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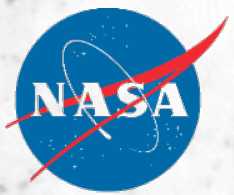
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How to calibrate your shear measurements

Eric Huff (JPL)

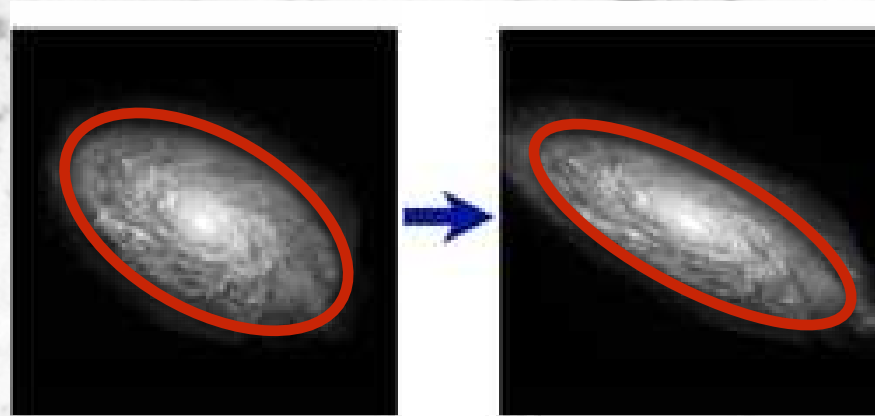
Rachel Mandelbaum (CMU), Erin Sheldon (BNL)

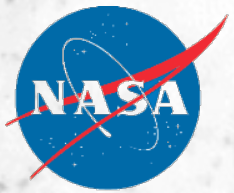


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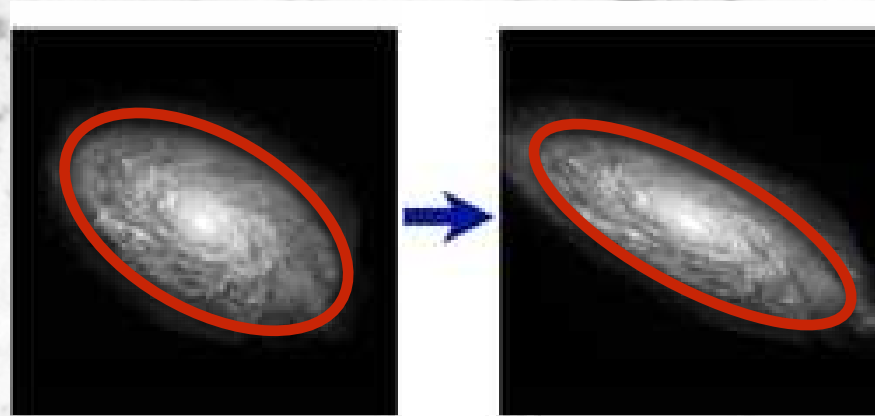
How do we get g from a galaxy image?



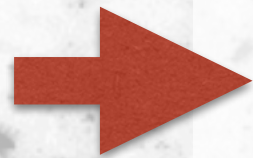


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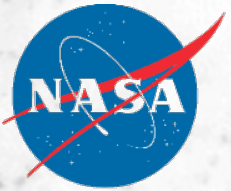
How do we get g from a galaxy image?



**ellipticity
suggests
2nd moments**



$$Q_{ij\dots k} = \int I(\theta) \theta_i \theta_j \dots \theta_k d^2\theta .$$
$$\chi = \frac{(Q_{11} - Q_{22}) + 2iQ_{12}}{Q_{11} + Q_{22}}$$

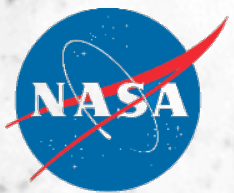


Kaiser, Squires, Broadhurst (1995)

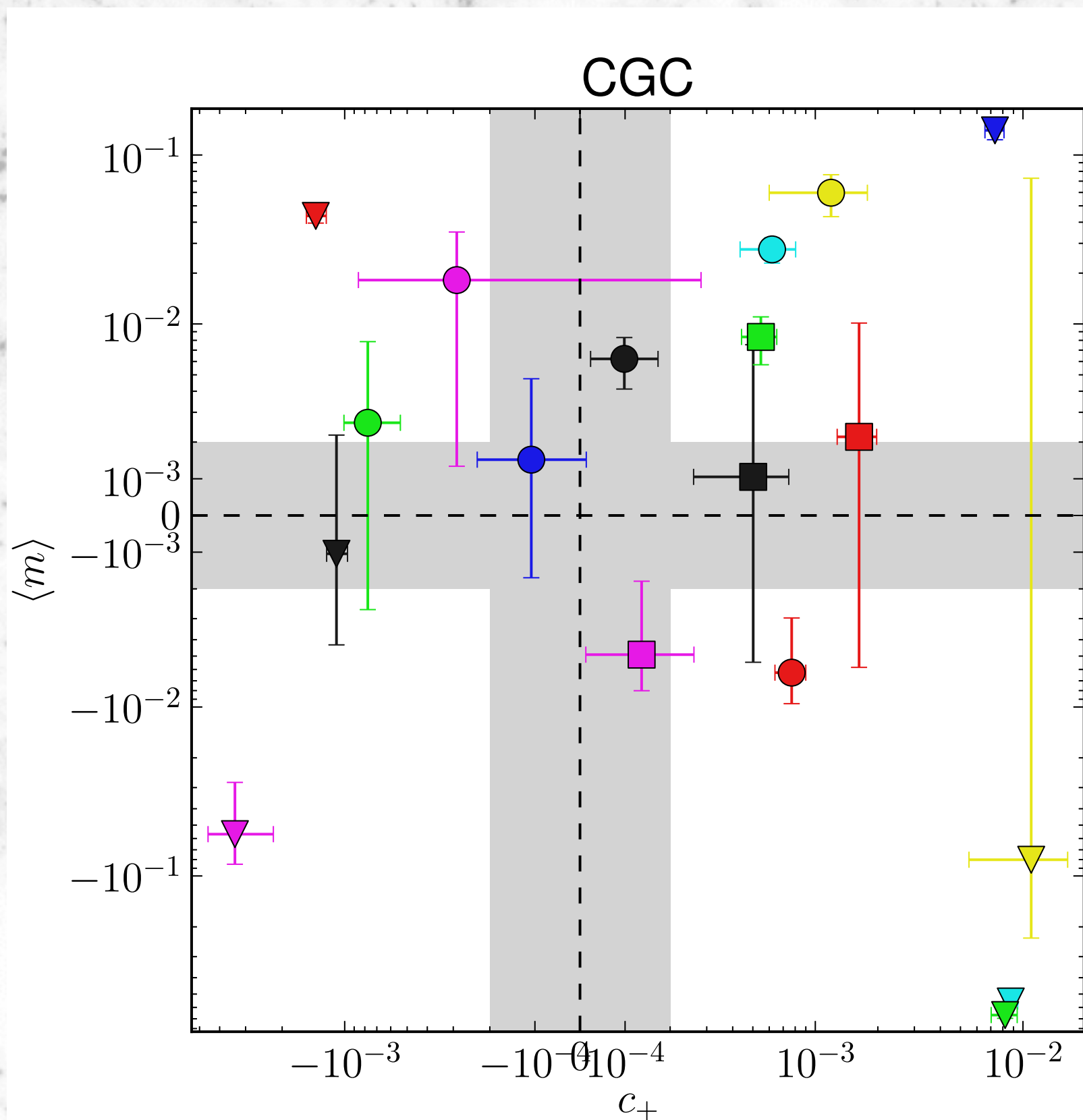
- 1. Compute second moments.**
- 2. Calculate the responses to shear (P_g) and PSF ellipticity (P^{sm}).**
- 3. Correct for PSF ellipticity (e^*).**

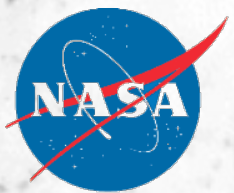
$$g = P_g^{-1} \left(e^{\text{obs}} - \frac{P^{\text{sm}}}{P^{\text{sm}*}} e^* \right)$$

This doesn't seem that bad.

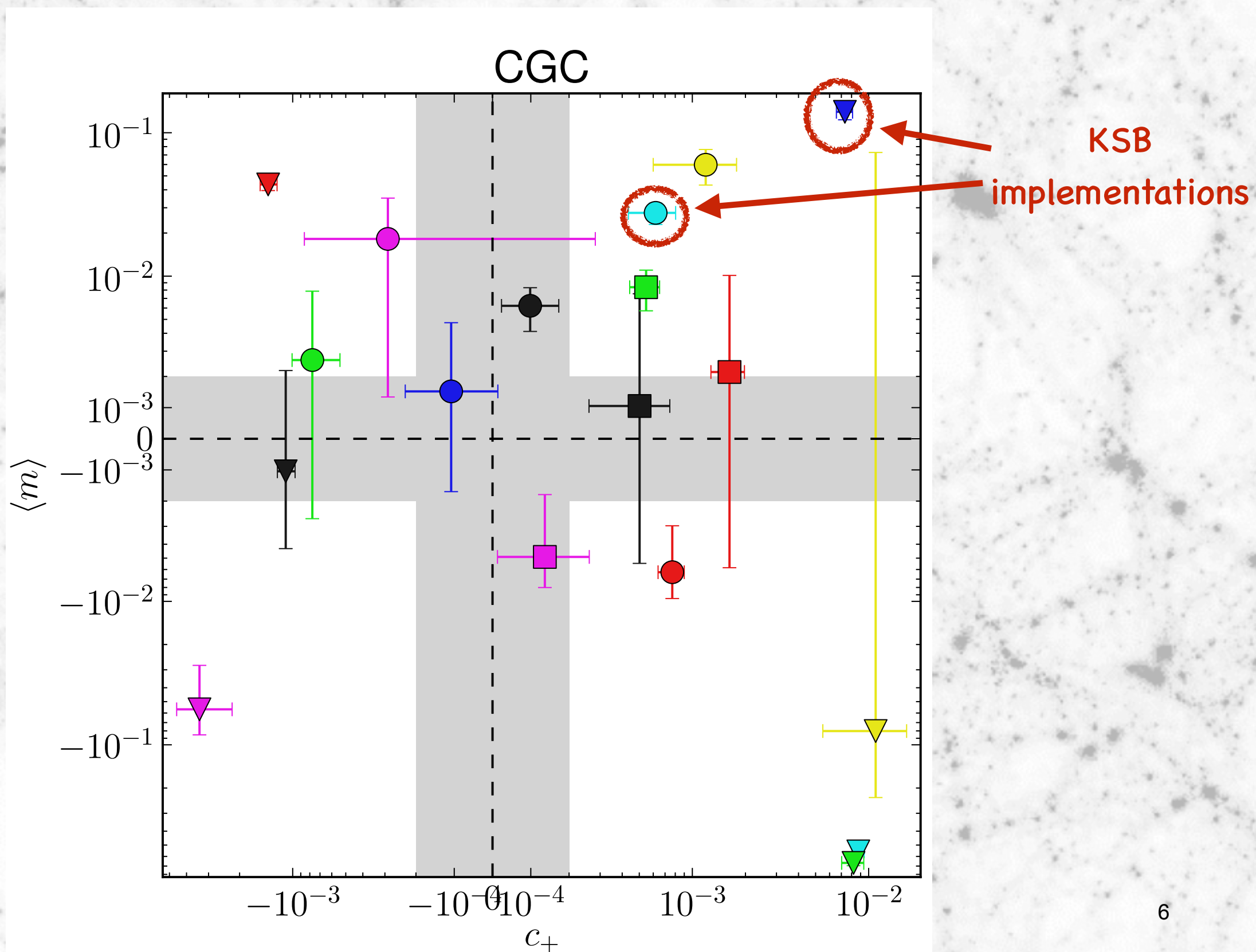


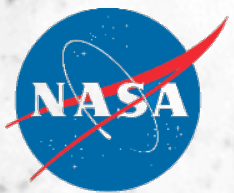
Broad dissensus in blind community challenges



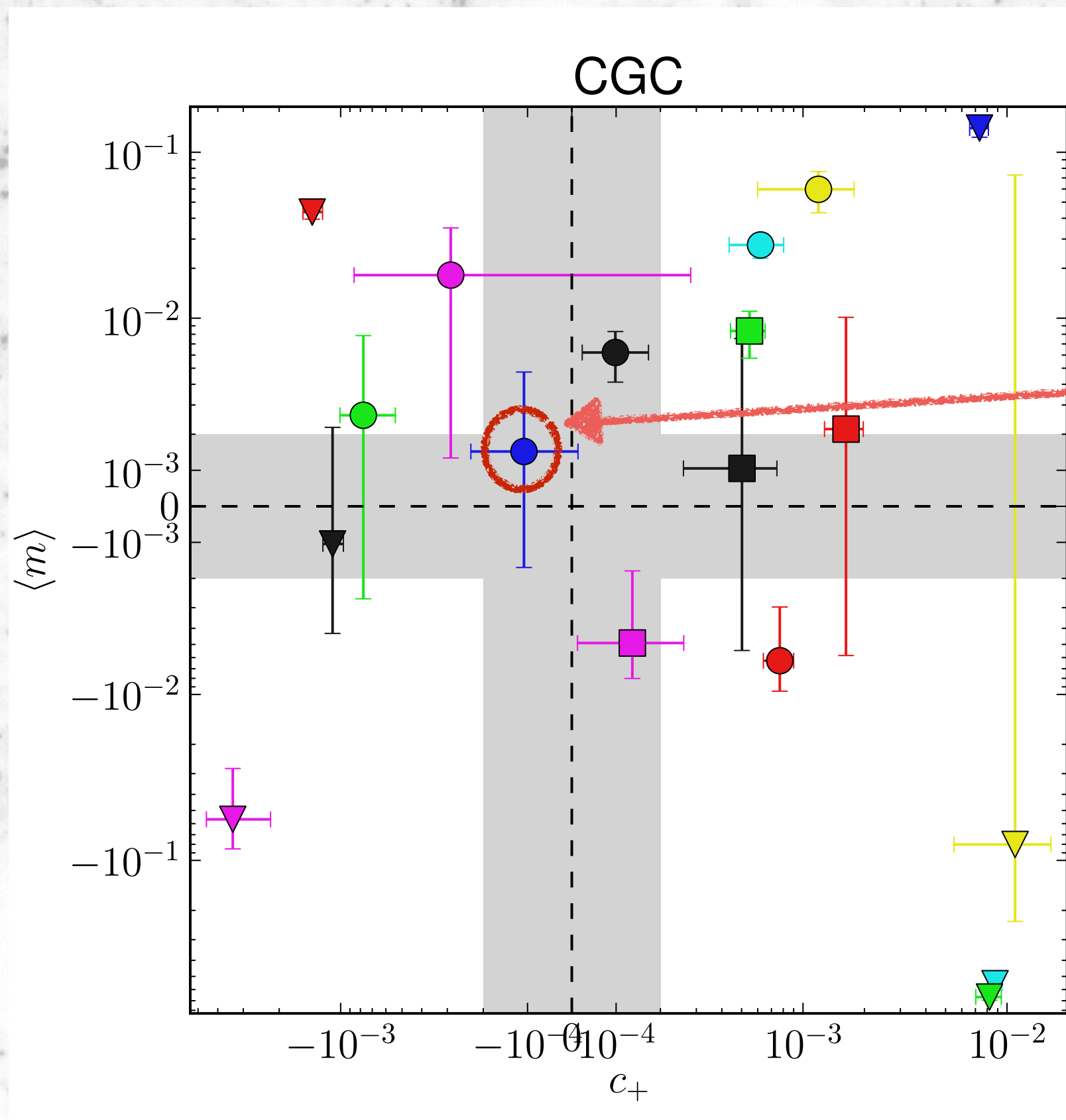


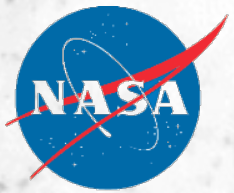
Broad dissensus in blind community challenges





Broad dissensus in blind community challenges

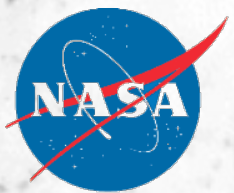




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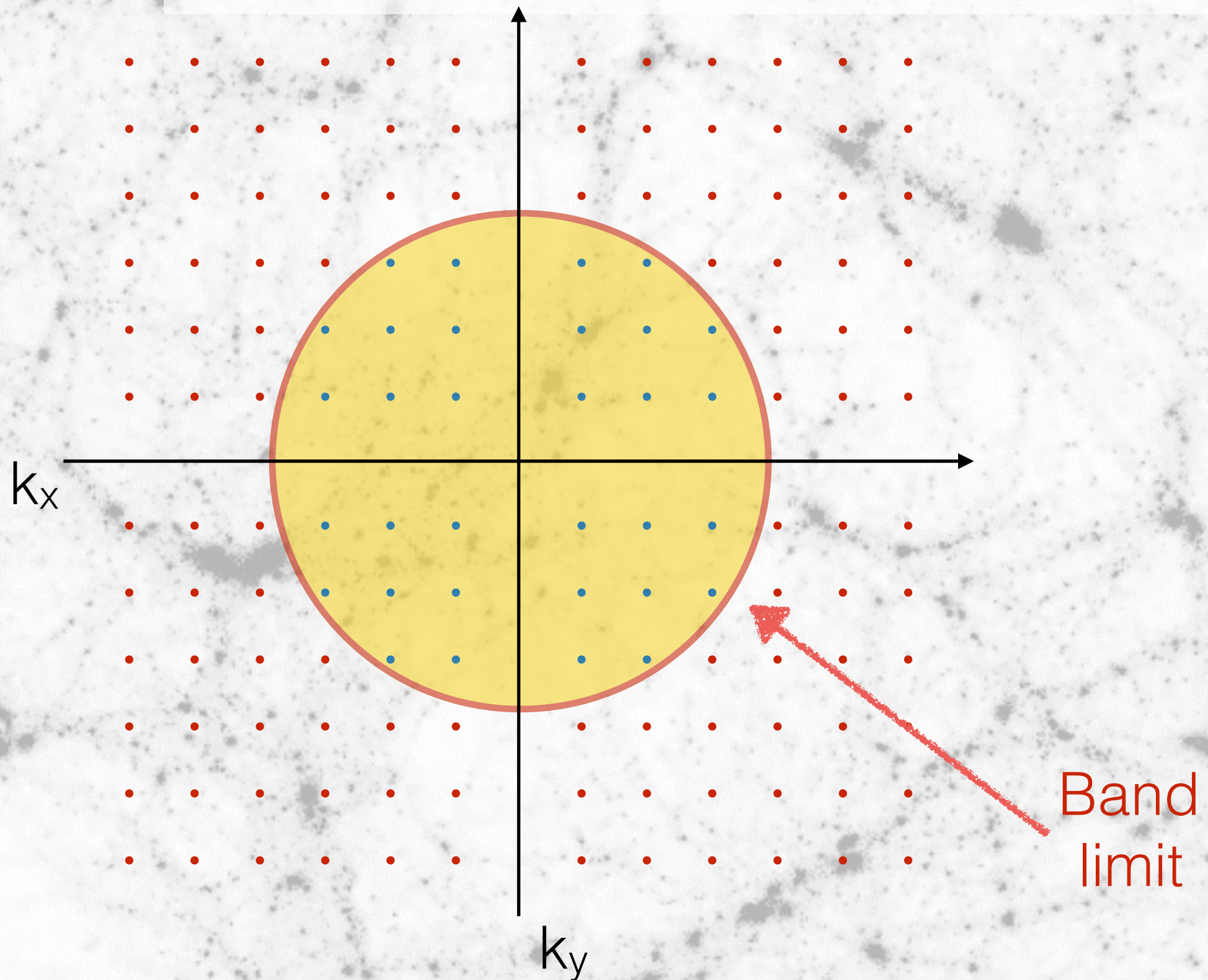
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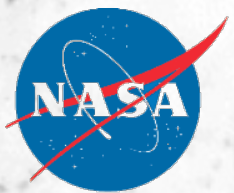
Why is this hard?



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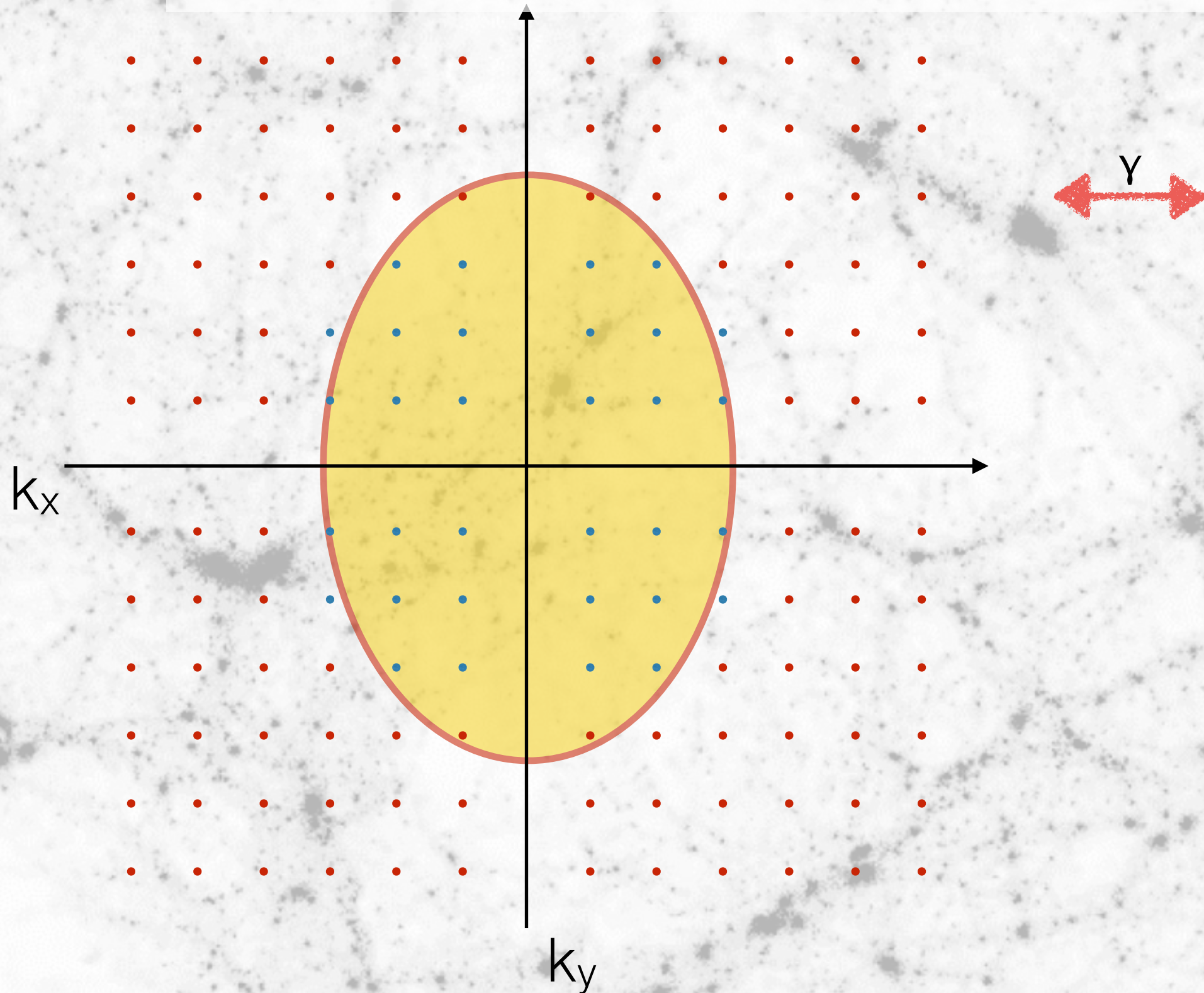
the full shear response depends on unresolved modes

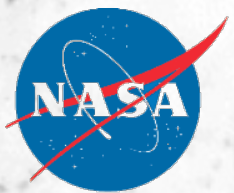




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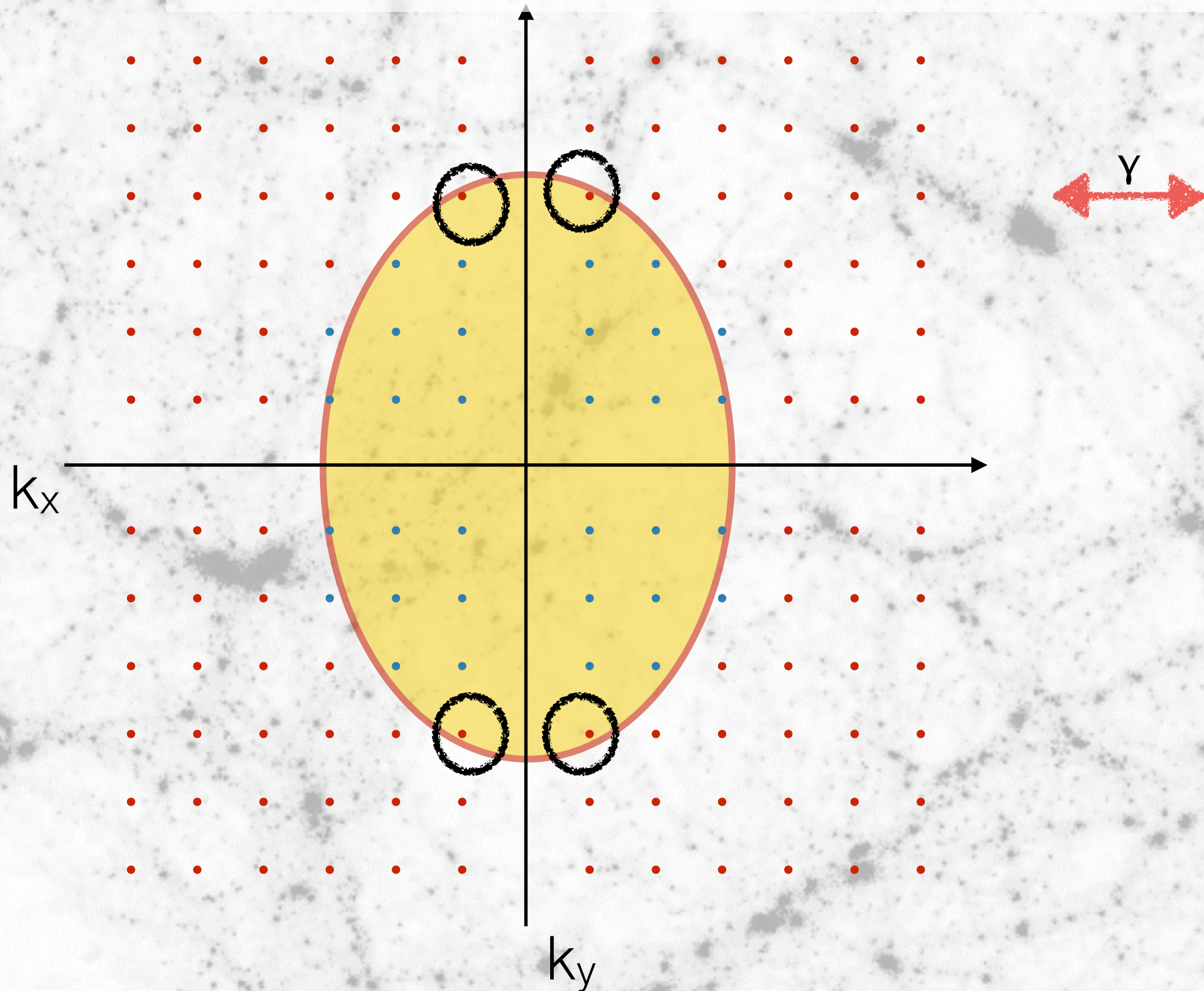
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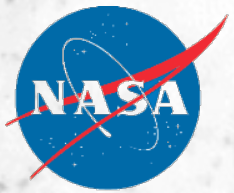




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the full shear response depends on unresolved modes



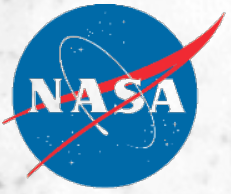


Selection biases matter

$$\begin{aligned}\langle R \rangle &= \int \frac{\partial S(e)P(e)e}{\partial \gamma} \Big|_{\gamma=0} de \\ &= \int \left[S(e) \frac{\partial P(e)e}{\partial \gamma} \Big|_{\gamma=0} + P(e)e \frac{\partial S(e)}{\partial \gamma} \Big|_{\gamma=0} \right] de\end{aligned}$$



**Shear selection biases:
calibration depends on sub-threshold
galaxy population**



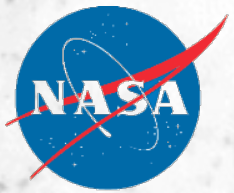
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The measurement process is complicated.

**but the images are simple
(linear)**

This suggests the following procedure:

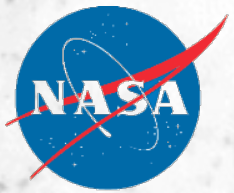


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Construct counterfactual images

$$I'(\mathbf{x}|\mathbf{g}) = P * (\hat{\mathbf{s}}_{\mathbf{g}}G)$$

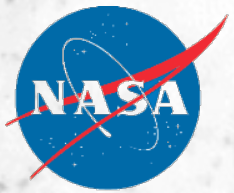


Construct counterfactual images

$$I'(\mathbf{x}|\mathbf{g}) = P * (\hat{\mathbf{s}}_{\mathbf{g}}G)$$

$$I'(x|g) = \Gamma * \left[\hat{s}_g(P^{-1} * I) \right]$$

remove the PSF, shear, and add a new PSF



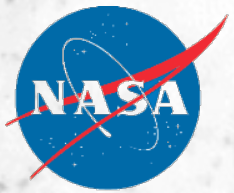
Construct counterfactual images

$$I'(\mathbf{x}|\mathbf{g}) = P * (\hat{\mathbf{s}}_{\mathbf{g}}G)$$

$$I'(x|g) = \Gamma * [\hat{s}_g(P^{-1} * I)]$$



we get to choose our final PSF



Construct counterfactual images

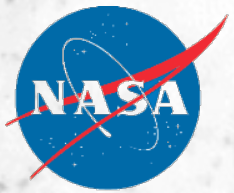
$$I'(\mathbf{x}|\mathbf{g}) = P * (\hat{\mathbf{s}}_{\mathbf{g}}G)$$

$$I'(x|g) = \Gamma * [\hat{s}_g(P^{-1} * I)]$$

$$e^+ = \hat{E} \{ I'(x|g^+) \}$$



any ~linear measurement algorithm



Construct counterfactual images

$$I'(\mathbf{x}|\mathbf{g}) = P * (\hat{\mathbf{s}}_{\mathbf{g}}G)$$

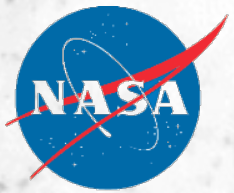
$$I'(x|g) = \Gamma * [\hat{s}_g(P^{-1} * I)]$$

$$e^+ = \hat{E} \{ I'(x|g^+) \}$$

$$1 + m = \frac{e^+ - e^-}{2\Delta g}$$

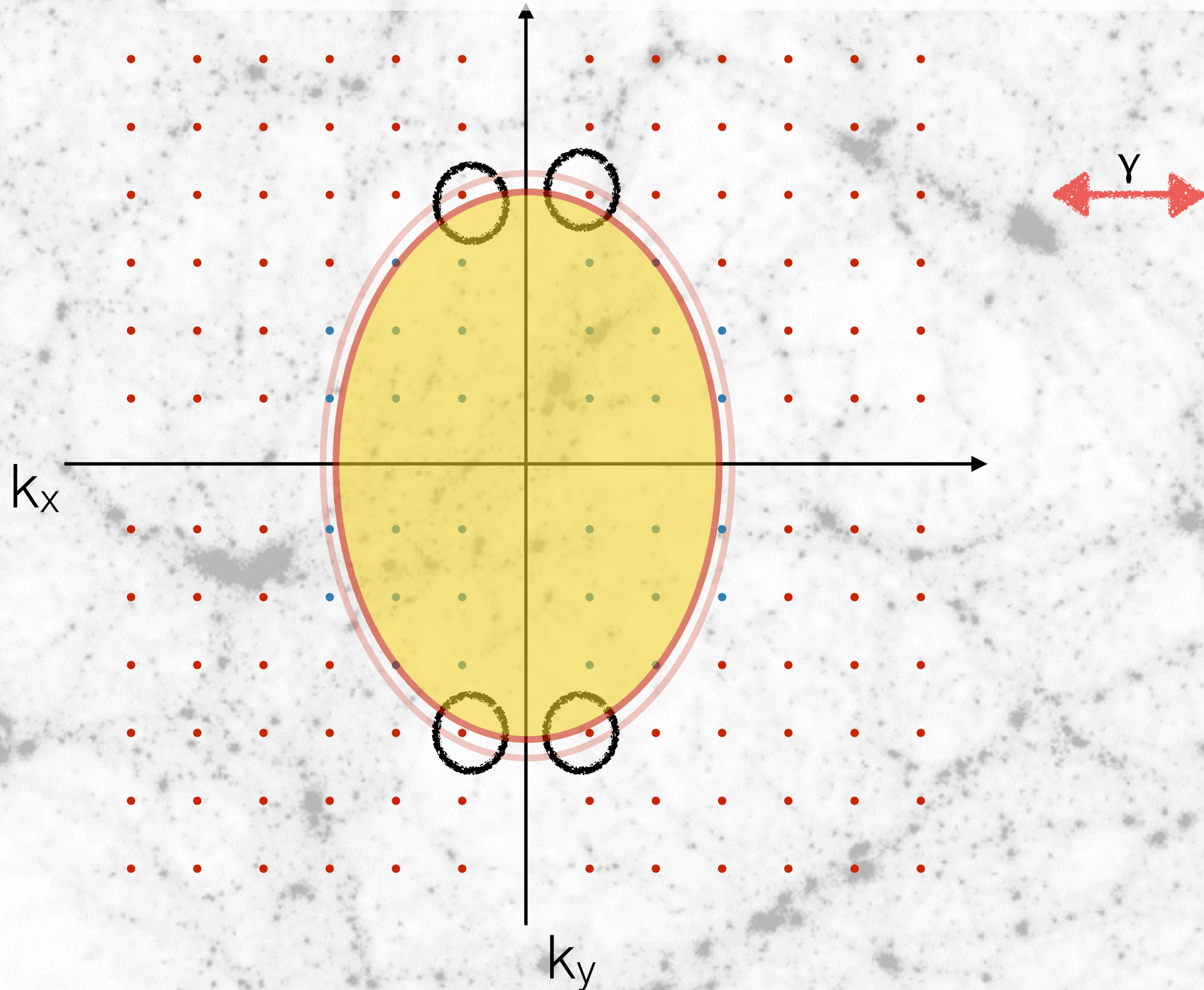


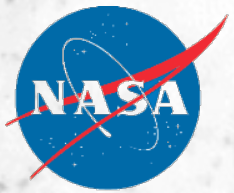
repeated measurement on counterfactuals



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smear the data further to hide the band-limited modes

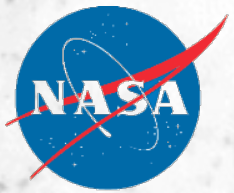




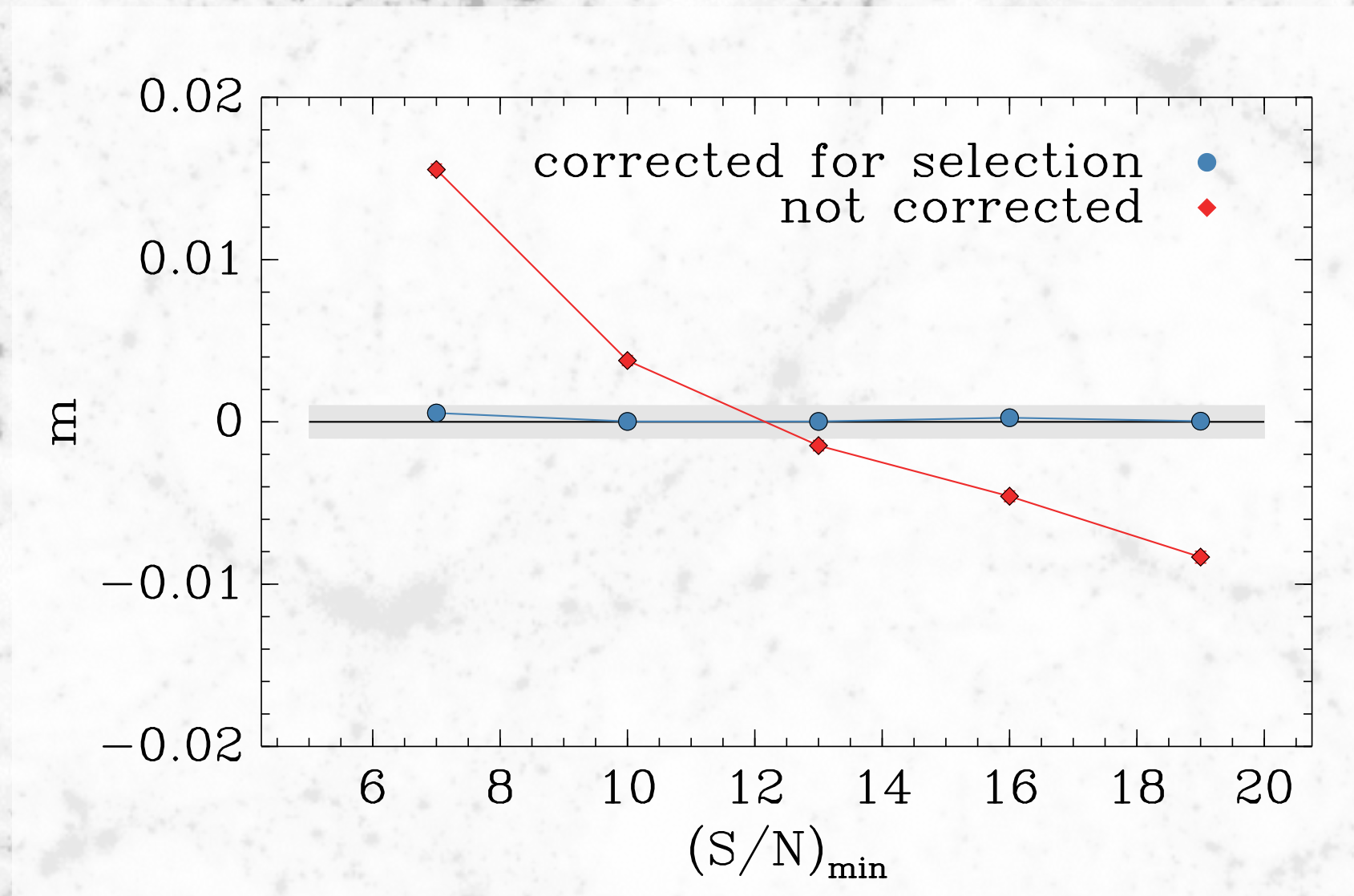
Correcting for selection effects:

$$\begin{aligned}\langle R \rangle &= \int \frac{\partial S(e) P(e) e}{\partial \gamma} \Big|_{\gamma=0} de \\ &= \int \left[S(e) \frac{\partial P(e) e}{\partial \gamma} \Big|_{\gamma=0} + P(e) e \frac{\partial S(e)}{\partial \gamma} \Big|_{\gamma=0} \right] de \\ &= \frac{\langle e_i^+ \rangle^S - \langle e_i^- \rangle^S}{\Delta \gamma_j} + \frac{\langle e_i \rangle^{S+} - \langle e_i \rangle^{S-}}{\Delta \gamma_j}\end{aligned}$$

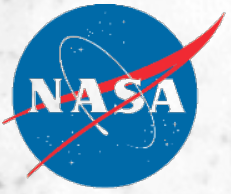
Apply a shear.
See how your measured shapes change.



**Selection effects are large,
but now effectively mitigated.**



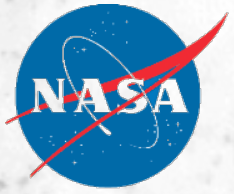
**There is no evidence
for any remaining calibration bias.**



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**This space
intentionally left blank.**



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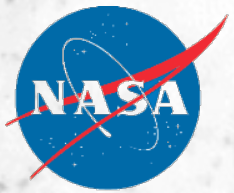
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There will be blending



Image courtesy of / stolen from
Peter Melchior

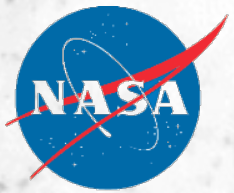


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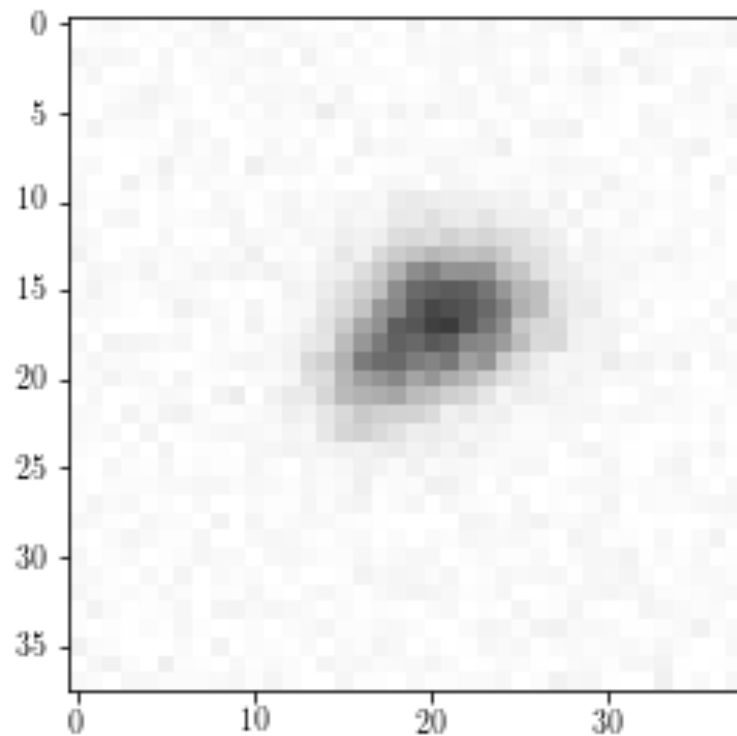
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summary of blending effects:

- 1. Reducing overall number density**
- 2. Mixing shear across redshifts**
- 3. Photo-z mis-estimation**
- 4. Density-dependent selection**



Consider the effects of blending on shear



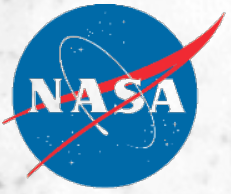
**Single blob,
imperfectly deblended,
2 or more galaxies**

$$R_1 = \frac{\partial e}{\partial \gamma_1} \quad R_2 = \frac{\partial e}{\partial \gamma_2}$$

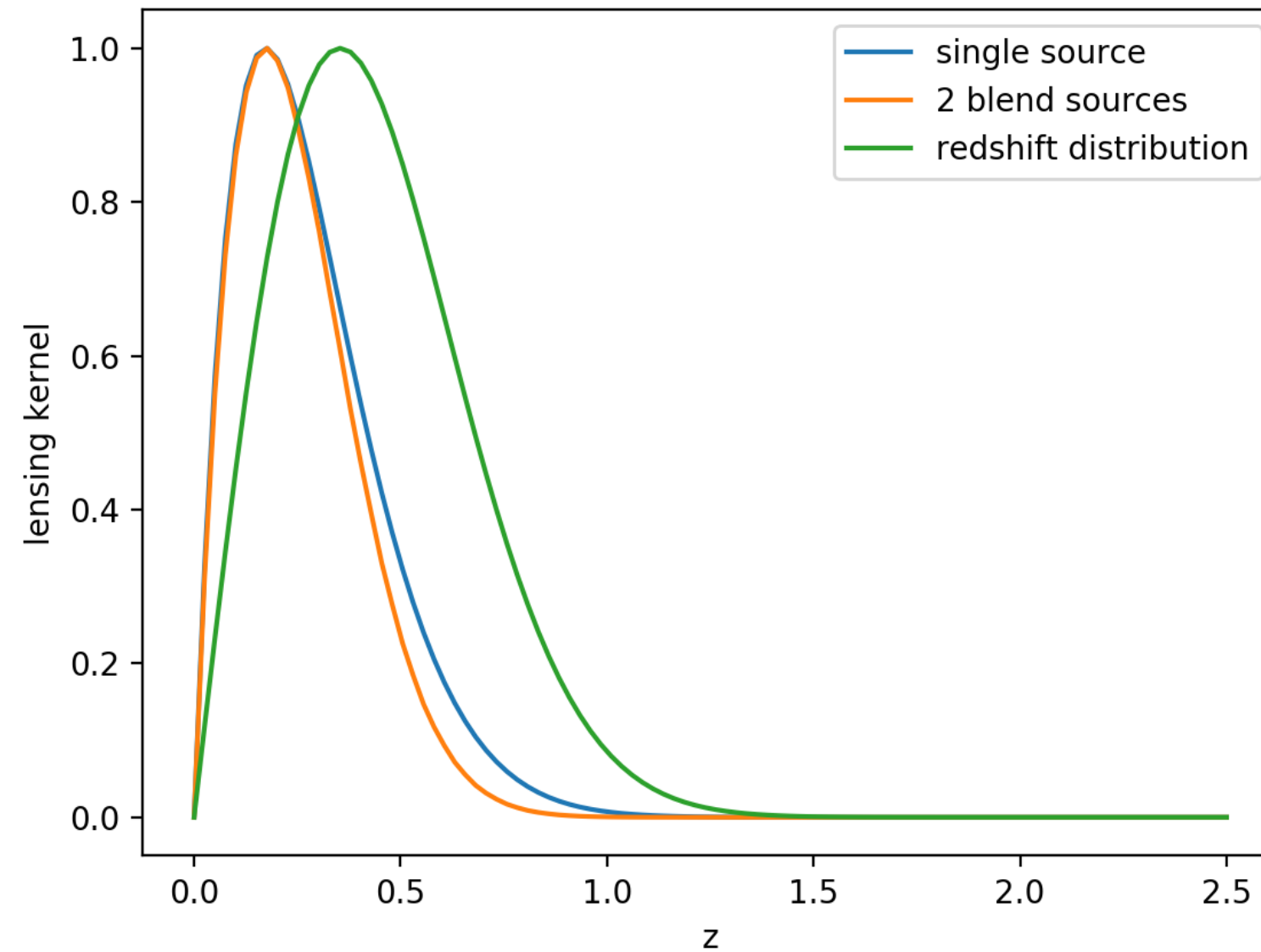
$$e_{\text{blob}} = R_1 \gamma_1 + R_2 \gamma_2$$

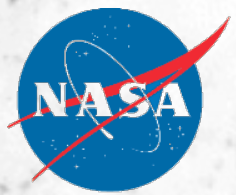
$$\langle e_i, e_{\text{blob}} \rangle = R_1 \langle e_i, \gamma_1 \rangle + R_2 \langle e_i, \gamma_2 \rangle$$

**Then we can write
effects on 2pt correlations**



Blending and shear

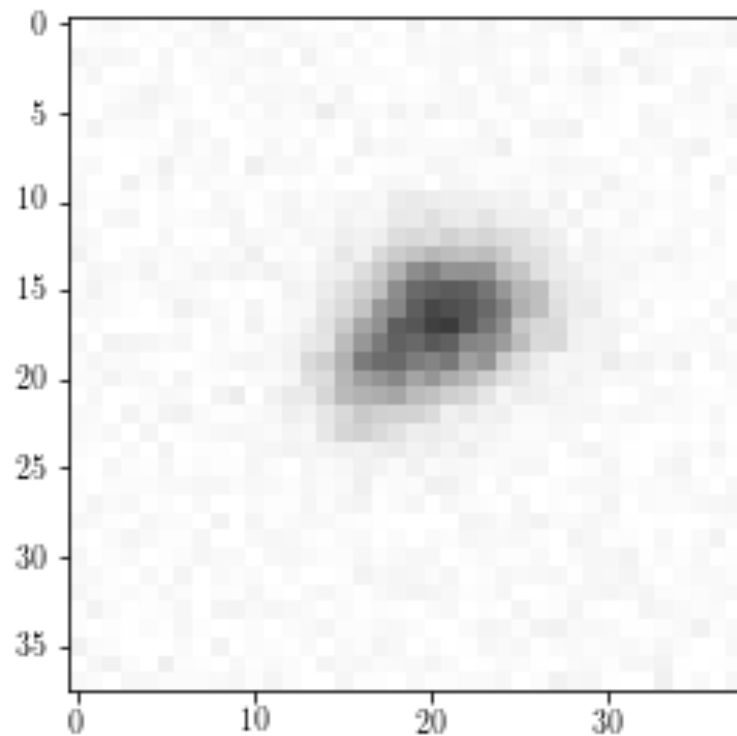




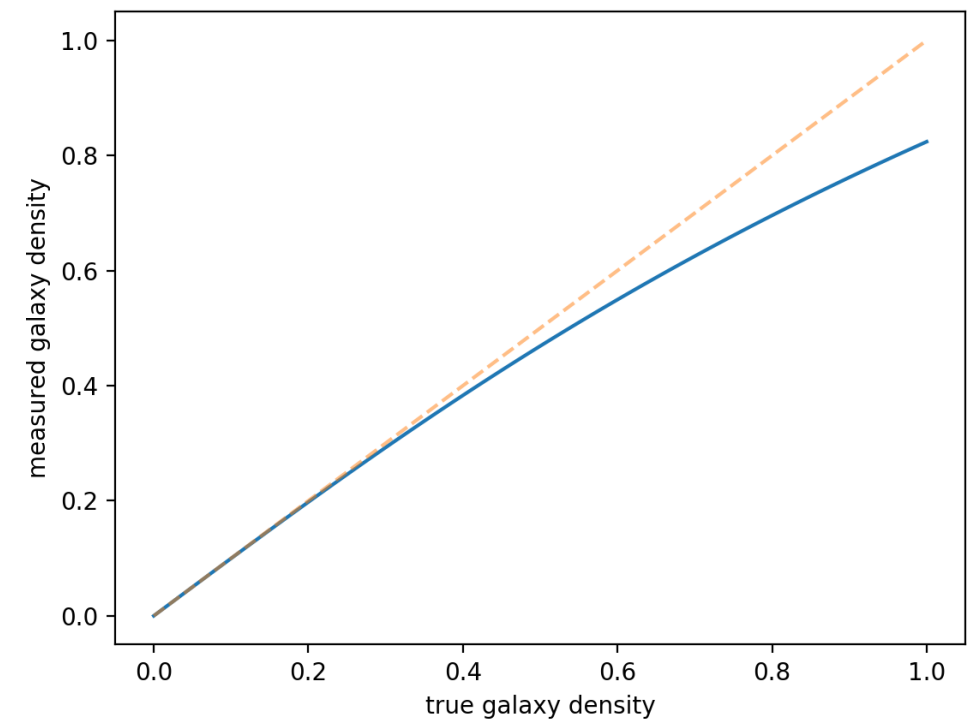
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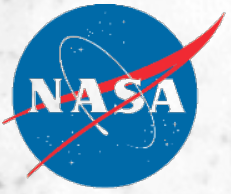
Consider the effects of blending on clustering



**Single blob,
imperfectly deblended,
2 or more galaxies**



**blending reduces completeness,
preferentially in overdense regions**



Impact on galaxy clustering

Lenses $z=0.45-0.55$, $m_z < 22$

Model for the impact of the neighbor
bias on the 2PCF:

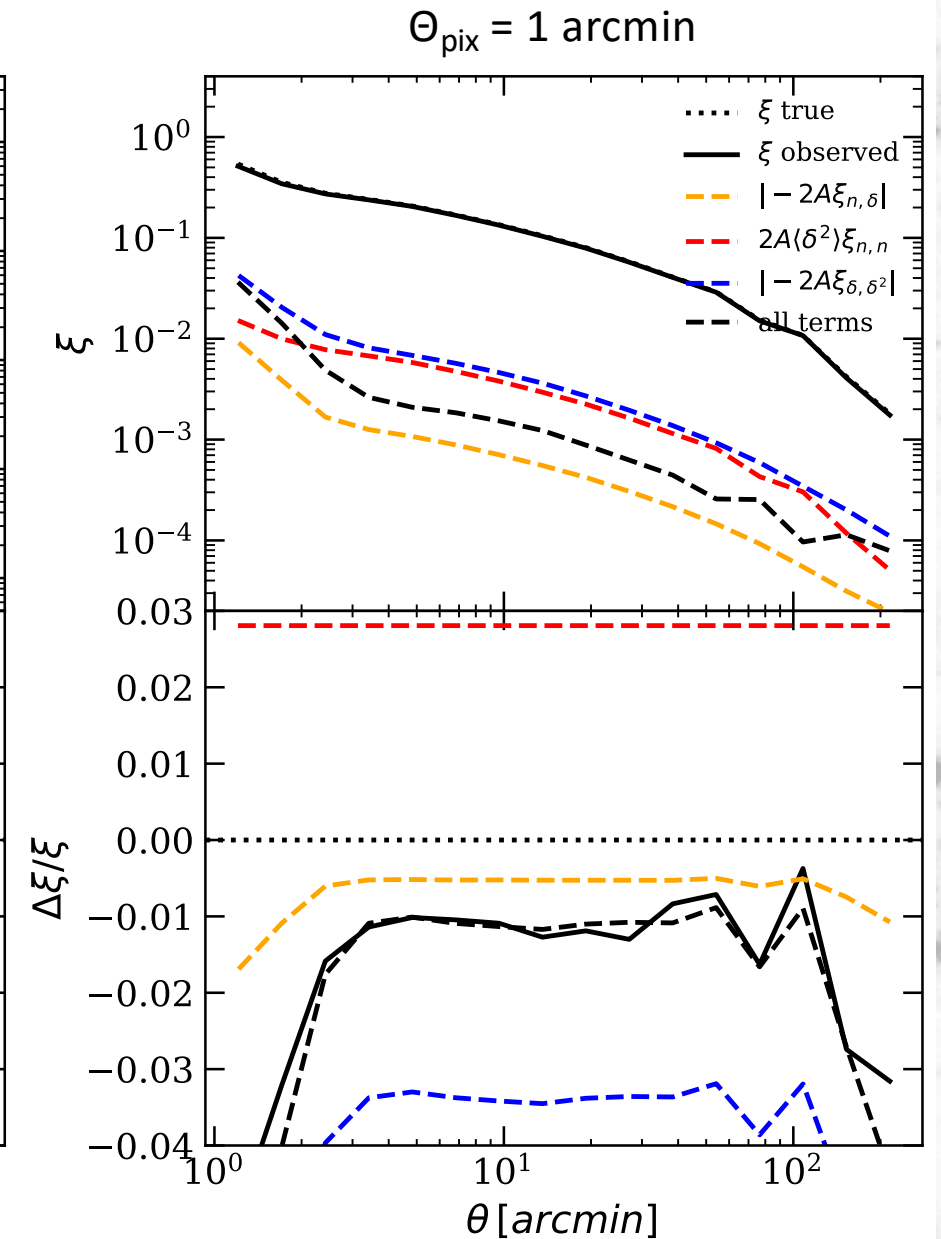
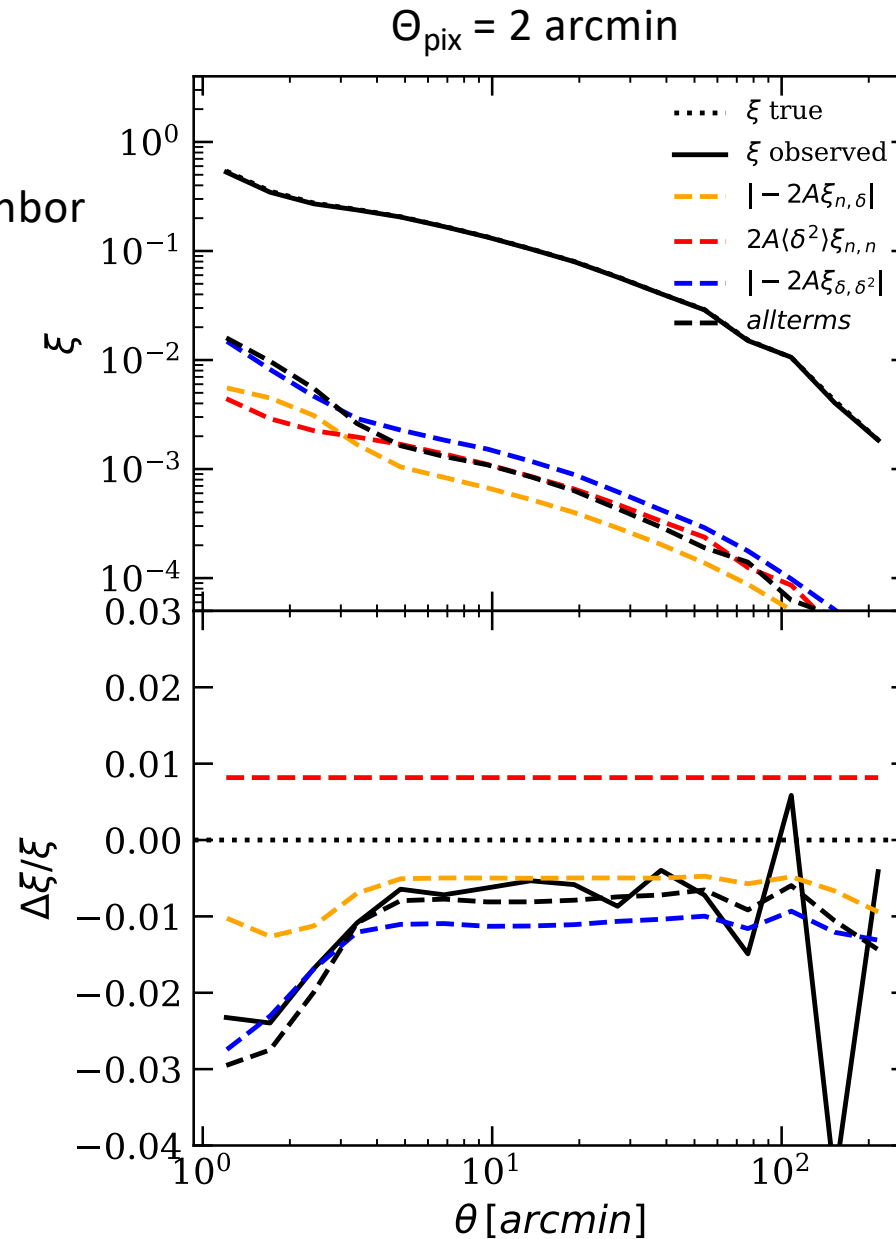
$$n_{\text{obs}} = \tilde{n}_t(1+\delta) W_s (1-\Gamma\delta)$$

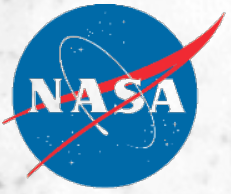
...

to first order in Γ

...

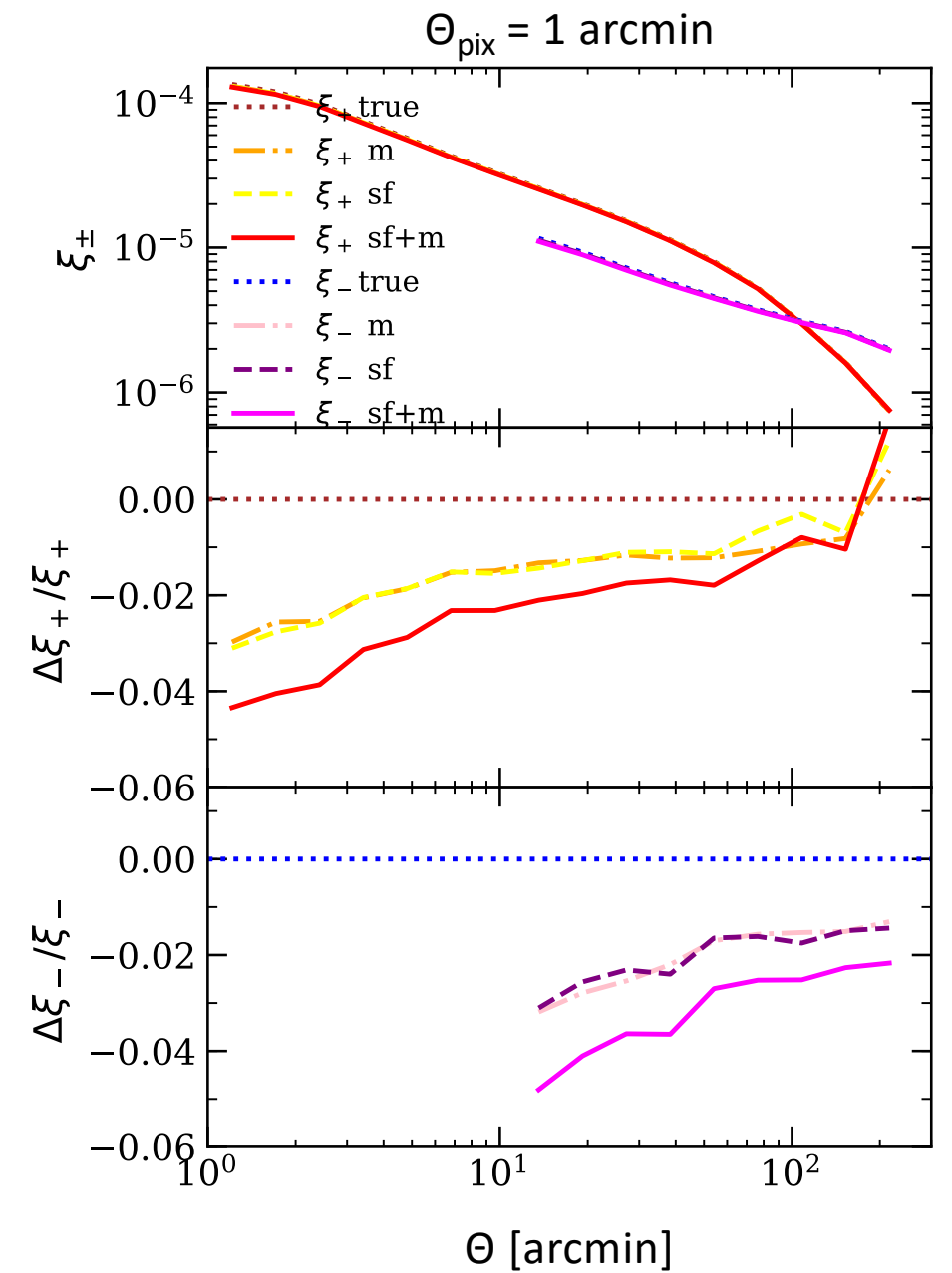
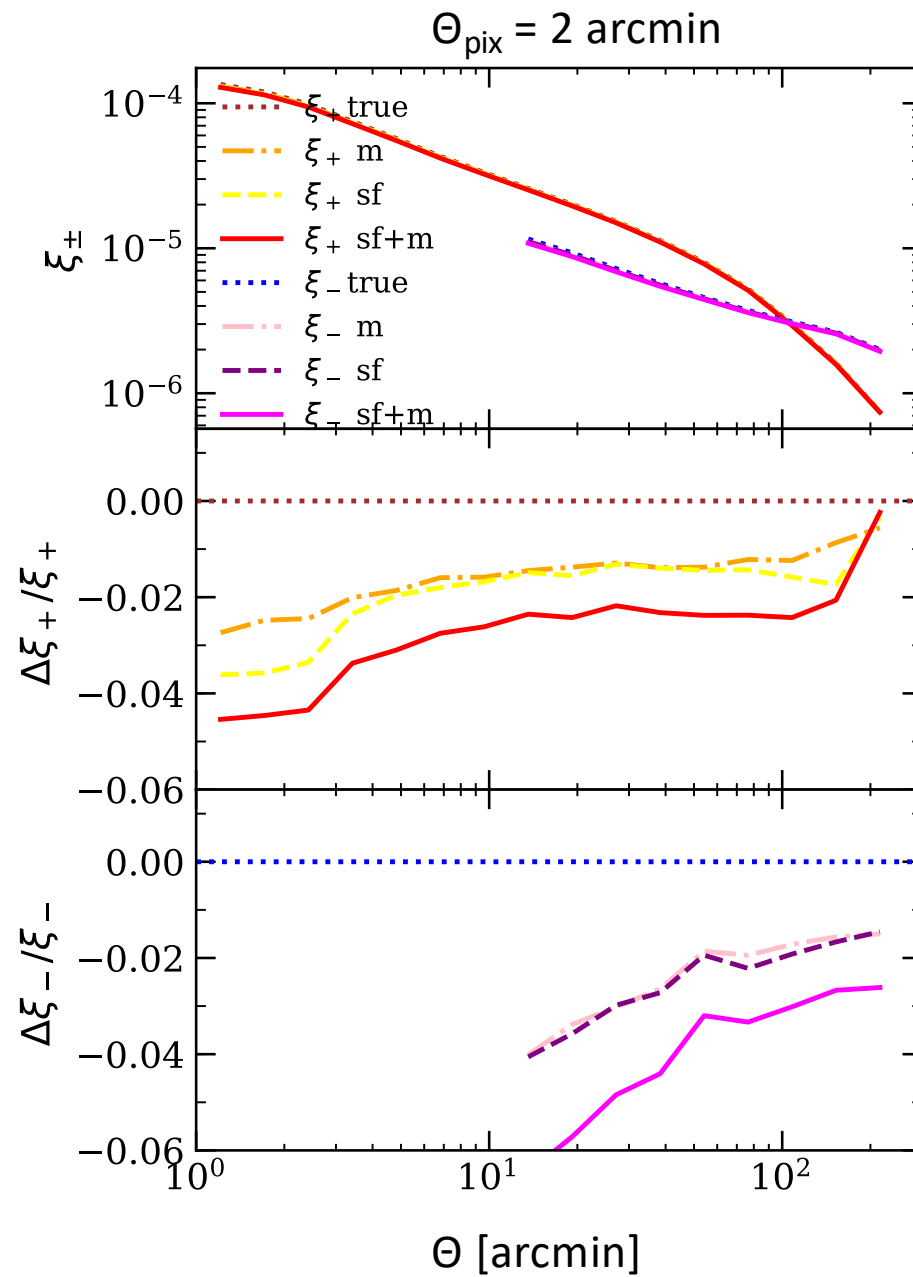
$$\langle n_o n_o' \rangle / \tilde{n}_o^2 - 1 = [1 + 2\Gamma(\langle \delta^2 \rangle - 1)]\xi - 2\Gamma\langle \delta\delta'^2 \rangle$$

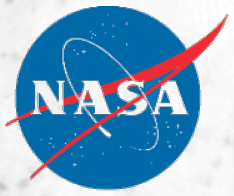




Impact on cosmic shear

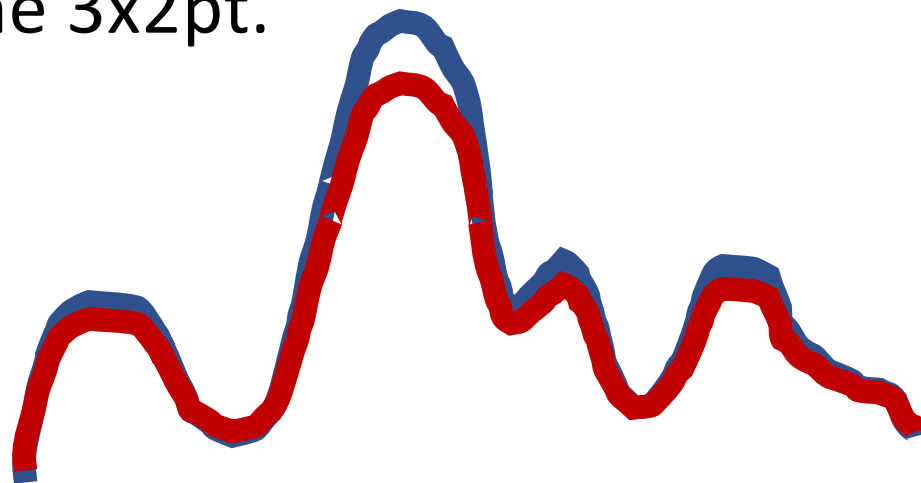
Sources $z=0.95-1.05$, $m_z < 24$

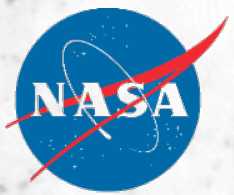




Summary

- We developed a general framework to capture any systematic that depends on the local number density of galaxies.
- It is a single parameter model: θ_{excl}
- All this is captured by Γ : $E \sim \Gamma \delta$ (Γ is related to θ_{excl})
- We have an analytic model for the impact on the 2PCF (it works!)
- Doing the same for all the 3x2pt.





Gratuitous omphaloskepsis

