



UNIVERSITY OF  
**OXFORD**



# Galactic-Scale Tests of Fundamental Physics

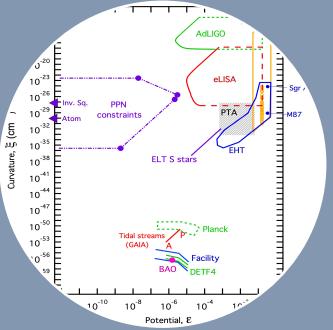
Cosmology/BCCP Seminar, Berkeley

12<sup>th</sup> October 2021

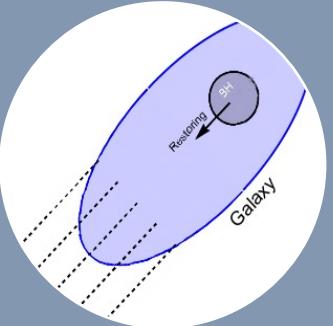
Deaglan Bartlett

Harry Desmond, Pedro Ferreira, Jens Jasche, Dexter Bergsdal, Guilhem Lavaux, Julien Devriendt, Adrienne Slyz

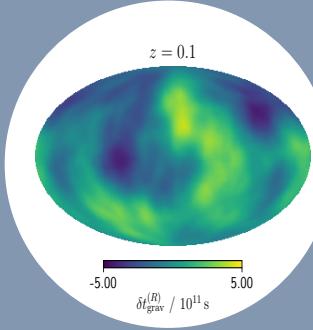
# Outline



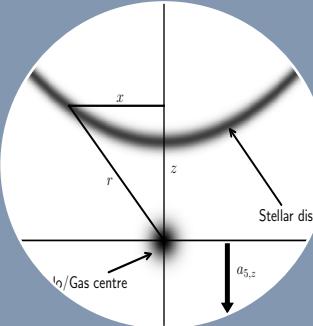
Why galaxies?



Black Holes and Galileons



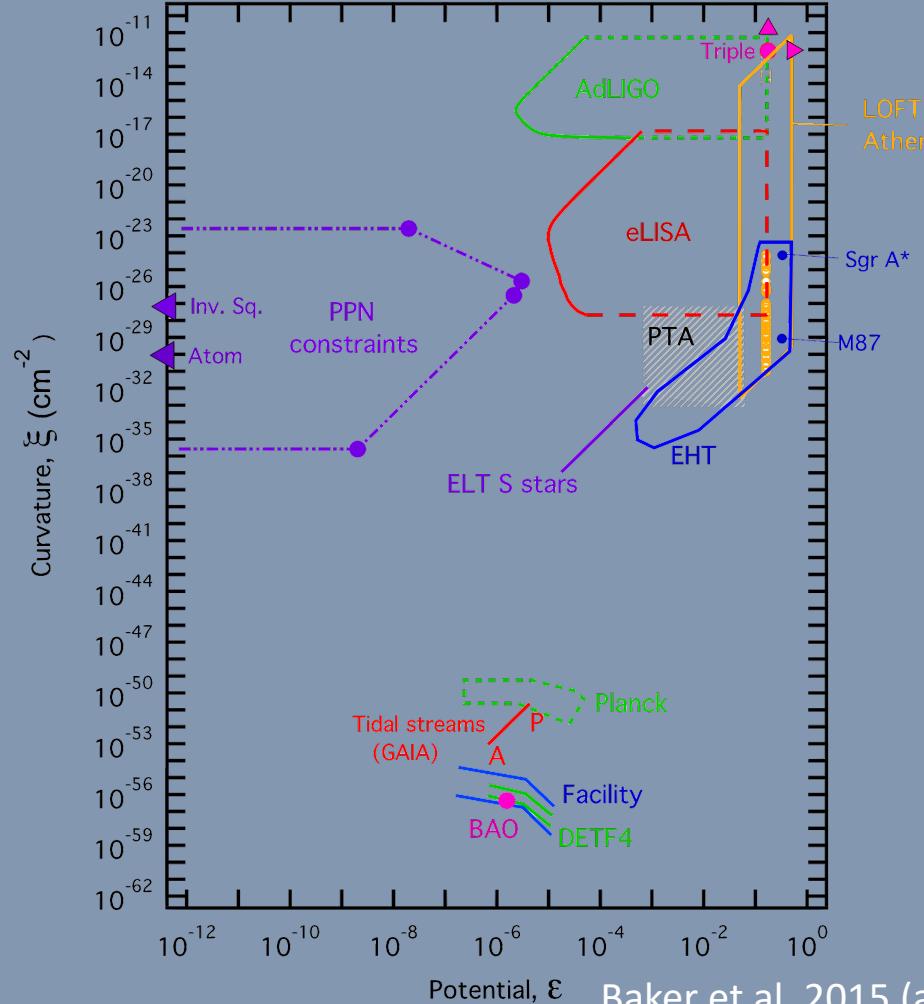
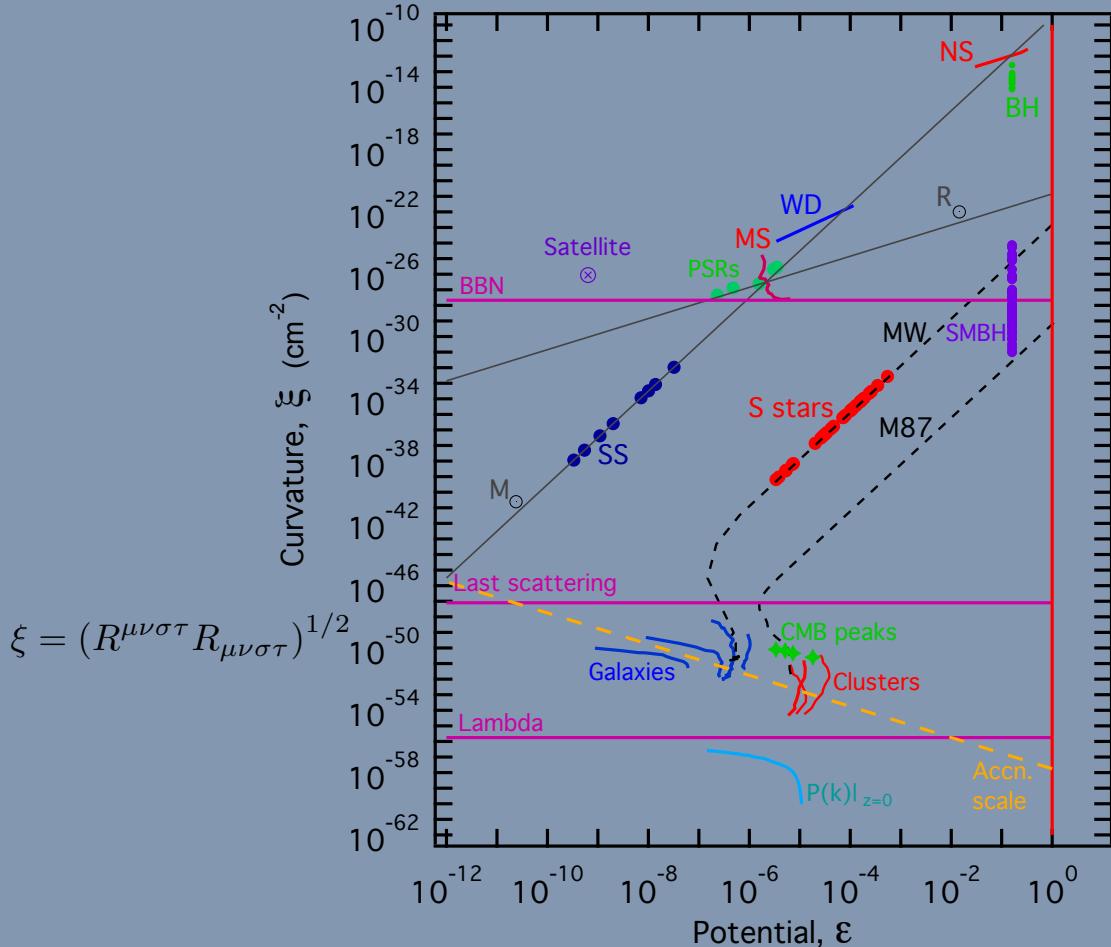
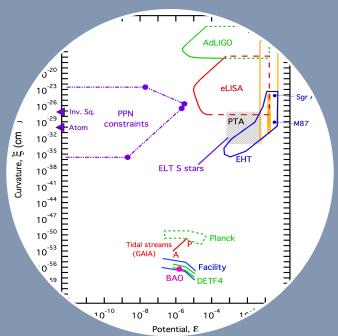
Photon dispersion relation



Modeling the noise

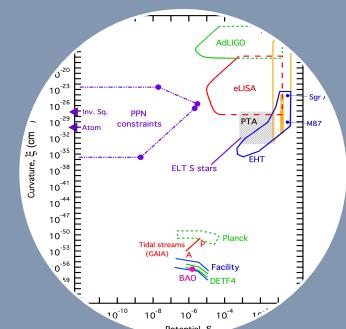
# Why galaxies?

# Astrophysics – An underexplored region of parameter space



Baker et al. 2015 (arXiv:1412.3455)

# Novel effects on astrophysical scales



- Fifth forces - screening
- Dark matter
- Lots of high quality data
- Usually assume a theory of gravity (GR or Newtonian) and then do astrophysics. Can we get a good enough handle on astrophysics to instead test the theory of gravity?

# Galaxy-by-galaxy forward model

$$\mathcal{L} = \prod_g \int \mathcal{L}_g(r_{\text{pred}}) \times \mathcal{L}_g(r_{\text{obs}} | r_{\text{pred}}) dr_{\text{pred}}$$

## “Predicted Signal”

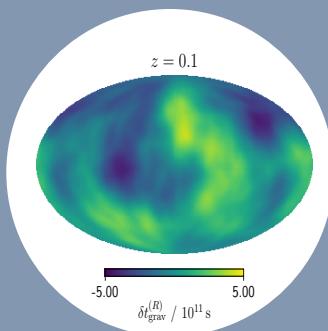
- “New physics” parameters
- Galactic properties
- Large scale structure (BORG)

## “Noise Model”

- “Noise” parameters
- Other astrophysical processes
- Observational effects

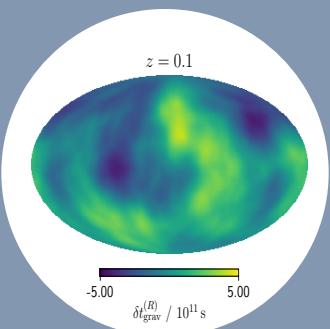
# Photon dispersion relation

# Quantum Gravity



Qualitative argument (Wheeler 1998):

- On timescales  $\Delta t \sim 1/E_P$ , expect quantum-gravitational fluctuations in spacetime of energy  $\Delta E \sim E_P$
- Spacetime would not look smooth on length scales  $\Delta x \sim 1/E_P$  - “foamy” structure
- Non-trivial refractive index for empty space? Would be important at high energies.



# Energy Dependent Arrival Times: Quantum Gravity

$$\vec{p}^2 = E^2 \left[ 1 + f \left( \frac{E}{E_{\text{QG}}} \right) \right] = E^2 \left[ 1 + \xi \frac{E}{E_{\text{QG}}} + \mathcal{O} \left( \frac{E^2}{E_{\text{QG}}^2} \right) \right]$$

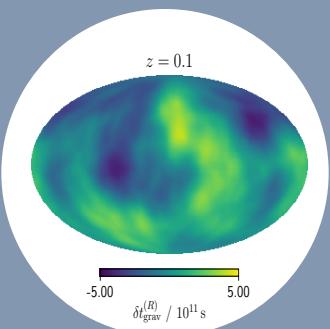
$$v = \frac{\partial E}{\partial p} \sim 1 - \xi \frac{E}{E_{\text{QG}}}$$

For two photons emitted simultaneously at redshift  $z$  and observed energies  $E_i$  and  $E_j$ ,

$$\Delta t_{ij}^{(\text{QG})} = \xi \frac{E_i - E_j}{E_{\text{QG}}} \int_0^z \frac{1 + z'}{H(z')} dz'$$

Could be observed at high energies.

Amelino-Camelia et al. 1998 (arXiv:9712103), Jacob & Piran 2008 (arXiv:0712.2170)

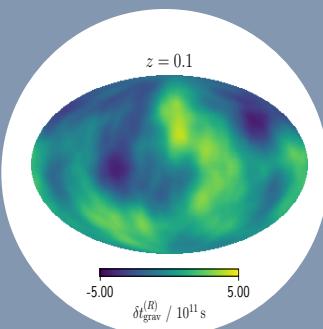


# Energy Dependent Arrival Times: Massive Photon

$$v = \sqrt{1 - \frac{m_\gamma^2}{E^2}} \approx 1 - \frac{1}{2} \frac{m_\gamma^2}{E^2}$$

$$\Delta t_{ij}^{(\text{MP})} = \frac{m_\gamma^2}{2} \left( \frac{1}{E_i^2} - \frac{1}{E_j^2} \right) \int_0^z \frac{1}{H(z') (1+z')^2} dz'$$

Would be important at low energies.

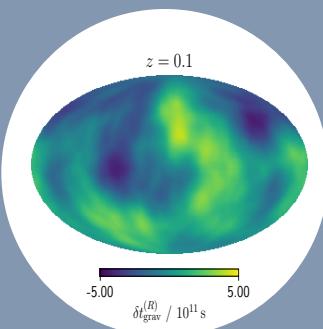


# GRBs are ideal for this

$$\Delta t_{ij}^{(\text{QG})} = \xi \frac{E_i - E_j}{E_{\text{QG}}} \int_0^z \frac{1 + z'}{H(z')} dz'$$

$$\Delta t_{ij}^{(\text{MP})} = \frac{m_\gamma^2}{2} \left( \frac{1}{E_i^2} - \frac{1}{E_j^2} \right) \int_0^z \frac{1}{H(z') (1 + z')^2} dz'$$

- Large  $z \Rightarrow$  large integral  $\Rightarrow$  tight constraints.
- Short duration and observed time delays  $\Rightarrow$  any novel effect must be small.
- QG (massive photon) best at high (low) energy.
- Radio frequencies - dispersion by electrons has same energy dependence as for a massive photon.



# Challenges to Overcome

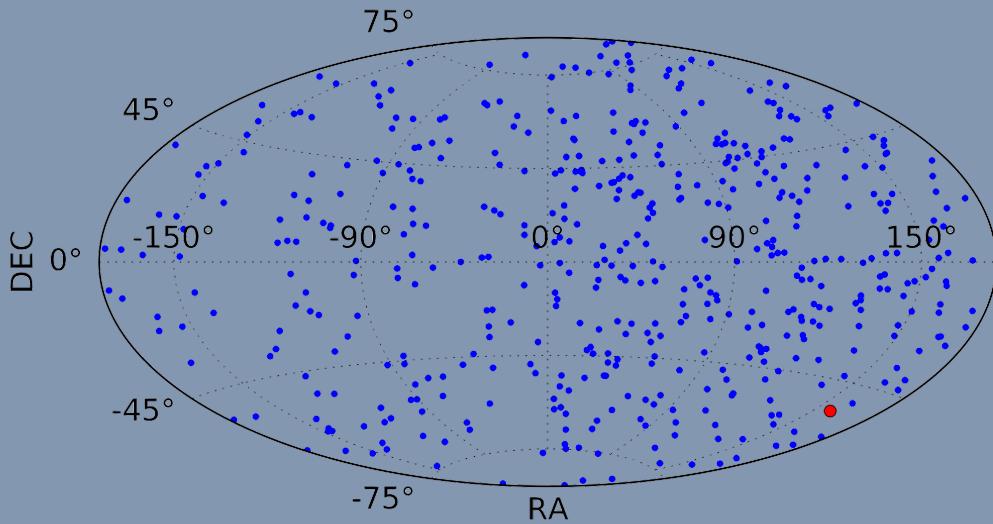
Some issues with previous constraints

- Do not propagate uncertainties in redshift of source.
- Simple (or no) model for other contributions to spectral lag.
- Only use a handful of GRBs.

How to overcome this

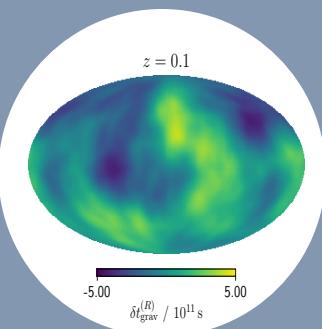
- Forward model time delays with Monte Carlo procedure.
- Test a variety of models - compare performance and sensitivity of constraint.
- Use a larger sample.

# GRB Data

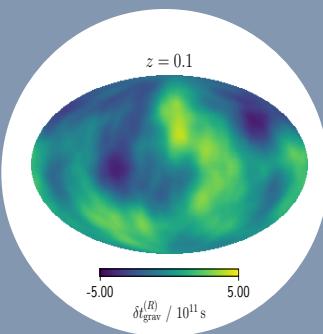


- 448 GRBs (BATSE)
- Time delay data for 4 energy channels (25 – 325 keV)
- Pseudo-redshifts (with uncertainty) from spectral peak energy-peak luminosity relation:

$$\log L \propto \log (E_p (1 + z))$$



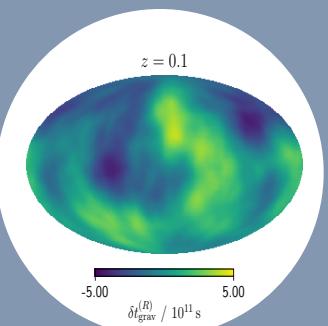
Yu, Xi & Wang 2018 (arXiv:1708.02396)



# Forward Model the Time Delays

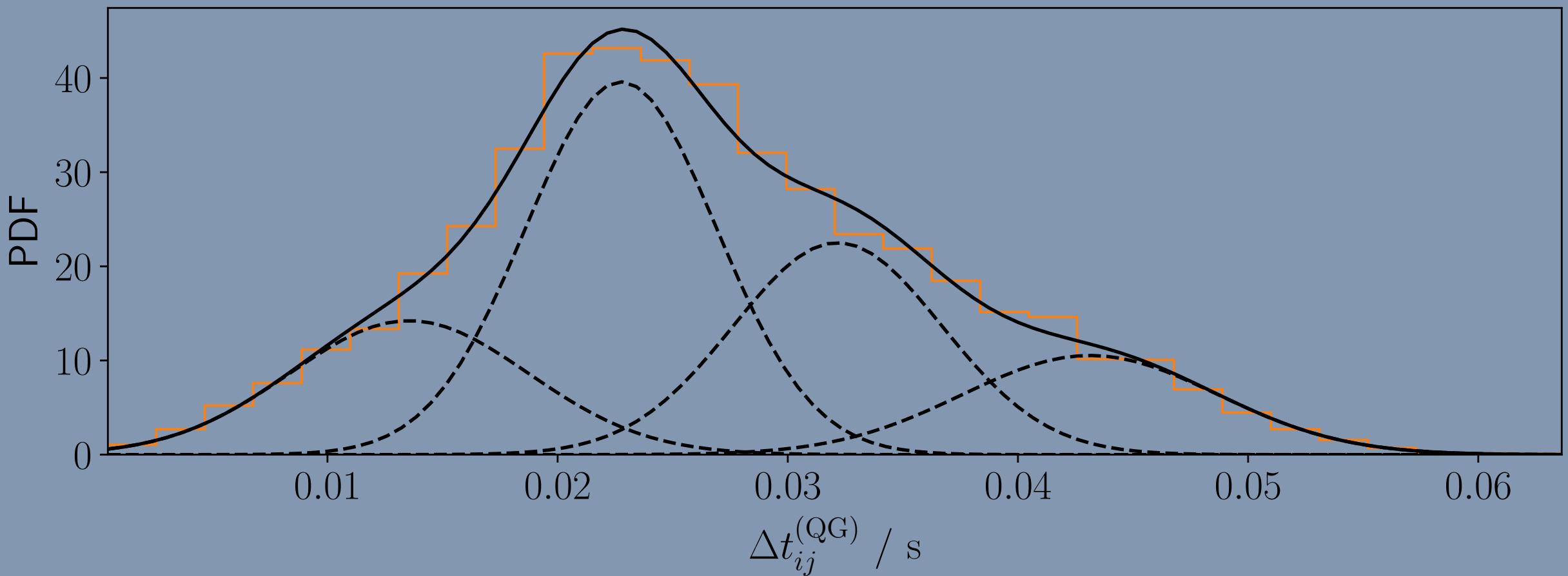
$$\Delta t_{ij}^{(\text{QG})} = \xi \frac{E_i - E_j}{E_{\text{QG}}} \int_0^z \frac{1 + z'}{H(z')} dz'$$

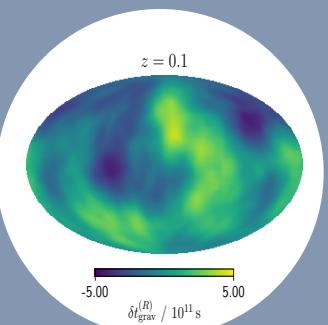
- Draw  $E_i$  and  $E_j$  from distribution proportional to spectrum between band's limits
- Draw redshift from two-sided Gaussian with upper and lower redshift errors
- Calculate  $\Delta t_{ij}^{(QG)}$  for some fiducial  $\ell_{QG} \equiv \xi/E_{QG}$



# Forward Model the Time Delays

Take samples and fit to a Gaussian Mixture Model



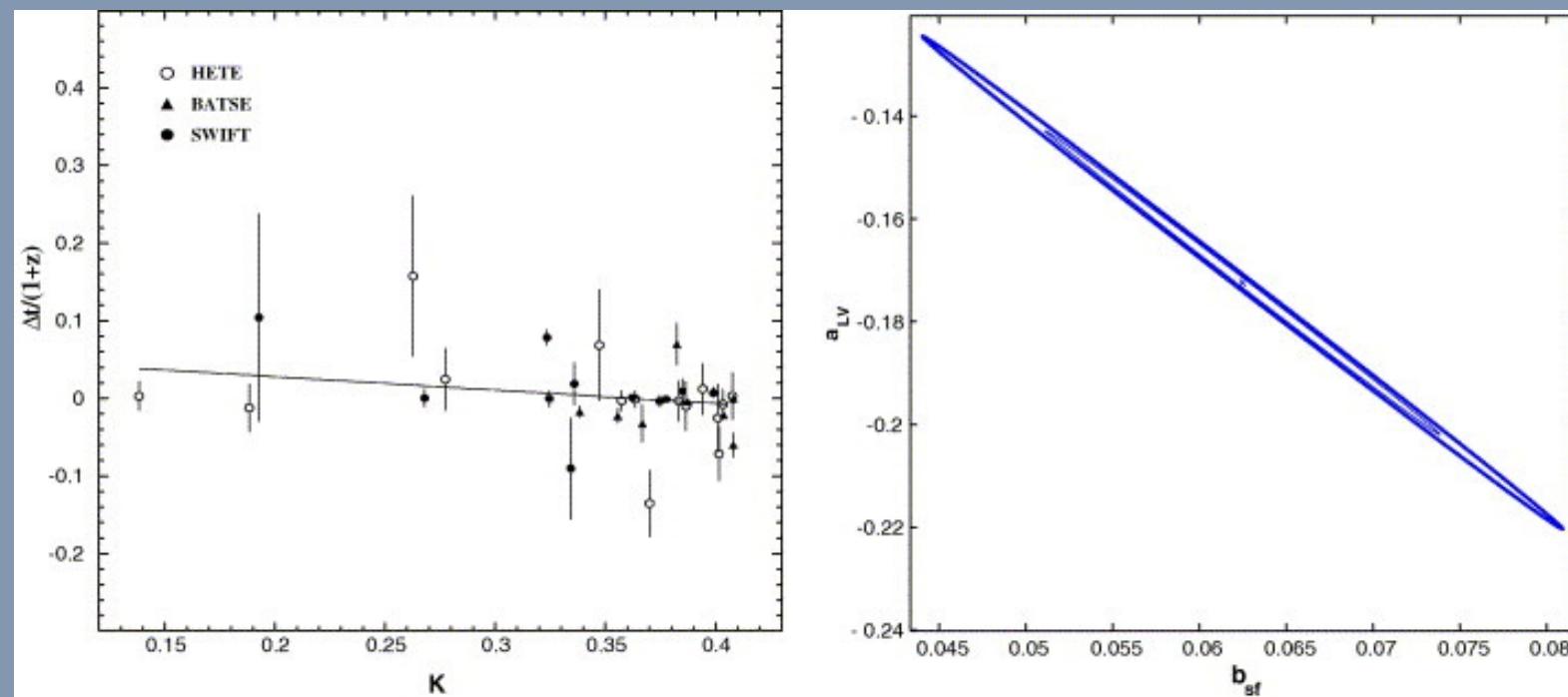


# Other contributions to the time delay

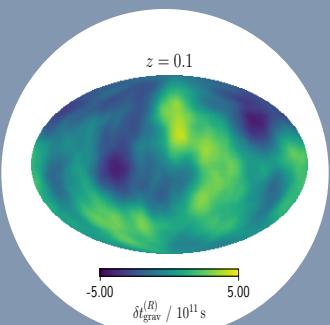
$$\Delta t_{ij}^{(\text{obs})} = \Delta t_{ij}^{(\text{th})} + B_{ij}$$

First approach: Suppose  $B_{ij}$  intrinsic to source and same for all GRBs

$$\mathcal{L}(B_{ij}) = \delta(B_{ij} - b_{ij}(1+z))$$



Ellis et al. 2006 (arXiv:0510172)

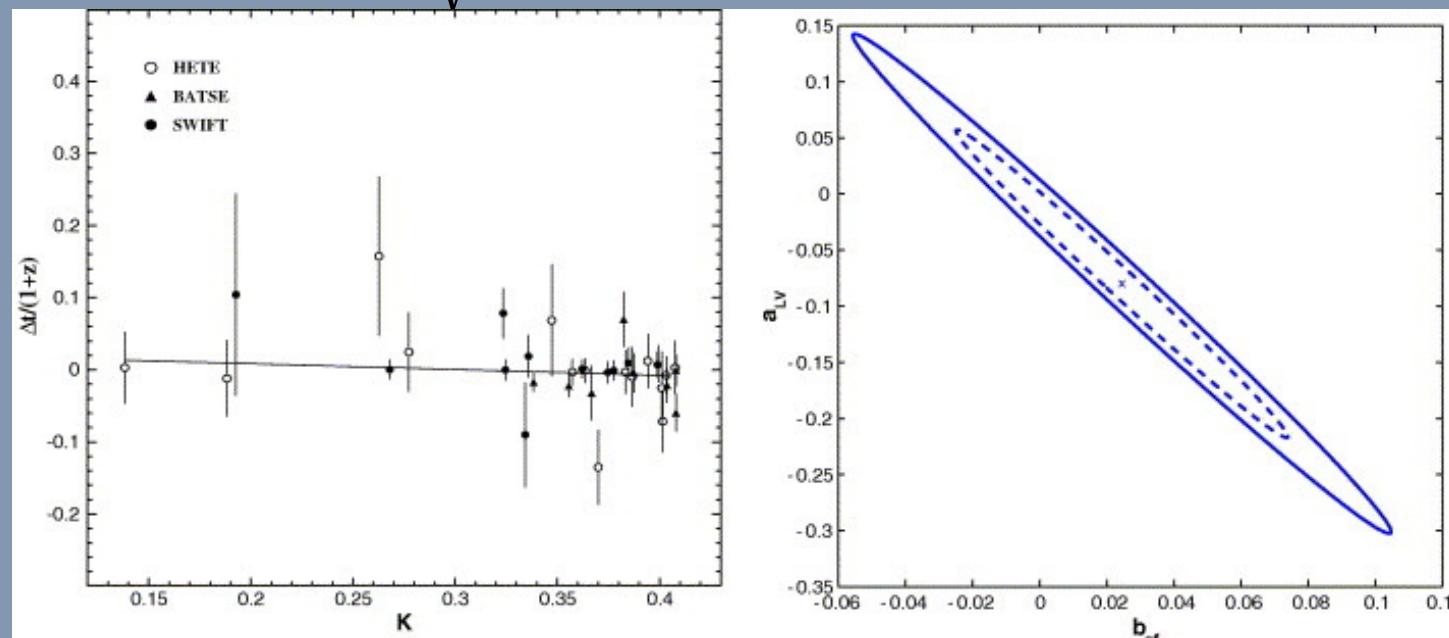


# Other contributions to the time delay

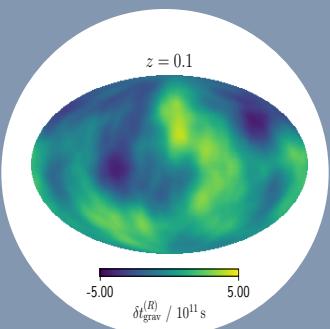
$$\Delta t_{ij}^{(\text{obs})} = \Delta t_{ij}^{(\text{th})} + B_{ij}$$

Second approach: Suppose  $B_{ij}$  intrinsic to source and stochastic

$$\mathcal{L}(B_{ij}) = \frac{1}{\sqrt{2\pi\sigma_{ij}^2}} \exp\left(-\frac{(B_{ij} - b_{ij}(1+z))^2}{2\sigma_{ij}^2}\right)$$



Ellis et al. 2006 (arXiv:0510172)



# Our approach – try several models

Ideally want

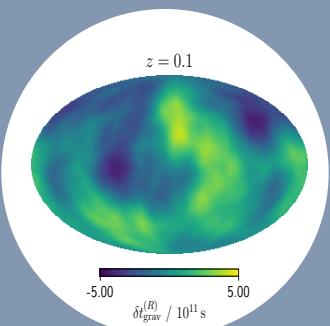
- Try different  $B_{ij}$  models - model comparison
- Results to be (relatively) insensitive to how one models the noise

For just observational noise:

$$\mathcal{L}(B_{ij}) = \sum_{\beta} \frac{\omega_{ij}^{(\beta)}}{\sqrt{2\pi}\sigma_{ij}^{(\beta)}} \exp \left[ -\frac{(B_{ij} - \mu_{ij}^{(\beta)})^2}{2\sigma_{ij}^{(\beta)}{}^2} \right]$$

Can add in intrinsic contributions by allowing some Gaussians to depend on redshift, e.g.

$$\mu_{ij}^{(\beta)} \rightarrow \mu_{ij}^{(\beta)} (1 + z_s), \quad \sigma_{ij}^{(\beta)} \rightarrow \sigma_{ij}^{(\beta)} (1 + z_s)$$



# Recap

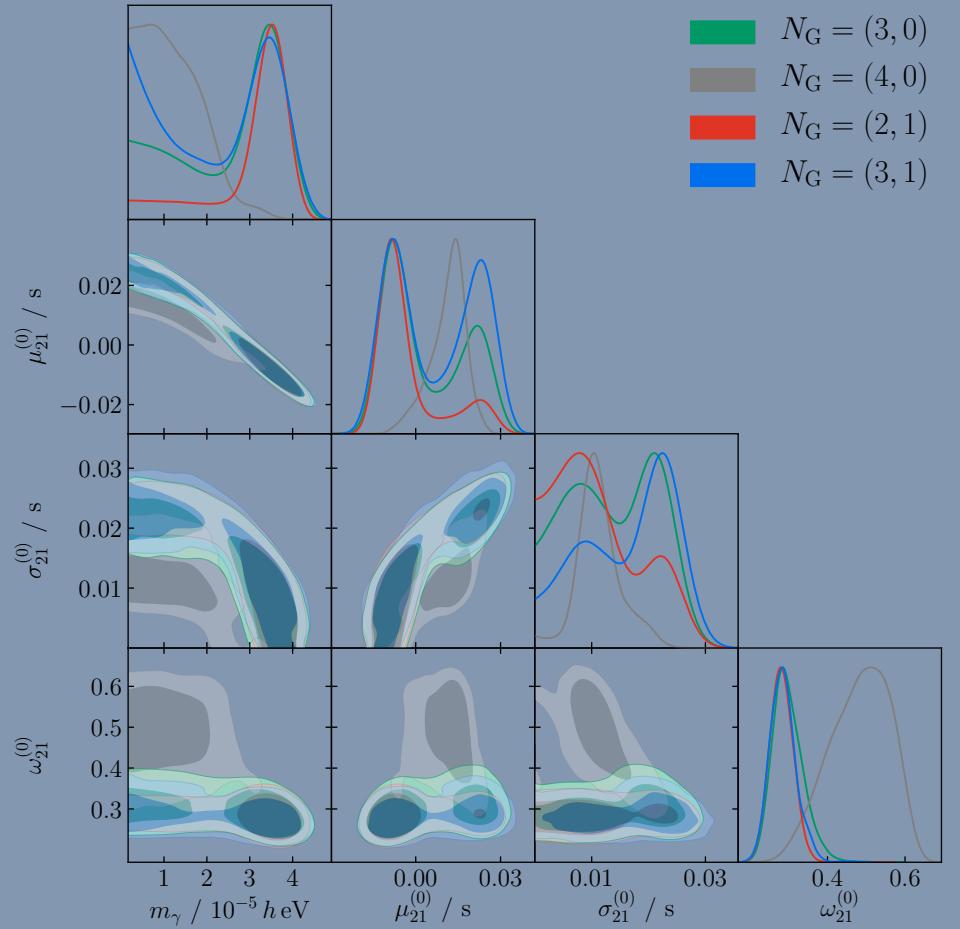
$$\Delta t_{ij}^{(\text{obs})} = \Delta t_{ij}^{(\text{th})} + B_{ij}$$

$$\mathcal{L}_s \left( \Delta t_{ij}^{(\text{QG})} \right) = \sum_{\alpha} \frac{w_{sij}^{(\alpha)}}{\sqrt{2\pi} \tau_{sij}^{(\alpha)}} \exp \left[ -\frac{\left( \Delta t_{ij}^{(\text{QG})} - \lambda_{sij}^{(\alpha)} \right)^2}{2 \tau_{sij}^{(\alpha)}{}^2} \right]$$

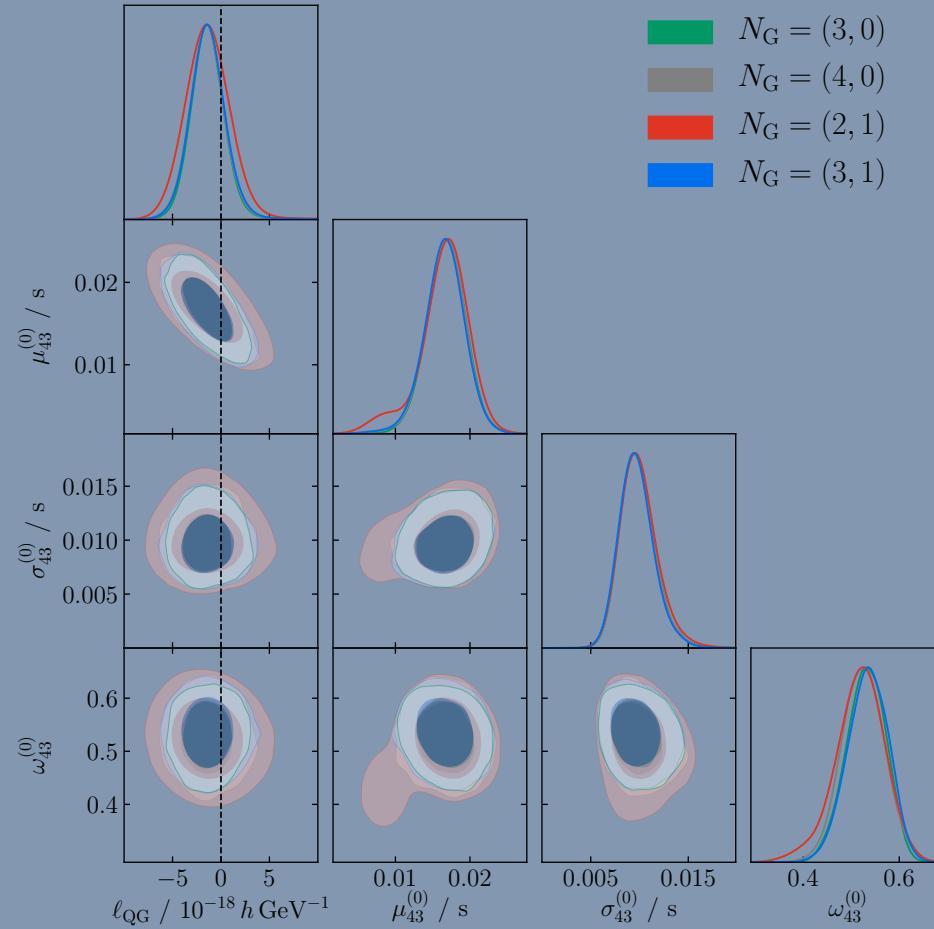
$$\mathcal{L} \left( B_{ij} \right) = \sum_{\beta} \frac{\omega_{ij}^{(\beta)}}{\sqrt{2\pi} \sigma_{ij}^{(\beta)}} \exp \left[ -\frac{\left( B_{ij} - \mu_{ij}^{(\beta)} \right)^2}{2 \sigma_{ij}^{(\beta)}{}^2} \right]$$

$$\mathcal{L} \left( \mathcal{D} | \vec{\theta} \right) = \prod_{sij} \mathcal{L}_s \left( \Delta t_{ij}^{(\text{obs})} \right)$$

# Results

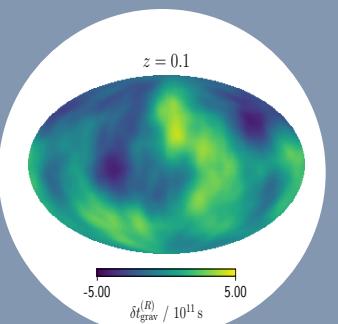


$m_\gamma < 4.0 \times 10^{-5} h \text{ eV}/c^2$  (95% confidence)

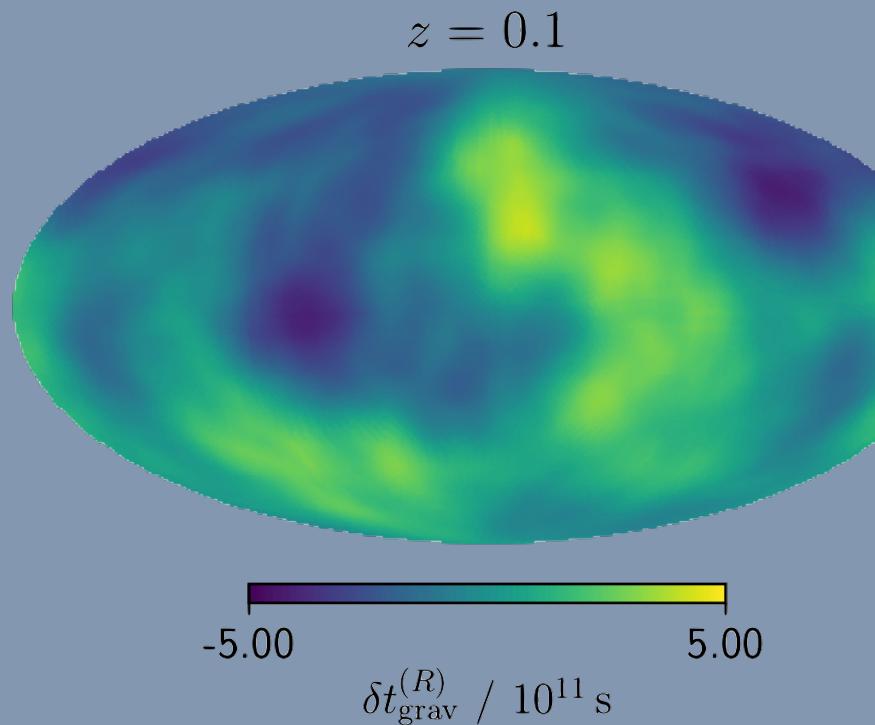


$|\ell_{\text{QG}}| < 5.3 \times 10^{-18} h \text{ GeV}^{-1}$  (95% confidence)

Bartlett, Desmond, Ferreira & Jasche 2021 (arXiv:2109.07850)



# Other tests?

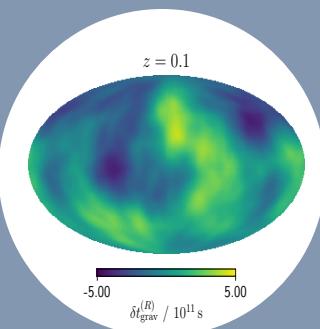


## Shapiro time delays

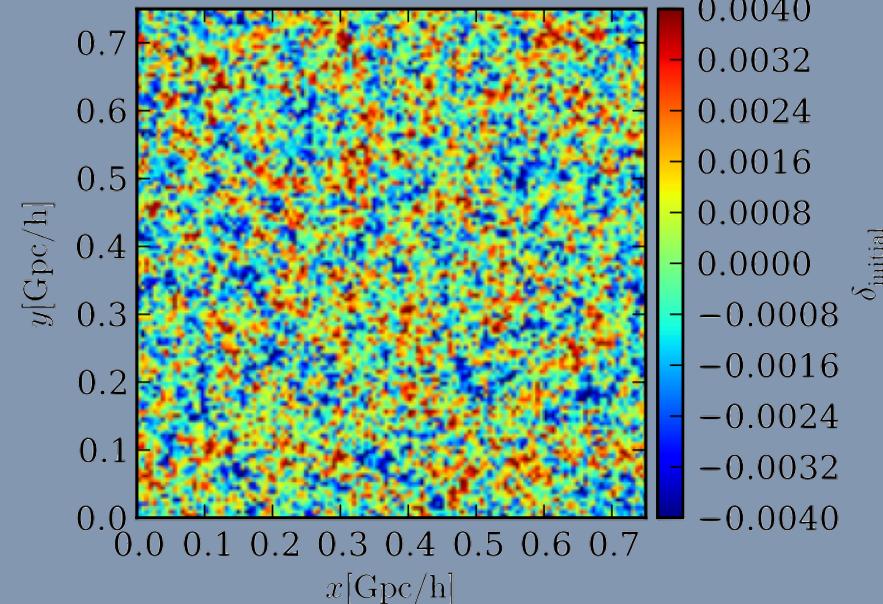
$$ds^2 = -(1 + 2\delta\phi) dt^2 + a^2(t) (1 - 2\gamma\delta\phi) d\vec{r}^2$$

$$t_{\text{grav}} = -(1 + \gamma) \int_0^{r_s} dr \delta\phi_0(\vec{r}) D(r)$$

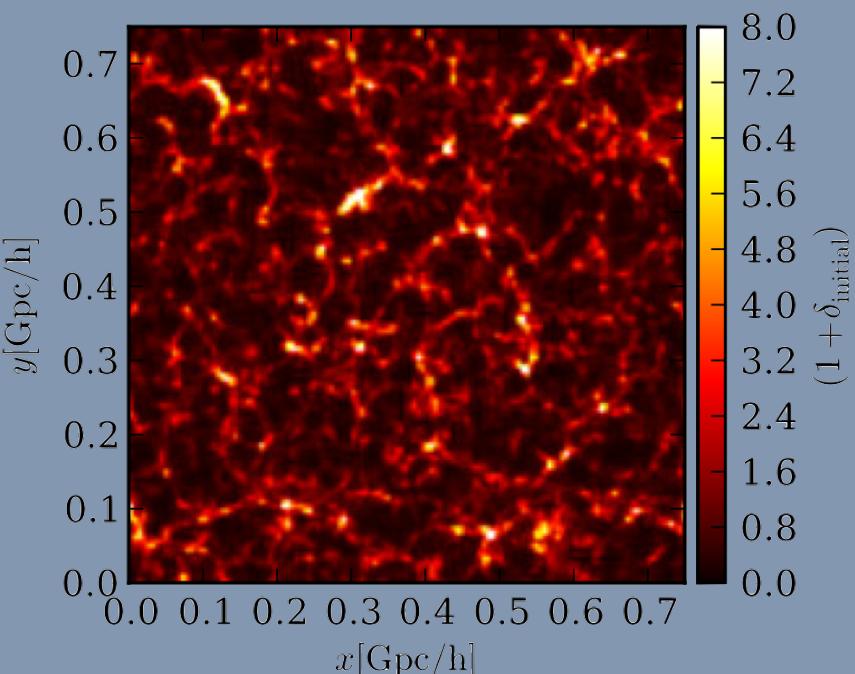
## Here can use spatial variation



# BORG - Bayesian Origin Reconstruction from Galaxies



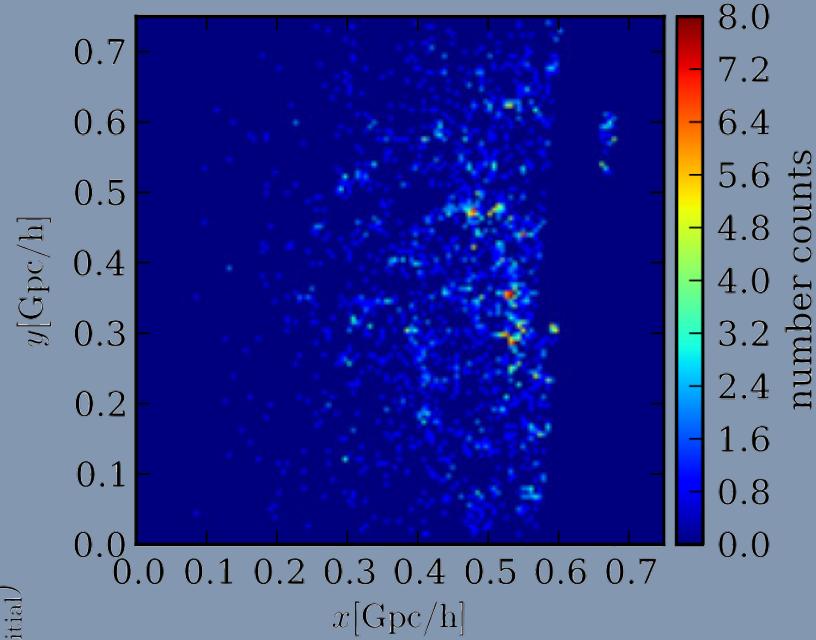
$$\mathcal{P}(\mathbf{s}|\mathbf{S}) = \frac{e^{-\frac{1}{2}\mathbf{s}^T\mathbf{S}^{-1}\mathbf{s}}}{\sqrt{\det(2\pi\mathbf{S})}}$$



$$\mathcal{P}(\delta|\mathbf{s}) = \prod_i \delta^D(\delta_i - G_i(\mathbf{s}))$$

$$\frac{d\vec{x}}{da} = \frac{\vec{p}}{\dot{a}a^2}, \quad \frac{d\vec{p}}{da} = -\frac{\nabla_{\vec{x}}\Phi}{aH(a)}$$

$$\nabla_{\vec{x}}^2\Phi = \frac{3}{2}H_0^2\Omega_{m,0}\frac{\delta_m(\vec{x})}{a}$$

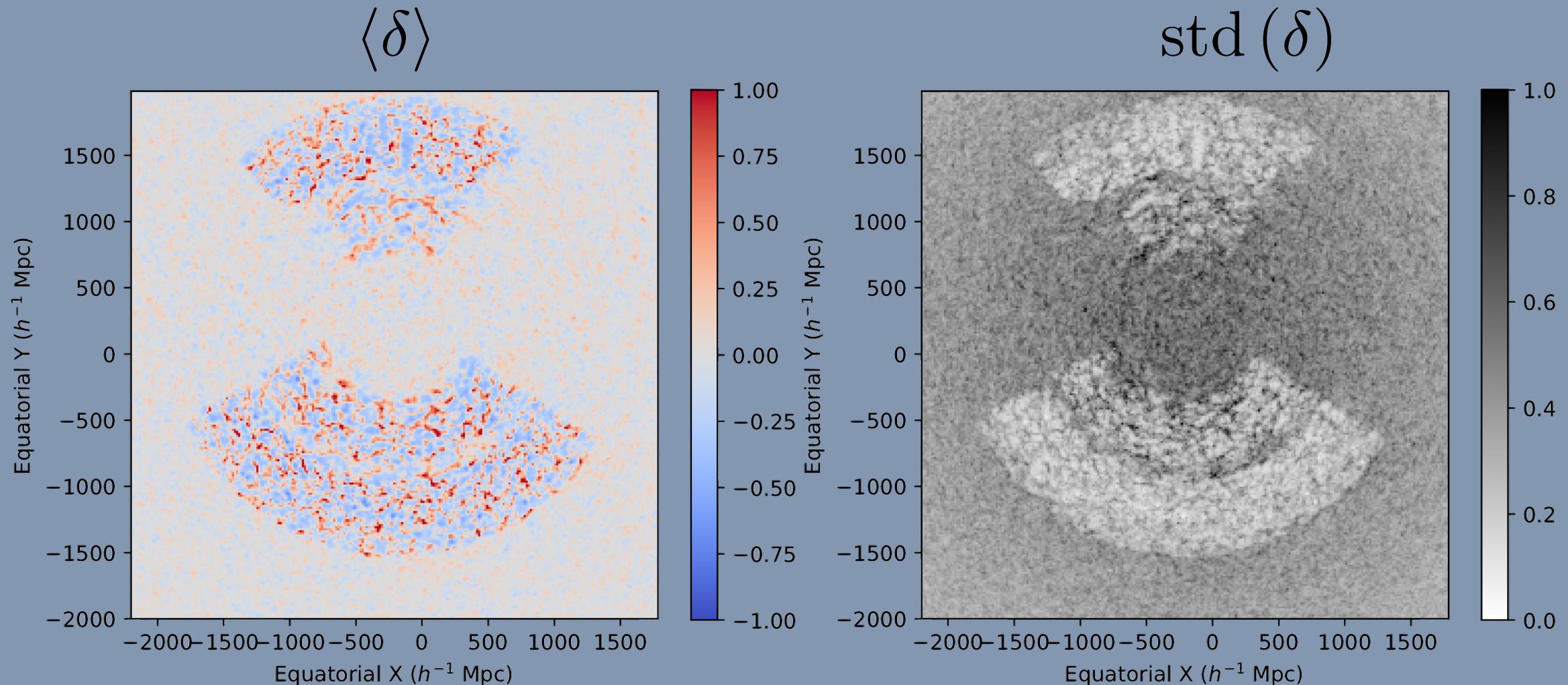


$$\mathcal{P}(\mathbf{N}|\lambda(\delta)) = \prod_i \frac{e^{-\lambda_i}\lambda_i^{N_i}}{N_i!}$$

+ galaxy bias model  $\lambda(\delta)$

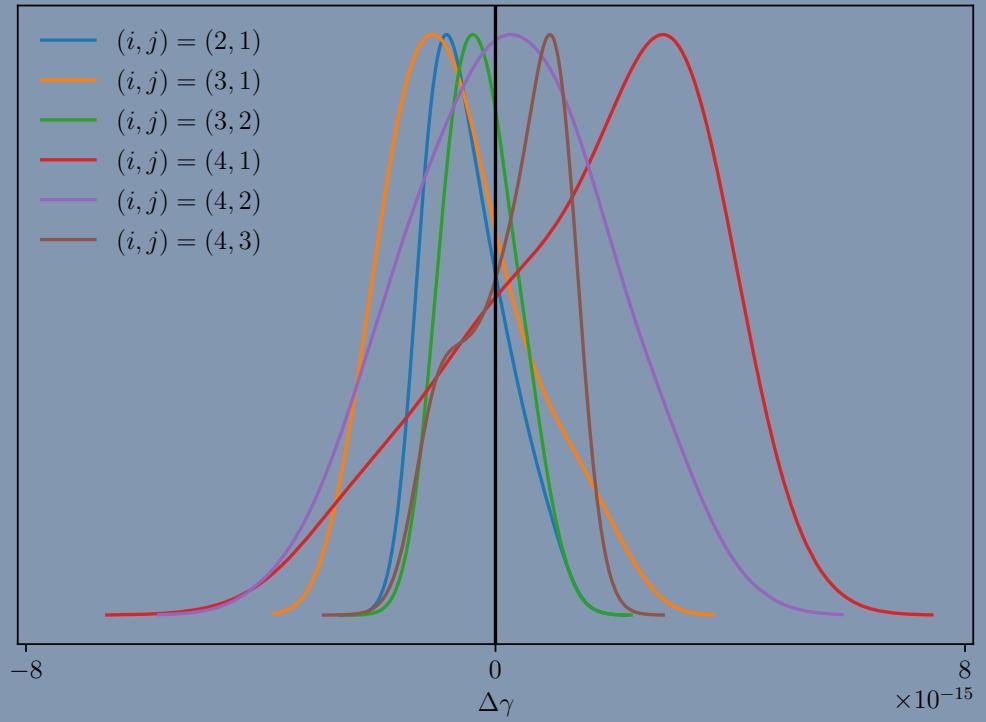
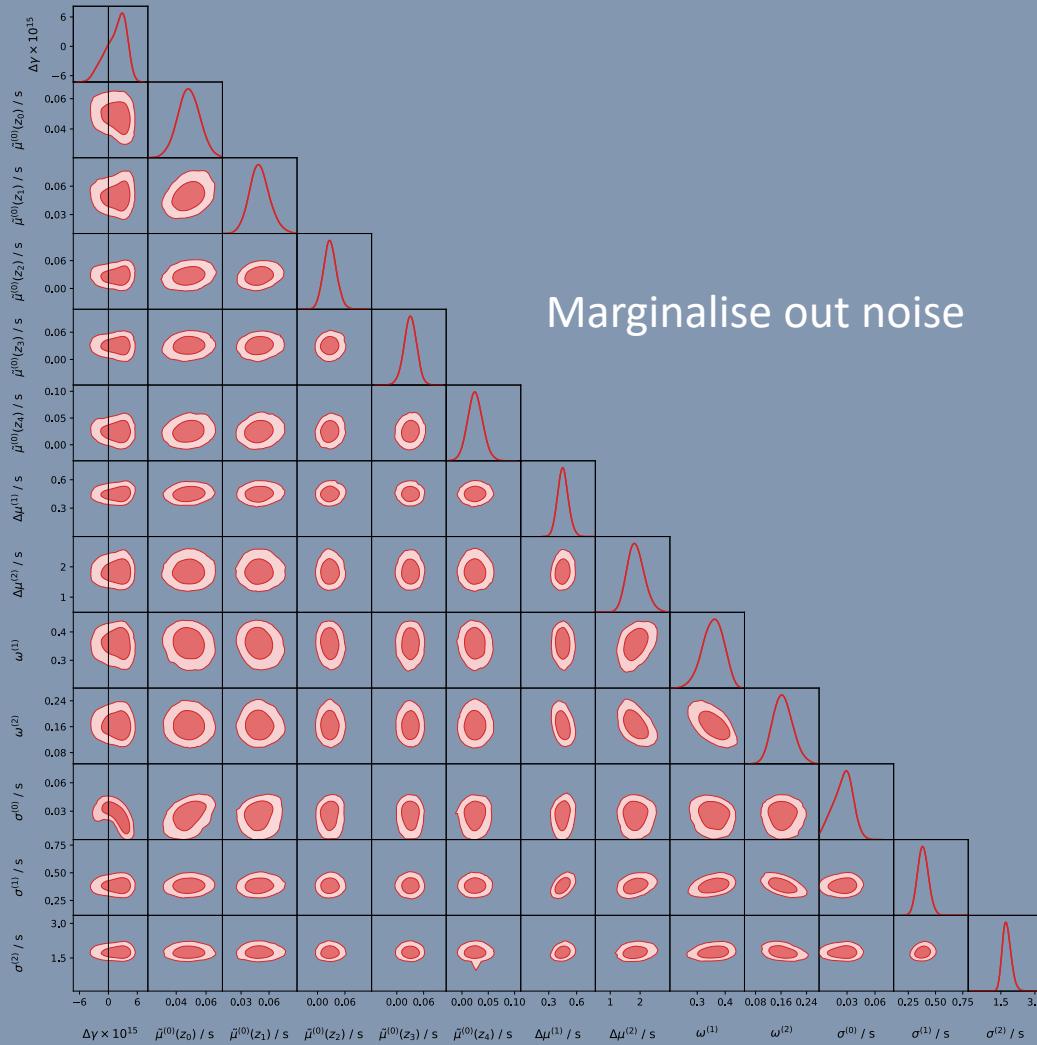
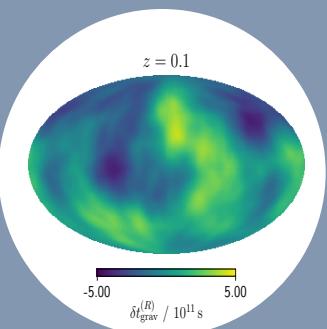
Jasche & Wandelt 2012 (arXiv: 1203.3639)  
Jasche & Lavaux 2019 (arXiv: 1806.11117)

# BORG - Bayesian Origin Reconstruction from Galaxies



Lavaux, Jasche & Leclercq 2019 (arXiv:1909.06396)

# Tightest constraints on $\Delta\gamma$

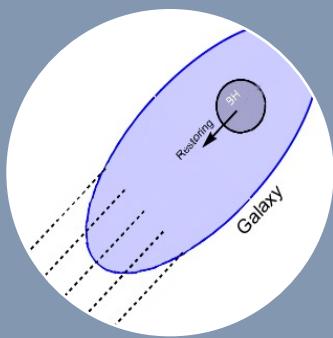


$E_{\text{QG}} \gtrsim 10^{11} \text{ GeV}$  (1 $\sigma$  confidence)

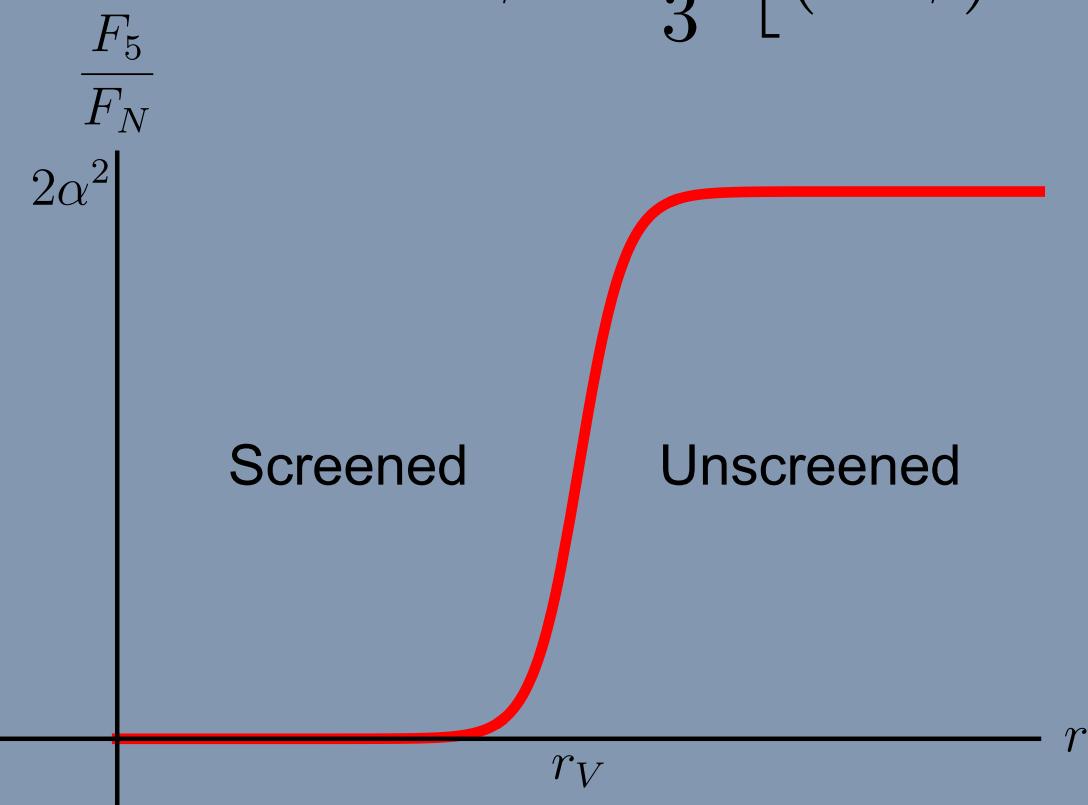
Bartlett, Bergsdal, Desmond, Ferreira & Jasche 2021 (arXiv:2106.15290)

# Black Holes and Galileons

# Galileons result in a screened fifth force



$$\nabla^2 \phi + \frac{r_C^2}{3} \left[ (\nabla^2 \phi)^2 - \nabla_i \nabla_j \phi \nabla^i \nabla^j \phi \right] = 8\pi \alpha G_N \bar{\rho} \Delta$$

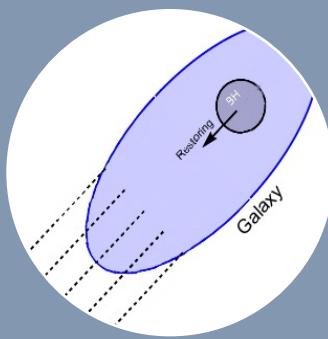


$$M \ddot{\vec{x}} = \vec{F}_N + \vec{F}_5 = -M \vec{\nabla} \Phi - \alpha Q \vec{\nabla} \phi$$

$$r_V = \left( \frac{4}{3} \alpha G_N M r_C^2 \right)^{\frac{1}{3}}$$

$$r_{V,\text{Sun}} \sim 100 \text{ pc}$$

# Large scale structure sources Galileon field



**Problem:** Black holes and stars within Vainshtein radius of host

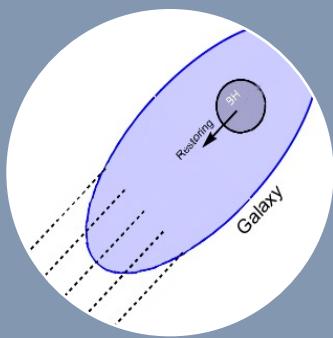
**Solution:** Galileon Symmetry

$$\phi \rightarrow \phi + b + c_\mu x^\mu$$

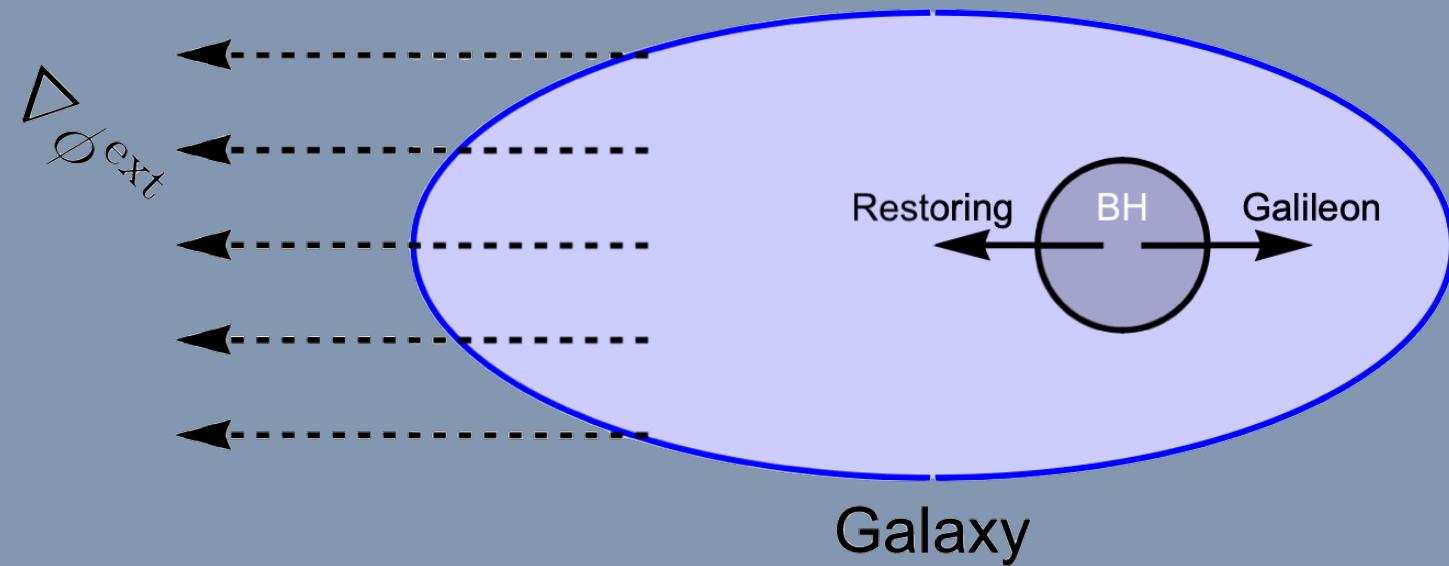
Can feel long wavelength modes, i.e. from large scale structure

$$M\ddot{\vec{x}} = -M\vec{\nabla}\Phi - \alpha Q\vec{\nabla}\phi_{\text{lss}}$$

# BHs do not feel a fifth force - offset BHs



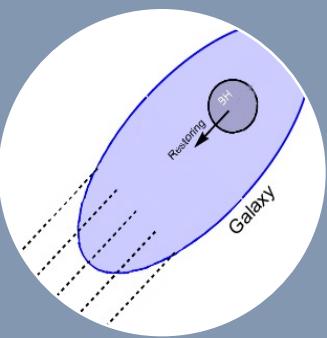
Stars, gas, dark matter:  $Q = M$       Black holes:  $Q = 0$



$$r_\bullet = 0.1 \text{ kpc} \left( \frac{2\alpha^2}{1} \right) \left( \frac{|g_{lss}|}{20 (\text{km s}^{-1})^2 \text{ kpc}^{-1}} \right) \left( \frac{0.01 \text{ M}_\odot \text{ pc}^{-3}}{\rho_0} \right)$$

Hui & Nicolis 2012 (arXiv:1201.1508)

# Properties, Environment & Astrophysics

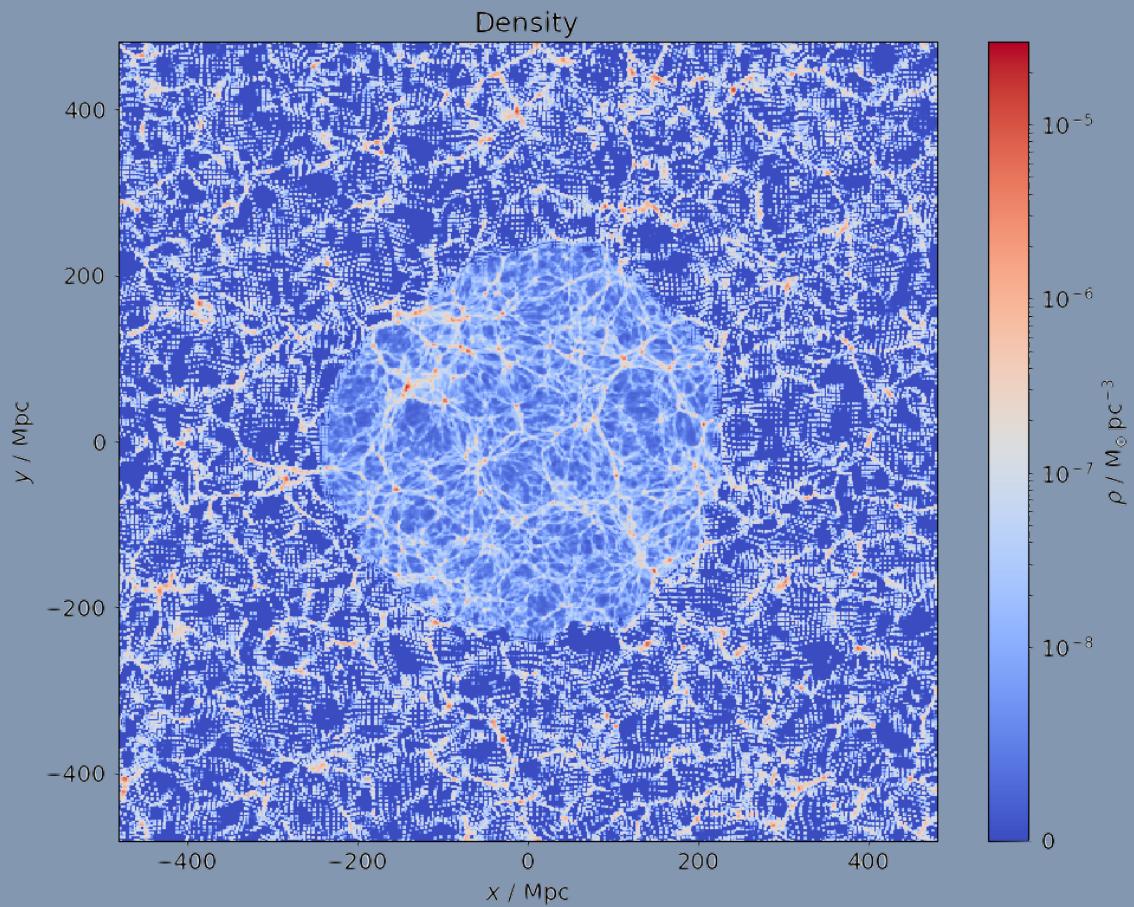


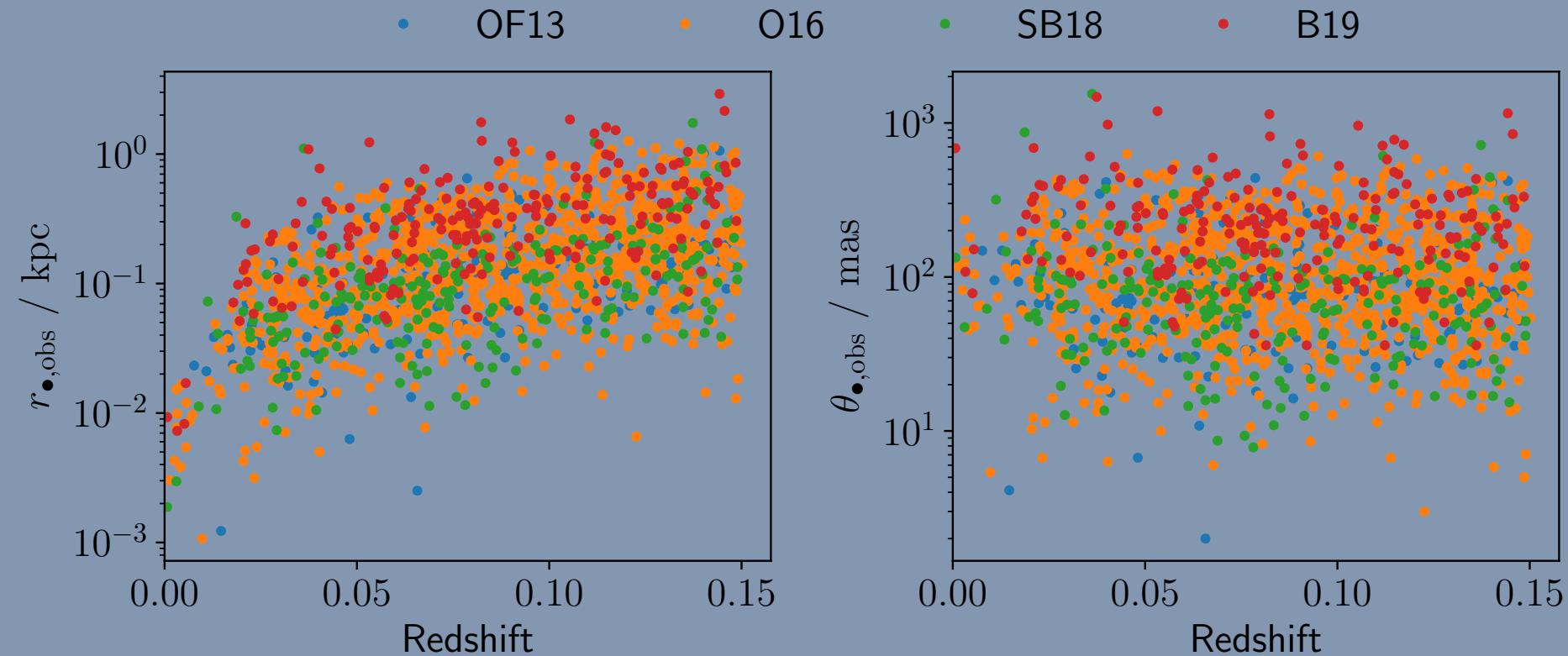
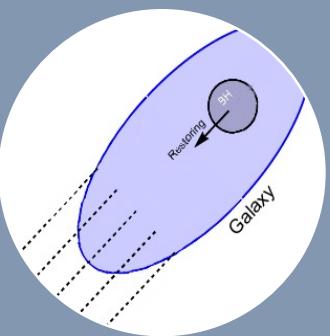
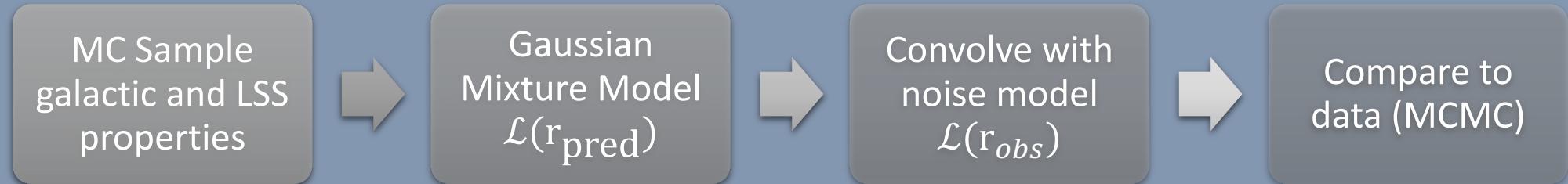
$$\ddot{r} = -\frac{GM(<r)}{r^2} - \alpha \vec{\nabla} \phi_{\text{lss}}$$

1. Galaxy properties -  $M(<r)$
2. Large scale structure -  $\phi_{\text{lss}}$
3. Knowledge of non-Galileon effects

# CSiBORG – Constrained Simulations in BORG

- Suite (100) of DM-only simulations - marginalize
- Initial conditions ( $z=69$ ) from BORG reconstruction of 2M++
- Mass resolution  $4.4 \times 10^8 M_{\odot}$
- Density field  $\rightarrow$  gravitational field
- Filter out high  $k$  modes to mimic Vainshtein mechanism ( $k_V = 2\pi/r_V$ )
- Add in unconstrained small  $k$  modes





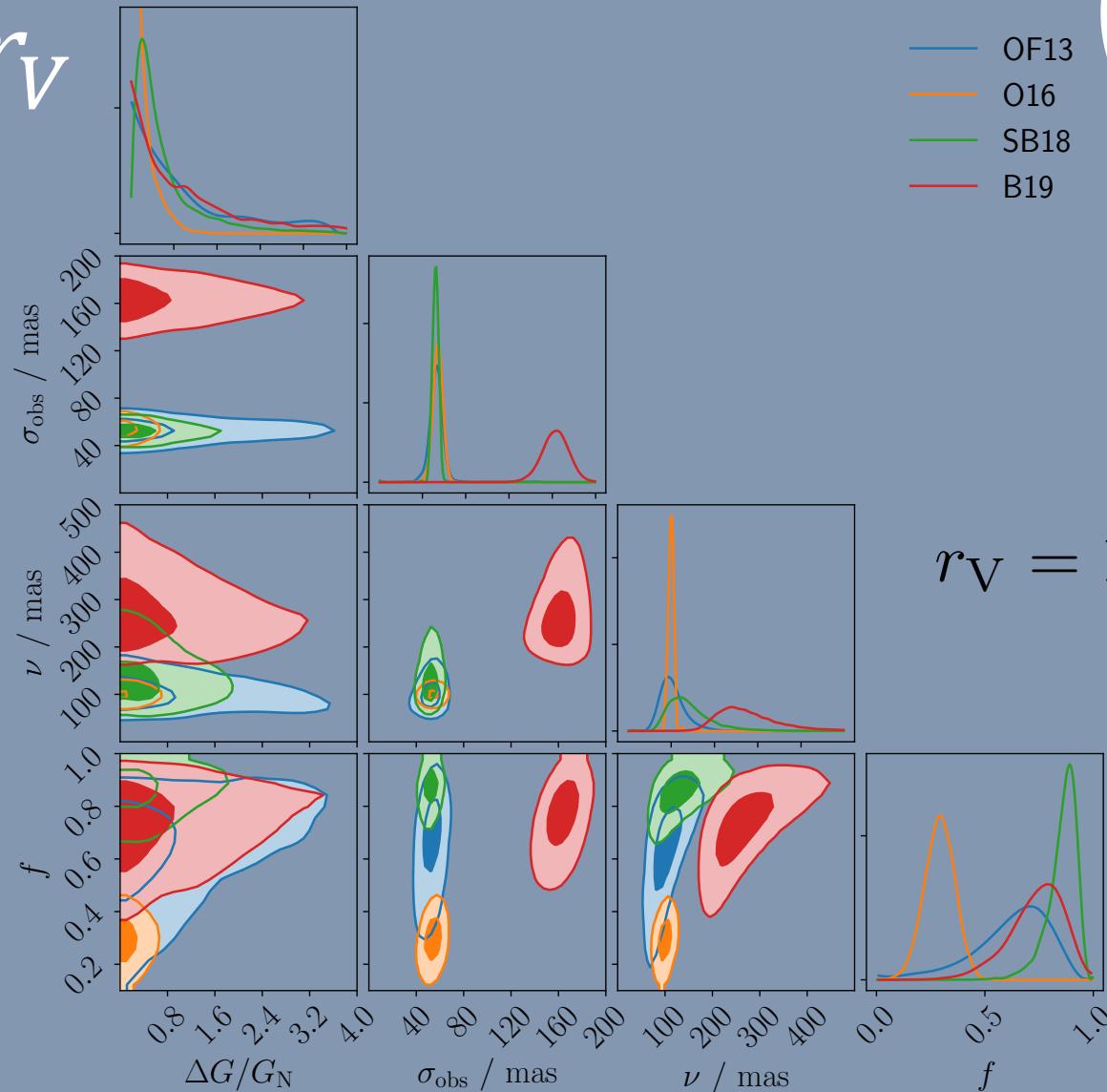
Bartlett, Desmond & Ferreira (arXiv:2010.05811)

# Constrain at fixed $r_V$

$$\mathcal{L}_g(\theta_{\alpha,\text{obs}}|\theta_{\alpha,\text{pred}}) = \frac{f}{\sqrt{2\pi}\sigma_{\text{obs}}} \exp\left(-\frac{(\theta_{\alpha,\text{obs}} - \theta_{\alpha,\text{pred}})^2}{2\sigma_{\text{obs}}^2}\right) + \frac{1-f}{2\nu} \exp\left(-\frac{|\theta_{\alpha,\text{obs}} - \theta_{\alpha,\text{pred}}|}{\nu}\right)$$

Can include intrinsic contribution:

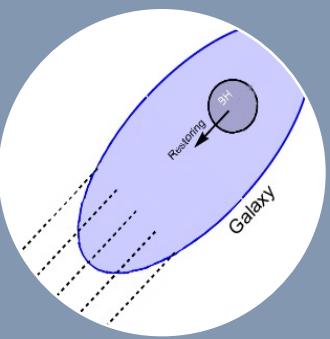
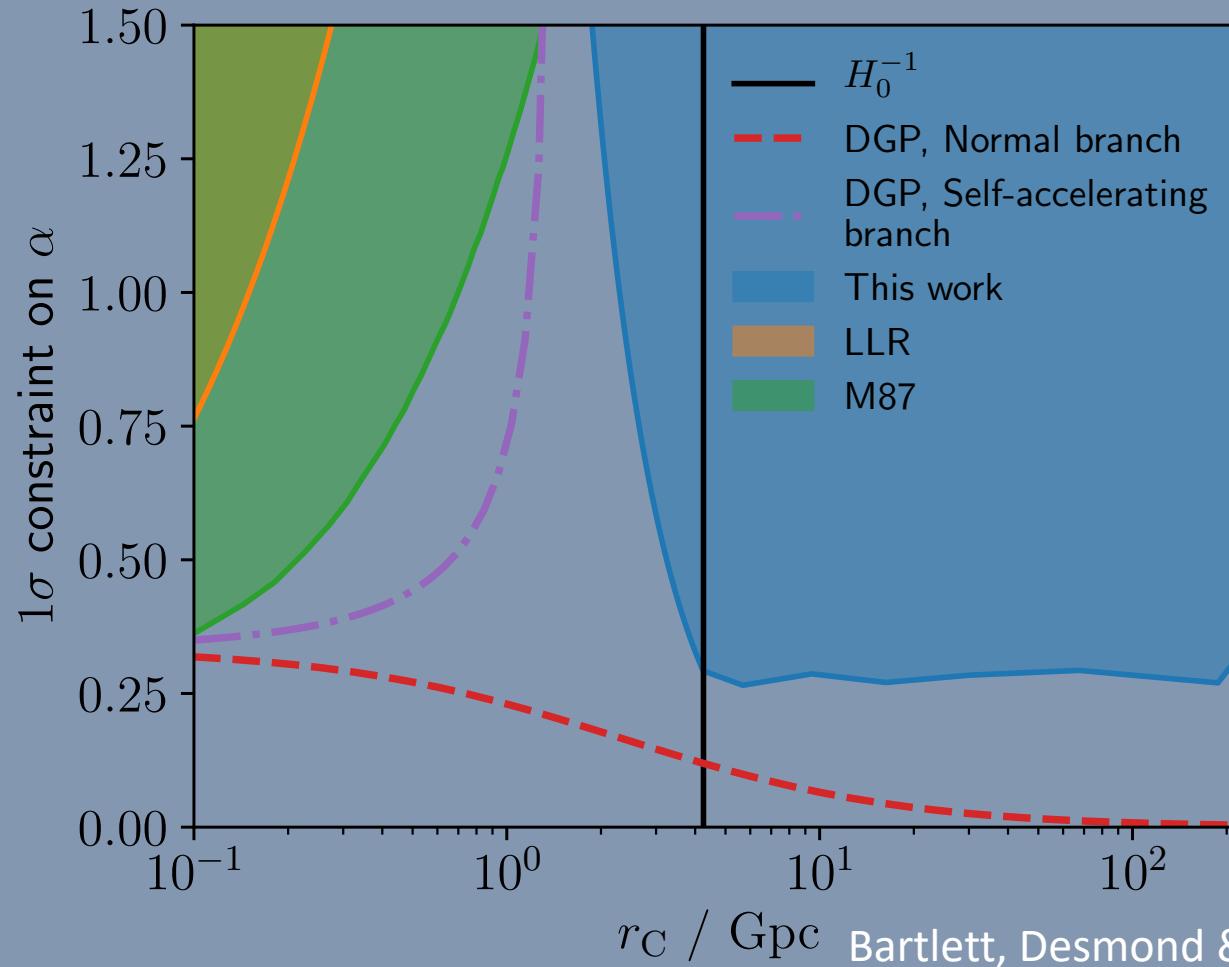
$$\sigma_{\text{obs}} \rightarrow \sqrt{\sigma_{\text{obs}}^2 + \left(\frac{\sigma_{\text{int}}}{d_A}\right)^2}$$



Bartlett, Desmond & Ferreira 2021 (arXiv:2010.05811)

Constrain  $\frac{\Delta G}{G_N} < 0.16$  at  $1\sigma$

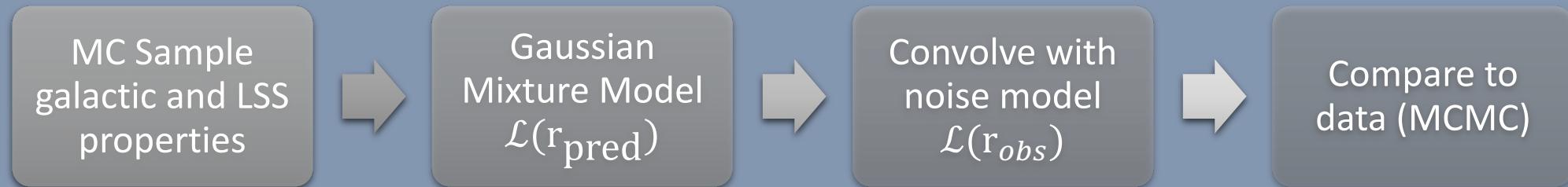
$$\frac{\Delta G}{G_N} \equiv 2\alpha^2$$



Bartlett, Desmond & Ferreira 2021 (arXiv:2010.05811)

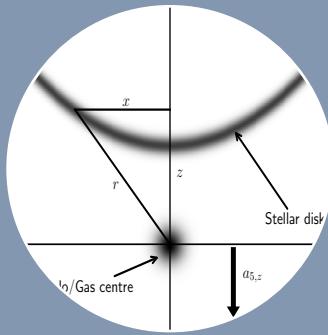
# Modeling the noise

# Recap



- Noise model – simple and empirical
- Astrophysics is messy – could be degenerate with my signal
- How do I know I can trust what I've done?

# What questions should I ask?

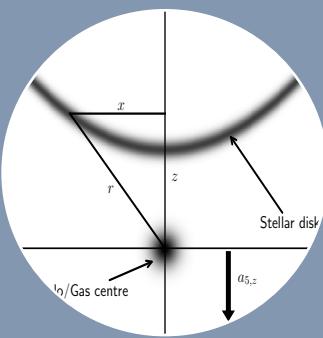


$$y_i = \bar{y}_i + \text{noise}$$

$$\mathcal{L}_i(y_i|\bar{y}_i) = \frac{1}{\sqrt{2\pi}\sigma_{y,i}} \exp\left(-\frac{(y_i - \bar{y}_i)^2}{2\sigma_{y,i}^2}\right)$$

1. Am I missing any correlations?
2. How significant are these correlations?
3. Is my model for  $\bar{y}_i$  sufficient?

Bartlett, Desmond & Ferreira 2021 (arXiv:2103.10356)



# Example – Screened Fifth Forces

$$S = \int d^4x \sqrt{-g} \frac{R + f(R)}{16\pi G_N} + S_m$$

Object unscreened if

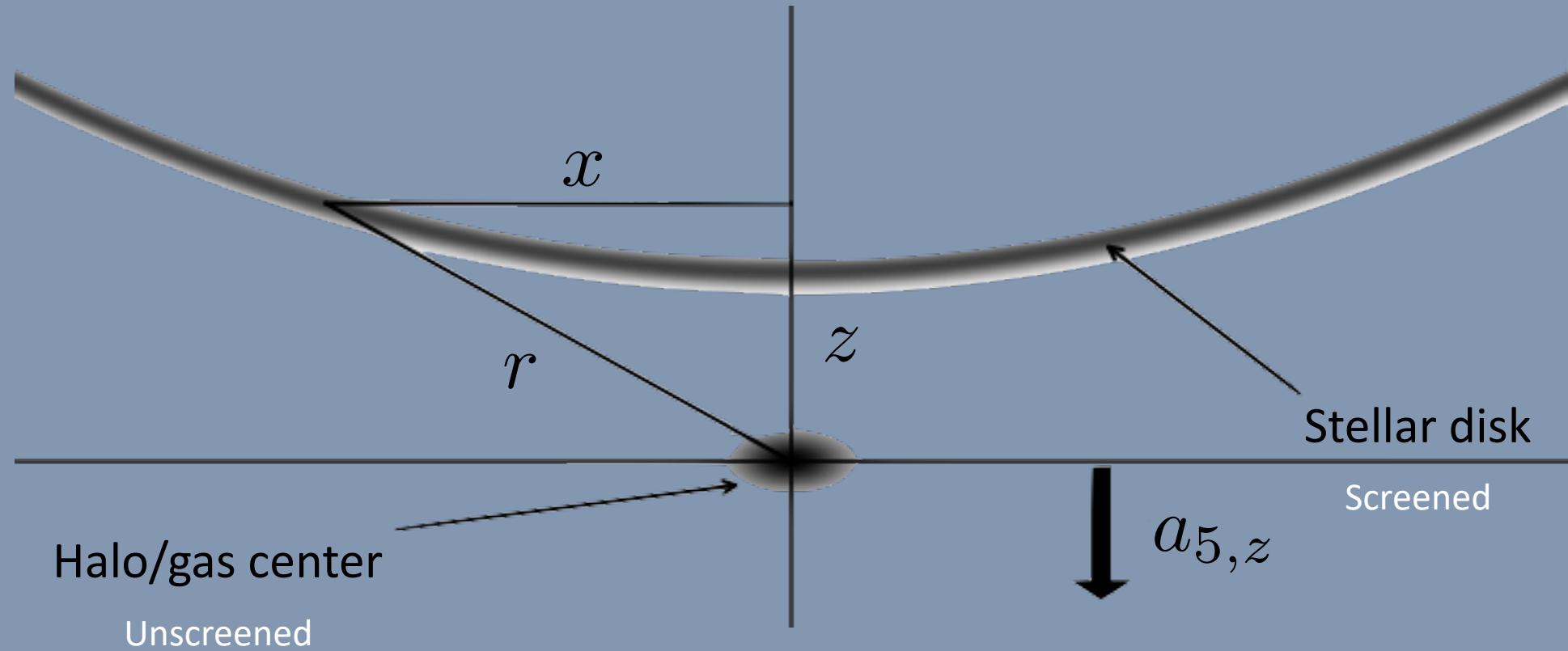
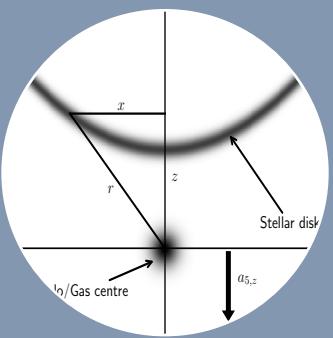
$$|\Phi| < \chi = \frac{3}{2} f_{R0} = \frac{3}{2} \times 10^{-8} \left( \frac{\lambda_c}{0.32 \text{ Mpc}} \right)^2$$

and thus feel fifth-force

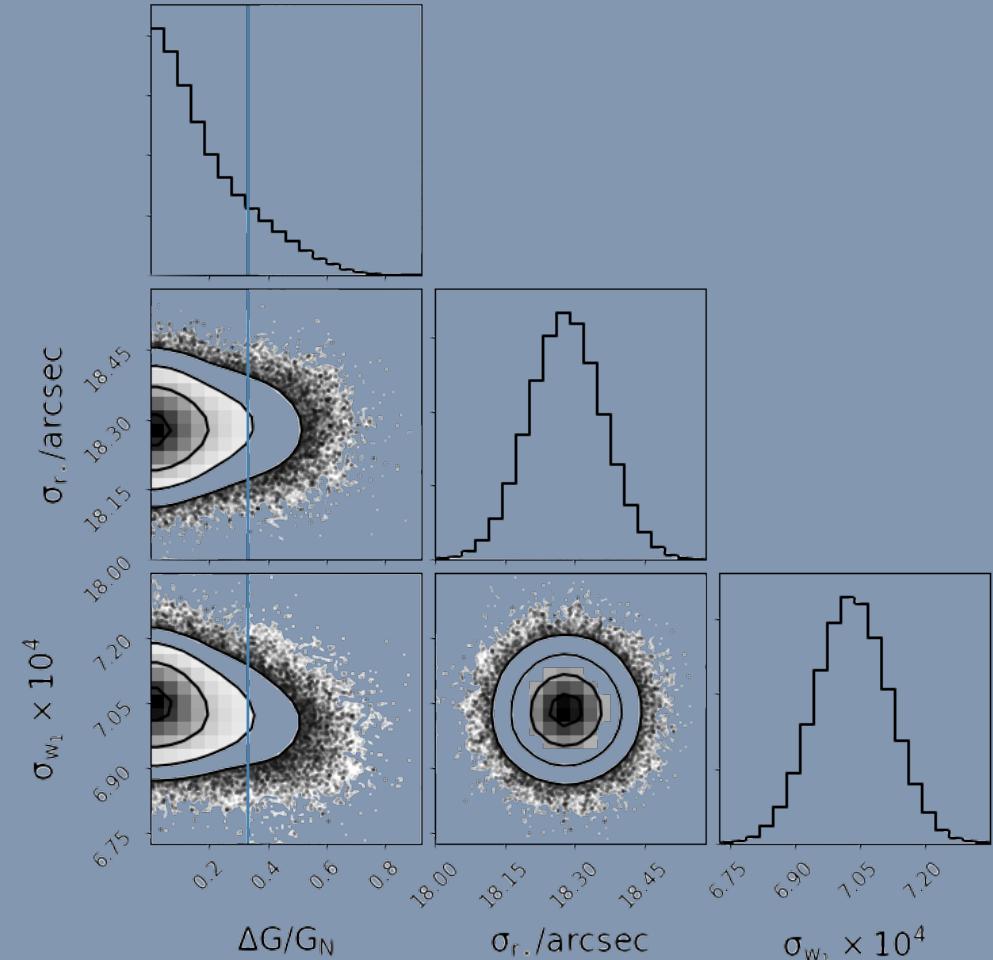
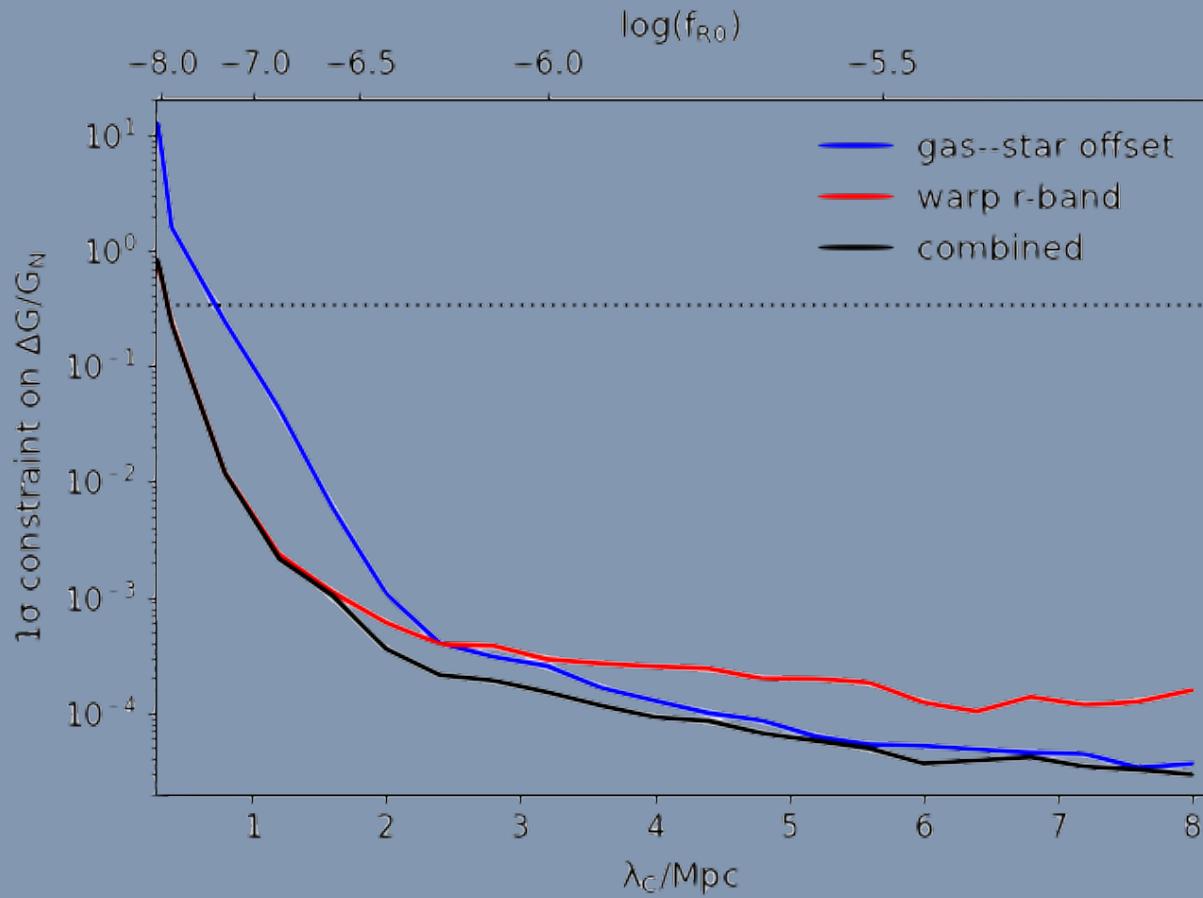
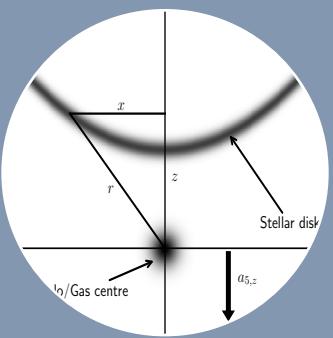
$$G_N \rightarrow G_N \left( 1 + \frac{\Delta G}{G_N} \right)$$

Stars will be “self-screened” for  $\chi \lesssim 10^{-6}$

# Gas-star offsets and galaxy warps

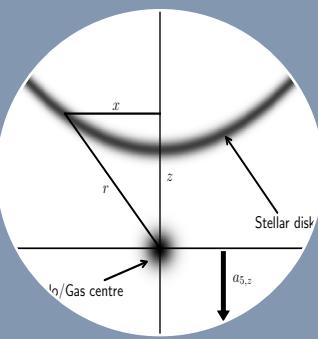


# Astrophysically Relevant Hu-Sawicki $f(R)$ ruled out

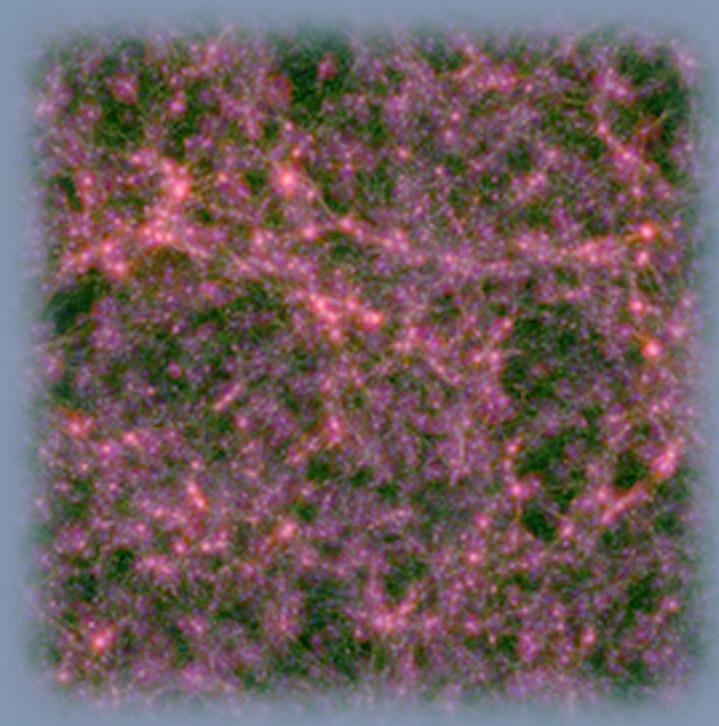


Desmond & Ferreira 2020 (arXiv:2009.08743)

# Use a hydrodynamical simulation to test model

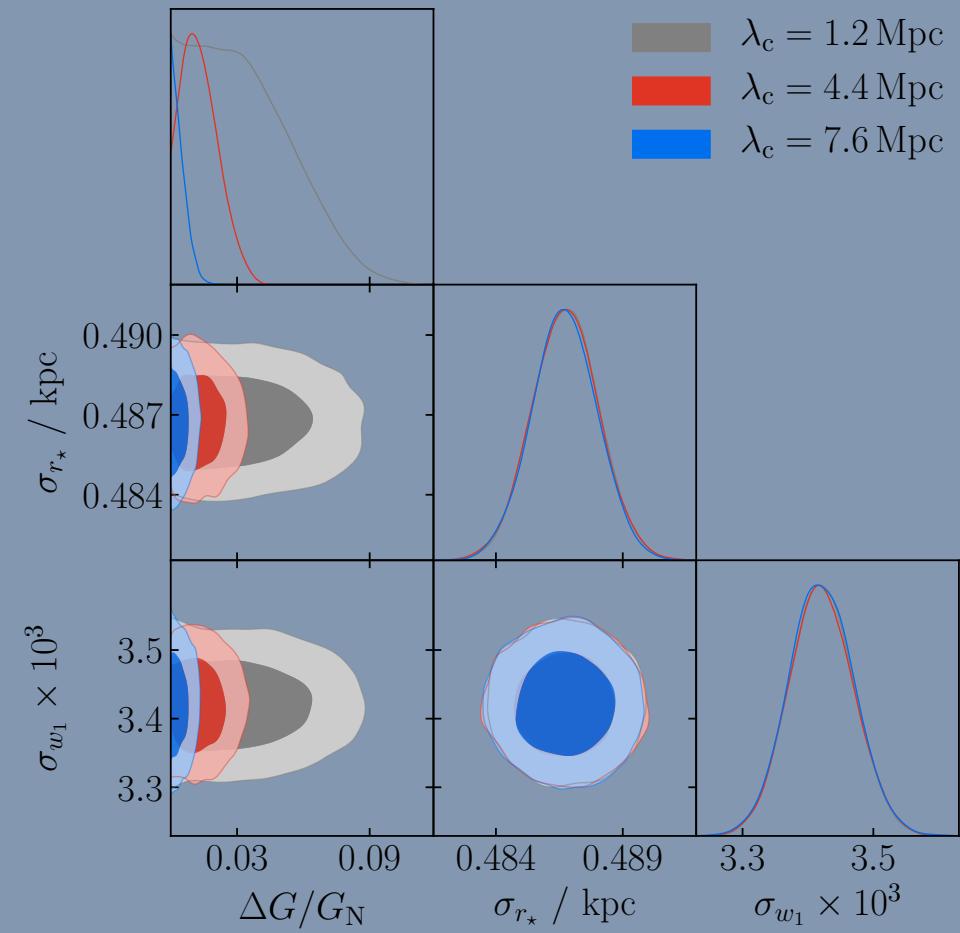
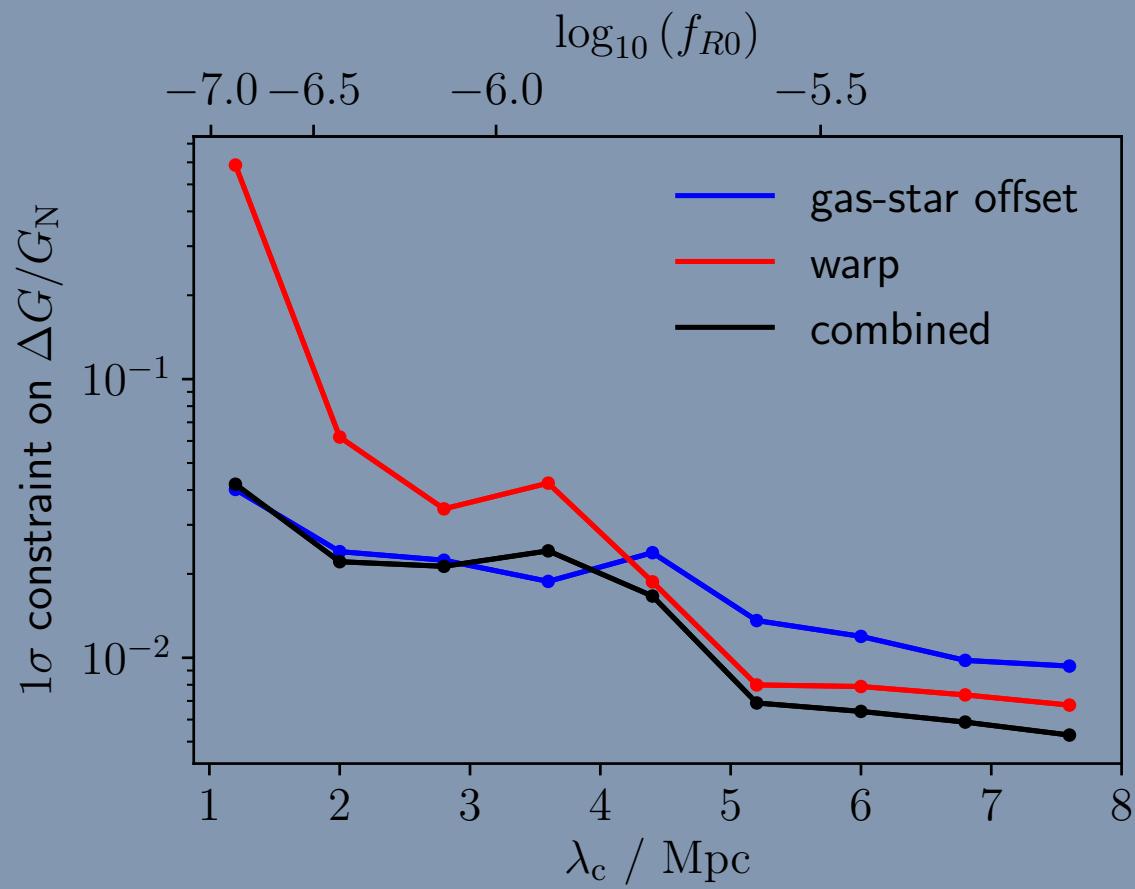
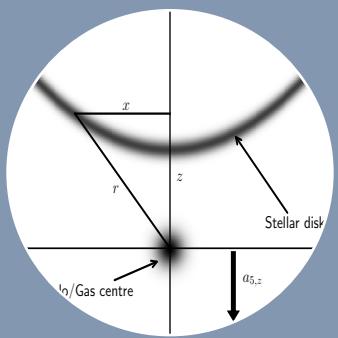


1. We know exactly what the theoretical parameters and implementation of baryonic physics are in the simulation
2. We have more information available there than we do observationally
  - Horizon-AGN simulation:  $(100 \text{ Mpc}/h)^3$ , 126,361 galaxies, baryonic effects
  - $\Lambda\text{CDM}$  – should get no fifth-force detection



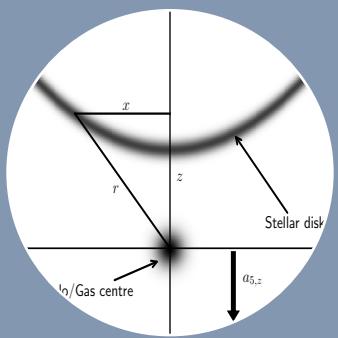
<http://www.horizon-simulation.org/about.html>

# First check you don't get a detection

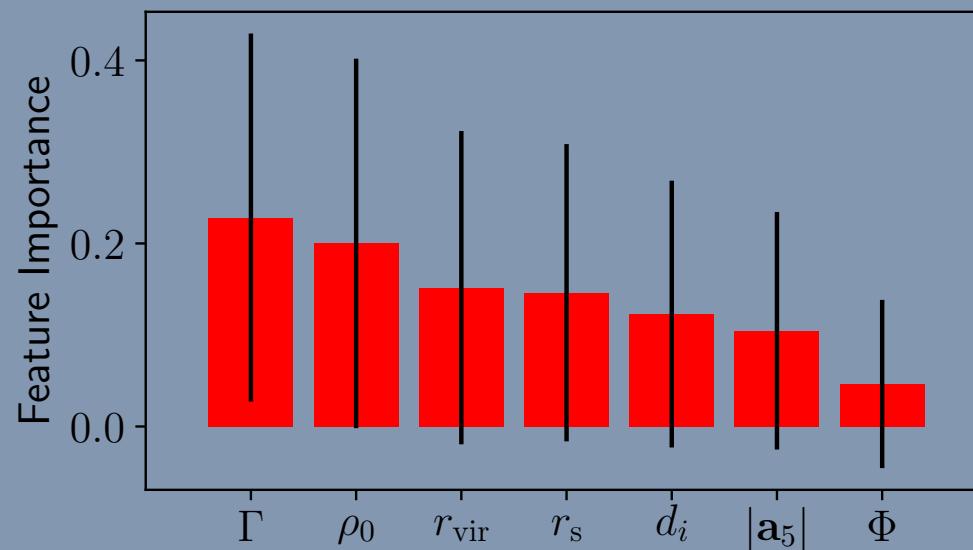


Bartlett, Desmond & Ferreira 2021 (arXiv:2103.10356)

# 1. Am I missing any correlations?

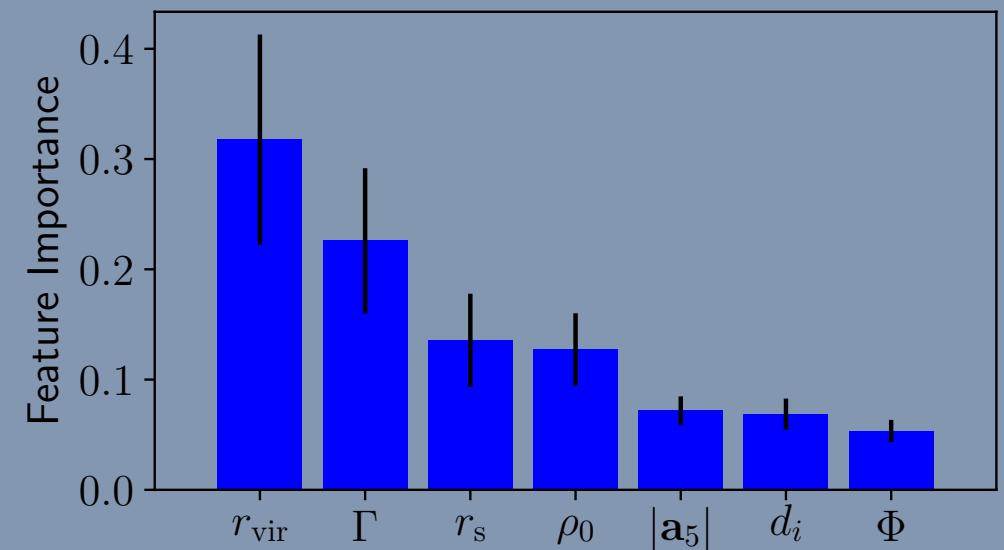


Galaxy warp



Cannot predict → Noise model should be sufficient

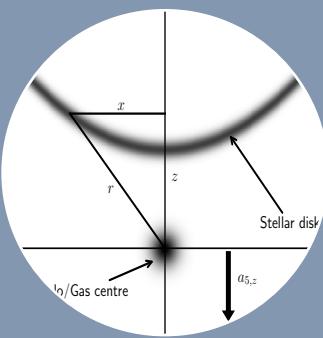
Gas-star offset



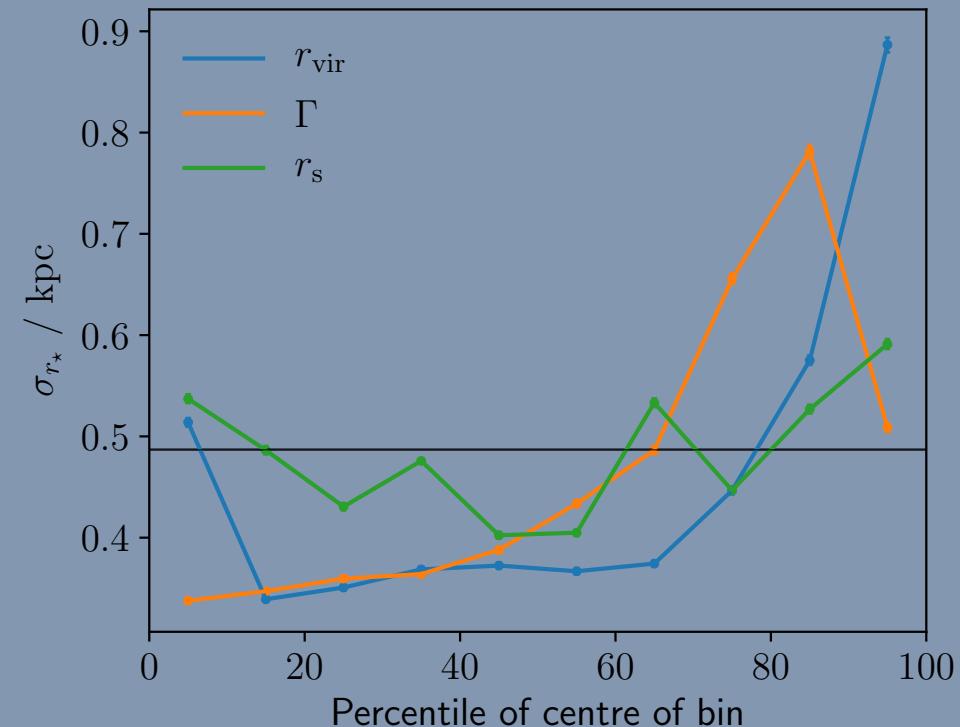
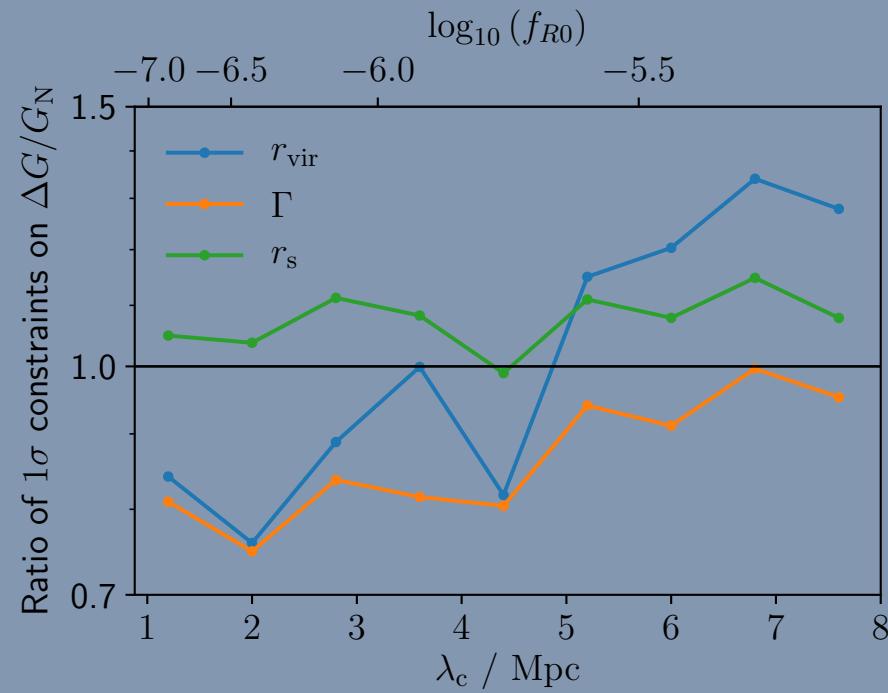
Can predict → Need to do more

Bartlett, Desmond & Ferreira 2021 (arXiv:2103.10356)

## 2. How significant are these correlations?

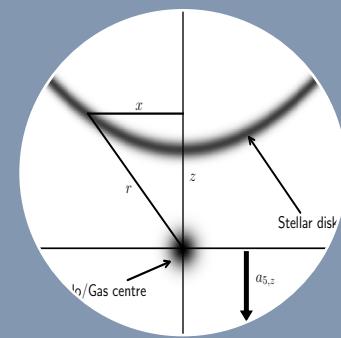


- Bin galaxies based on properties
- Assign different noise parameters to each bin

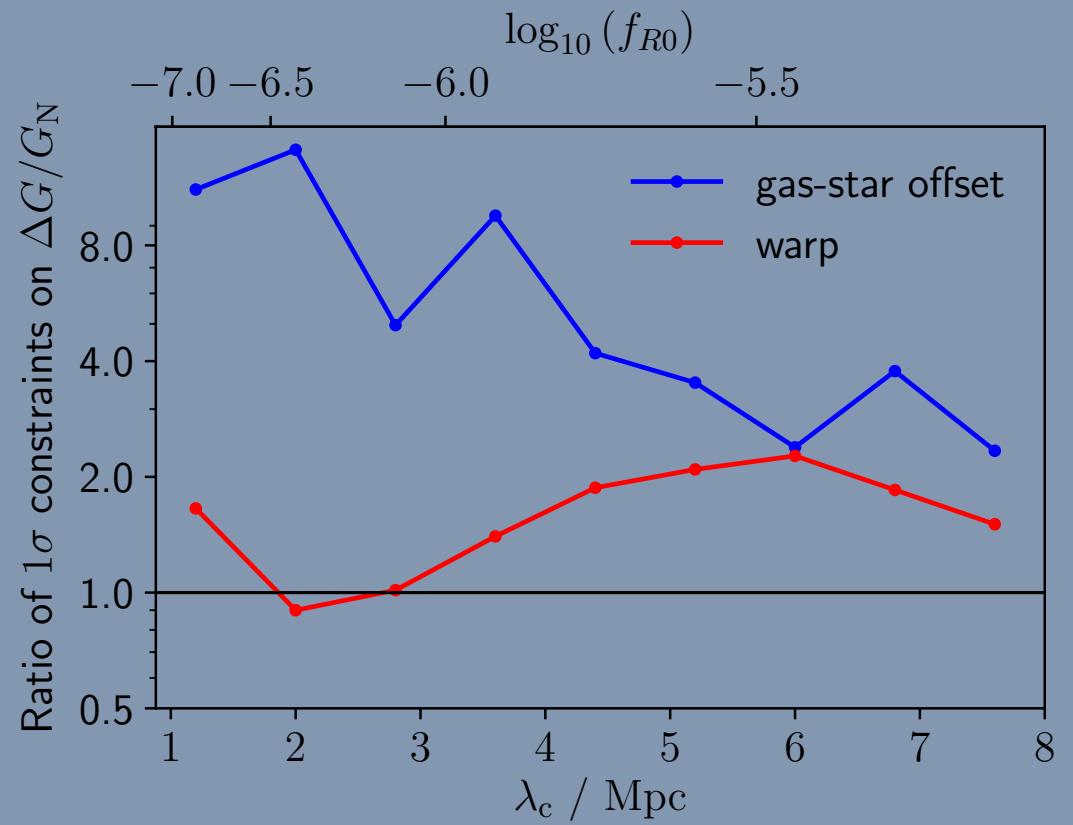


Bartlett, Desmond & Ferreira 2021 (arXiv:2103.10356)

### 3. Is my model for $\bar{y}_i$ sufficient?

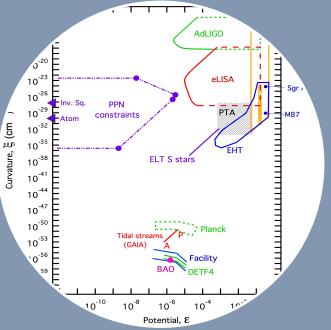


- With real data, had to assume the inner power-law slope of halo density profile
- We can measure this in the simulation
- Gas-star offsets far more sensitive than warps
- Conclusions depend on how much you trust your simulation

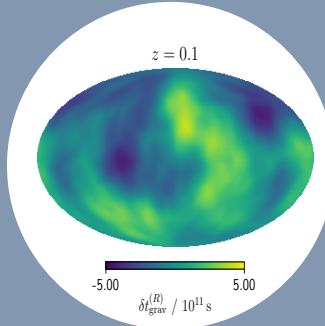


Bartlett, Desmond & Ferreira 2021 (arXiv:2103.10356)

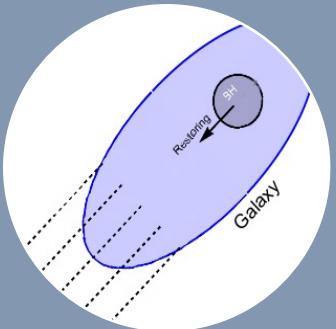
# Summary



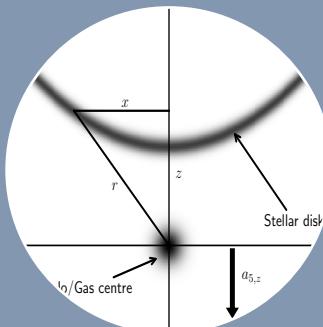
- Under-explored region of parameter space
- Lots of data



- Importance of noise
- $\Delta\gamma$  constraint 30 times stronger than previous attempts



- Tight constraints on fifth-force
- Improve non-linear modeling



- Astrophysics messy – should test your noise model