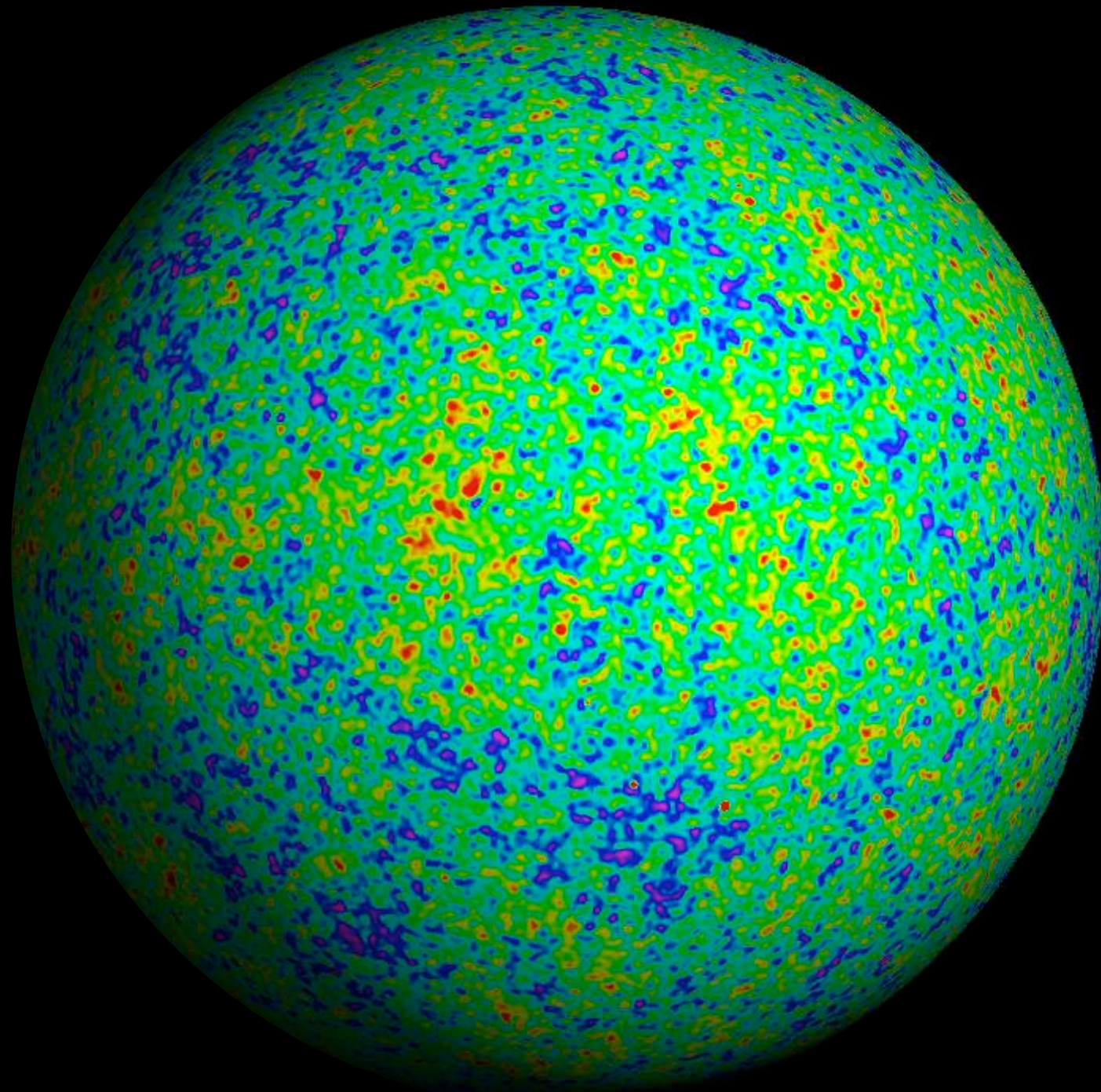


Recent Progress in 21 cm Cosmology with HERA

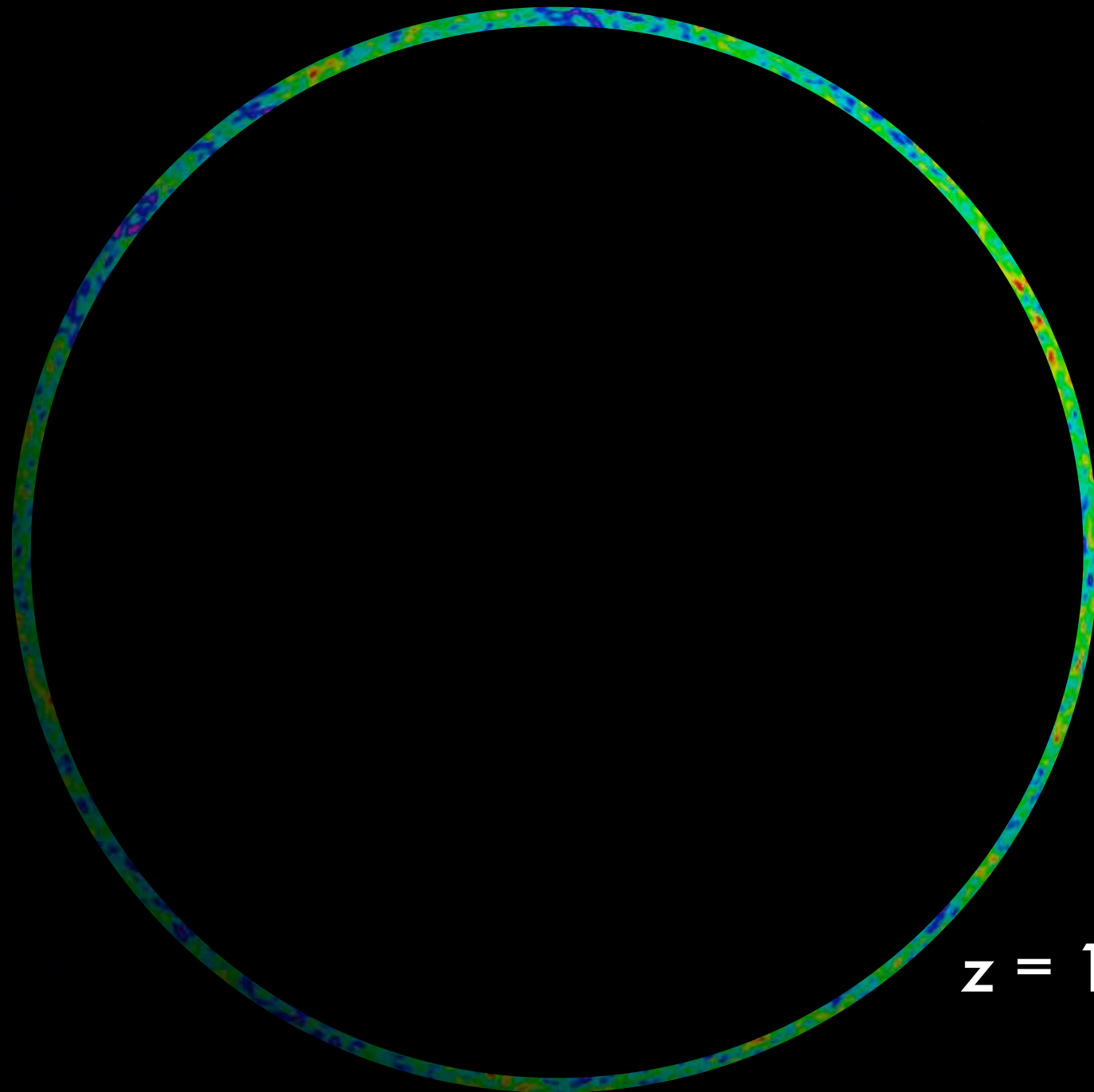
Josh Dillon
UC Berkeley

How can we map out
our whole universe?

With the Cosmic Microwave Background...



...we only get a thin shell at high redshift.



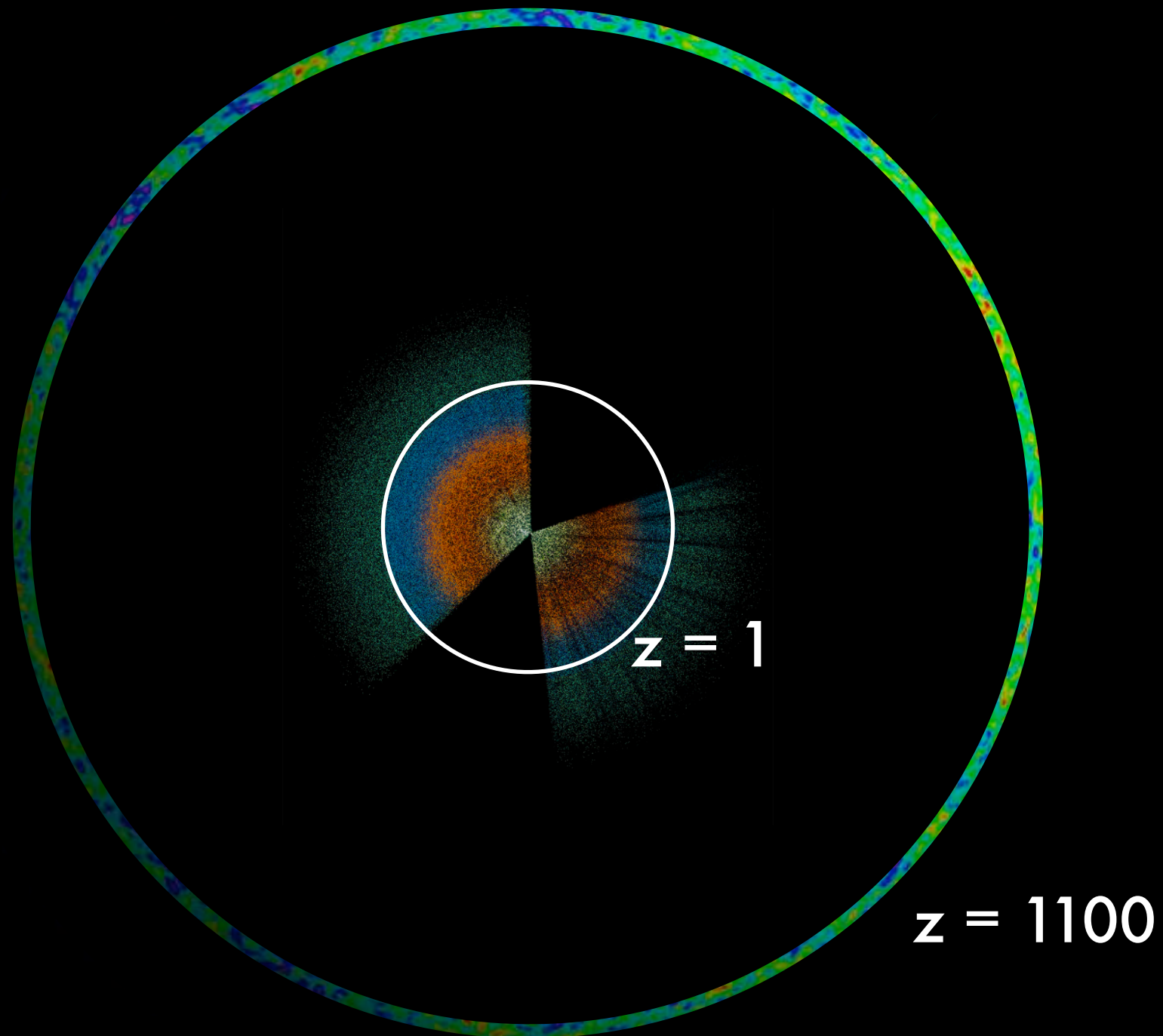
$z = 1100$

Even with huge surveys of galaxies...

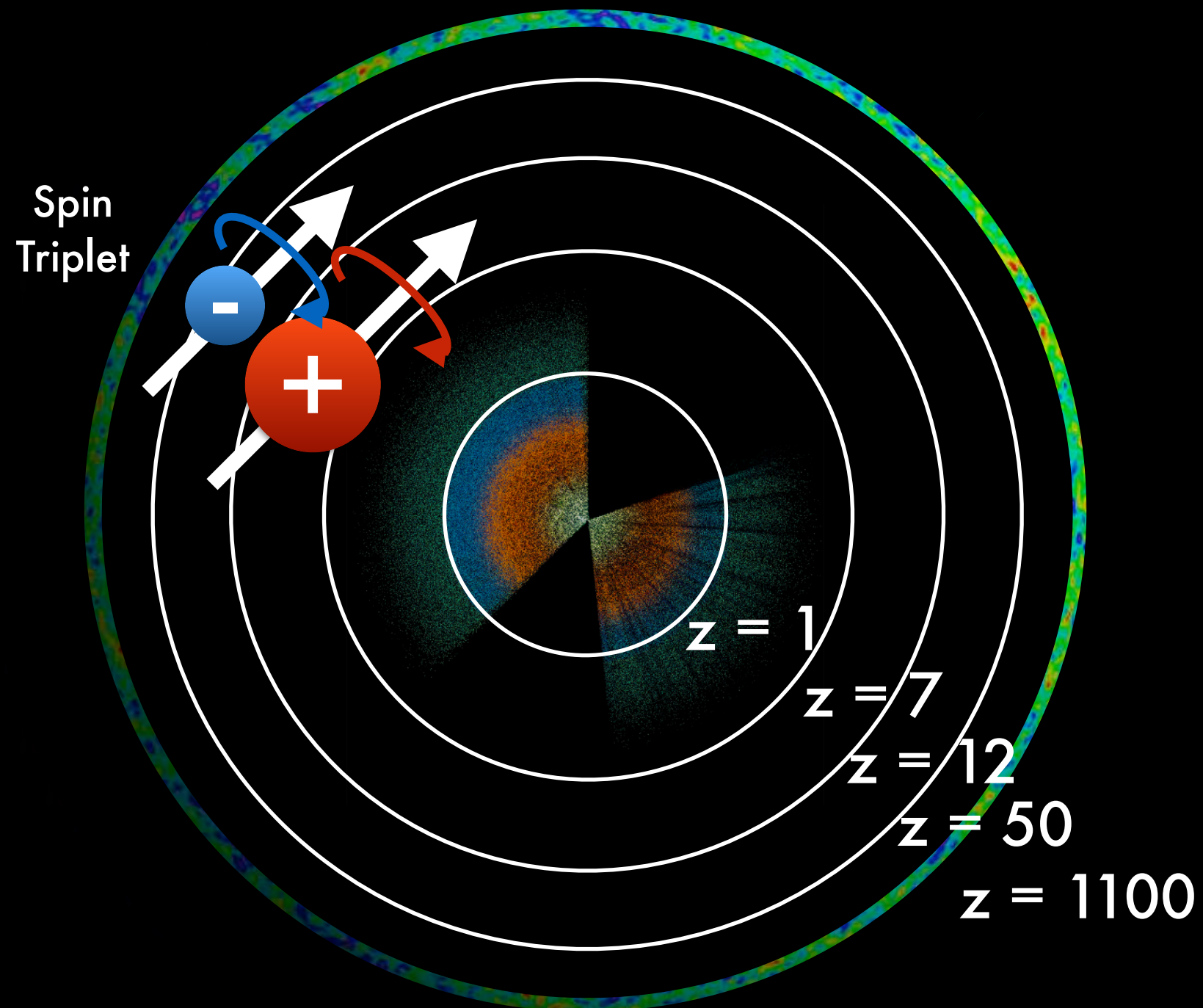


$z = 1100$

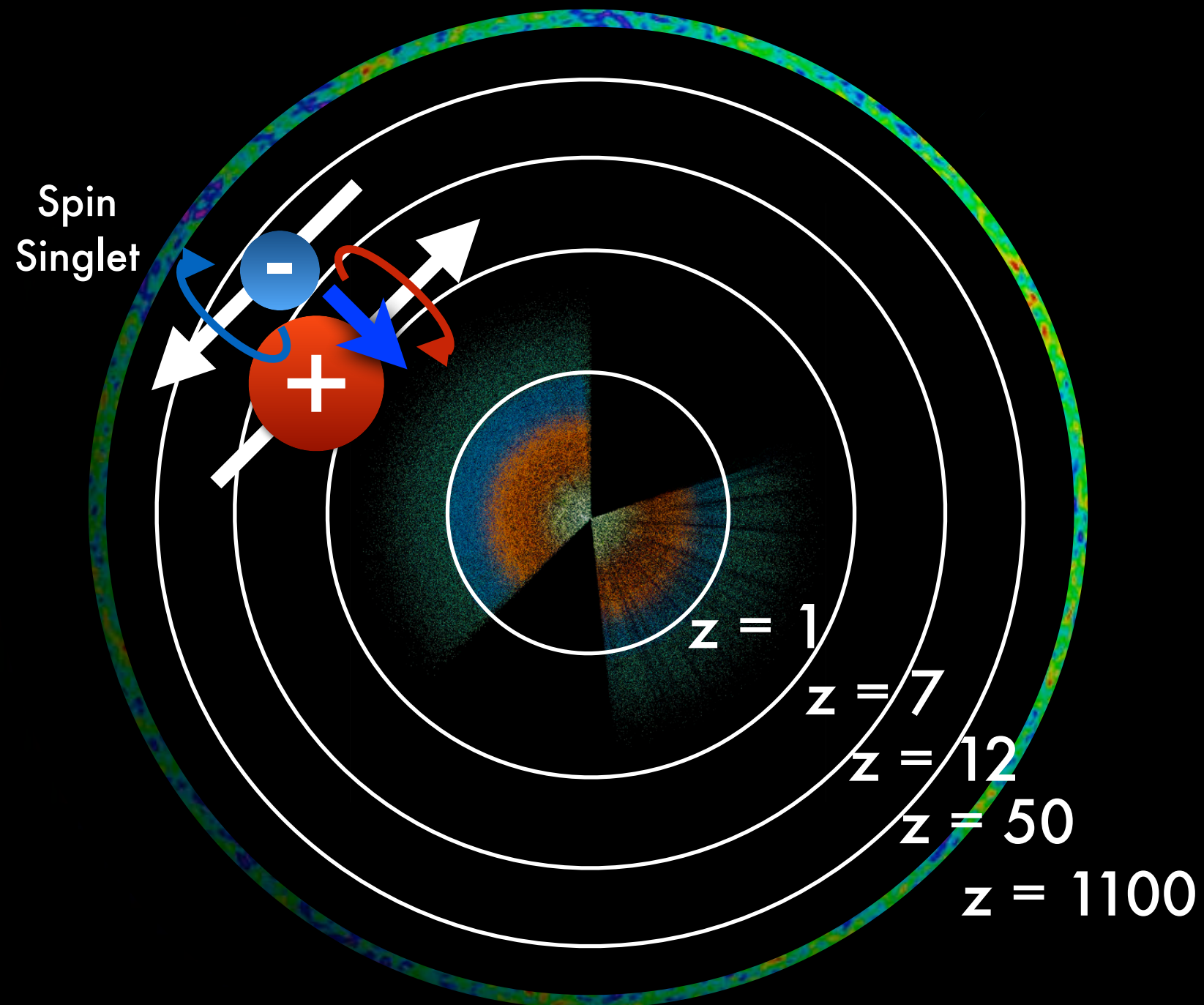
...we only map out the local universe.



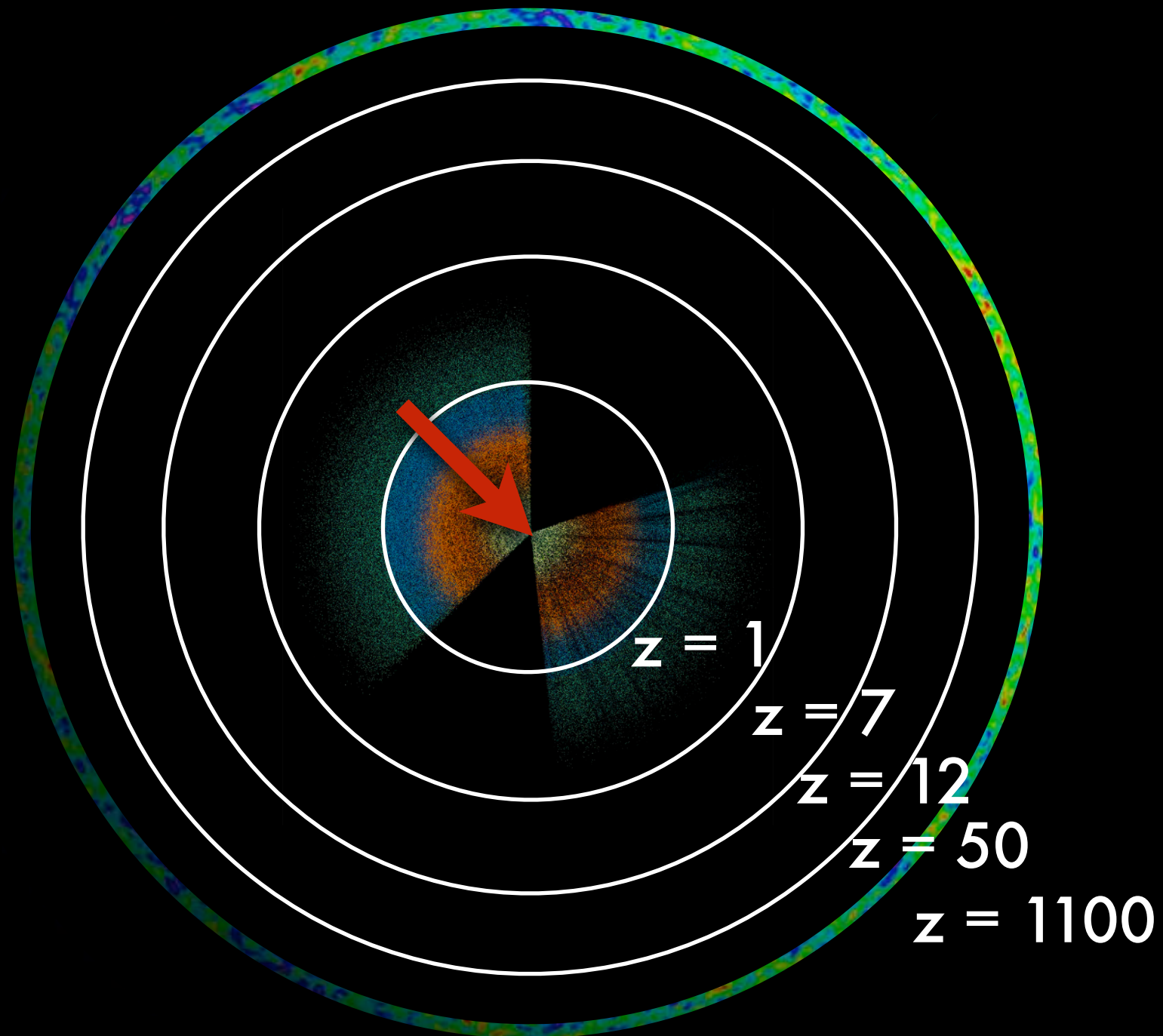
But using the 21 cm hydrogen line...



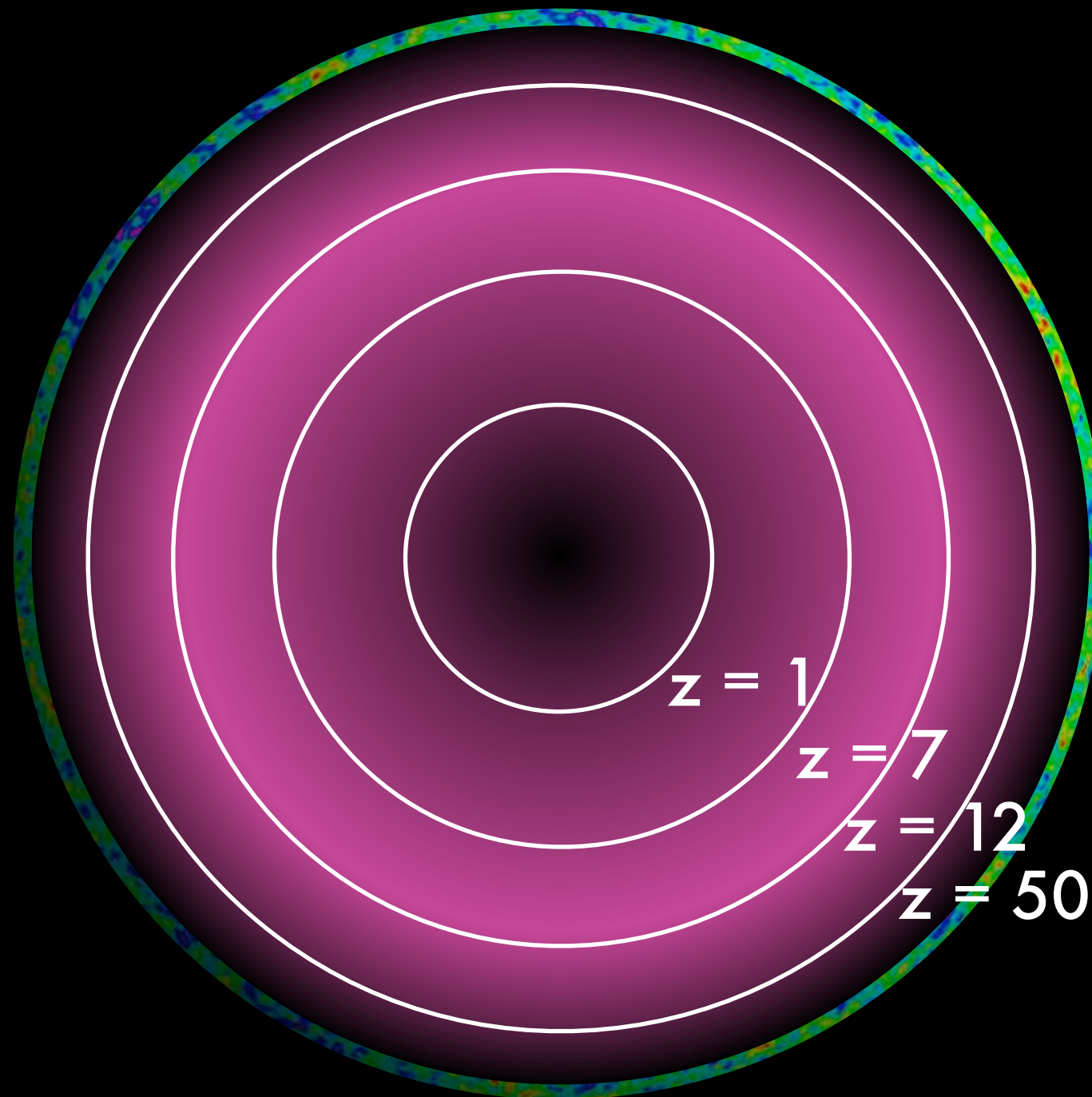
But using the 21 cm hydrogen line...



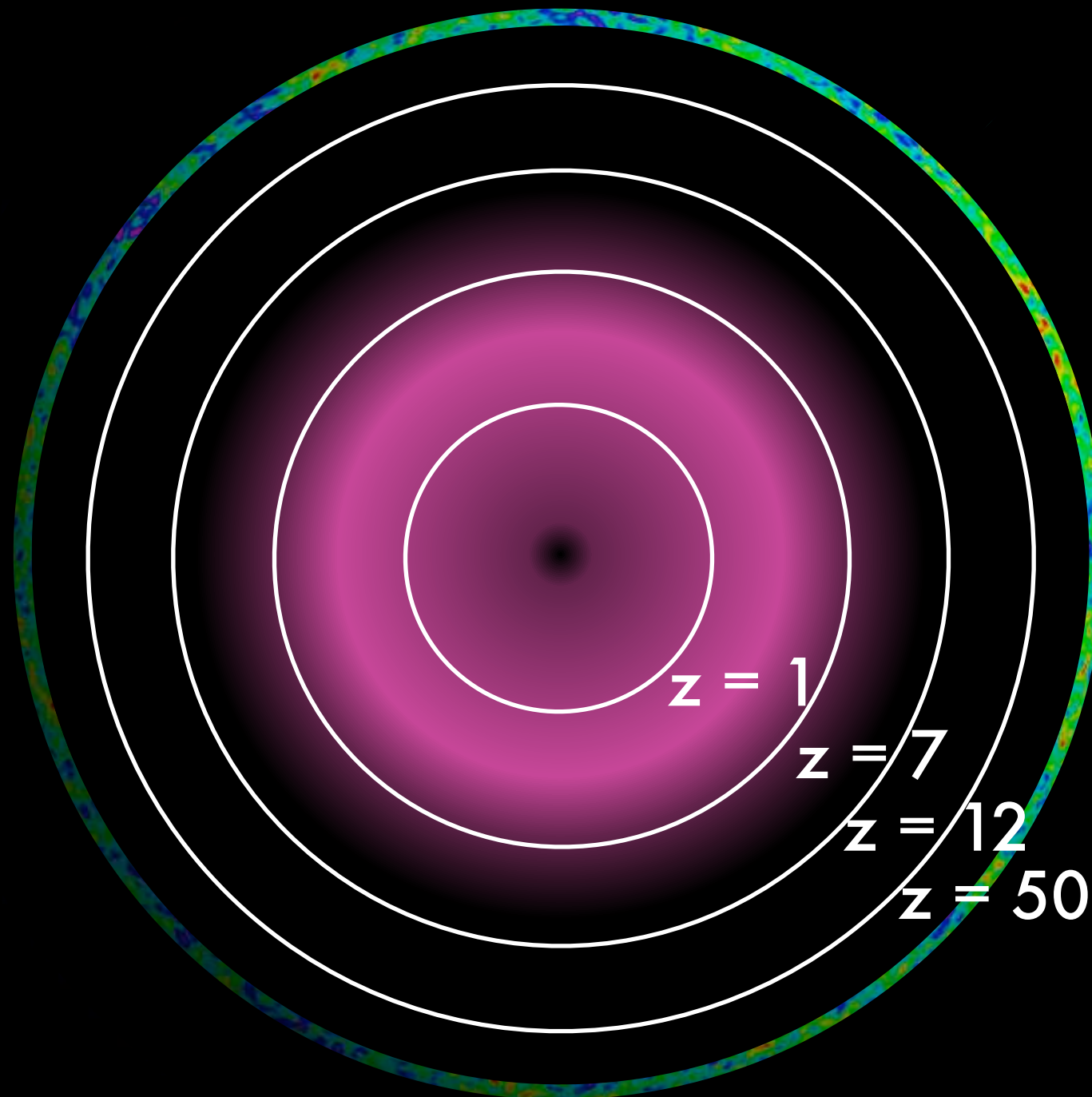
But using the 21 cm hydrogen line...



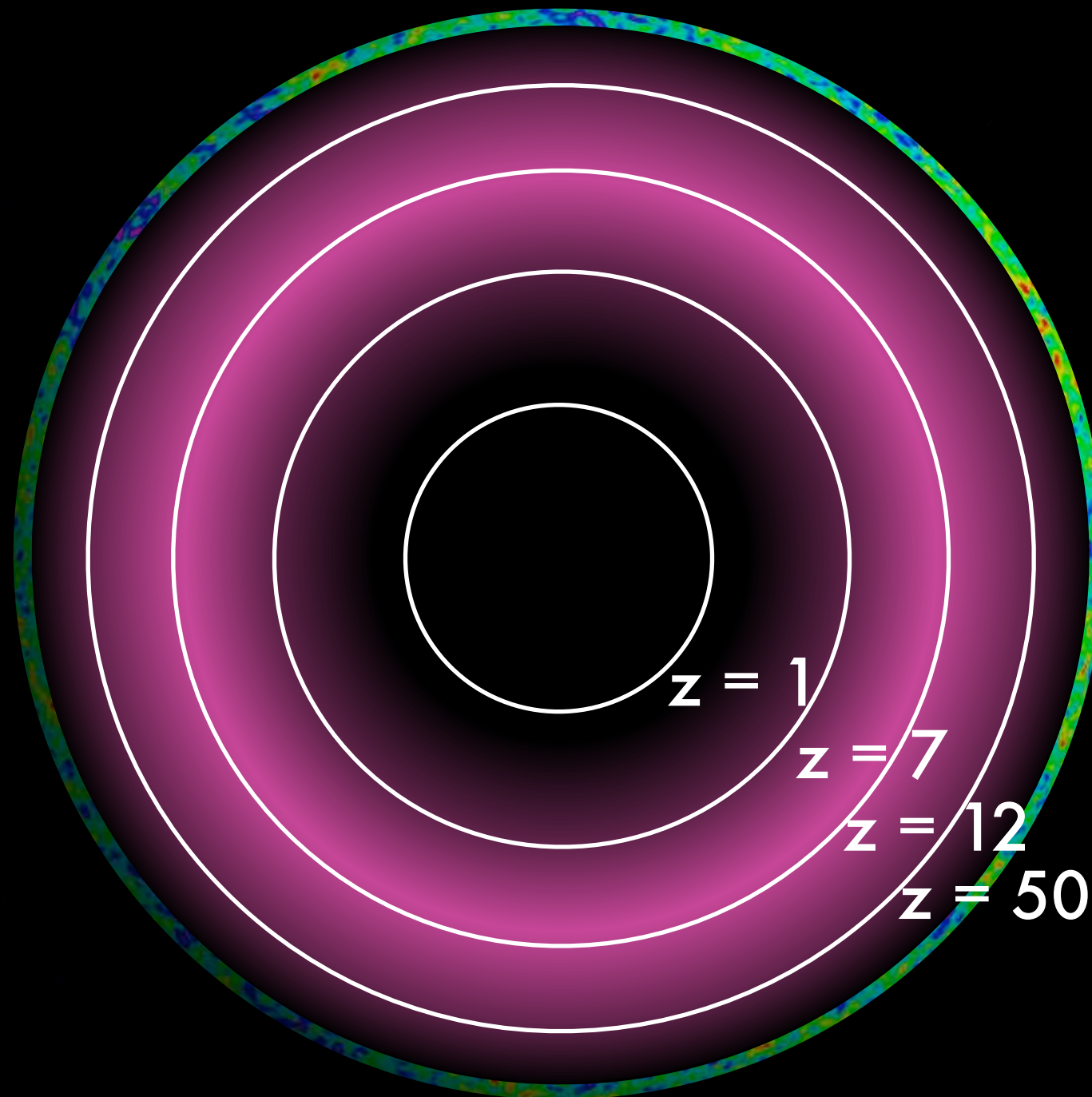
...a huge volume of the universe
can be directly probed ($z \lesssim 200$).



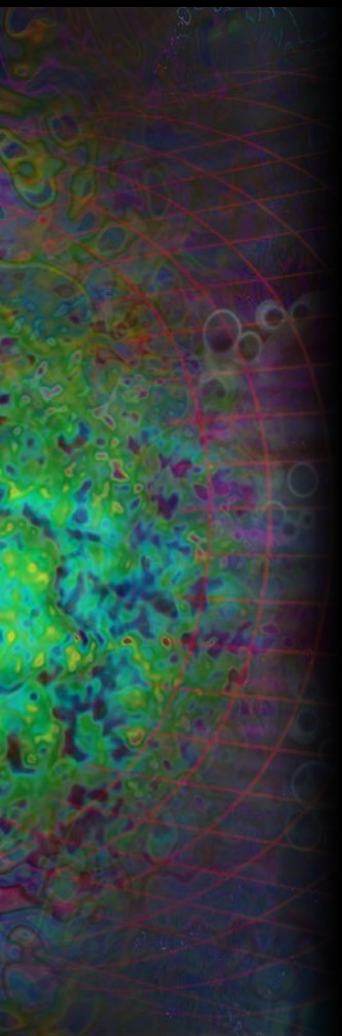
At “low” redshift, 21 cm can probe large-scale structure by tracing atomic gas in galaxies.



At $z \gtrsim 6$, we can map the universe as it undergoes a dramatic transformation.

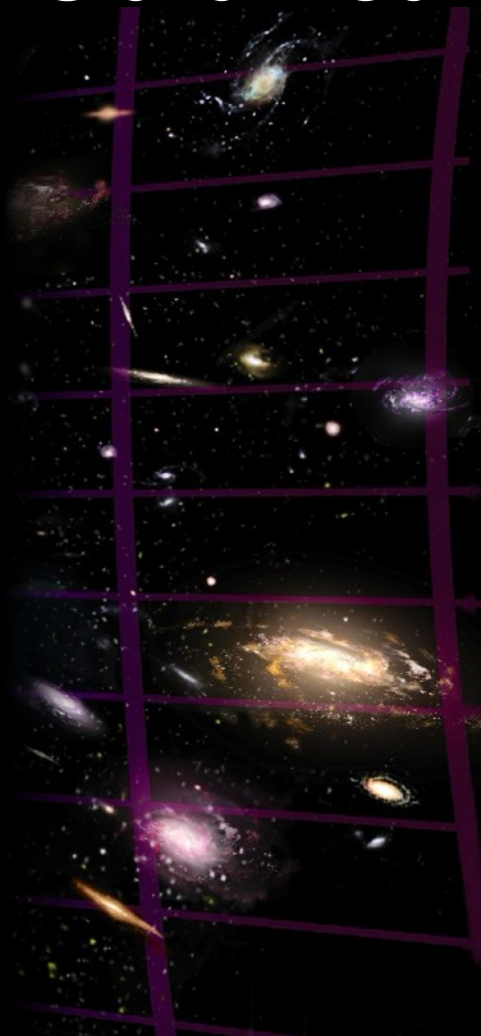


CMB



$z = 1100$

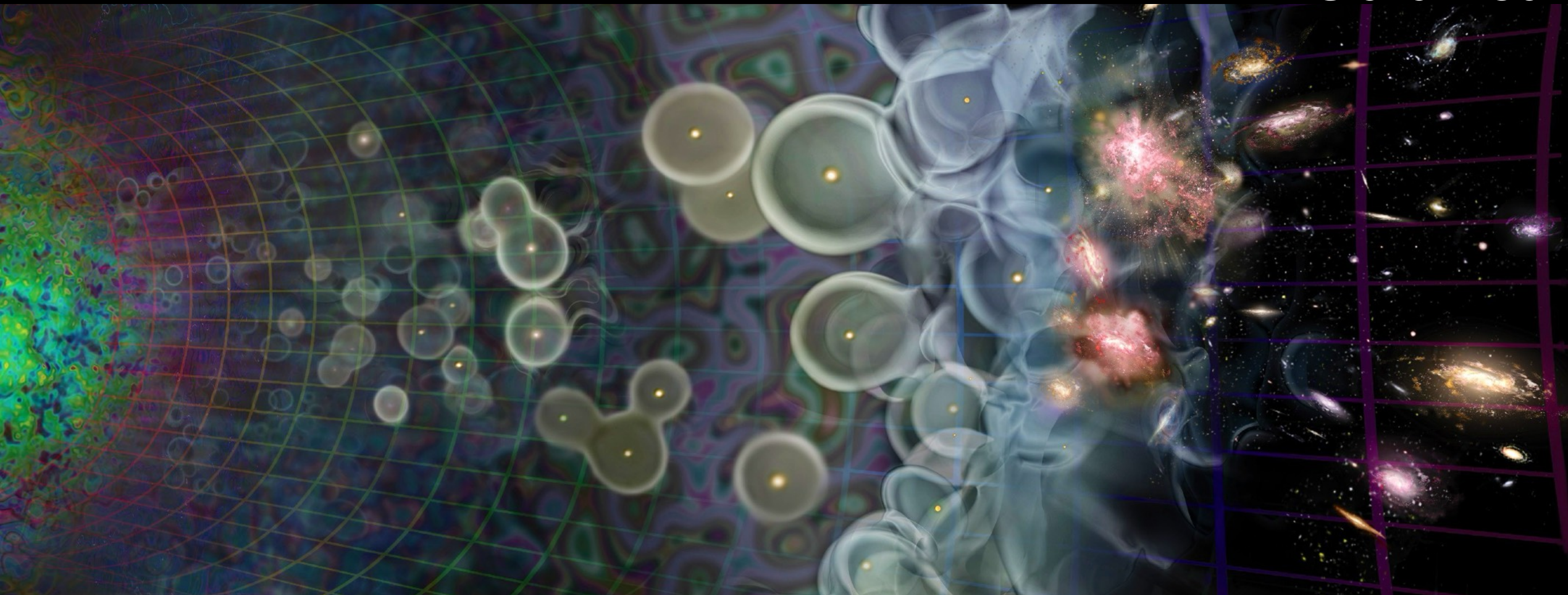
Modern
Galaxies



$z < 6$

CMB

Modern
Galaxies



$z = 1100$

$z < 6$

Image: Avi Loeb

CMB

Modern
Galaxies

Dark Ages

$z = 1100$

$z < 6$

Image: Avi Loeb



CMB

Modern
Galaxies

Dark Ages

First Stars

$z = 1100$

$z \approx 50$

$z < 6$

Image: Avi Loeb

CMB

Modern
Galaxies

Dark Ages

First Black Holes

First Stars

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$z \approx 50$

$z < 6$

Image: Avi Loeb

CMB

Modern
Galaxies

Dark Ages

First Black Holes

First Stars

The Epoch of
Reionization

$z = 1100$

$z \approx 50$

$z \approx 8$

$z < 6$

Image: Avi Loeb

The background of the slide is a composite image representing the Cosmic Dawn. On the left, there is a colorful, swirling pattern in shades of green, blue, and purple, overlaid with a grid of thin, intersecting lines in red, green, and blue. This pattern transitions into a dark, deep space on the right, filled with numerous galaxies of various shapes and sizes, some appearing as bright, glowing clouds of gas and dust. The overall effect is a sense of vastness and the early stages of cosmic evolution.

First Black Holes

The Cosmic Dawn

First Stars

The Epoch of
Reionization

We already have some
clues about reionization.

Quasar Lyman- α spectra tell us that reionization ended around redshift 6.

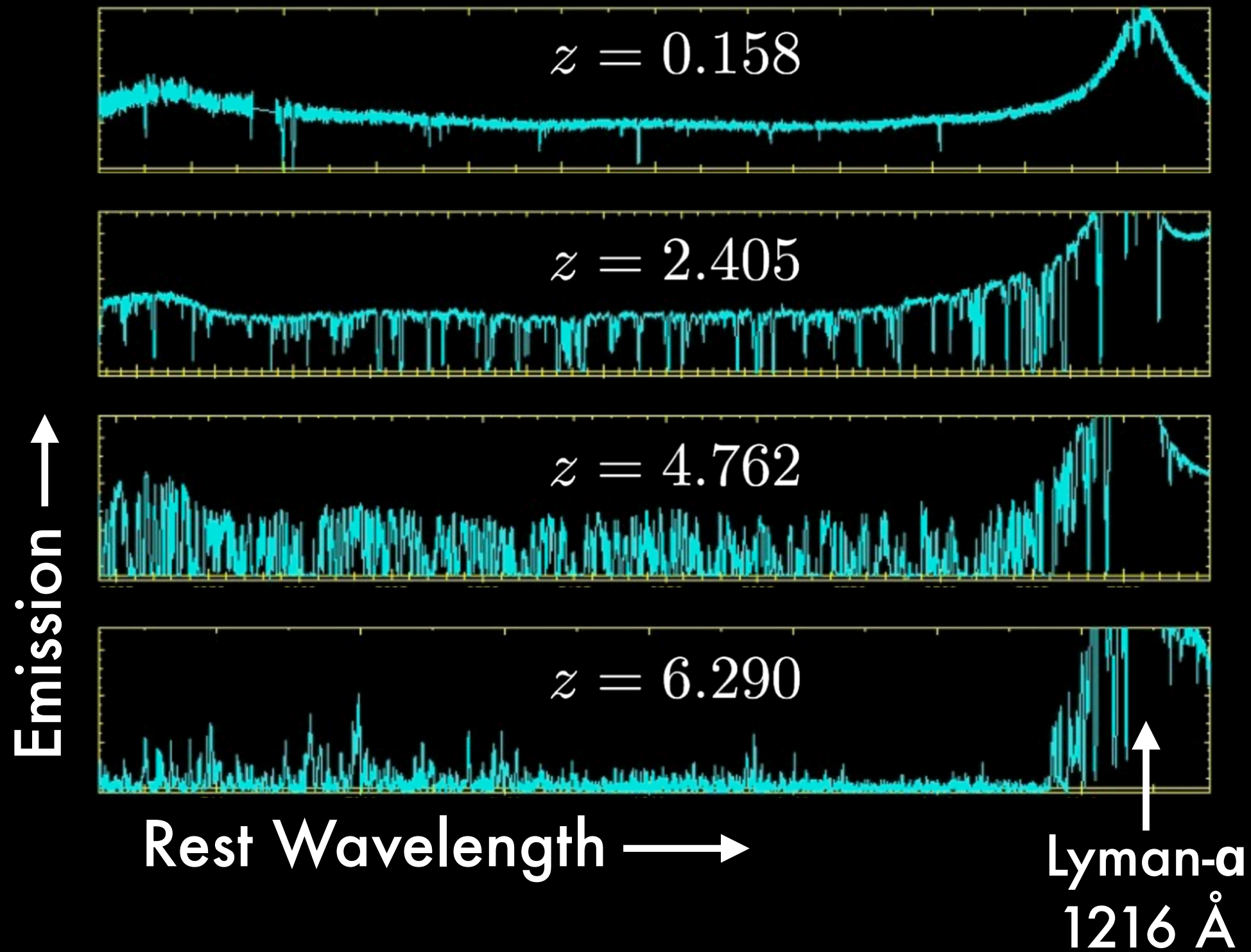
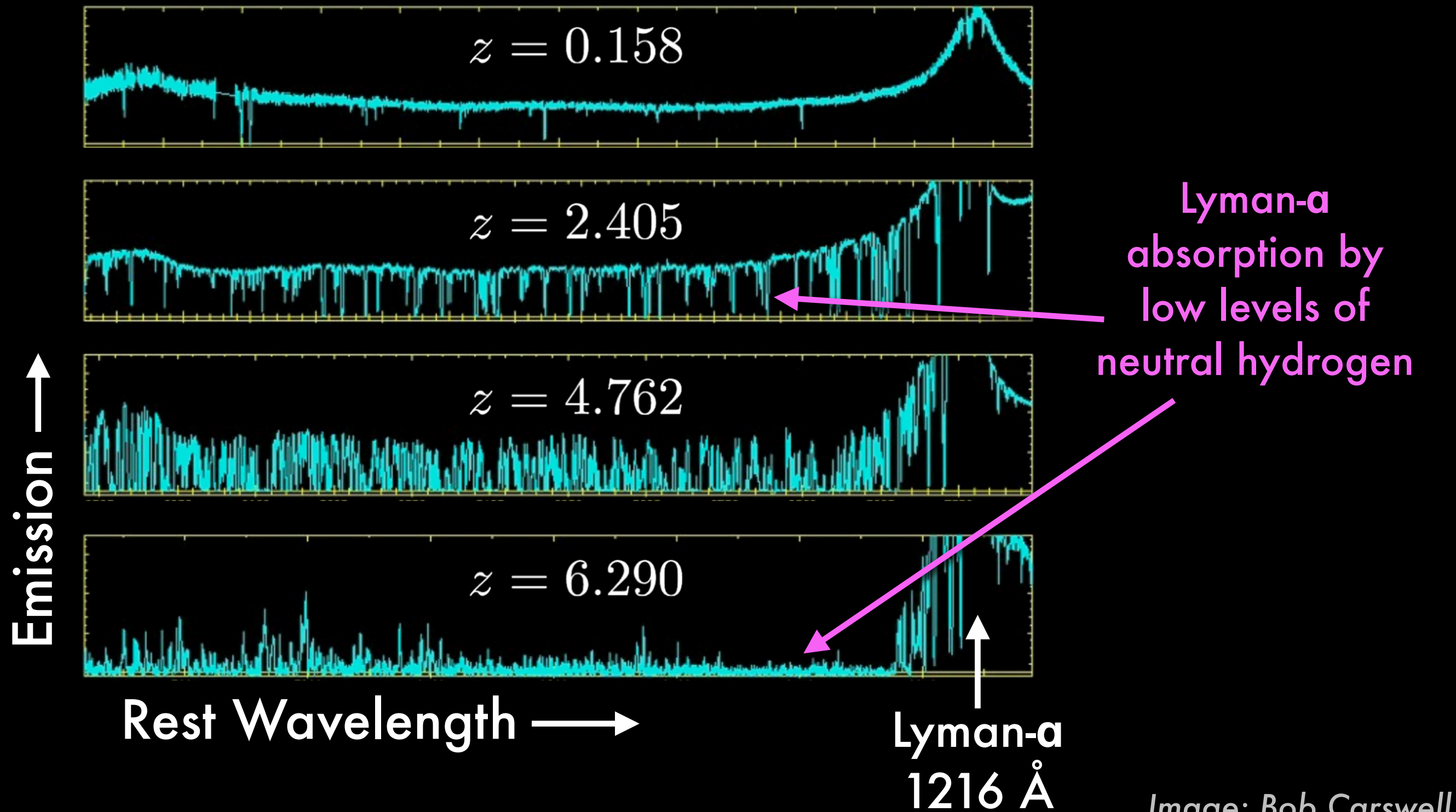


Image: Bob Carswell

Quasar Lyman- α spectra tell us that reionization ended around redshift 6.



We also get an integral constraint on reionization from CMB polarization.

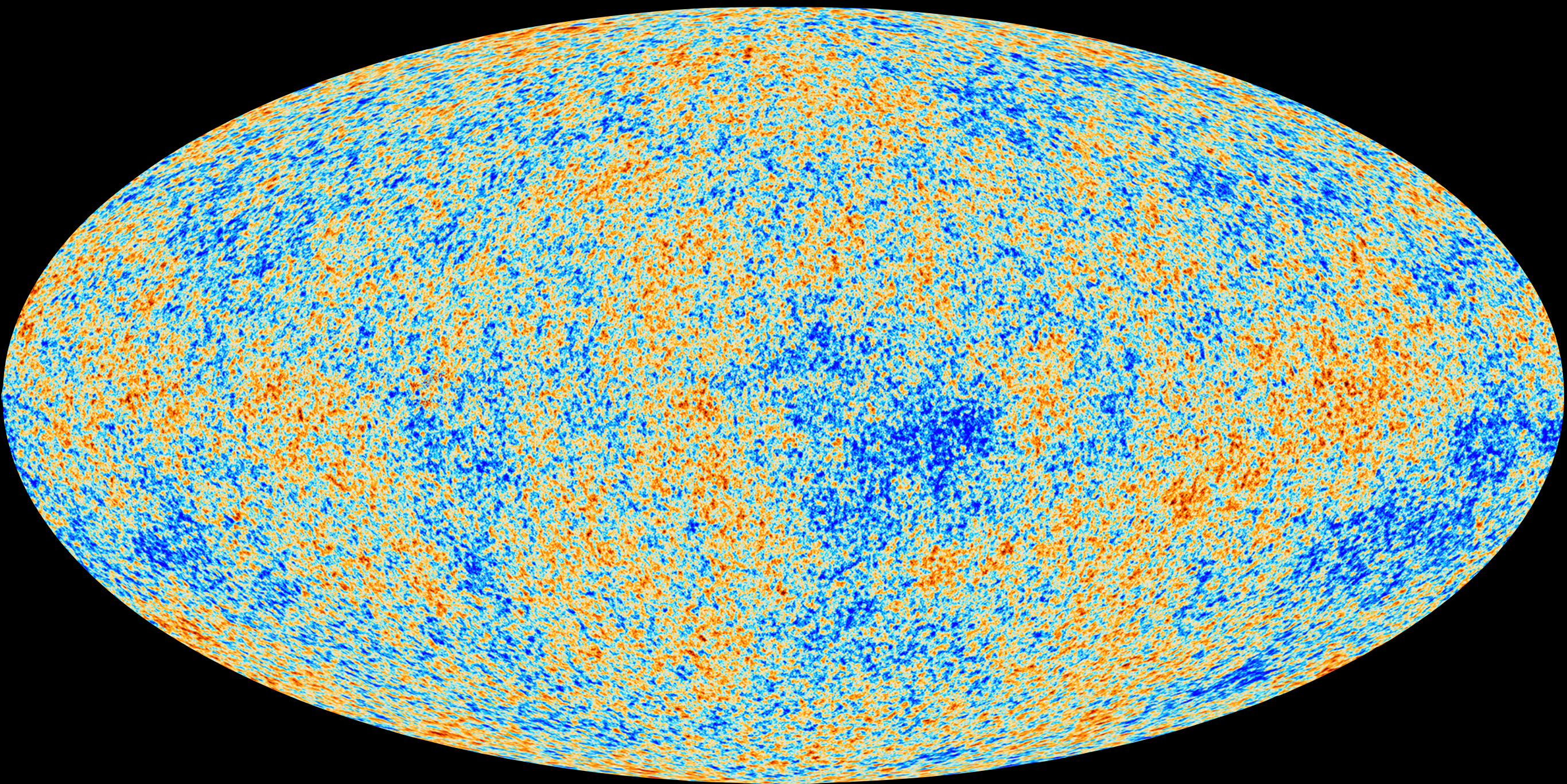
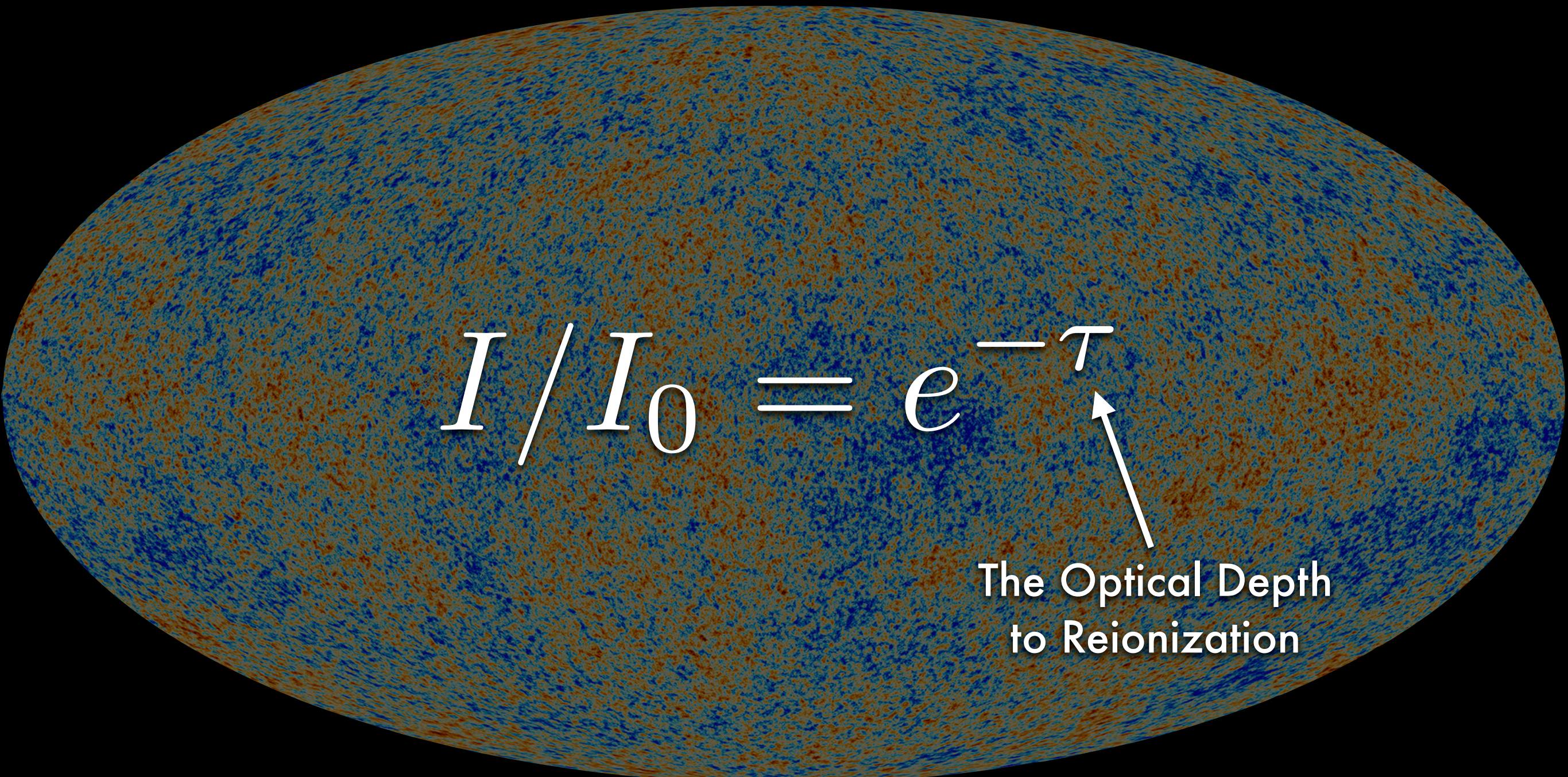


Image: Planck Collaboration

We also get an integral constraint on reionization from CMB polarization.


$$I/I_0 = e^{-\tau}$$

The Optical Depth
to Reionization

So we think reionization
looked something like this...

So we think reionization
looked something like this...

$$\delta T_{21 \text{ cm}} \propto (1 + \delta) \left[1 - \frac{T_{\text{CMB}}}{T_s} \right] x_{\text{HI}}$$

↑
21 cm Brightness
Temperature

Overdensity of
Hydrogen

21 cm Brightness
Temperature

$$\delta T_{21 \text{ cm}} \propto (1 + \delta) \left[1 - \frac{T_{\text{CMB}}}{T_s} \right] x_{\text{HI}}$$

The diagram shows the equation for the 21 cm brightness temperature fluctuation, $\delta T_{21 \text{ cm}}$, which is proportional to the overdensity of hydrogen, $(1 + \delta)$, multiplied by a bracketed term $\left[1 - \frac{T_{\text{CMB}}}{T_s}\right]$, and then multiplied by the neutral hydrogen fraction, x_{HI} . Three pink arrows point to specific parts of the equation: one points up to $\delta T_{21 \text{ cm}}$ from the label '21 cm Brightness Temperature', another points down to δ from the label 'Overdensity of Hydrogen', and a third points up to T_s from the label 'Spin Temperature'.

$$\delta T_{21 \text{ cm}} \propto (1 + \delta) \left[1 - \frac{T_{\text{CMB}}}{T_s} \right] x_{\text{HI}}$$

21 cm Brightness Temperature

Overdensity of Hydrogen

Spin Temperature

Overdensity of Hydrogen

Neutral Fraction

$$\delta T_{21 \text{ cm}} \propto (1 + \delta) \left[1 - \frac{T_{\text{CMB}}}{T_s} \right] x_{\text{HI}}$$

21 cm Brightness Temperature

Spin Temperature

The diagram illustrates the equation for the 21 cm brightness temperature. The equation is $\delta T_{21 \text{ cm}} \propto (1 + \delta) \left[1 - \frac{T_{\text{CMB}}}{T_s} \right] x_{\text{HI}}$. Four variables are highlighted with arrows: δ is labeled 'Overdensity of Hydrogen', x_{HI} is labeled 'Neutral Fraction', T_s is labeled 'Spin Temperature', and $\delta T_{21 \text{ cm}}$ is labeled '21 cm Brightness Temperature'.

The brightness temperature probes different physics at different times.

Dark Ages

First Black Holes

$$\delta T_{21\text{ cm}} \propto (1 + \delta) \left[1 - \frac{T_{\text{CMB}}}{T_s} \right] x_{\text{HI}}$$

First Stars

The Epoch of
Reionization

$z = 1100$

$z \approx 50$

$z \approx 8$

$z < 6$

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First Stars

The Epoch of Reionization

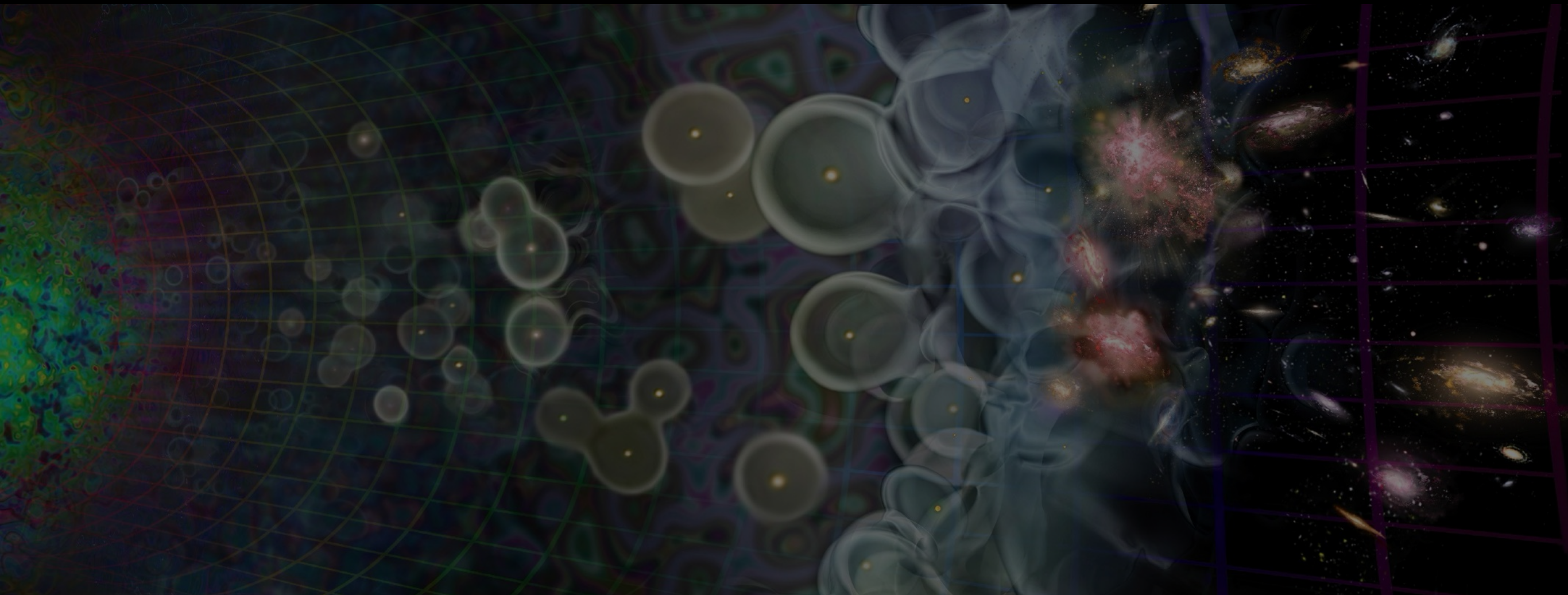
$z = 1100$

$z \approx 50$

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$z < 6$

There's still a lot of open
astrophysical questions.



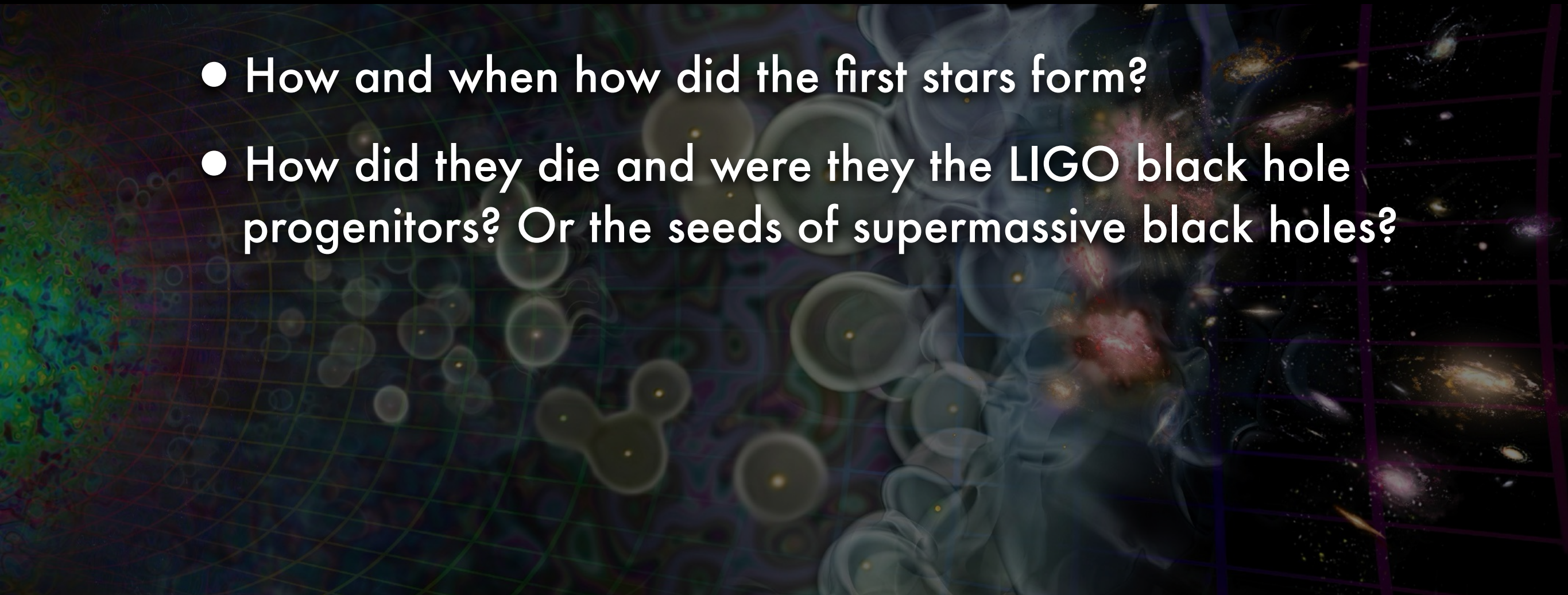
There's still a lot of open astrophysical questions.

- How and when how did the first stars form?



There's still a lot of open astrophysical questions.

- How and when how did the first stars form?
- How did they die and were they the LIGO black hole progenitors? Or the seeds of supermassive black holes?



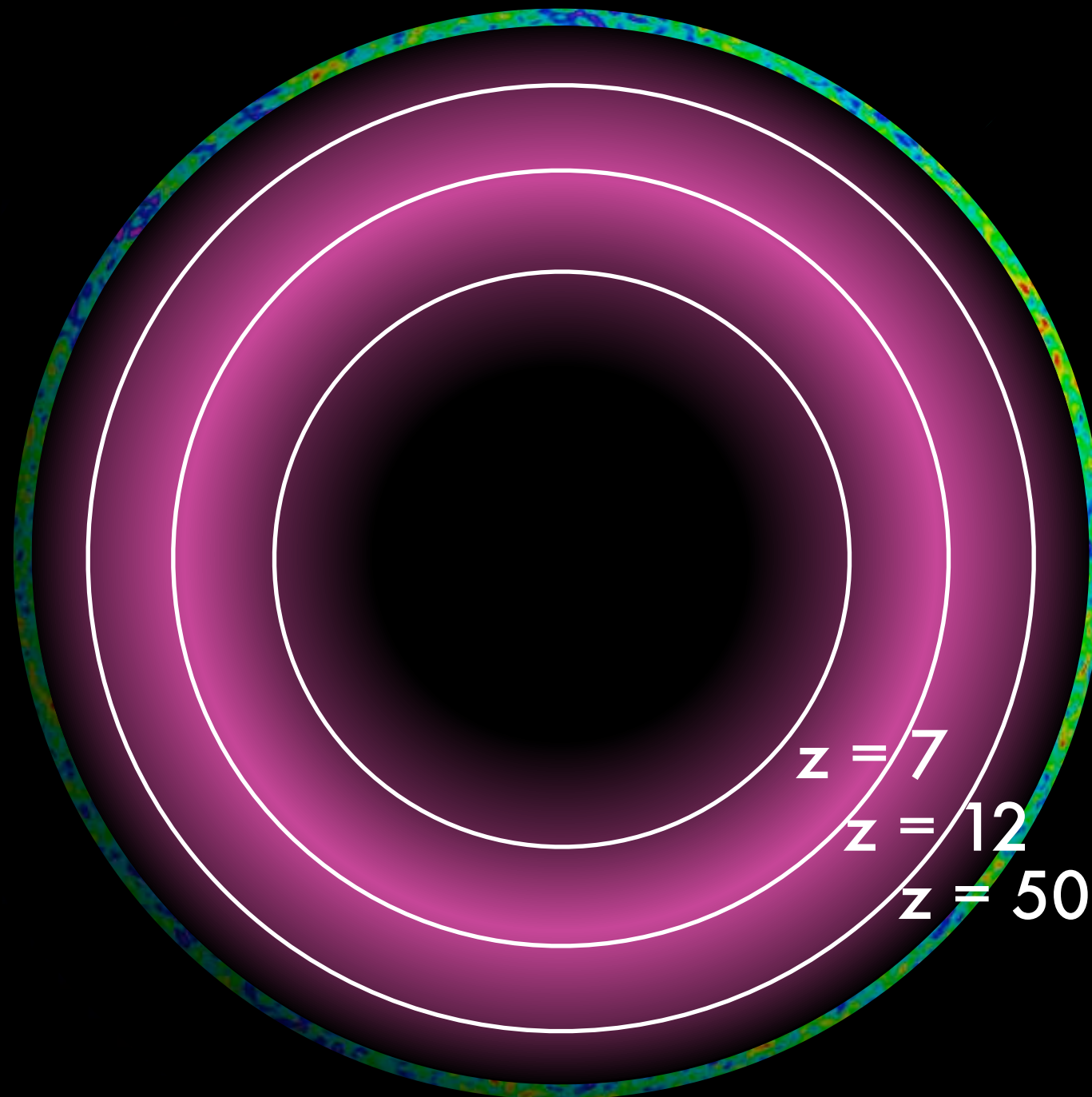
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There's still a lot of open astrophysical questions.

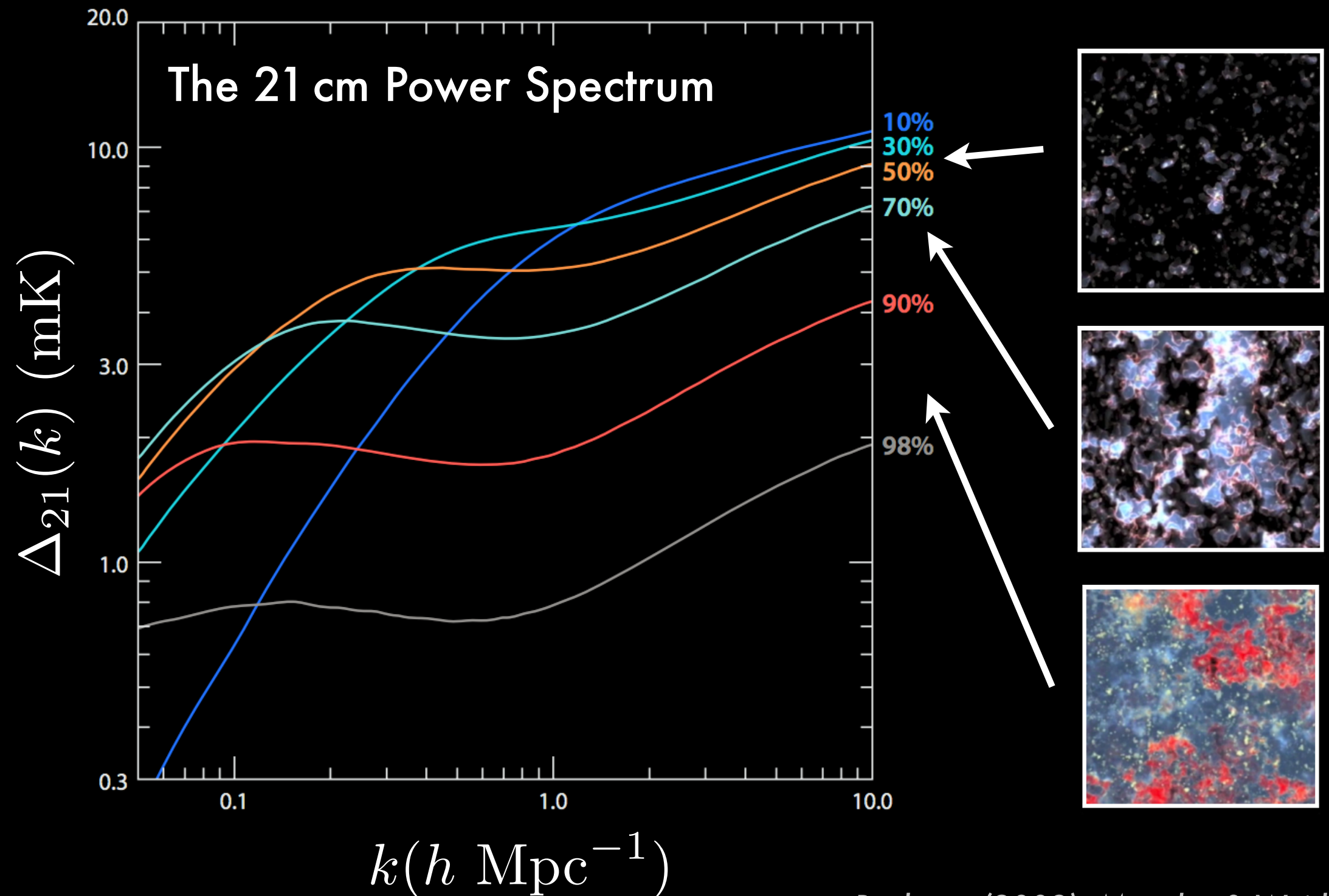
- How and when how did the first stars form?
- How did they die and were they the LIGO black hole progenitors? Or the seeds of supermassive black holes?
- What determined the thermal history of the intergalactic medium? Are there new physics at play?
- What reionized the universe and when? What role did very high redshift galaxies found by JWST play?

The Cosmic Dawn is roughly half of the comoving volume of the observable universe.



**The first detection
will be statistical.**

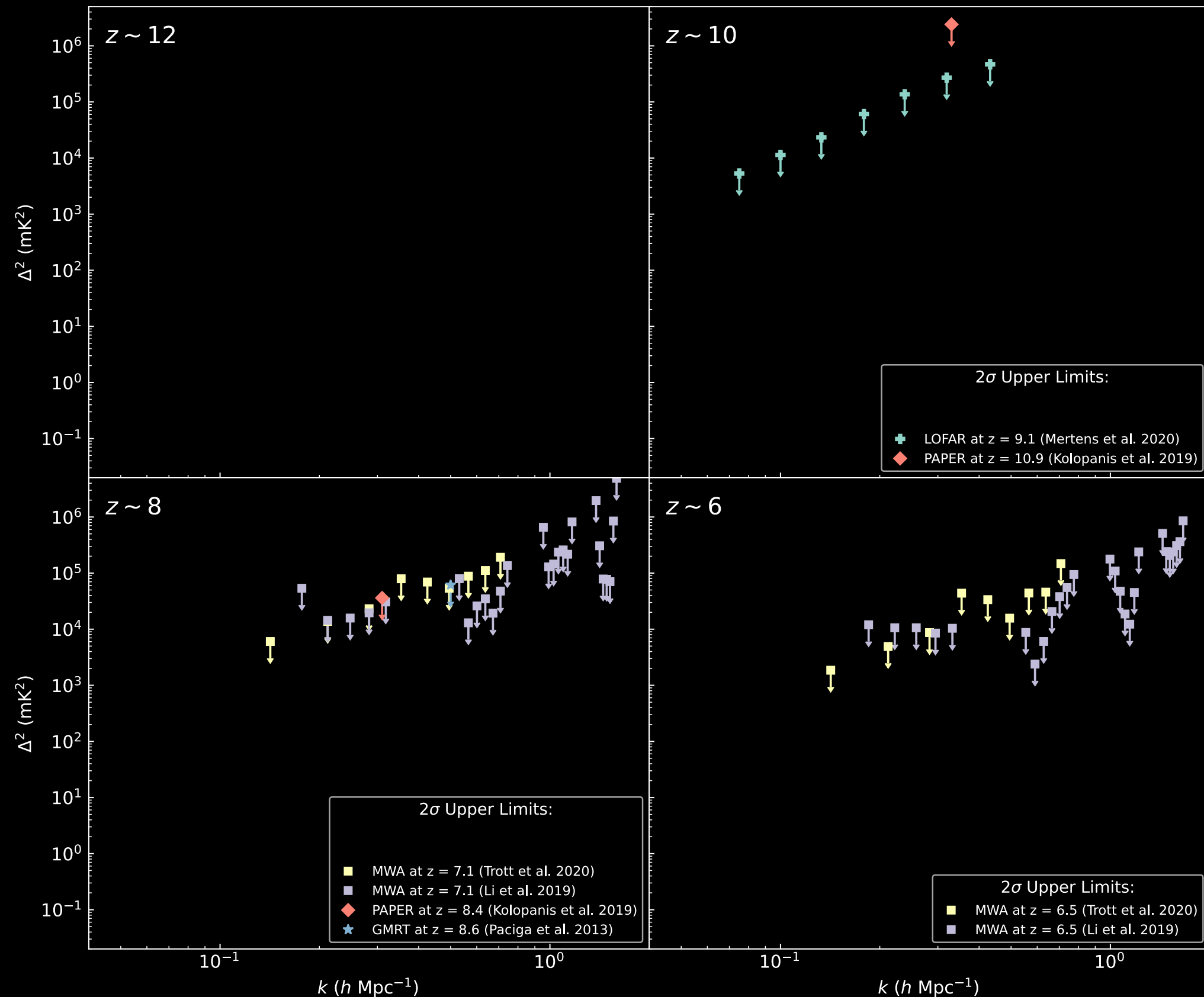
Our best probe is the power spectrum – the evolution of brightness temperature fluctuations as a function of time and spatial scale.



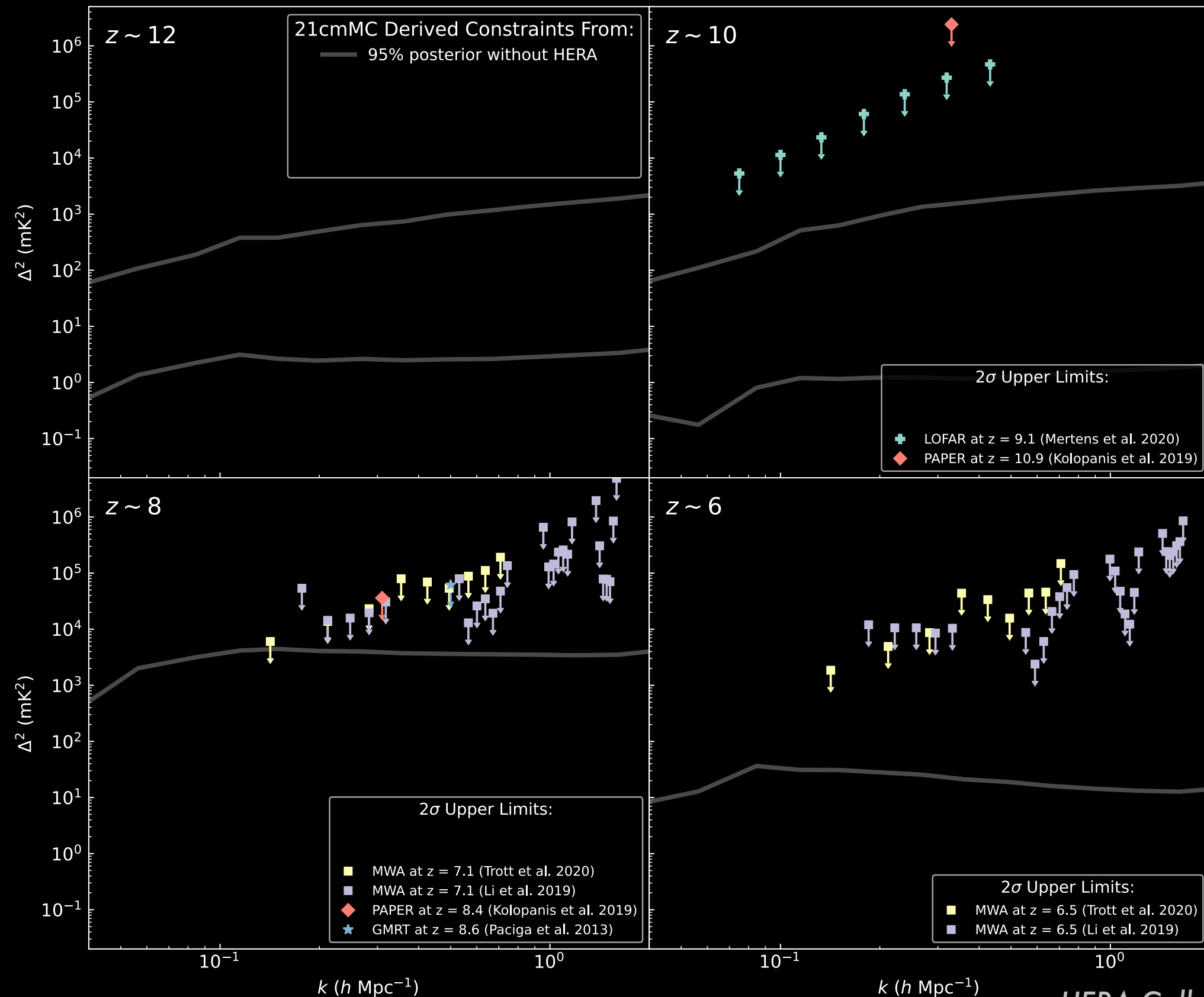
The first generation of interferometers for 21 cm cosmology got us started, deploying different strategies.



Before HERA, the state of the field was upper limits still quite far from the expected EoR signal.



A wide range of power spectra were still possible, even with CMB and galaxy luminosity function constraints.





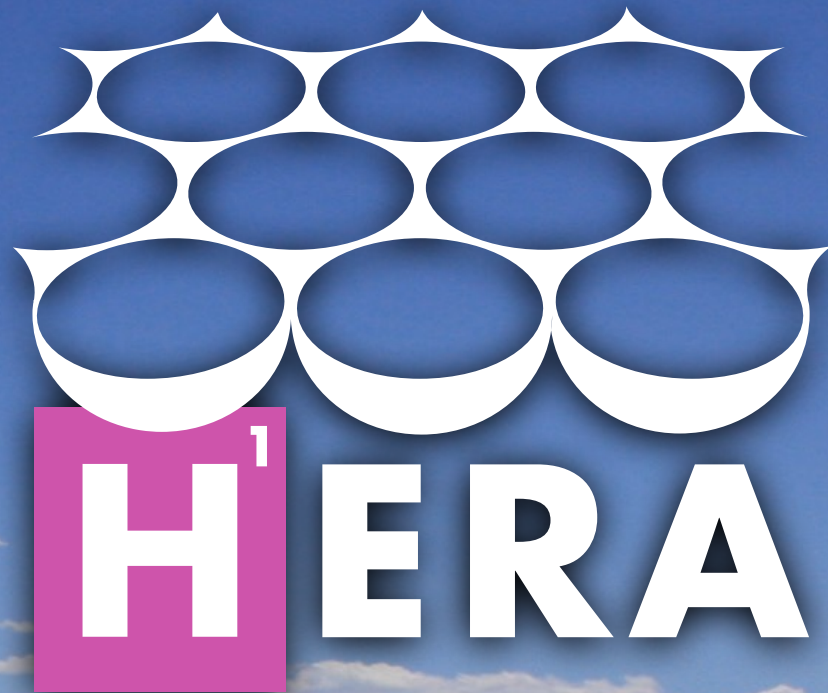
Google Earth
Data SIO, NOAA, U.S. Navy

So we went bigger...

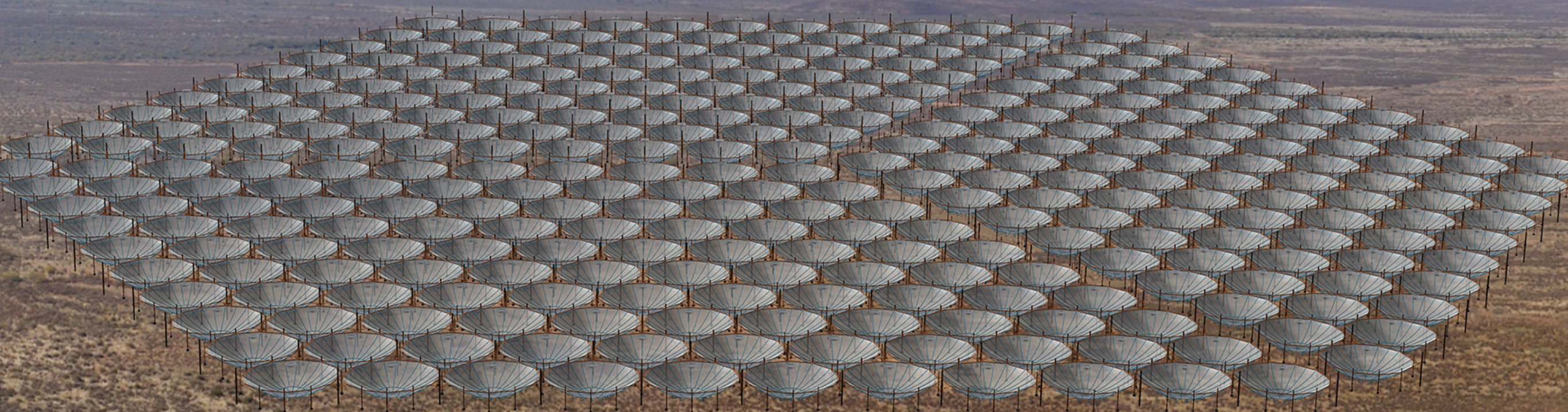


Google Earth
Data SIO, NOAA, U.S. Navy

So we went bigger...

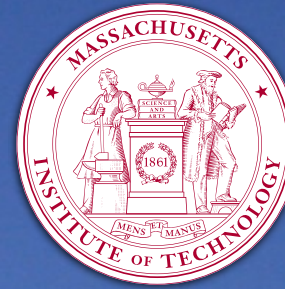


The Hydrogen Epoch of Reionization Array





CAL POLY POMONA



SCUOLA
NORMALE
SUPERIORE



UNIVERSITY OF
CAMBRIDGE



Penn
UNIVERSITY of PENNSYLVANIA



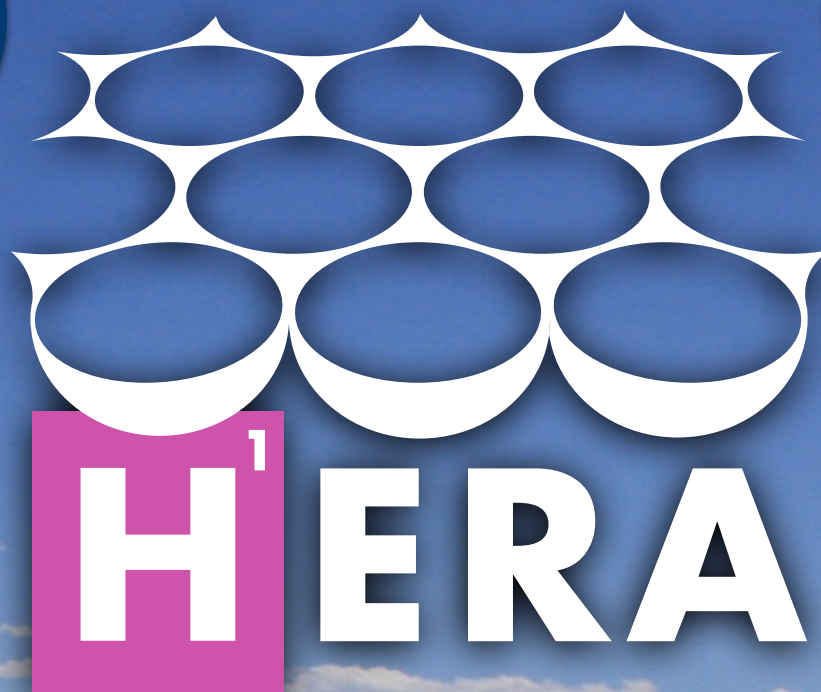
McGill
UNIVERSITY



Queen Mary
University of London

SARAO

South African Radio
Astronomy Observatory



BROWN



SKA AFRICA
SQUARE KILOMETRE ARRAY



The Hydrogen Epoch of Reionization Array



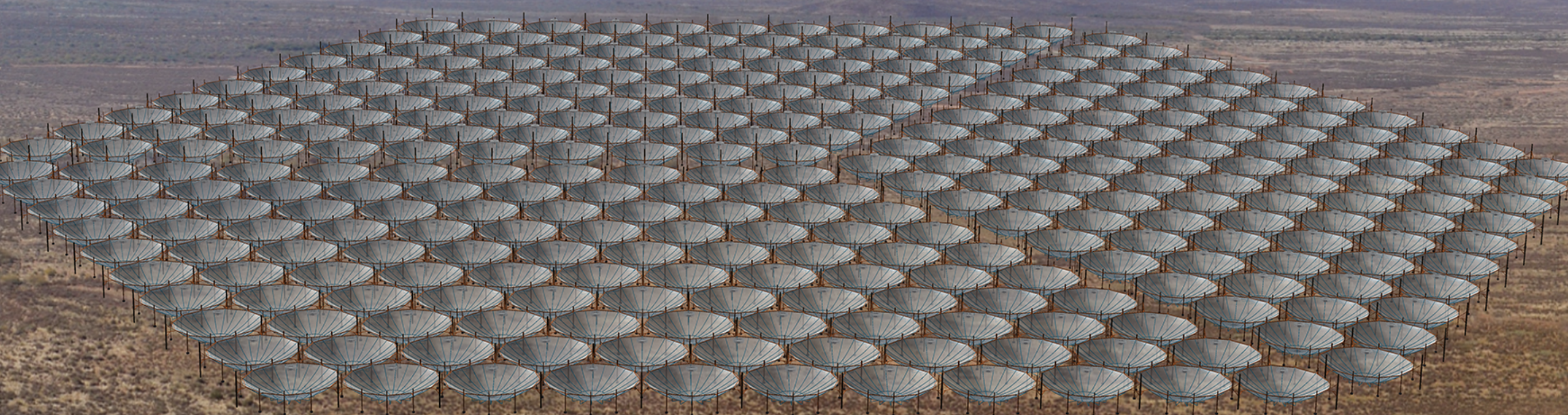
BROWN

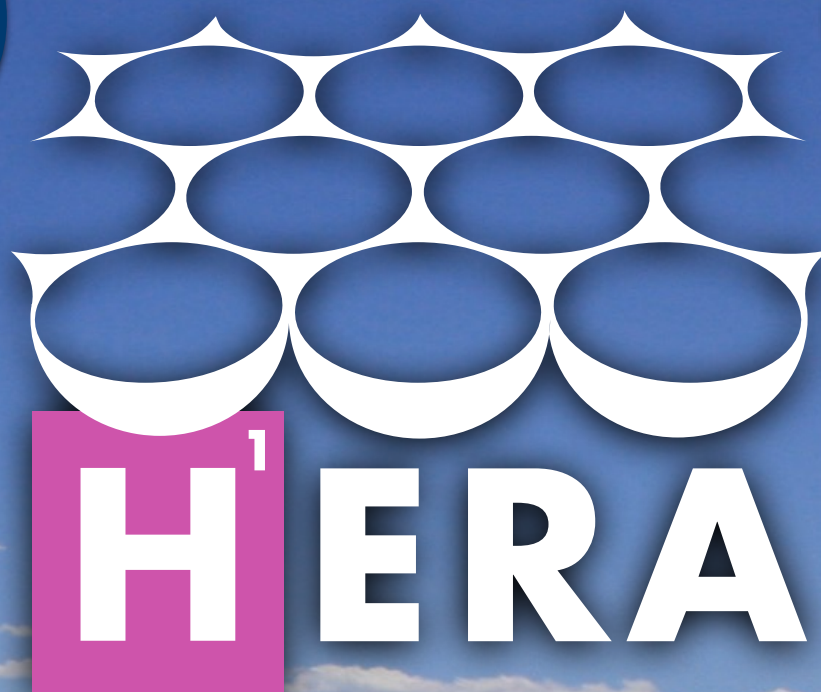
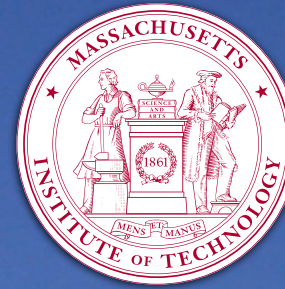


RHODES UNIVERSITY
Where leaders learn



UNIVERSITY of the
WESTERN CAPE





The Hydrogen Epoch of Reionization Array



GORDON AND BETTY
MOORE
FOUNDATION



The 21 cm signal is faint,
so HERA is huge.

← 350 14-m diameter dishes →

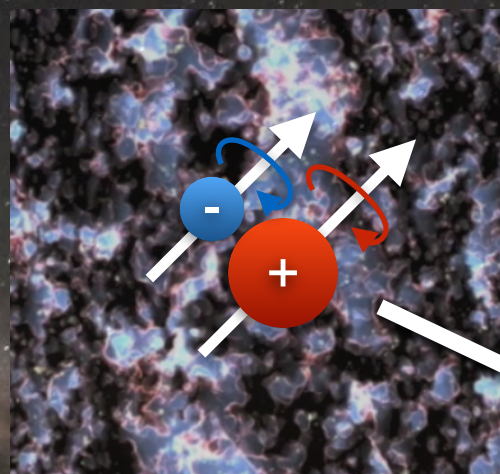


The image is a radio sky map, likely a Mosaic-1 map, showing a large portion of the sky in a Mollweide projection. The background is a dark red color, representing the radio continuum emission. A prominent, bright, yellowish-white arc, known as the HERA Stripe, is visible, curving from the upper right towards the lower right. This stripe represents a region of constant declination. The text "The HERA Stripe" is overlaid in white on the left side of the image. Below the text, there are two horizontal white lines that span the width of the image, highlighting the stripe's position.

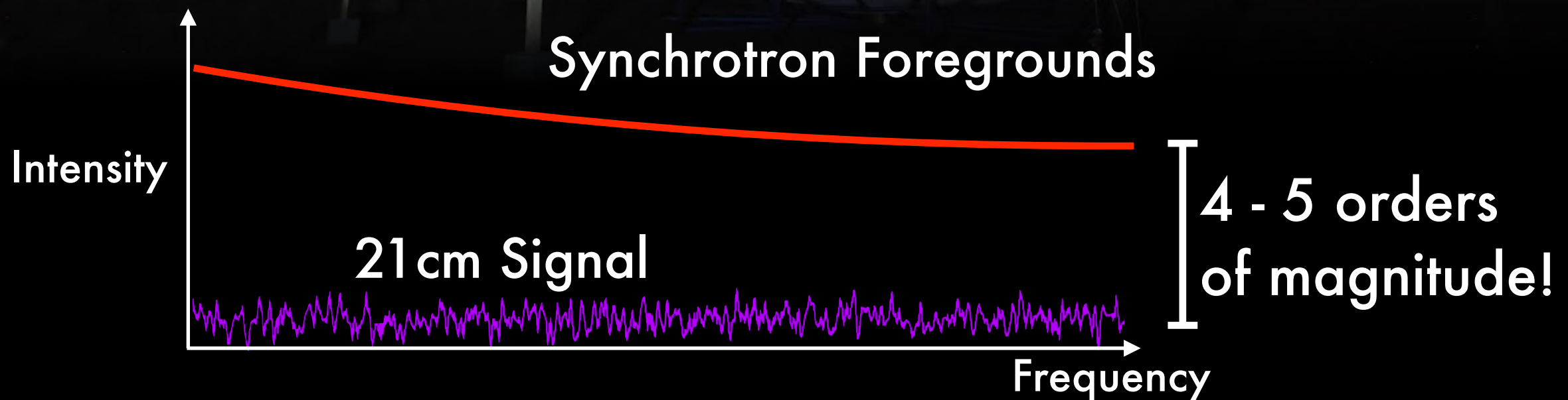
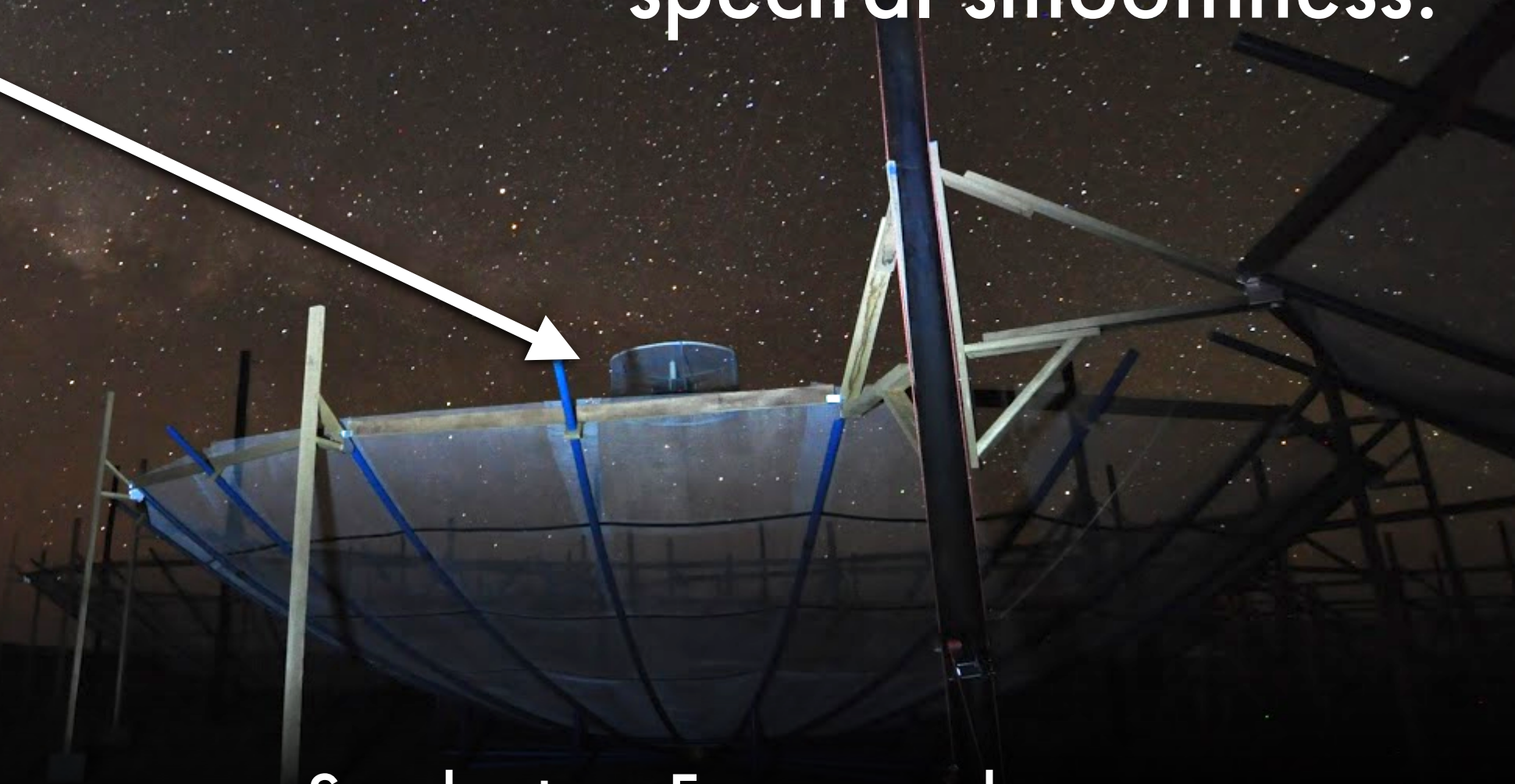
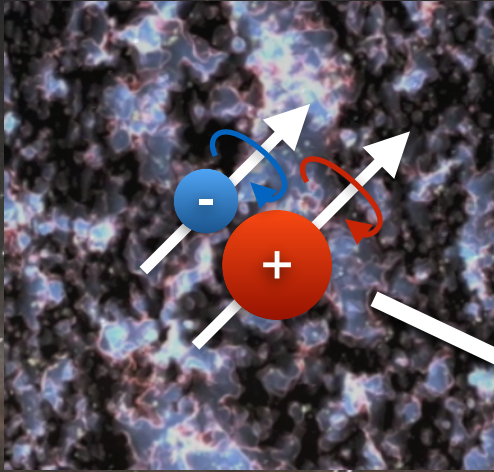
The HERA Stripe

HERA is a drift scan instrument that maps out a stripe of constant declination.

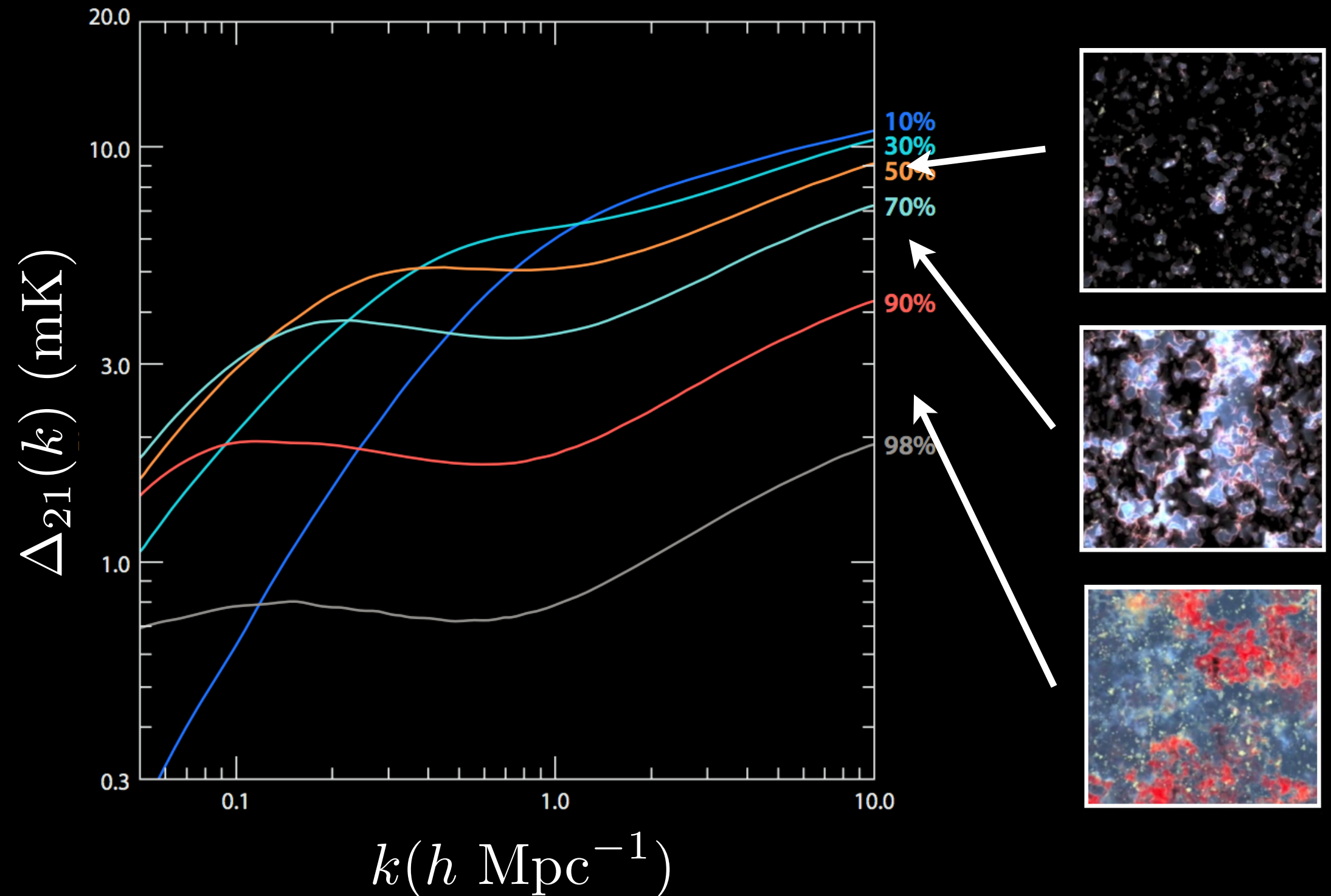
Our biggest problem
is foregrounds.



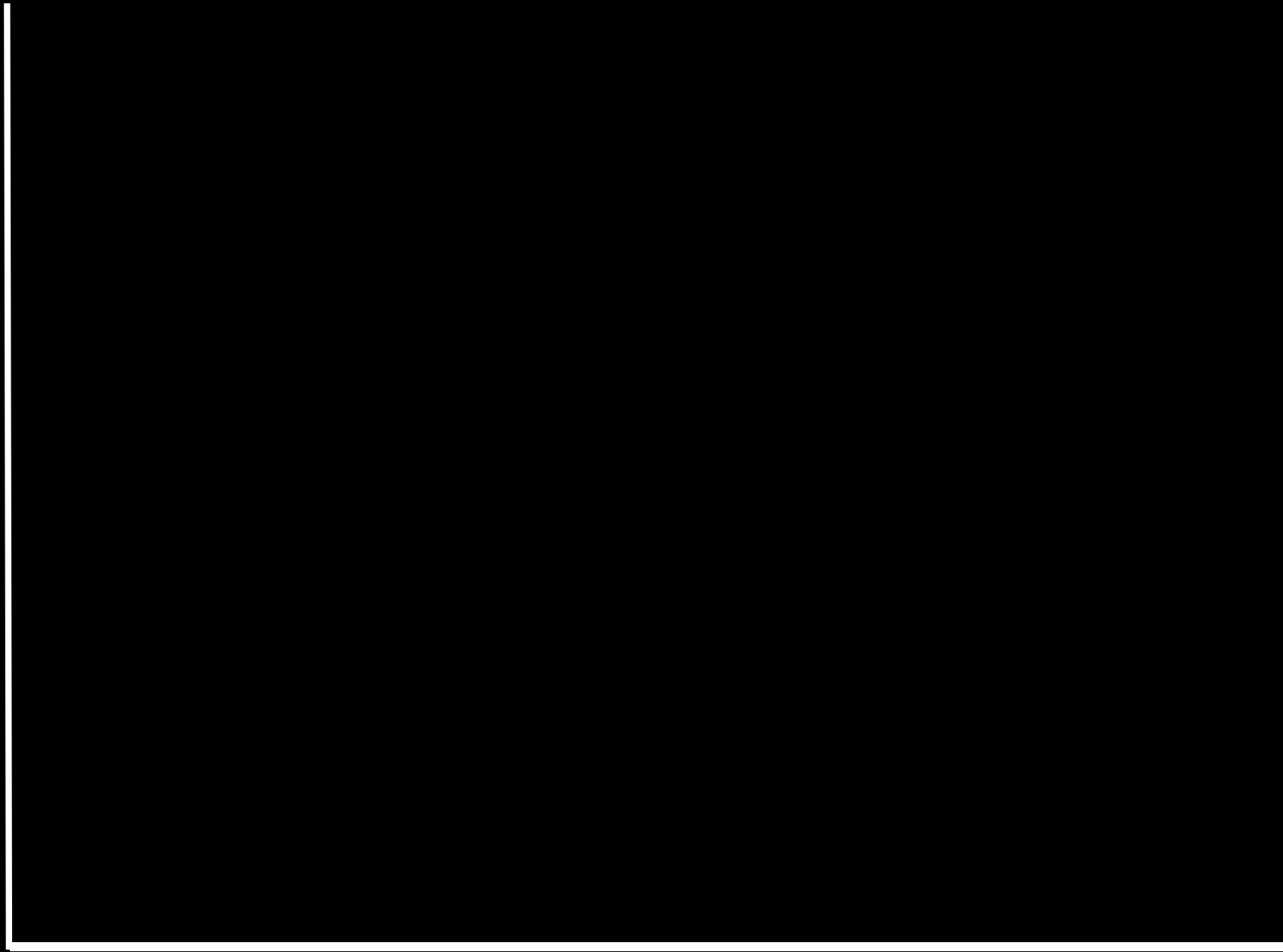
The key to separating out foregrounds is their spectral smoothness.



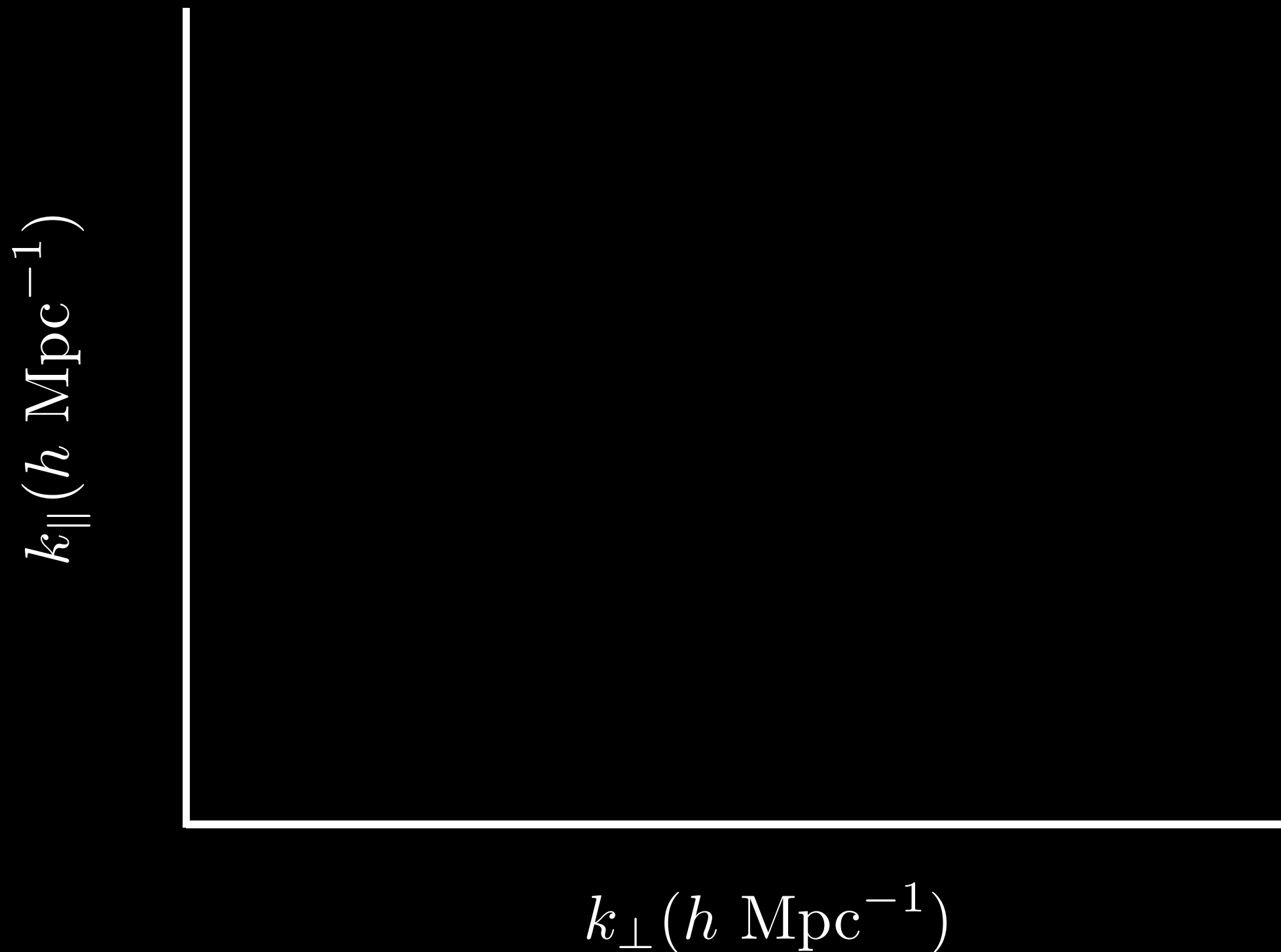
So instead of spherically averaged Fourier space...



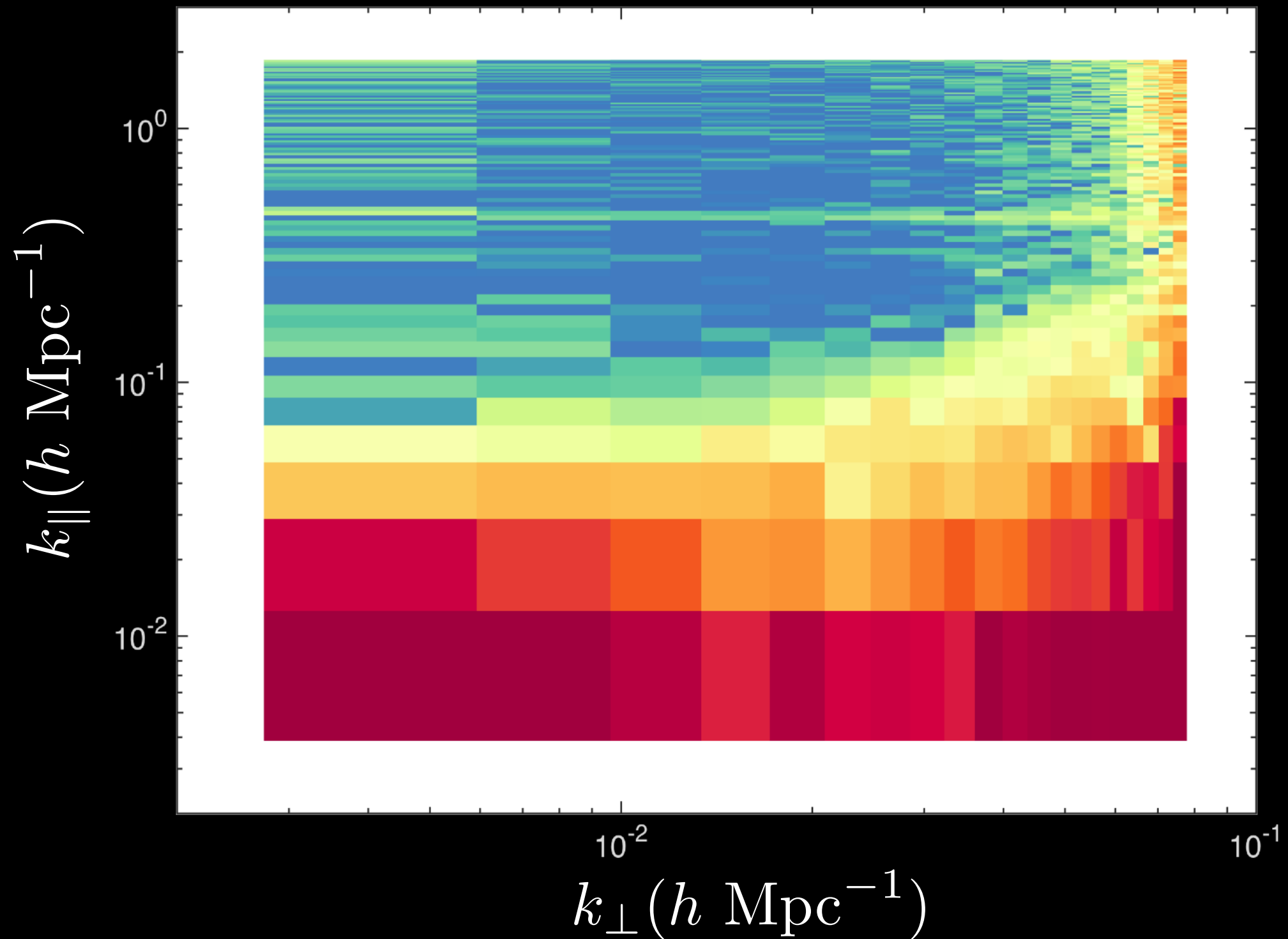
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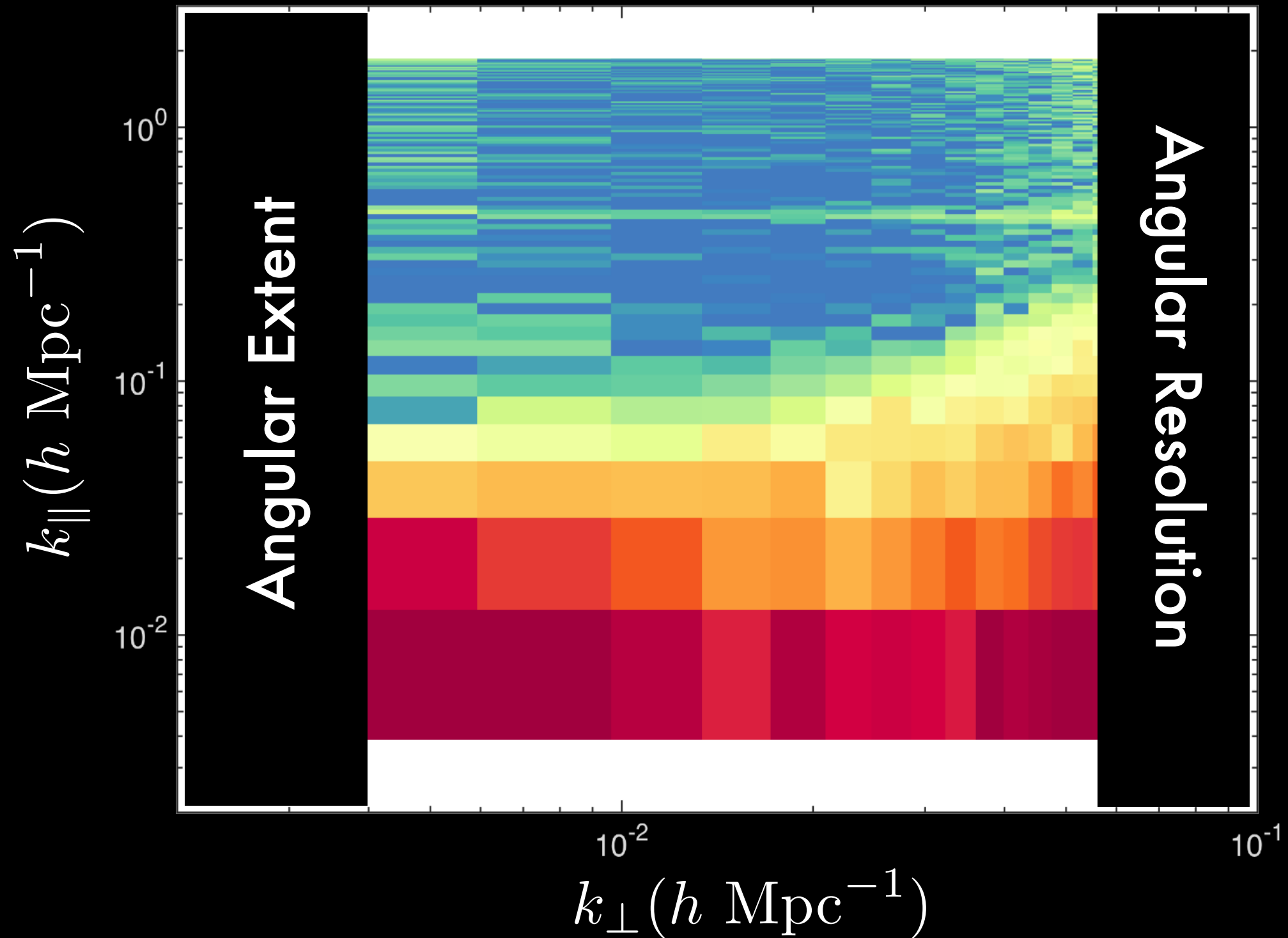
We separate out Fourier modes parallel and perpendicular to the line of sight.



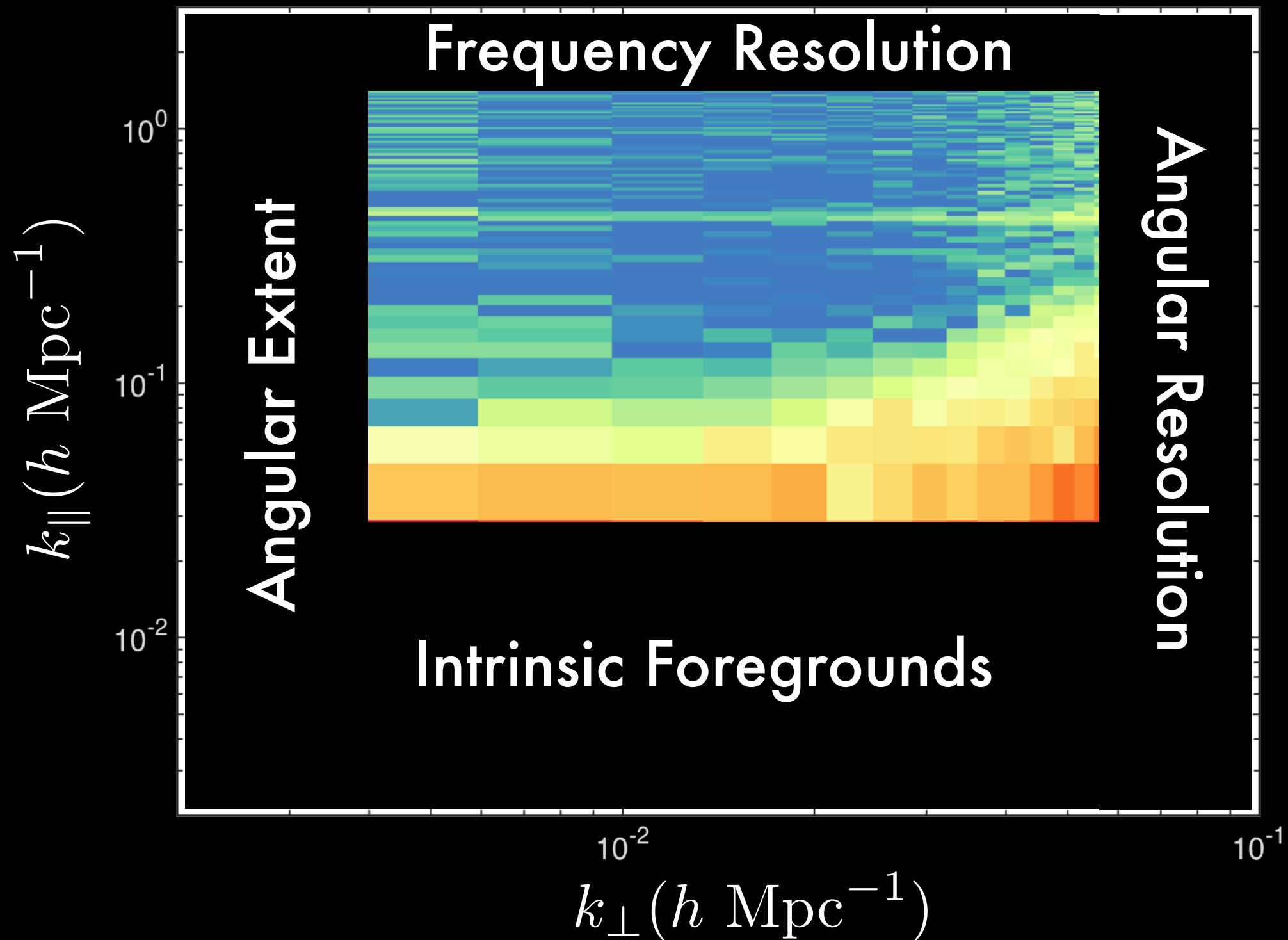
And we find a “window.”



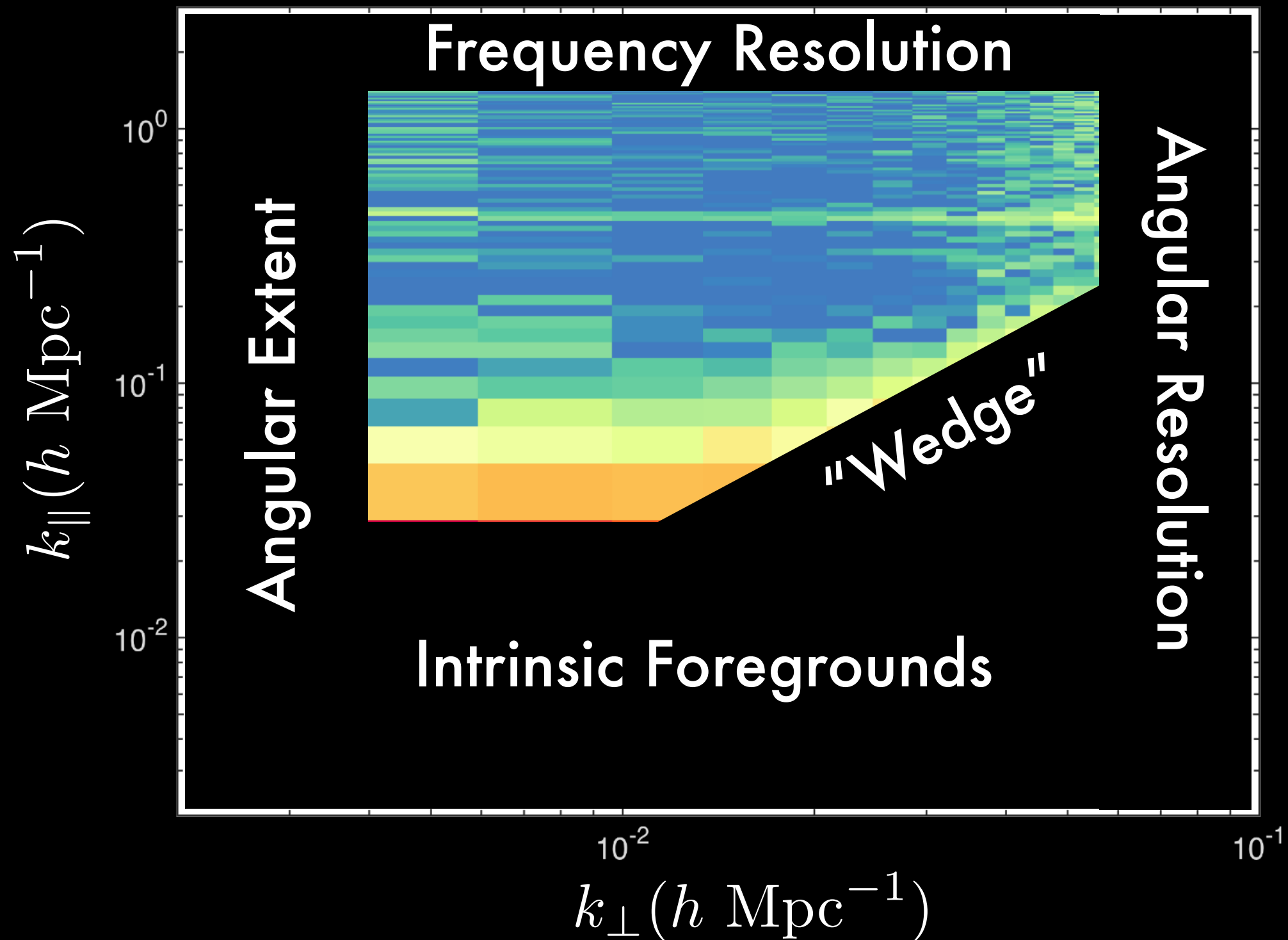
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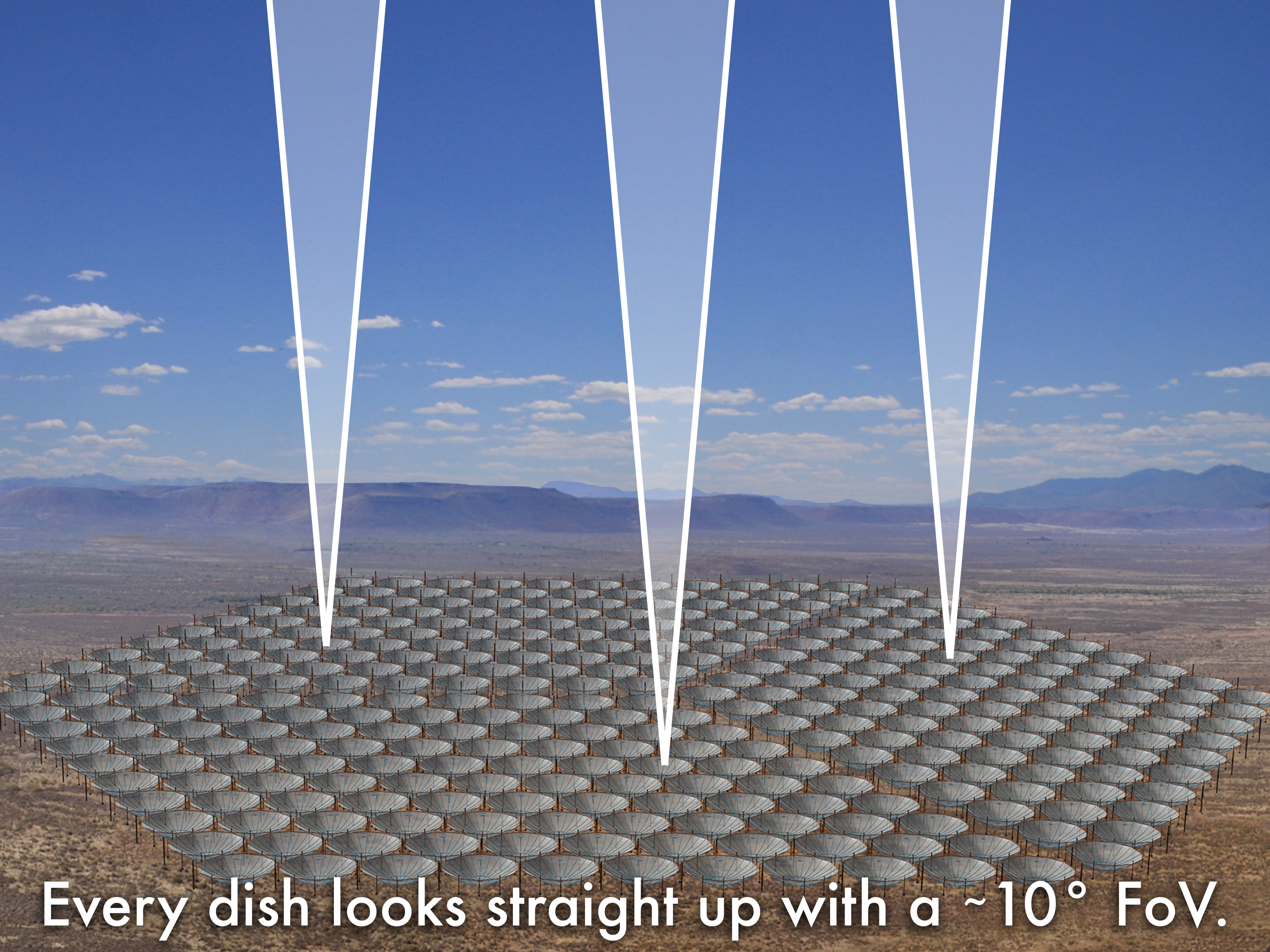


And we find a "window."



What does HERA actually measure?

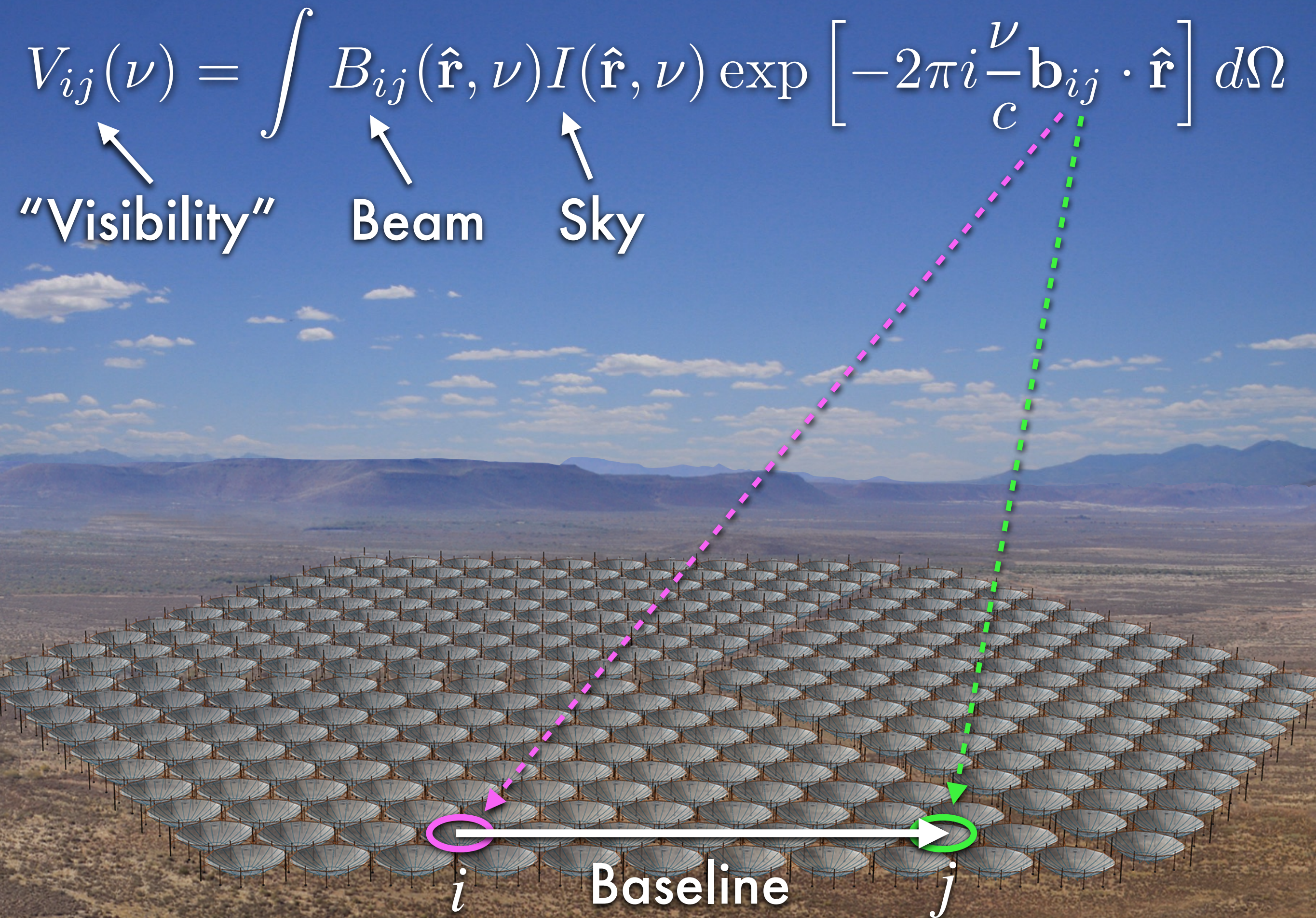




Every dish looks straight up with a $\sim 10^\circ$ FoV.

Interferometers measure Fourier modes
on the sky, which we call “visibilities.”



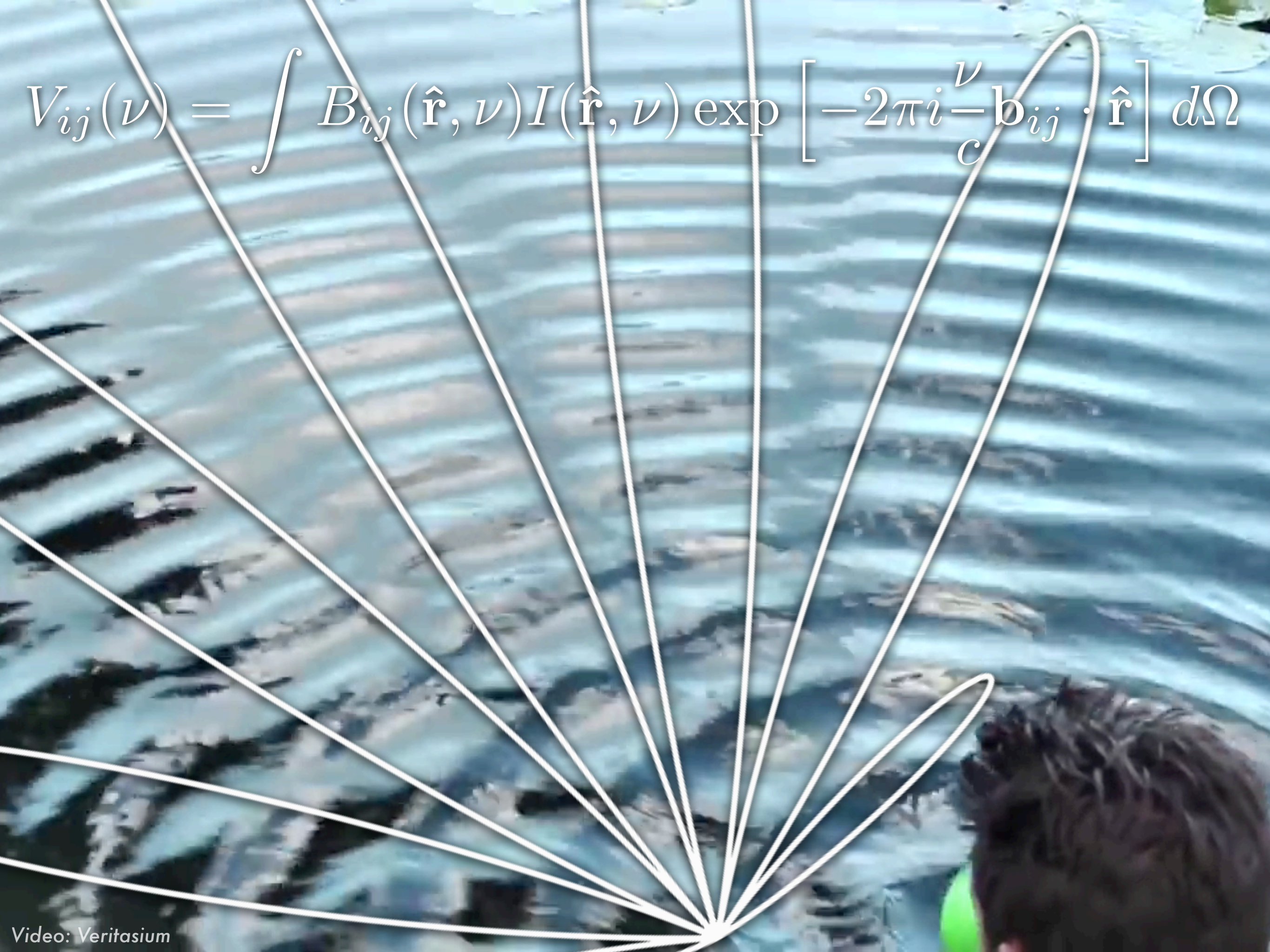




$$V_{ij}(\nu) = \int B_{ij}(\hat{\mathbf{r}}, \nu) I(\hat{\mathbf{r}}, \nu) \exp \left[-2\pi i \frac{\nu}{c} \mathbf{b}_{ij} \cdot \hat{\mathbf{r}} \right] d\Omega$$



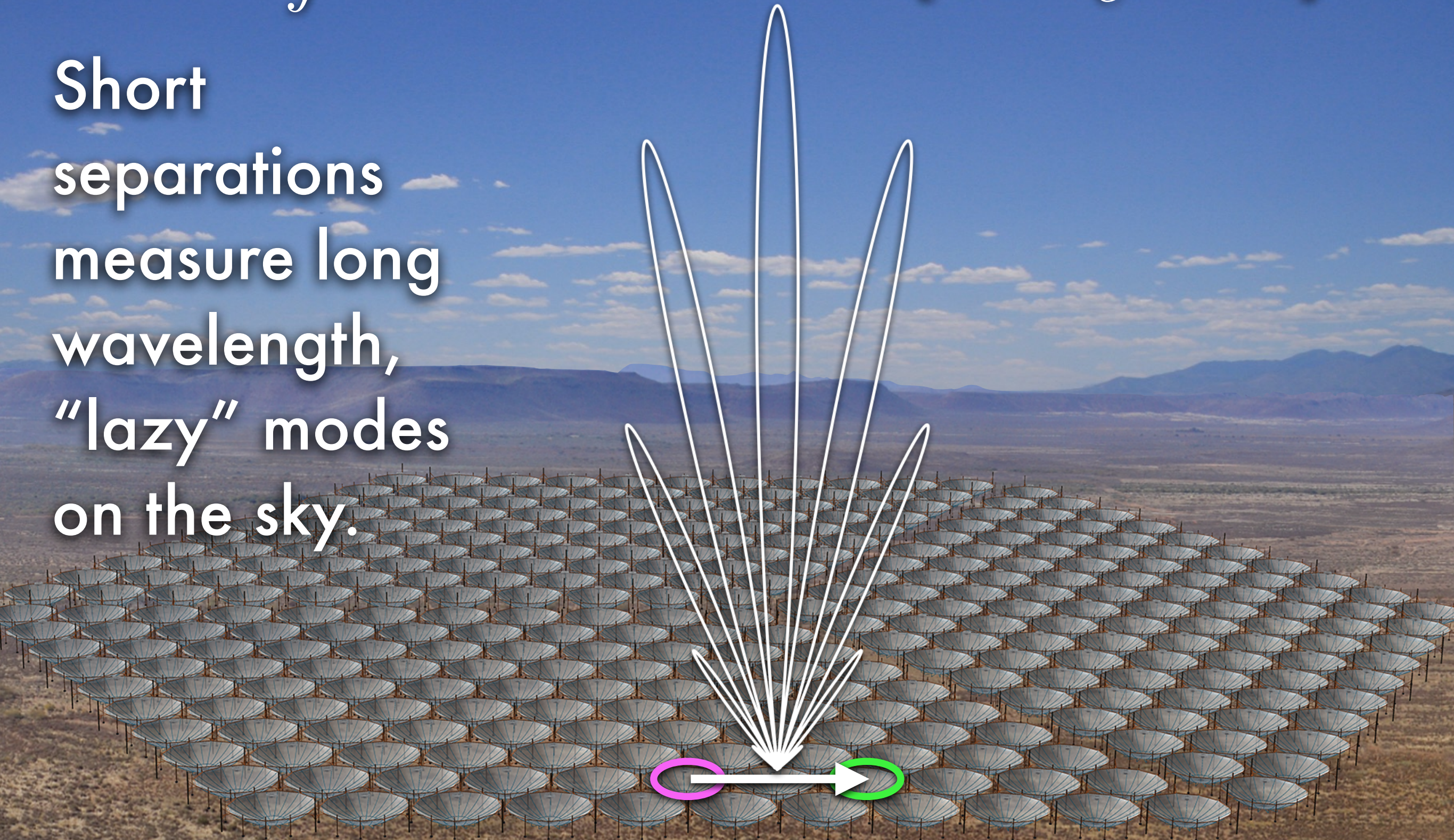
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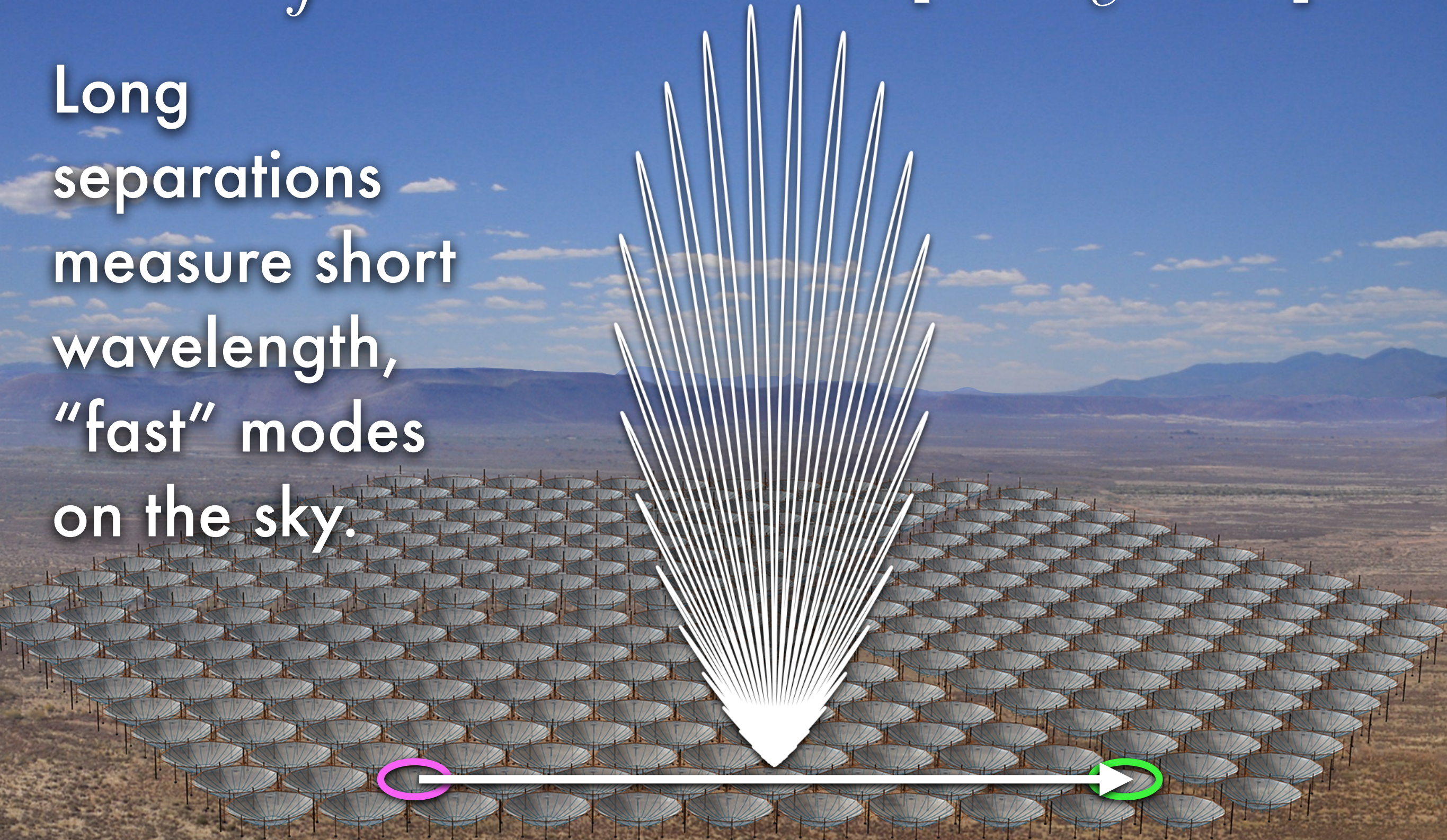
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Short
separations
measure long
wavelength,
“lazy” modes
on the sky.

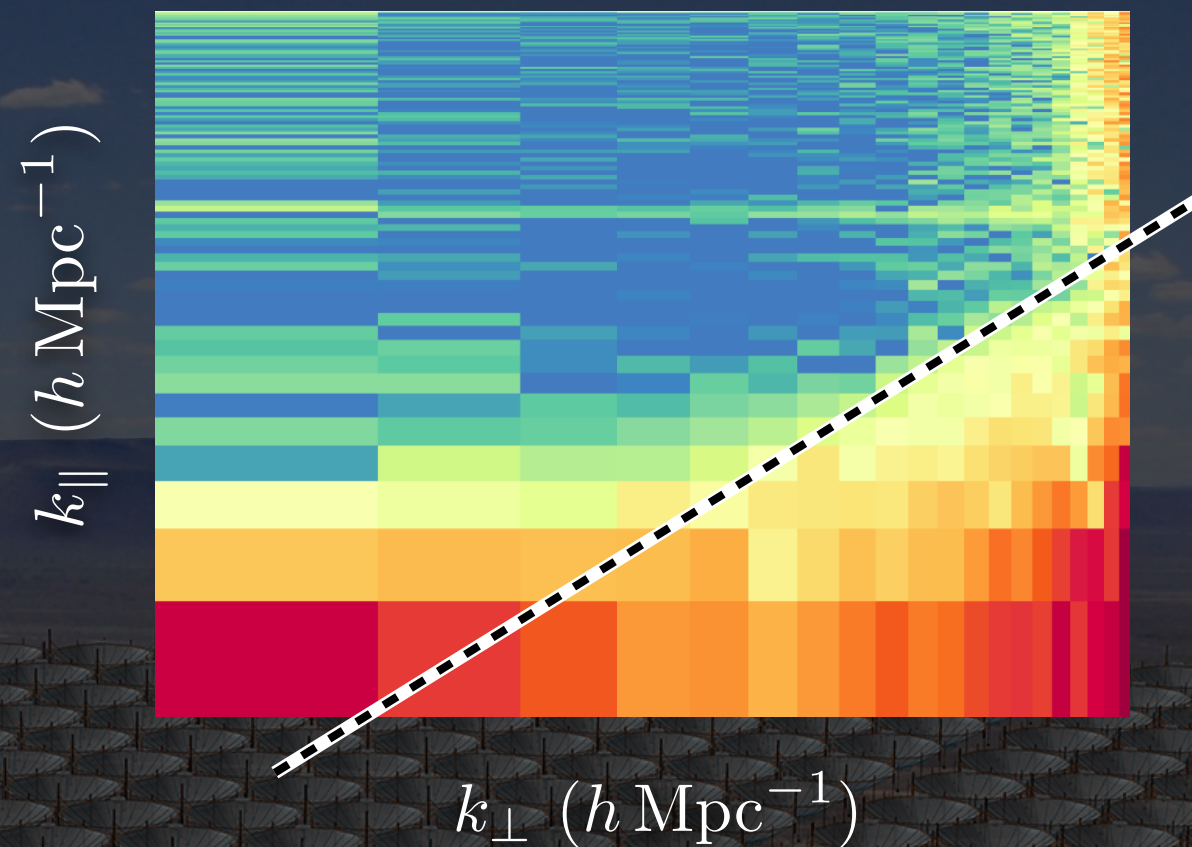


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Long
separations
measure short
wavelength,
“fast” modes
on the sky.

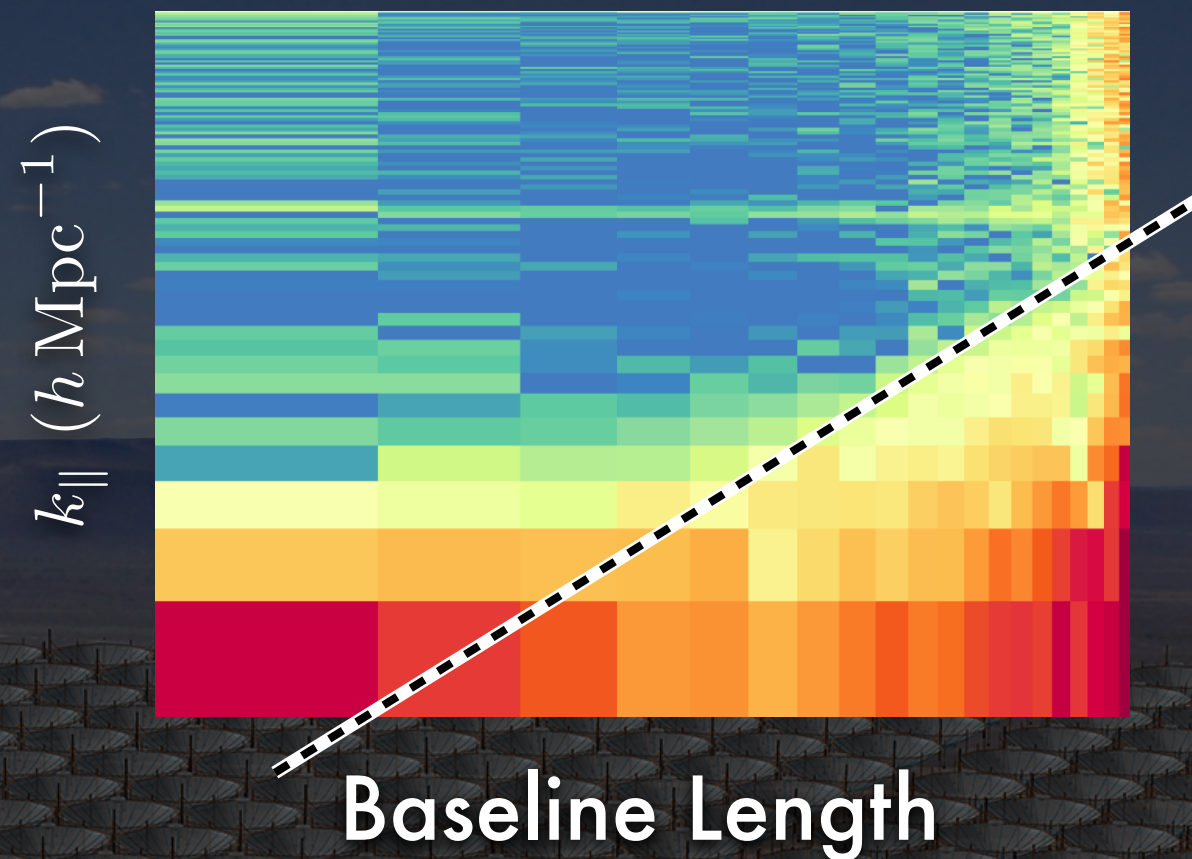


$$V_{ij}(\nu) = \int B_{ij}(\hat{\mathbf{r}}, \nu) I(\hat{\mathbf{r}}, \nu) \exp \left[-2\pi i \frac{\nu}{c} \mathbf{b}_{ij} \cdot \hat{\mathbf{r}} \right] d\Omega$$



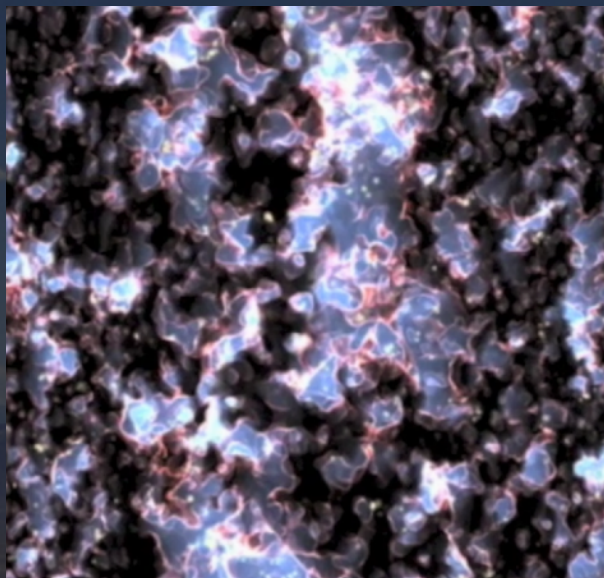
\mathbf{k}_{\perp} is effectively baseline length.

$$V_{ij}(\nu) = \int B_{ij}(\hat{\mathbf{r}}, \nu) I(\hat{\mathbf{r}}, \nu) \exp \left[-2\pi i \frac{\nu}{c} \mathbf{b}_{ij} \cdot \hat{\mathbf{r}} \right] d\Omega$$

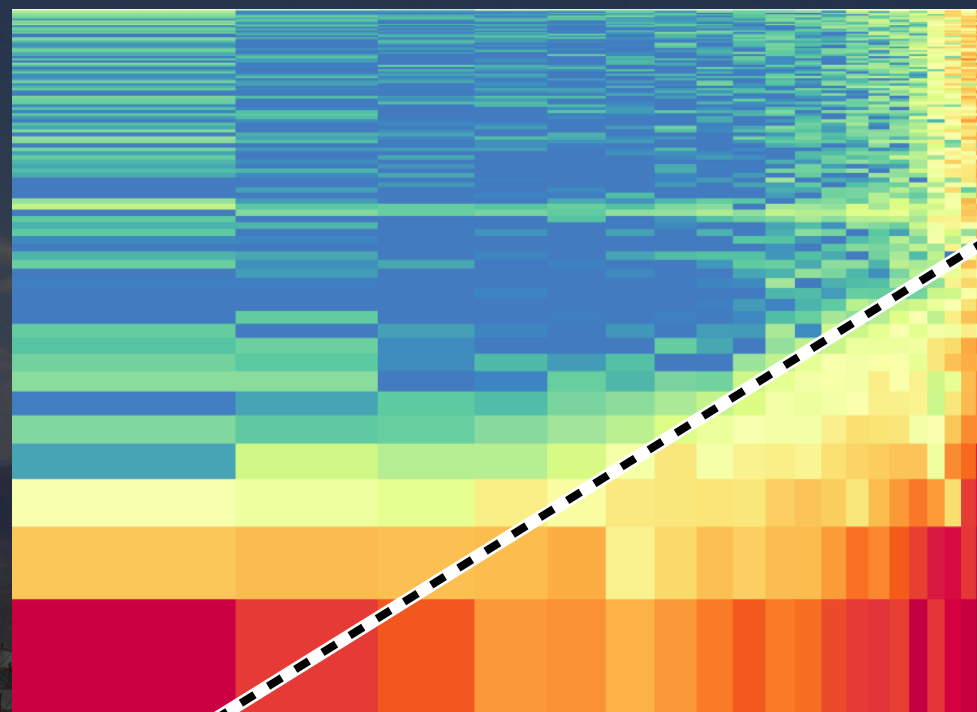


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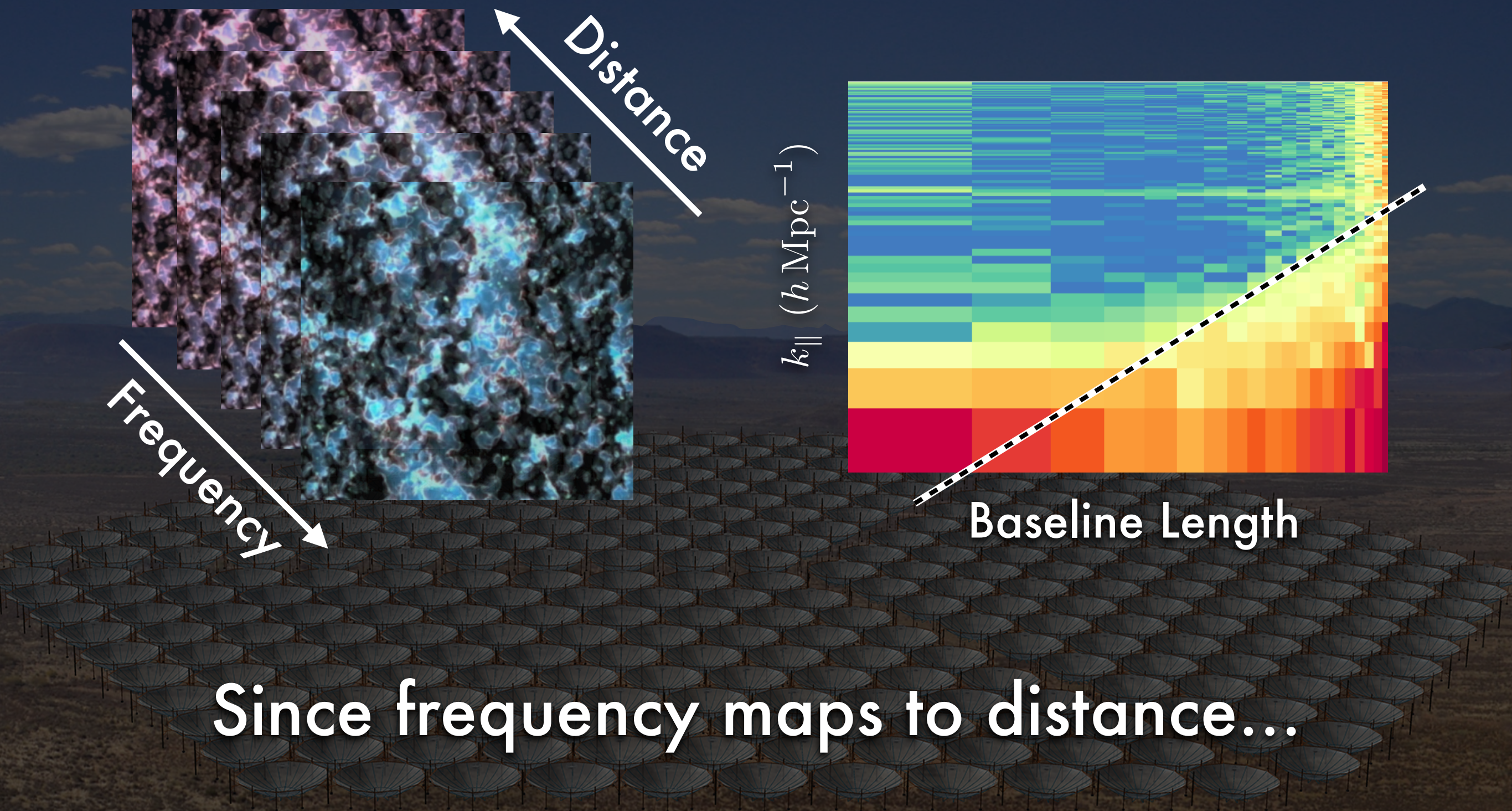
$k_{\parallel} \text{ (} h \text{ Mpc}^{-1} \text{)}$



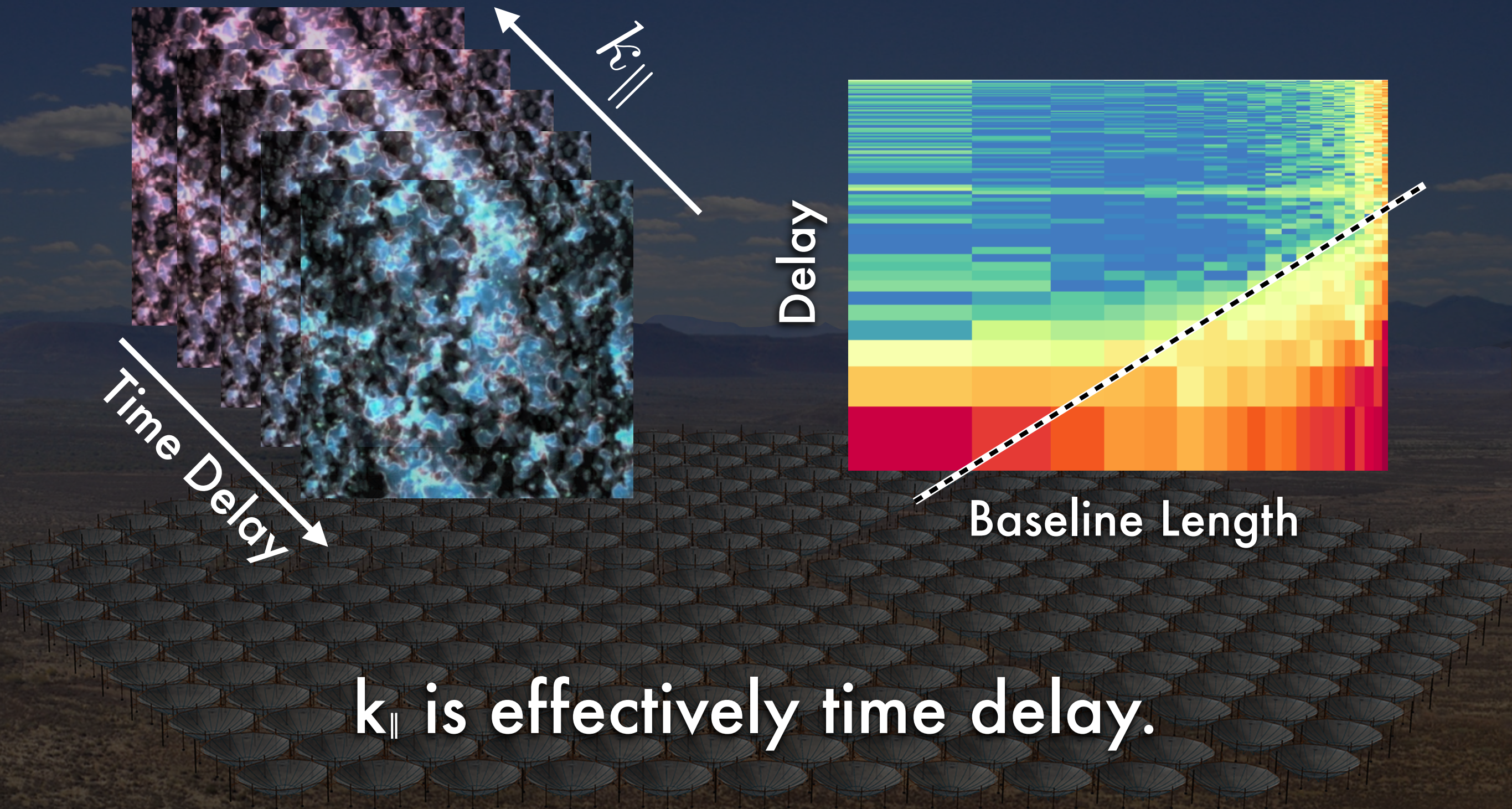
Baseline Length

Since frequency maps to distance...

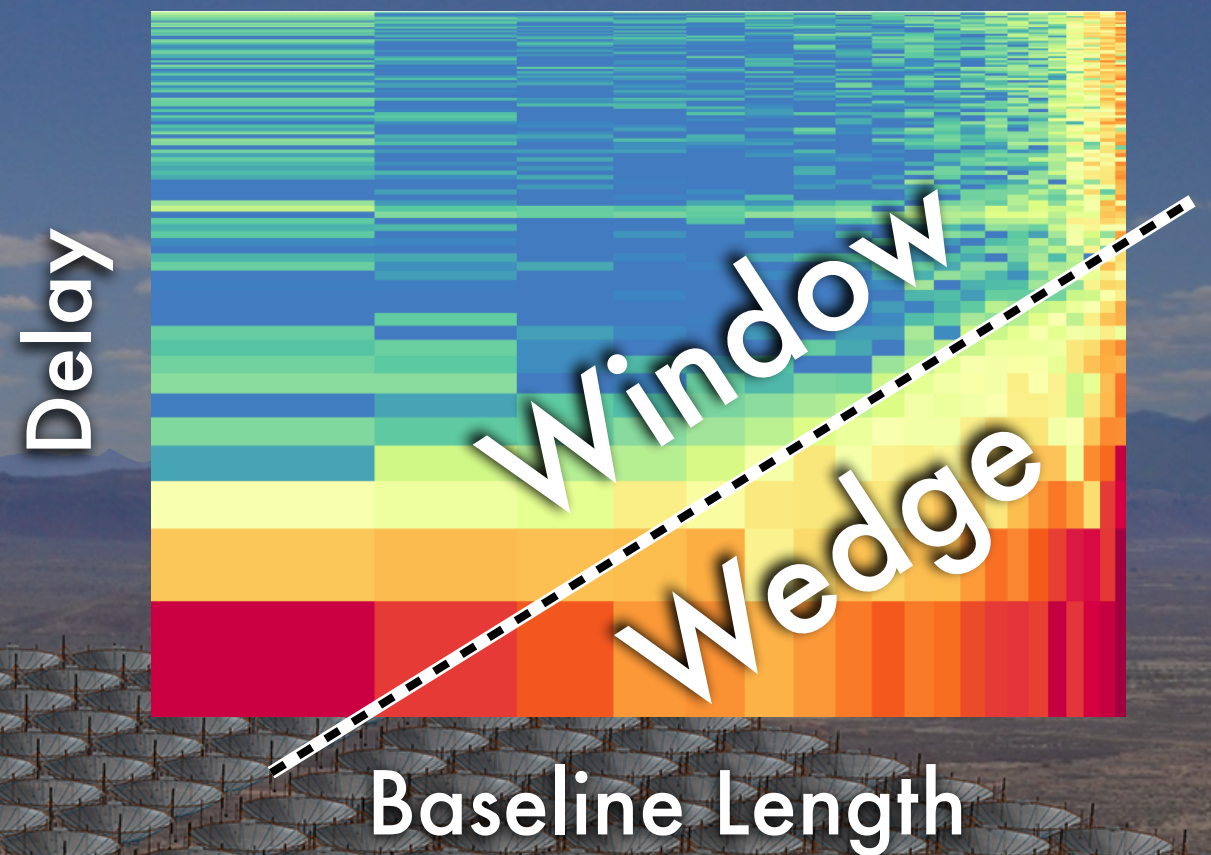
$$V_{ij}(\nu) = \int B_{ij}(\hat{\mathbf{r}}, \nu) I(\hat{\mathbf{r}}, \nu) \exp \left[-2\pi i \frac{\nu}{c} \mathbf{b}_{ij} \cdot \hat{\mathbf{r}} \right] d\Omega$$



$$V_{ij}(\nu) = \int B_{ij}(\hat{\mathbf{r}}, \nu) I(\hat{\mathbf{r}}, \nu) \exp \left[-2\pi i \frac{\nu}{c} \mathbf{b}_{ij} \cdot \hat{\mathbf{r}} \right] d\Omega$$



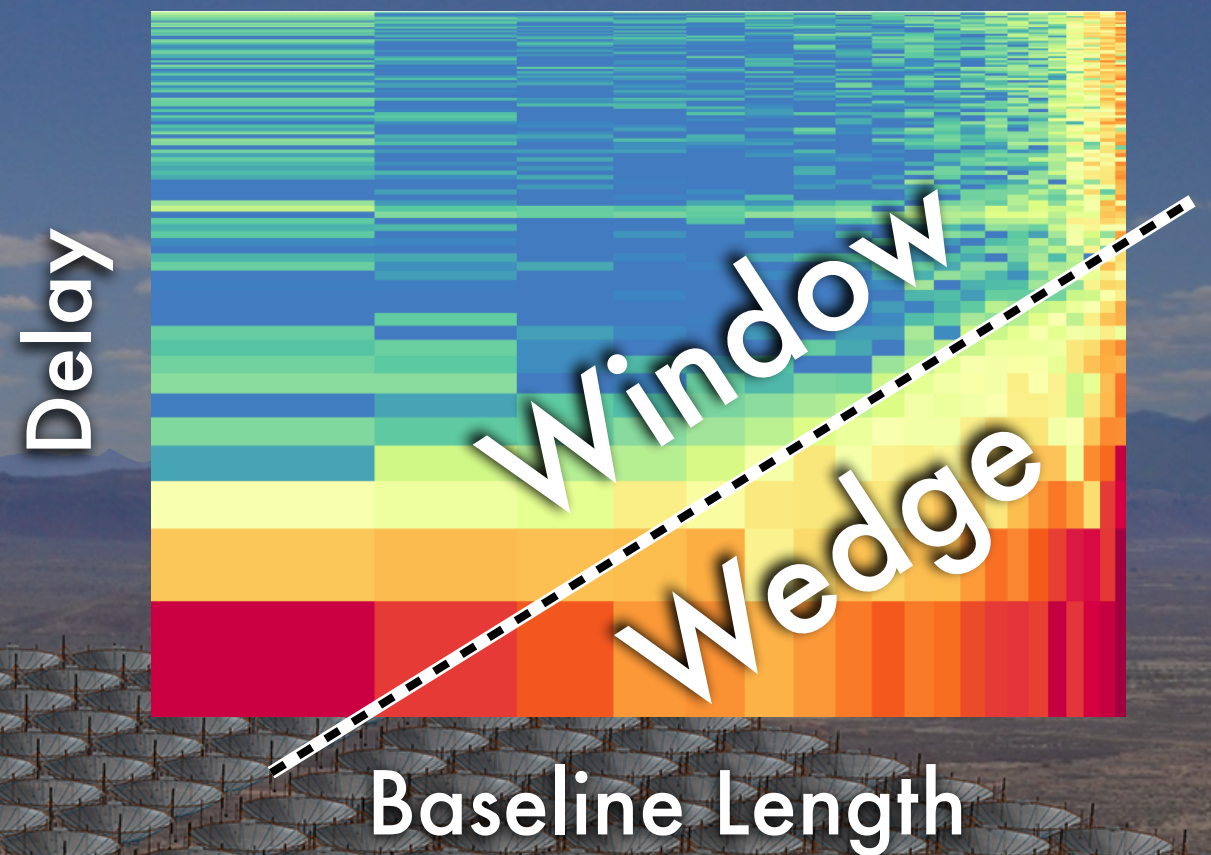
The maximum delay of foregrounds for a baseline is simply the light travel time.



$$\Delta t_{\max} = |\mathbf{b}|/c$$

$$\Delta t_{\max} = |\mathbf{b}|/c$$

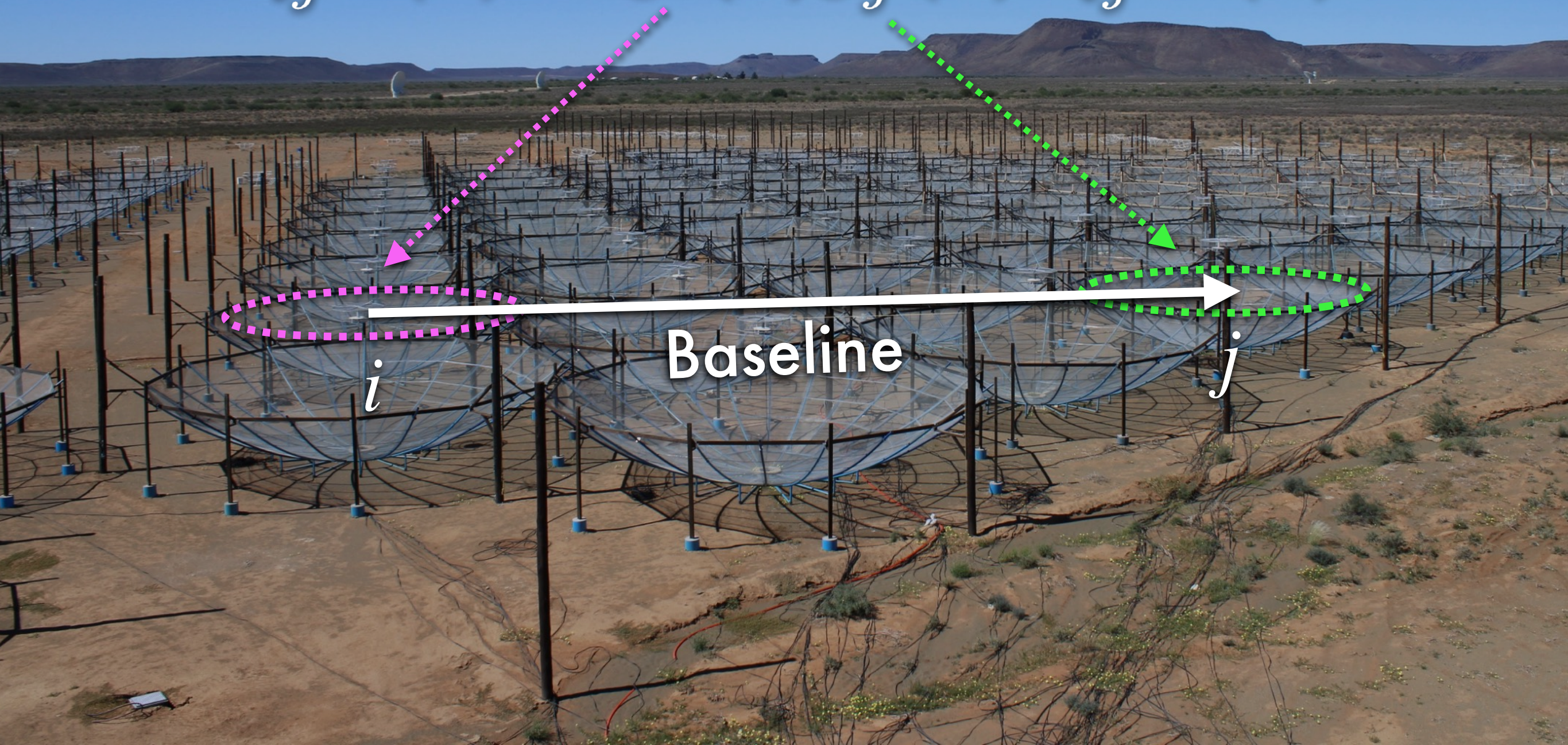
Our design for
HERA's configuration
maximizes sensitivity
on short baselines.



$$\Delta t_{\max} = |\mathbf{b}|/c$$

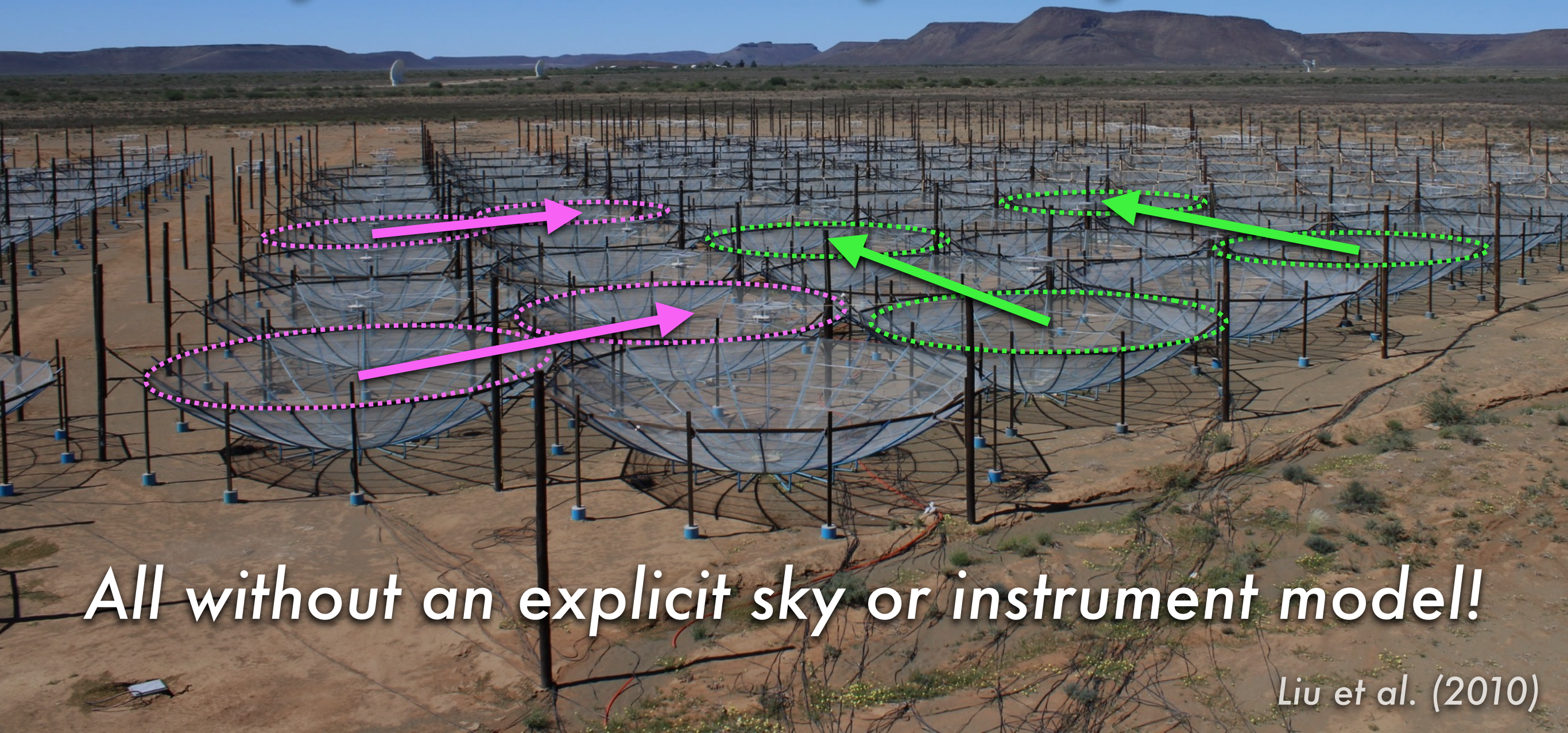
Foreground avoidance won't work without precision calibration.

$$V_{ij}^{\text{obs}}(\nu) = g_i(\nu)g_j^*(\nu)V_{ij}^{\text{true}}(\nu)$$



HERA was designed to be calibrated using the internal consistency of redundant baselines.

$$V_{ij}^{\text{obs}}(\nu) = g_i(\nu)g_j^*(\nu)V_{ij}^{\text{true}}(\nu)$$

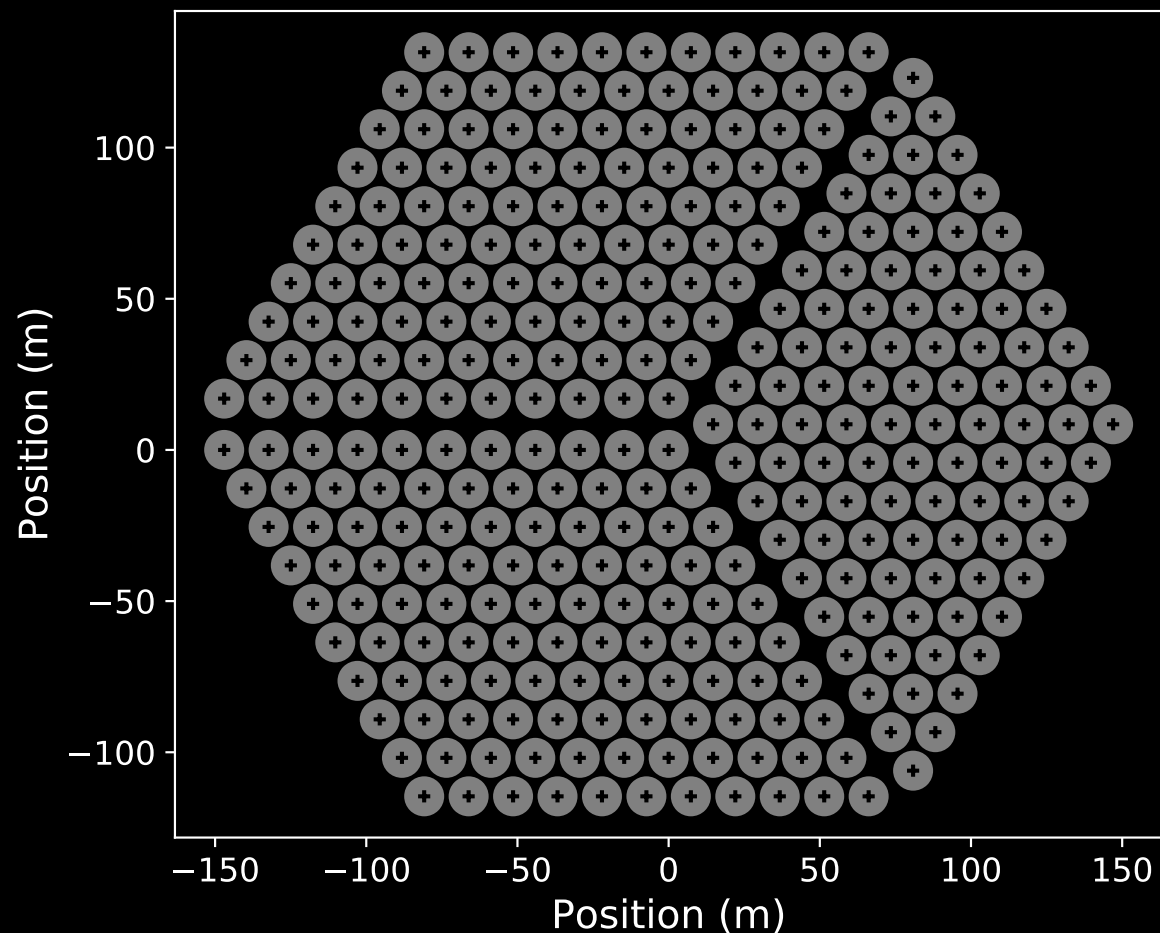


All without an explicit sky or instrument model!

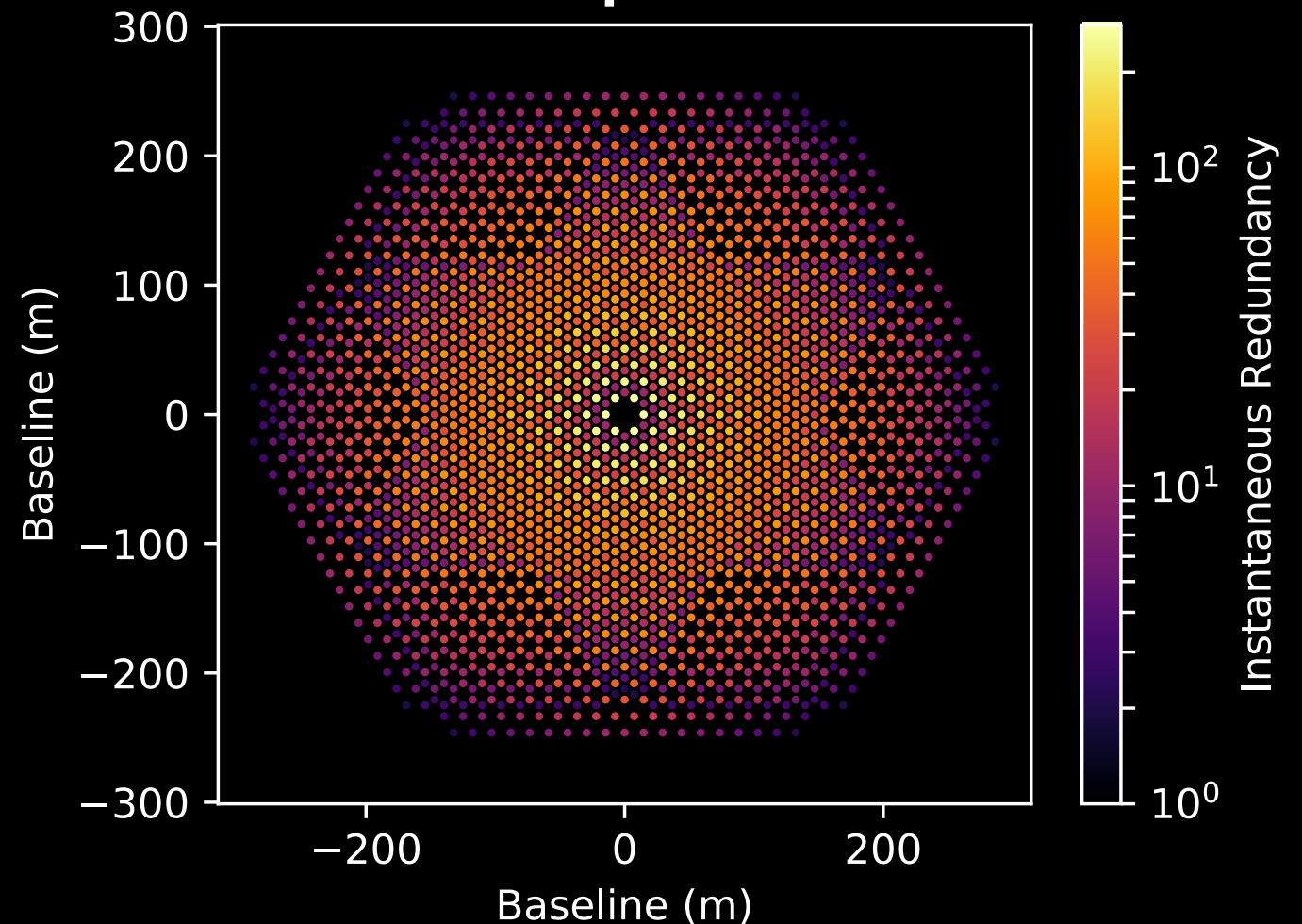
$$V_{ij}^{\text{obs}}(\nu) = g_i(\nu)g_j^*(\nu)V_{ij}^{\text{true}}(\nu)$$

$$V_{ij}^{\text{obs}}(\nu) = g_i(\nu)g_j^*(\nu)V_{ij}^{\text{true}}(\nu)$$

320 Antenna Gains



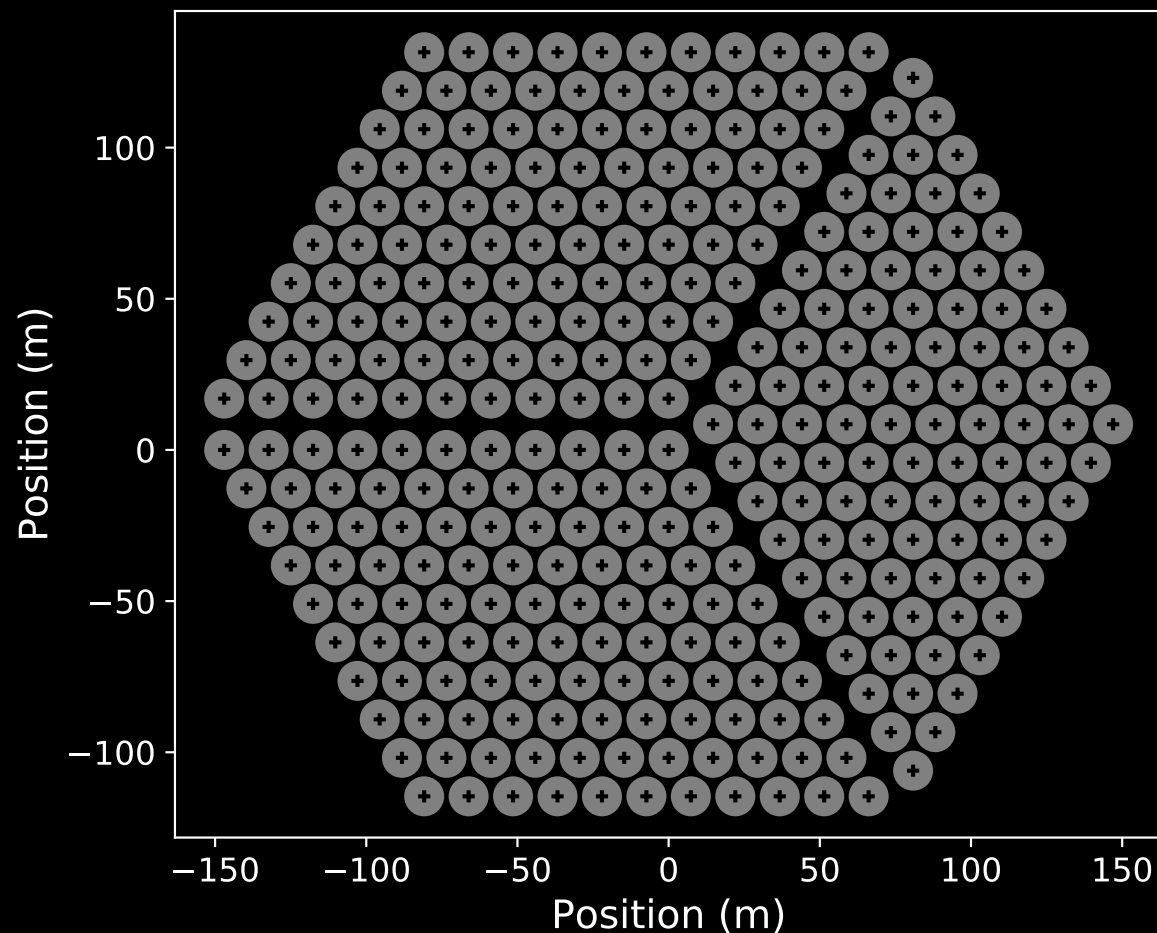
1,501 Unique Visibilities



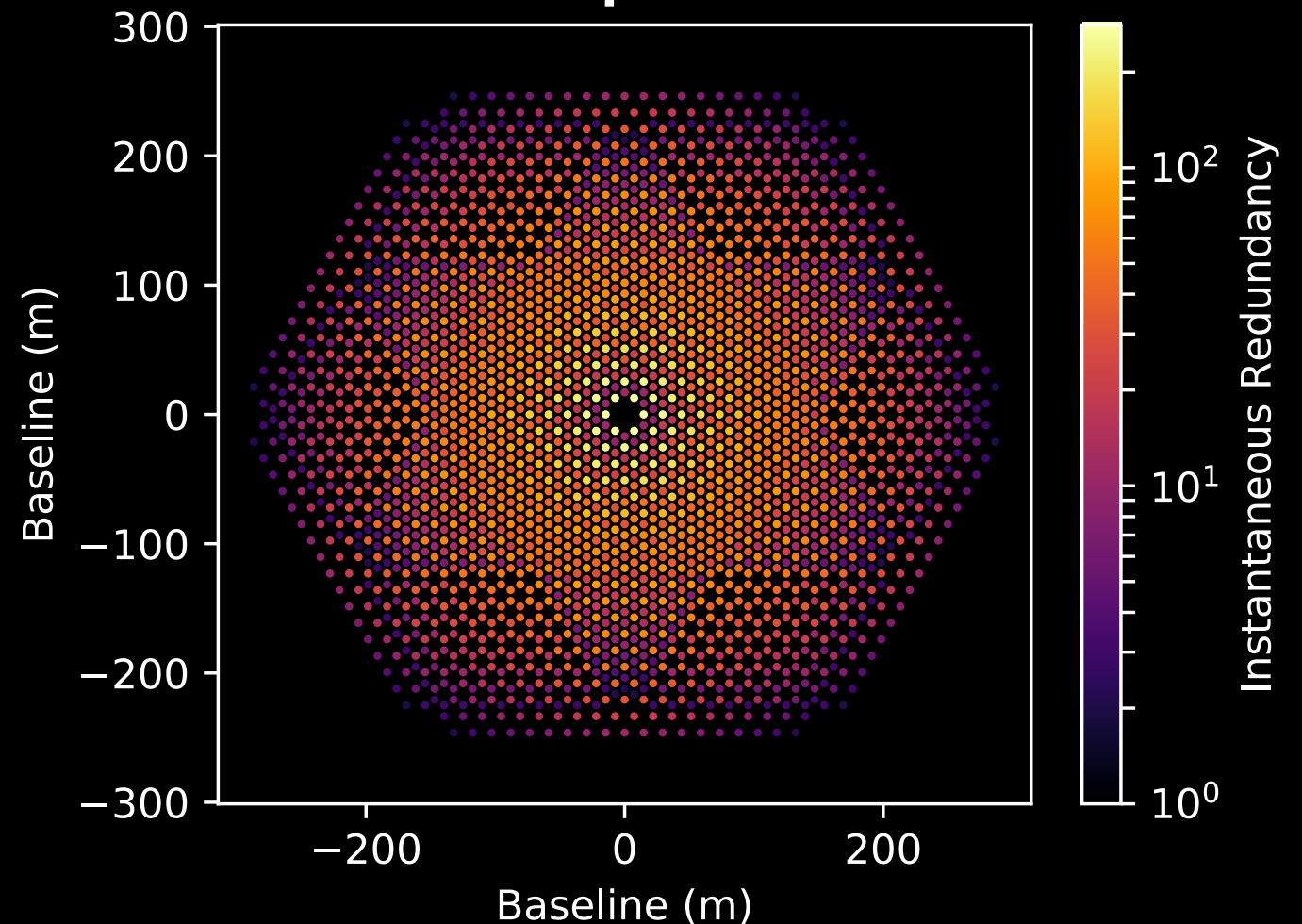
51,040 Total Measurements

$$V_{ij}^{\text{obs}}(\nu) = g_i(\nu)g_j^*(\nu)V_{ij}^{\text{true}}(\nu)$$

320 Antenna Gains



1,501 Unique Visibilities

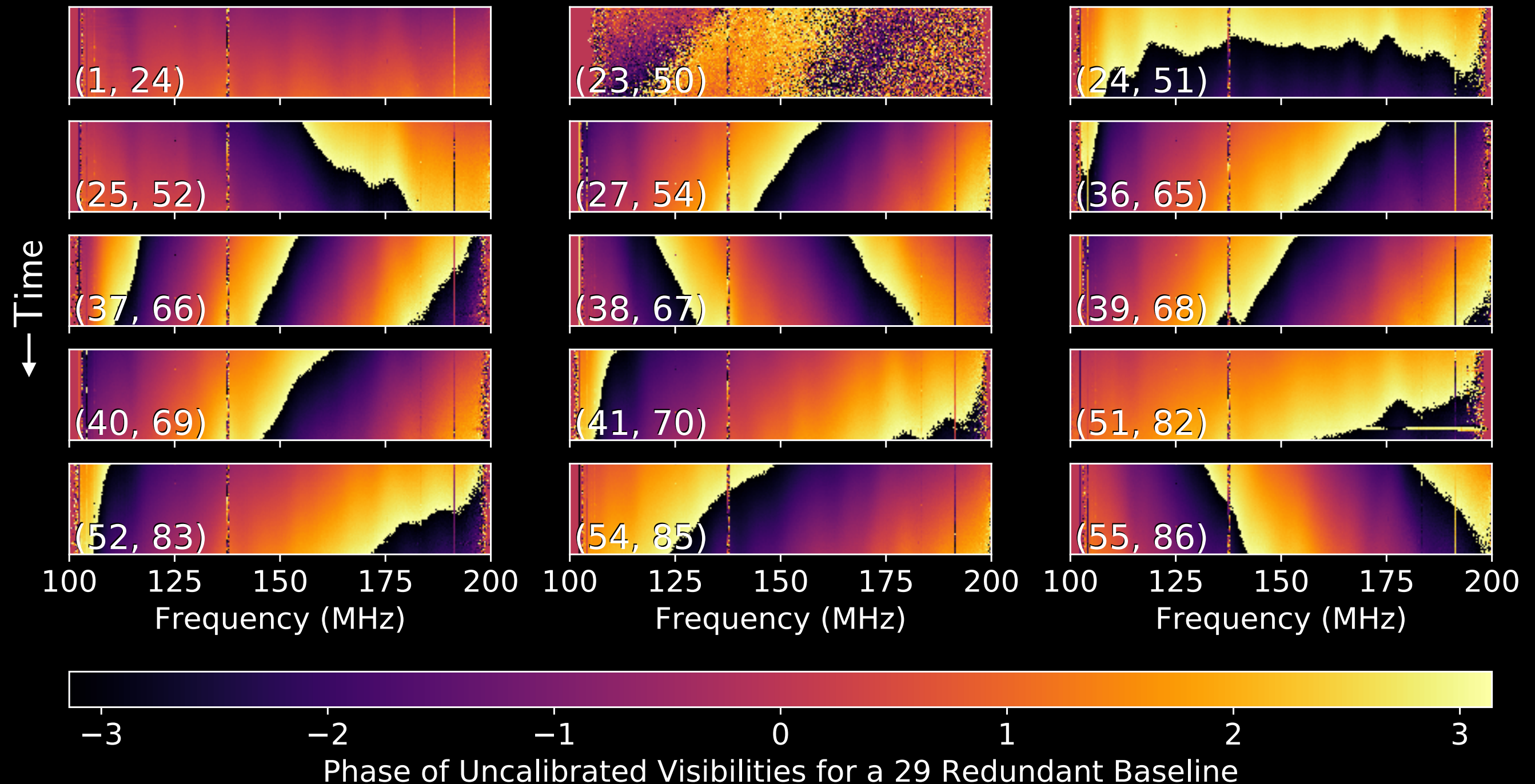


51,040 Total Measurements

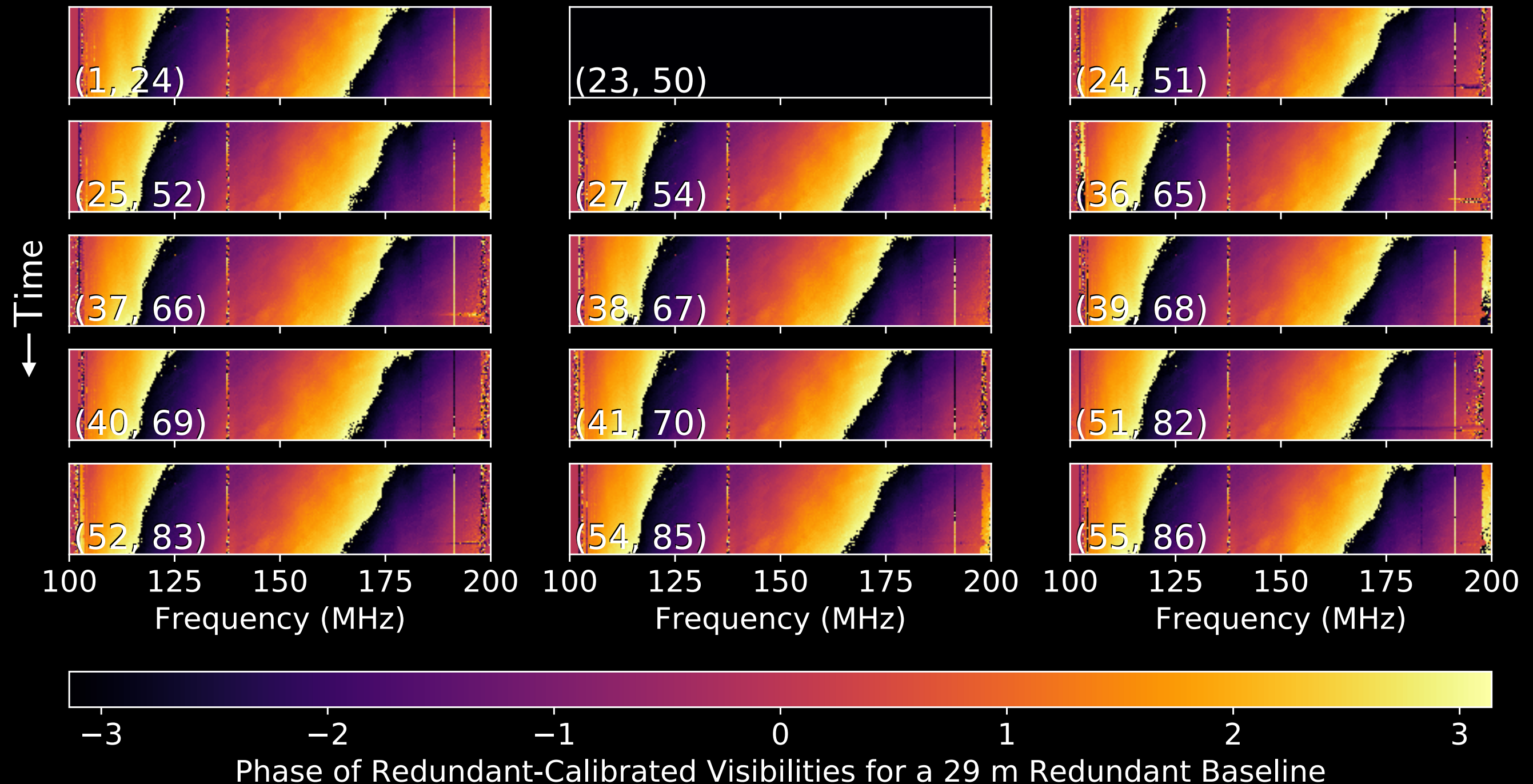
Goal: Minimize $\chi^2 \equiv \sum \frac{|V_{ij}^{\text{obs}} - g_i g_j^* V_{i-j}^{\text{sol}}|^2}{\sigma_{ij}^2}$

Redundant calibration
is working quite well.

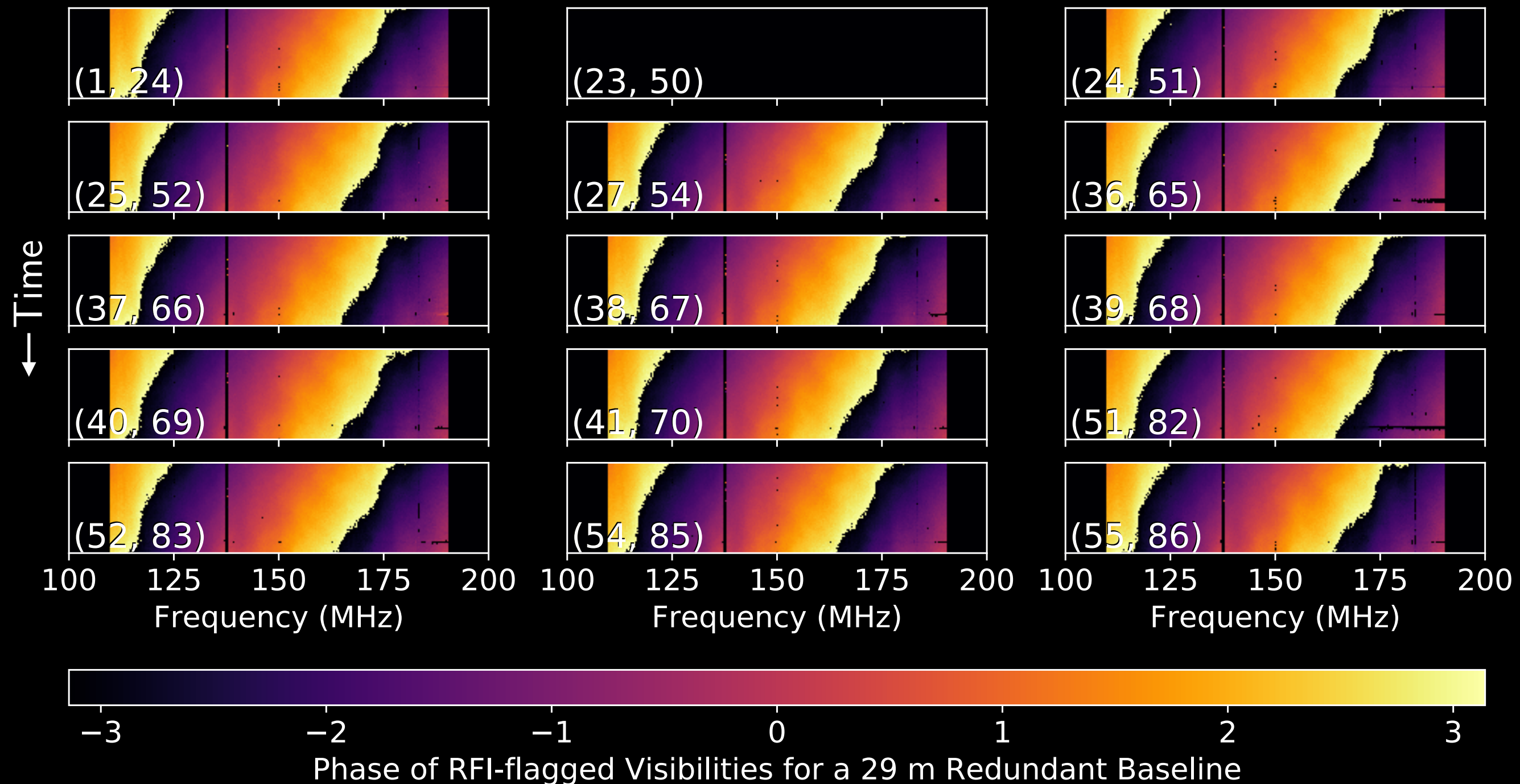
Example raw HERA data for a single redundant baseline group.



Next we flag bad antennas and impose the redundancy constraint to solve for all gains.

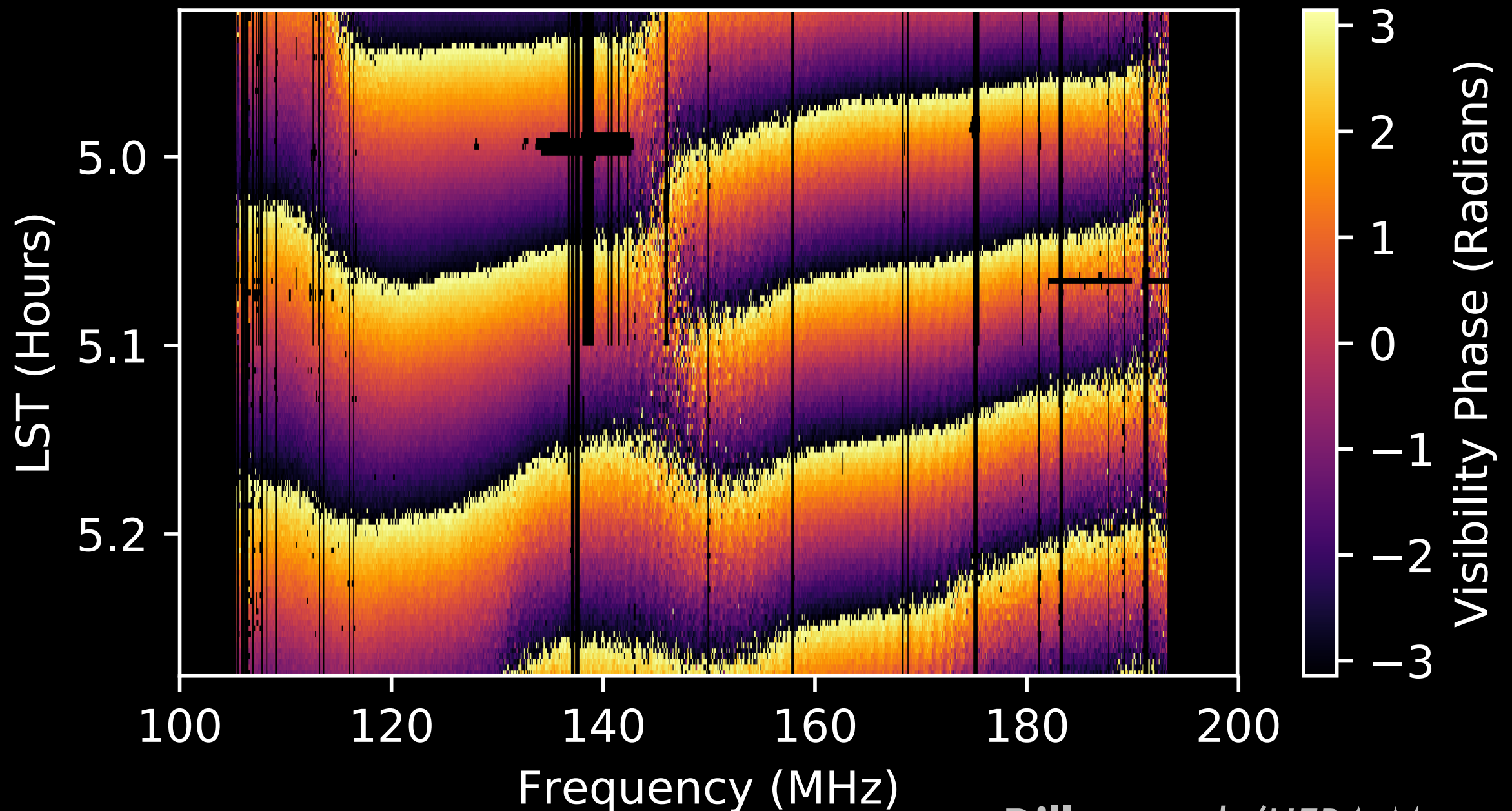


Then we mask-out radio-frequency interference (RFI).



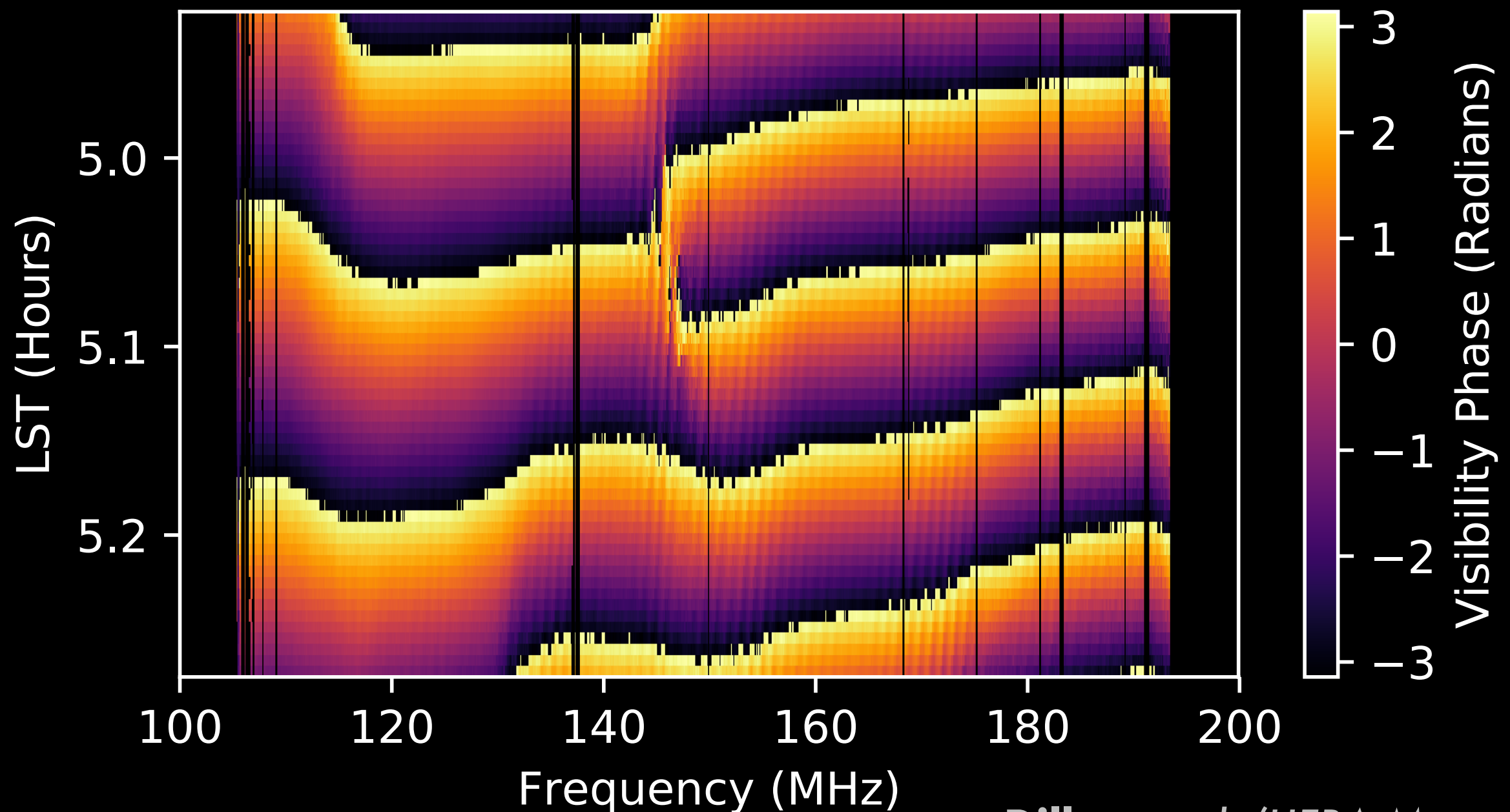
The instrument looks stable from day to day...

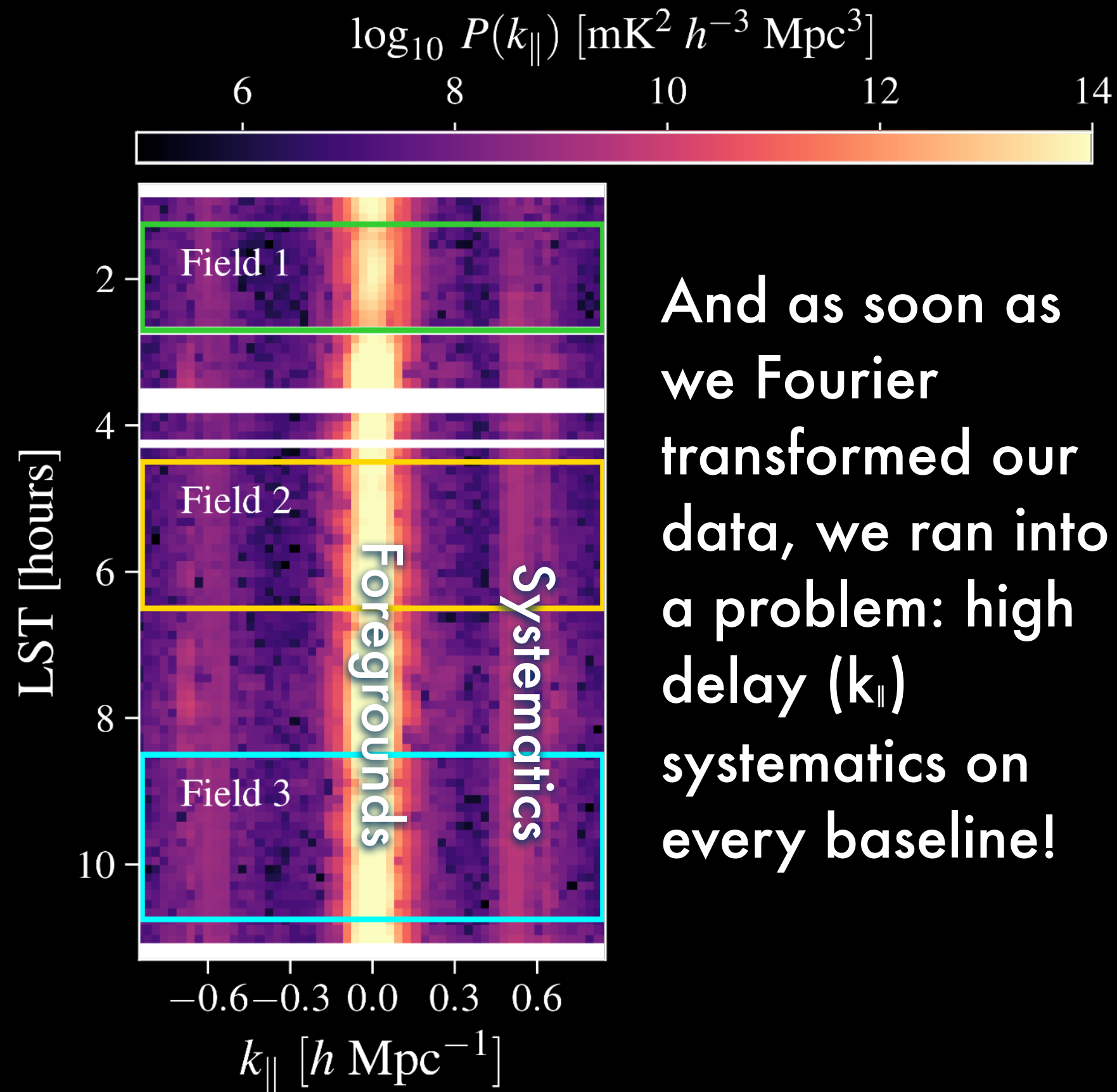
(65, 71) on 2458098



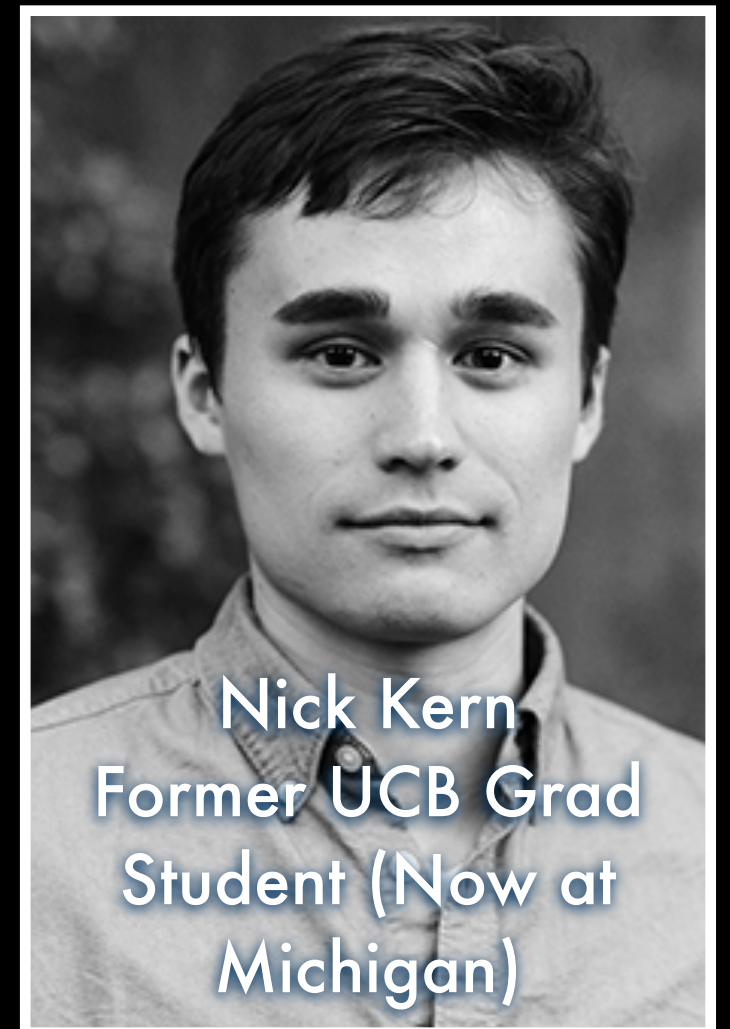
So we can keep integrating down to maximize sensitivity.

(65, 71) LST-Binned

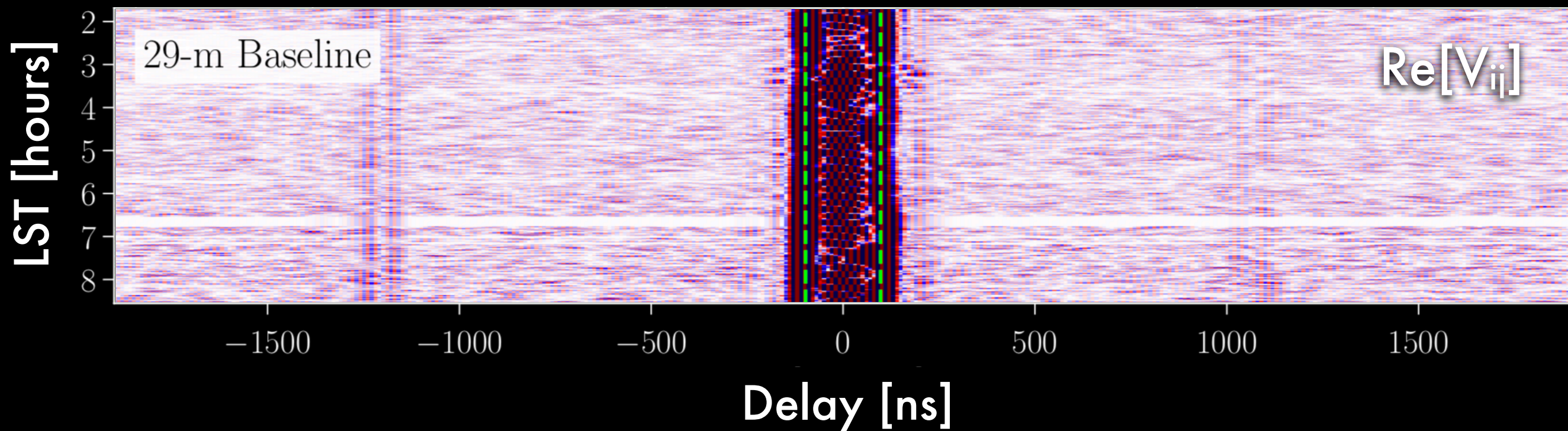




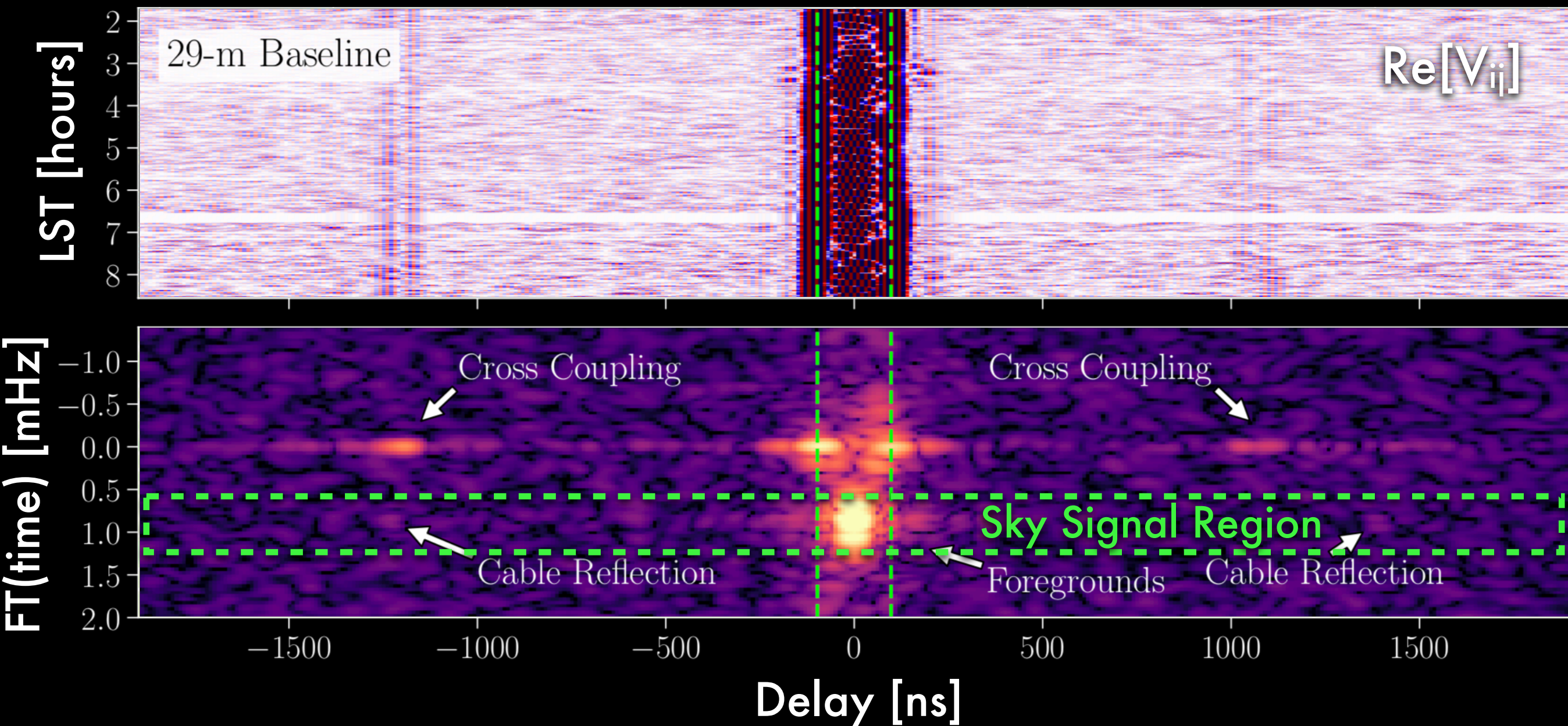
And as soon as
we Fourier
transformed our
data, we ran into
a problem: high
delay (k_{\parallel})
systematics on
every baseline!



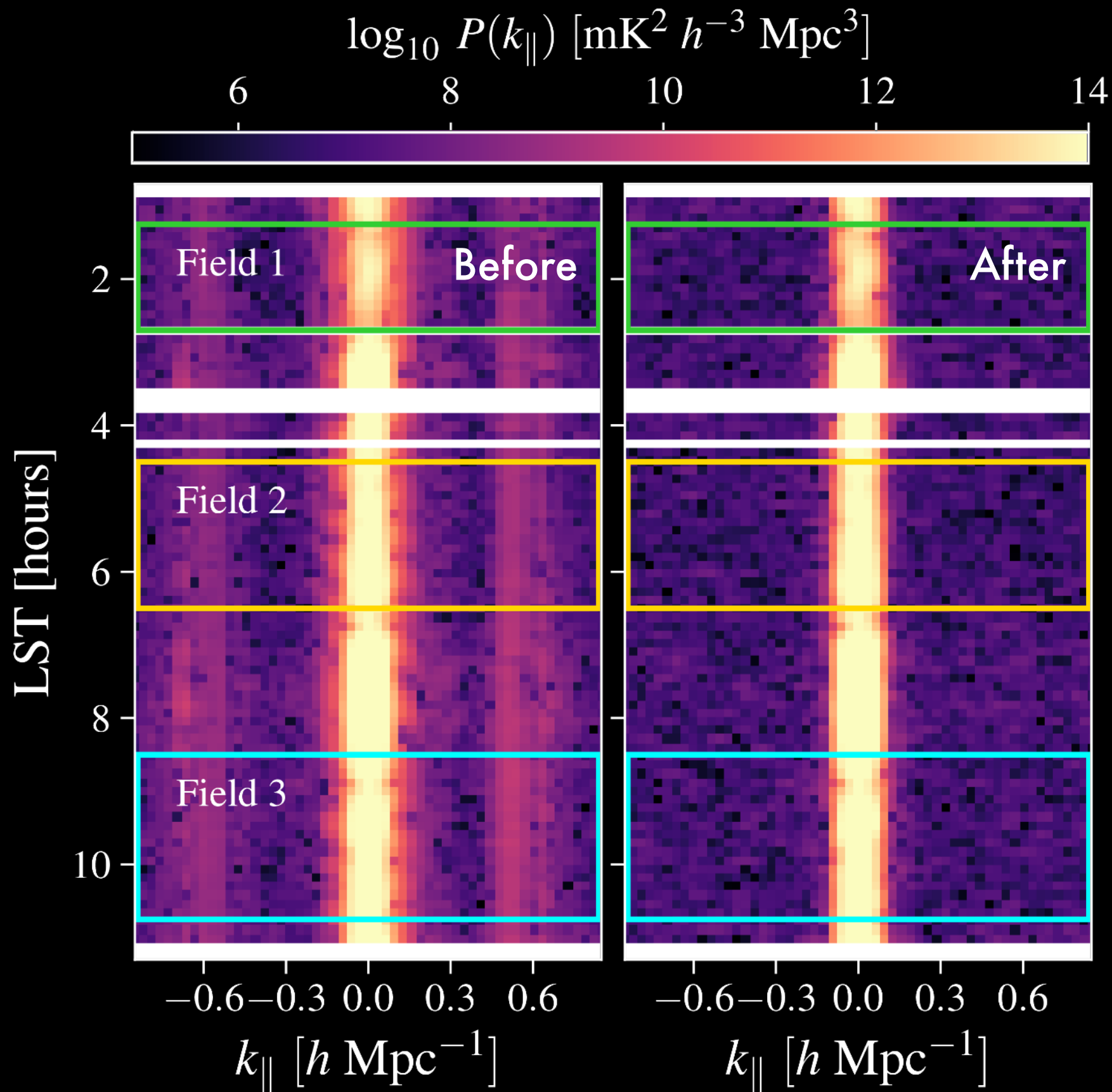
To understand this effect, we have to examine the temporal structure of the foregrounds and the systematics—how fast they “fringe.”



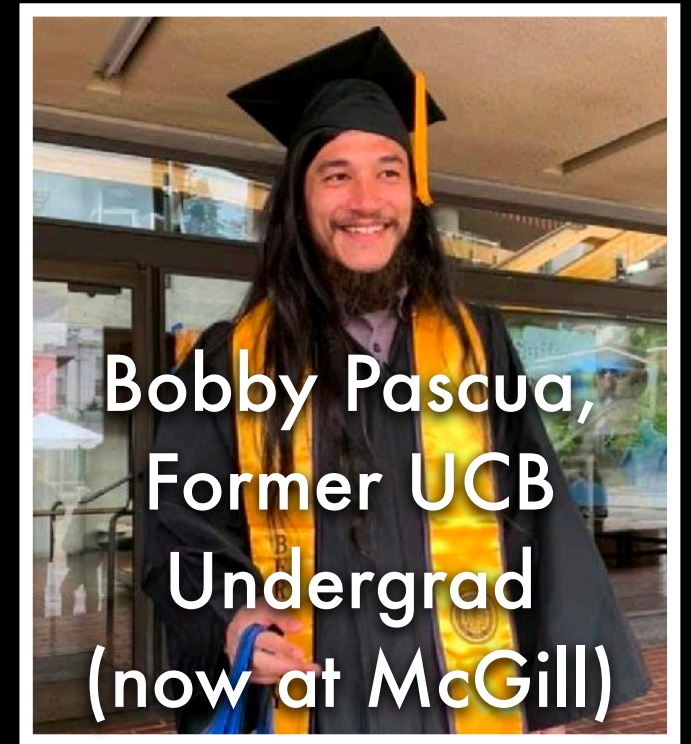
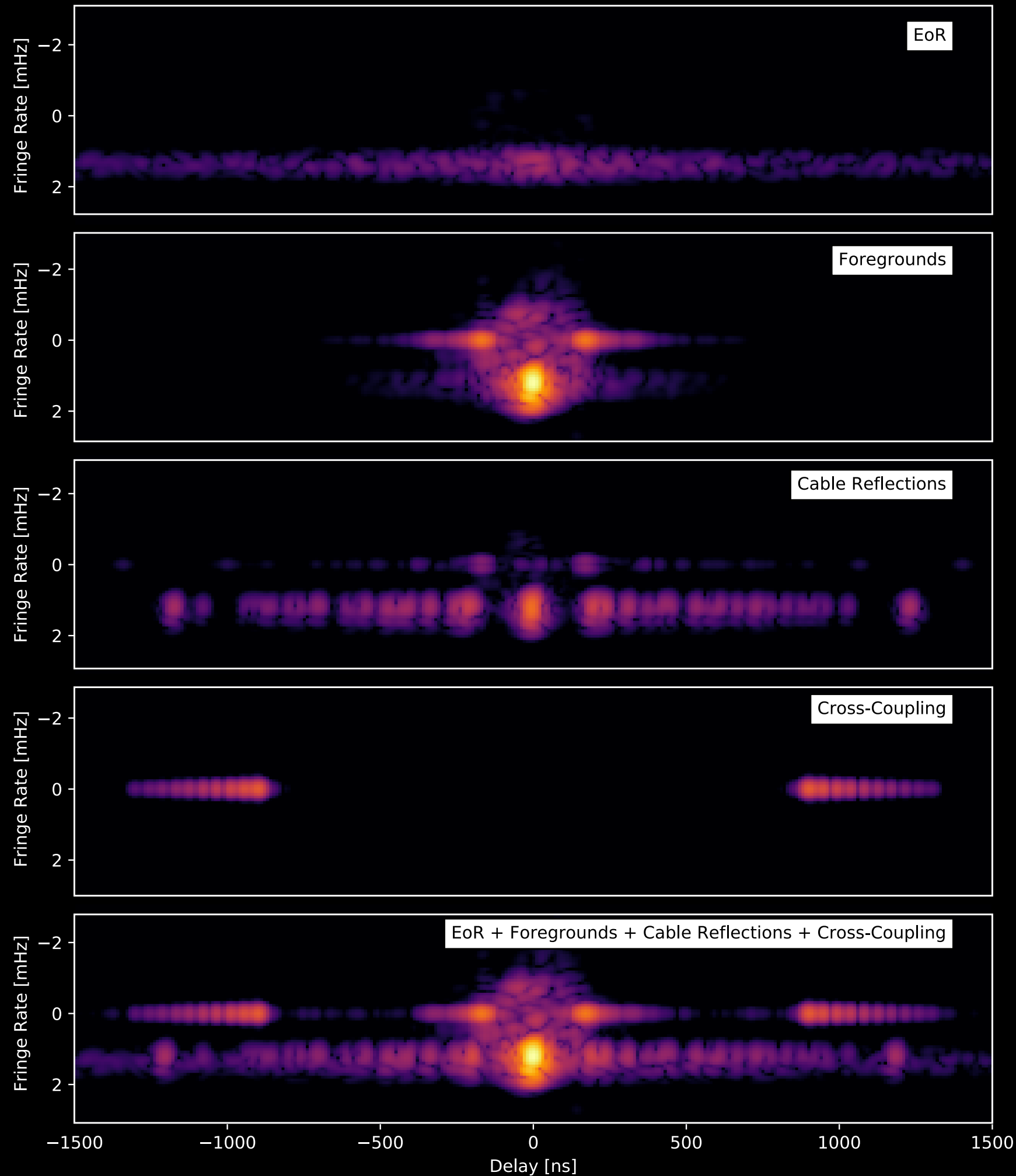
To understand this effect, we have to examine the temporal structure of the foregrounds and the systematics—how fast they “fringe.”



With our techniques for relatively lossless systematics removal, we're getting very close to the thermal noise limit.



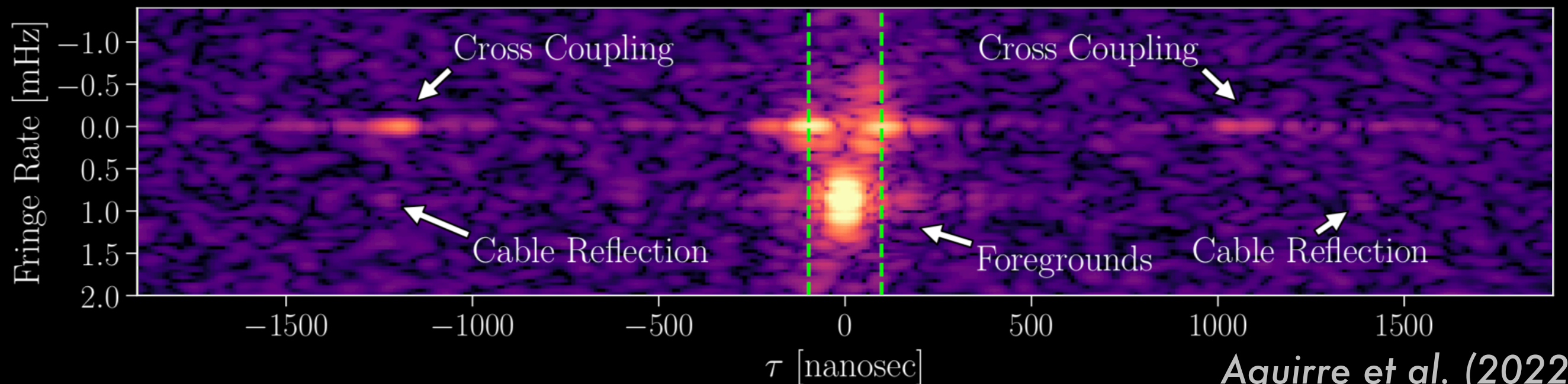
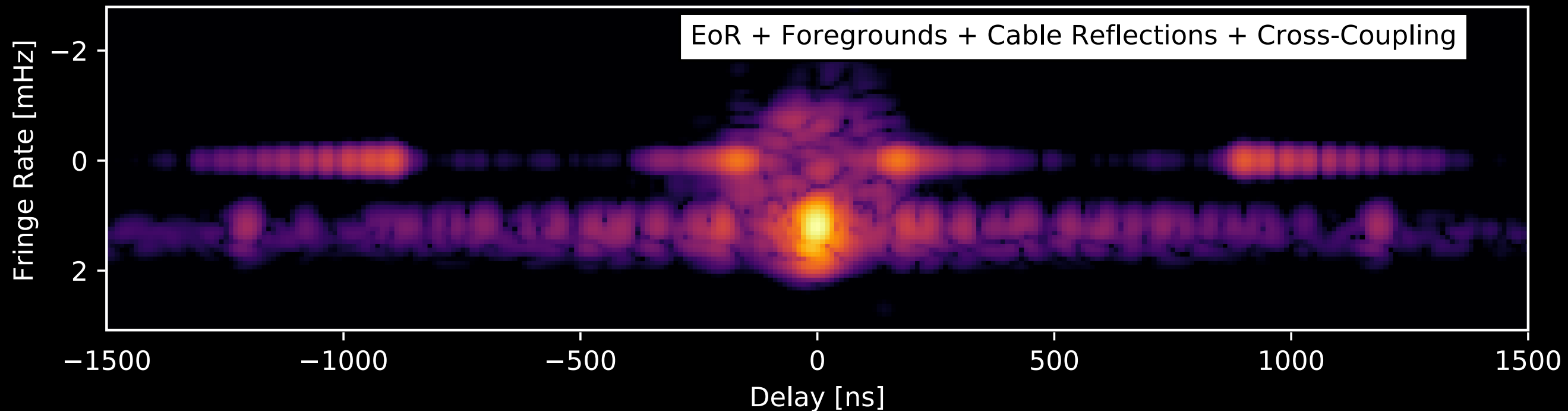
How are we building
confidence in our techniques
and quantifying signal loss?



We built end-to-end tests of analysis pipeline with simulated EoR, foregrounds, and systematics.

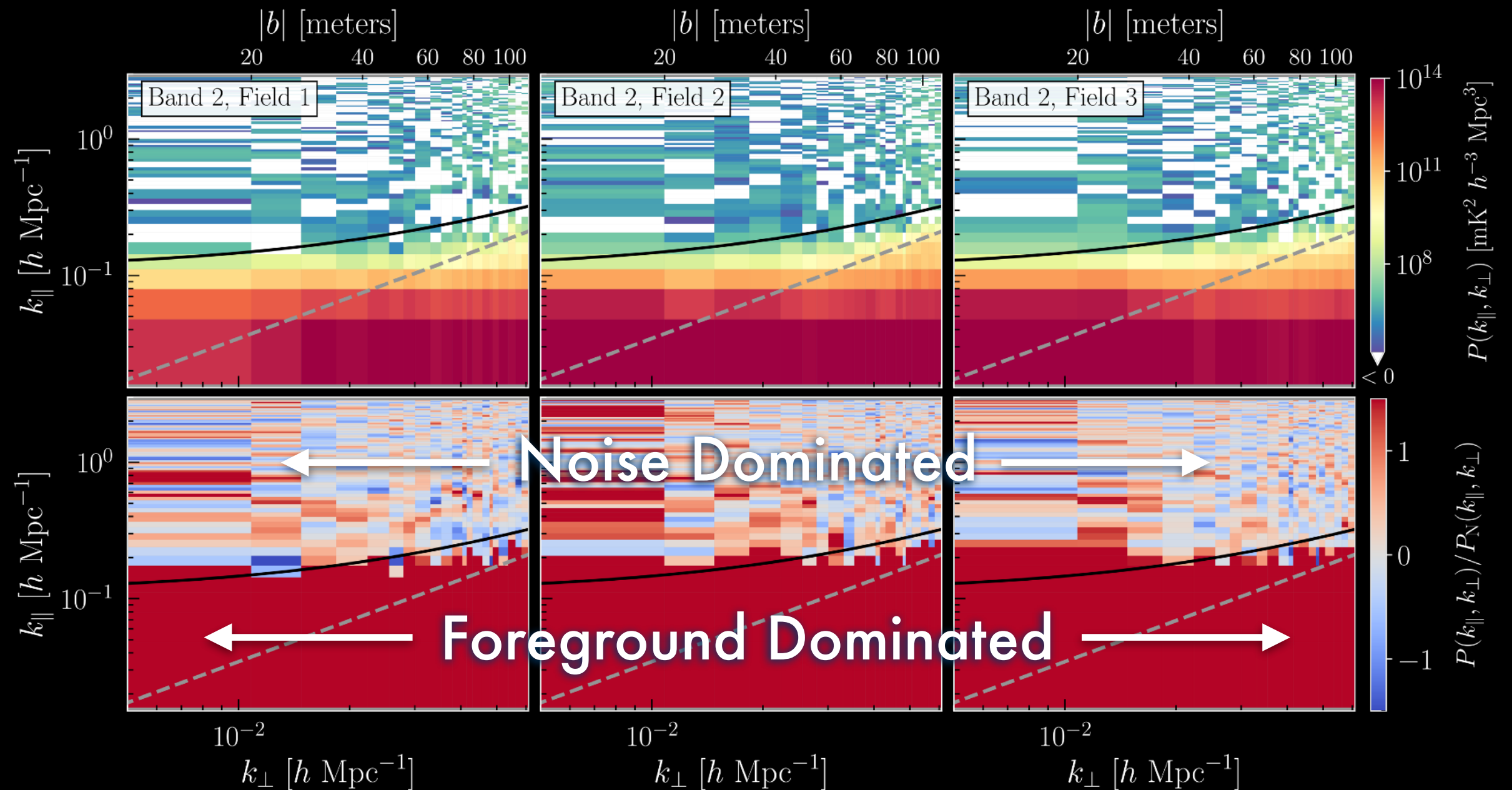
Aguirre et al. (2022)

The simulation is really starting to reflect the complexity of real data.

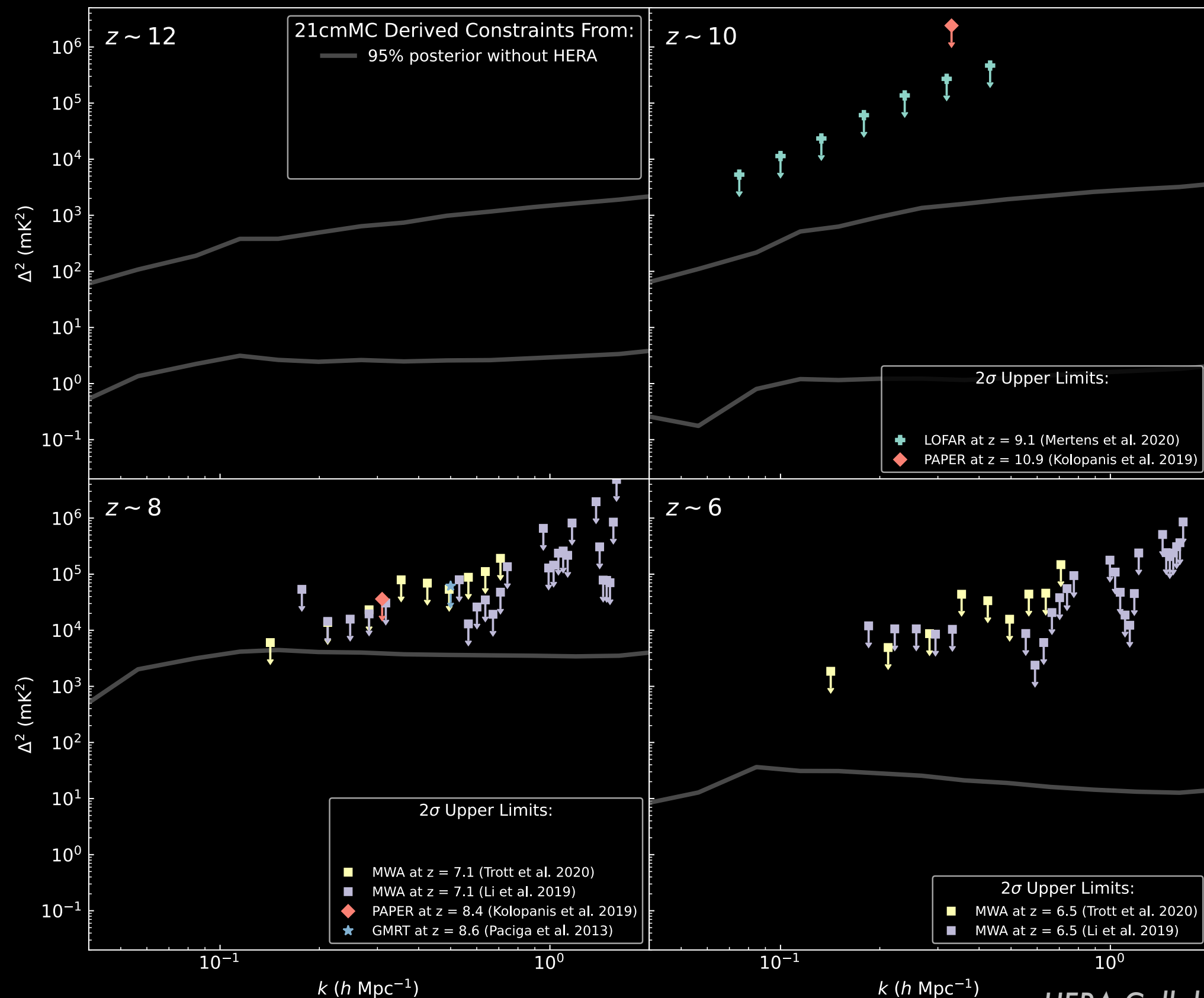


Aguirre et al. (2022)

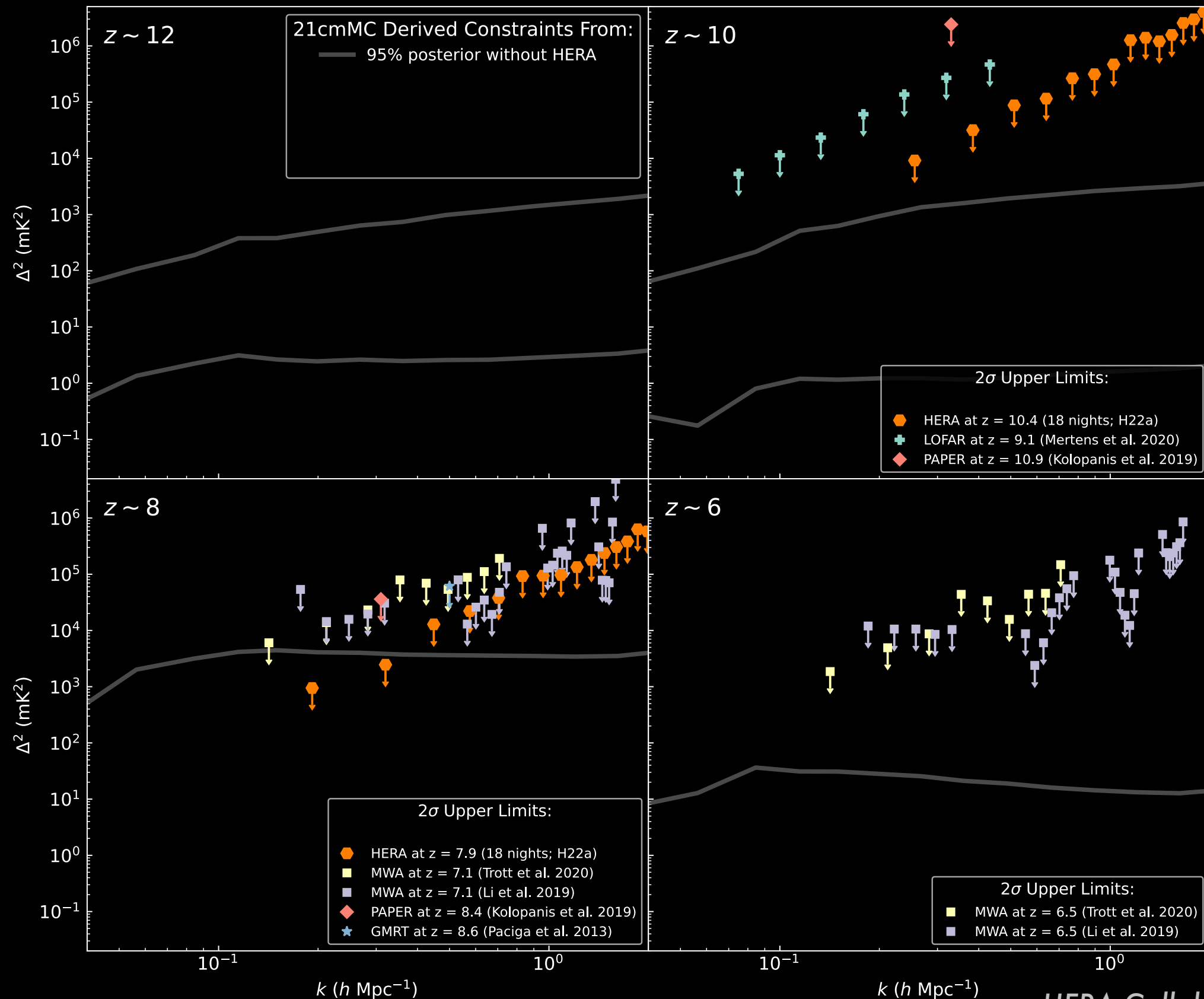
With high-delay systematics mitigated, we can finally form power spectra with those 18 nights.



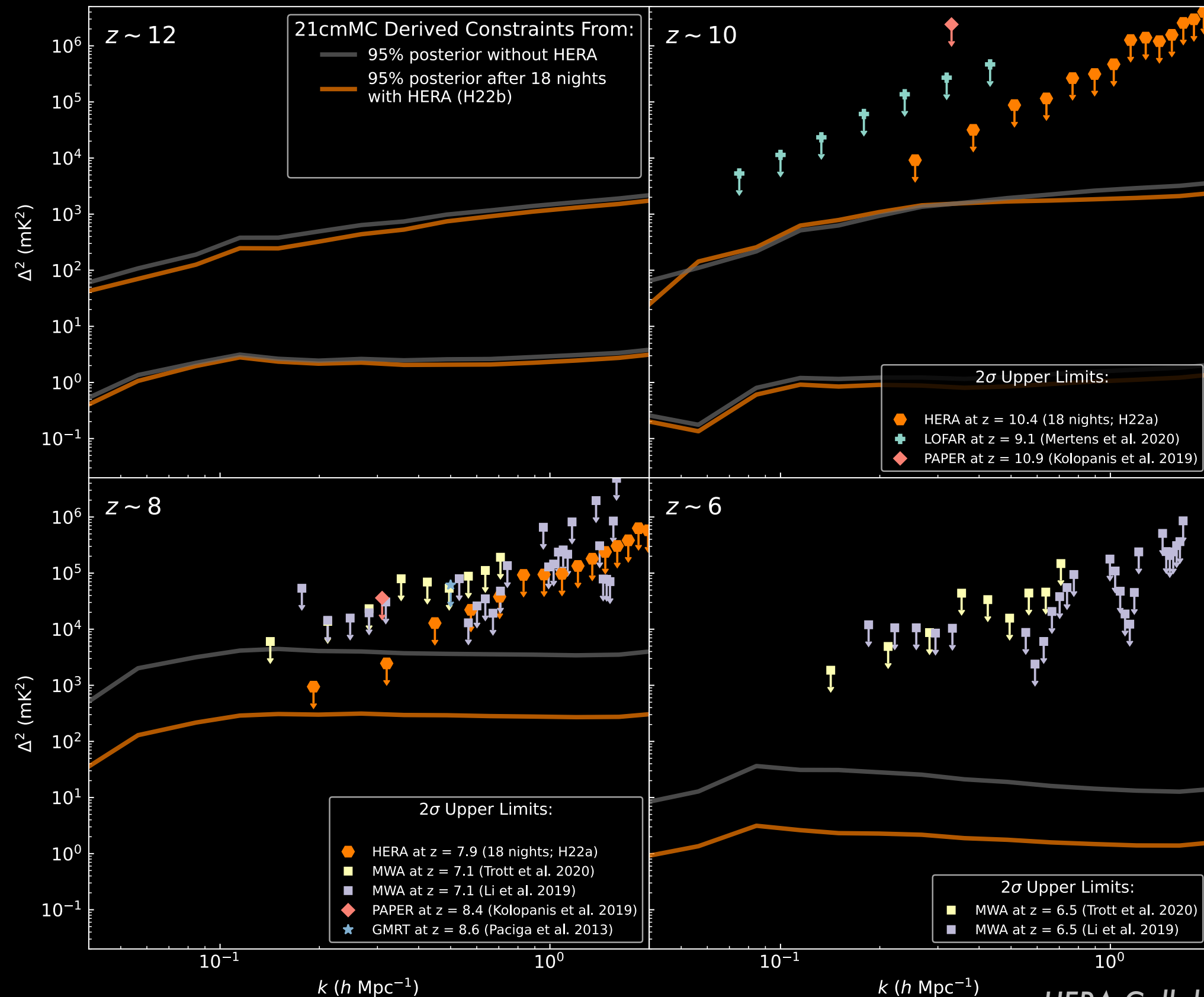
So again, here's where we were before HERA.



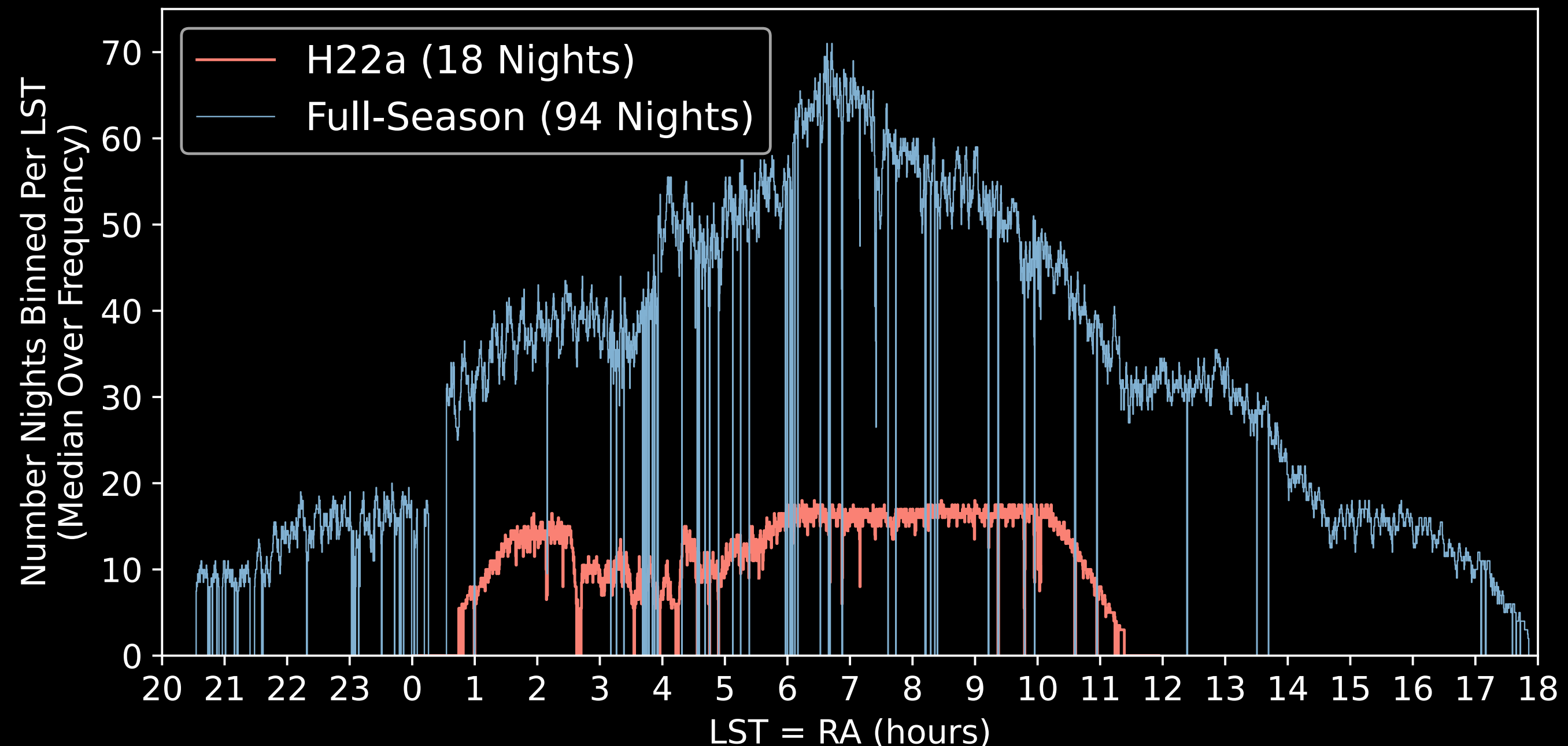
And here are HERA's first, world-leading upper limits on the 21 cm power spectrum.



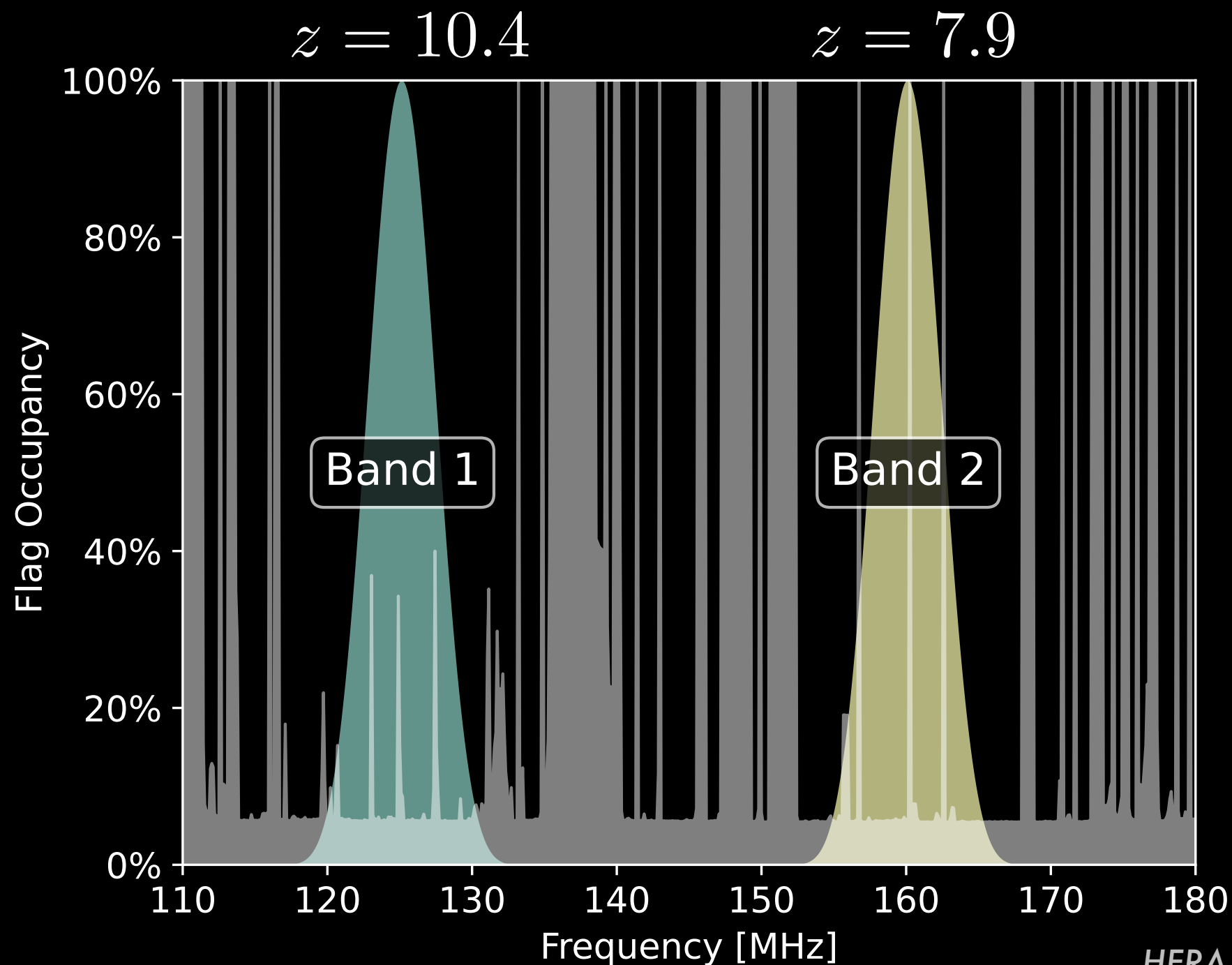
Which constrained the space of models, largely by ruling out an IGM unheated by X-rays at $z=8$.



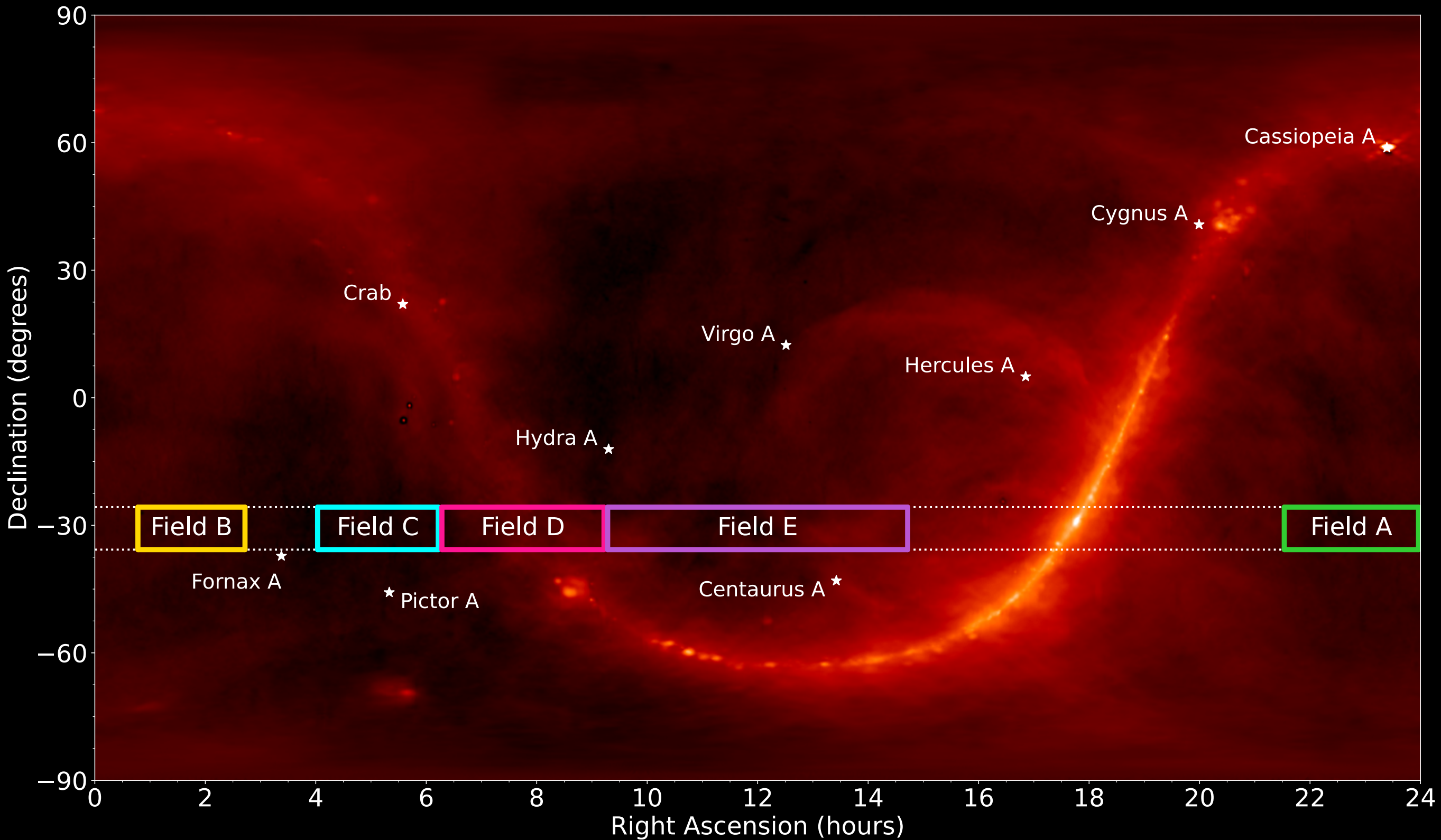
But all that was with only 18 nights of data. We had 94 good nights from HERA's first season.



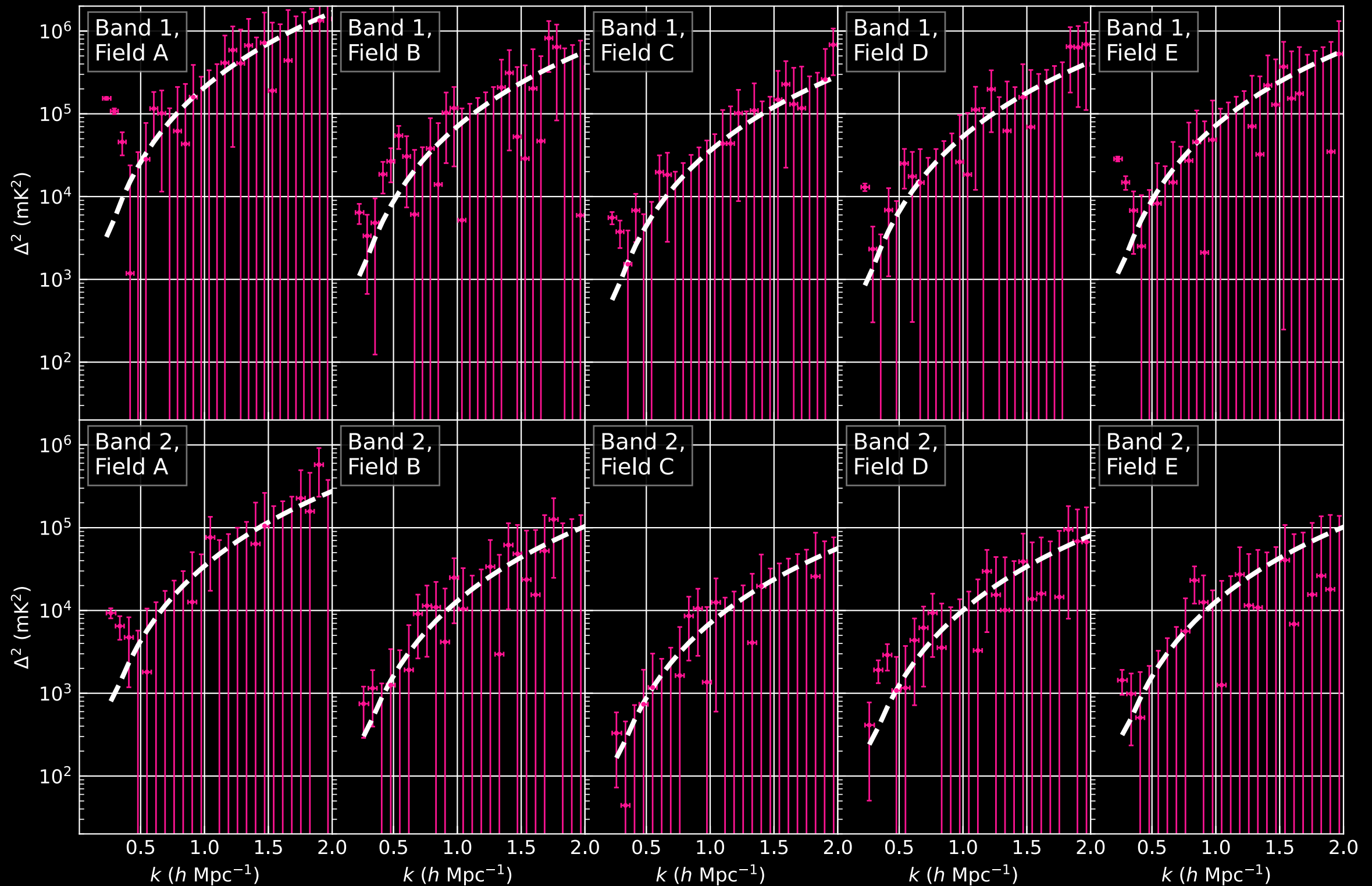
Adapting the same analysis techniques to the larger data set, we picked two redshift bands with minimal RFI contamination.



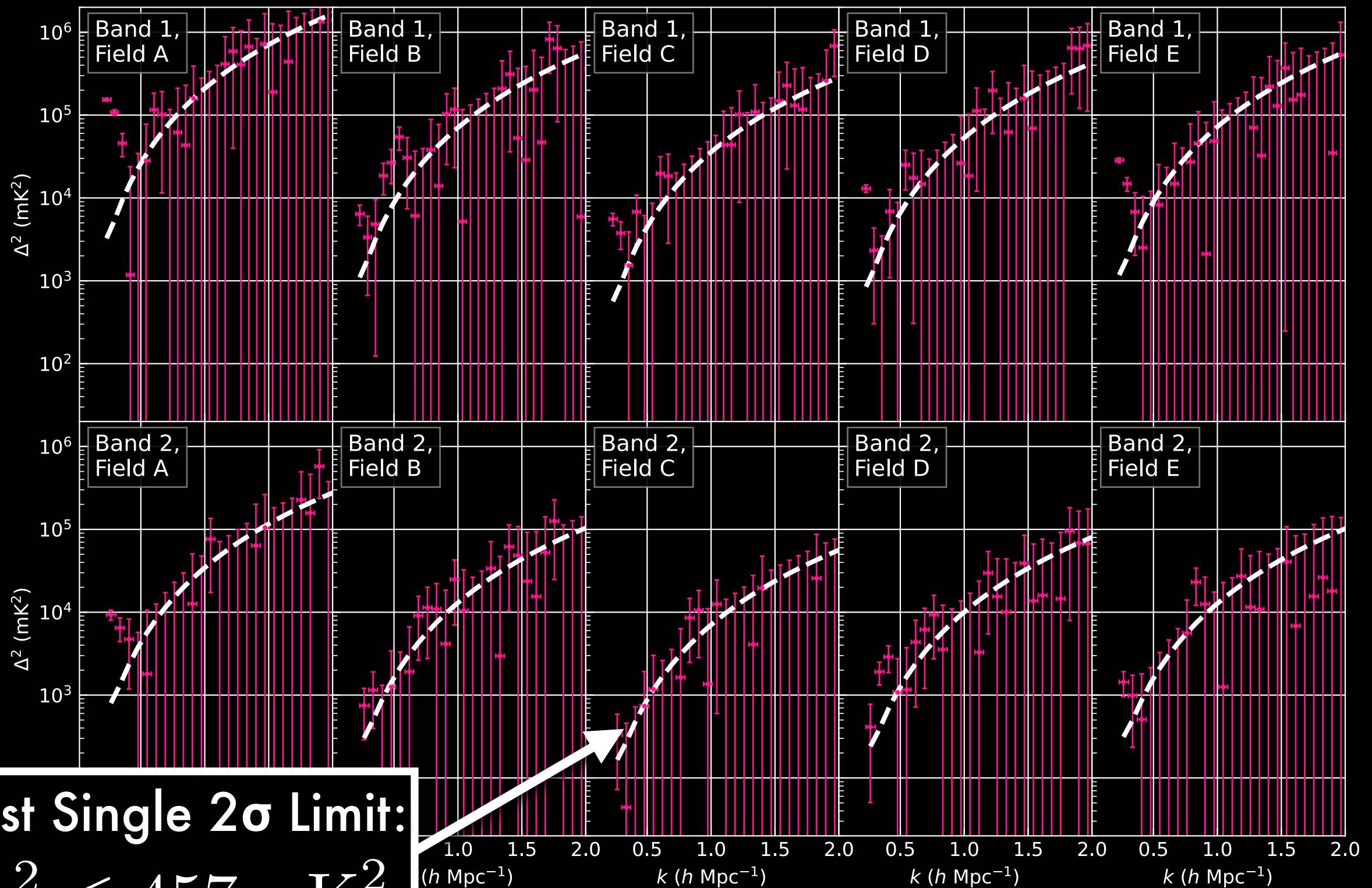
We divided our LSTs observed into five fields.



And thus set power spectrum upper limits across bands and fields.



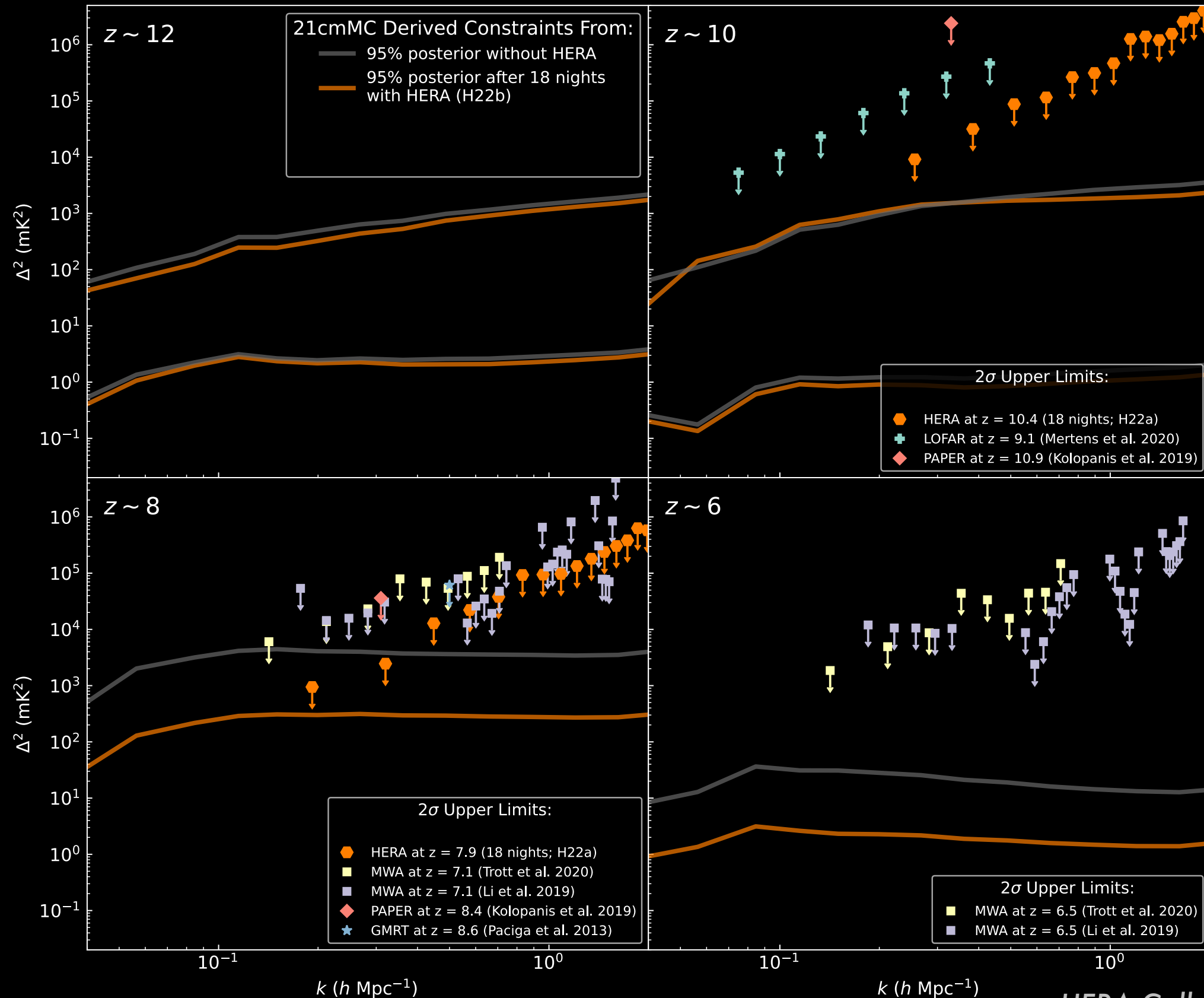
And thus set power spectrum upper limits across bands and fields.



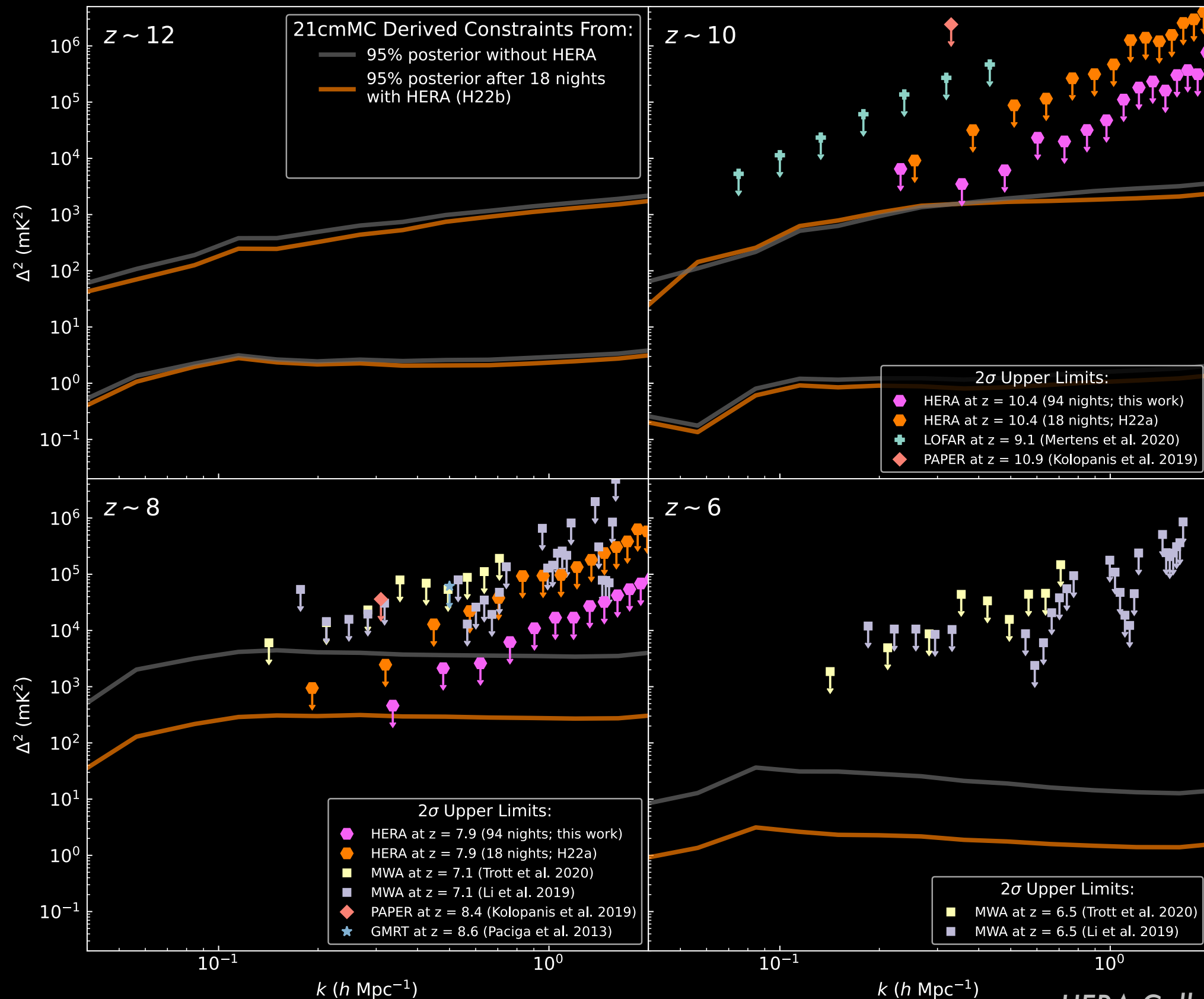
Best Single 2σ Limit:

$$\Delta^2 \leq 457 \text{ mK}^2$$

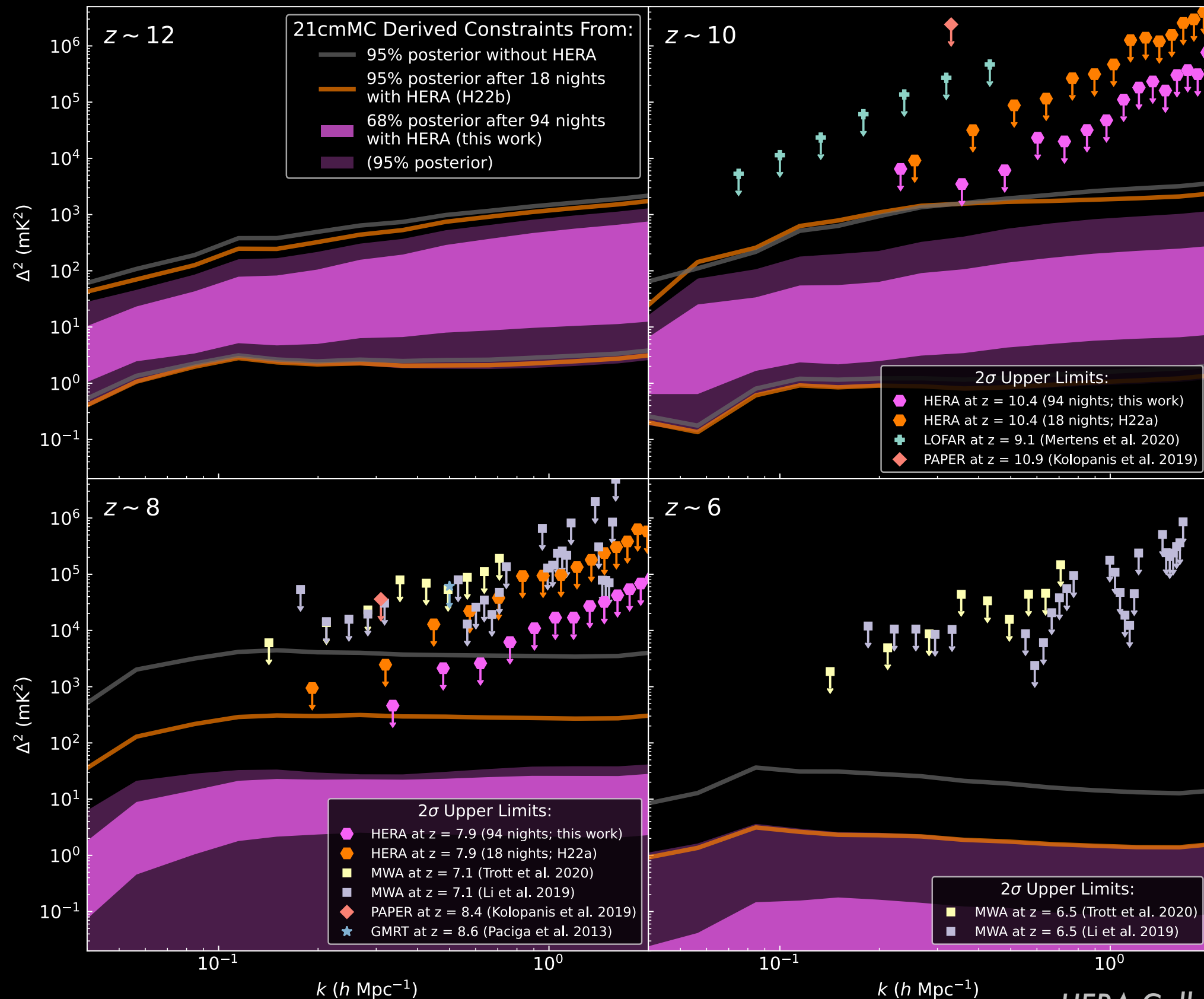
So here's where we were again...



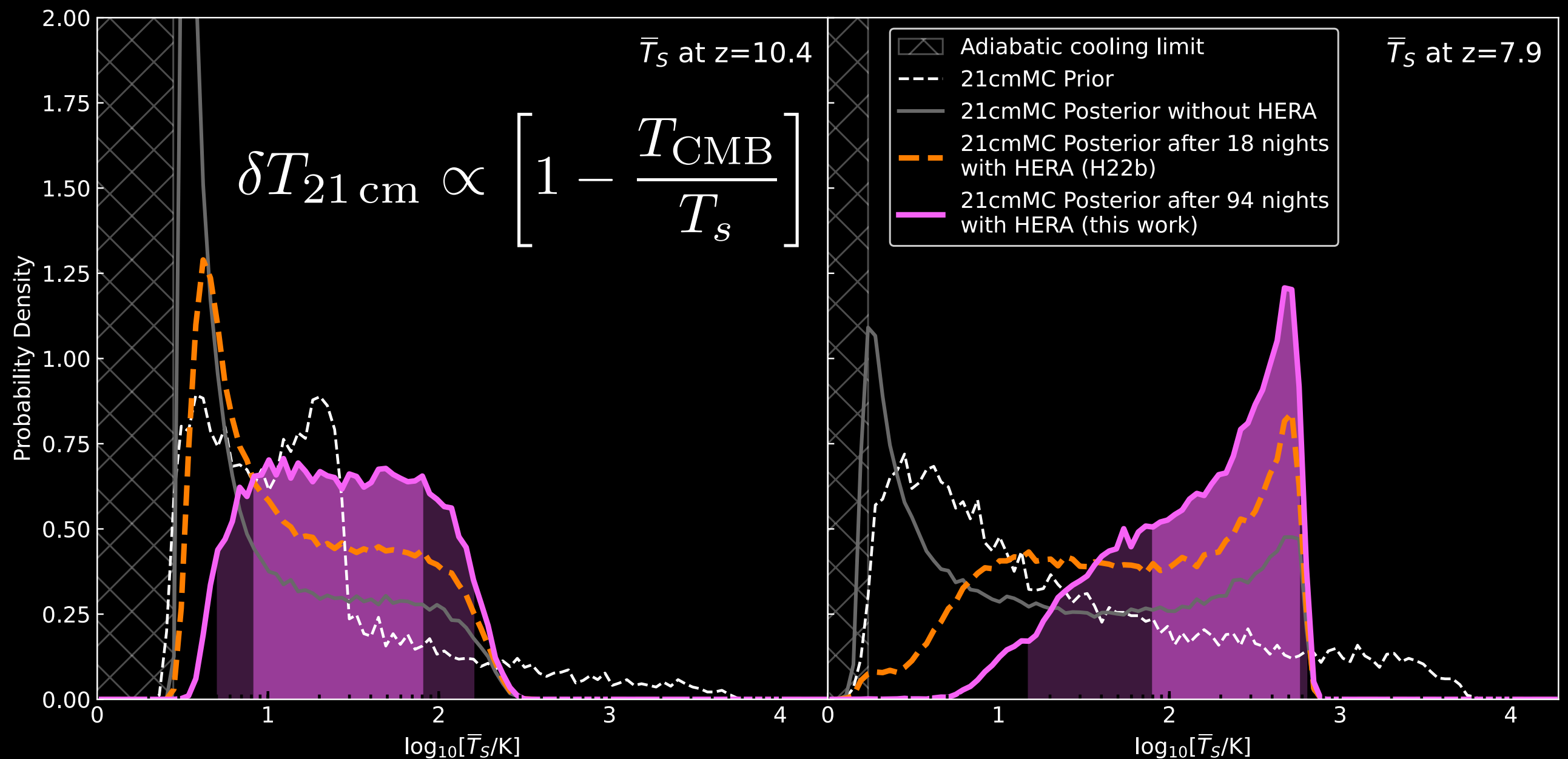
With a full season, our limits come down by more than a factor of 2 at both redshifts.



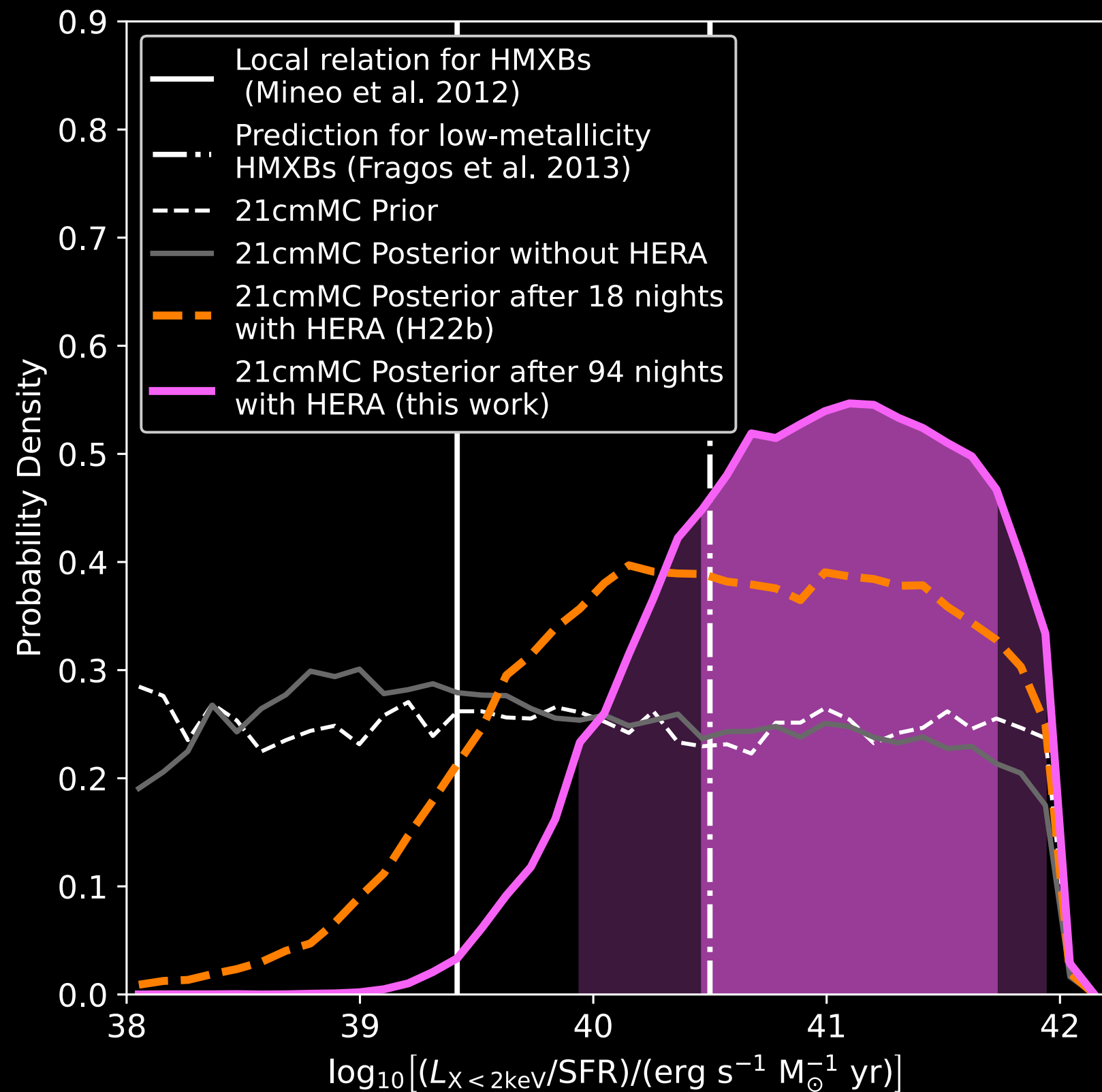
Our posterior for the power spectrum with 21cmMC tightens substantially.



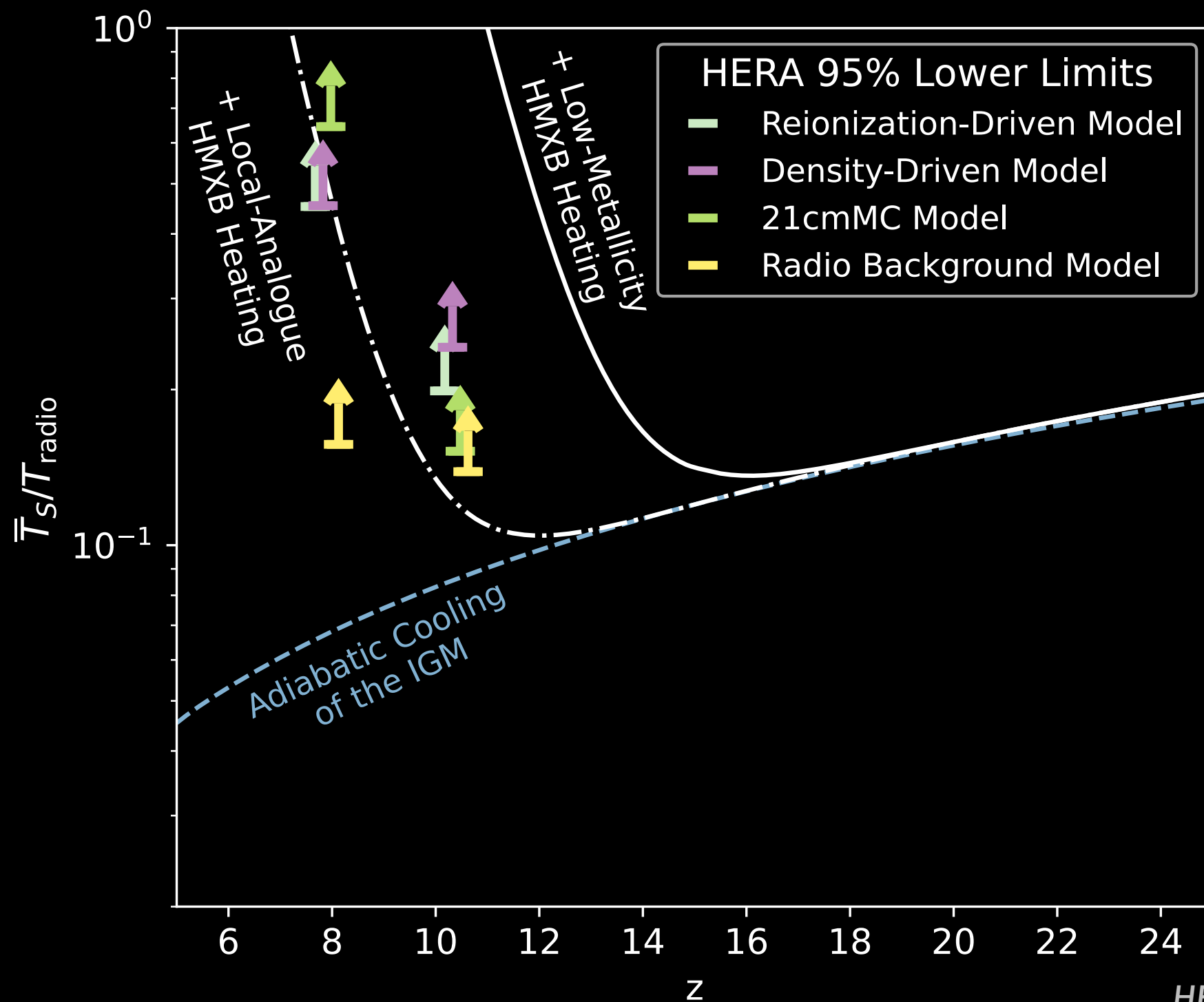
The big shift comes from showing the IGM was heated by $z=10.4$, since a cold IGM produces a bright 21 cm signal.



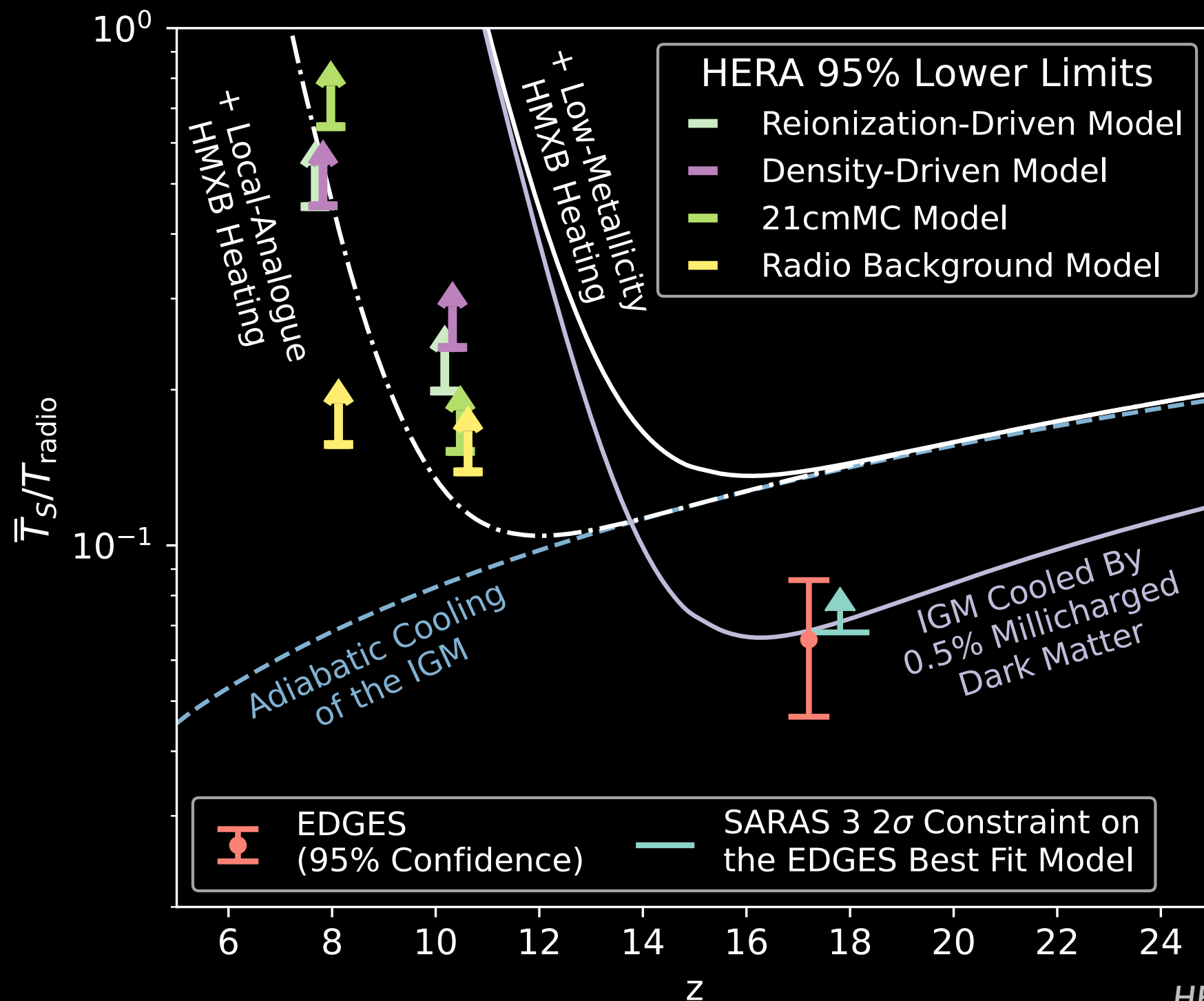
If the IGM was heated by high-mass X-ray binaries — as is generally believed — this result rules out high-metallicity HMXBs (which are less X-ray-efficient per unit SFR) and thus requires heating driven by evolved low-metallicity stars.



Four independent theoretical models agree:
the IGM was heated before $z = 10.4$.



However, we are not yet able to say much about the tension between EDGES and SARAS or the exotic models invoked to explain EDGES.



What's next for HERA?



HERA

Array Core



Phase I
94 Nights
~40 Good Antennas

100 m



HERA

Array Core

Phase II

~550 good nights (so far)

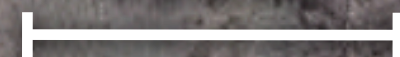
~190 good antennas
(on average)

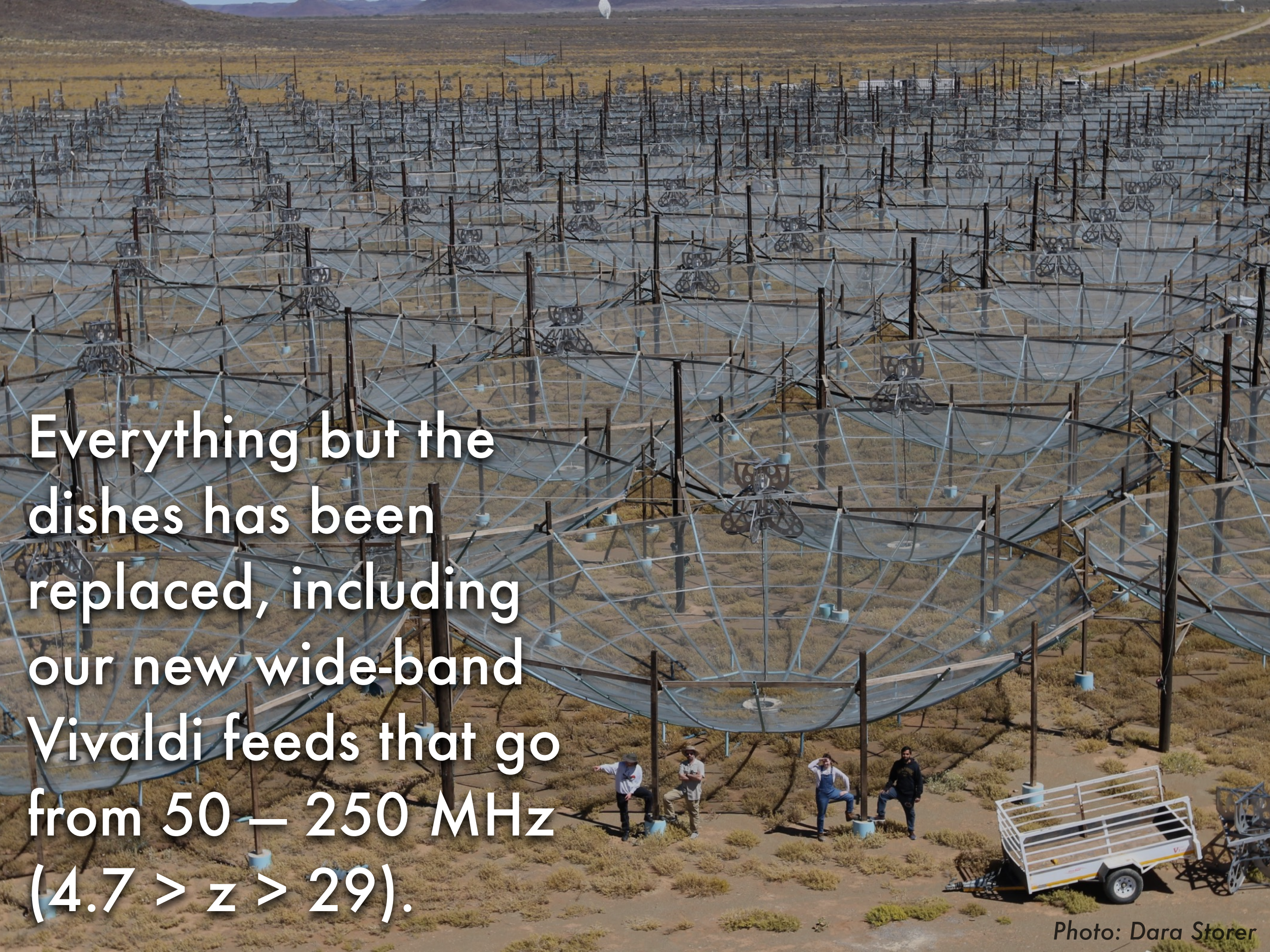
Phase I

94 Nights

~40 Good Antennas

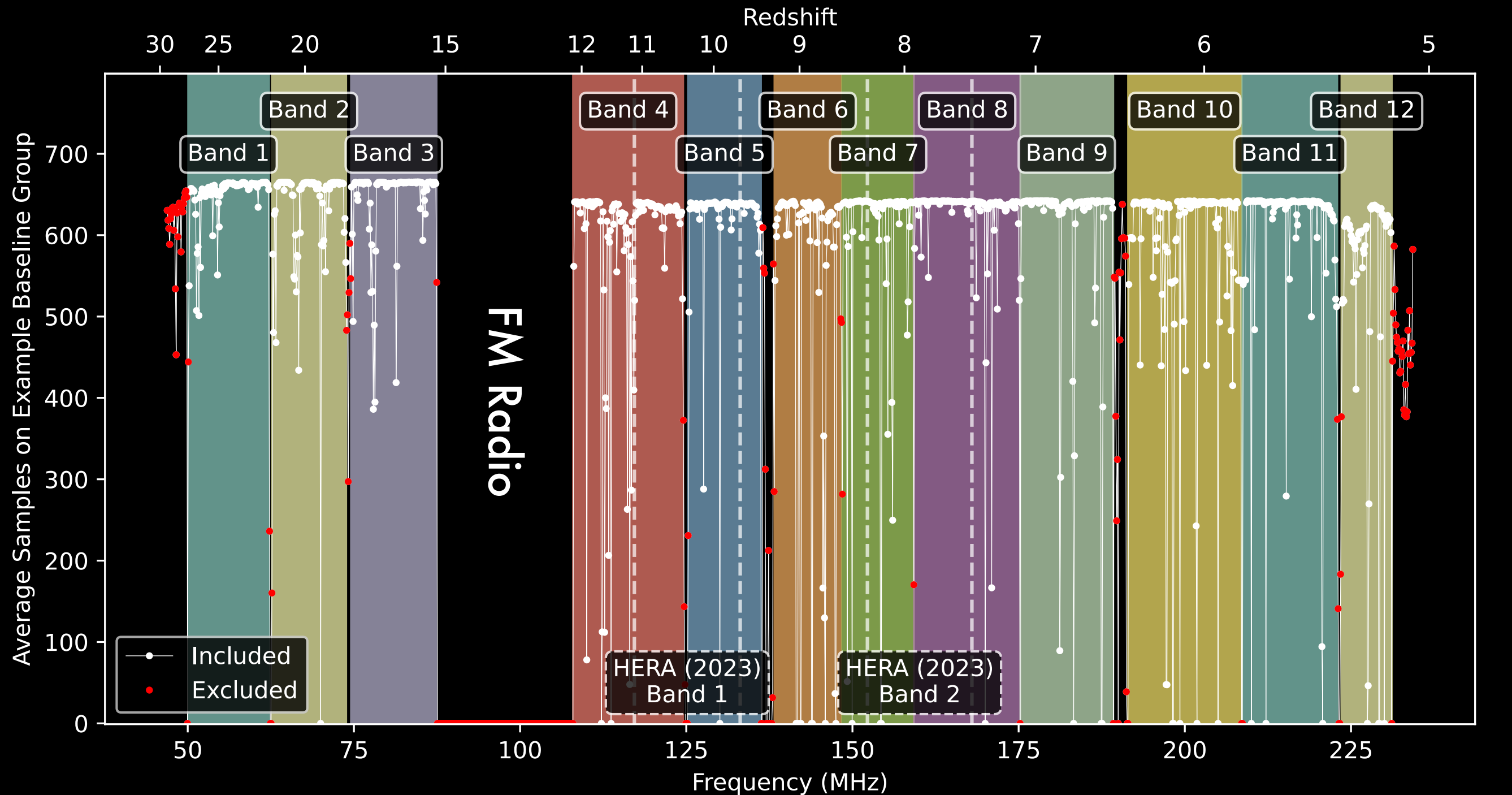
100 m



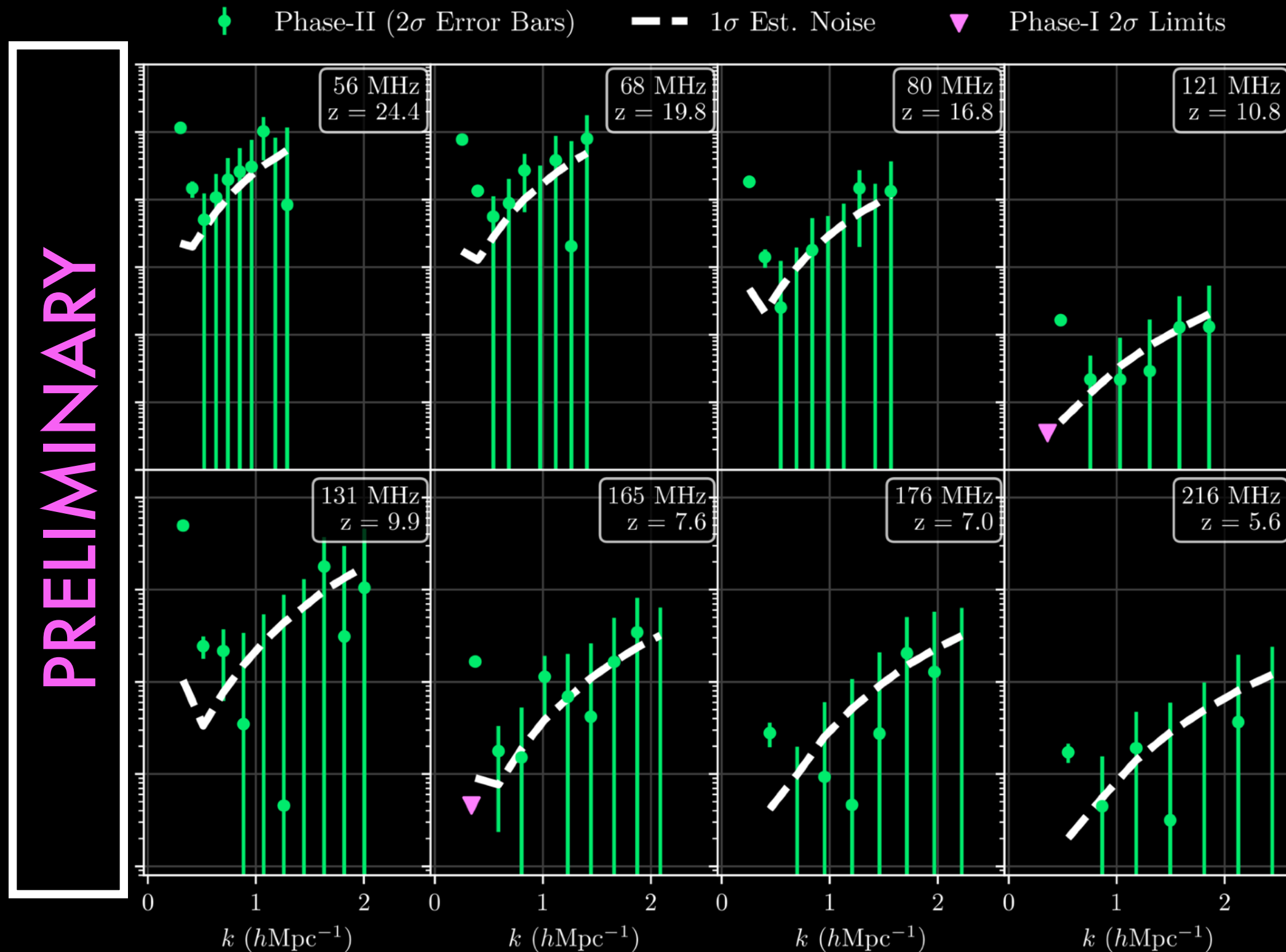


Everything but the
dishes has been
replaced, including
our new wide-band
Vivaldi feeds that go
from 50 – 250 MHz
($4.7 > z > 29$).

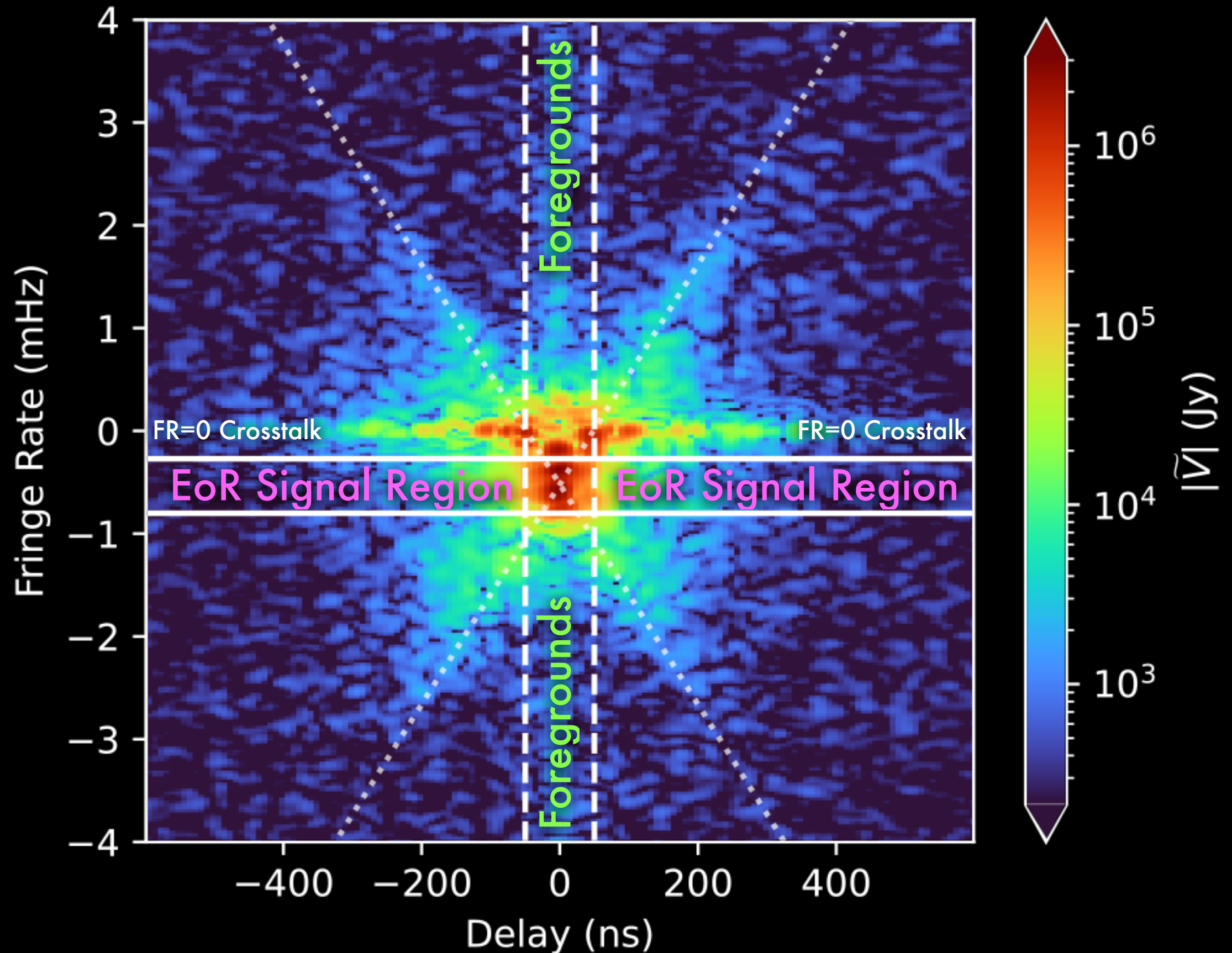
With this data set, we're looking at a much wider range of redshifts, including in the pre-reionization epoch.



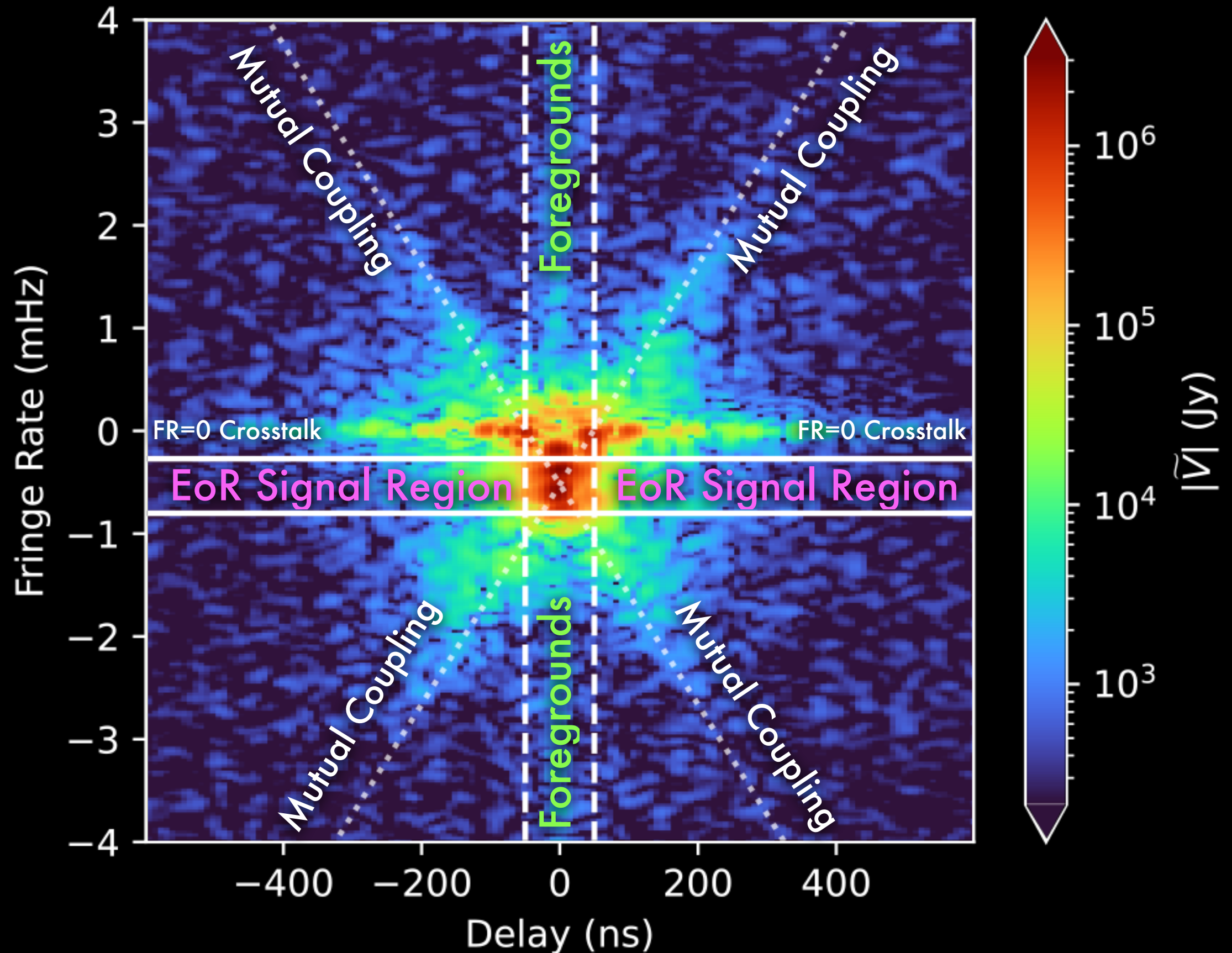
We're now writing up limits from the first 14 nights of Phase II data, testing our brand new pipeline.



But the new array presents new systematics...



But the new array presents new systematics...

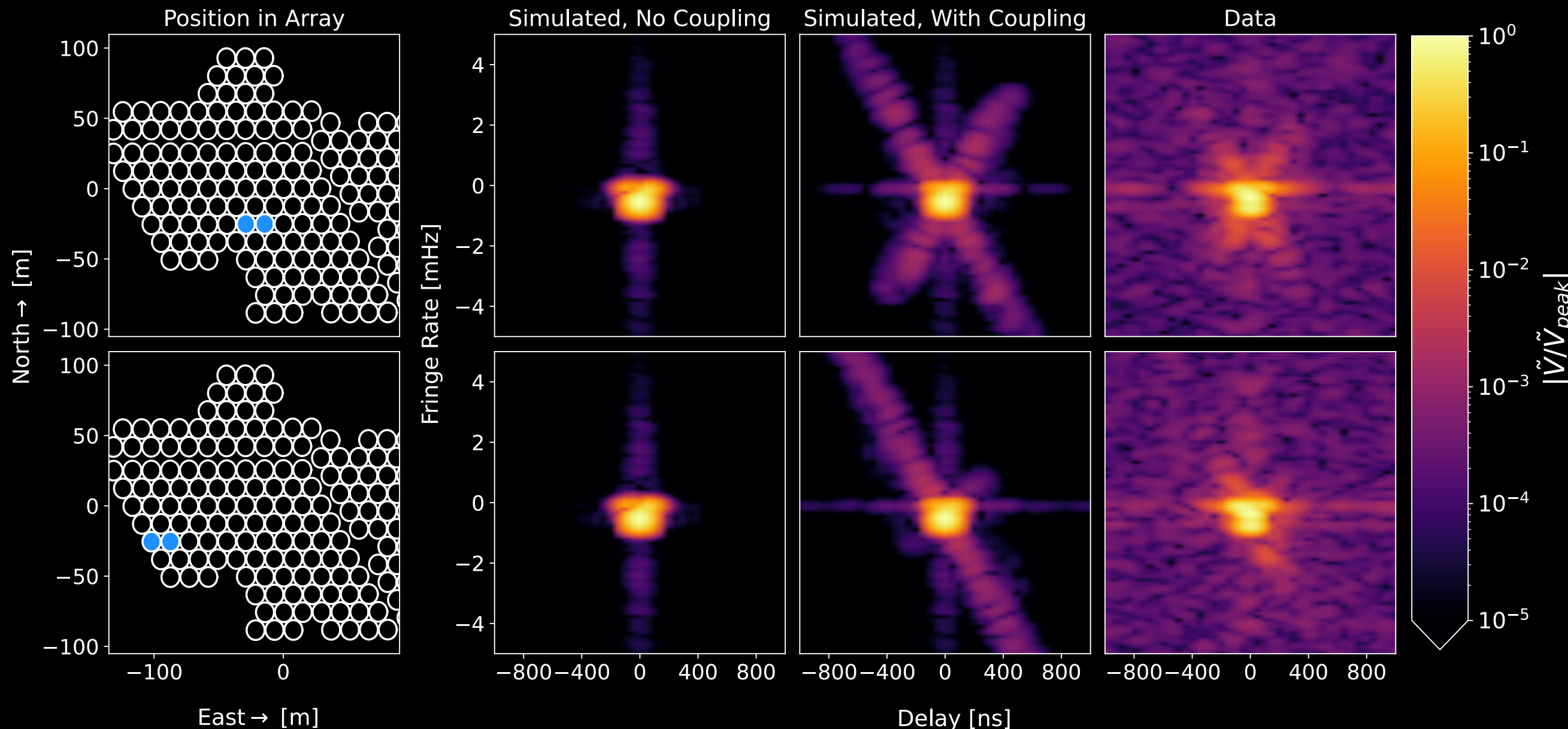




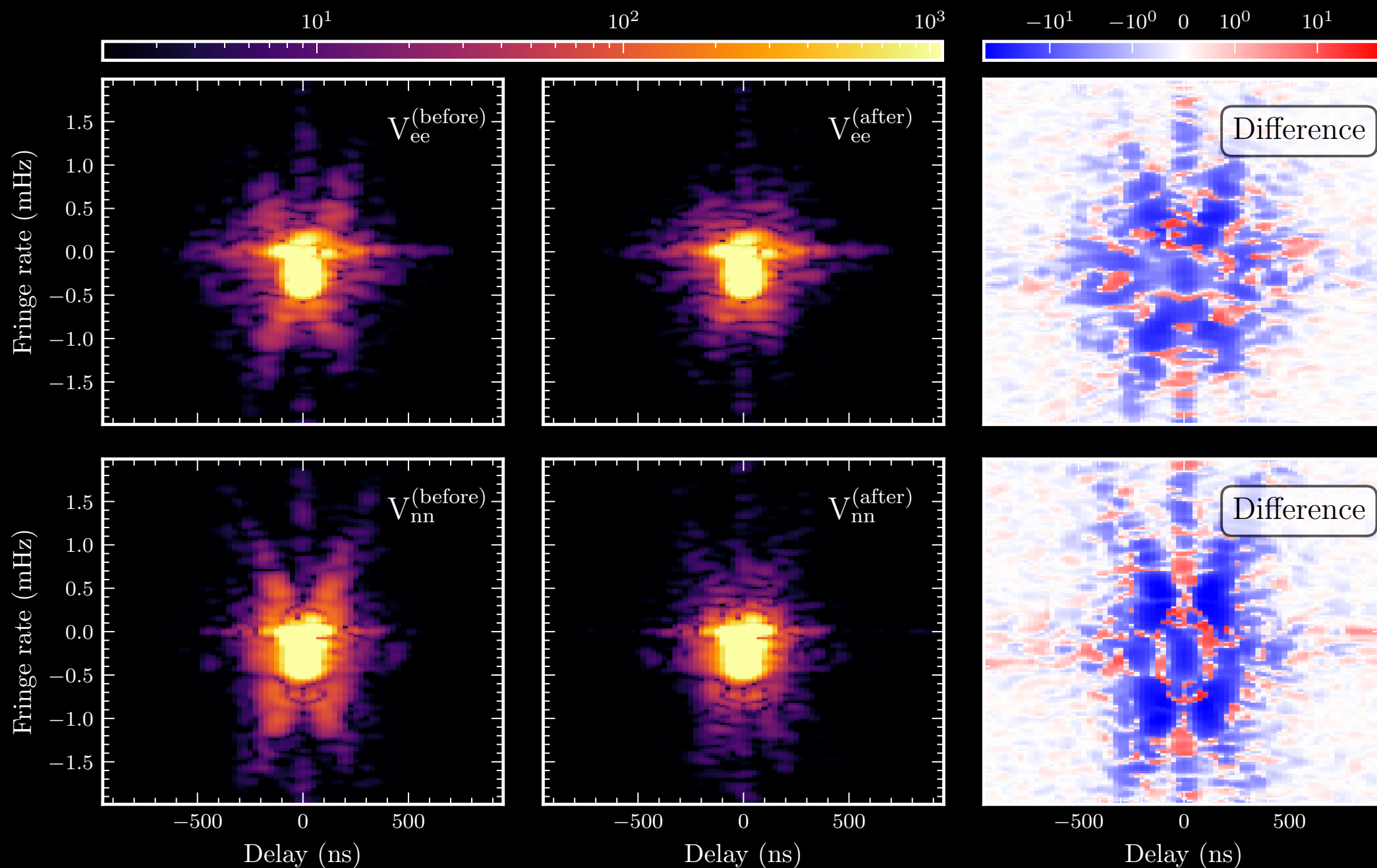
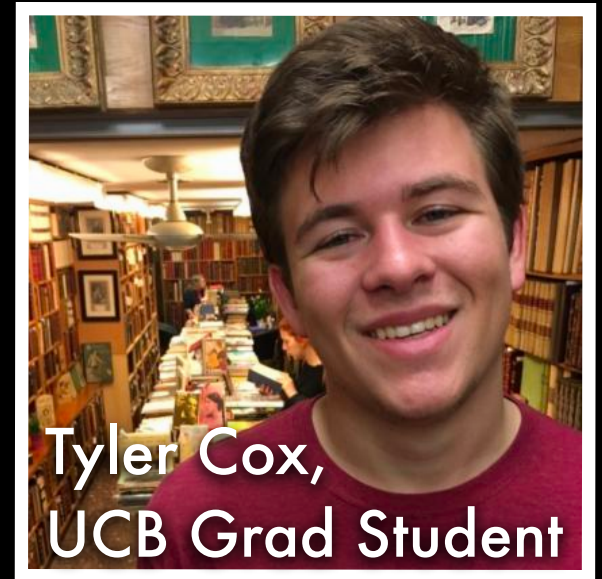
HERA's dishes aren't the most electromagnetically isolated elements....

Photo: Dara Storer

We now have a good qualitative picture for how leakage of visibilities into one another via mutual coupling causes this “X.”



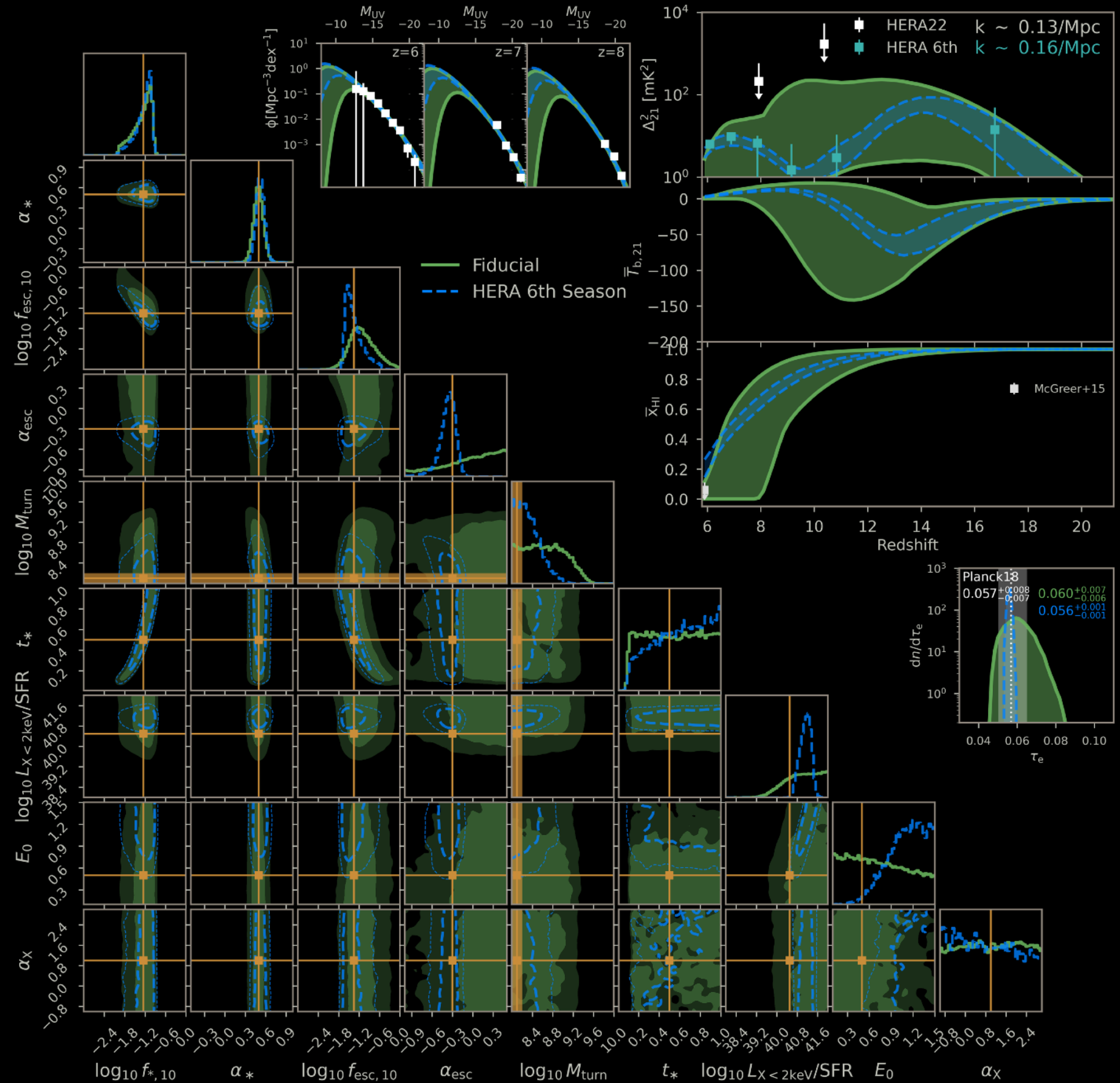
We're making good progress modeling and suppressing the contribution from mutual coupling on our first 14 nights of data.



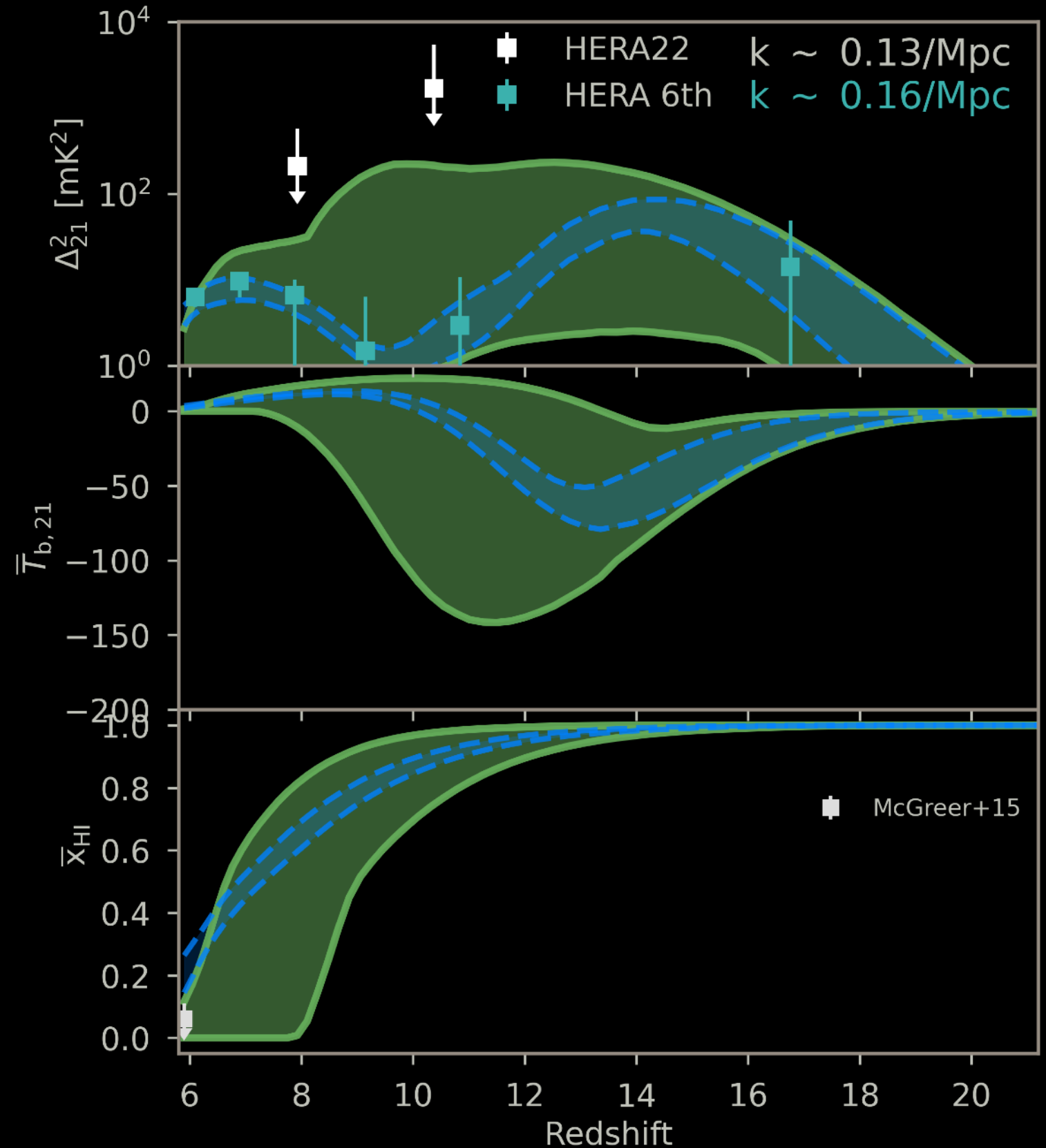
PRELIMINARY

**What do we expect to
learn from Phase II?**

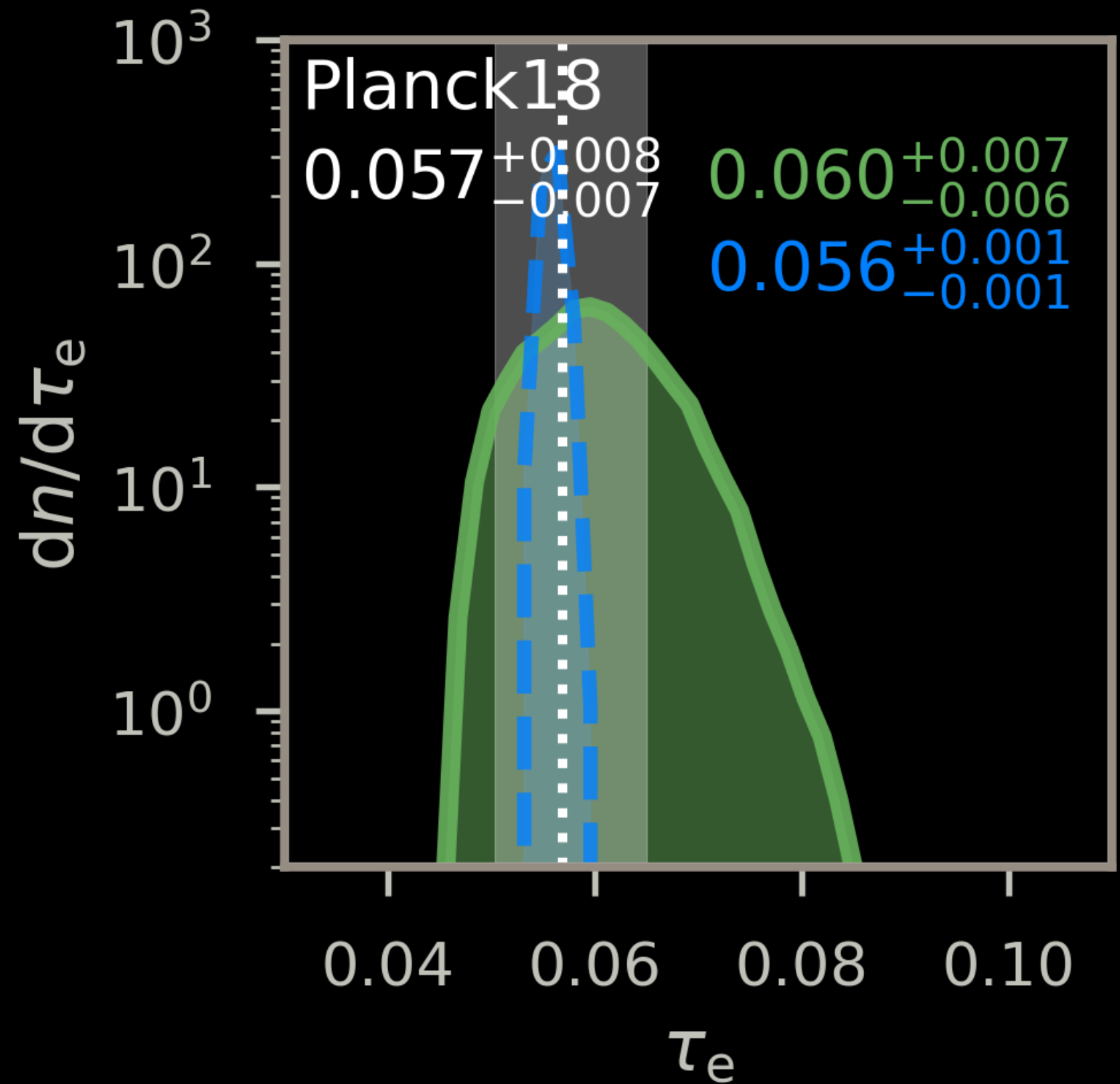
The HERA theory team has already spun up on their new 21CMFAST emulator to forecast HERA's sensitivity to a fiducial model for IGM heating and reionization.



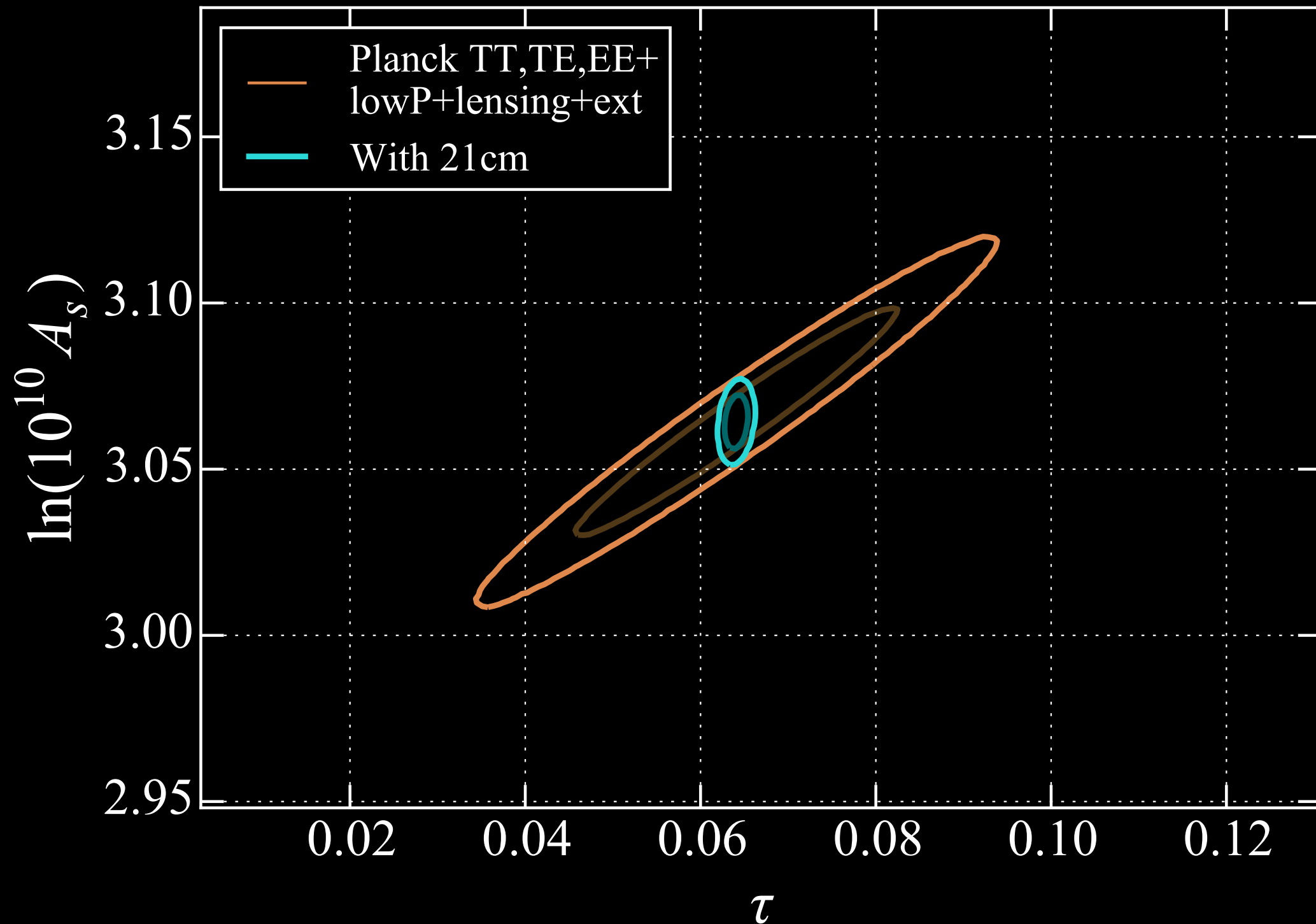
If we can get as close to the wedge as we did in the published limits, the data in hand is sensitive enough to make a $\sim 5\sigma$ detection of a fiducial EoR power spectrum.



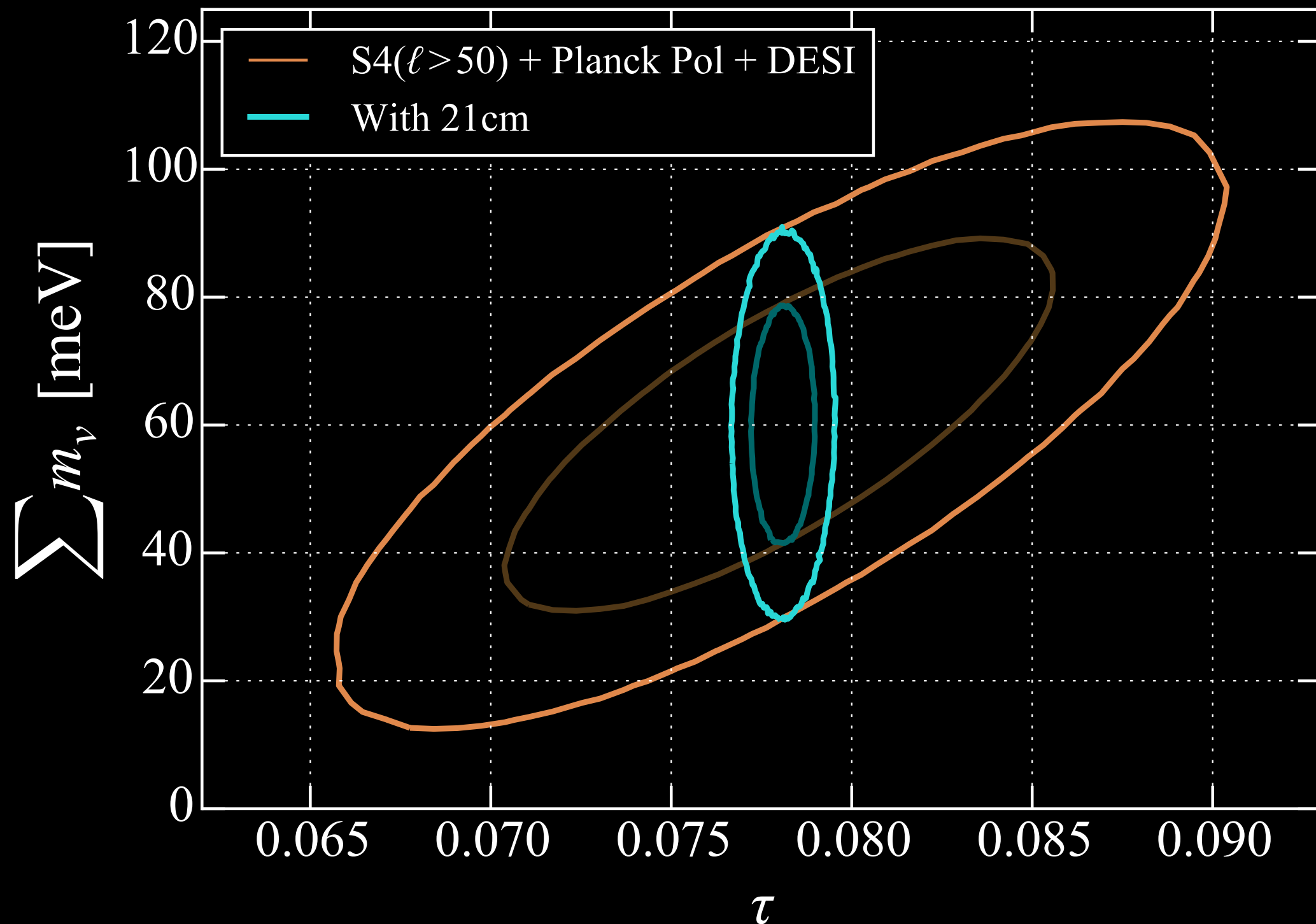
And we'll make
our first major
contribution to
cosmology: a ~6
times tighter
constraint on τ ,
improving CMB
constraints of more
“fundamental”
cosmological
parameters.



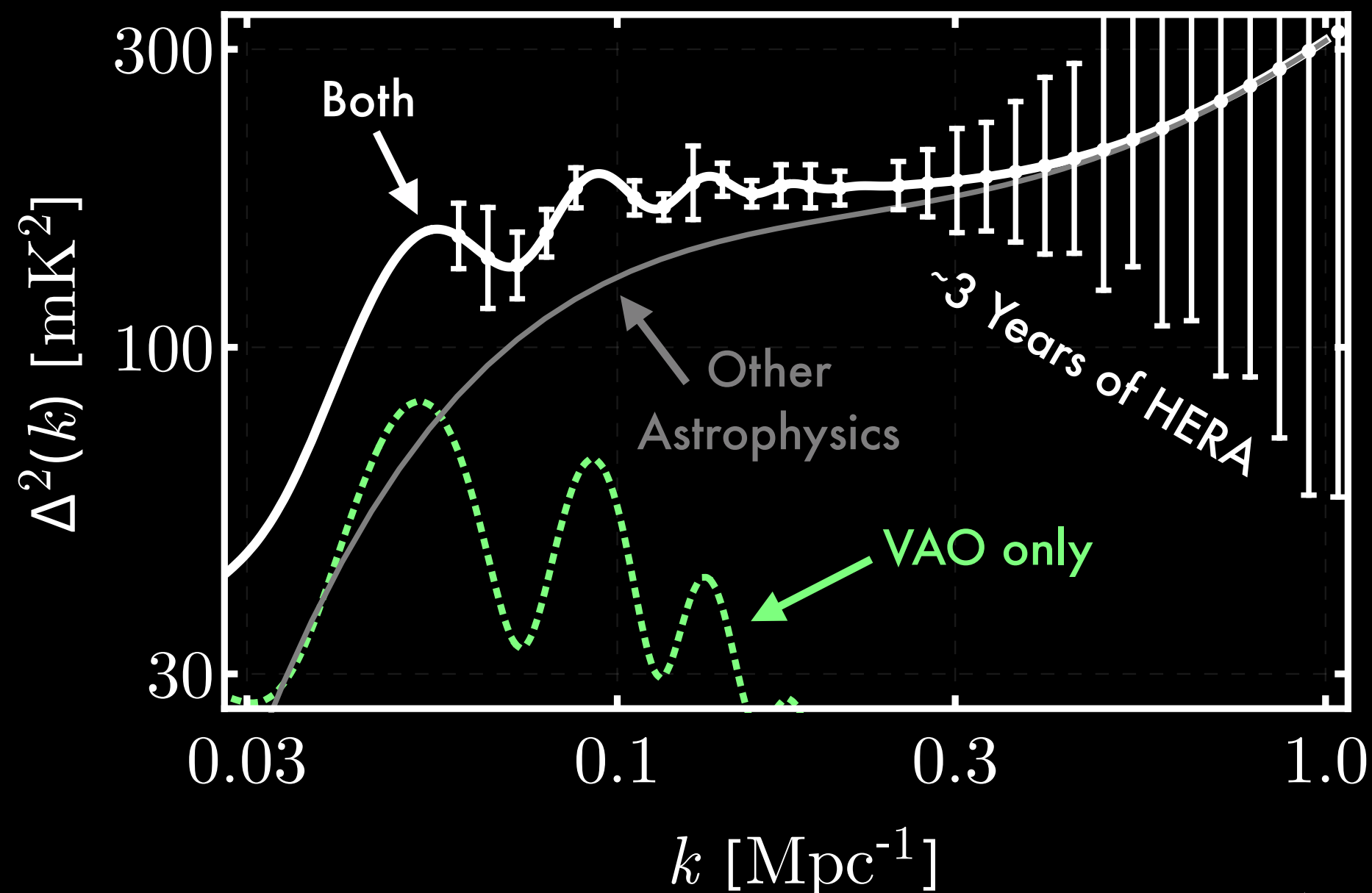
With our full data set, we'll eliminate τ as a CMB nuisance parameter, improving A_s errors by a factor of 4.



Combined with CMB-S4, we can likely detect a non-zero Σm_ν , distinguishing between neutrino hierarchies.

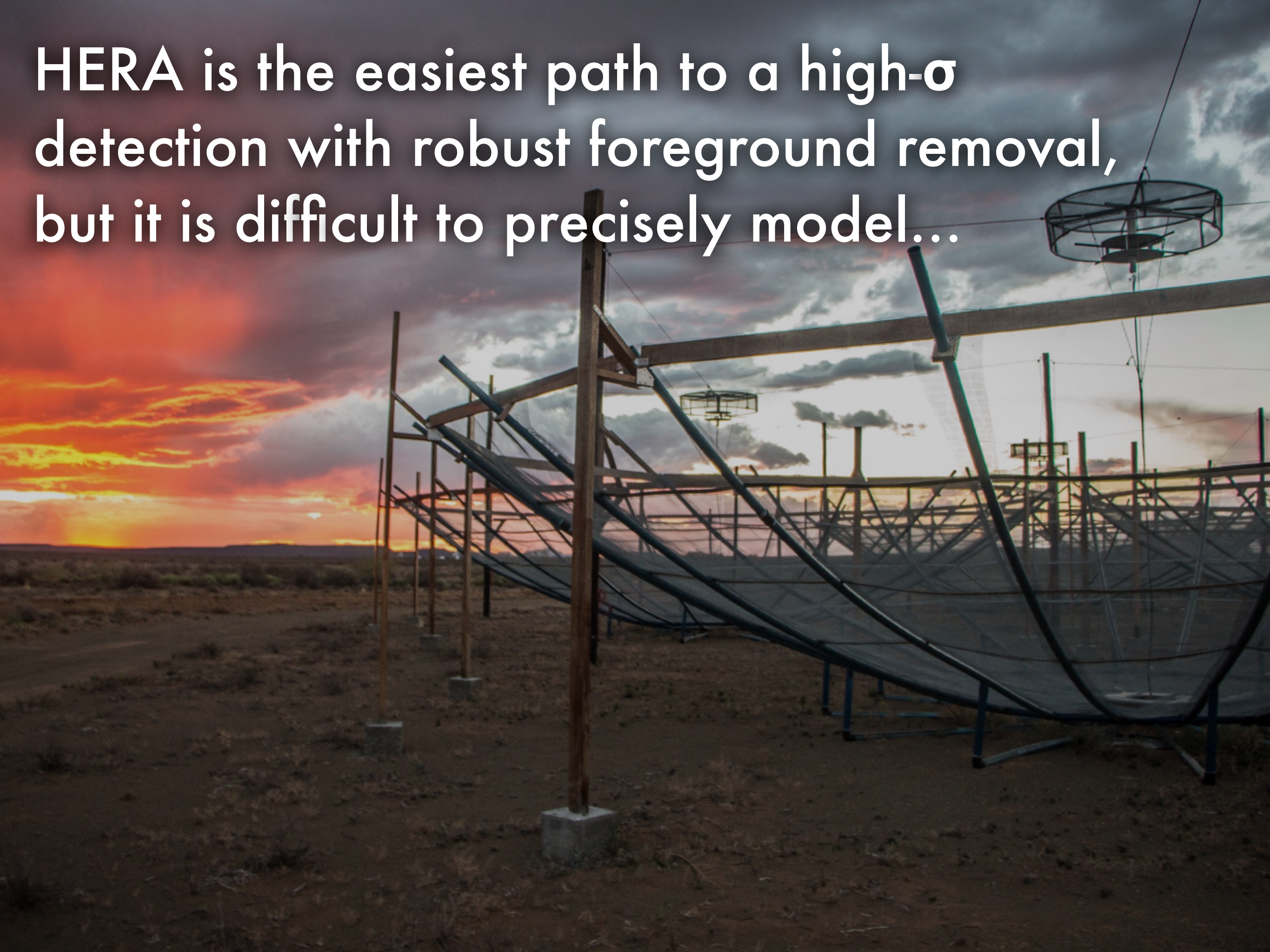


With a few years of observing, we may
detect velocity acoustic oscillations,
providing a new standard ruler at $z \approx 16$.



What comes after HERA?

HERA is the easiest path to a high- σ detection with robust foreground removal, but it is difficult to precisely model...



HERA is the easiest path to a high- σ detection with robust foreground removal, but it is difficult to precisely model...

...a *bigger* array of *smaller, simpler* antennas with *less coupling* and larger fields of view is likely the way forward.

There's a problem with how we measure visibilities.

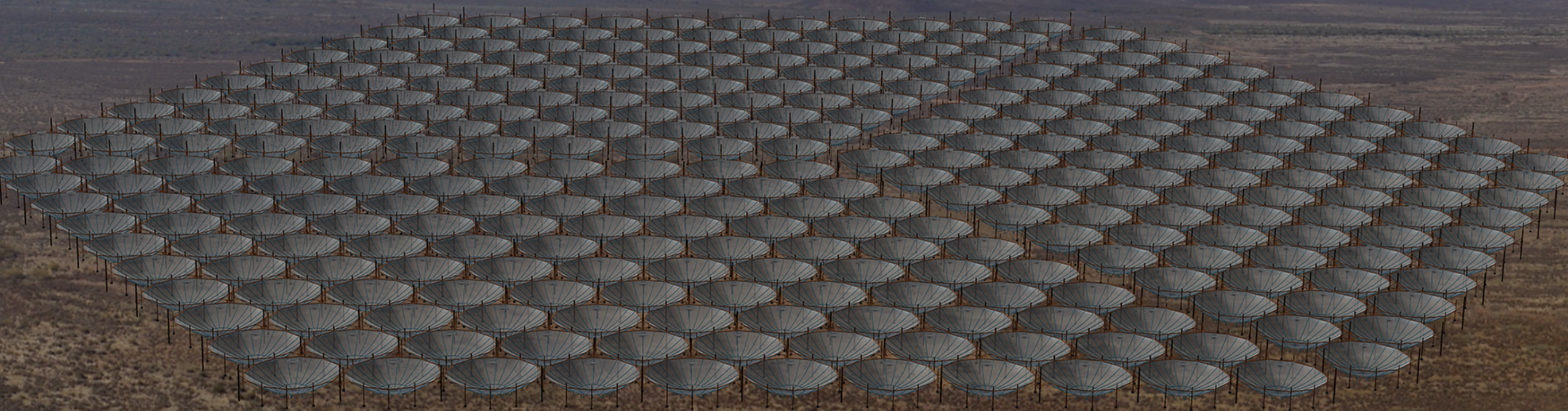
$$V_{ij}(\nu) = \int B_{ij}(\hat{\mathbf{r}}, \nu) I(\hat{\mathbf{r}}, \nu) \exp \left[-2\pi i \frac{\nu}{c} \mathbf{b}_{ij} \cdot \hat{\mathbf{r}} \right] d\Omega$$



There's a problem with how we measure visibilities.

$$V_{ij}(\nu) = \int B_{ij}(\hat{\mathbf{r}}, \nu) I(\hat{\mathbf{r}}, \nu) \exp \left[-2\pi i \frac{\nu}{c} \mathbf{b}_{ij} \cdot \hat{\mathbf{r}} \right] d\Omega$$

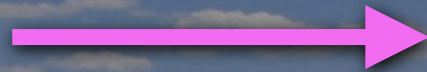
Measure antenna
voltages $v_i(t)$.



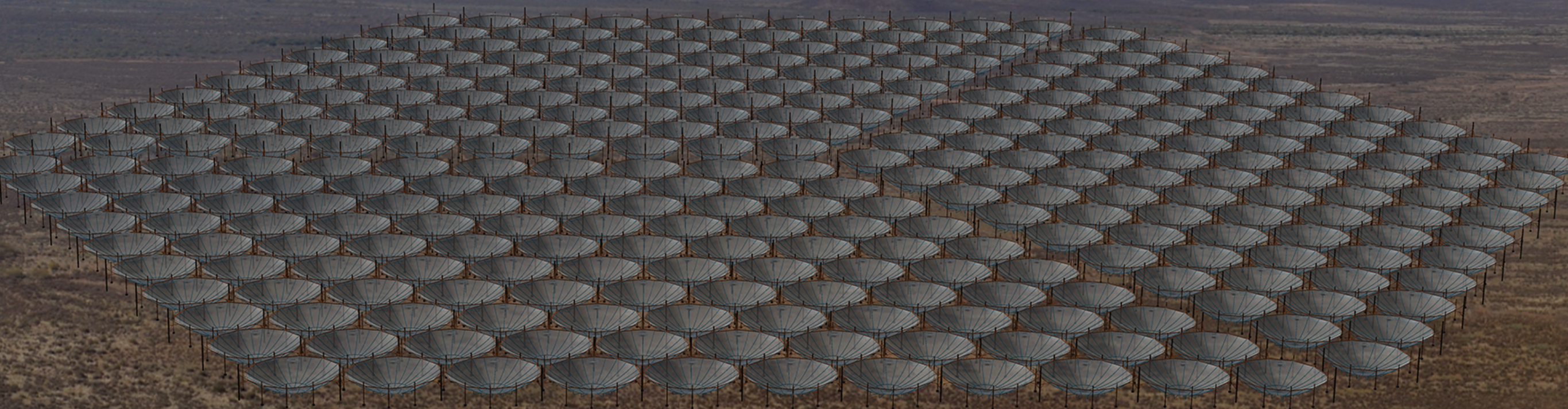
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Measure antenna
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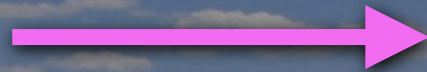
Fourier transform
to frequency: $\tilde{v}_i(\nu)$



There's a problem with how we measure visibilities.

$$V_{ij}(\nu) = \int B_{ij}(\hat{\mathbf{r}}, \nu) I(\hat{\mathbf{r}}, \nu) \exp \left[-2\pi i \frac{\nu}{c} \mathbf{b}_{ij} \cdot \hat{\mathbf{r}} \right] d\Omega$$

Measure antenna
voltages $v_i(t)$.

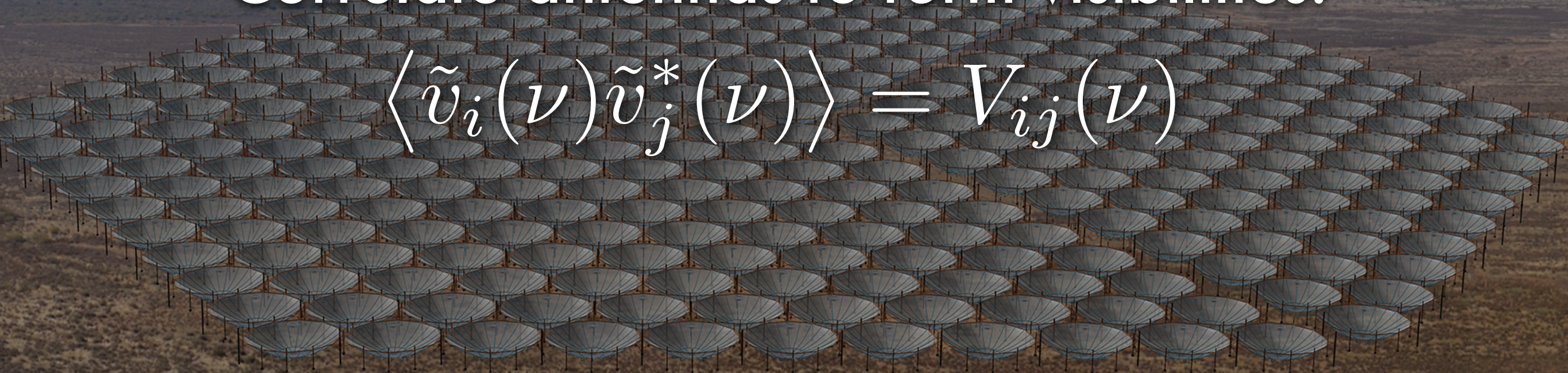


Fourier transform
to frequency: $\tilde{v}_i(\nu)$



Correlate antennas to form visibilities:

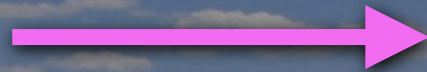
$$\langle \tilde{v}_i(\nu) \tilde{v}_j^*(\nu) \rangle = V_{ij}(\nu)$$



There's a problem with how we measure visibilities.

$$V_{ij}(\nu) = \int B_{ij}(\hat{\mathbf{r}}, \nu) I(\hat{\mathbf{r}}, \nu) \exp \left[-2\pi i \frac{\nu}{c} \mathbf{b}_{ij} \cdot \hat{\mathbf{r}} \right] d\Omega$$

Measure antenna
voltages $v_i(t)$.



Fourier transform
to frequency: $\tilde{v}_i(\nu)$

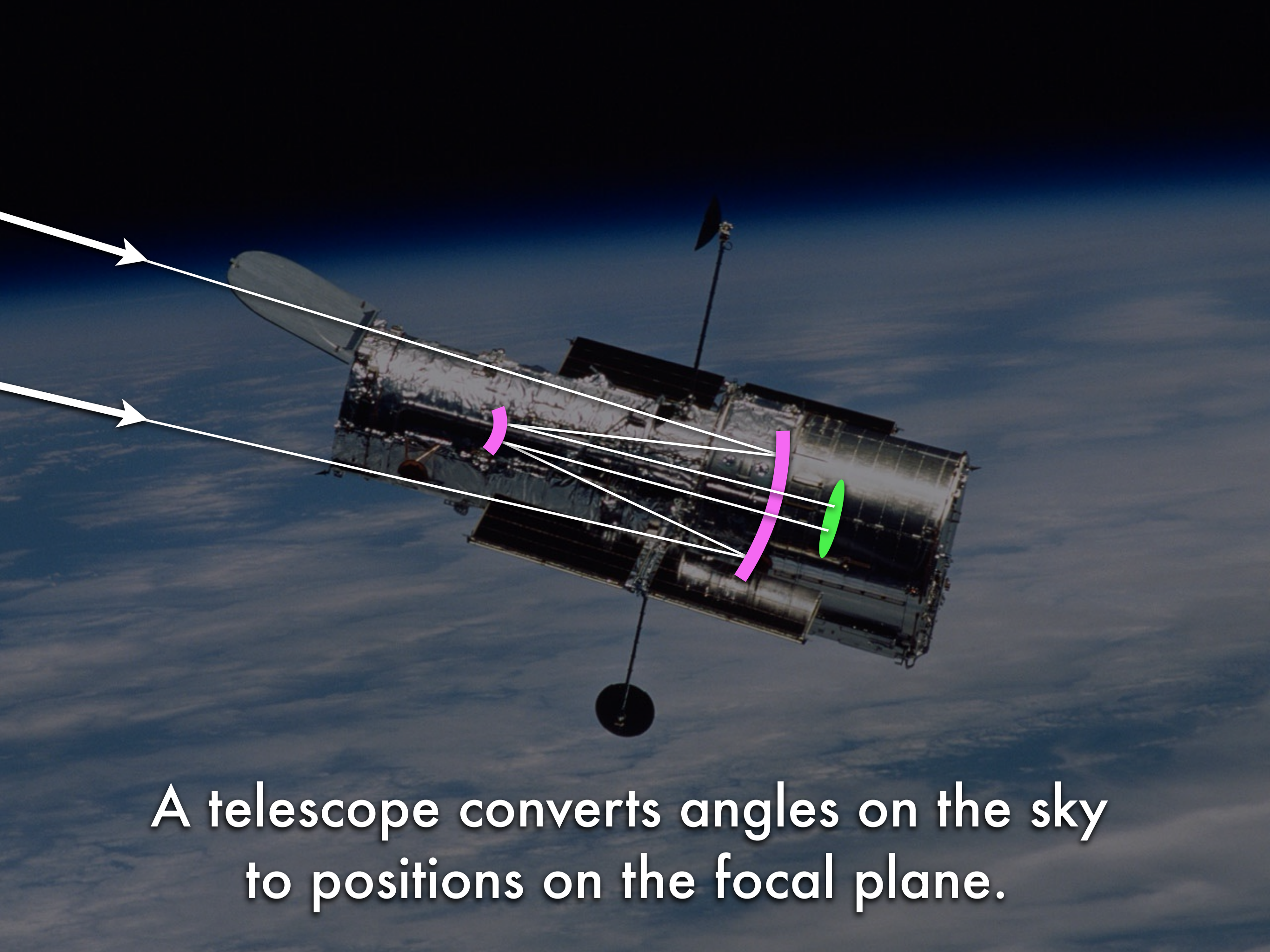


Correlate antennas to form visibilities:

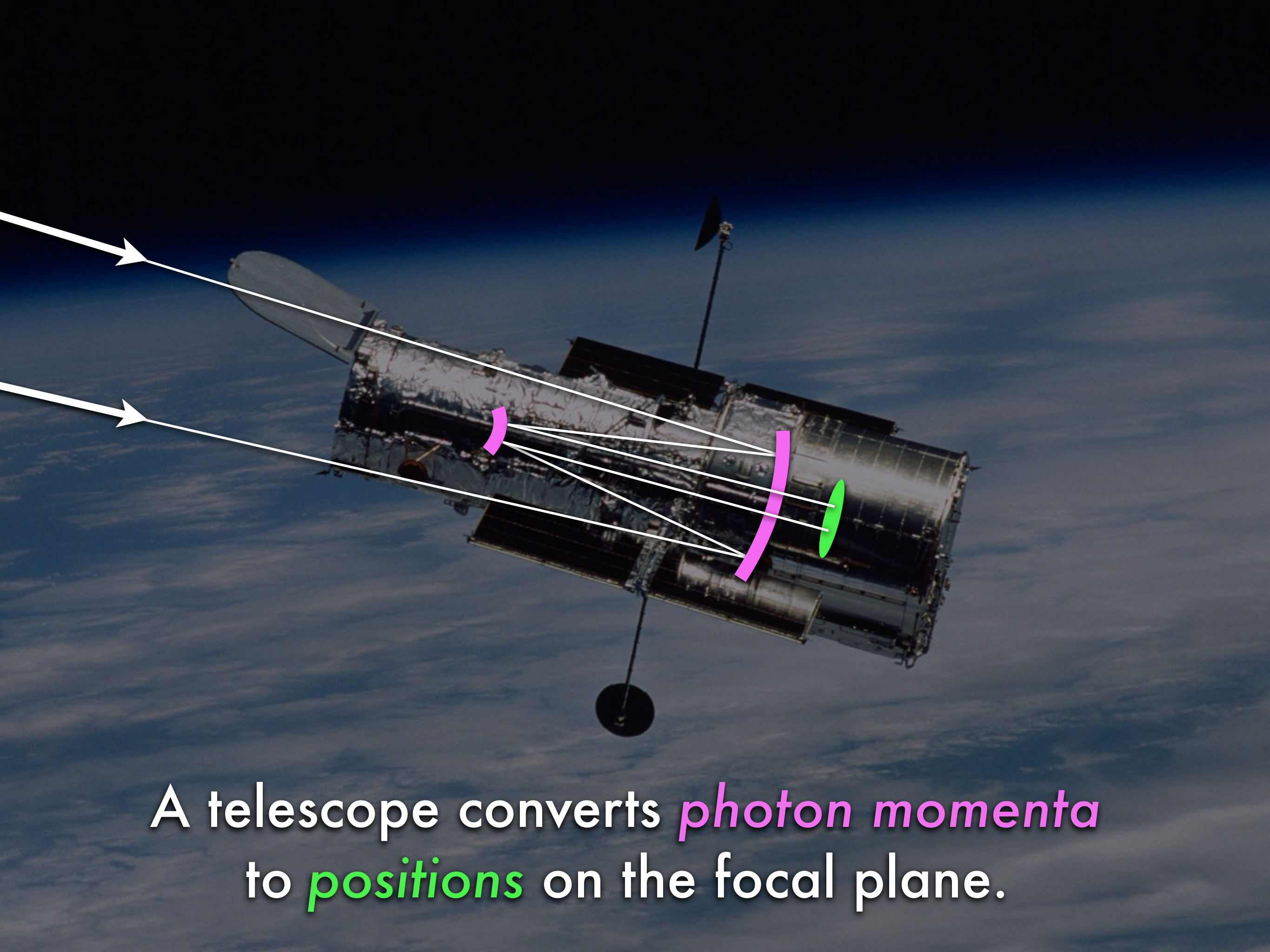
$$\langle \tilde{v}_i(\nu) \tilde{v}_j^*(\nu) \rangle = V_{ij}(\nu)$$

This scales like $O(N^2)$!

All telescopes are
Fourier transformers.



A telescope converts angles on the sky
to positions on the focal plane.

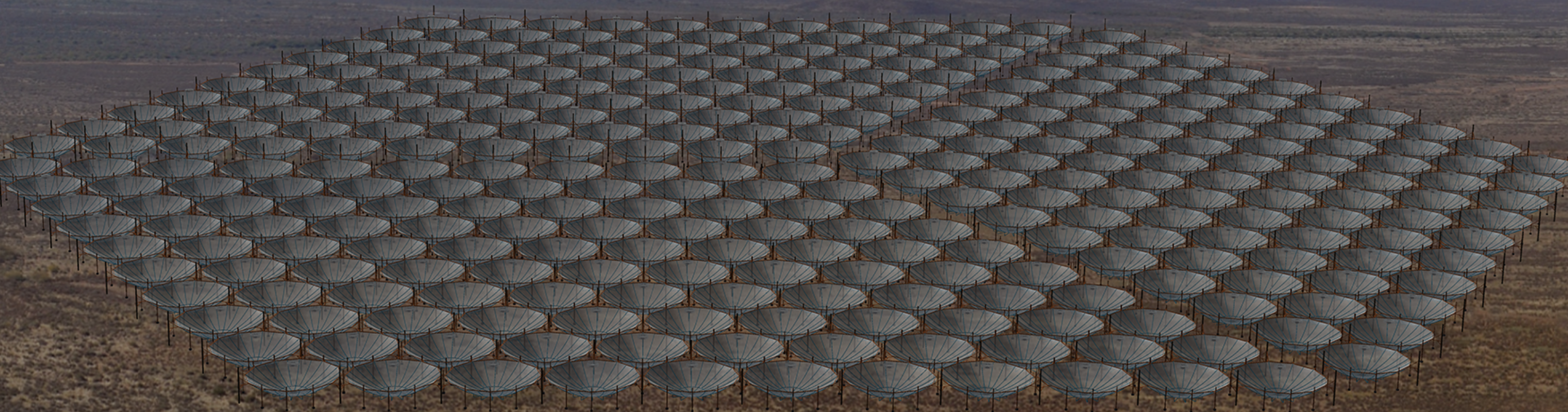


A telescope converts *photon momenta*
to *positions* on the focal plane.

$$V_{ij}(\nu) = \int B_{ij}(\hat{\mathbf{r}}, \nu) I(\hat{\mathbf{r}}, \nu) \exp \left[-2\pi i \frac{\nu}{c} \mathbf{b}_{ij} \cdot \hat{\mathbf{r}} \right] d\Omega$$

can be rewritten suggestively as...

$$\langle \tilde{v}_i(k) \tilde{v}_j^* \rangle = \int B(\mathbf{k}) I(\mathbf{k}) \exp [i\mathbf{k} \cdot (\mathbf{x}_i - \mathbf{x}_j)] d\Omega$$



Tegmark & Zaldarriaga (2009)

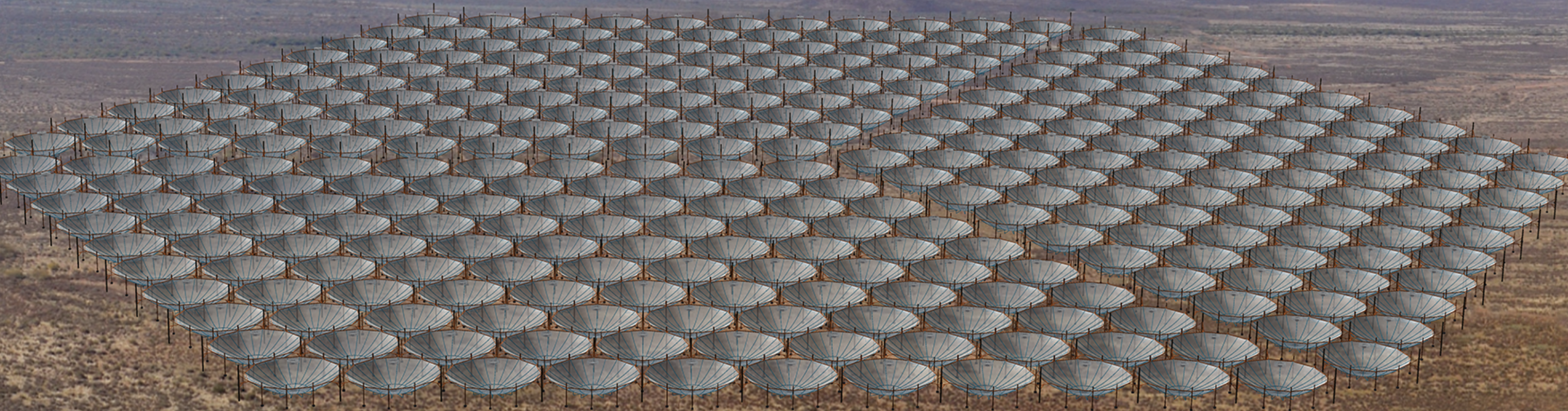
$$V_{ij}(\nu) = \int B_{ij}(\hat{\mathbf{r}}, \nu) I(\hat{\mathbf{r}}, \nu) \exp \left[-2\pi i \frac{\nu}{c} \mathbf{b}_{ij} \cdot \hat{\mathbf{r}} \right] d\Omega$$

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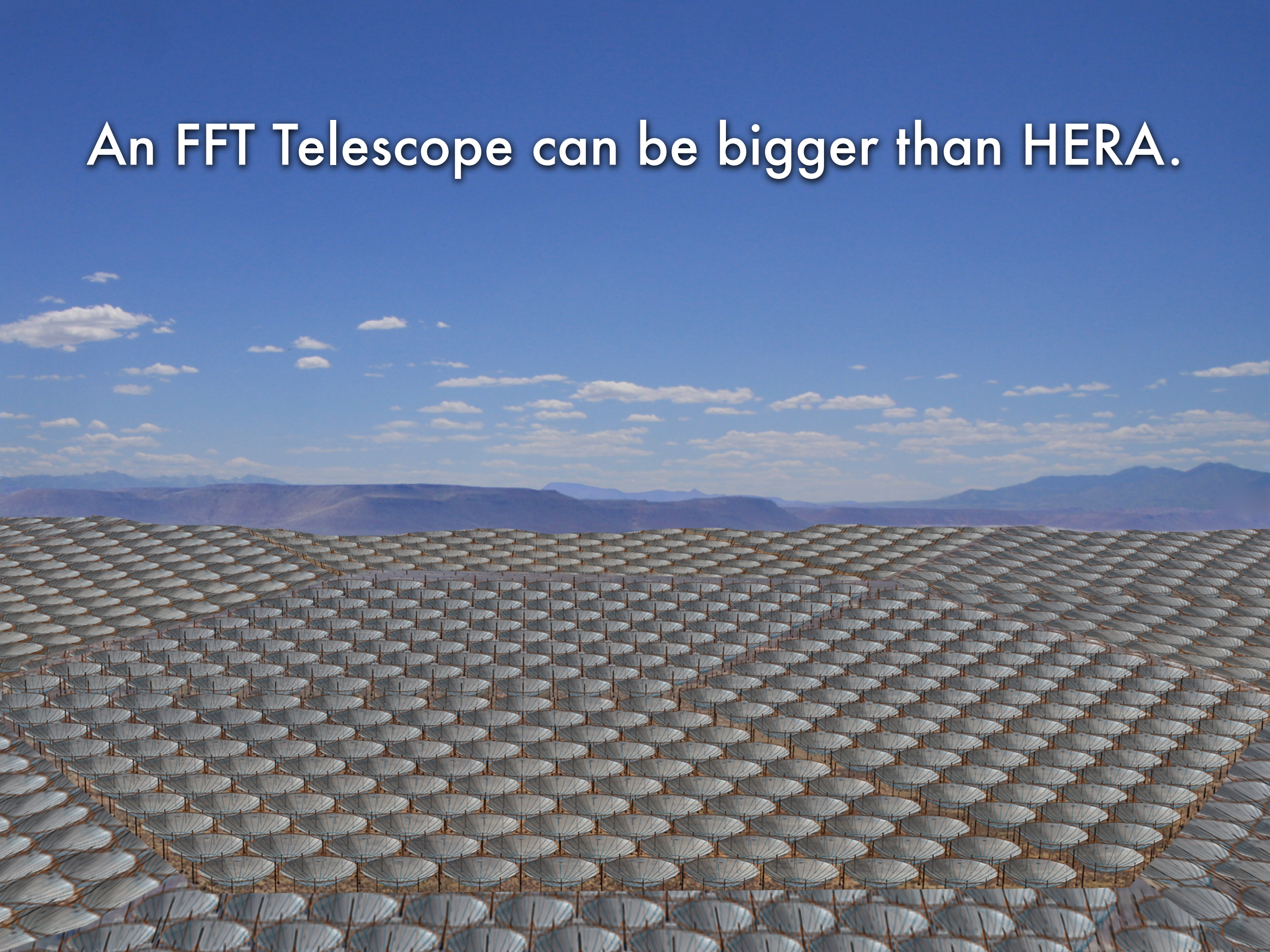
$$\langle \tilde{v}_i(k) \tilde{v}_j^* \rangle = \int B(\mathbf{k}) I(\mathbf{k}) \exp [i\mathbf{k} \cdot (\mathbf{x}_i - \mathbf{x}_j)] d\Omega$$

If antenna positions \mathbf{x}_i are on a regular grid,
we can directly sample the electric field, FFT,
and square to get beam-weighted maps...
effectively correlating in $O(N \log N)$!

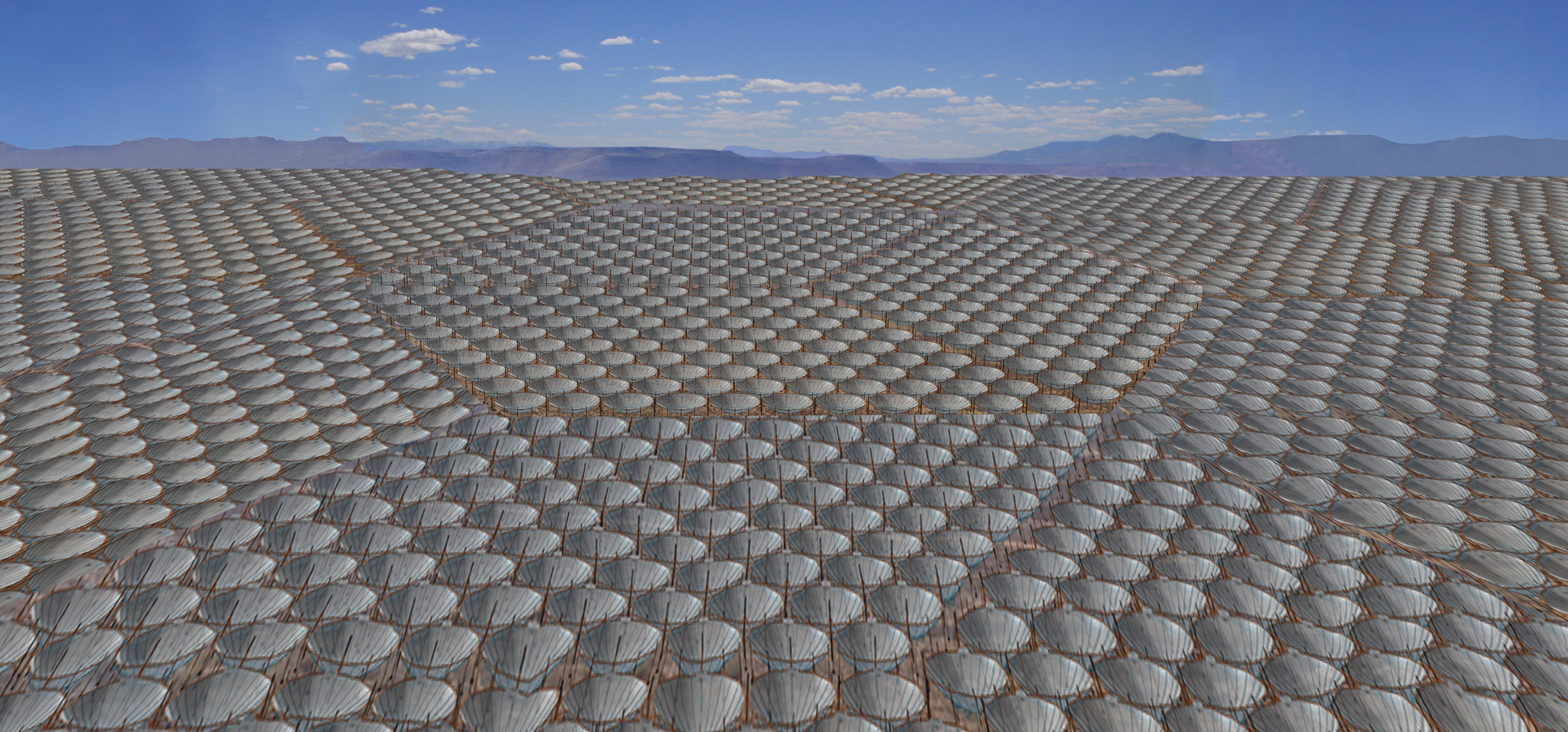
An FFT Telescope can be bigger than HERA.



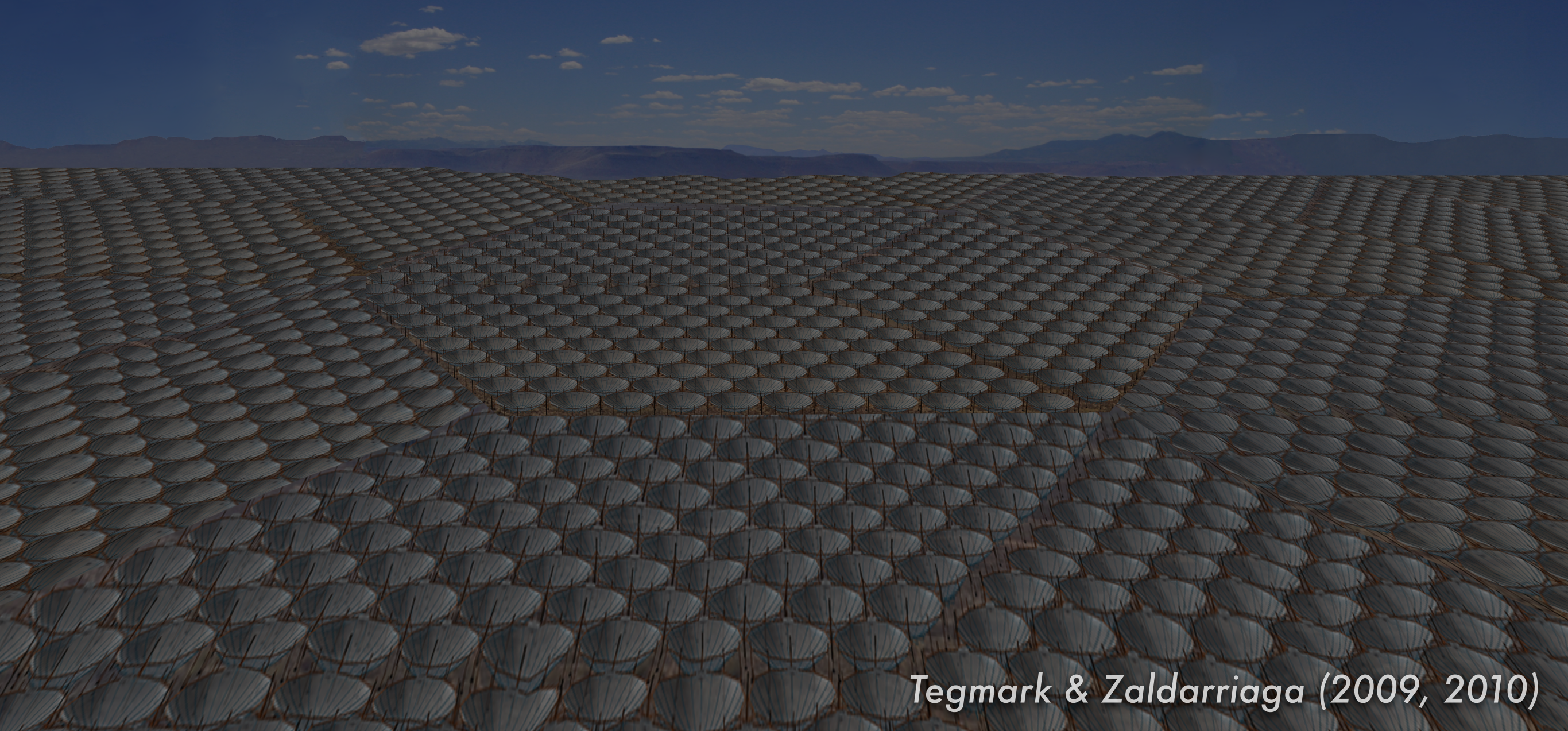
An FFT Telescope can be bigger than HERA.



An FFT Telescope can be bigger than HERA.
Much, much bigger.



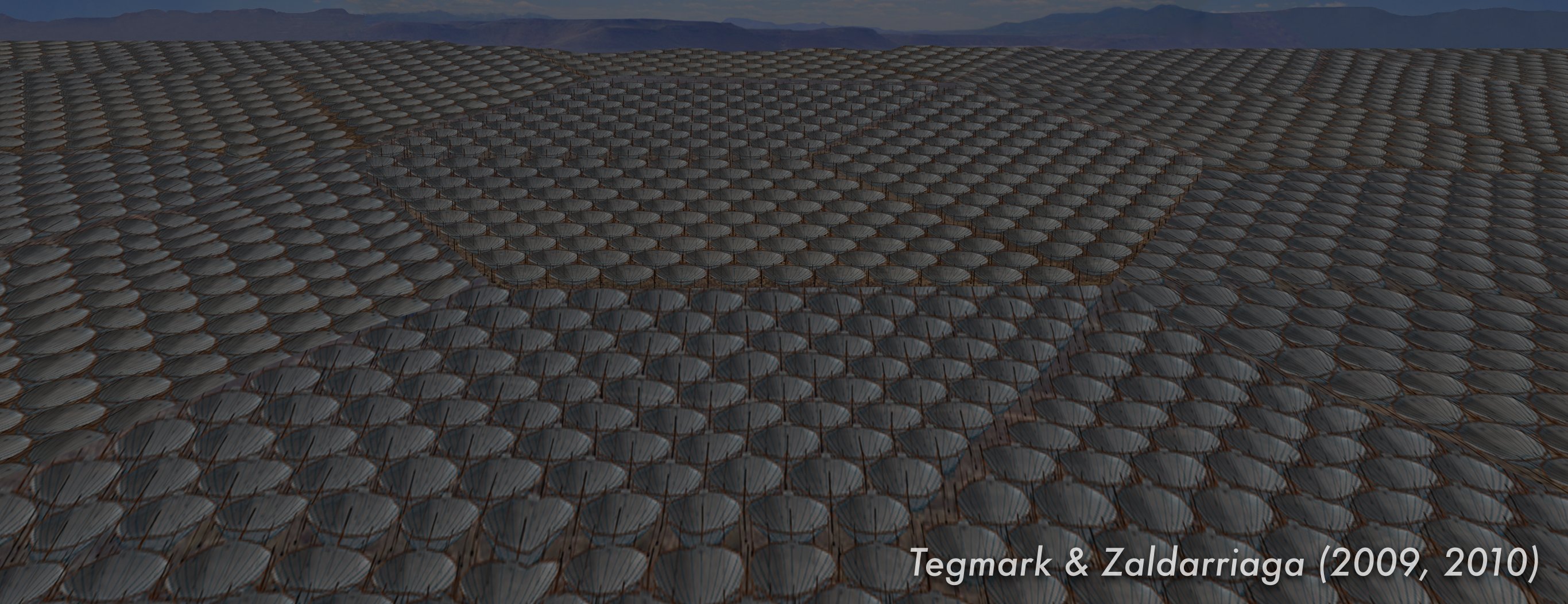
An FFT Telescope needs to be...



Tegmark & Zaldarriaga (2009, 2010)

An FFT Telescope needs to be...

- Co-planar.



Tegmark & Zaldarriaga (2009, 2010)

An FFT Telescope needs to be...

- Co-planar.
- Made up of identical antenna elements with identical beams (ideally with low mutual coupling).

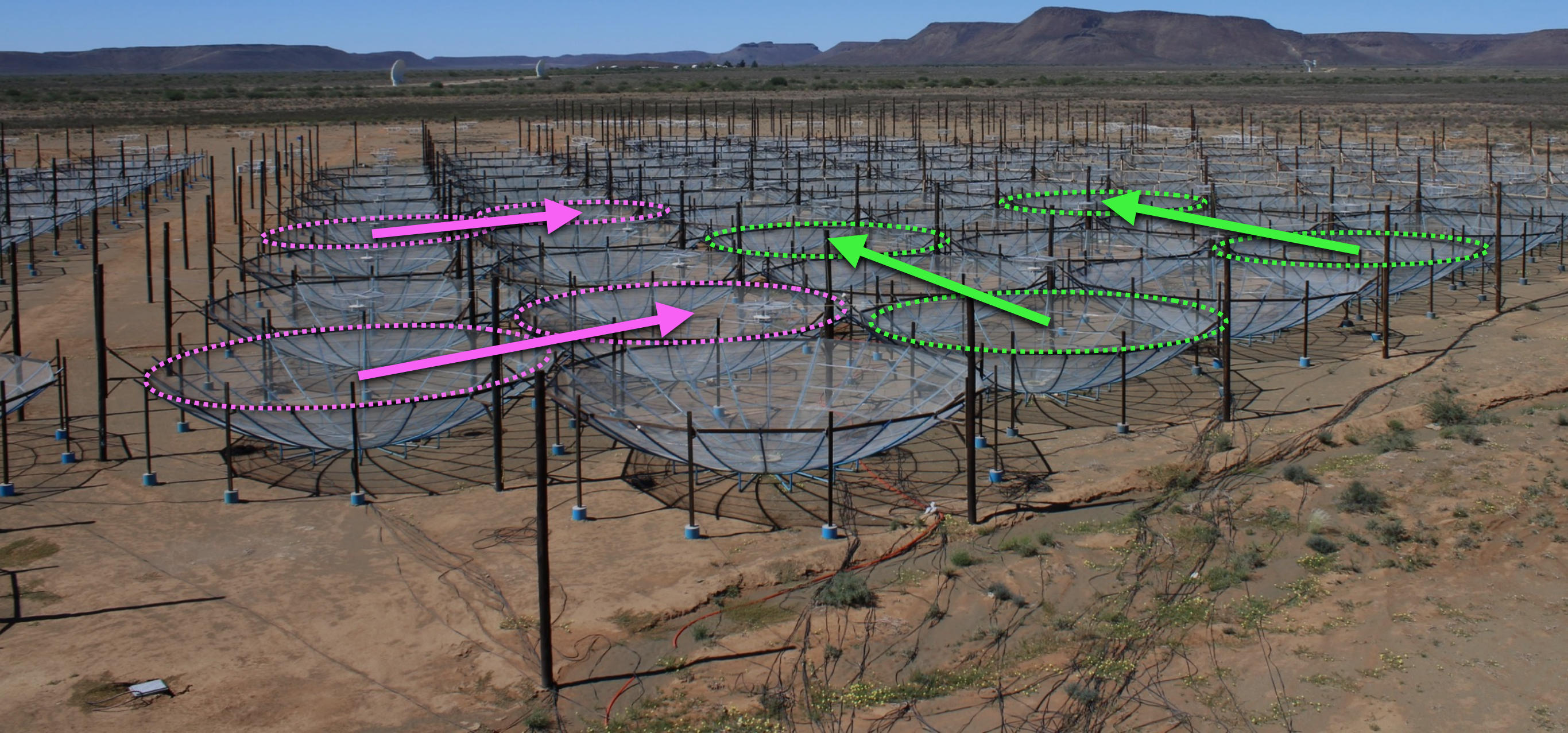
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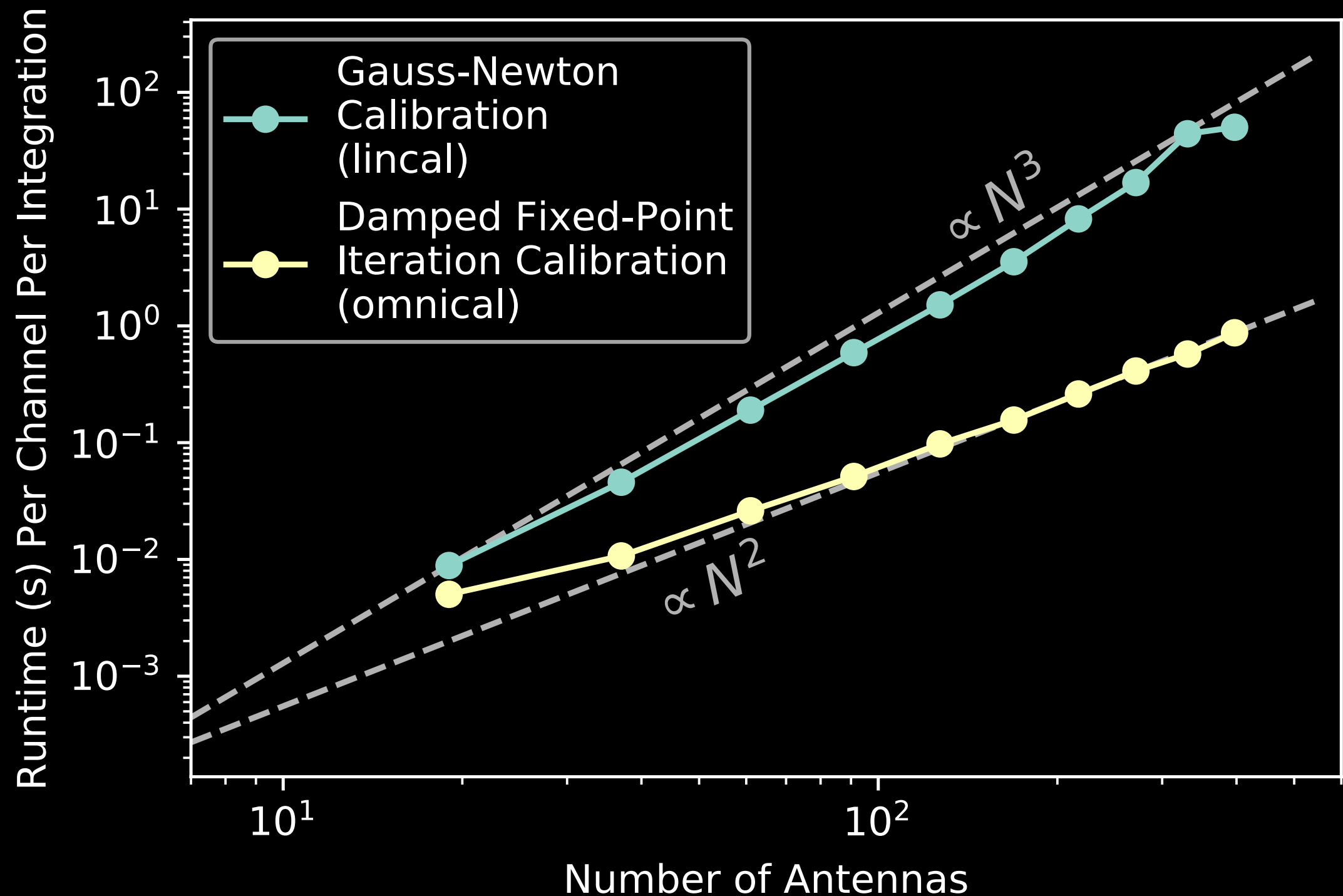
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- Calibrated in real time.

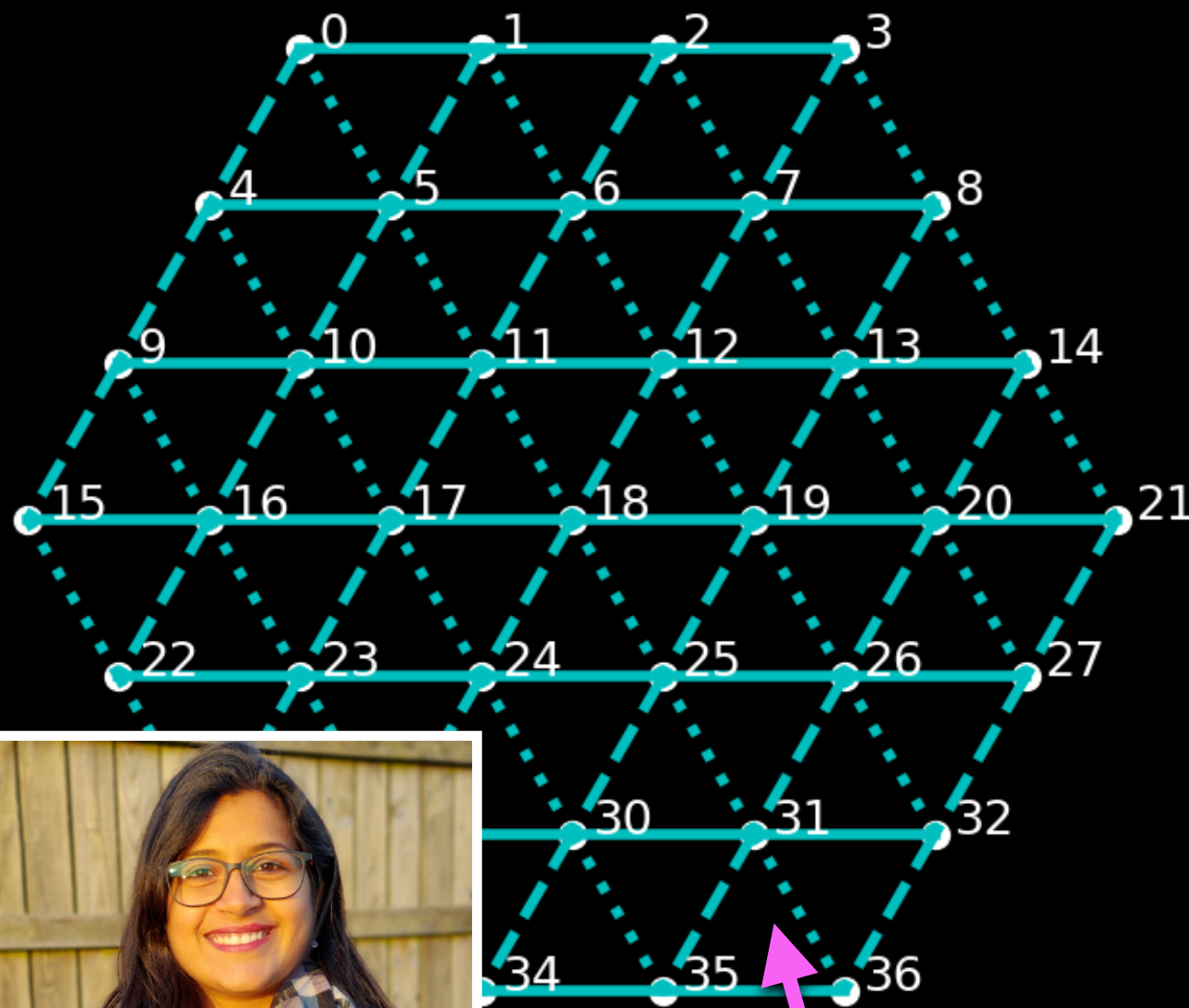
Real-time redundant-baseline calibration
of regular arrays is precisely what we're
learning to do with HERA!



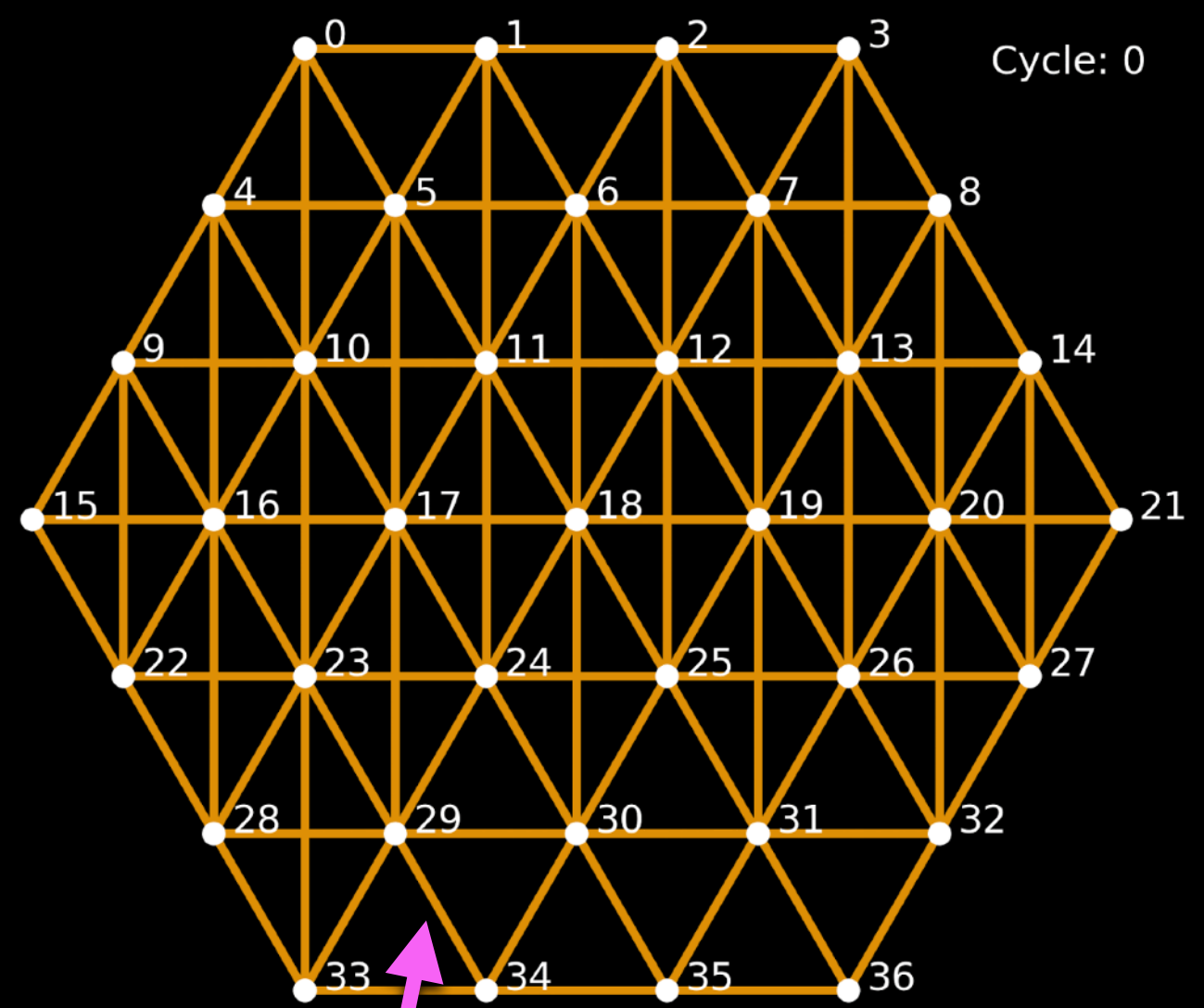
We showed how to speed up redundant calibration from $O(N^3)$ to $O(N^2)$.



And how to use a subset of the data to reduce calibration from $O(N^2)$ to $O(N \log N)$.



Subset-Redundant
Calibration

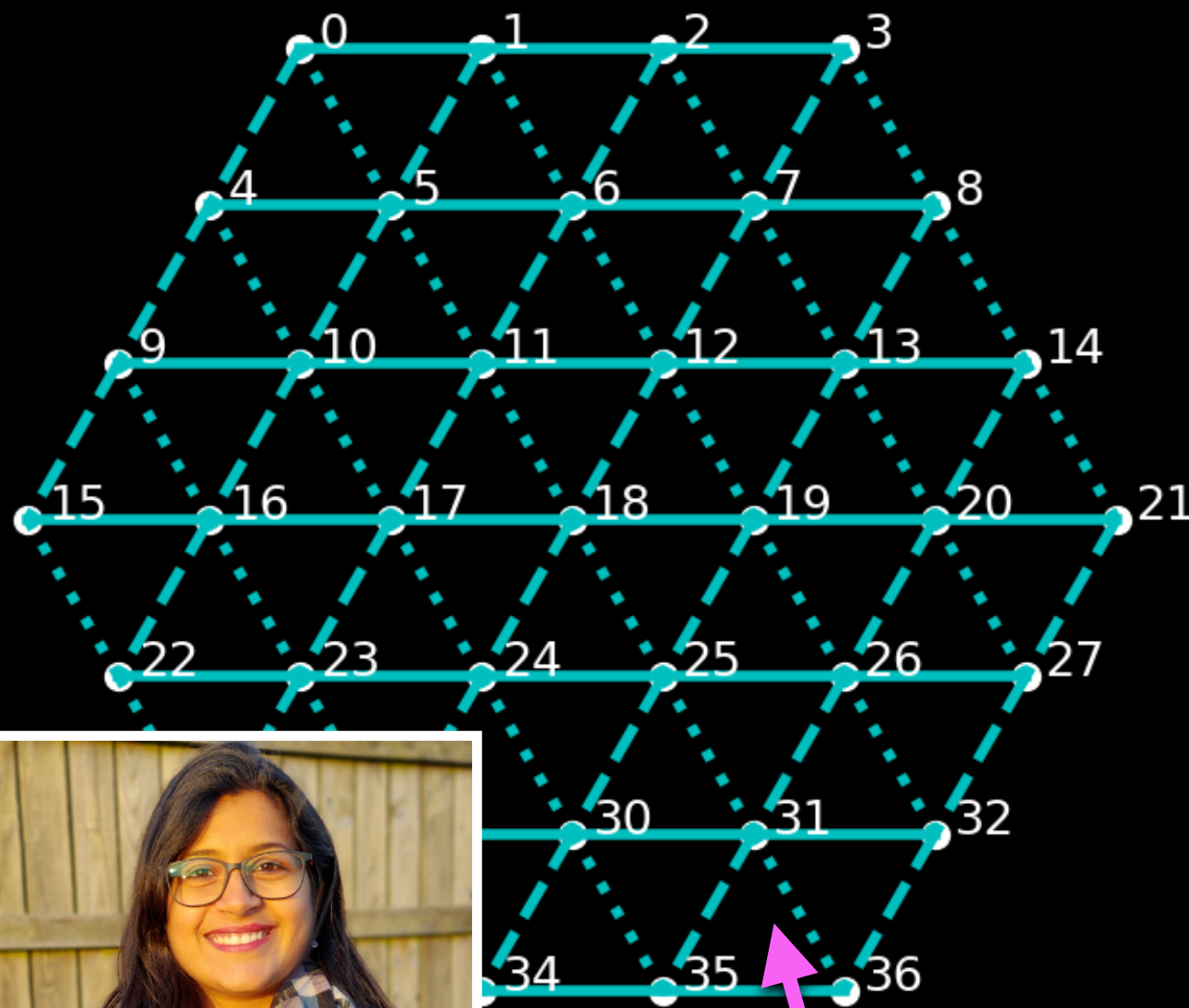


Low-Cadence
Calibration

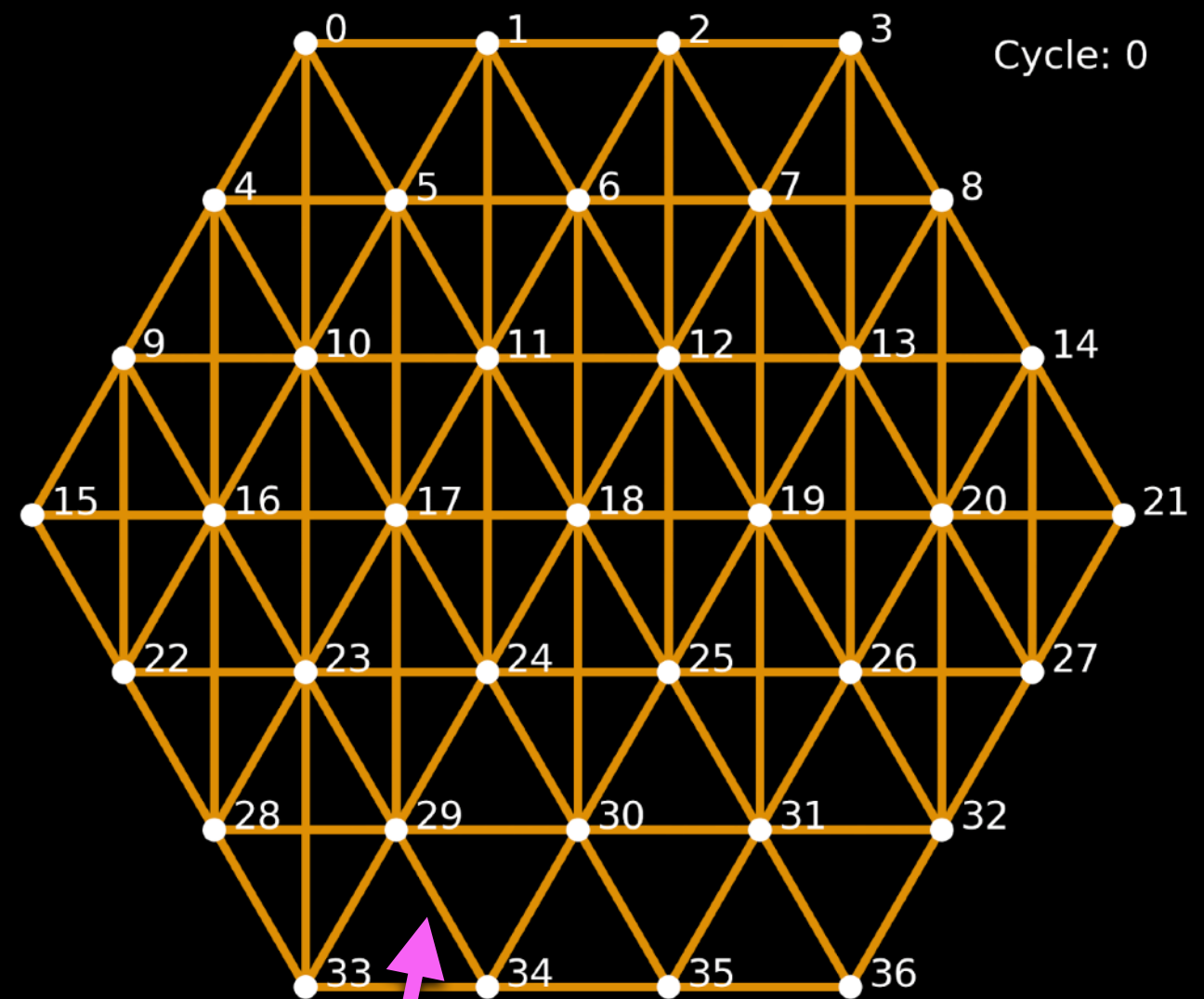
vs.

Deepthi Gorthi
Former UCB
Grad Student

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Subset-Redundant
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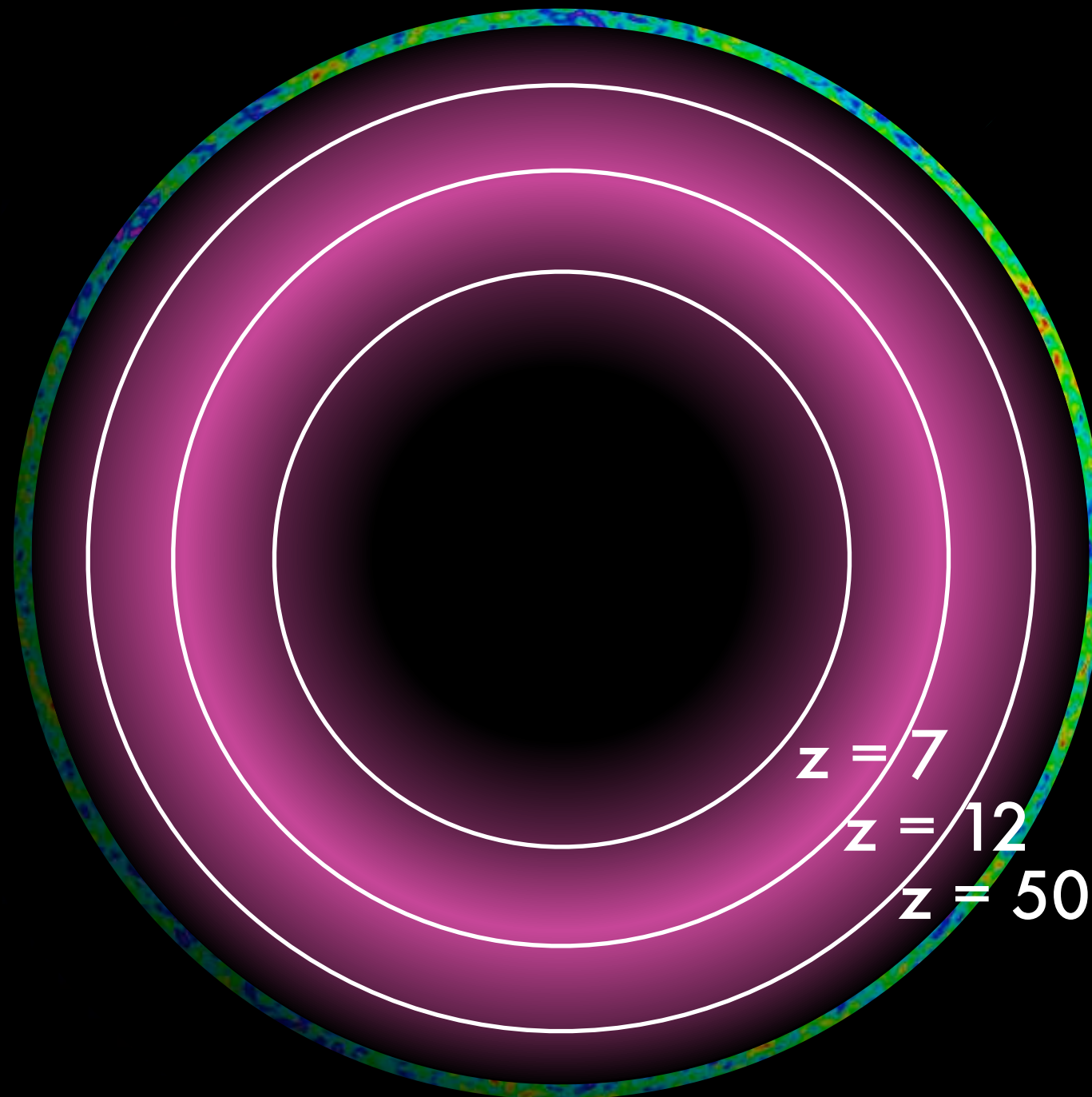


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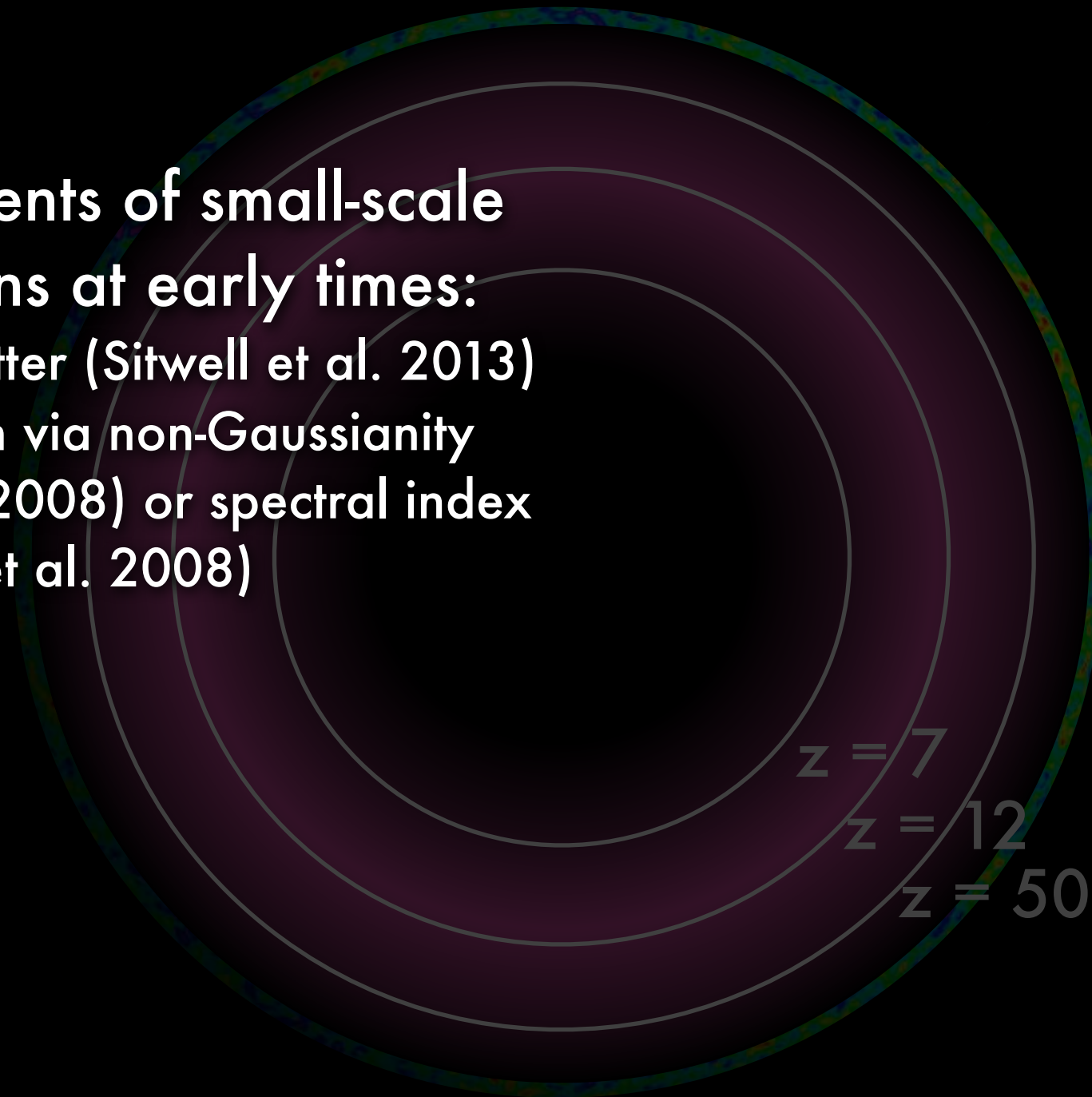
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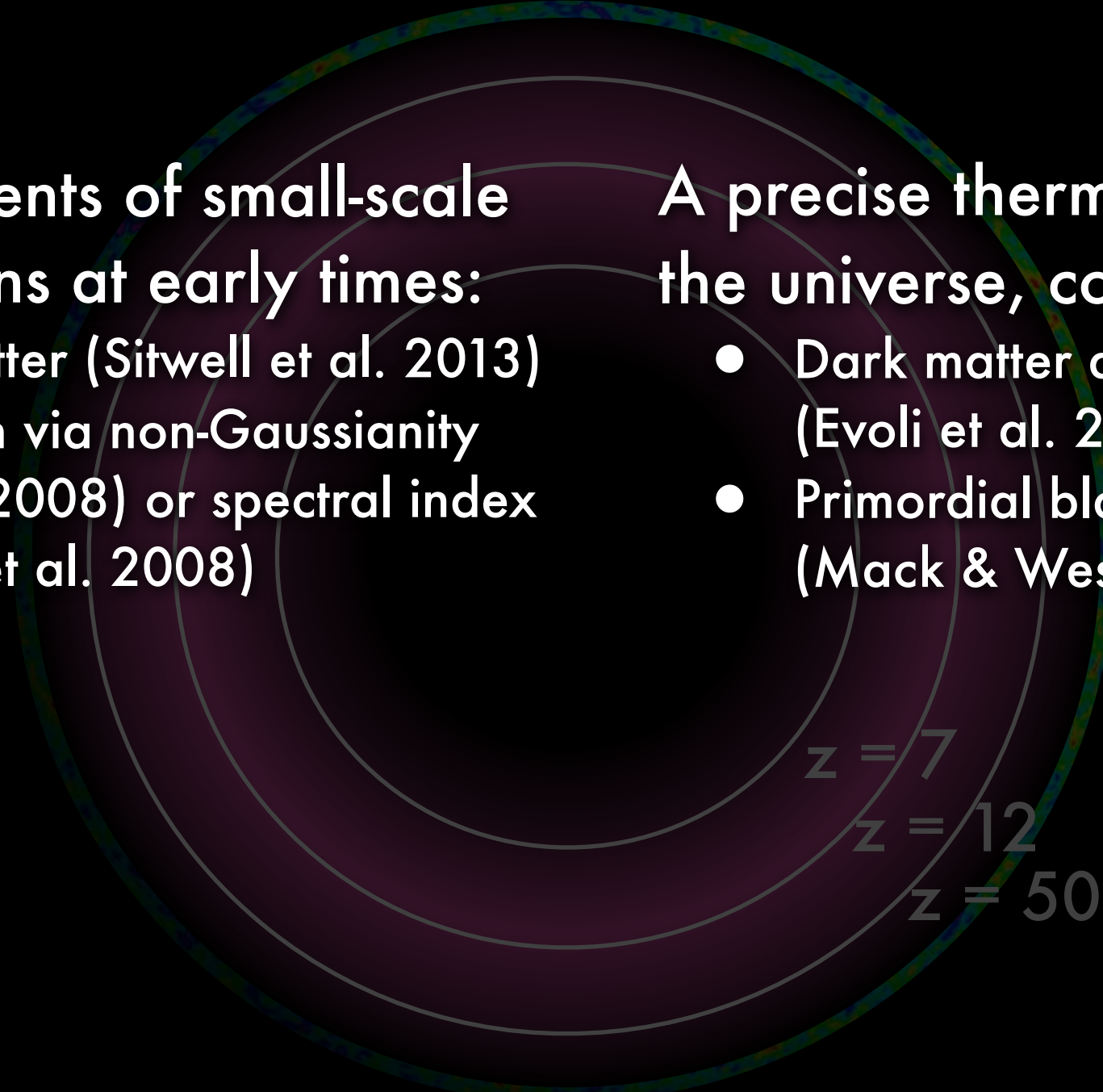
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$z = 7$
 $z = 12$
 $z = 50$

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Unprecedented constraints on the standard model of cosmology:

- Orders of magnitude better than Planck, e.g. $\Delta\Omega_k \approx .0002$ and $\Delta\Sigma v \approx 7$ meV (Mao et al. 2008)

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Photo: Dara Storer

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- HERA is elucidating key systematics as we scale to instruments that can realize the potential of 21 cm cosmology