Analytic covariance of the redshift-space galaxy power spectrum using PT

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with Roman Scoccimarro (arXiv 1910.02914)



UC Berkeley, Oct 2019

Motivation: Power Spectrum Covariance • We estimate cosmological parameters (θ) from galaxy surveys using

the power spectrum

$$p(P_d|\theta) = \frac{1}{(2\pi)^{n/2}\sqrt{\det \mathbf{C}}} \exp\left[-\frac{1}{2}(P_d - P(\theta))^T \mathbf{C}^{-1}(P_d - P(\theta))\right]$$

$$\boldsymbol{C}(\boldsymbol{k_1},\boldsymbol{k_2}) = \langle \hat{P}(k_1)\hat{P}(k_2)\rangle - \langle \hat{P}(k_1)\rangle\langle \hat{P}(k_2)\rangle$$

0 estimation

Robust knowledge of the covariance matrix is crucial for parameter

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Covariance from mock catalogs

Need to simulate mock surveys (~ thousands) to compute the covariance matrix





 $C_{1,2} = \frac{1}{N_{\text{mock}} - 1} \sum_{i=1}^{N_{\text{mock}}} (\hat{P}_i(k_1) - \bar{P}(k_1)) (\hat{P}_i(k_2) - \bar{P}(k_2))$

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Covariance from mock catalogs

Need to simulate mock surveys (~ thousands) to compute the covariance matrix







Covariance from mock catalogs

As survey volume increases, mock catalogs become tougher to simulate (LSST, DESI, Euclid and others) $C_{1,2} = \frac{1}{N_{\text{mock}} - 1} \sum_{i=1}^{N_{\text{mock}}} (\hat{P}_i(k_1) - \bar{P}(k_1))(\hat{P}_i(k_2) - \bar{P}(k_2))$

Mocks suffer from sampling noise (error $\propto 1/\sqrt{N_{mocks}}$) (noise can make matrix harder to invert)

$$p(P_d|\theta) = \frac{1}{(2\pi)^{n/2}\sqrt{\det C(\theta)}} \exp(\frac{1}{(2\pi)^{n/2}} \sqrt{\det C(\theta)})$$

Dependence of covariance on cosmology and bias parameters is computationally prohibitive

 $\exp\left[-\frac{1}{2}(P_d - P(\theta))^T C(\theta)^{-1}(P_d - P(\theta))\right]$

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What are the challenges for an analytic method?

Challenge I: Highly non-trivial survey window

 $\delta_W(\mathbf{x}) \equiv W(\mathbf{x})\delta(\mathbf{x})$



 \mathcal{Z}







Survey window enters Covariance

$$C(k_1, k_2) = \langle \hat{P}(k_1) \hat{P}(k_2) \rangle - \langle \hat{P}(k_1) \rangle \langle \hat{P}(k_2) \rangle$$

$$\begin{split} \langle \hat{P}(k_1)\hat{P}(k_2)\rangle = &\frac{1}{V_2^2} \int_{\hat{\mathbf{k}}_1, \hat{\mathbf{k}}_2} \langle \boldsymbol{\delta}_{\boldsymbol{W}}(\mathbf{k}_1)\boldsymbol{\delta}_{\boldsymbol{W}}(-\mathbf{k}_1) \\ = &\frac{1}{V^2} \int_{\hat{\mathbf{k}}_1, \hat{\mathbf{k}}_2, \mathbf{p}_1, \mathbf{p}_1', \mathbf{p}_2, \mathbf{p}_2'} W(\mathbf{k}_1) \\ & \times \langle \boldsymbol{\delta}(\mathbf{p}_1) \rangle \end{split}$$

18 dimensional integral



- $_{1})\delta_{W}(\mathbf{k}_{2})\delta_{W}(-\mathbf{k}_{2})
 angle$
- $egin{aligned} \mathbf{k}_1 \mathbf{p}_1) & W(-\mathbf{k}_1 \mathbf{p}_1') W(\mathbf{k}_2 \mathbf{p}_2) W(-\mathbf{k}_2 \mathbf{p}_2') \ & \mathbf{p}_1) \delta(\mathbf{p}_1') \delta(\mathbf{p}_2) \delta(\mathbf{p}_2')
 angle \end{aligned}$



Challenge II: 4-point function (Trispectrum)

Need to calculate non-linear 0 structure formation in PT

- Need to model: 0
- RSD
- Shot Noise 2.
- Bias (Linear, non-linear and non-local) 3.
- Effect of Super-Survey modes 4.



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$$\langle k_1, k_2 \rangle = \langle \delta(k_1)\delta(-k_1)\delta(k_2)\delta(-k_2) \rangle$$

- $\langle \delta(k_1)\delta(-k_1) \rangle \langle \delta(k_2)\delta(-k_2) \rangle$

$$C^{\mathbf{G}}(k_1, k_2) \simeq 2 \left\langle \delta(k_1) \delta(-k_2) \right\rangle \left\langle \delta(k_2) \delta(-k_2) \right\rangle$$

$$N^{\mathbf{G}}(k_1, k_2) = \left\langle \delta(k_1) \delta(-k_2) \delta(k_2) \delta(-k_1) \right\rangle$$



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Is it necessary to model the NG part?

NG part causes 40% leakage of info. at k=0.3 h/Mpc







Diagonal auto-covariance

--- Our analytic method

Patchy Mocks
 (state-of-the-art mocks used for
 SDSS BOSS parameter estimation)

DW & Scoccimarro 19

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Diagonal auto-covariance

 Patchy Mocks (~ MONTHS) (state-of-the-art mocks used for SDSS BOSS parameter estimation)

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Off-diagonal auto-covariance

(2 rows compared in fig)

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Off-diagonal cross-covariance

High sampling noise in mocks

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I. Gaussian Covariance

 $C(k_1, k_2) = \langle \hat{P}(k_1) \hat{P}(k_2) \rangle - \langle \hat{P}(k_1) \rangle \langle \hat{P}(k_2) \rangle$

$\mathbf{C}^{\mathrm{G}}(k_1, k_2) \simeq 2 \left\langle \delta_{W}(k_1) \delta_{W}(-k_2) \right\rangle \left\langle \delta_{W}(k_2) \delta_{W}(-k_1) \right\rangle$



I. Gaussian Covariance

Contains all dependence on cosmology and bias parameters

 $\mathbf{C}_{\ell_1\ell_2}^{\rm G}(k_1,k_2) \simeq \sum_{\ell'=\ell'} P_{\ell'_1}(k_2) P_{\ell'_2}(k_1) \left\{ \frac{(2\ell_1+1)(2\ell_2+1)}{\mathbf{I}_{22}^2} \int_{\hat{\mathbf{k}}_1,\hat{\mathbf{k}}_2,\mathbf{x}_1,\mathbf{x}_2} W_{22}(\mathbf{x}_1) W_{22}(\mathbf{x}_2) e^{-i(\mathbf{x}_1-\mathbf{x}_2)\cdot(\mathbf{k}_1-\mathbf{k}_2)} \right\}$

{Computed from survey} random catalog

 $\times \mathcal{L}_{\ell_1}(\hat{\mathbf{x}}_1 \cdot \hat{\mathbf{k}}_1) \mathcal{L}_{\ell'_1}(\hat{\mathbf{x}}_1 \cdot \hat{\mathbf{k}}_1) \mathcal{L}_{\ell'_2}(\hat{\mathbf{x}}_2 \cdot \hat{\mathbf{k}}_2) \left[\mathcal{L}_{\ell_2}(\hat{\mathbf{x}}_2 \cdot \hat{\mathbf{k}}_2) + \mathcal{L}_{\ell_2}(\hat{\mathbf{x}}_1 \cdot \hat{\mathbf{k}}_1) \right] \Big\}$

See also: Li et al. 2019

2. Non-Gaussian covariance

$$\begin{aligned} \mathbf{C}^{\mathrm{T}}(k_{1},k_{2}) = & \frac{1}{\mathbf{I}_{22}^{2}} \int_{\hat{\mathbf{k}}_{1},\hat{\mathbf{k}}_{2},\boldsymbol{\epsilon},\mathbf{p}_{1},\mathbf{p}_{2}} W_{11}(\mathbf{p}_{1}) W_{11}(\mathbf{p}_{2}) W_{11}(\boldsymbol{\epsilon}-\mathbf{p}_{1}) W_{11}(-\boldsymbol{\epsilon}-\mathbf{p}_{2}) \\ & \times T(\mathbf{k}_{1}-\mathbf{p}_{1},-\mathbf{k}_{1}-\boldsymbol{\epsilon}+\mathbf{p}_{1},\mathbf{k}_{2}-\mathbf{p}_{2},-\mathbf{k}_{2}+\boldsymbol{\epsilon}+\mathbf{p}_{2}) . \end{aligned}$$

can be split into an ordinary (T_0) and a long-mode contribution (T_{BC})



Beat Coupling effect: Hamilton, Rimes & Scoccimarro (2006)



Beat Coupling (BC)/ Super-sample (SSC)

Large-scale (super-survey) modes 0 couple to small-scale fluctuations



Beat Coupling (BC)/ Super-sample (SSC)

Large-scale (super-survey) modes 0 couple to small-scale fluctuations



k_1 k_2



Beat Coupling (BC)/ Super-sample (SSC)

- Large-scale (super-survey) modes couple to small-scale fluctuations
- Causes substantial contribution to non-diagonal covariance C(k₁,k₂)











Beat Coupling (BC)

Contribution of "beat-modes" (ϵ) to the covariance:

 $T(\mathbf{k}_1 - \mathbf{p}_1, -\mathbf{k}_1 - \boldsymbol{\epsilon} + \mathbf{p}_1, \mathbf{k}_2 - \mathbf{p}_2, -\mathbf{k}_2 - \boldsymbol{\epsilon} + \mathbf{p}_2) = 4P(\boldsymbol{\epsilon})[P_L(\mathbf{k}_1 - \mathbf{p}_1)F_2(\mathbf{k}_1 - \mathbf{p}_1, \boldsymbol{\epsilon})]$



Super-sample covariance approach (SSC)

$$T(\mathbf{k}_{1} - \mathbf{p}_{1}, -\mathbf{k}_{1} - \boldsymbol{\epsilon} + \mathbf{p}_{1}, \mathbf{k}_{2} - \mathbf{p}_{2}', -\mathbf{k}_{2} - \boldsymbol{\epsilon} + \mathbf{p}_{2}) = 4P(\boldsymbol{\epsilon})[P_{L}(\mathbf{k}_{1} - \mathbf{p}_{1})F_{2}(\mathbf{k}_{1} - \mathbf{p}_{1}, \boldsymbol{\epsilon})] + P_{L}(-\mathbf{k}_{1} - \mathbf{p}_{1}')F_{2}(-\mathbf{k}_{1} - \mathbf{p}_{1}', \boldsymbol{\epsilon})] \times [P_{L}(\mathbf{k}_{2} - \mathbf{p}_{2})F_{2}(\mathbf{k}_{2} - \mathbf{p}_{2}, -\boldsymbol{\epsilon}) + P_{L}(-\mathbf{k}_{2} - \mathbf{p}_{2}')F_{2}(-\mathbf{k}_{2} - \mathbf{p}_{2}', -\boldsymbol{\epsilon})]$$

Contribution of beat-mode covariance can be written as product of power spectrum responses

$$C^{\rm SSC}(k_1, k_2) = \sigma_W^2 \left[\frac{\partial P(k_1)}{\partial \delta_b} \right] \left[\frac{\partial P(k_2)}{\partial \delta_b} \right]$$

 (σ_w^2)

RMS fluctuations over window volume

Takada & Hu (2013)





asVw

Local Average Effect

Fluctuations in galaxy surveys 0 are normalized by the survey mean density (NOT true mean) \Rightarrow Beat-mode effect on covariance damped by ~90%

 $\frac{\partial}{\partial \delta_b} P(k) \to \frac{\partial}{\partial \delta_b} \frac{P(k)}{(1+\delta_b)^2} = \frac{\partial P(k)}{\partial \delta_b} -2P(k)$

de Putter et al. (2012)

 $\hat{\delta} \equiv \frac{\Delta \rho}{(\bar{\rho})_{\text{survey}}} \equiv \hat{\delta} \equiv \frac{\delta(\mathbf{k})}{1 + \delta_b}$



 $C^{\rm SSC}(k_1, k_2) = \sigma_W^2 \frac{\partial P(k_1)}{\partial s} \frac{\partial F}{\partial s}$ $\partial \delta_b$

 $\partial \delta_b$

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Final 'simplified' NG covariance

$$\begin{split} \mathbf{C}_{\ell_{1}\ell_{2}}^{\mathrm{T}}(k_{1},k_{2}) = & \frac{1}{\mathbf{I}_{22}^{2}} \int_{\epsilon}^{P} \mathbf{L}(\epsilon) \bigg\{ 4 \int_{\hat{\mathbf{k}}_{\epsilon_{1}},\mathbf{p}_{1}}^{W_{11}} W_{11}(\mathbf{p}_{1}) W_{11}(\epsilon - \mathbf{p}_{1}) P_{\mathbf{L}}(|\mathbf{k}_{1} - \mathbf{p}_{1}|) Z_{1}(\mathbf{k}_{1} - \mathbf{p}_{1}) Z_{2}(\mathbf{k}_{1} - \mathbf{p}_{1},\epsilon) \bigg\} \\ & \times \bigg\{ 4 \int_{\hat{\mathbf{k}}_{\epsilon_{2}},\mathbf{p}_{2}}^{W_{11}} W_{12}(\mathbf{p}_{2}) W_{11}(-\epsilon - \mathbf{p}_{2}) P_{\mathbf{L}}(|\mathbf{k}_{2} - \mathbf{p}_{2}|) Z_{1}(\mathbf{k}_{2} - \mathbf{p}_{2}) Z_{2}(\mathbf{k}_{2} - \mathbf{p}_{2},-\epsilon) \bigg\} \\ & + \frac{1}{\mathbf{I}_{22}^{2}} \int_{\hat{\mathbf{k}}_{\epsilon_{1}},\hat{\mathbf{k}}_{\epsilon_{2}},\epsilon}^{W_{11}} |W_{22}(\epsilon)|^{2} \bigg\{ \bigg[8 P_{\mathbf{L}}^{2}(\mathbf{k}_{1}) Z_{1}^{2}(\mathbf{k}_{1}) P_{\mathbf{L}}(\mathbf{k}_{1} + \mathbf{k}_{2}) Z_{2}^{2}(-\mathbf{k}_{1},\mathbf{k}_{1} + \mathbf{k}_{2}) + (\mathbf{k}_{1} \leftrightarrow \mathbf{k}_{2}) \bigg] \\ & + 16 P_{\mathbf{L}}(\mathbf{k}_{1}) Z_{1}(\mathbf{k}_{1}) P_{\mathbf{L}}(\mathbf{k}_{2}) Z_{1}(\mathbf{k}_{2}) P_{\mathbf{L}}(\mathbf{k}_{1} + \mathbf{k}_{2}) Z_{2}(-\mathbf{k}_{1},\mathbf{k}_{1} + \mathbf{k}_{2}) Z_{2}(-\mathbf{k}_{2},\mathbf{k}_{1} + \mathbf{k}_{2}) \bigg\} \\ & + 16 P_{\mathbf{L}}(\mathbf{k}_{1}) Z_{1}(\mathbf{k}_{2}) P_{\mathbf{L}}(\mathbf{k}_{2}) Z_{1}(\mathbf{k}_{2}) P_{\mathbf{L}}(\mathbf{k}_{1} + \mathbf{k}_{2}) Z_{2}(-\mathbf{k}_{1},\mathbf{k}_{1} + \mathbf{k}_{2}) Z_{2}(-\mathbf{k}_{2},\mathbf{k}_{1} + \mathbf{k}_{2}) \bigg\} \\ & + 16 P_{\mathbf{L}}(\mathbf{k}_{1}) Z_{1}(\mathbf{k}_{2}) P_{\mathbf{L}}(\mathbf{k}_{2}) Z_{1}(\mathbf{k}_{2}) P_{\mathbf{L}}(\mathbf{k}_{1} + \mathbf{k}_{2}) Z_{2}(-\mathbf{k}_{1},\mathbf{k}_{1} + \mathbf{k}_{2}) Z_{2}(-\mathbf{k}_{2},\mathbf{k}_{1} + \mathbf{k}_{2}) \bigg\} \\ & + 16 P_{\mathbf{L}}(\mathbf{k}_{1}) P_{\mathbf{L}^{2}}(\mathbf{k}_{2}) Z_{1}(\mathbf{k}_{2}) P_{\mathbf{L}}(\mathbf{k}_{2}) Z_{1}(\mathbf{k}_{1} - \mathbf{k}_{1}, \mathbf{k}_{2}) \bigg\} \\ & + \left[12 Z_{1}^{2}(\mathbf{k}_{1}) P_{\mathbf{L}^{2}}(\mathbf{k}_{2}) Z_{1}(\mathbf{k}_{2}) Z_{1}(\mathbf{k}_{2}) Z_{1}(\mathbf{k}_{1} + \mathbf{k}_{2}) Z_{2}(-\mathbf{k}_{2},\mathbf{k}_{1}, \mathbf{k}_{2}) \bigg\} \\ & - \frac{P(\mathbf{k}_{1}) P_{\mathbf{k}^{2}}(\mathbf{k}_{2}) Z_{1}^{2}}{\mathbf{k}} \int_{\epsilon}^{W_{10}(-\epsilon) W_{22}(\epsilon) [P(\epsilon) Z_{1}(\epsilon) Z_{21}(\mathbf{k}_{2}, \ell_{2}, \hat{\epsilon} \cdot \hat{\mathbf{n}}) + P_{\ell_{1}}(\mathbf{k}_{1}) b_{2}] \bigg\} \\ & - \frac{P(\mathbf{k}_{2}) P_{\mathbf{k}}(\mathbf{k}_{1})}{\mathbf{1}_{0} \mathbf{1}_{22}} \int_{\epsilon}^{W_{10}(-\epsilon) W_{22}(\epsilon) [P(\epsilon) Z_{1}(\epsilon) Z_{21}(\mathbf{k}_{2}, \ell_{2}, \hat{\epsilon} \cdot \hat{\mathbf{n}}) + P_{\ell_{2}}(\mathbf{k}_{2}) b_{2}] \bigg\} \\ & + \frac{1}{\mathbf{1}_{10}^{2}} \bigg\{ \bigg\{ \bigg\{ \sum_{i}^{N_{g}} - \alpha_{i} \sum_{j}^{N_{g}} \bigg\} \bigg\} \bigg\} \\ & + \frac{1}{\mathbf{1}_{22}^{2}} \int_{\hat{\mathbf{k}}_{\ell_{1}}, \hat{\mathbf{k}}_{\ell_{2}}} \bigg\} \bigg\{ \bigg\{ \bigg\{ \sum_{i}^{N_{g}} - \alpha_{i} \sum_{j}^{N$$

Beat-Coupling

-Ordinary Trispectrum (T₀)

Local Average effect

Shot noise contribution

DW & Scoccimarro 19





Comparison of Non-Gaussian effects

Ordinary trispectrum (T_0) in redshift-space is relatively more important at low-k



DW & Scoccimarro 19

See also: Mohammed et al. 2016

Even if we just use linear and tree-level PT (no Loops or FOG), our analytic approach works very well at high-k, why!?









Shot Noise

 $< \frac{1}{\overline{n}}$

- We observe particles
 (galaxies) instead of fields
- Affects results at high-k:

P(k)



Why does analytic work so well?

Shot noise (dashed) dominates at high-k

$$\mathbf{C}(k_1, k_2) = \langle \delta(k_1)\delta(-k_1)\delta(k_2)\delta(-k_2) \rangle$$
$$- \langle \delta(k_1)\delta(-k_1)\rangle \langle \delta(k_2)\delta(-k_2) \rangle$$
$$\mathbf{C}^{\mathbf{G}}(k_1, k_2) \simeq 2 \langle \delta(k_1)\delta(-k_2)\rangle \langle \delta(k_2)\delta(-k_1) \rangle$$
$$\mathbf{C}^{\mathbf{NG}}(k_1, k_2) = \langle \delta(k_1)\delta(-k_2)\delta(k_2)\delta(-k_1) \rangle_c$$



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Why does analytic work so well?

Shot noise (dashed) dominates at high-k

$$\mathbf{C}(k_1, k_2) = \langle \delta(k_1)\delta(-k_1)\delta(k_2)\delta(-k_2) \rangle$$
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$$\mathbf{C}^{\mathbf{NG}}(k_1, k_2) = \langle \delta(k_1)\delta(-k_2)\delta(k_2)\delta(-k_1)\rangle_c$$



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Why does analytic work so well?

Leakage due to anisotropic survey window & radial LOS

Monopole (\Rightarrow shot noise) leaks into multipole covariance at high-k



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Shot noise level for Ag future surveys

Redshift surveys typically designed with $\bar{n}P(k_{\rm BAO}) \simeq {\rm few}$

> Font-Ribera et al. 2014



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Shot noise level for $\hat{10}^1$ future surveys

Redshift surveys typically designed with $\bar{n}P(k_{\rm BAO}) \simeq {\rm few}$

> Font-Ribera et al. 2014

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Important differences between our approach and rest of literature

I. Estimator motivated by simulations

$$\hat{\delta} \equiv \frac{\delta(\mathbf{k})}{1+\delta_b}$$

Our

I. FKP estimator (motivated from surveys)

$$\hat{\delta}^{\text{FKP}}(\mathbf{x}) \equiv \frac{w(\mathbf{x})[n_g(\mathbf{x}) - \alpha n_s(\mathbf{x})]}{[\alpha \int d^3 x \, w^2(\mathbf{x}) \bar{n}(\mathbf{x}) n_s(\mathbf{x})]}$$

where $\alpha = N_g / N_s$ and simplifies to

$$\hat{\delta}^{\rm FKP} \simeq \frac{1}{\sqrt{I_{22}}} \frac{W(\mathbf{x})\delta(\mathbf{x})}{(1+\delta_{\rm Ng})^{1/2}}$$



I. Estimator motivated by simulations

 $\hat{\delta} \equiv$

2. Normalized by a matter mode Tresponses for tides, bias, shot noise 7 need to be added separately

Our Approach I. FKP estimator (motivated from surveys)

$$\hat{\delta}^{\rm FKP} \simeq \frac{1}{\sqrt{I_{22}}} \frac{W(\mathbf{x})\delta(\mathbf{x})}{(1+\delta_{\rm N_g})^{1/2}}$$

2. Normalized by a galaxy mode includes bias, shot noise _ and non-linearities like tides



I. Estimator motivated by simulations

$$\hat{\delta} \equiv \frac{\delta(\mathbf{k})}{1+\delta_b}$$

2. Normalized by a matter mode Tresponses for tides, bias, shot noise 7 need to be added separately

3. Stronger normalization of long mode effects

$$\frac{\partial}{\partial \delta_b} P(k) \to \frac{\partial}{\partial \delta_b} \frac{P(k)}{(1+\delta_b)^2} = \frac{\partial P(k)}{\partial \delta_b} - 2P(k)$$

~0.7 P(k) at k=0.1 h/Mpc

Our Approach I. FKP estimator (motivated from surveys)

$$\hat{\delta}^{\mathrm{FKP}} \simeq \frac{1}{\sqrt{\mathrm{I}_{22}}} \frac{W(\mathbf{x})\delta(\mathbf{x})}{(1+\delta_{\mathrm{Ng}})^{1/2}}$$

2. Normalized by a galaxy mode includes bias, shot noise and non-linearities like tides

3. Weaker normalization of long mode effects

$$\frac{\partial \hat{P}^{\text{FKP}}}{\partial \delta_{\text{Ng}}} \rightarrow \frac{\partial}{\partial \delta_{\text{Ng}}} \frac{P}{(1+\delta_{\text{Ng}})} = \frac{\partial P}{\partial \delta_{\text{Ng}}} - \frac{P(k)}{P(k)}$$

~1.7 P(k) at k=0.1 h/Mpc



I. Estimator motivated by simulations

$$\hat{\delta} \equiv \frac{\delta(\mathbf{k})}{1+\delta_b}$$

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$$\hat{\delta}^{\mathrm{FKP}} \simeq \frac{1}{\sqrt{\mathrm{I}_{22}}} \frac{W(\mathbf{x})\delta(\mathbf{x})}{(1+\delta_{\mathrm{Ng}})^{1/2}}$$

2. Normalized by a galaxy mode includes bias, shot noise and non-linearities like tides

3. Weaker normalization of long mode effects

Effect of long mode on covariance is 5 x stronger than literature



SSC approach assumes a top-hat window

$$C^{\rm SSC}(k_1, k_2) = \sigma_W^2 \left[\frac{\partial P(k_1)}{\partial \delta_b} \right] \left[\frac{\partial P(k_2)}{\partial \delta_b} \right]$$

$$\sigma_W^2 = \frac{1}{V_W^2} \int \frac{d^3 \epsilon}{(2\pi)^3} P(\epsilon) |W(\epsilon)|^2$$

In the presence of a local-average effect, we get three different types of σ_w^2 0

$$\frac{1}{\mathbf{I}_{22}^2} \int_{\boldsymbol{\epsilon}} P(\boldsymbol{\epsilon}) |W_{22}(\boldsymbol{\epsilon})|^2$$

$$\frac{1}{\mathrm{N_g I_{22}}} \int_{\boldsymbol{\epsilon}} P(\boldsymbol{\epsilon}) W_{22}(\boldsymbol{\epsilon}) \bar{n}(-\boldsymbol{\epsilon}) W_{22}($$

 $W_{22}(\mathbf{x}) = \bar{n}^2(\mathbf{x}) w_{\rm FKP}^2(\mathbf{x})$

These differ by up to 35% for the SDSS window

RMS fluctuations ver window volume

$$\langle \delta_{N_{g}}^{2} \rangle \equiv \frac{1}{I_{10}^{2}} \int_{\boldsymbol{\epsilon}} P(\boldsymbol{\epsilon}) |\bar{n}(\boldsymbol{\epsilon})|^{2}$$

Summary

- Analytic covariance is an excellent 0 alternative to mocks
 - Monopole Leakage simplifies RSD modeling
 - Shot noise dominates at high-k 2. (Dominates over loops and FOG)



Mocks spend computer time calculating a 0 slightly noisy covariance plus shot noise



Summary

- Analytic covariance is an excellent alternative to mocks
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Mocks spend computer time calculating a slightly noisy covariance plus shot noise

Next

- Bispectrum covariance 0 (Simulations are computationally prohibitive)
- Redo BOSS analysis with parameter-0 dependent analytic covariance



Analytic covariance is crucial for BISPECTRUM

- Number of triangles to estimate ~ 6000 !! 0
- Number of mock simulations: O(1000)0
- Bottleneck for BOSS 0 - Gil-Marin et al 17 could only use ~800 triangles





True vs FKP shot noise

$$\hat{P}_0(k) = \frac{1}{\mathbf{I}_{22}} \left[\int \frac{d\Omega_k}{4\pi} |F_0(\mathbf{k})|^2 - \mathbf{N}_0 \right]$$
$$F_0(\mathbf{k}) = \left(\sum_{j=1}^{N_g} -\alpha \sum_{j=1}^{N_r} \right) w_j e^{i\mathbf{k}\cdot\mathbf{x}_j} \quad \alpha = \frac{N_g}{N_r}$$

$$\mathbf{N_0^{true}} \equiv \left(\sum_{j=1}^{\mathbf{N_g}} + \alpha^2 \sum_{j=1}^{\mathbf{N_r}}\right) w_j^2 \qquad \mathbf{N_0^{FKP}} \equiv \left(\alpha \sum_{j=1}^{\mathbf{N_r}} + \alpha^2 \sum_{j=1}^{\mathbf{N_r}}\right) w_j^2 = \mathbf{N_0^{FKP}} = \mathbf{$$

100% change in covariance at k~0.4 !





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