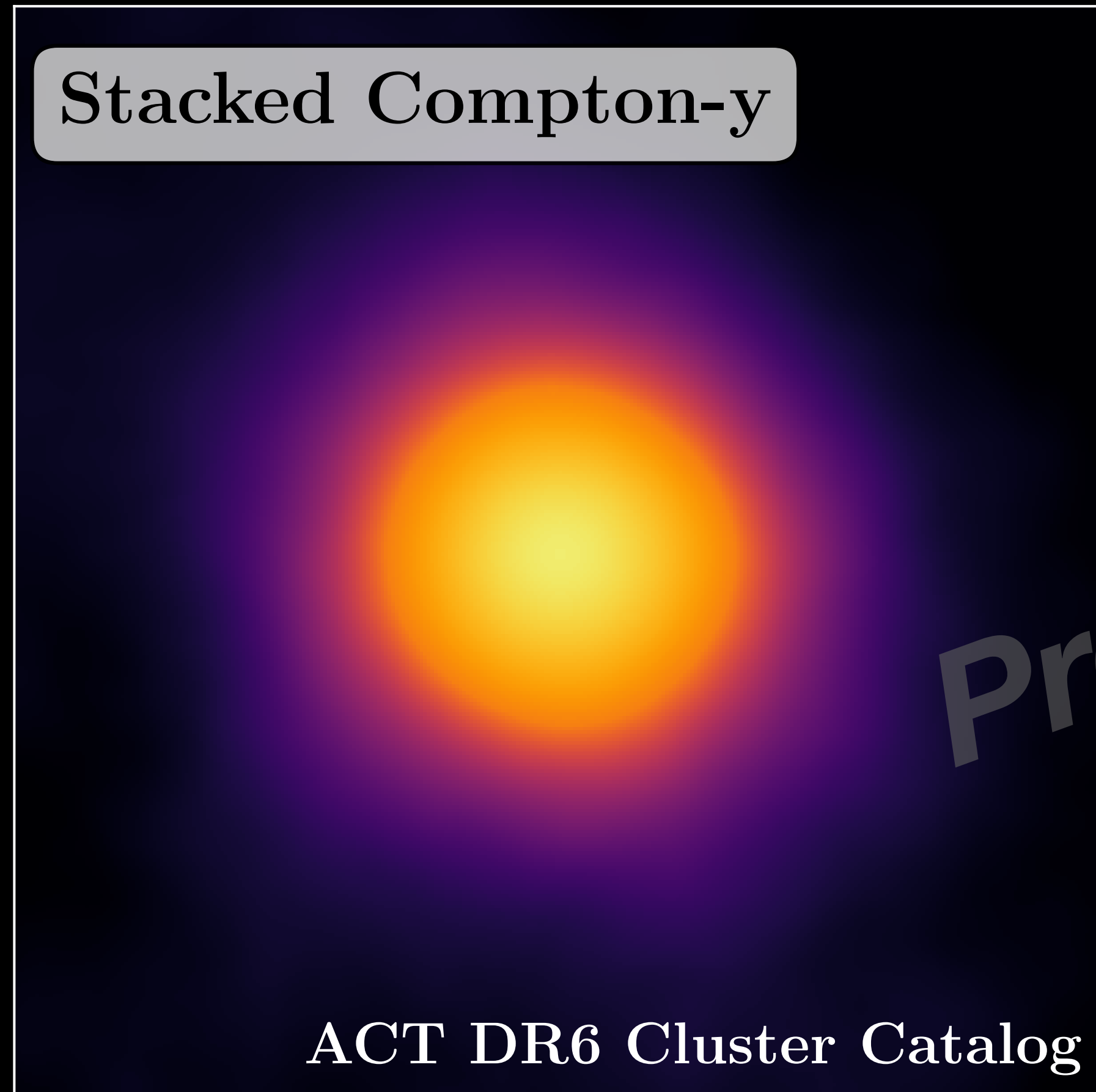


Signal-dominated T and E -mode maps from < 200 hours of data

Credit: Simons Observatory Collaboration

Goldstein++ (in prep.)

Future outlook: improvements from the CMB side



- Already have a significant detection of kSZ (*thanks also to DESI*)

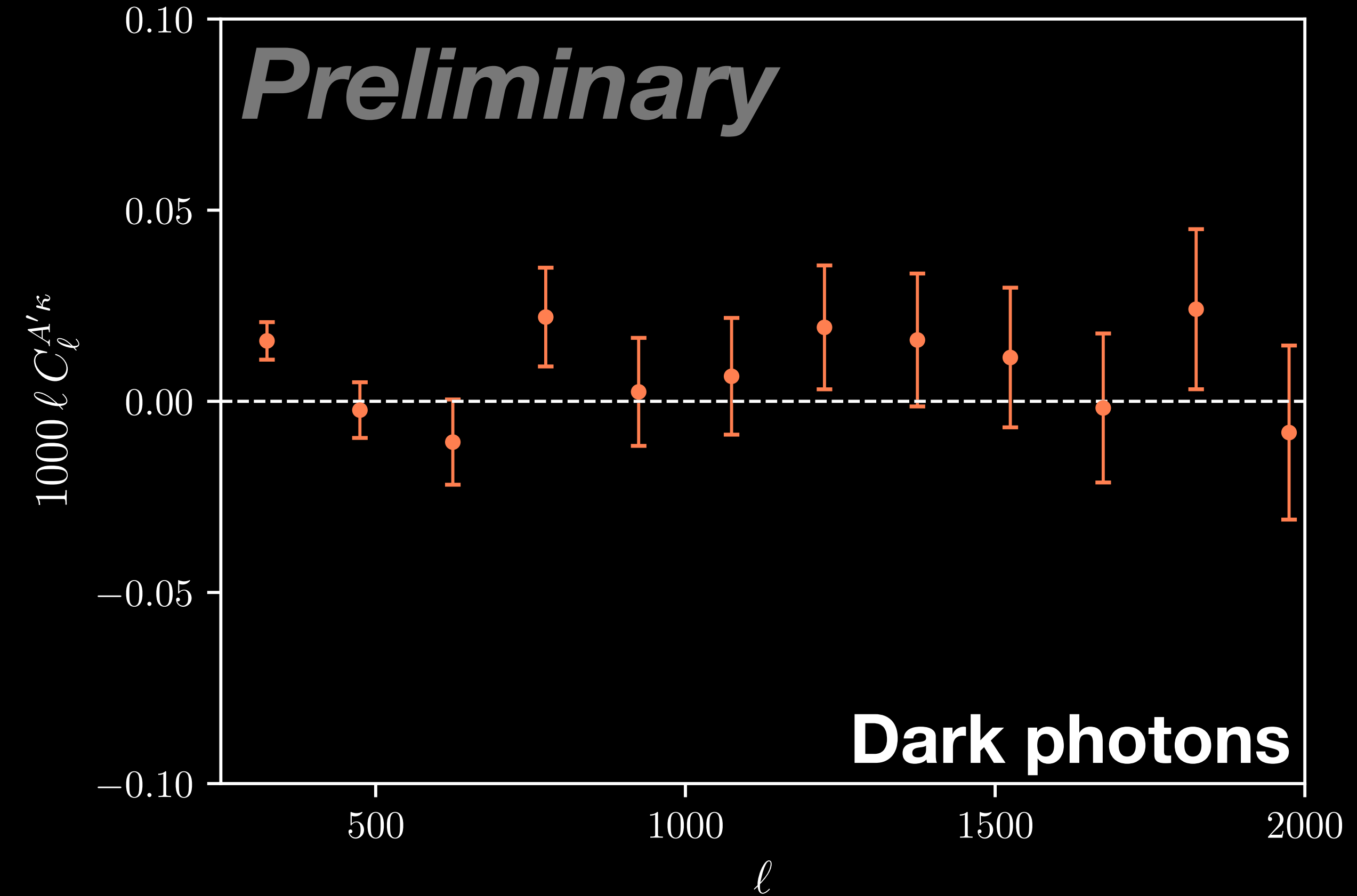
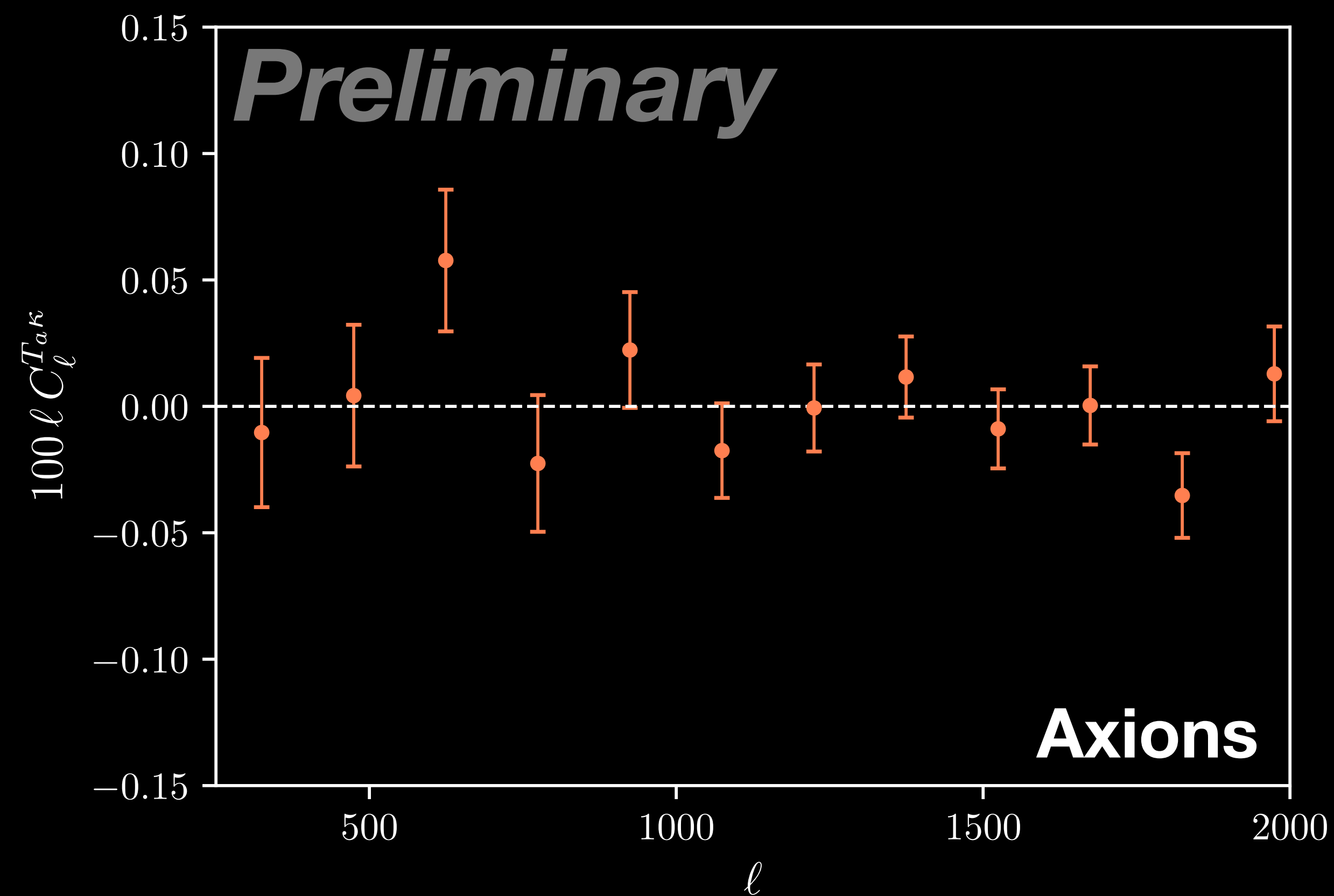
Credit: Simons Observatory Collaboration

Goldstein++ (in prep.)

Future outlook: improvements from the LSS side

- *unWISE* is great for cross-correlations, but broad and photometric
- Can do better with DESI
 - Spectroscopy $\rightarrow$ tomography: wider range of m_a and better control of CIB and HOD
 - DESI Legacy Imaging Survey DR11: 4 billion sources of 75% of the sky
- Large overlap between DESI-II and SO

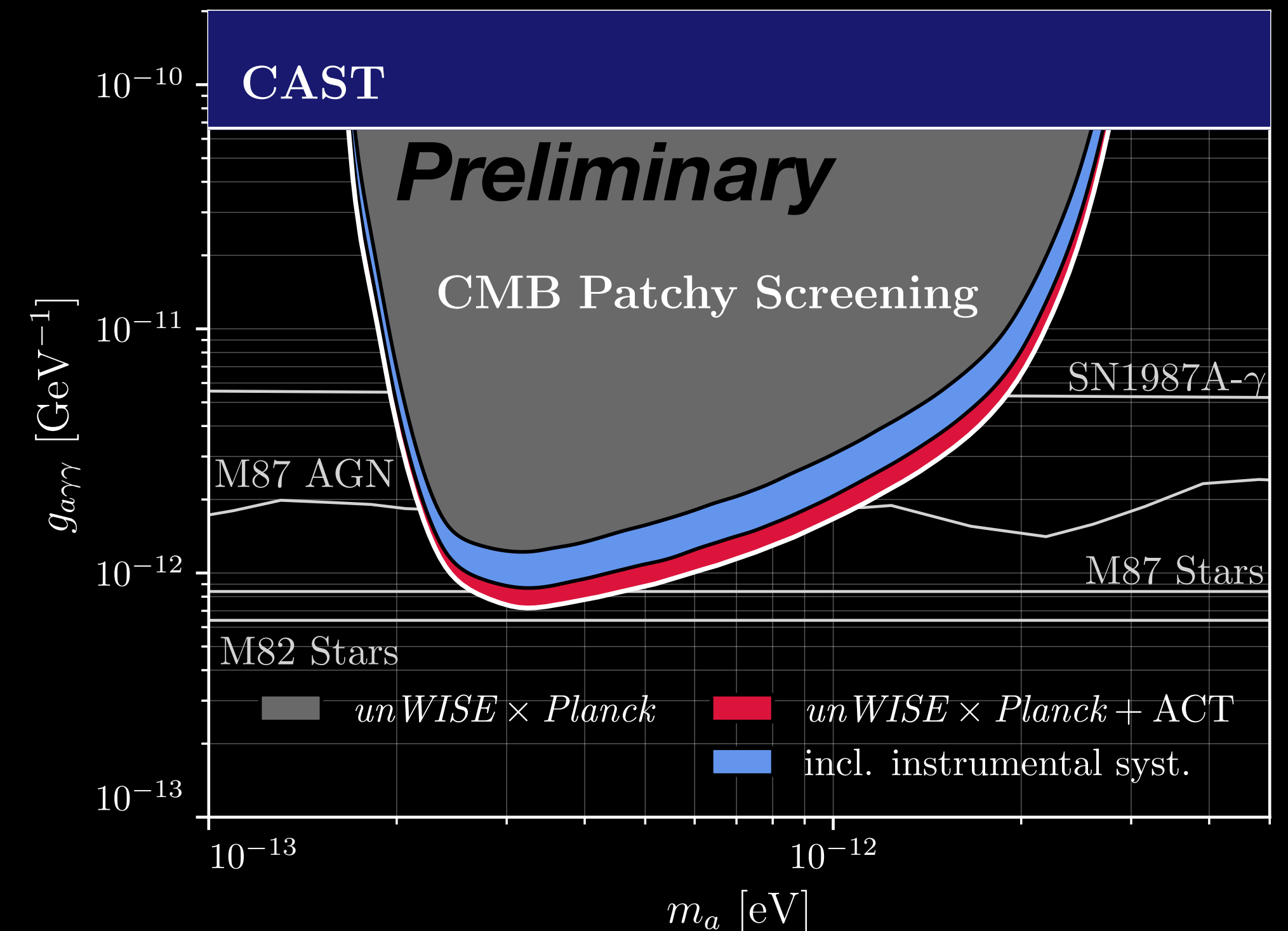
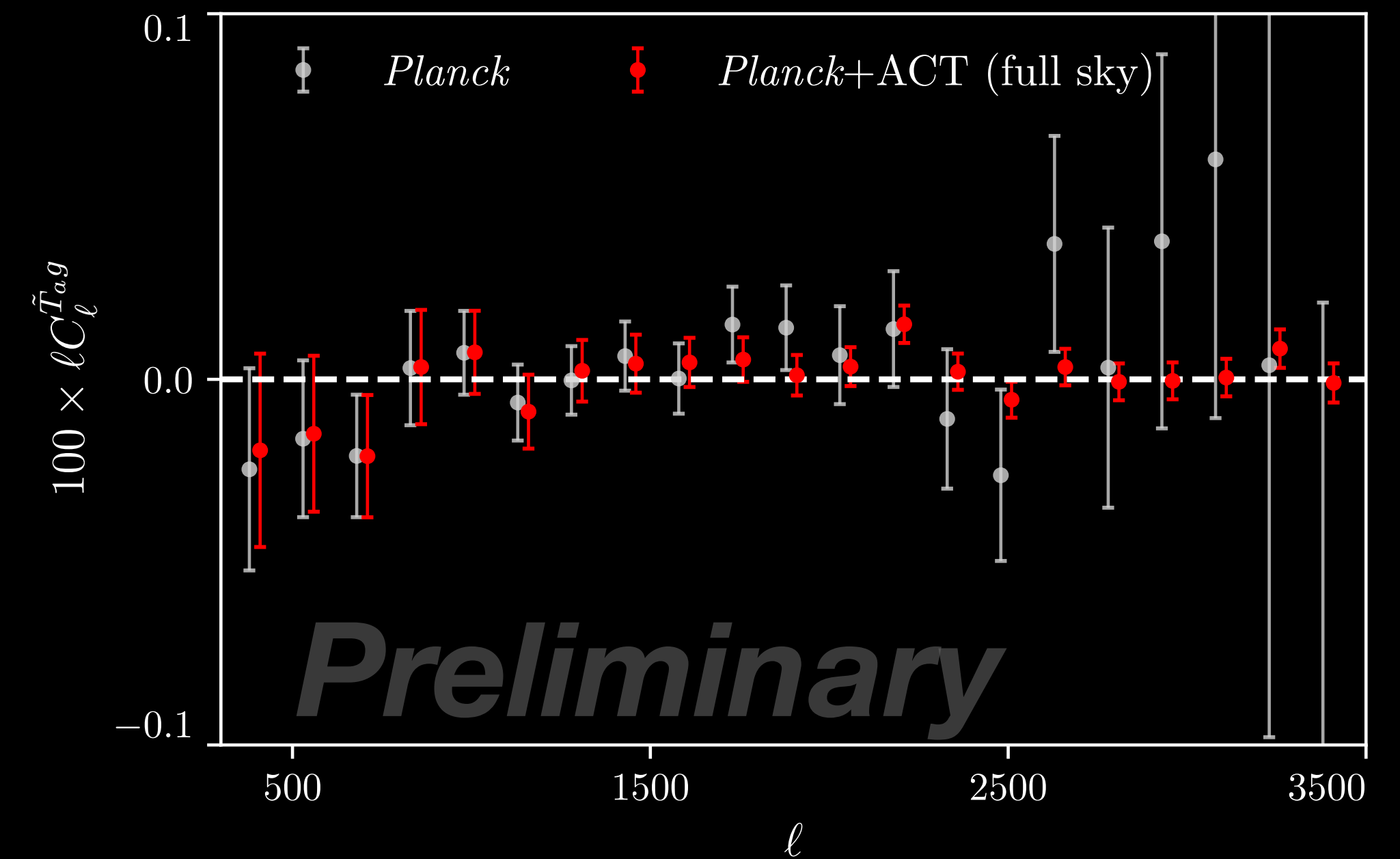
Future outlook: new tracers



- Cross-correlate with ACT DR6 CMB lensing map ([Madhavacheril++23](#))
- Sensitive to lower mass axions and dark photons

Conclusion

- CMB secondaries are a powerful probe of the dark sector
- No evidence for axion-induced screening in *unWISE* × *Planck* + ACT
- Constrains $g_{a\gamma\gamma} \lesssim 8 \times 10^{-13} \text{ GeV}^{-1}$ at $m_a \sim 3 \times 10^{-13} \text{ eV}$
- ***Along the way:*** public NILC pipeline for Simons Observatory, new foreground-mitigation techniques, and high-resolution *Planck* + ACT NILC maps

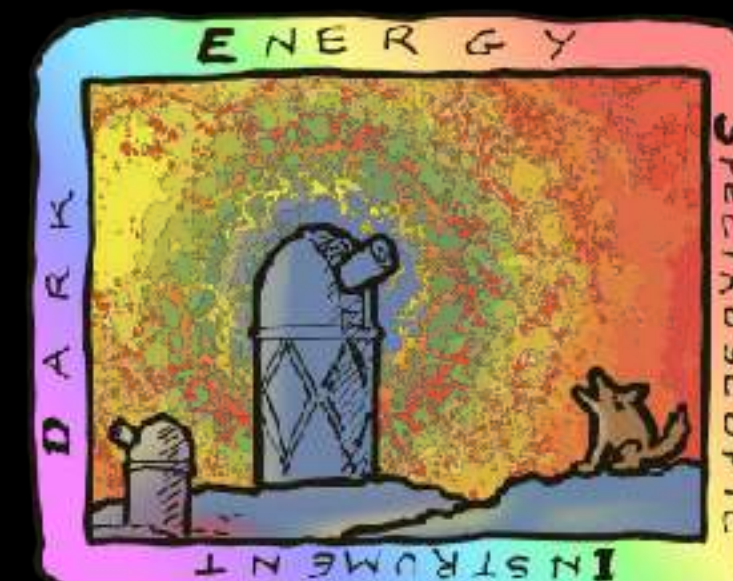


Part ii: uncovering the physics of inflation with galaxy surveys



U.S. DEPARTMENT OF
ENERGY

Office of Science
SCGSR Program




BERKELEY LAB



Edmond
Chaussidon



Simone
Ferraro

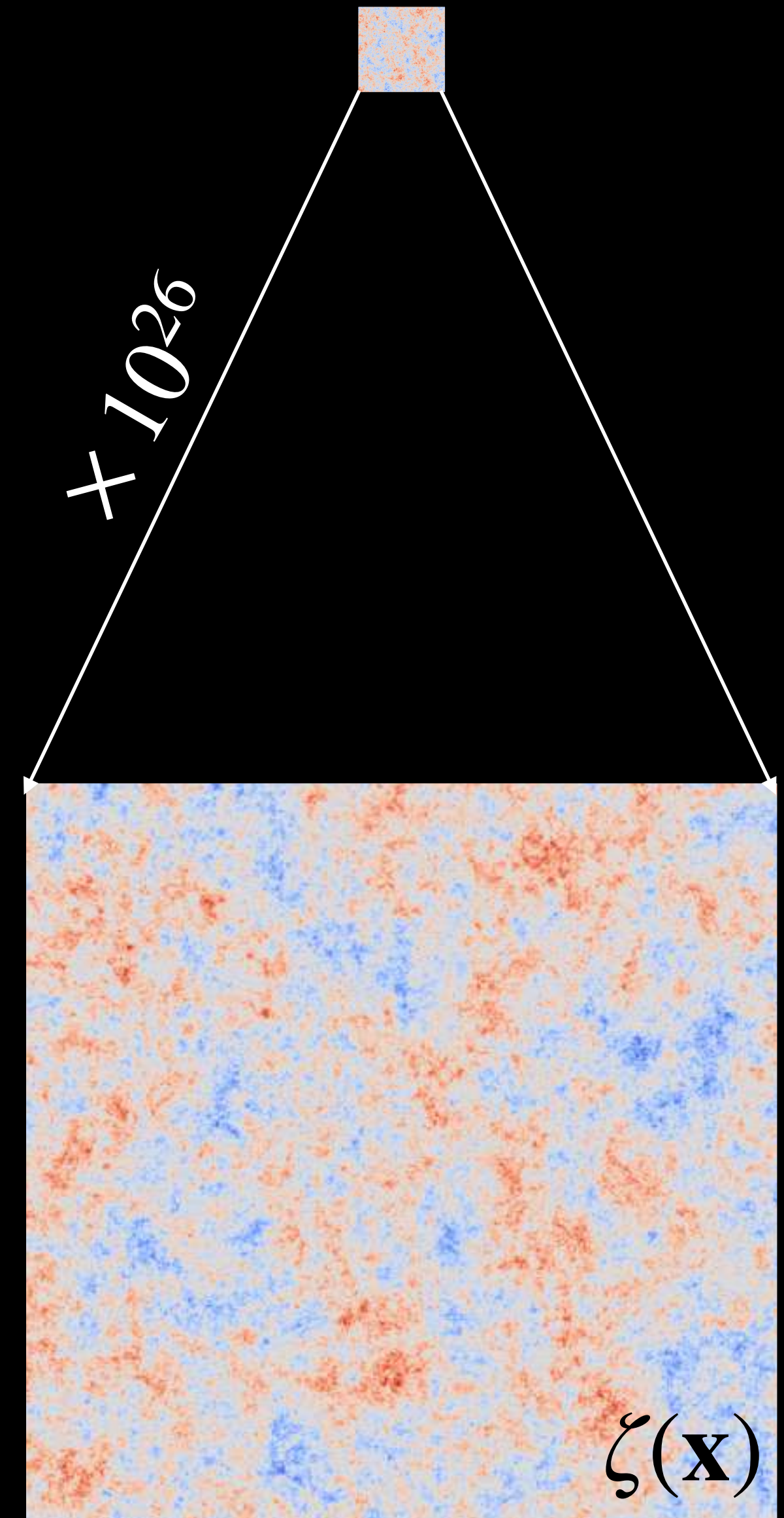


Noah
Sailer

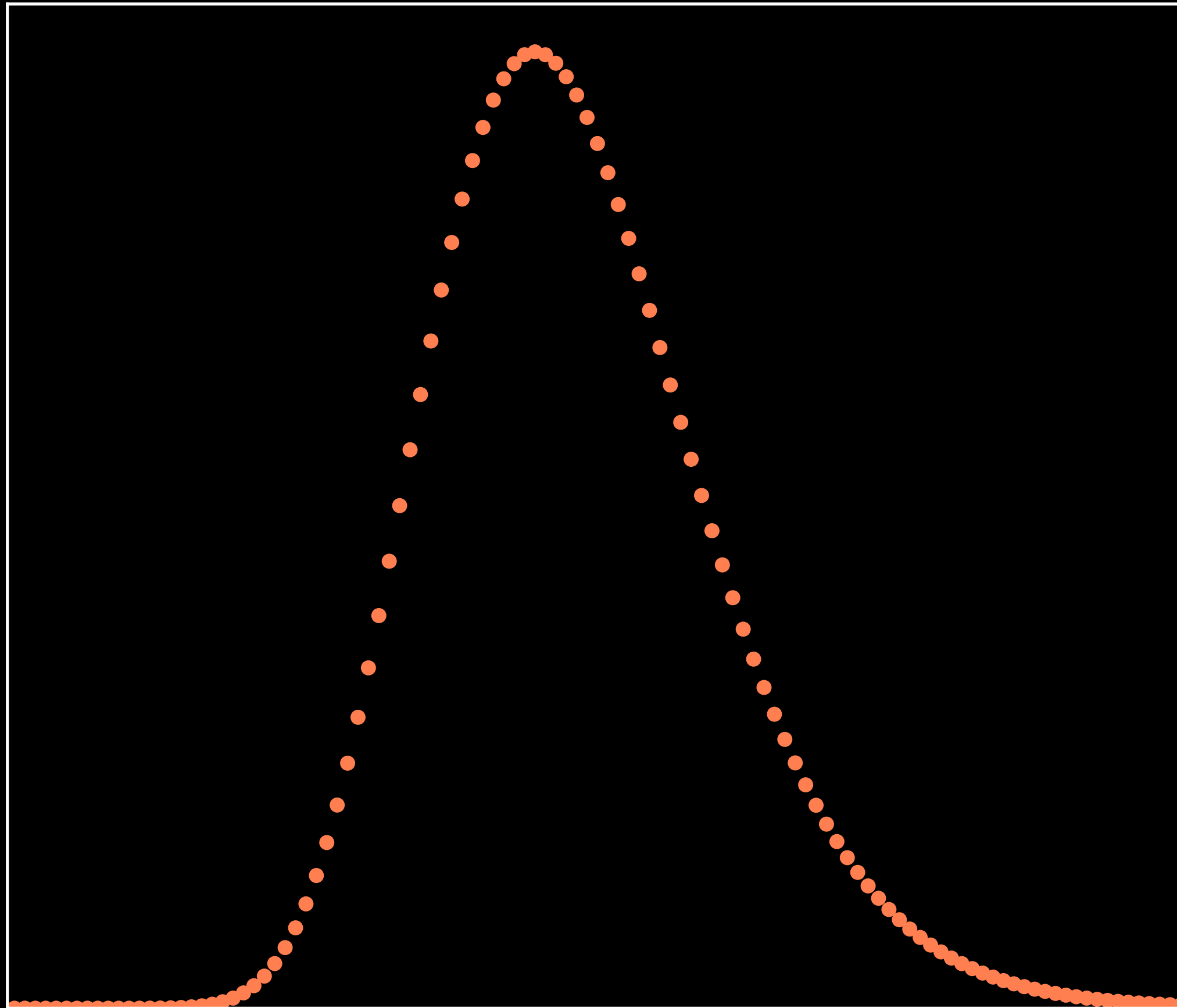
Cosmic Inflation

- Period of accelerated expansion in early Universe ($t \approx 10^{-34}$ seconds)
- **Quantum fluctuations seed structure**
Stats. of fluctuations $\leftrightarrow$ inflationary physics
 $\langle \zeta \dots \zeta \rangle$ $\mathcal{L}_{\text{inf}}(\phi, \chi, \psi_S)$
- Simplest models: **Gaussian** perturbations, $\zeta(\mathbf{x})$

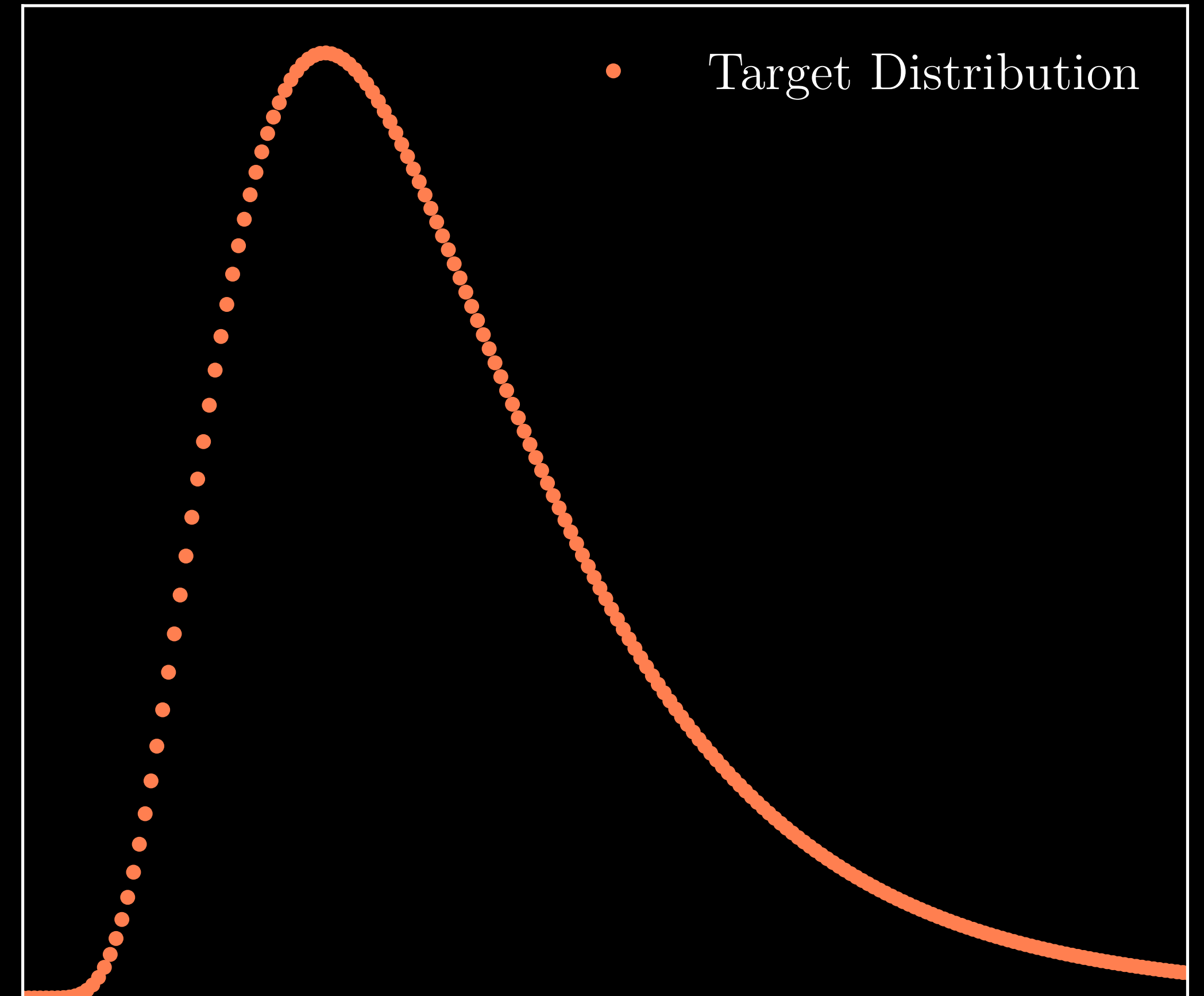
What about primordial non-Gaussianity (PNG)?



1). Weak non-Gaussianity

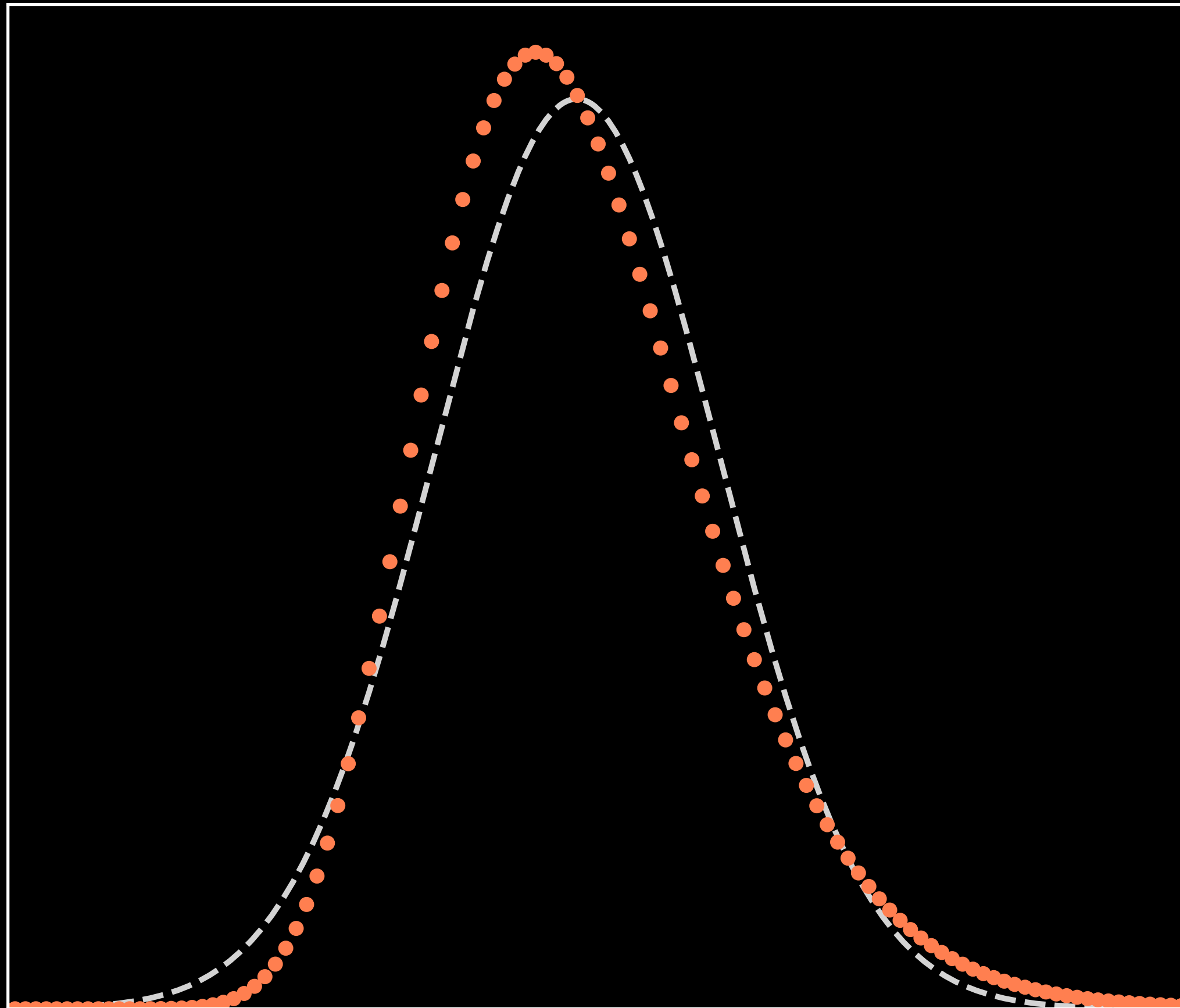


2). Strong non-Gaussianity

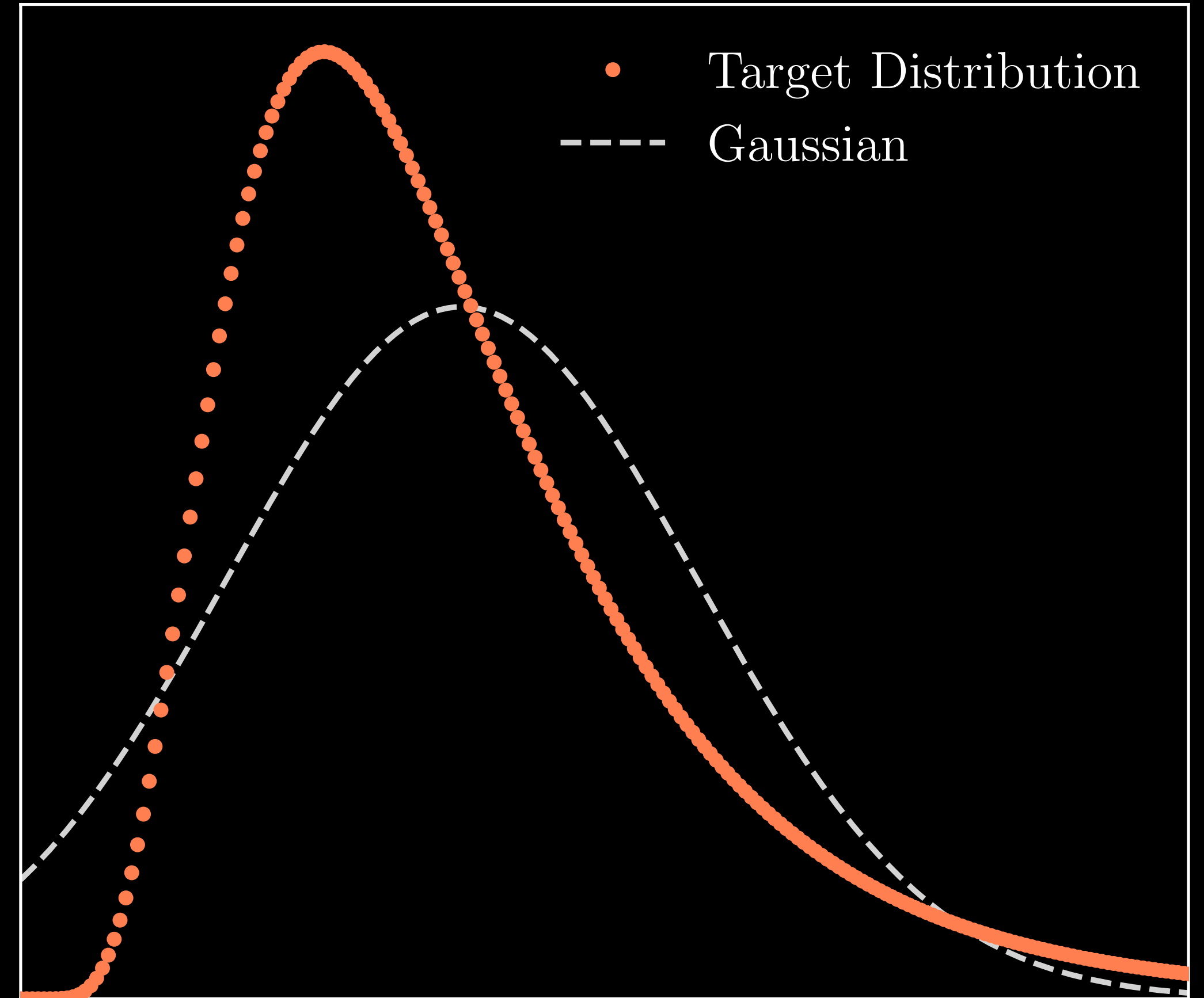


The two regimes of non-Gaussianity

1). Weak non-Gaussianity



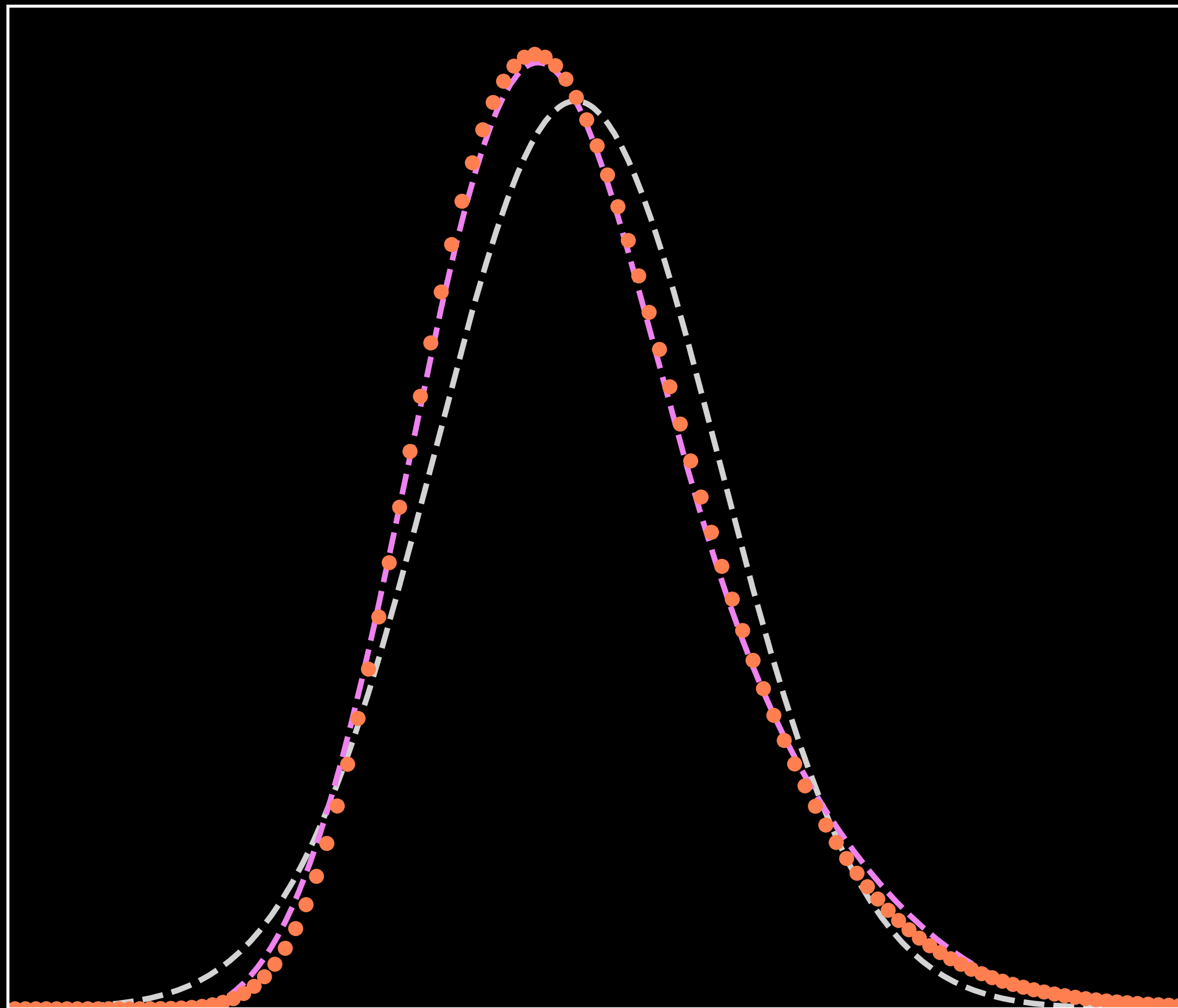
2). Strong non-Gaussianity



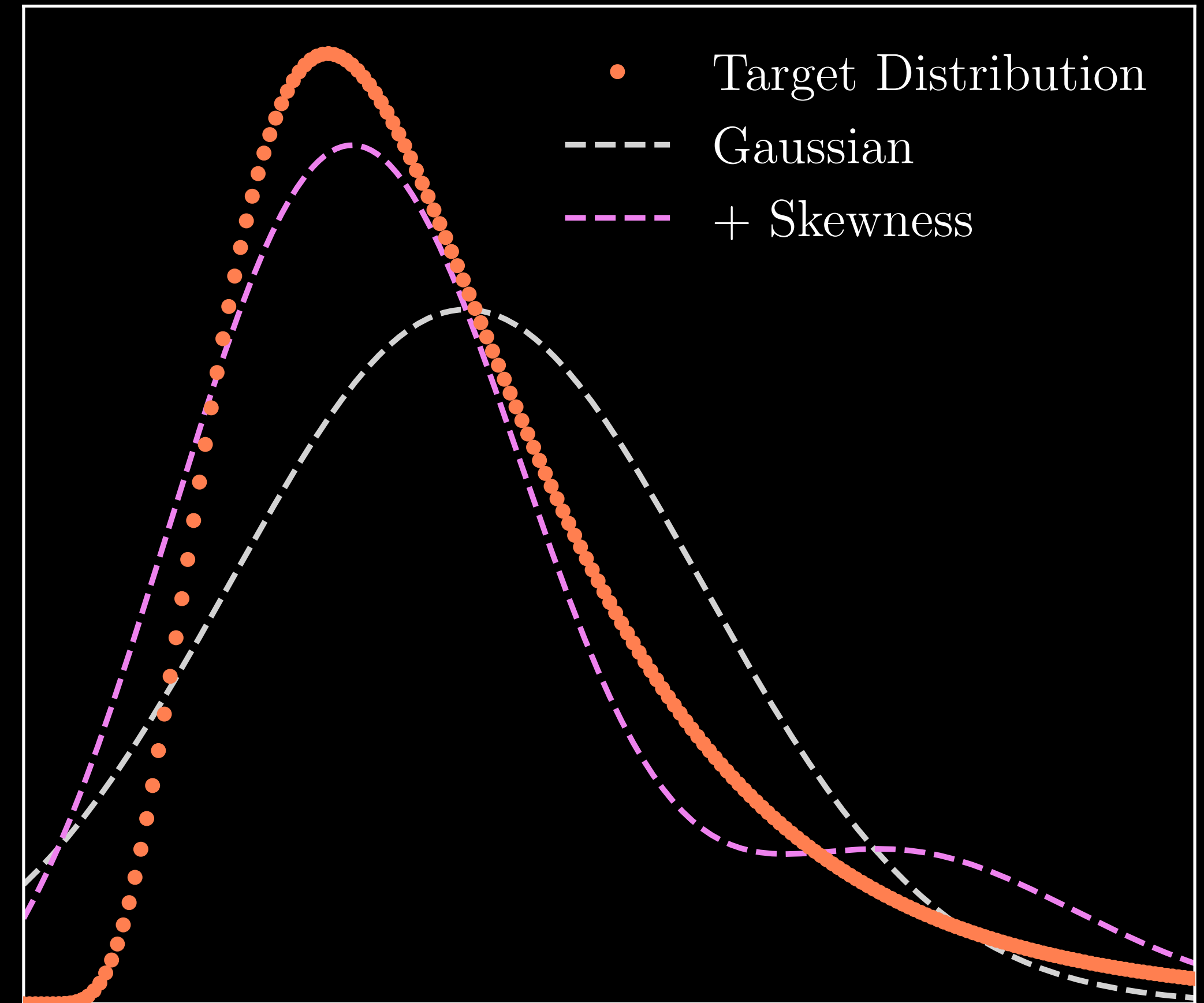
Start with a Gaussian

$$P(x) \approx \phi \left(\tilde{x} \equiv \frac{x - \mu}{\sigma} \right)$$

1). Weak non-Gaussianity



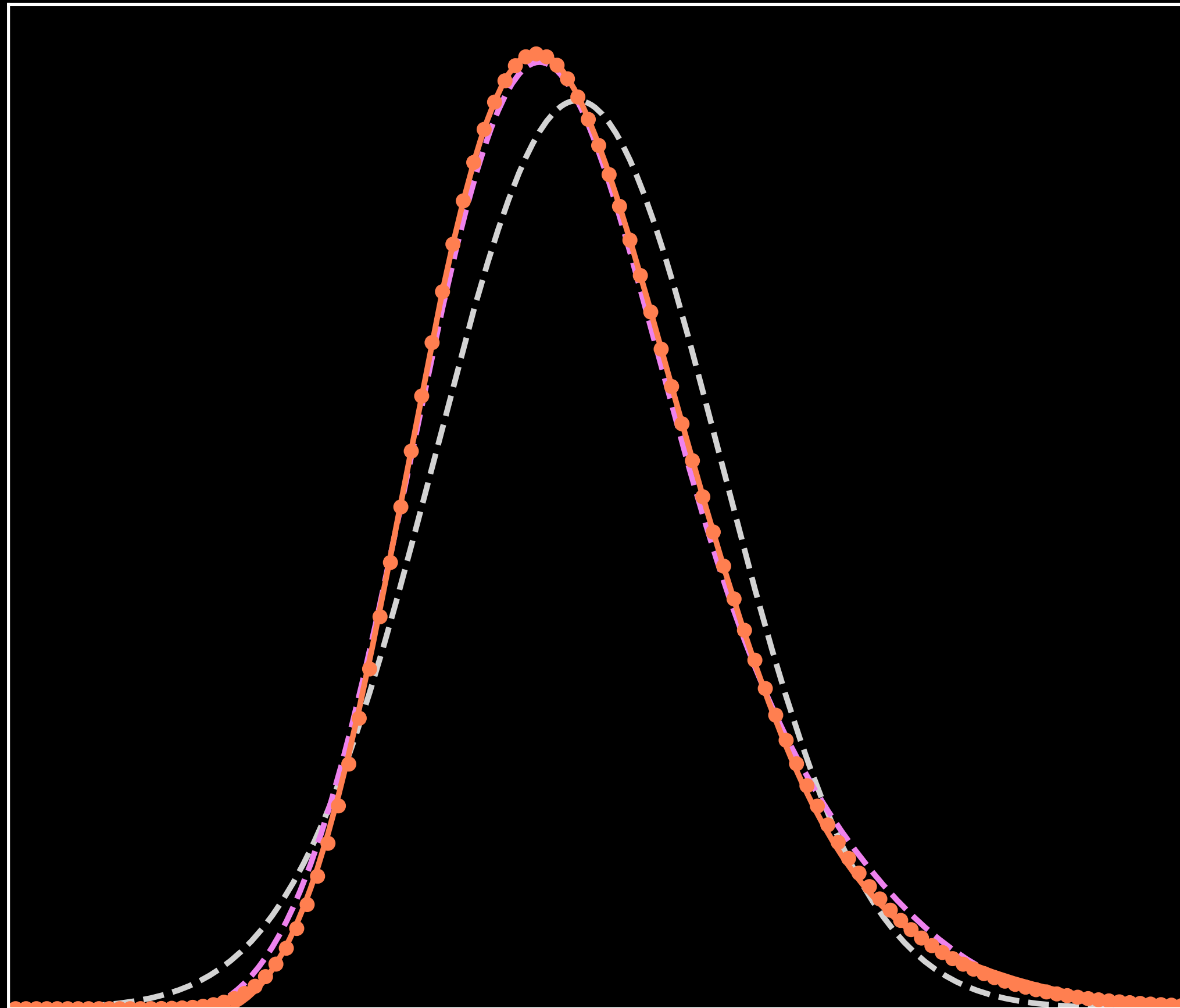
2). Strong non-Gaussianity



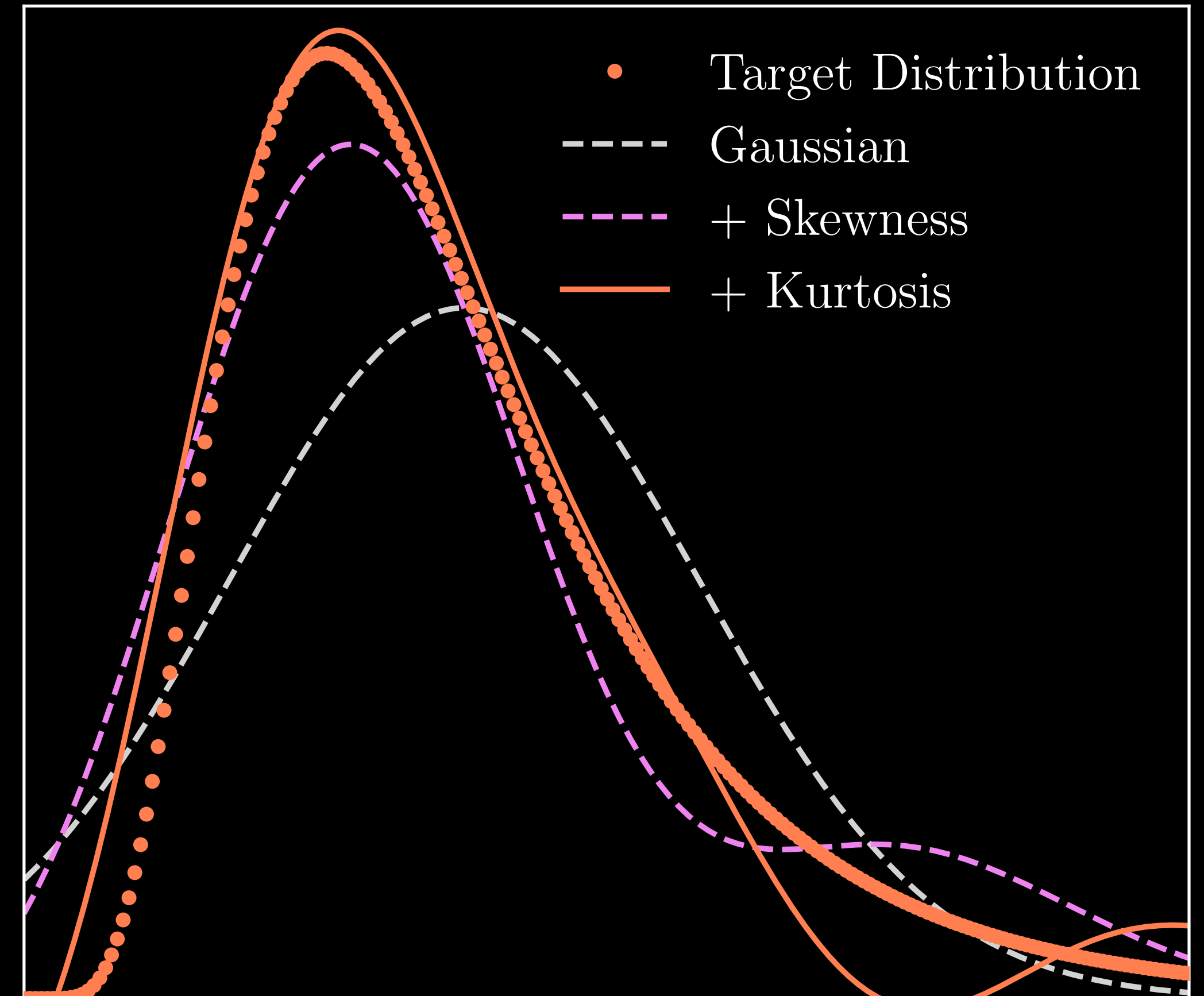
Add in skewness

$$P(x) \approx \phi(\tilde{x}) \left[1 + \frac{\gamma_1}{6} \text{He}_3(\tilde{x}) \right] \quad \text{“Edgeworth expansion”}$$

1). Weak non-Gaussianity



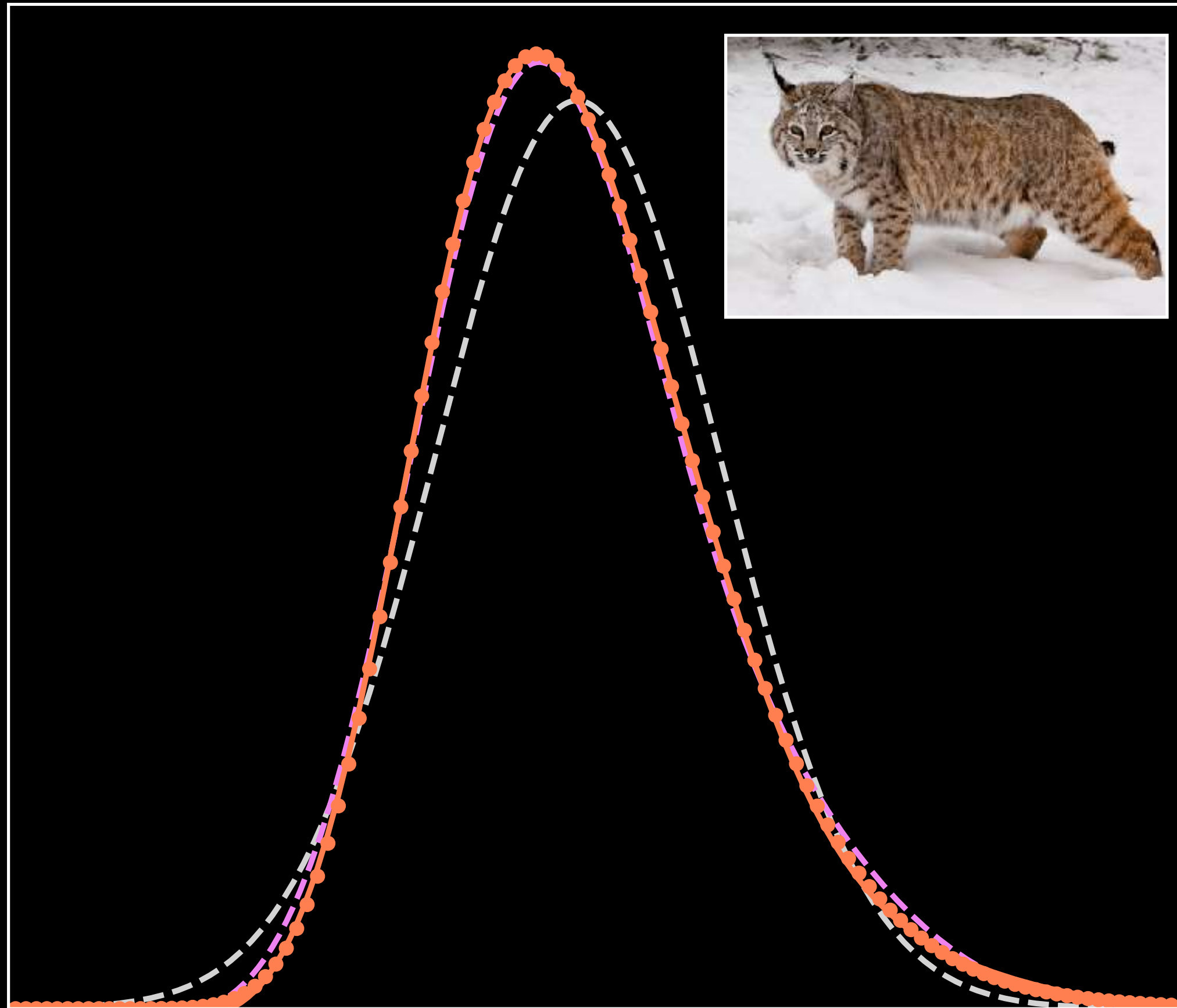
2). Strong non-Gaussianity



Add in kurtosis

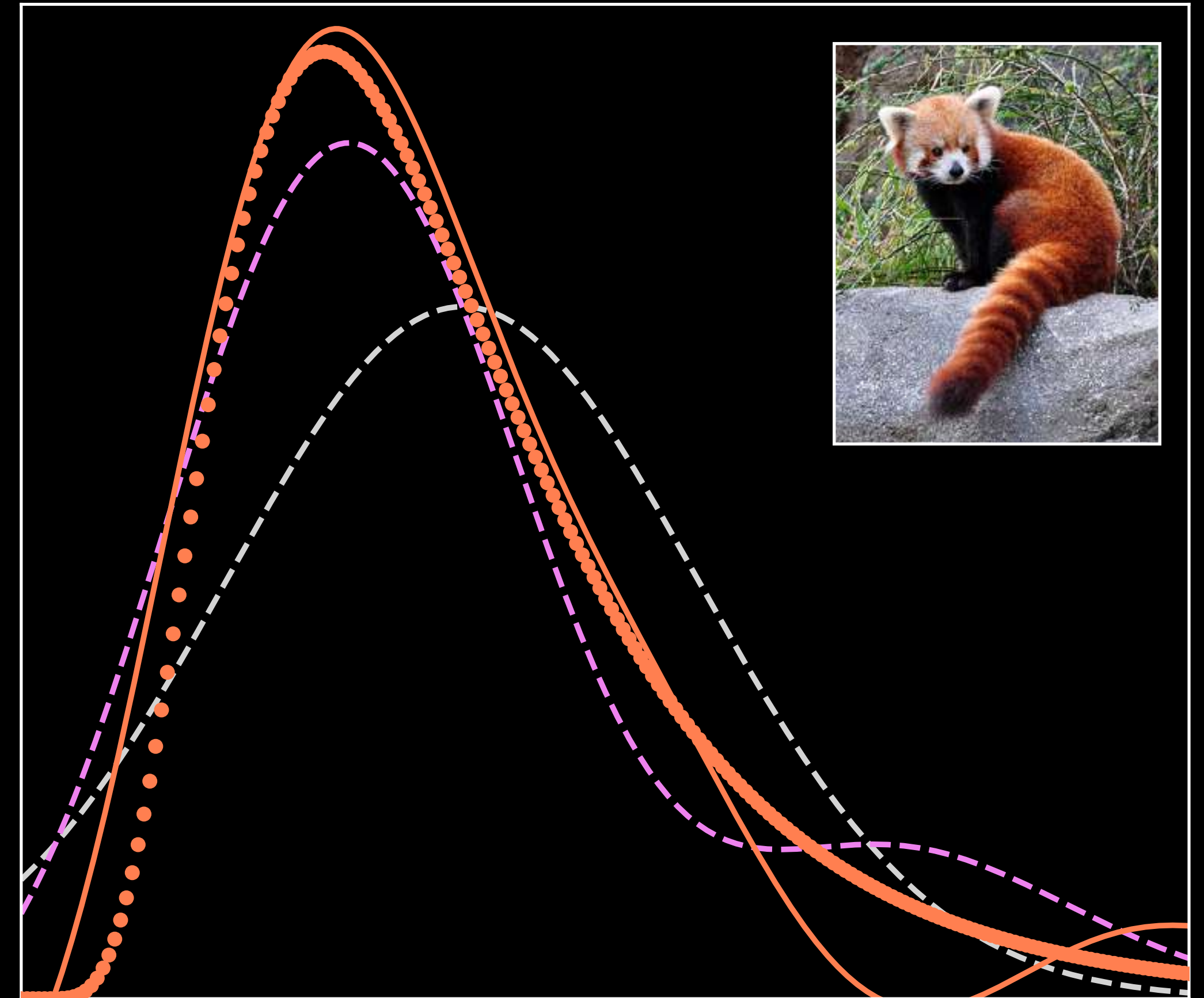
$$P(x) \approx \phi(\tilde{x}) \left[1 + \frac{\gamma_1}{6} \text{He}_3(\tilde{x}) + \frac{\gamma_2}{24} \text{He}_4(\tilde{x}) + \frac{\gamma_1^2}{72} \text{He}_6(\tilde{x}) \right] \quad \text{“Edgeworth expansion”}$$

1). Weak non-Gaussianity



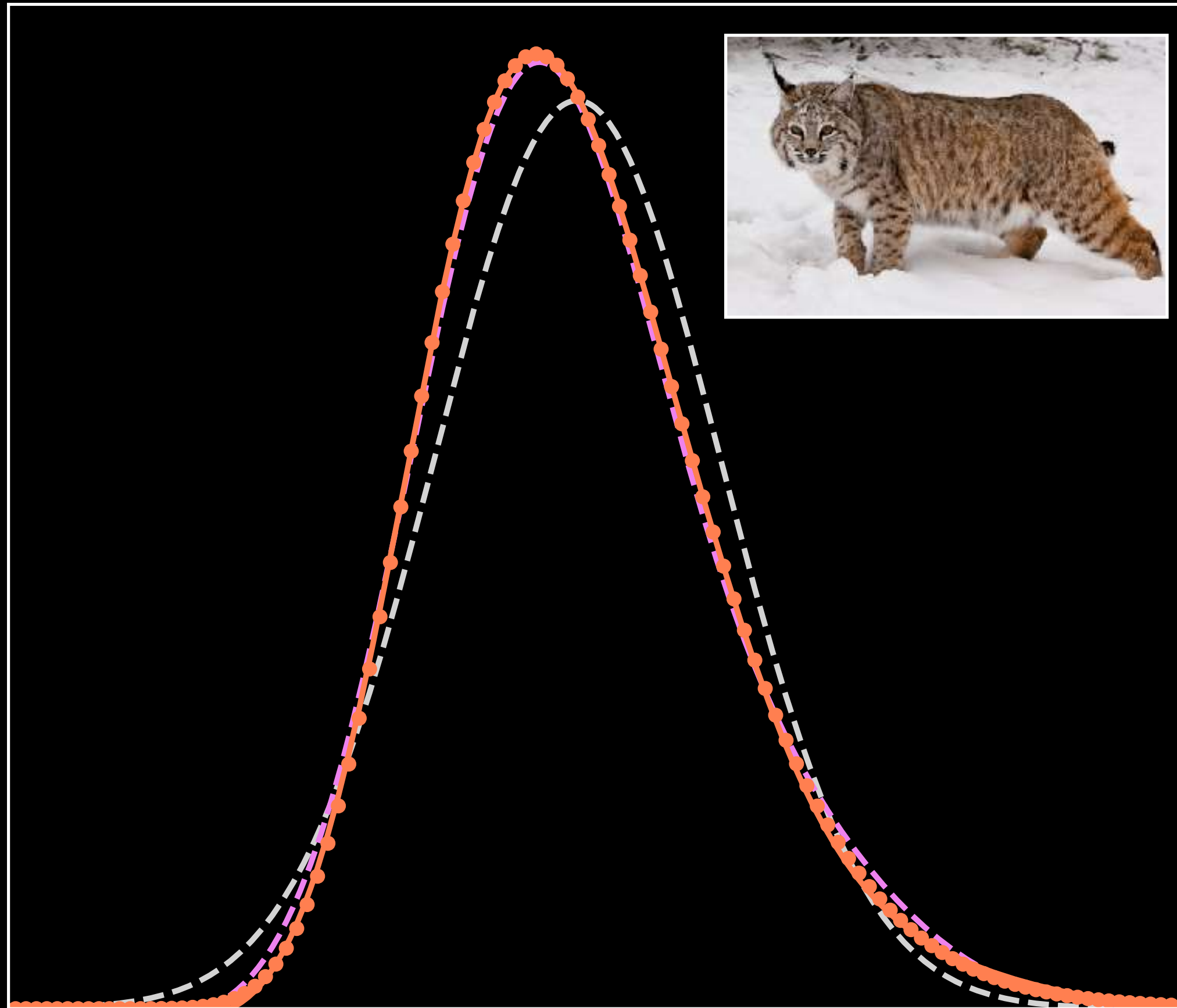
- **Optimal summary statistic:** N-point correlation functions with $N \sim 3, 4, \dots$

2). Strong non-Gaussianity

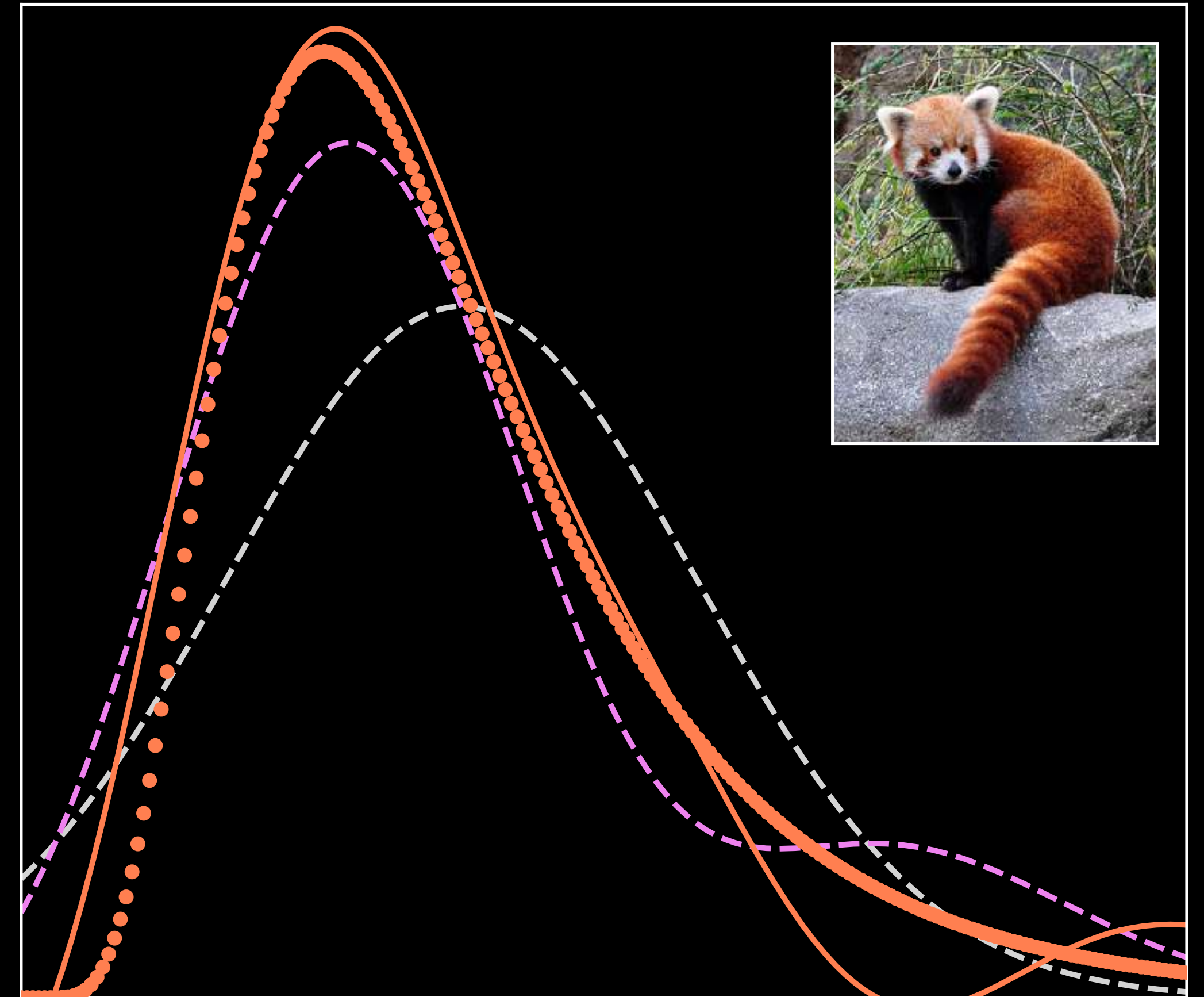


- **Optimal summary statistic:** matched filter, ask a neural network

1). Weak non-Gaussianity



2). Strong non-Gaussianity



This talk

- **Optimal summary statistic:** N-point correlation functions with $N \sim 3, 4, \dots$

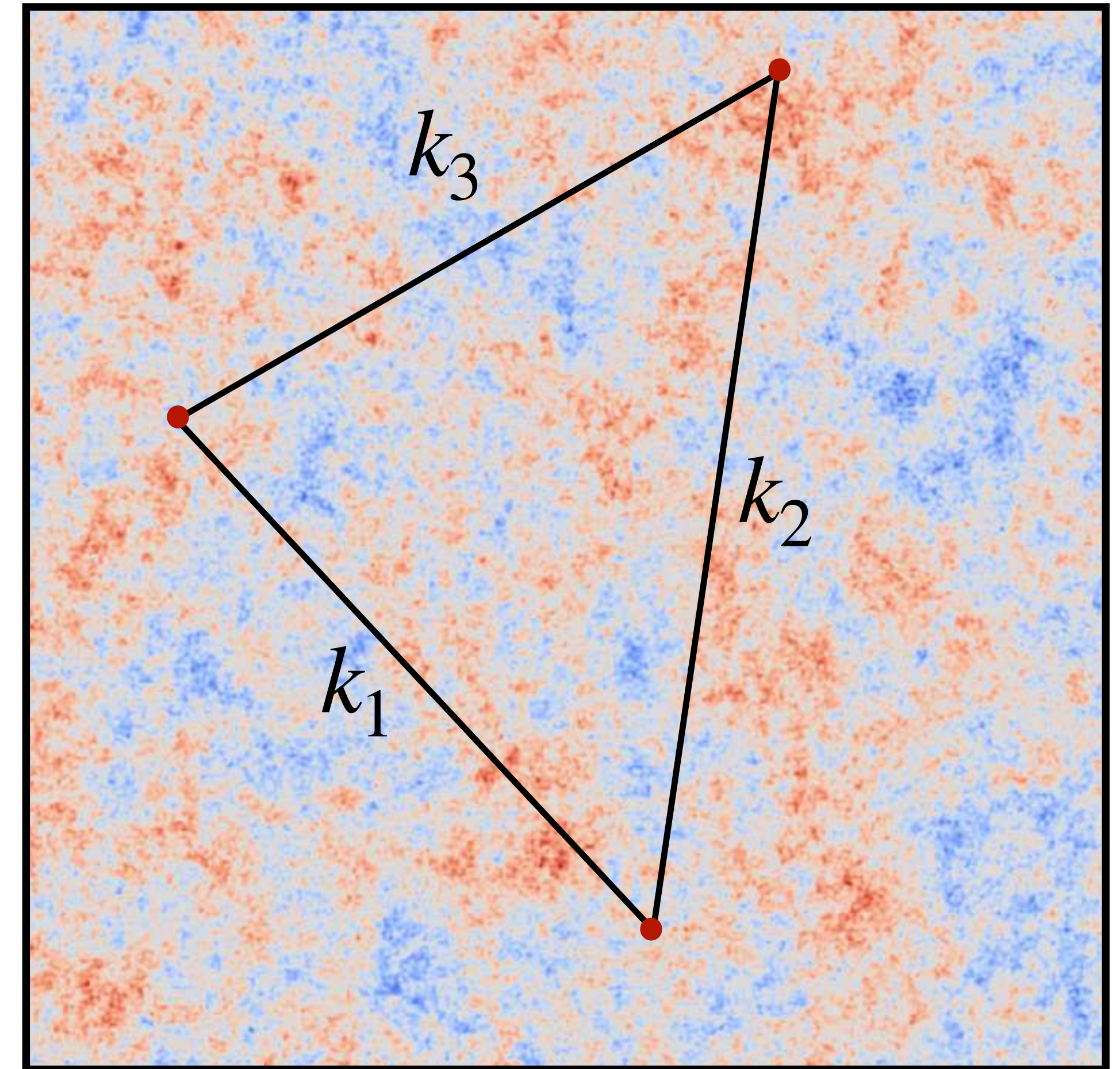
- **Optimal summary statistic:** matched filter, ask a neural network

The 3-point function (bispectrum)

- Natural higher-order statistic to consider

$$\langle \zeta(\mathbf{k}_1) \zeta(\mathbf{k}_2) \zeta(\mathbf{k}_3) \rangle \equiv (2\pi)^3 \delta_D(\mathbf{k}_1 + \mathbf{k}_2 + \mathbf{k}_3) B_\zeta(k_1, k_2, k_3)$$

- Fully described by three side lengths of the triangle
- Is this an *interesting/useful* statistic...



The 3-point function (bispectrum)

- Natural higher-order statistic to consider

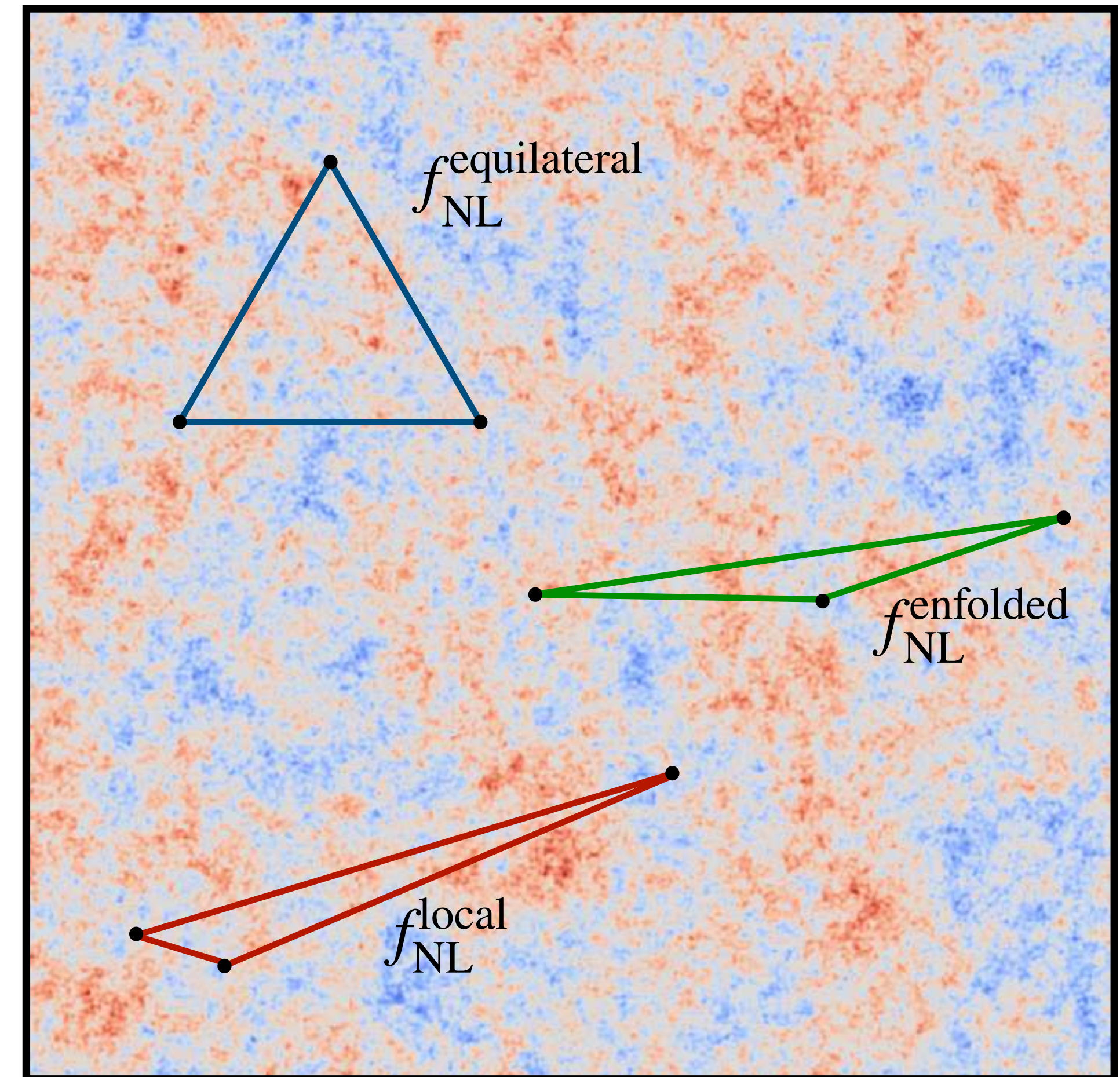
$$\langle \zeta(\mathbf{k}_1) \zeta(\mathbf{k}_2) \zeta(\mathbf{k}_3) \rangle \equiv (2\pi)^3 \delta_D(\mathbf{k}_1 + \mathbf{k}_2 + \mathbf{k}_3) B_\zeta(k_1, k_2, k_3)$$

- Fully described by three side lengths of the triangle
- Is this an *interesting/useful* statistic... **YES**
- Triangle shapes probe various early universe models!

$$B_\zeta(k_1, k_2, k_3) \sim f_{\text{NL}} \times S(k_1, k_2, k_3)$$

Amplitude (zero for GRF)

Shape



The 3-point function (bispectrum)

- Natural higher-order statistic to consider

$$\langle \zeta(\mathbf{k}_1) \zeta(\mathbf{k}_2) \zeta(\mathbf{k}_3) \rangle \equiv (2\pi)^3 \delta_D(\mathbf{k}_1 + \mathbf{k}_2 + \mathbf{k}_3) B_\zeta(k_1, k_2, k_3)$$

- Fully described by three side lengths of the triangle

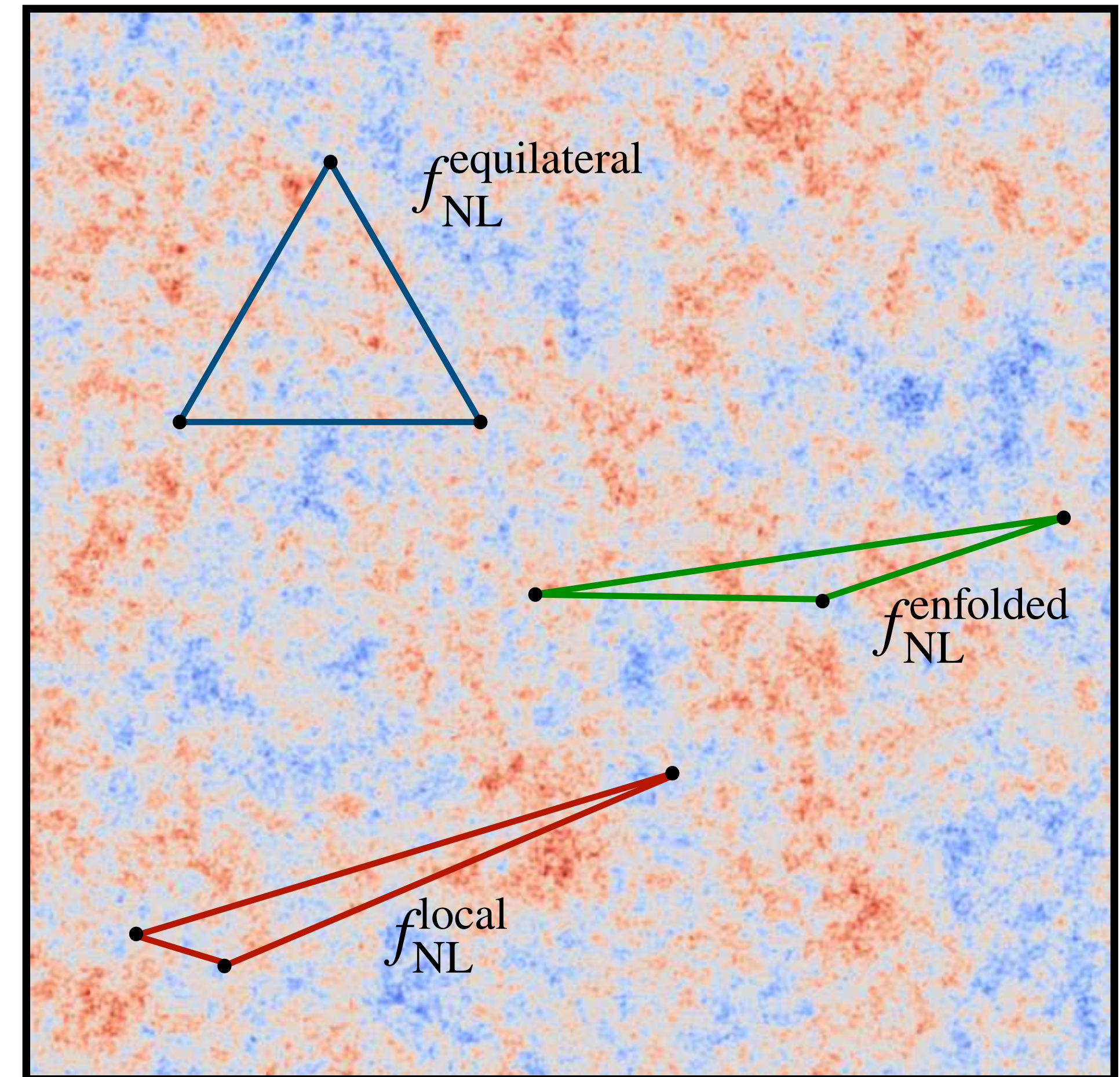
- Is this an *interesting/useful* statistic... **YES**

- $f_{\text{NL}}^{\text{eq.}}$: Details of particle sourcing inflation

- $f_{\text{NL}}^{\text{enf.}}$: Initial quantum state of inflation

- $f_{\text{NL}}^{\text{loc.}}$: **Multiple fields**

$$\Phi(\mathbf{x}) = \phi_G(\mathbf{x}) + f_{\text{NL}}^{\text{loc.}} (\phi_G^2(\mathbf{x}) - \langle \phi^2 \rangle)$$



The 3-point function (bispectrum)

- Natural higher-order statistic to consider

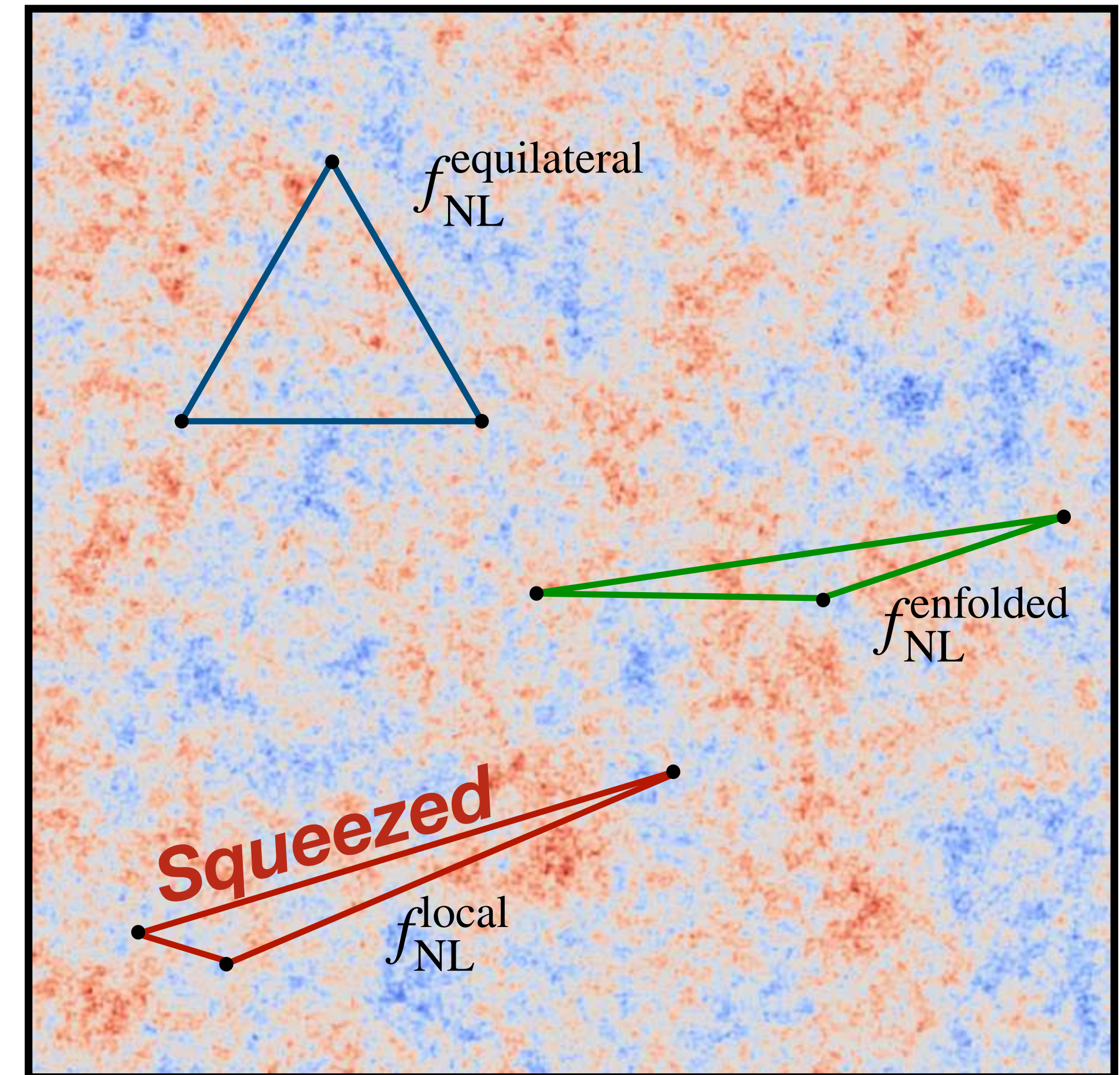
$$\langle \zeta(\mathbf{k}_1) \zeta(\mathbf{k}_2) \zeta(\mathbf{k}_3) \rangle \equiv (2\pi)^3 \delta_D(\mathbf{k}_1 + \mathbf{k}_2 + \mathbf{k}_3) B_\zeta(k_1, k_2, k_3)$$

- Fully described by three side lengths of the triangle

- Is this an *interesting/useful* statistic... **YES**

- $f_{\text{NL}}^{\text{eq.}}$: Details of particle sourcing inflation
- $f_{\text{NL}}^{\text{enf.}}$: Initial quantum state of inflation
- $f_{\text{NL}}^{\text{loc.}}$: **Multiple fields**

$f_{\text{NL}}^{\text{loc.}} \neq 0$ would rule out single-field inflation



The 4-point function (trispectrum)

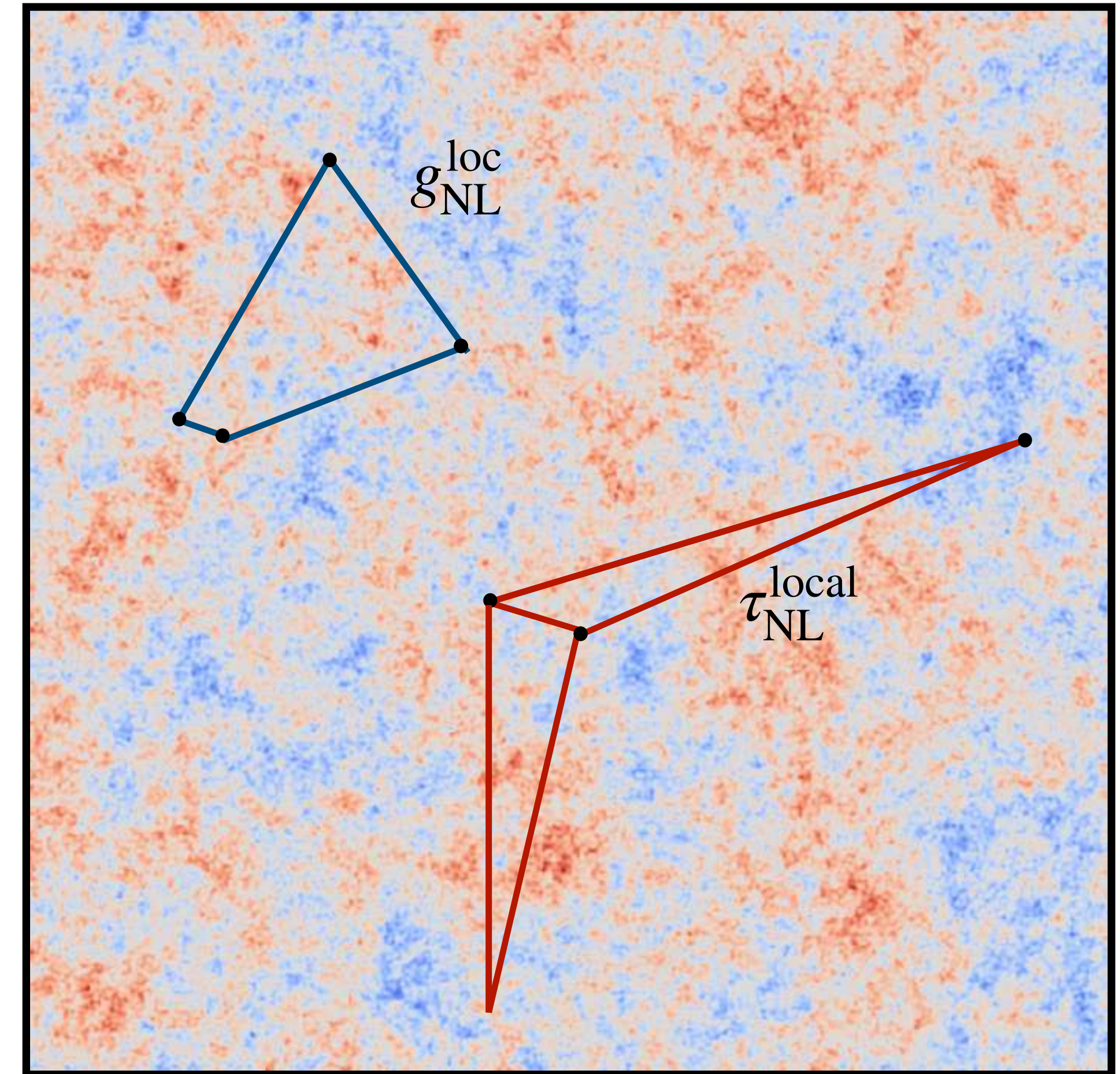
- Natural higher-order statistic to consider

$$\langle \zeta(\mathbf{k}_1)\zeta(\mathbf{k}_2)\zeta(\mathbf{k}_3)\zeta(\mathbf{k}_4) \rangle \equiv (2\pi)^3 \delta_D(\mathbf{k}_1 + \mathbf{k}_2 + \mathbf{k}_3 + \mathbf{k}_4) T_\zeta(k_1, k_2, k_3, k_4, k_{12}, k_{14})$$

- Fully described by four side lengths and two diagonals

- Is this an *interesting/useful* statistic... **YES**

- Parity violation
- $g_{\text{NL}}^{\text{loc}}$: cubic non-Gaussianity
- $\tau_{\text{NL}}^{\text{loc}}$: direct probe of # of fields responsible for ζ



The 4-point function (trispectrum)

- Natural higher-order statistic to consider

$$\langle \zeta(\mathbf{k}_1)\zeta(\mathbf{k}_2)\zeta(\mathbf{k}_3)\zeta(\mathbf{k}_4) \rangle \equiv (2\pi)^3 \delta_D(\mathbf{k}_1 + \mathbf{k}_2 + \mathbf{k}_3 + \mathbf{k}_4) T_\zeta(k_1, k_2, k_3, k_4, k_{12}, k_{14})$$

- Fully described by four side lengths and two diagonals

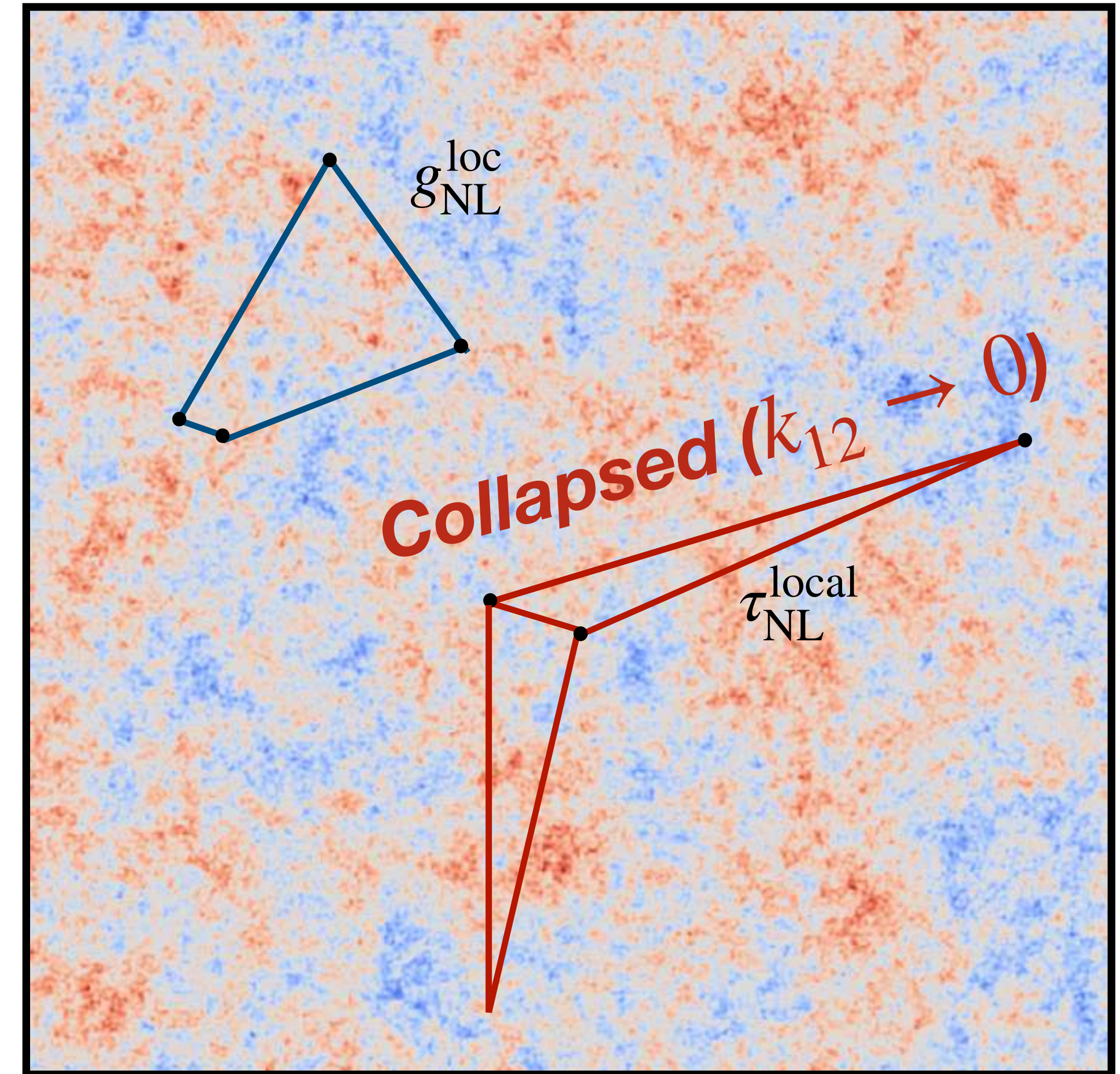
- Is this an *interesting/useful* statistic... **YES**

- Parity violation (e.g., [Cahn++21](#))
- $g_{\text{NL}}^{\text{loc}}$: cubic non-Gaussianity
- $\tau_{\text{NL}}^{\text{loc}}$: direct probe of # of fields responsible for ζ

Suyama-Yamaguchi (SY) inequality

$$\tau_{\text{NL}}^{\text{loc}} > \left(\frac{6}{5} f_{\text{NL}}^{\text{loc}} \right)^2 \text{ would rule out single-source inflation}$$

[Byrnes++2006](#), [Suyama & Yamaguchi, 2008](#), [Smith++ 2010](#)

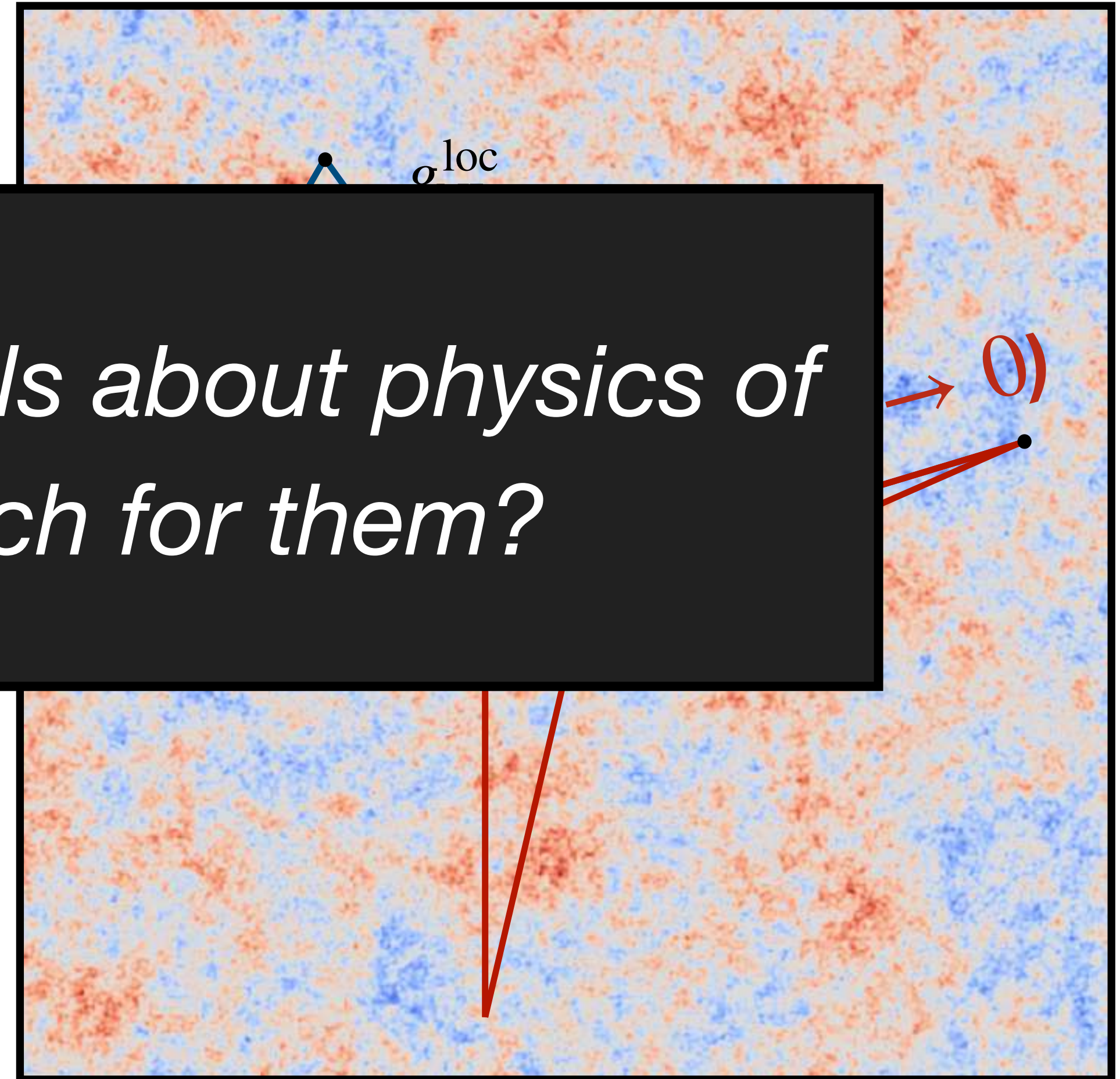


The 4-point function (trispectrum)

- Natural higher-order statistic to consider

$$\langle \zeta(\mathbf{k}_1)\zeta(\mathbf{k}_2)\zeta(\mathbf{k}_3)\zeta(\mathbf{k}_4) \rangle \equiv (2\pi)^3 \delta_D(\mathbf{k}_1 + \mathbf{k}_2 + \mathbf{k}_3 + \mathbf{k}_4) T_\zeta(k_1, k_2, k_3, k_4, k_{12}, k_{14})$$

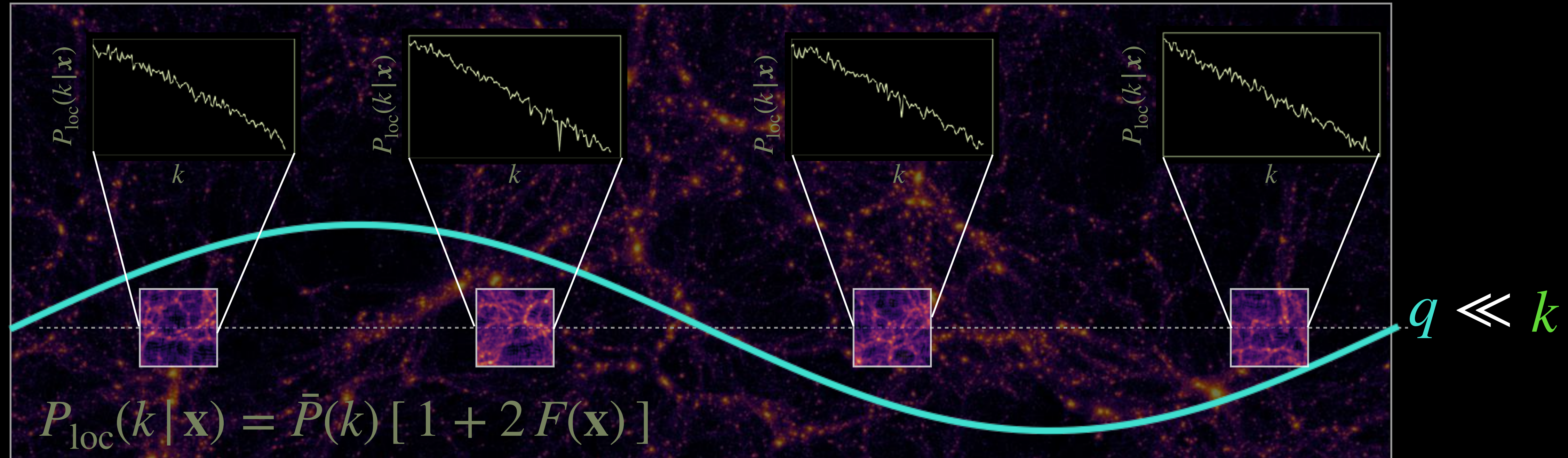
- Fully determined by $f_{\text{NL}}^{\text{loc}}$ and $\tau_{\text{NL}}^{\text{loc}}$ encode crucial details about physics of inflation. How can we search for them?
- Is this consistent with single-field inflation?
- Parity-odd
- $g_{\text{NL}}^{\text{loc}}$: cubic non-Gaussianity
- $\tau_{\text{NL}}^{\text{loc}}$: direct probe of # of fields responsible for ζ



Suyama-Yamaguchi (SY) inequality

$$\tau_{\text{NL}}^{\text{loc}} > \left(\frac{6}{5} f_{\text{NL}}^{\text{loc}} \right)^2 \text{ would rule out single-source inflation}$$

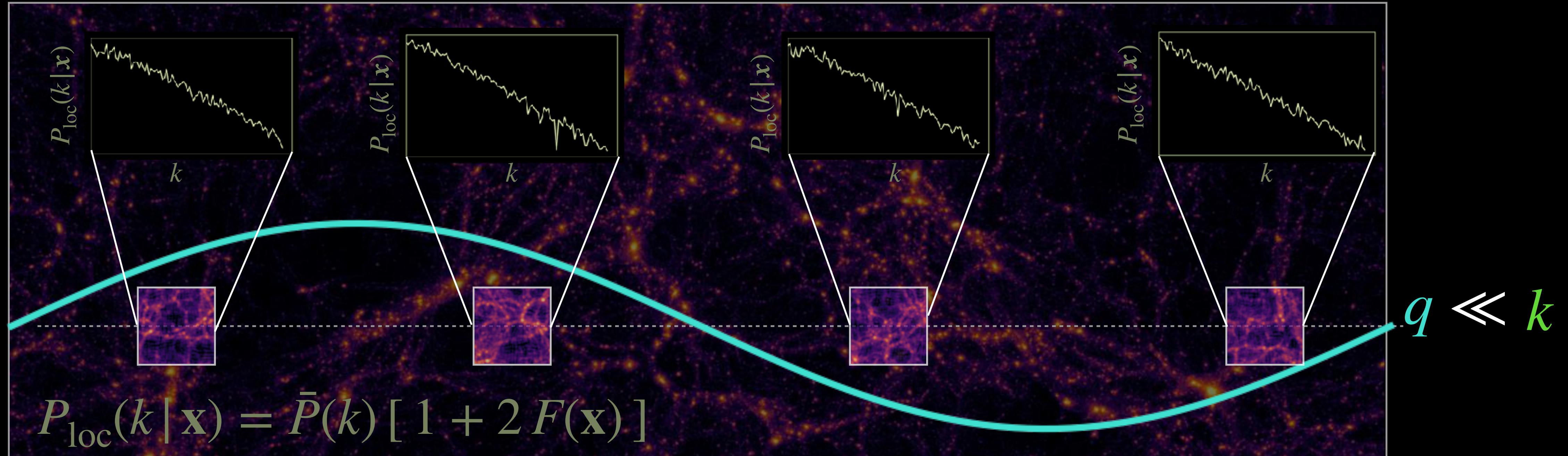
What are $f_{\text{NL}}^{\text{loc}}$ and $\tau_{\text{NL}}^{\text{loc}}$ **REALLY**?



“Squeezed/soft limits $\rightarrow$ long-wavelength modulations of small-scale (local) power spectra” — Dr. Soft Limit

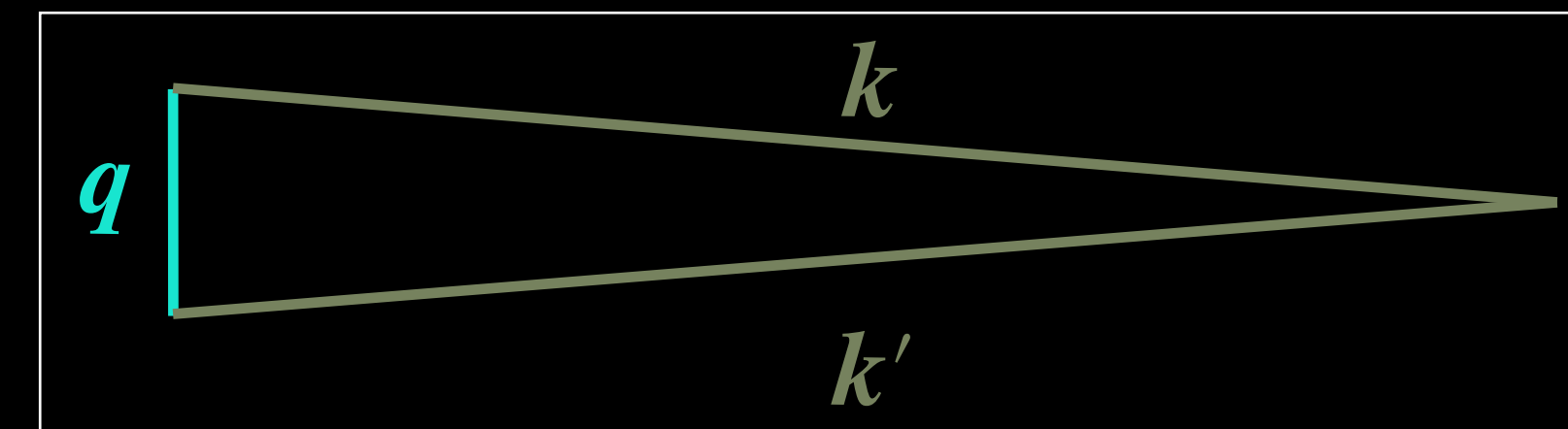


What are $f_{\text{NL}}^{\text{loc}}$ and $\tau_{\text{NL}}^{\text{loc}}$ **REALLY**?



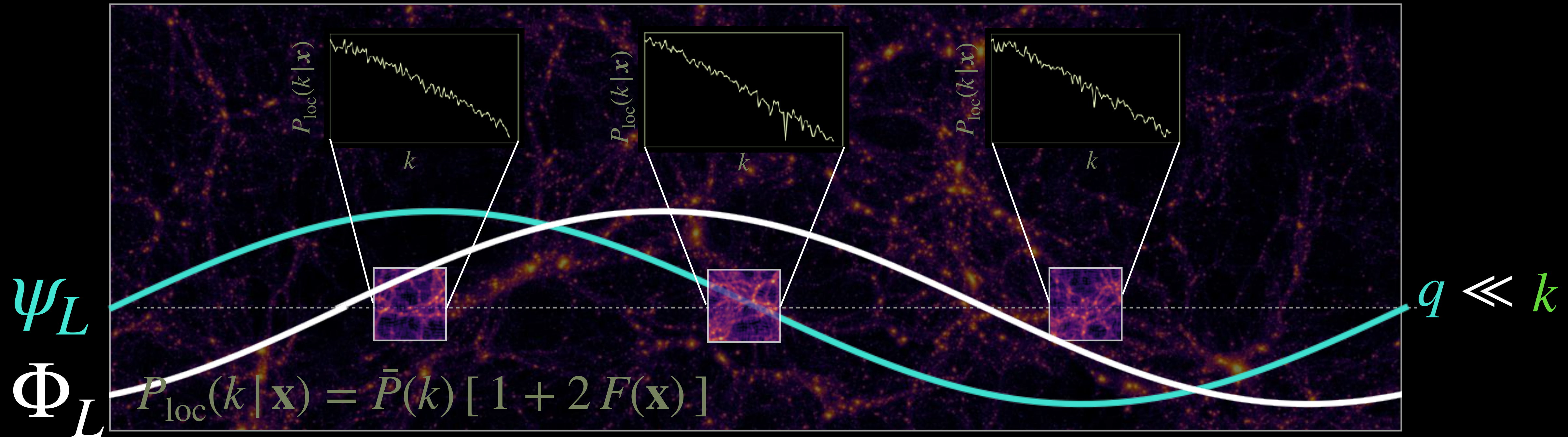
$$f_{\text{NL}}^{\text{loc}}: \Phi \equiv \Phi + f_{\text{NL}}^{\text{loc}} \Phi^2 \implies \Phi_S = \phi_S (1 + 2 f_{\text{NL}}^{\text{loc}} \Phi_L)$$

$$F(\mathbf{x}) \propto f_{\text{NL}}^{\text{loc}} \Phi_L(\mathbf{x})$$



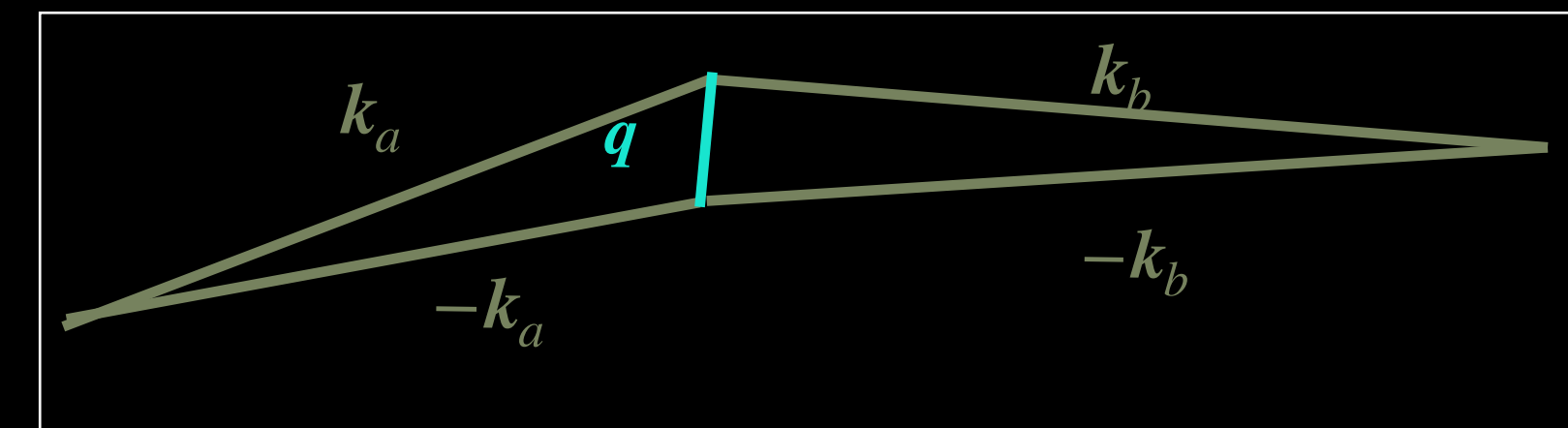
Squeezed bispectrum $\sim \langle \Phi_L F \rangle$

What are $f_{\text{NL}}^{\text{loc}}$ and $\tau_{\text{NL}}^{\text{loc}}$ **REALLY**?



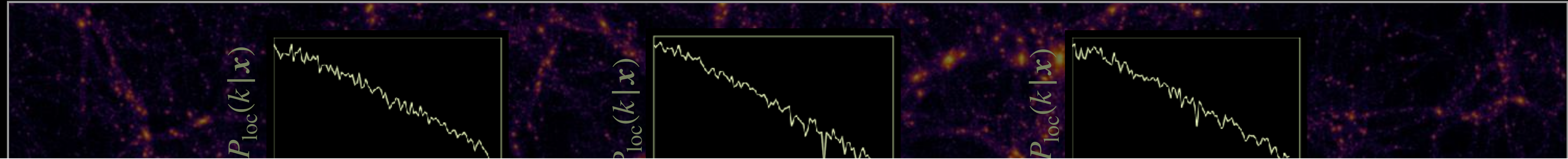
$\tau_{\text{NL}}^{\text{loc}}$: power spectrum of the modulating field

$$F(\mathbf{x}) = f_{\text{NL}}^{\text{loc}} \Phi_L(\mathbf{x}) + \sqrt{\tau_{\text{NL}}^{\text{loc}} - \left(\frac{6}{5} f_{\text{NL}}^{\text{loc}}\right)^2} \psi_L(\mathbf{x}),$$



Collapsed trispectrum $\sim \langle F F \rangle$

What are $f_{\text{NL}}^{\text{loc}}$ and $\tau_{\text{NL}}^{\text{loc}}$ **REALLY**?

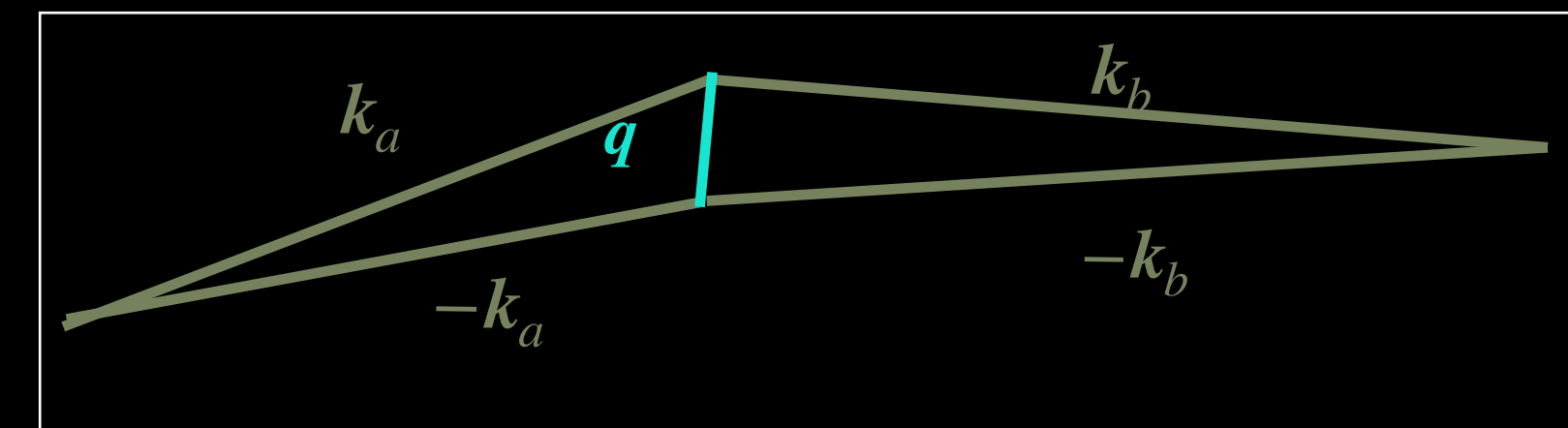


This has significant implications for the large-scale galaxy distribution (bias)

$q \ll k$

$\tau_{\text{NL}}^{\text{loc}}$: power spectrum of the modulating field

$$F(\mathbf{x}) = f_{\text{NL}}^{\text{loc}} \Phi_L(\mathbf{x}) + \sqrt{\tau_{\text{NL}}^{\text{loc}} - \left(\frac{6}{5} f_{\text{NL}}^{\text{loc}}\right)^2} \psi_L(\mathbf{x}),$$



Collapsed trispectrum $\sim \langle F F \rangle$

Scale-dependent galaxy bias

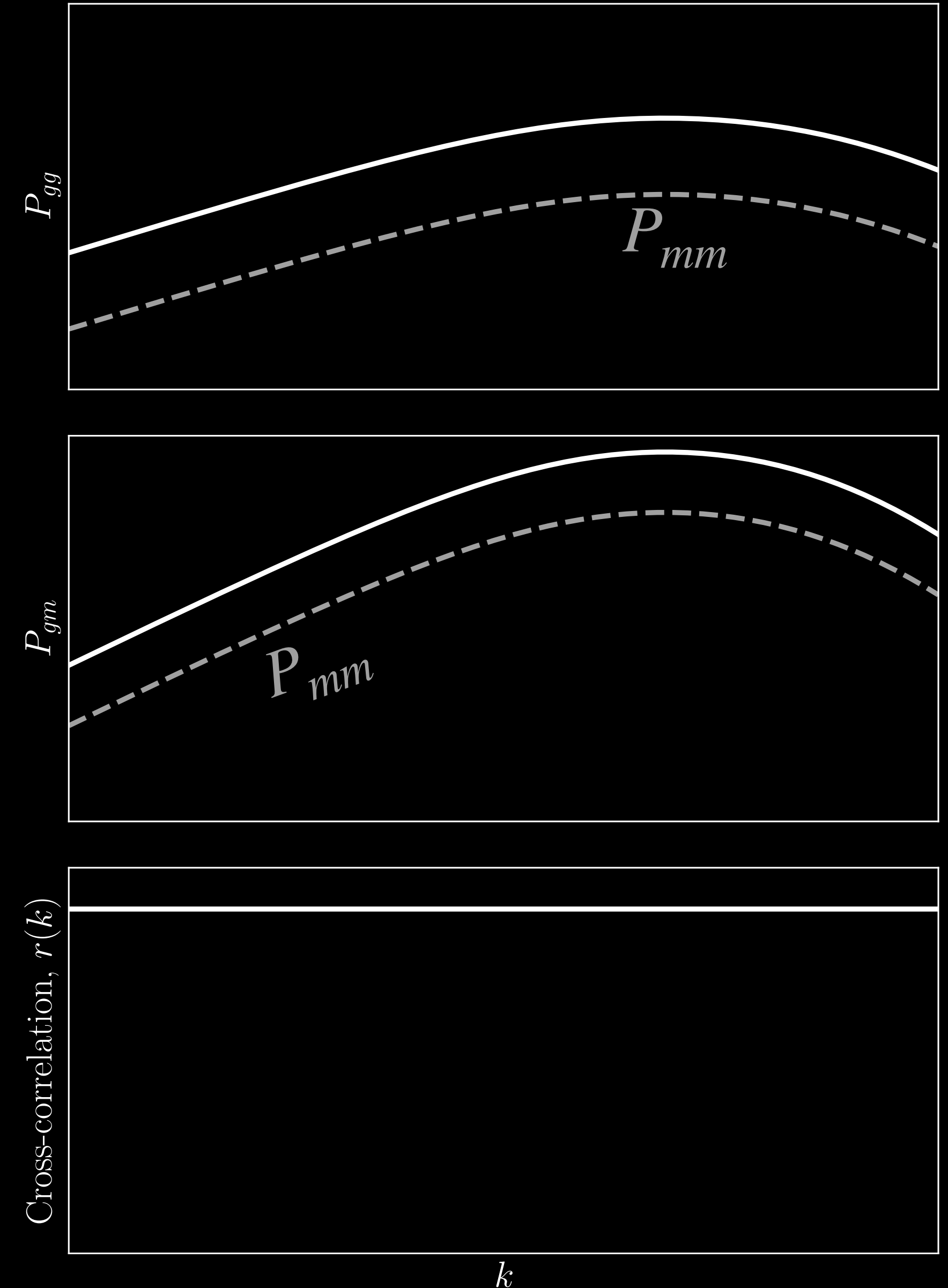
- **Gaussian initial conditions**

$$\delta_g(\mathbf{x}) = b\delta_m(\mathbf{x}) \sim b \nabla^2 \Phi(\mathbf{x})$$

$$\Rightarrow P_{gg} = b^2 P_{mm}, \quad P_{gm} = b P_{mm}$$

- **Galaxies trace matter: deterministic***

$$r \equiv P_{gm} / \sqrt{P_{gg} P_{mm}} = 1$$



*Ignoring shotnoise for convenience here

Scale-dependent galaxy bias

- $f_{\text{NL}}^{\text{loc}}$: galaxies respond to modulating field Φ

$$\delta_g(\mathbf{x}) = b\delta_m(\mathbf{x}) + b_{\phi}^* f_{\text{NL}} \Phi(\mathbf{x})$$

Poisson Eq. $\Phi(\mathbf{k}) = \delta_m(\mathbf{k})/\alpha(k)$

$$\alpha(k) = \frac{2k^2 T(k) D(z)}{3\Omega_m H_0^2} \propto k^2$$

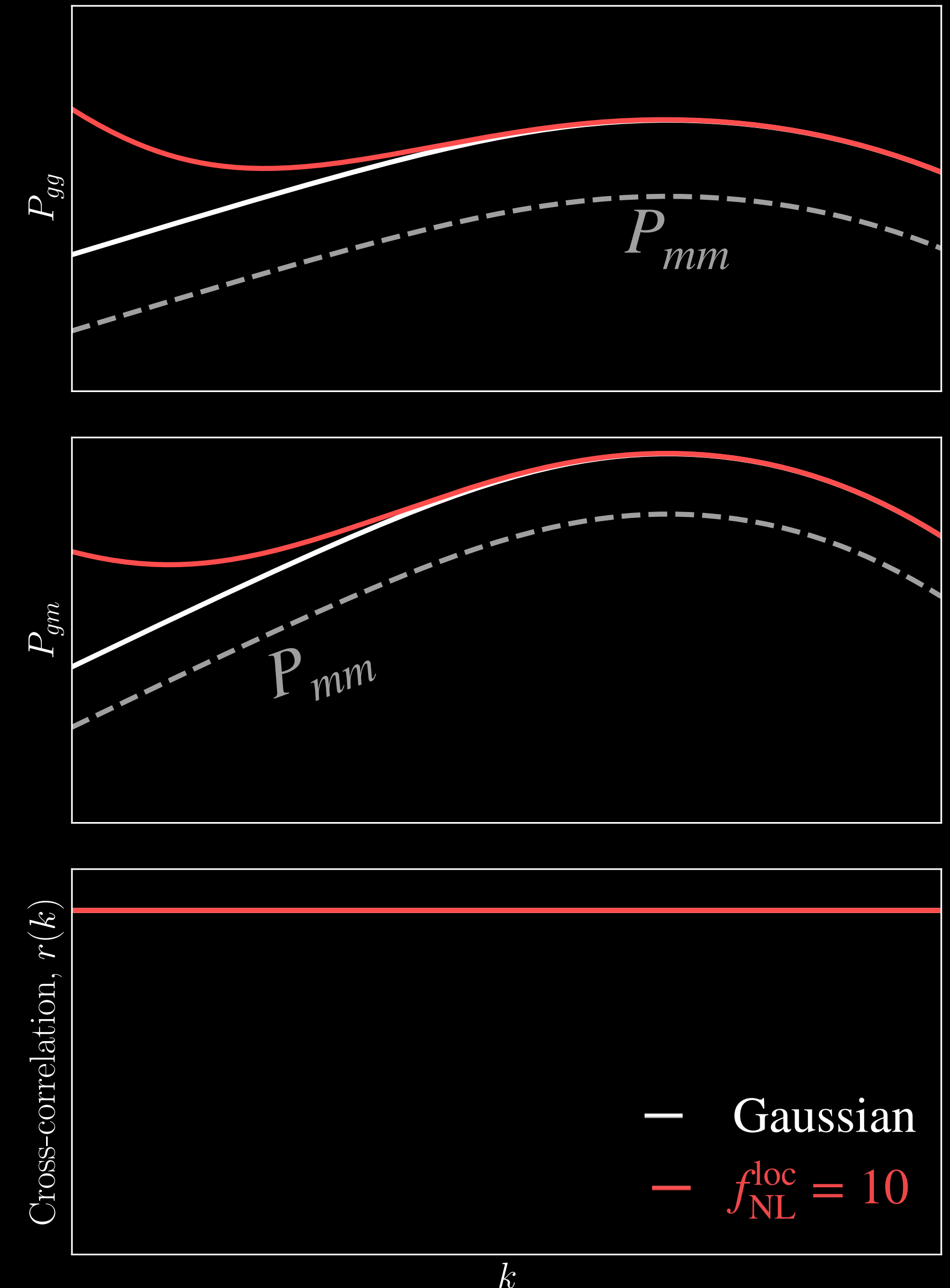
$$P_{gg} = \left[b + \frac{b_{\phi} f_{\text{NL}}}{\alpha(k)} \right]^2 P_{mm}$$

$$P_{gm} = \left[b + \frac{b_{\phi} f_{\text{NL}}}{\alpha(k)} \right] P_{mm}$$

- **Galaxies still trace matter:** deterministic

$$r \equiv P_{gm} / \sqrt{P_{gg} P_{mm}} = 1$$

*Here, we assume $b_{\phi} = 2\delta_c(b_1 - p)$ with $p = 1.4$ for QSO



Scale-dependent galaxy bias

- $\tau_{\text{NL}}^{\text{loc}}$: a 2nd field $\psi \perp \Phi$ also modulates power

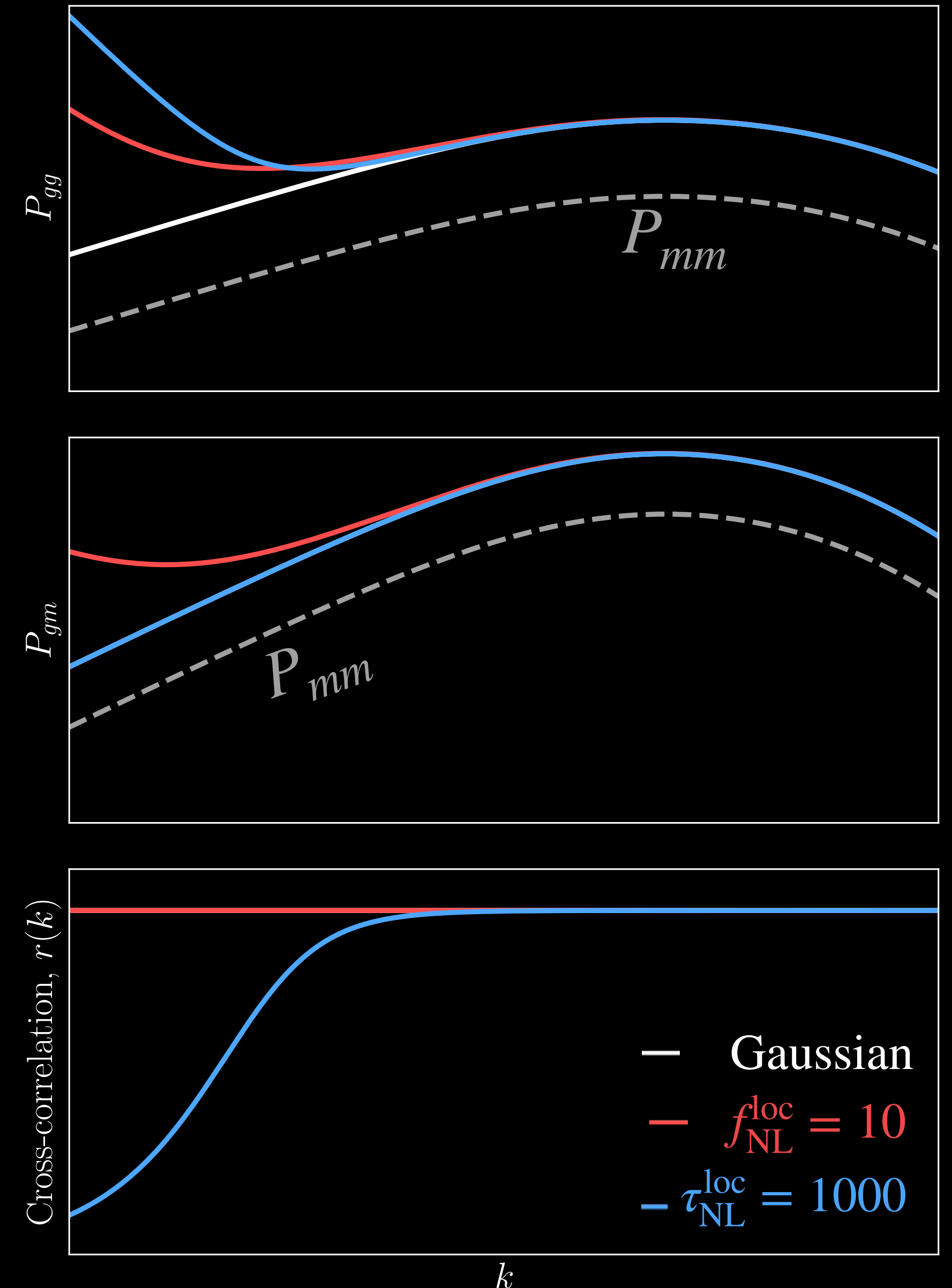
$$\delta_g(\mathbf{x}) = b \delta_m(\mathbf{x}) + b_\phi \left[f_{\text{NL}} \Phi(\mathbf{x}) + \sqrt{\tau_{\text{NL}} - \left(\frac{6}{5} f_{\text{NL}}\right)^2} \psi(\mathbf{x}) \right]$$

$$P_{gg} = \left[b + \frac{b_\phi f_{\text{NL}}}{\alpha(k)} \right]^2 P_{mm} + b_\phi^2 \frac{\tau_{\text{NL}} - \left(\frac{6}{5} f_{\text{NL}}\right)^2}{\alpha^2(k)} P_{mm}$$

$$P_{gm} = \left[b + \frac{b_\phi f_{\text{NL}}}{\alpha(k)} \right] P_{mm}$$

- **Galaxies no longer fully trace matter: stochastic**

$$r \equiv P_{gm} / \sqrt{P_{gg} P_{mm}} < 1$$



Stochastic bias: [Baumann++12](#), [Ferraro & Smith14](#)

Observational status of PNG: LSS is the future

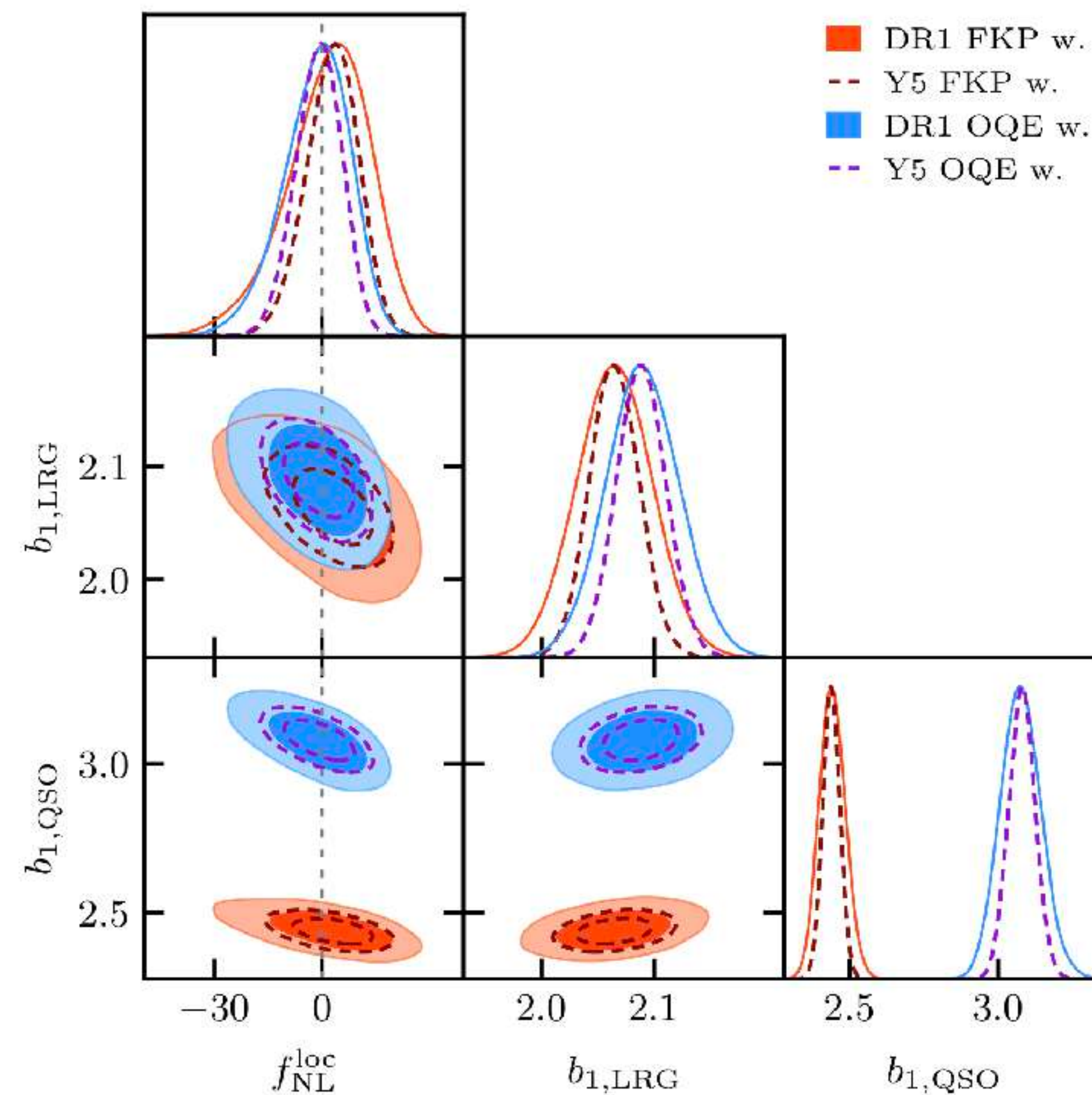
- Tightest bounds on PNG are currently from **CMB**

$$f_{\text{NL}}^{\text{loc}} = -0.9 \pm 5.1 \quad \tau_{\text{NL}}^{\text{loc}} < 1700 \quad (95\%)$$

[Planck 2018](#)

[Marzouk++22, Philcox++25](#)

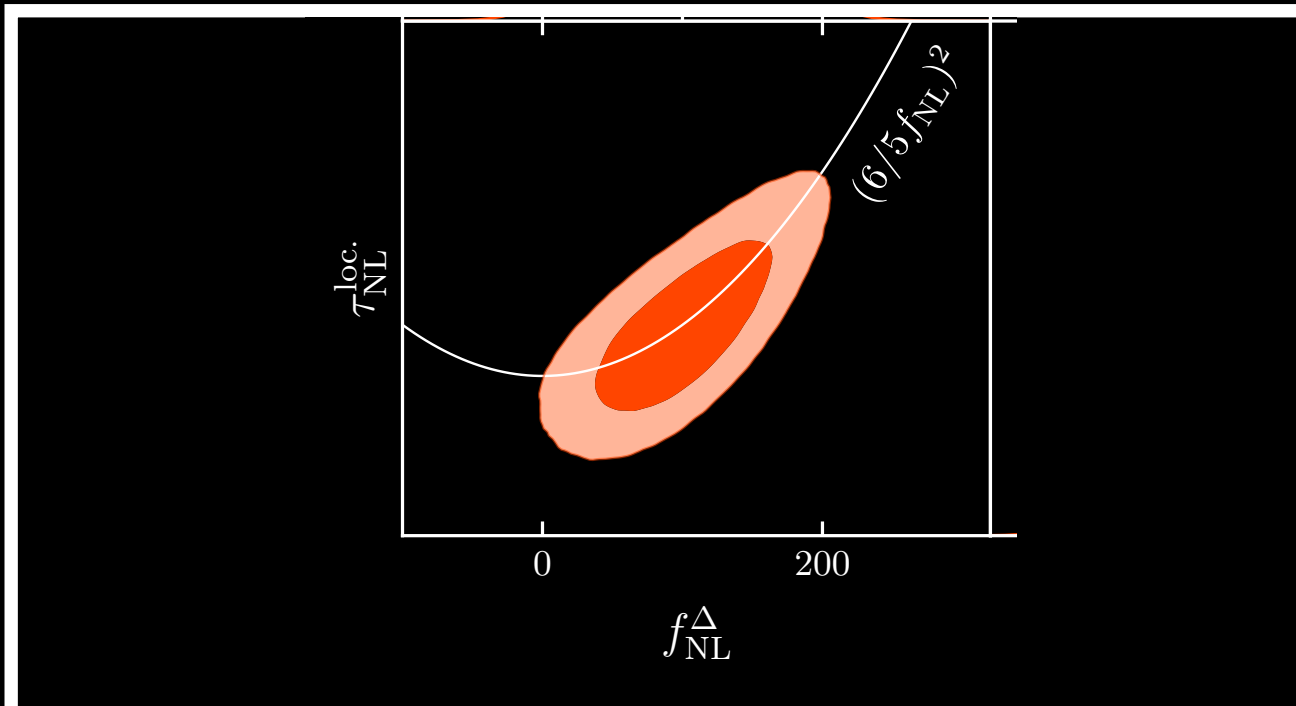
- Limits are close to cosmic variance limit
- CMB is 2D field: running out of modes
- **Galaxy surveys** can do much better
 - 3D volume, more modes, scale-dep bias
 - DESI already DR1 reached $\sigma(f_{\text{NL}}^{\text{loc}}) \sim 9$; DR2 will improve; SPHERE-X $\sigma(f_{\text{NL}}^{\text{loc}}) < 1$
 - *Have never been used for τ_{NL}*



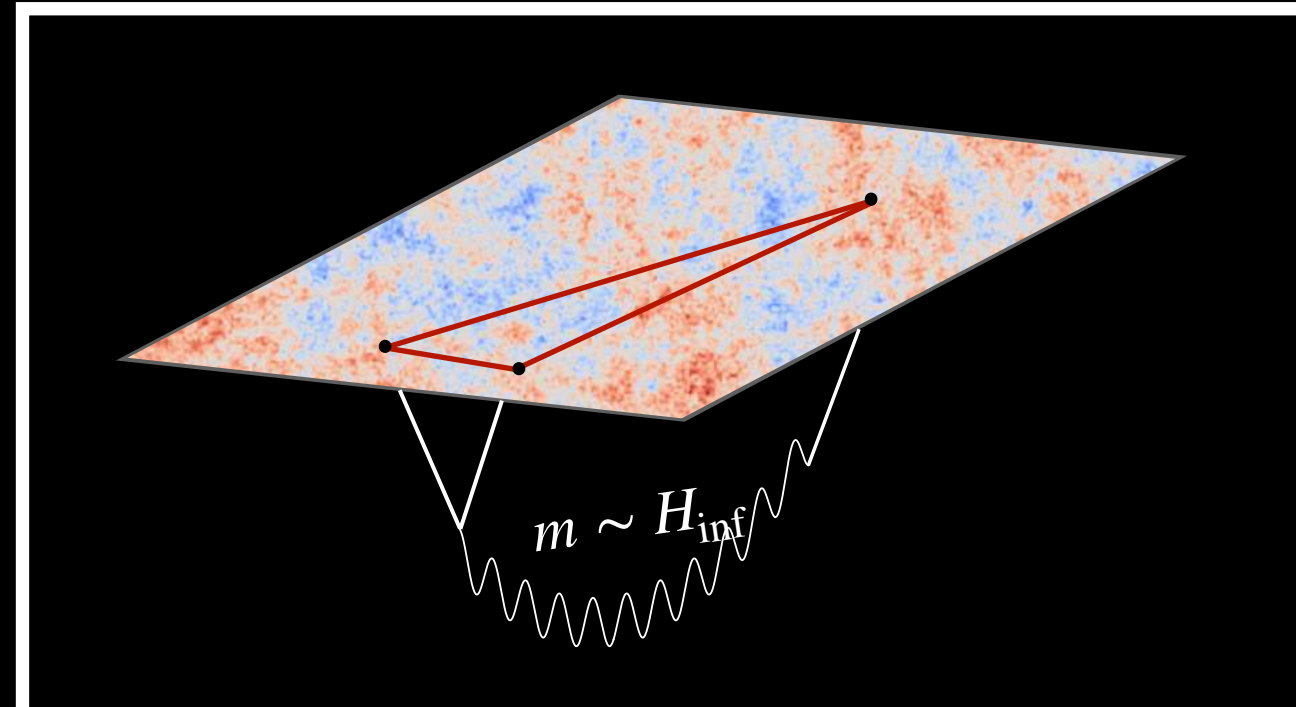
DESI DR1 constraints on $f_{\text{NL}}^{\text{loc}}$

[***Chaussidon++24***](#)

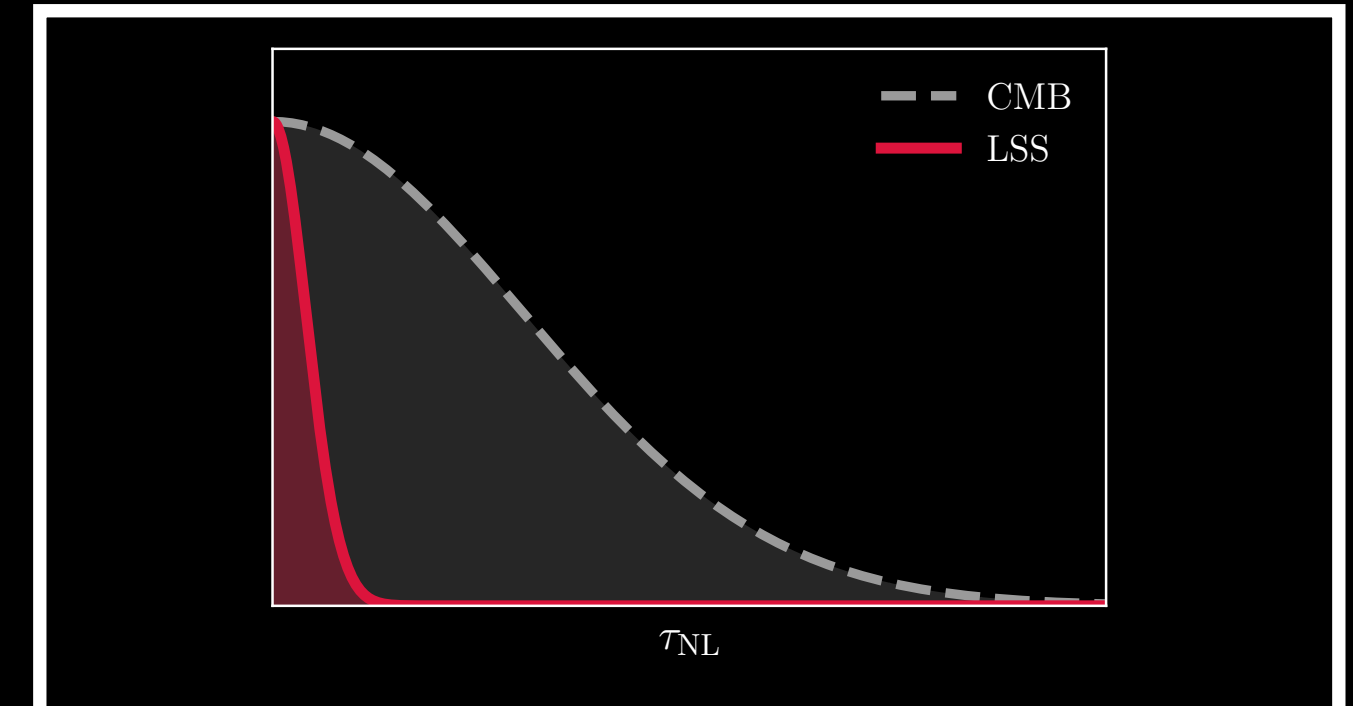
We should search for τ_{NL} with LSS



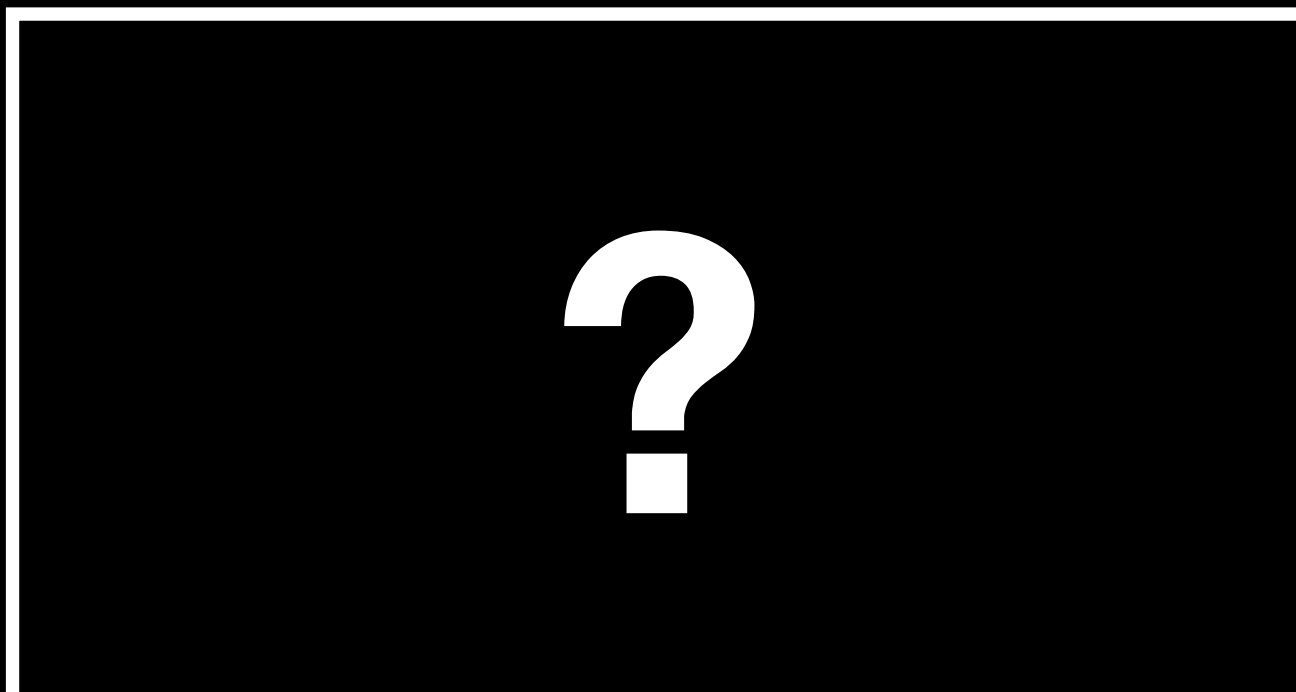
Theoretical motivation:
Multi-field inflation



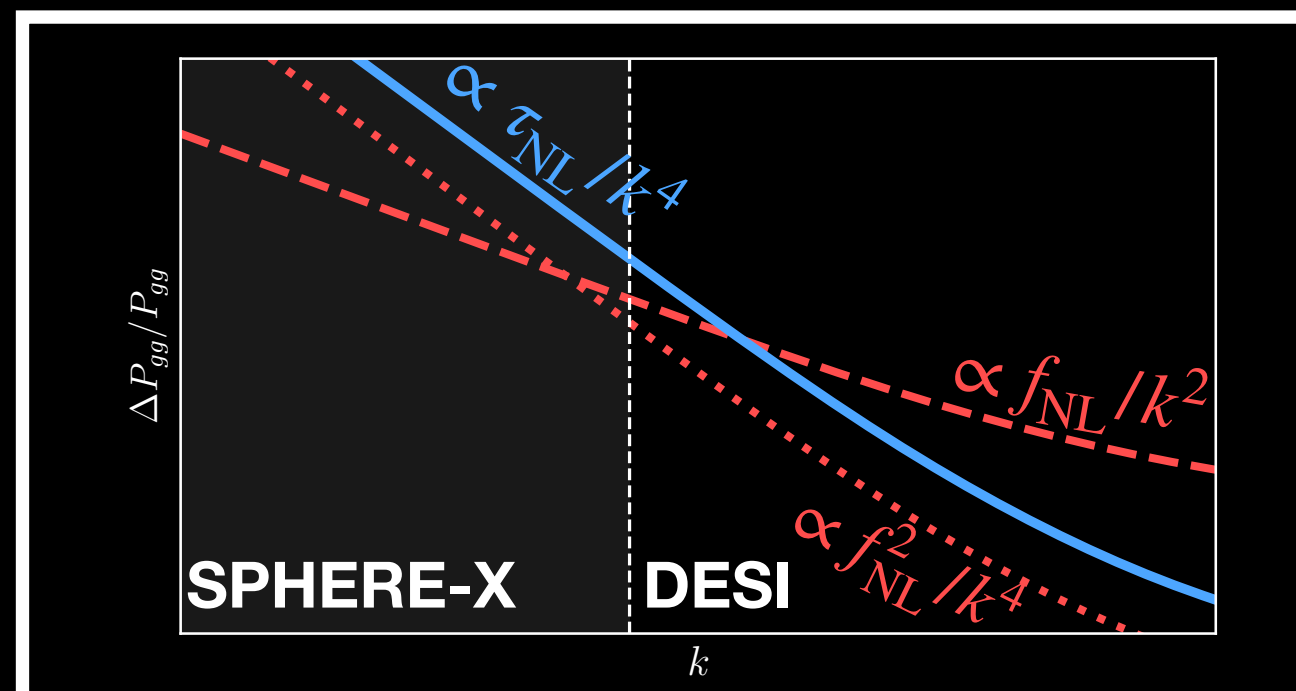
Theoretical motivation:
Cosmological collider



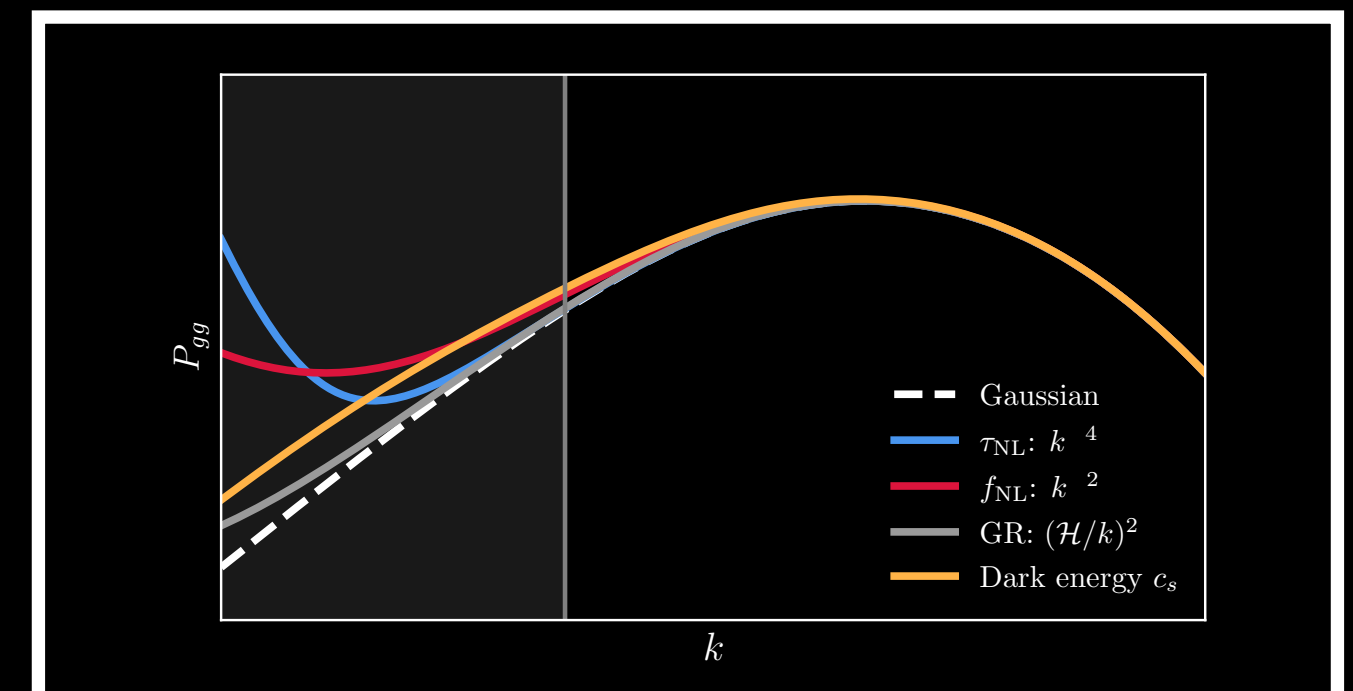
Near-term LSS surveys **10-100x**
more sensitive than CMB



Never before constrained with
LSS



τ_{NL} analysis of today is the f_{NL}
analysis of tomorrow



τ_{NL} pipeline can be used for
ultra-large scale $P(k)$

3-step plan to constraining τ_{NL} with DESI

```
tnl_exc_loc = params['tnl_exc_loc']  
if bphiX is not None:  
    pkmu = pkmu + jac * fog * (25. / 36.) * tnl_exc_loc * bphiX * bphiY * alpha**2 * pk_obs  
self.power = self.to_poles(pkmu)
```

Step 1: Add τ_{NL} to DESI PNG pipeline (5 lines)

3-step plan to constraining τ_{NL} with DESI



Step 2: Run analysis on DESI DR2 blinded quasars $\tau_{\text{NL}} = 80 \pm 240$

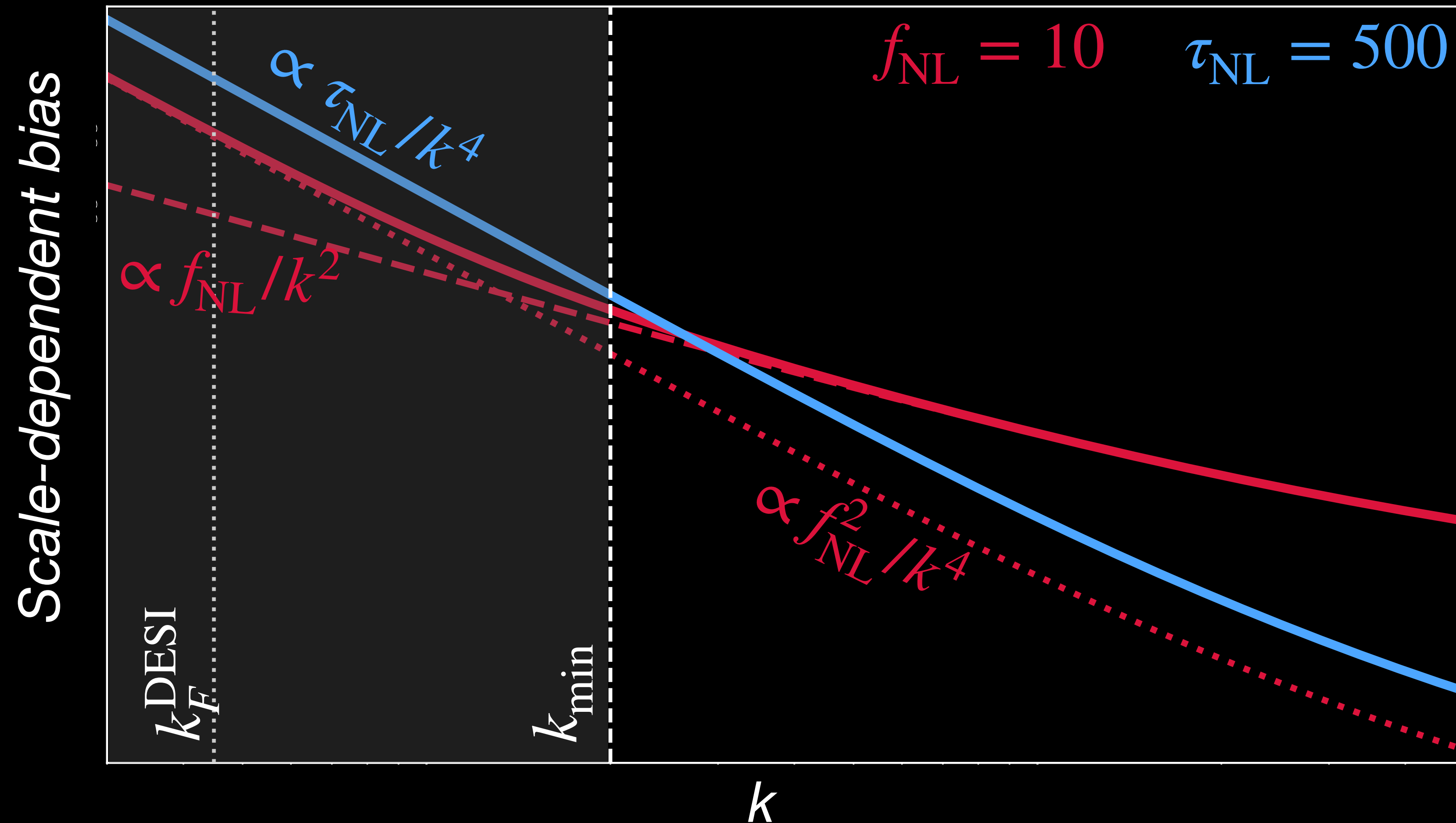
3-step plan to constraining τ_{NL} with DESI



$$\sigma < \sigma_{\text{Cramér-Rao}}$$

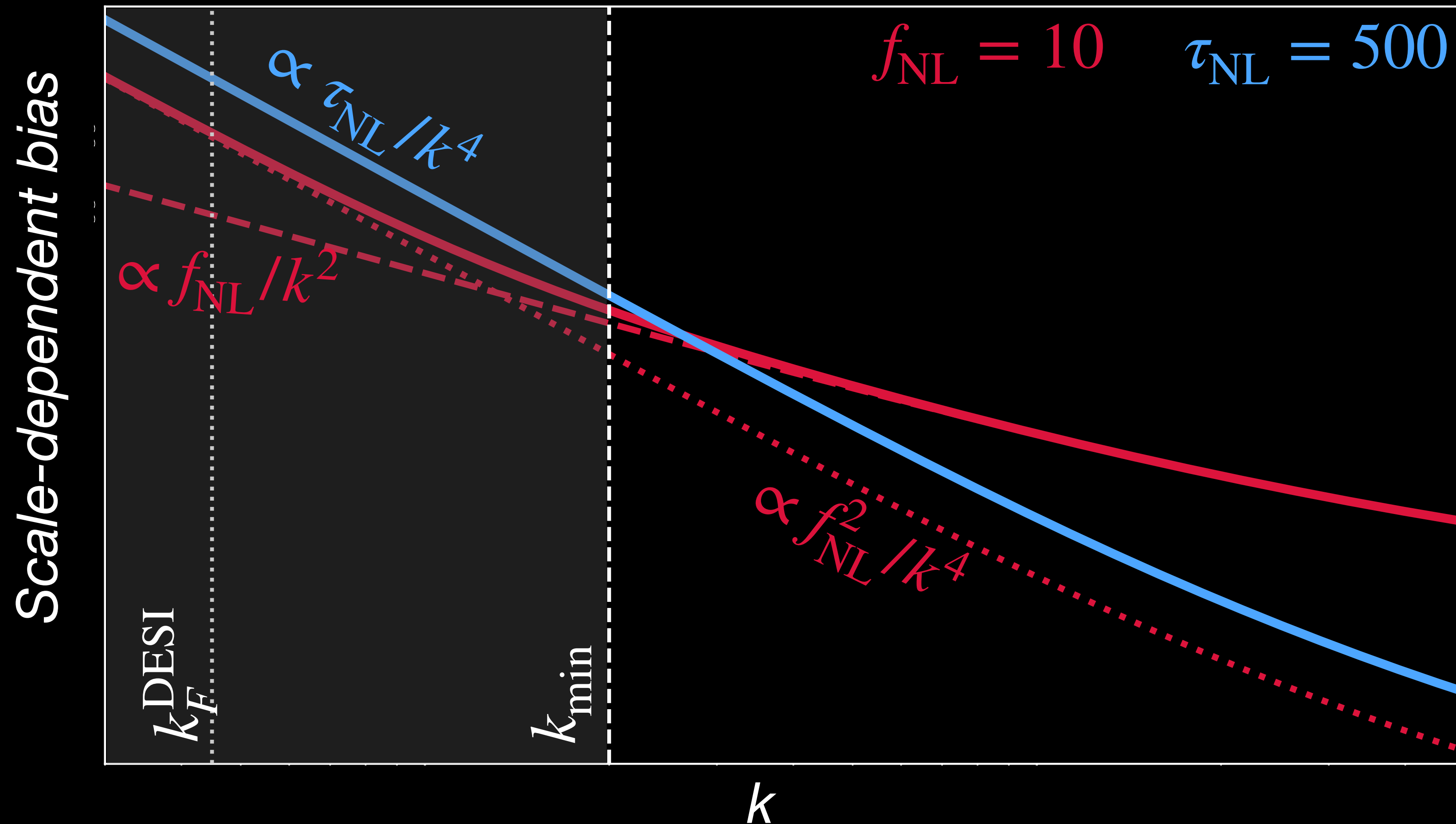
Step 3: find constraints that are tighter than a Fisher forecast with similar settings

Why is it difficult to constrain τ_{NL} with DESI?



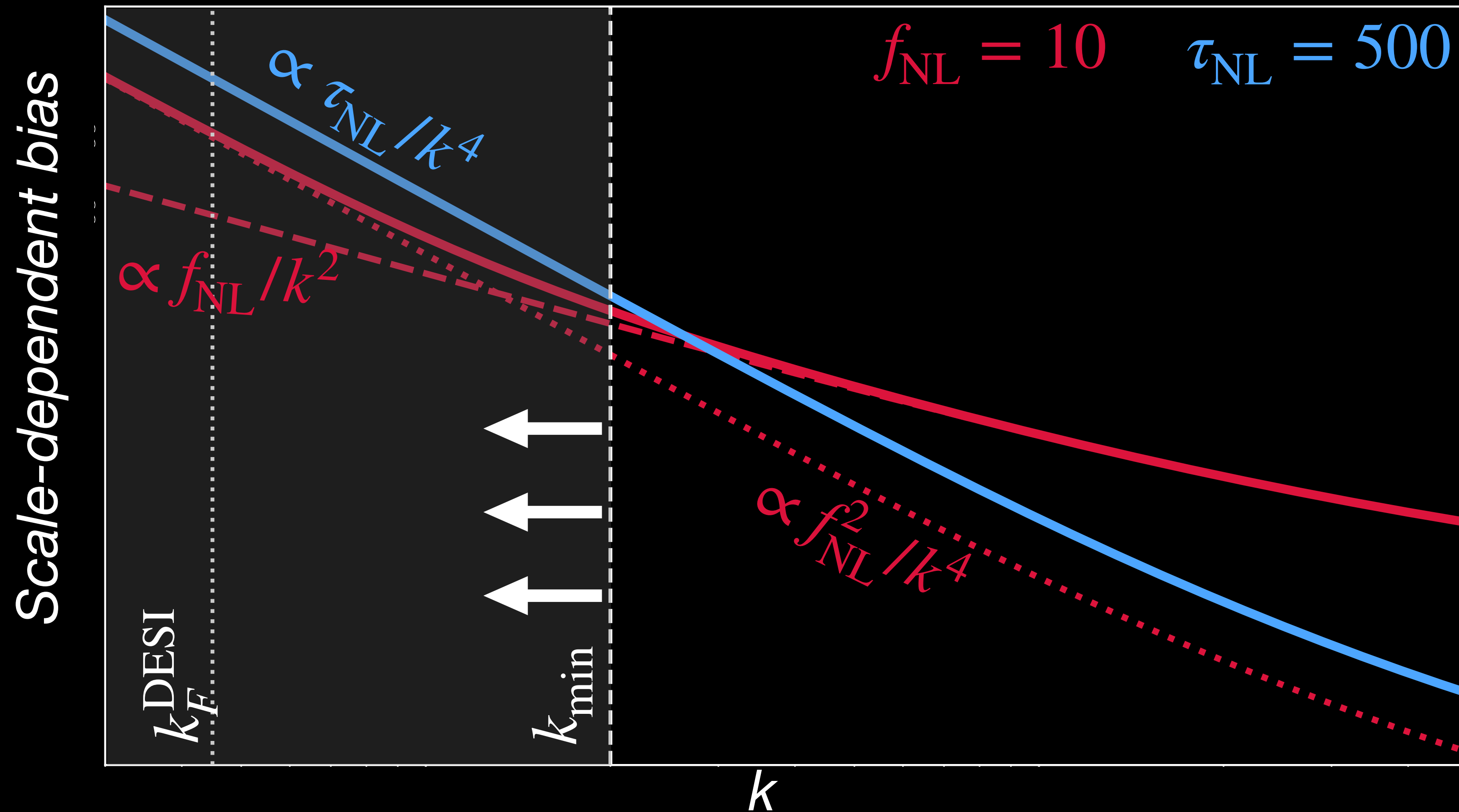
τ_{NL} : *all* information is in the largest-scale modes

Why is it difficult to constrain τ_{NL} with DESI?



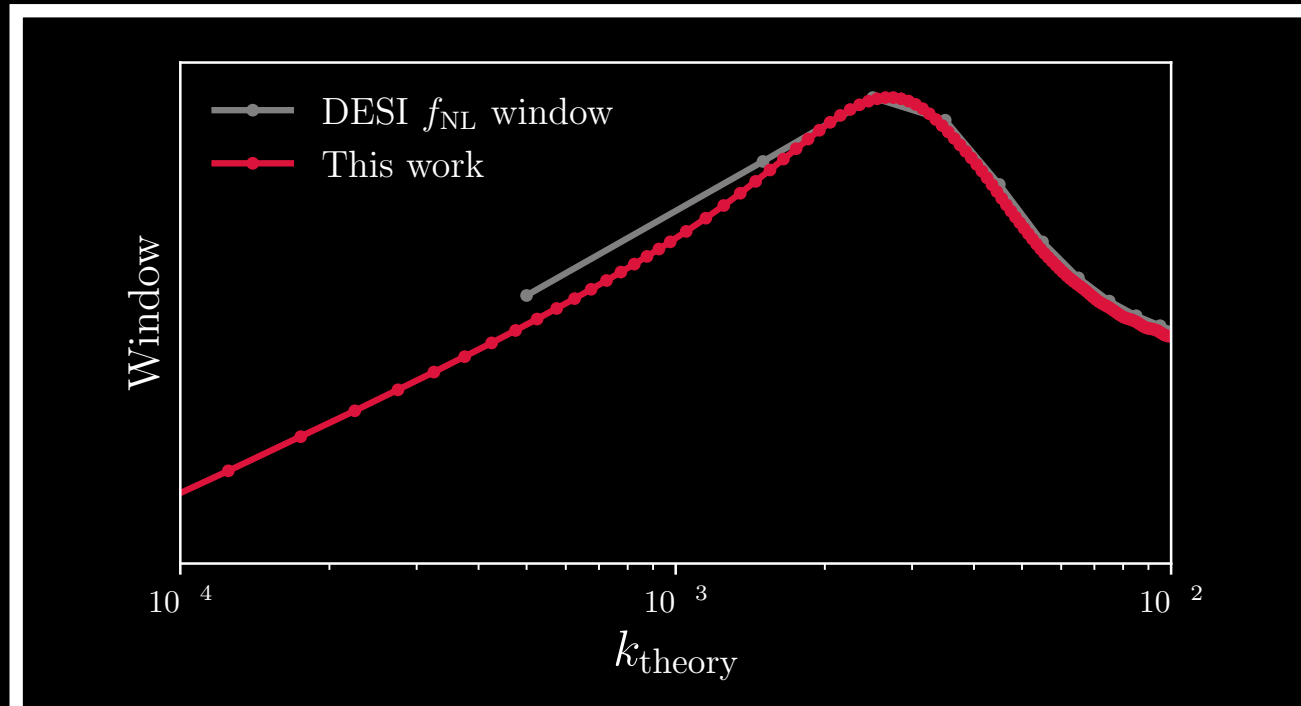
A robust f_{NL} pipeline is not necessarily a robust τ_{NL} pipeline

Why is it difficult to constrain τ_{NL} with DESI?

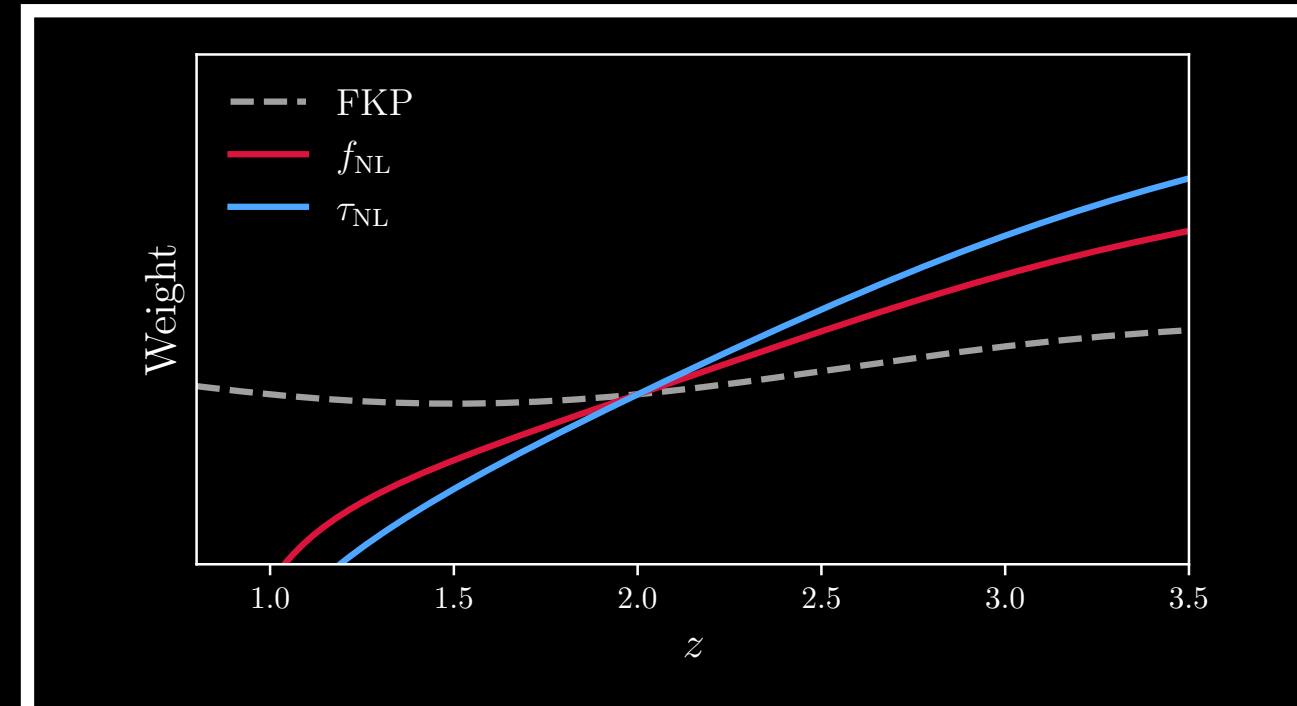


At low k , local f_{NL} $\Delta b \propto f_{\text{NL}}^2/k^4$. τ_{NL} with DESI is a precursor for all future searches for f_{NL} using the scale-dependent bias

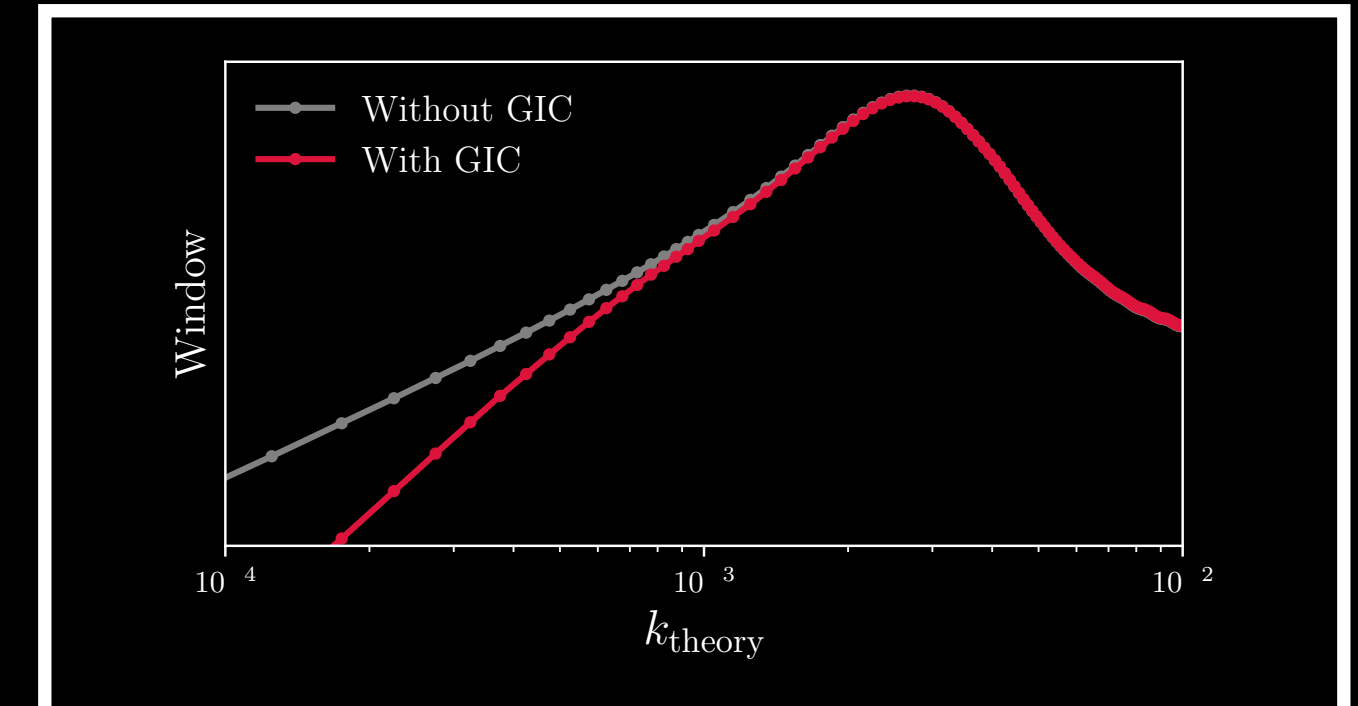
Creating a robust τ_{NL} pipeline for DESI quasars



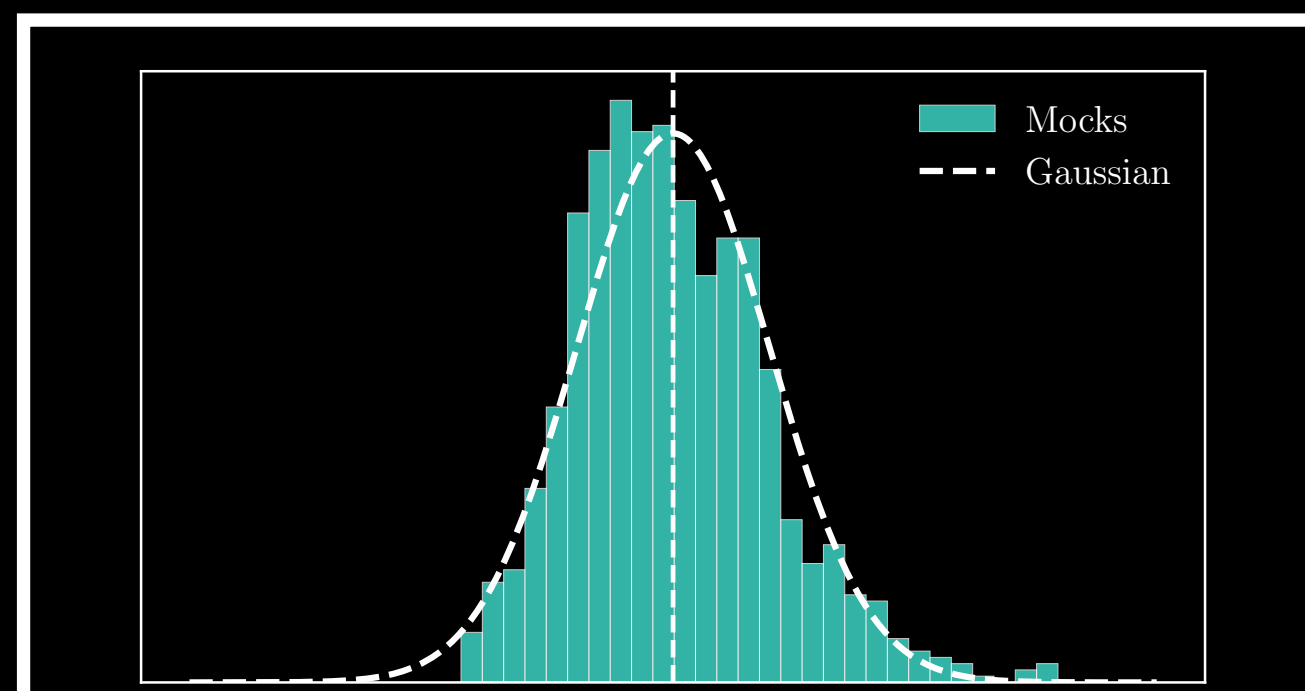
New window functions



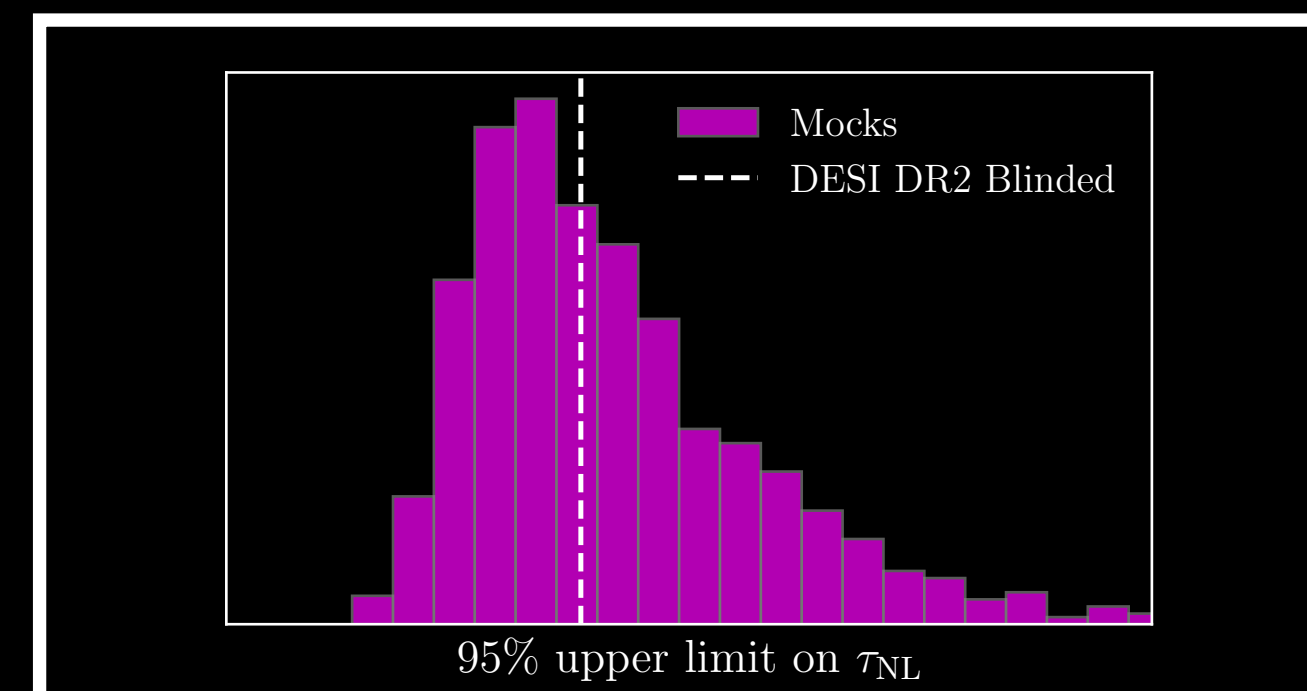
τ_{NL} optimal weights



Improved integral constraints

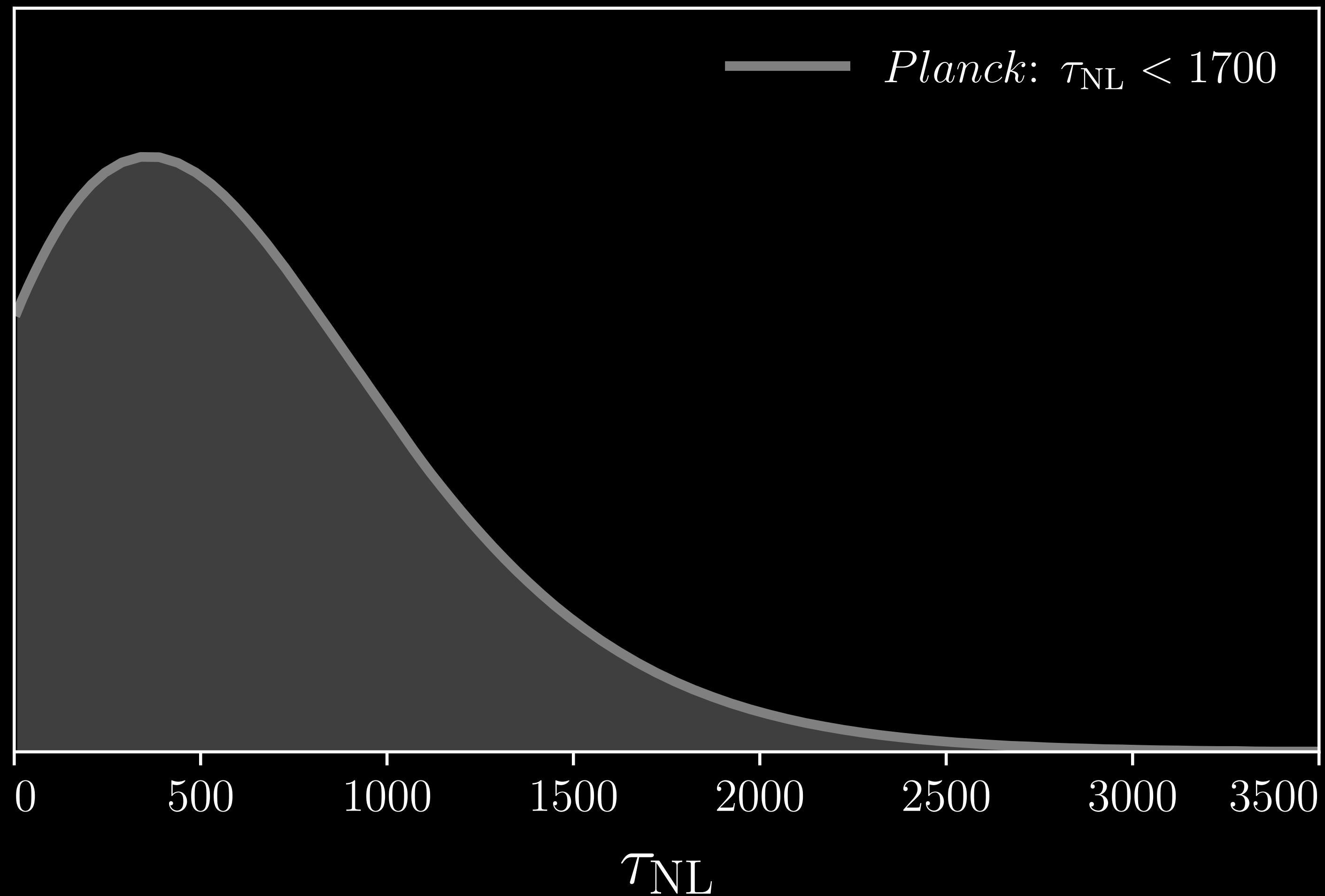


Non-Gaussian likelihood



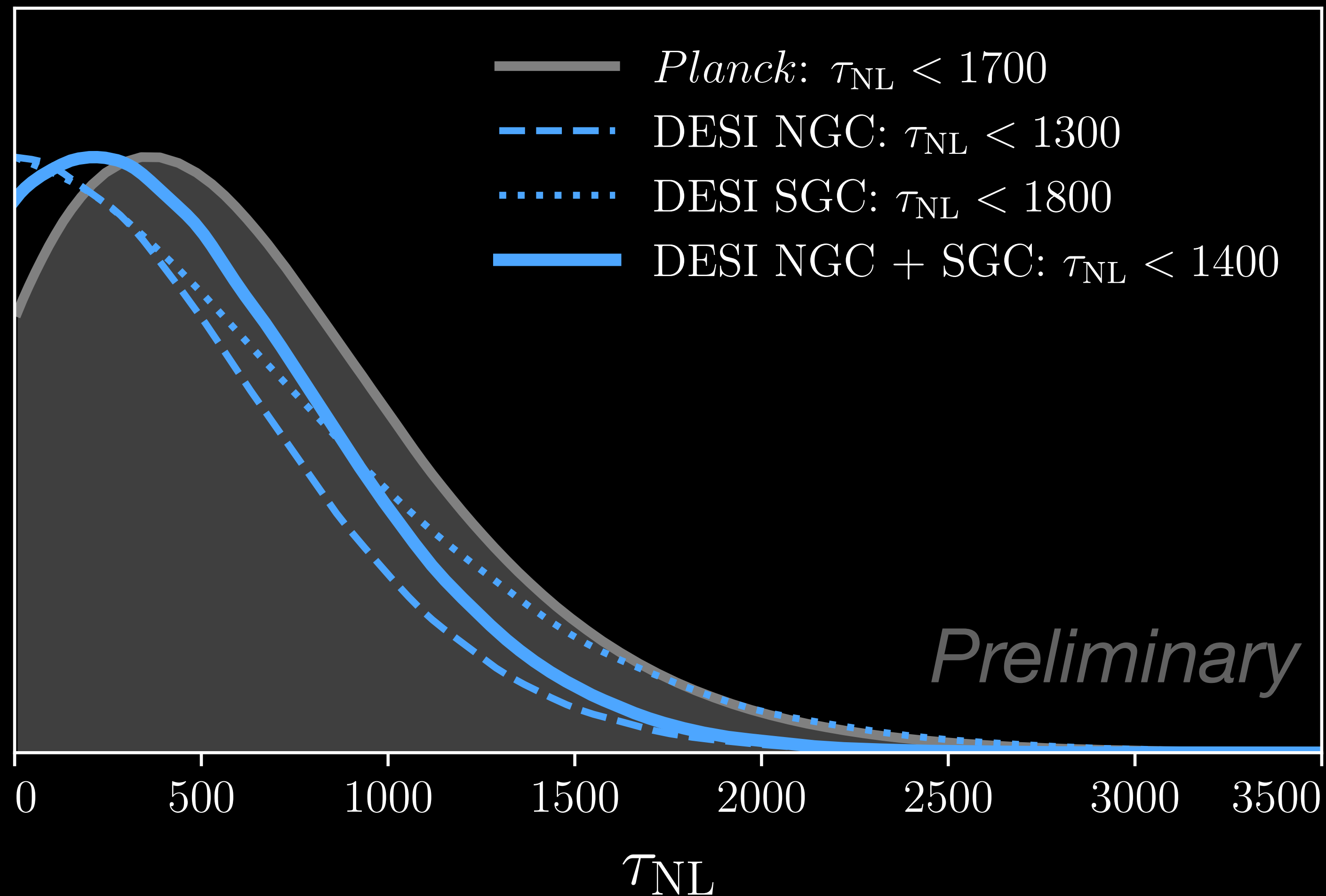
Systematics tests

Main result: first LSS constraints on τ_{NL}



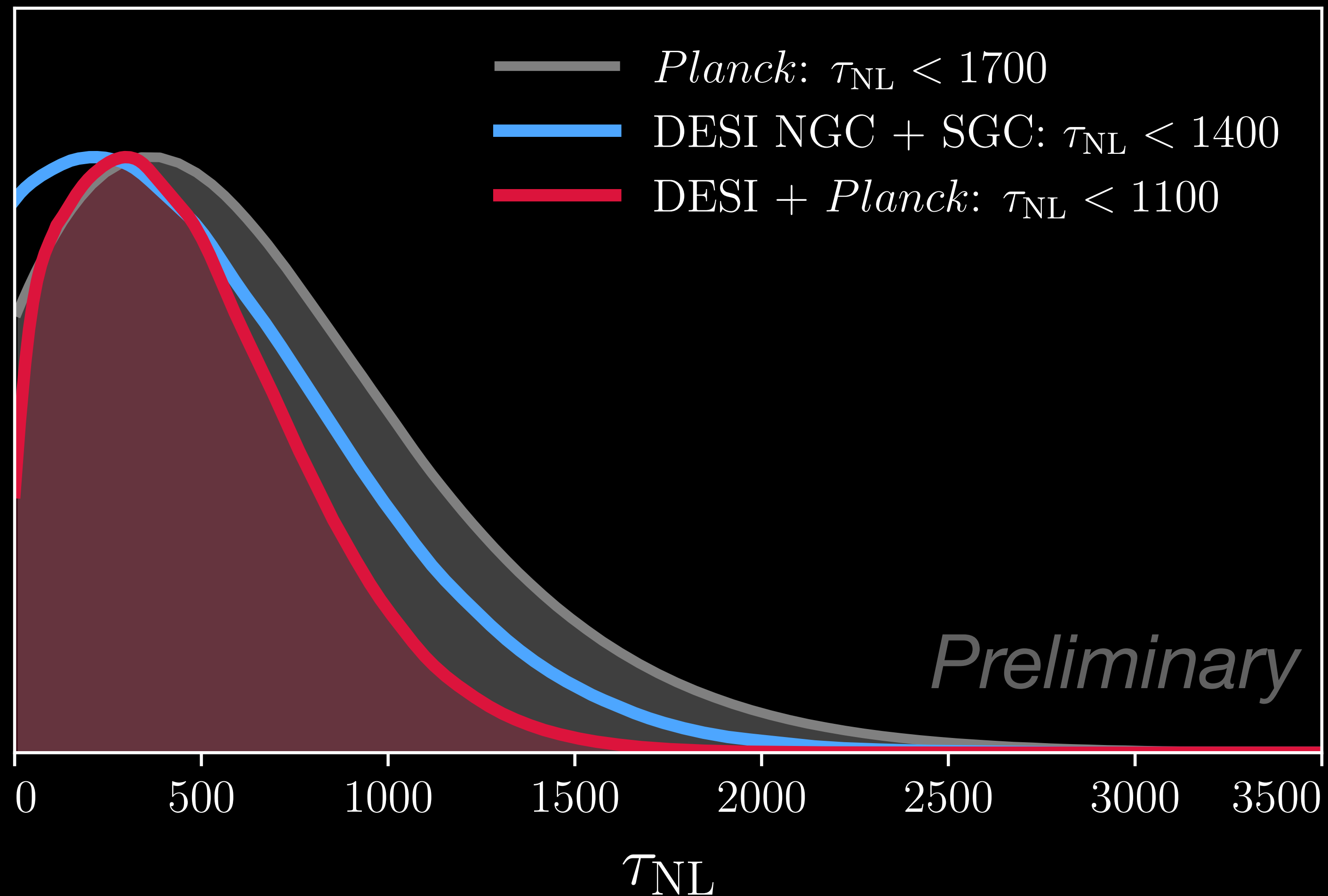
*Constraints from DESI DR2 **blinded** QSO: DESI DR2 data are already significantly more sensitive to τ_{NL} than the primary CMB trispectrum*

Main result: first LSS constraints on τ_{NL}



Constraints from DESI DR2 **blinded** QSO: DESI DR2 data are already significantly more sensitive to τ_{NL} than the primary CMB trispectrum

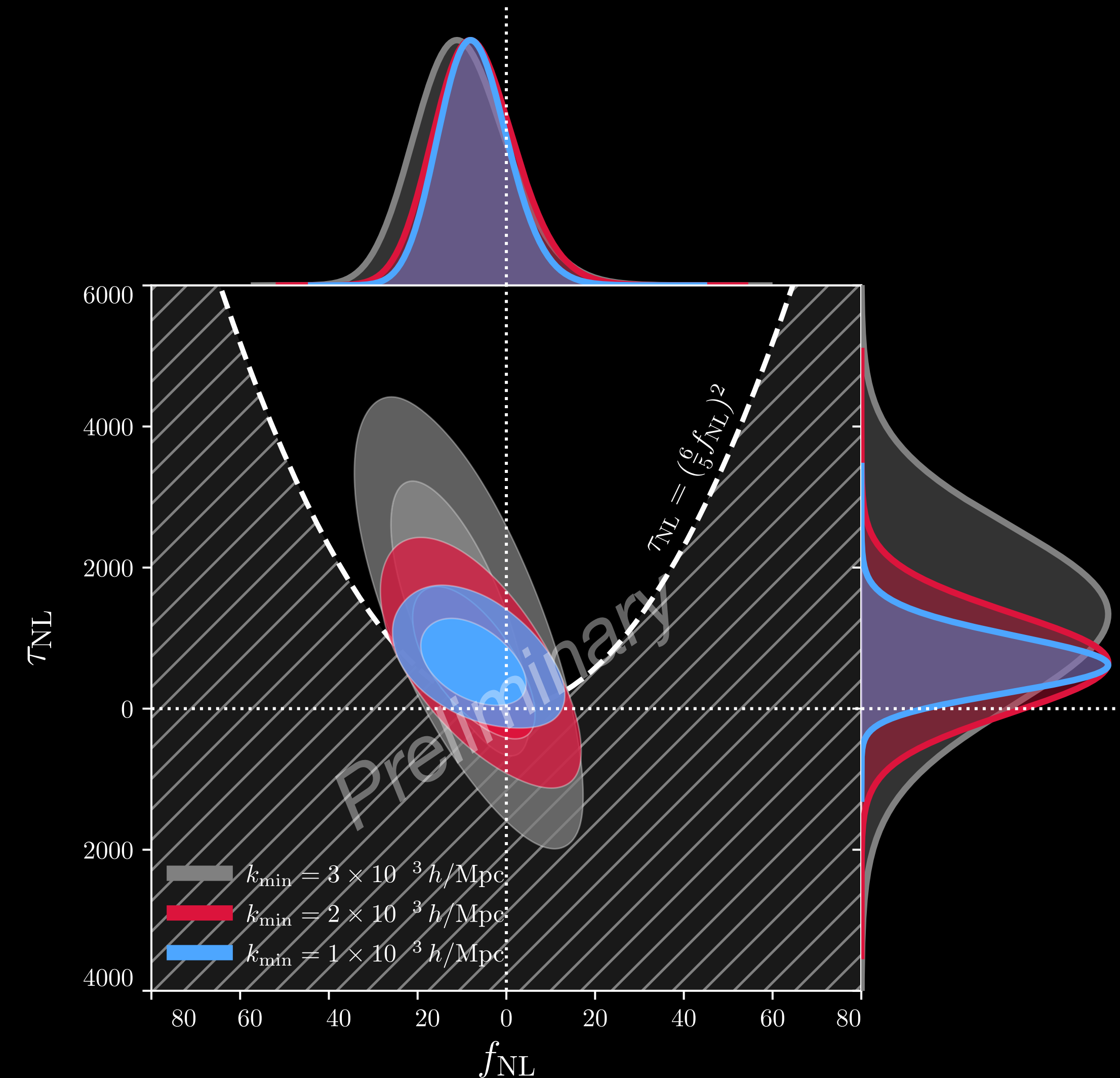
Main result: first LSS constraints on τ_{NL}



Constraints from DESI DR2 **blinded** QSO: DESI DR2 data are already significantly more sensitive to τ_{NL} than the primary CMB trispectrum

Varying f_{NL} and τ_{NL} jointly

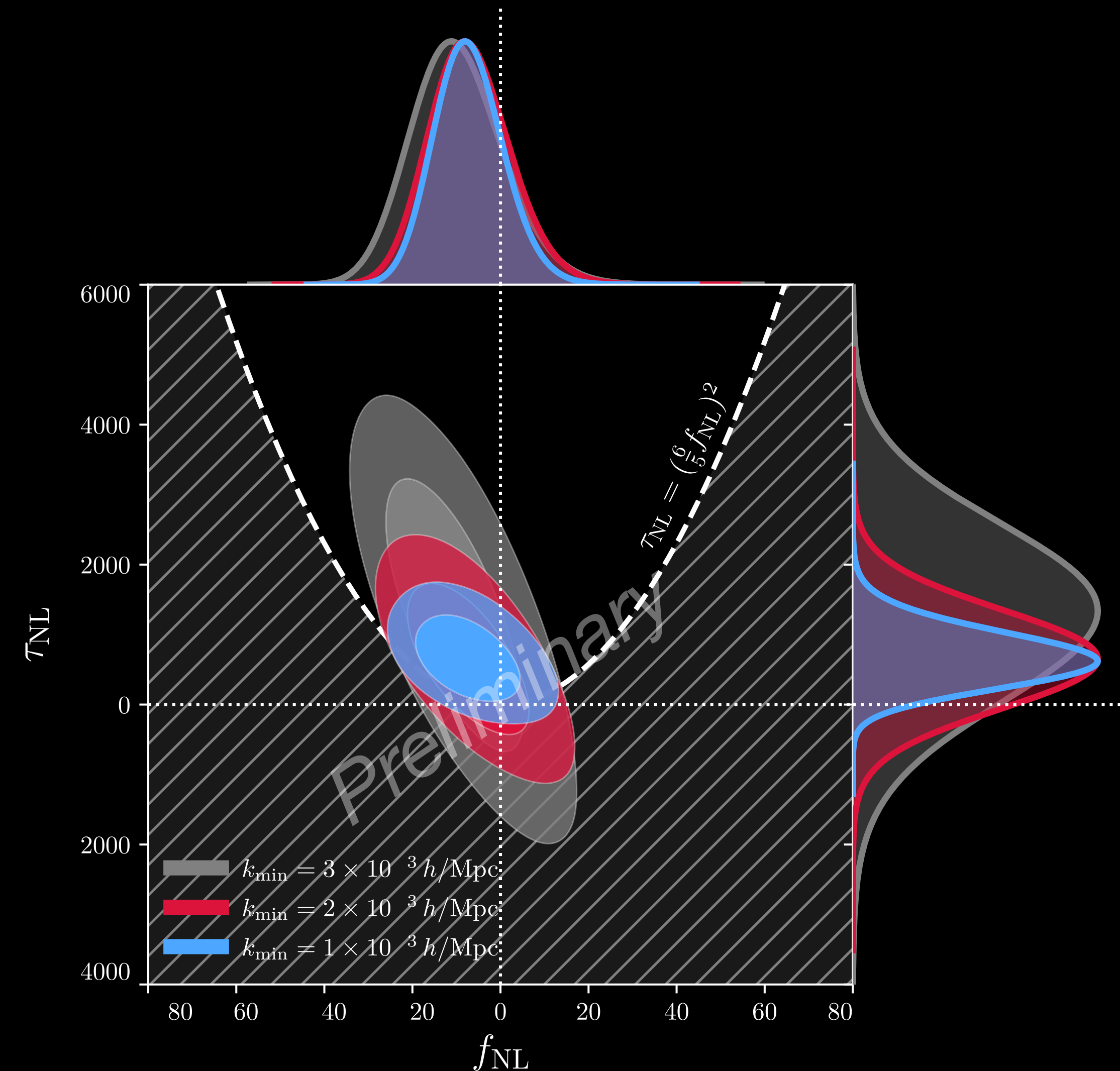
- Best test of multi-field inflation to date
 - ζ generated by one source: $\tau_{\text{NL}} = \left(\frac{6}{5}f_{\text{NL}}\right)^2$
- Powerful test of DESI QSO data
 - Imaging systematics/target selection ([Chaussidon++22](#))



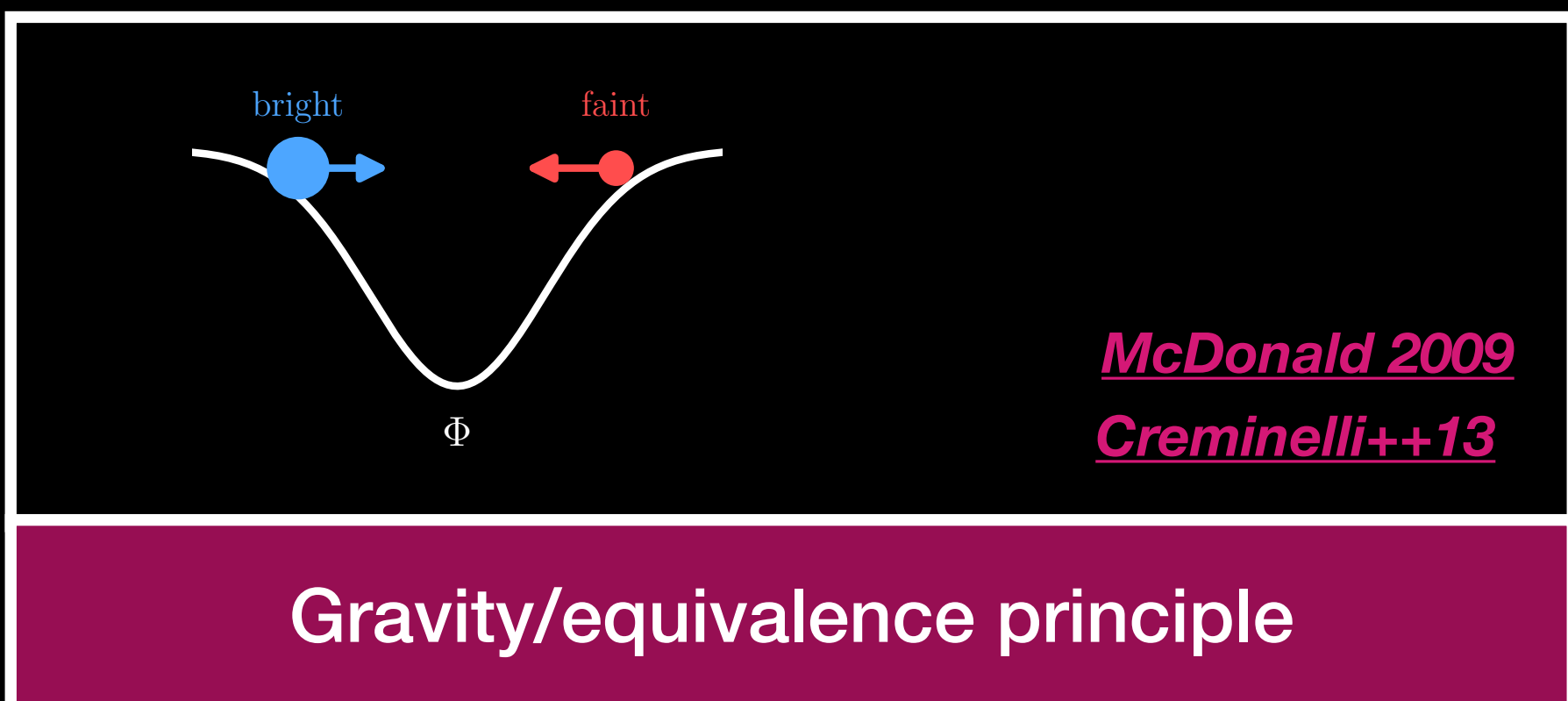
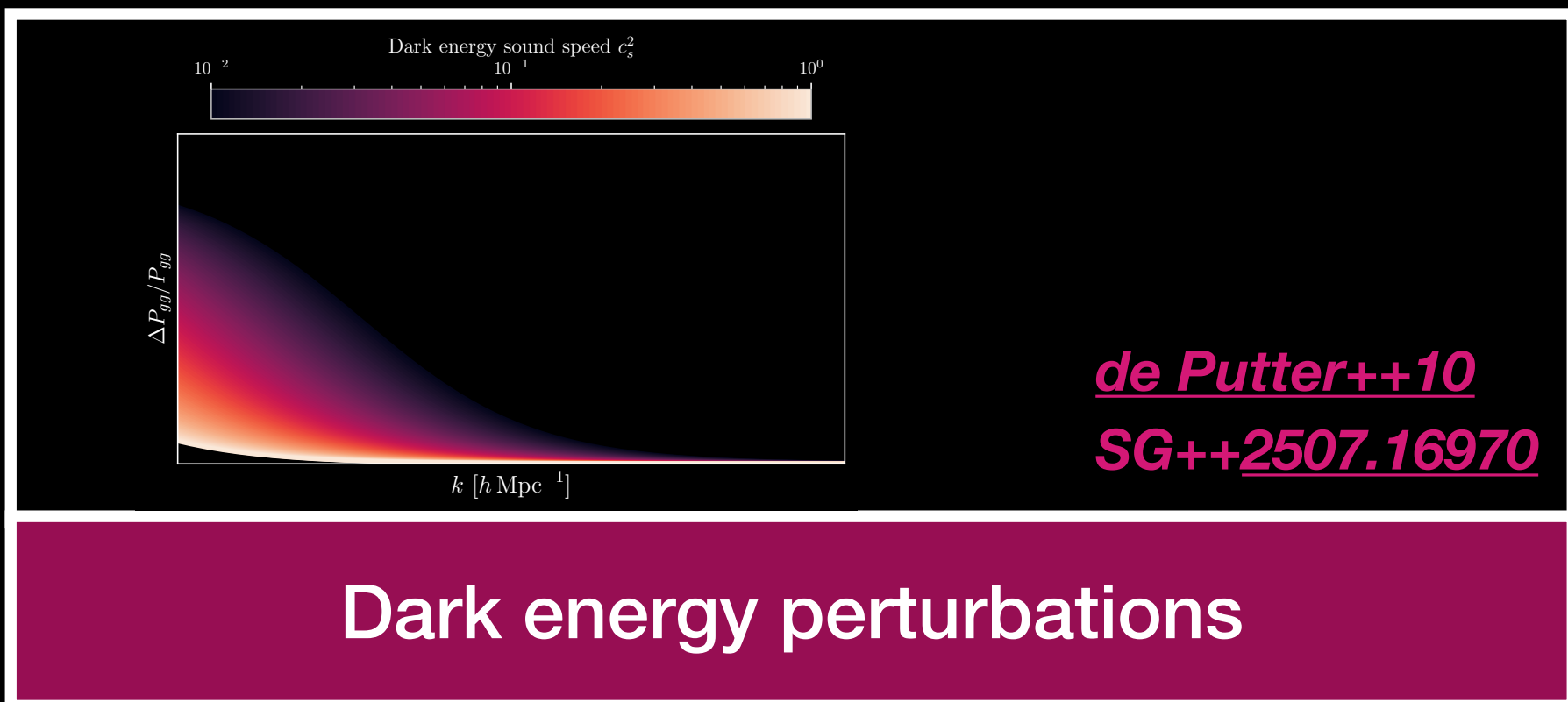
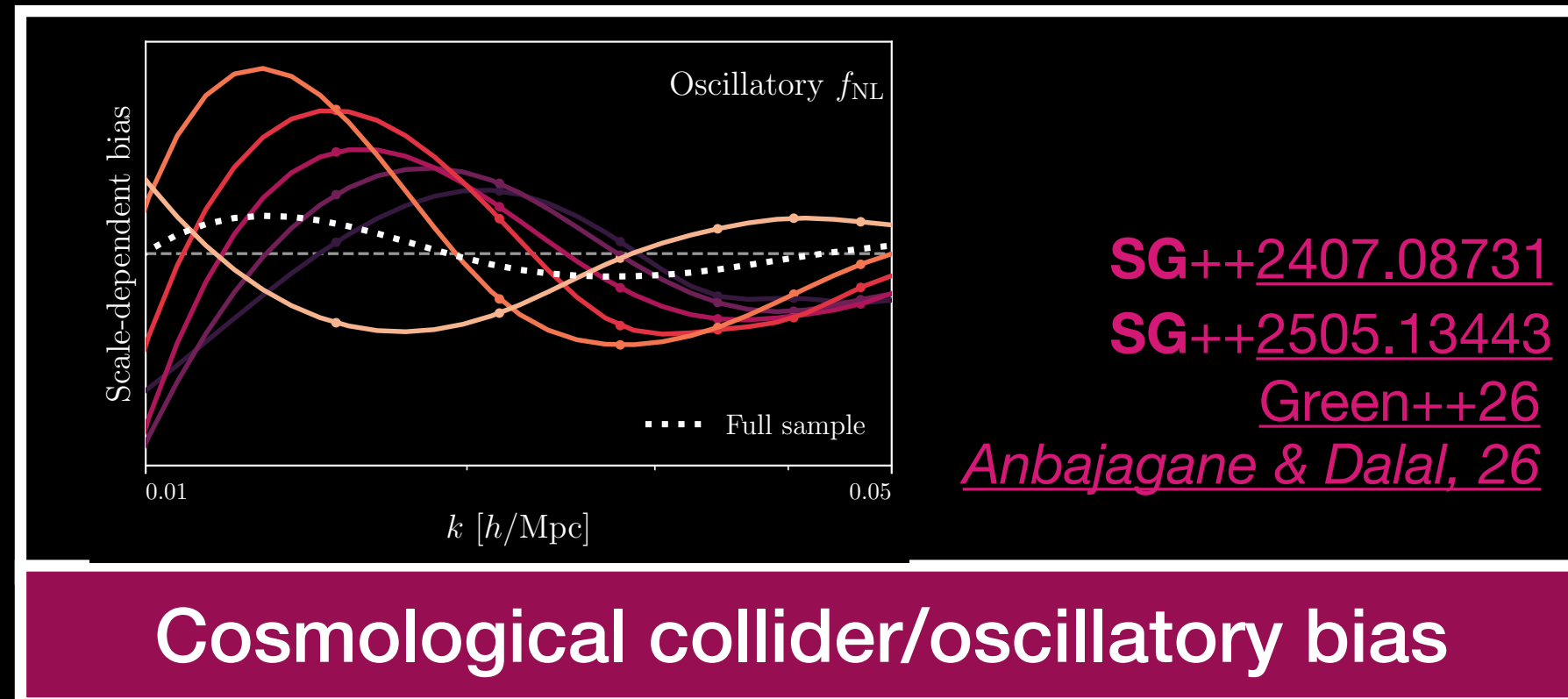
Varying f_{NL} and τ_{NL} jointly

- Best test of multi-field inflation to date
 - ζ generated by one source: $\tau_{\text{NL}} = \left(\frac{6}{5}f_{\text{NL}}\right)^2$
- Powerful test of DESI QSO data
 - Imaging systematics/target selection ([Chaussidon++22](#))

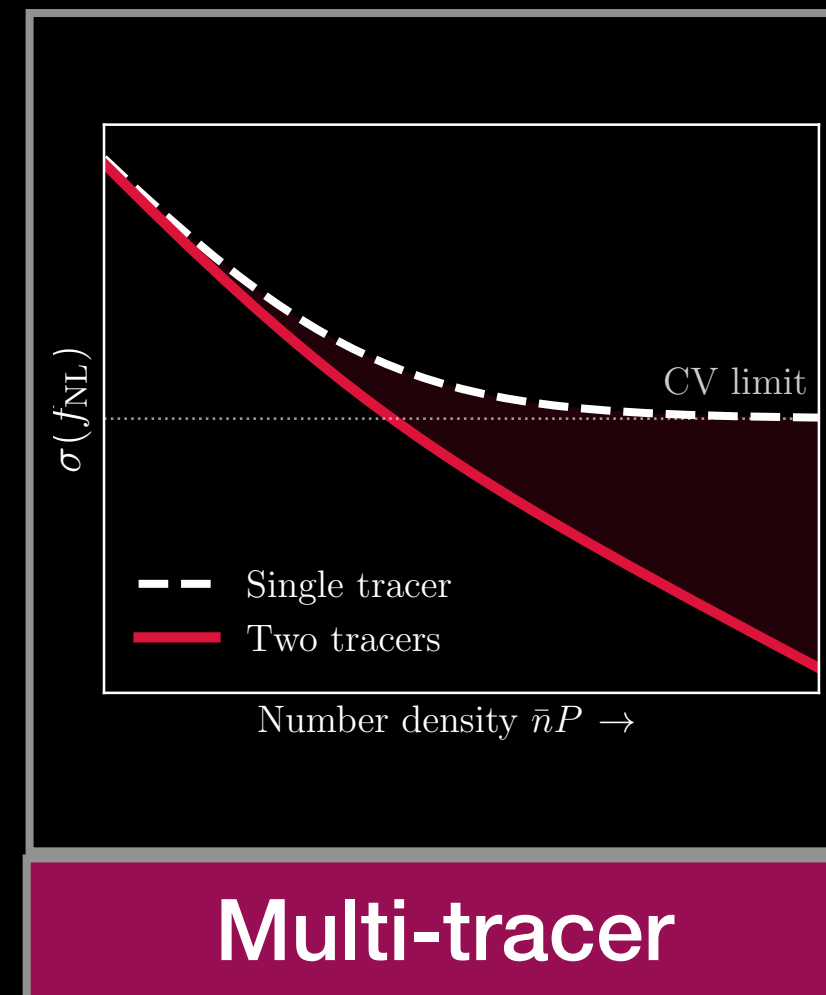
What's next?



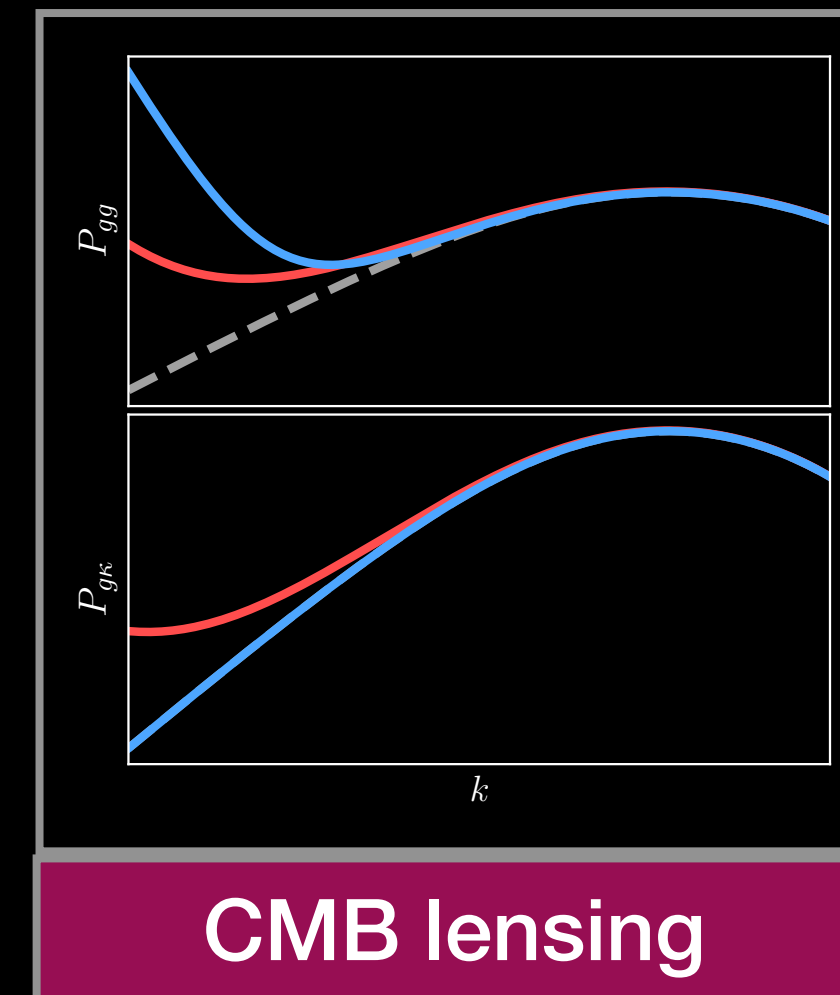
Near term: new physics with τ_{NL} pipeline



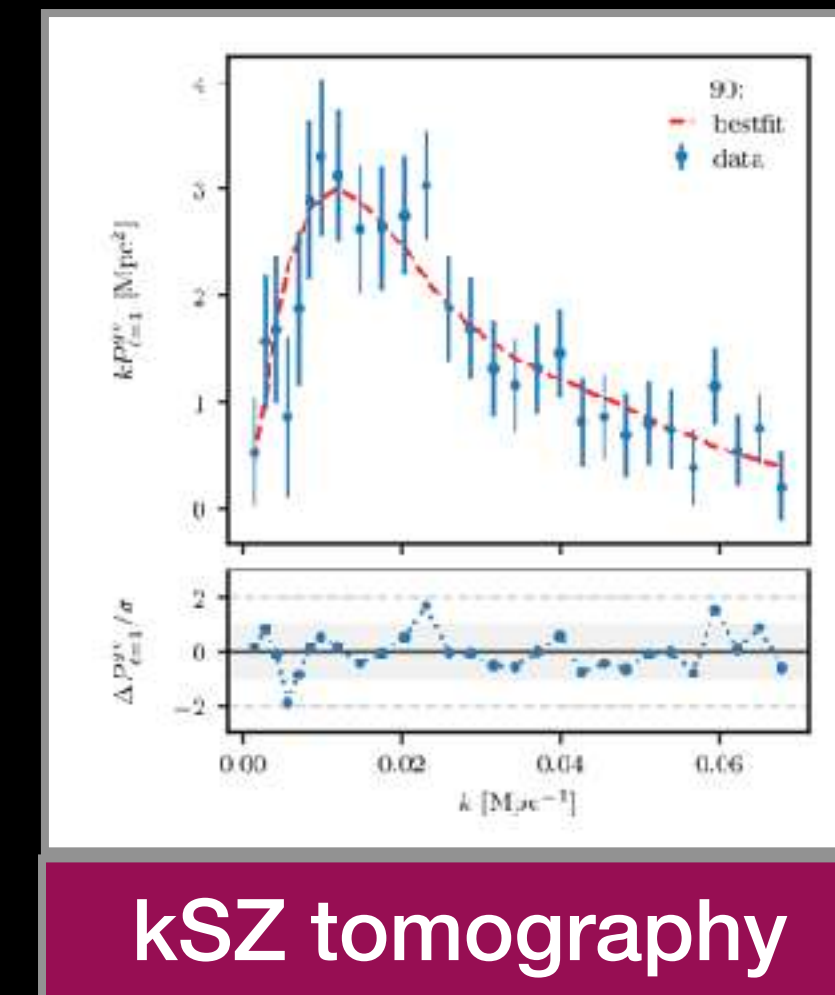
Medium term: multi-tracer and cross-correlations



See [Seljak 08](#), [McDonald & Seljak 08](#), [Ferraro & Smith 14](#)

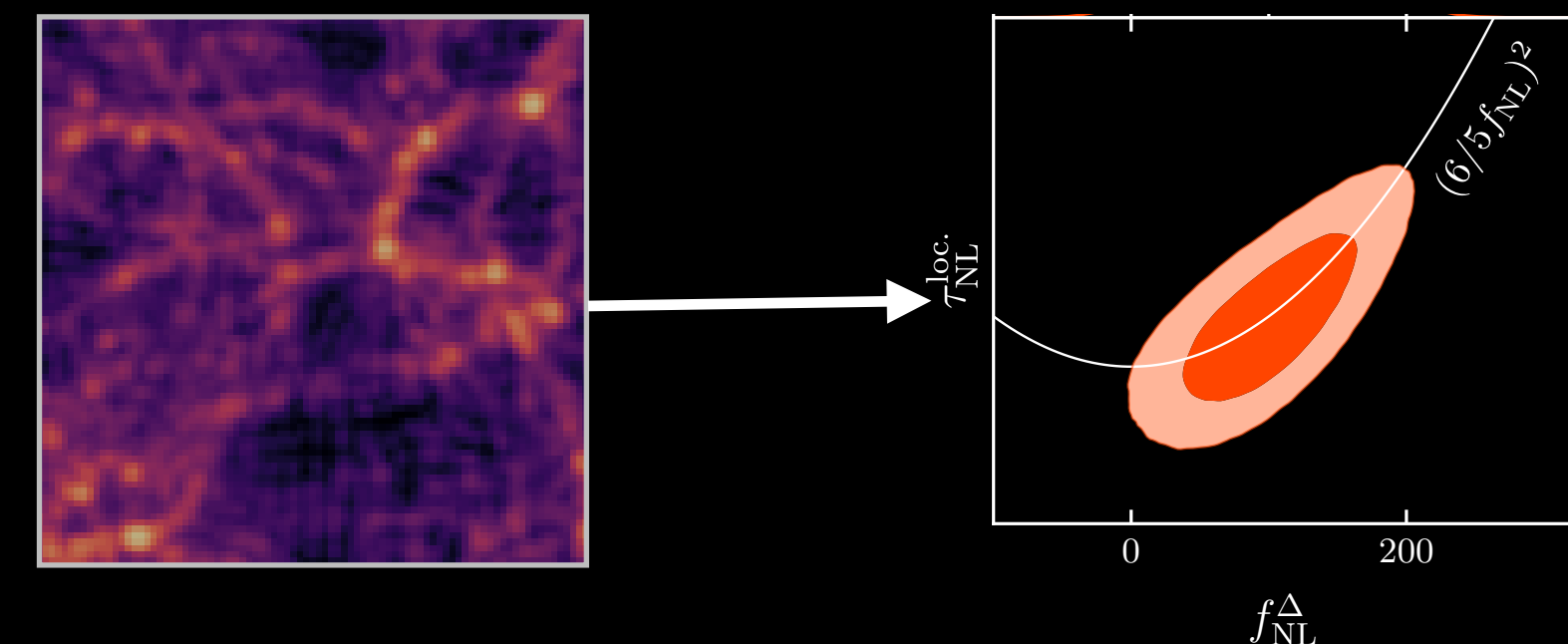


See [Schmittfull+Seljak, 17](#)



See [Anil Kumar++22](#), [Chaussidon++26](#)

Long term: field-level inference/SBI

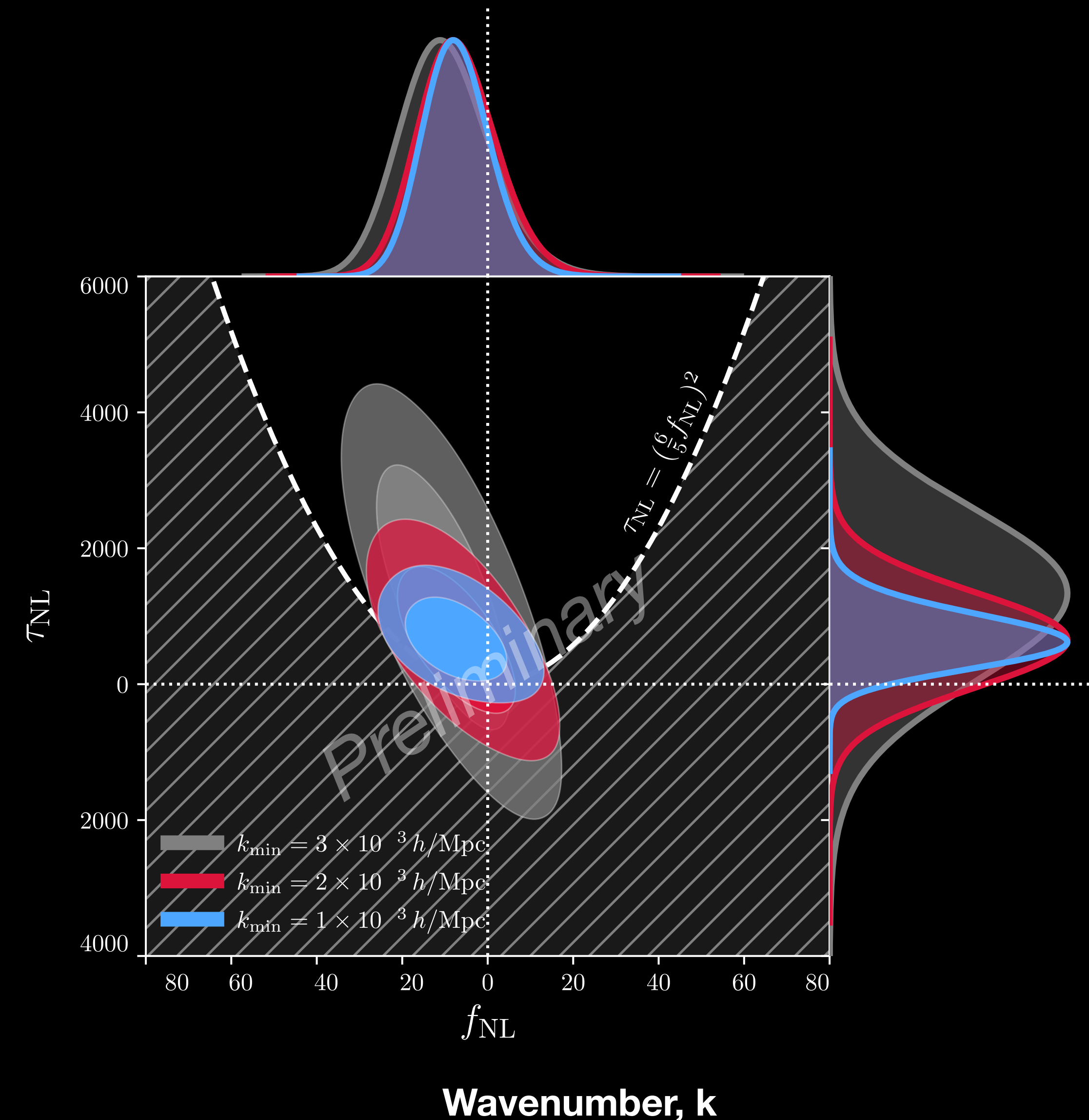


Seljak++17
Schmidt++18
Parker++25
Bayer++26

- Non-Gaussian likelihood
- Low dimensional dataset
- Optimal estimator -> complex weighting
- Real-space systematics (windows, integral constraints, etc.)

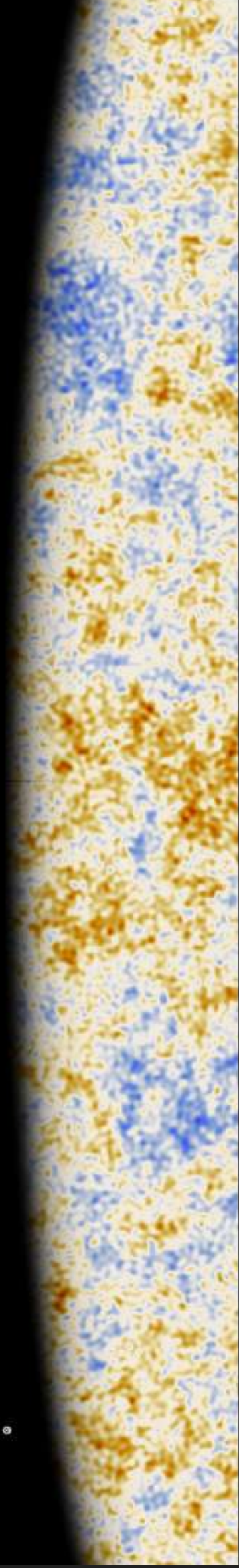
Conclusion

- Scale-dependent bias is more than $f_{\text{NL}}^{\text{loc}}$
- First LSS search for τ_{NL}
 - DESI DR2 QSO analysis: still blinded, but expecting most sensitive search to date
- τ_{NL} pipeline is useful for range of BSM scenarios
- Multifield inflation, dark energy perturbations, large-scale behavior of gravity
- **DESI:** first time LSS is more constraining than CMB for inflation



Overall conclusion

- Cosmology is invaluable probe of BSM physics
- Searching for BSM signals requires developing new pipelines that have applications even if no new physics is detected
 - **Example 1:** searching for axions led to stringent constraints on axions as well as component-separation pipeline with new methods for foreground cleaning that is now being used to produce component-separated maps for SO
 - **Example 2:** constraining τ_{NL} with DESI has led to the most sensitive search for multifield inflation to date. Analysis is still blinded, but in absence of detection our new pipeline is starting point for range of BSM searches



Thank you!

Sam Goldstein