Gappy Weak Lensing: KL interpolation of shear fields

Jake VanderPlas May 6, 2011

Outline

Challenges of Weak Lensing cluster searches

Shear Peak statistics: a solution?

KL interpolation of shear fields

Results and future prospects

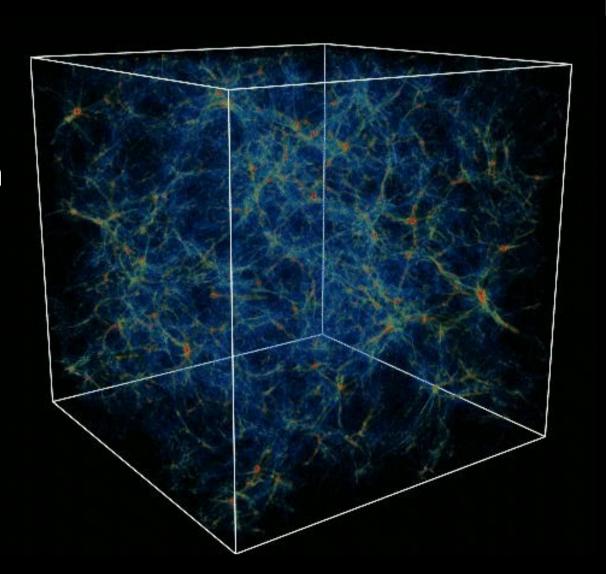
Motivation:

Galaxy Clusters probe cosmology:

Number counts with redshift

2-point correlation function

Mass profiles & Substructure



Finding Clusters: Optical & IR

Most straightforward method: search for groups of galaxies (e.g. SDSS, 2MASS)

Problem: how well do optical/IR sources trace the mass distribution?



Perseus Cluster: 2MASS

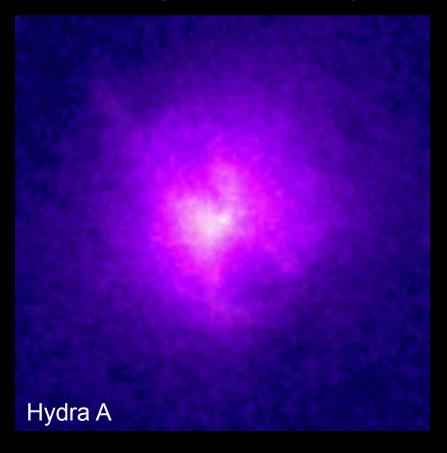


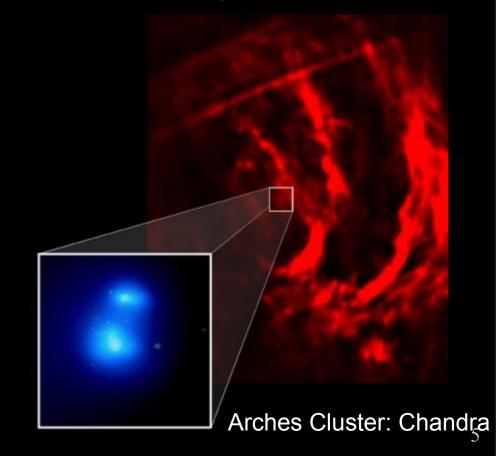
Coma Cluster: SDSS

Finding Clusters: X-ray

Look for X-ray signatures of Intracluster gas

Problem: uncertain conversions from X-ray flux to gas density to dark matter density

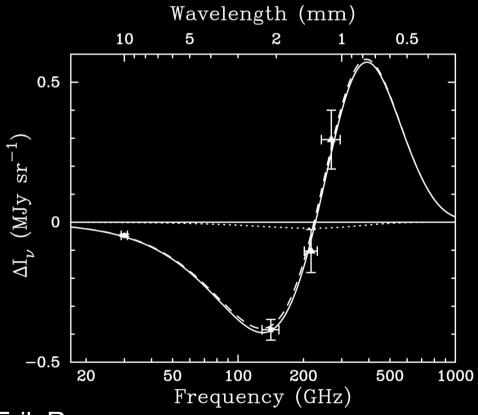


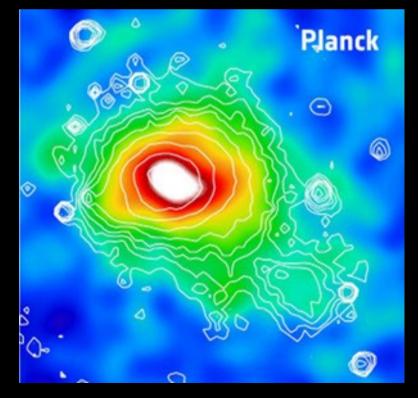


Finding Clusters: SZ

Sunyaev and Zeldovich's bright idea: look for scattering of CMB off hot gas!

Problem: gas density to mass conversion

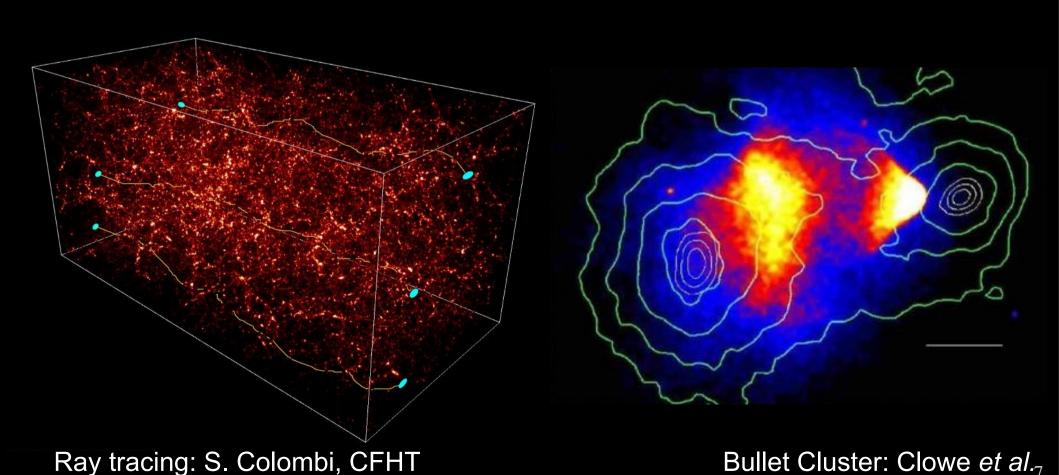




Planck & ROSAT: Coma Cluster

Finding Clusters: Weak Lensing

Look for gravitational distortion of background galaxies: towards a robust mass-selected cluster catalog?

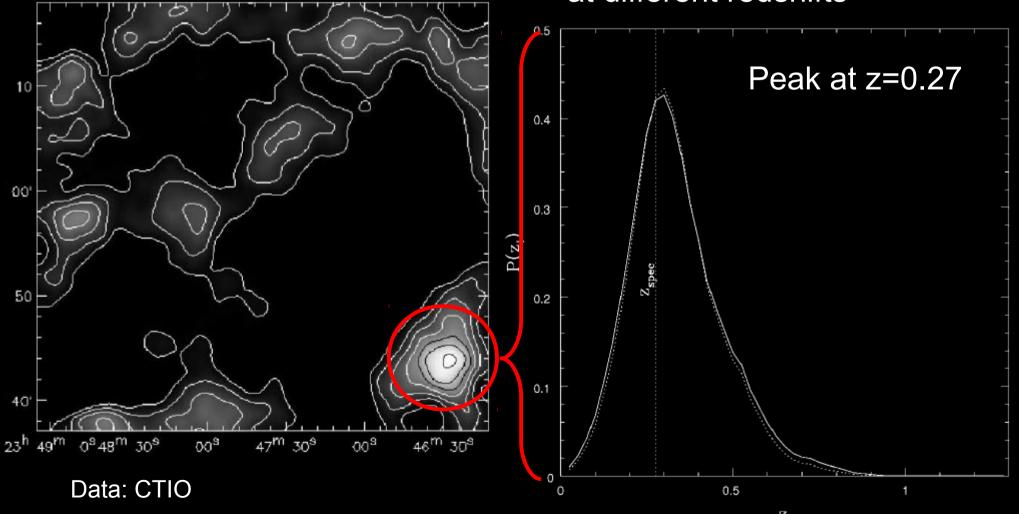


Finding Clusters in 3D

Parametric methods:

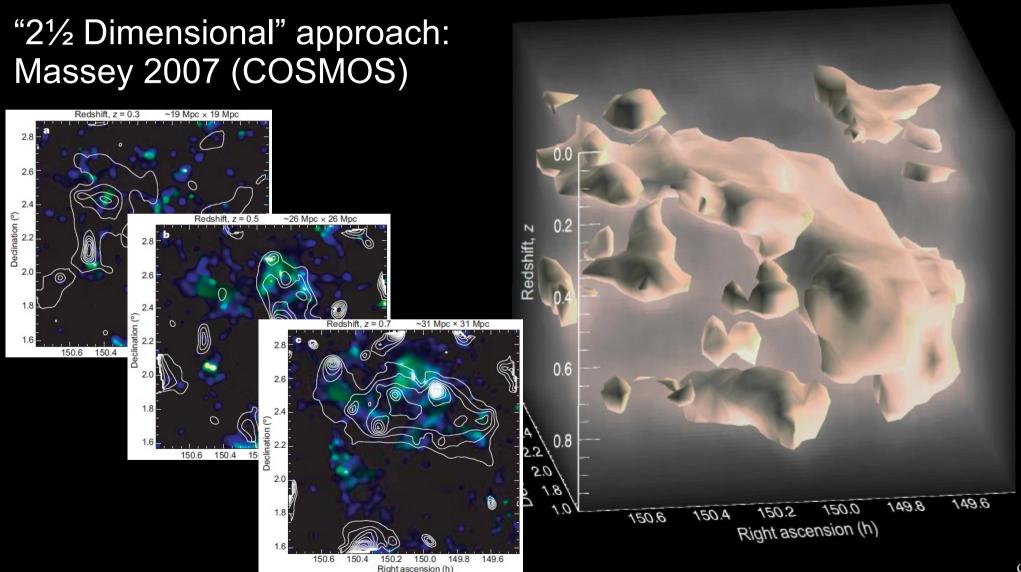
e.g. Wittman et al. 2001

Fit SIS and NFW profiles at different redshifts

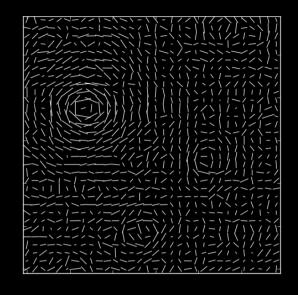


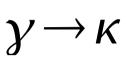
Finding Clusters in 3D

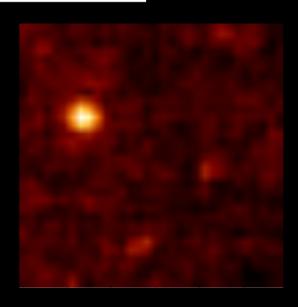
Nonparametric methods:



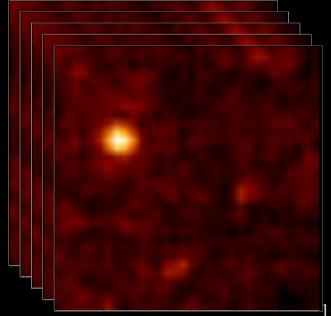
Toward a full 3D reconstruction:







$$\kappa \! \to \! \delta$$



Toward a full 3D reconstruction:

Vanderplas et al 2011: SVD filtering

$$\gamma \rightarrow \kappa$$
: $\gamma(\vec{\theta}) = \int d\vec{\theta}^{'2} \kappa(\vec{\theta}') D(\vec{\theta} - \vec{\theta}')$

 $\vec{\gamma} = P_{\gamma\kappa} \vec{\kappa}$: operates in each source plane

$$\kappa \to \delta: \qquad \kappa(\chi_S) = \frac{3H_0^2\Omega_M}{2} \int_0^{\chi_S} \frac{\chi(\chi_S - \chi)}{\chi_S} \frac{1 + \delta(\chi)}{a(\chi)} d\chi$$

 $\vec{\kappa} = Q_{\kappa\delta} \vec{\delta}$: operates in each line-of-sight

Final Result:
$$\rightarrow \vec{y} = M \vec{\delta}$$

Hu&Keeton 2002 Simon et al 2009

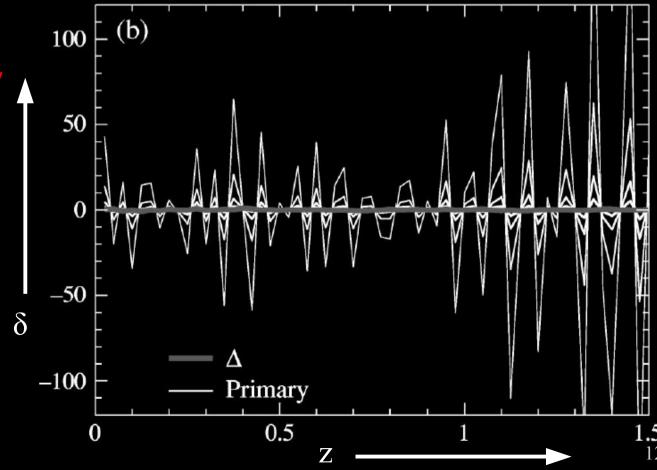
3D Lensing is simply a linear inversion:

$$\vec{\gamma} = M \, \vec{\delta} + \vec{n_{\gamma}}$$

Best estimator, via Aitken (1935):

$$\hat{\delta} = \left(M^T N_{\gamma \gamma}^{-1} M \right)^{-1} M^T N_{\gamma \gamma}^{-1} \vec{\gamma}$$

Problem: Noise can obscure the signal by several orders of magnitude!



(Hu & Keeton 2002)

Tracing the source of the problem:

$$\vec{y} = M \, \vec{\delta} + n_{y}$$

Singular Value Decomposition (SVD)

$$N_{\gamma\gamma}^{-1/2} M = U \Sigma V^T$$

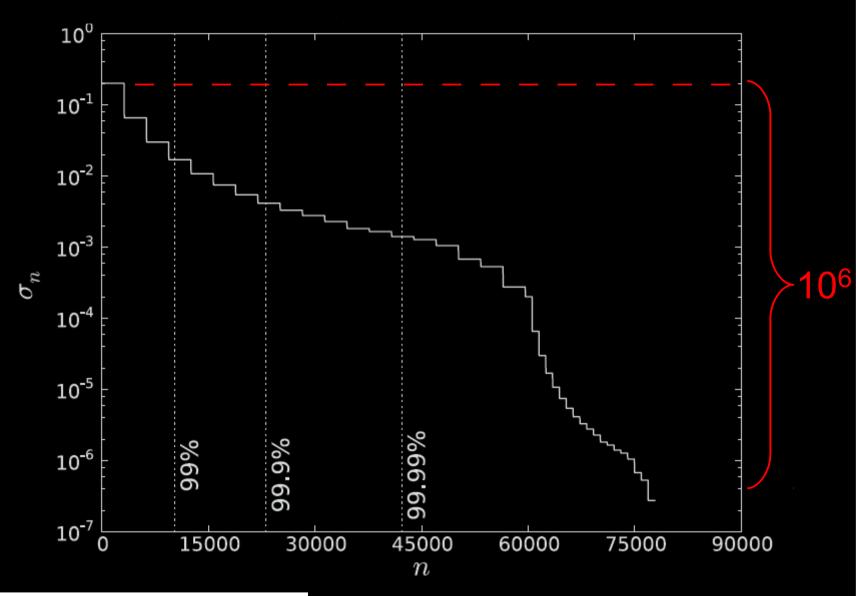
$$\begin{cases} U^T U = V^T V = I \\ \Sigma = \text{diagonal} \end{cases}$$

Aitken estimator becomes:

$$\hat{\delta} = V \sum^{-1} U^T N_{\gamma\gamma}^{-1/2} \vec{\gamma}$$

Small singular values lead to large noise in δ !

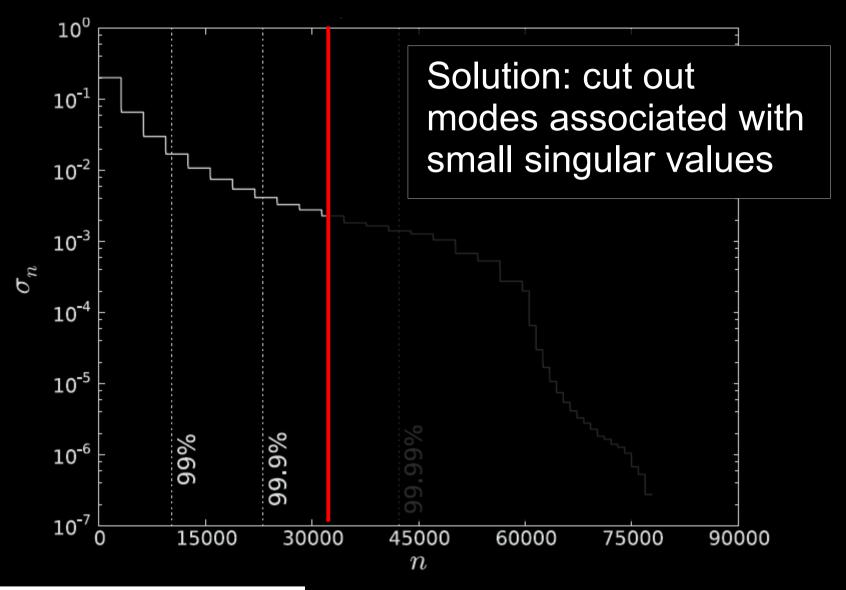
Singular Value Profile:



$$\hat{\delta} = V \Sigma^{-1} U^T N_{\gamma\gamma}^{-1/2} \vec{\gamma}$$

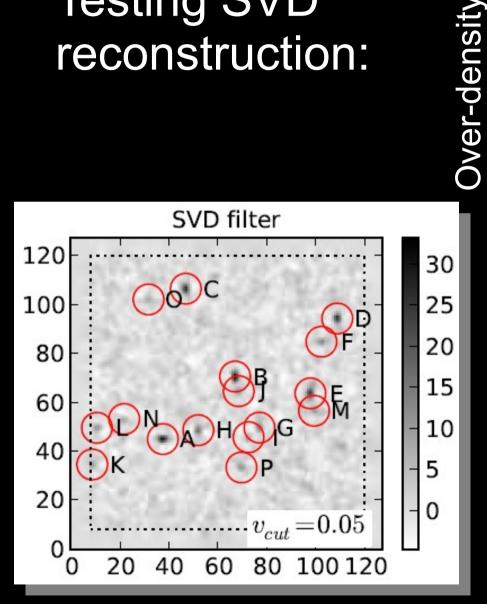
Noise amplified by a factor of 10⁶!

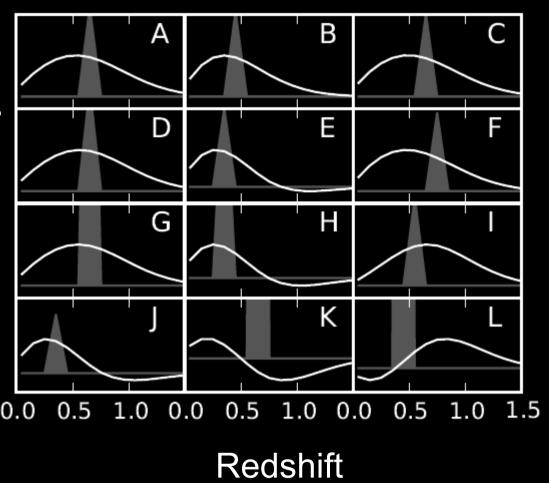
Singular Value Profile:



$$\hat{\delta} = V \sum^{-1} U^T N_{\gamma\gamma}^{-1/2} \vec{\gamma}$$

Testing SVD reconstruction:





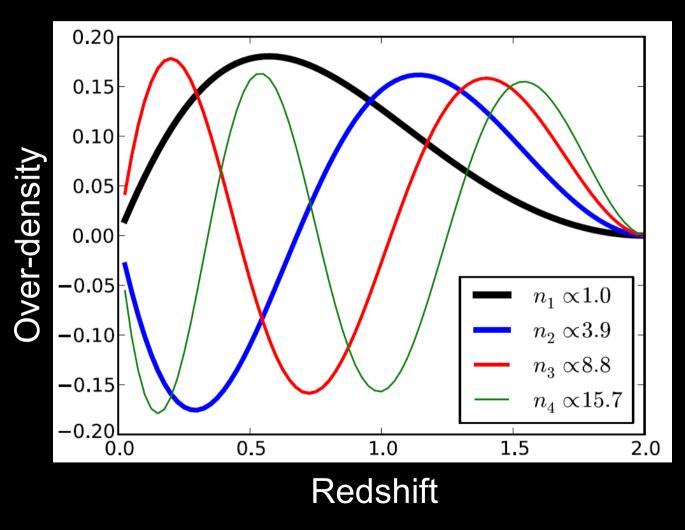
Spread and bias in redshift direction

Vanderplas et al. 2011

3D non-parametric weak lensing: Fundamentally Limited?

Noise in mode n scales as ~n²

Typical surveys can constrain first ~ few modes



Vanderplas et al. 2011

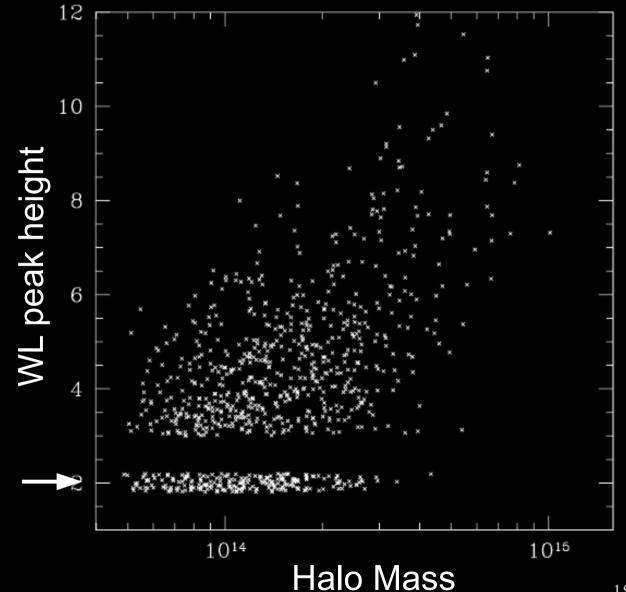
Cluster Mass Calibration

Difficulty with WL mass calibration

- Bias from shape noise

Also...

- Correlated Projections
- Unresolved Substructure
- Halo asymmetry
- Redshift dependence



Undetected peaks here

Hamana et al. 2004

Problem:

Data and theory are difficult to compare

Weak lensing yields projected mass,

Theory gives 3D mass.

Solution:

Rather than force-fitting data to theory, let's work toward a theory that naturally fits the data.



Shear Peak Statistics

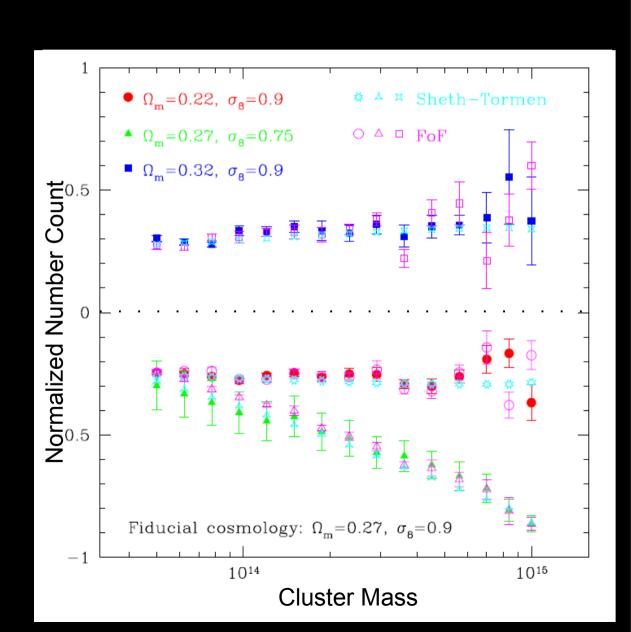
Pioneering Work: Marian et al. 2009, Wang et al. 2009

Marian et al:

How do *projected* mass peaks scale with cosmology?

Slab Thickness: 50Mpc (Δz~0.1 at z=0.6)

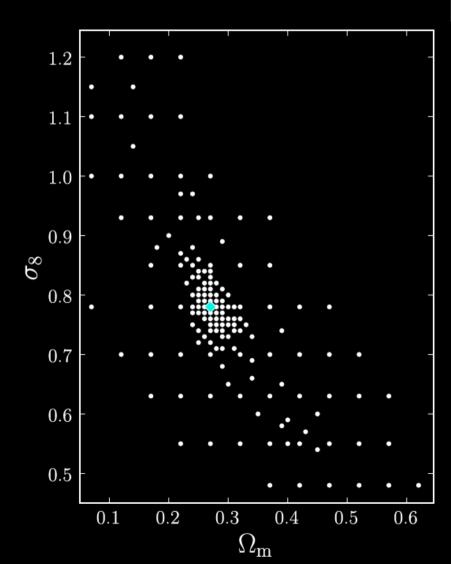
Marian *et al* 2010: results extended to different thicknesses



Shear Peak Statistics

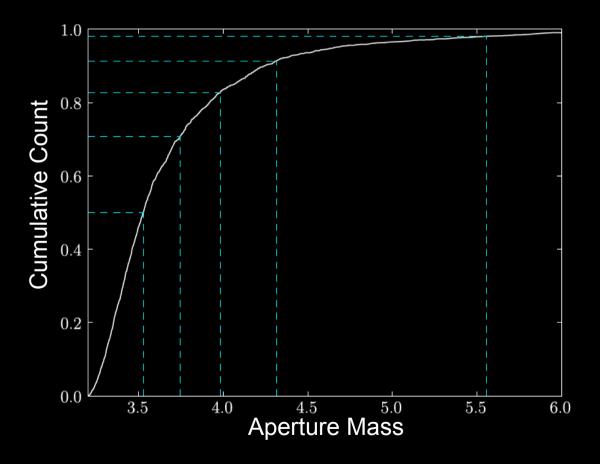
What information is in the correlated projections?

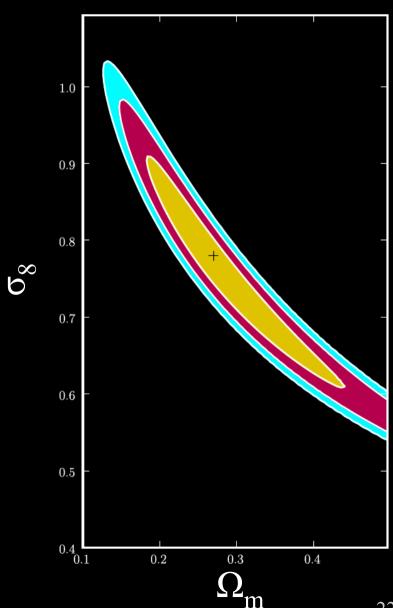
Dietrich & Hartlap 2010: Test cosmological information content of shear peaks



Cosmology Constraints

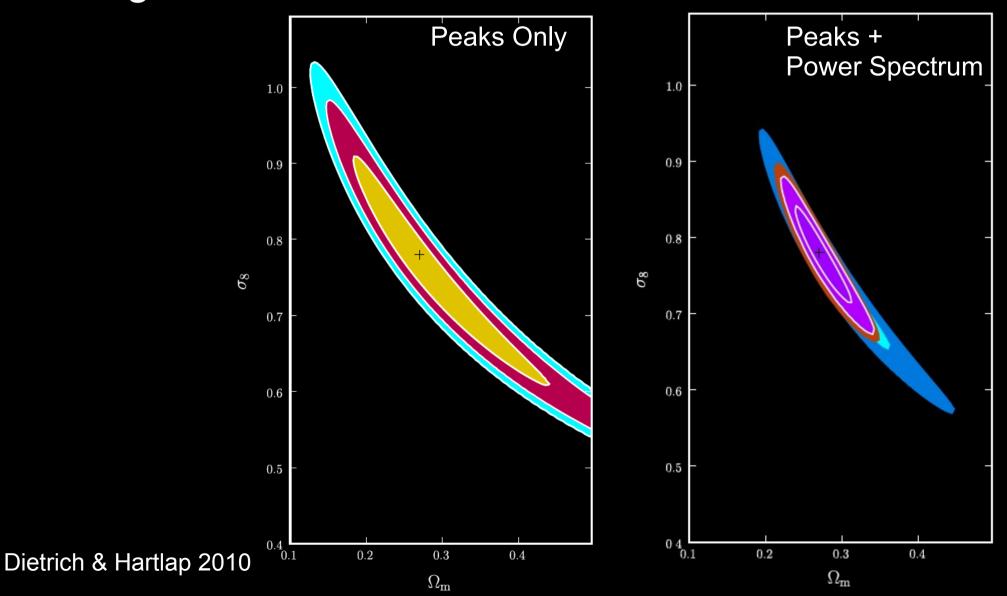
Cumulative distribution of peak heights can be used to constrain cosmology



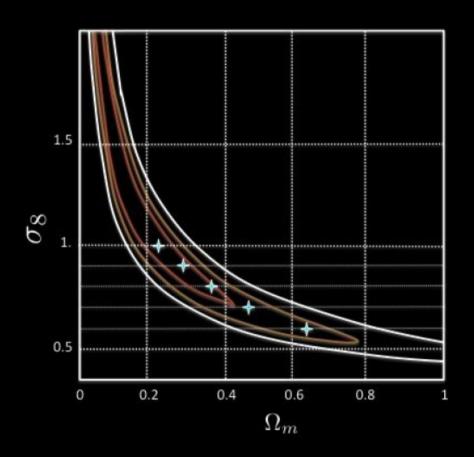


Shear Peak Statistics

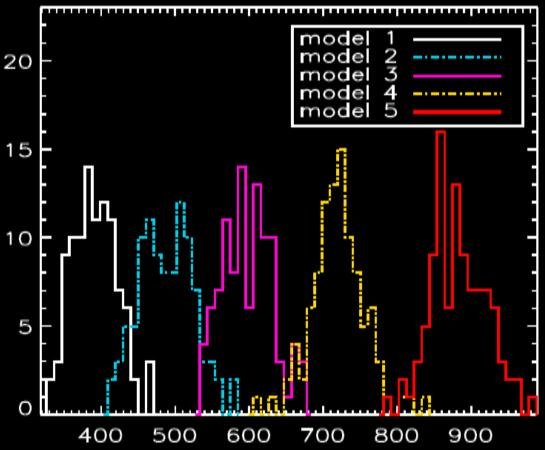
Peak functions probe nonlinear structure: higher-order information!



Shear Peaks: Higher order information



Using a "wavelet transform" filter gives similar discriminatory power



Pires et al. 2009

Shear Peak Statistics

Summary:

Shear peak counts scale with cosmology in predictable ways (Marian 2009, 2010)

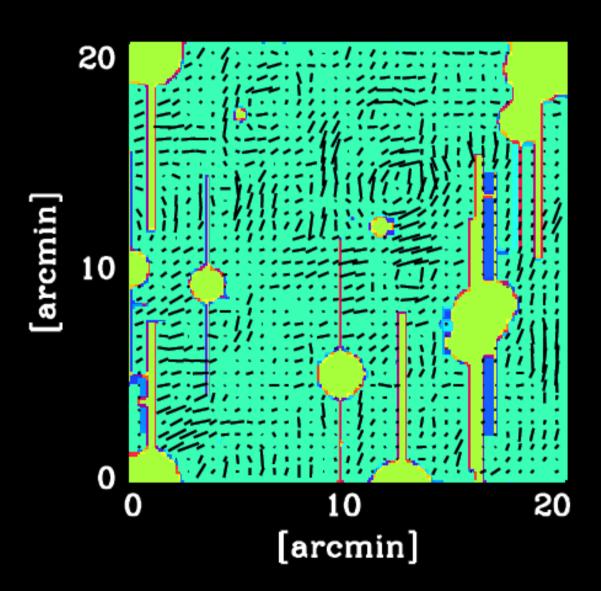
Projected structure encodes cosmological information (Wang 2009, Kratochvil 2010)

Peak distributions contain complementary information to 2-point analyses (Dietrich 2010, Pires 2009. See also Maturi 2011)

Methodology is in its infancy: the ideal filtering and peak-finding method needs to be explored (but see Pires 2009, Schmidt & Rozo 2010)

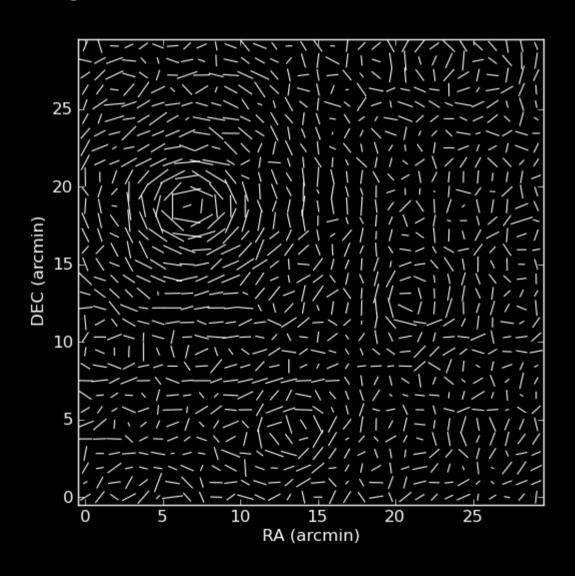
Shear Peaks: The Problem

Shear surveys are subject to masking effects: what sort of bias will this create?



Our Solution: KL analysis

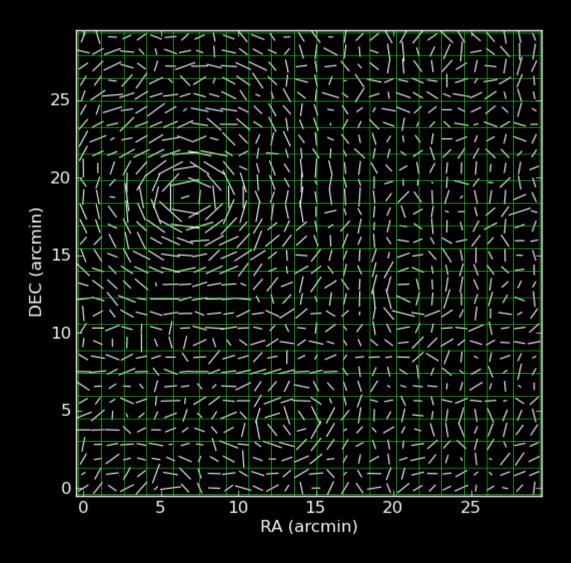
Use the *theoretical* 2-point correlation function to reconstruct the missing information



Our Solution: KL analysis

Use the *theoretical* 2-point correlation function to reconstruct the missing information

1. Pixelize the shear



Our Solution: KL analysis

Use the *theoretical* 2-point correlation function to reconstruct the missing information

- 1. Pixelize the shear
- 2. Compute the correlation of the shear between pixels, using the expected nonlinear matter power spectrum.

(We use Smith *et al* 2003)

$$\begin{split} \boldsymbol{\xi}_{ij} &= \langle \gamma_i \gamma_j^* \rangle + \langle n_i n_j^* \rangle \\ &= \left[\int_{A_i} d^2 x_i \int_{A_j} d^2 x_j \boldsymbol{\xi}_+ (|\boldsymbol{x_i} - \boldsymbol{x_j}|) \right] + \delta_{ij} \frac{\sigma_\epsilon^2}{\bar{n}} \\ & \\ \boldsymbol{\xi}_+(\theta) &= \frac{1}{2\pi} \int_0^\infty d\ell \; \ell P_\gamma(\ell) J_0(\ell\theta) \\ & \\ P_\gamma(\ell) &= \int_0^{\chi_s} d\chi W^2(\xi) \chi^{-2} P_\delta \left(k = \frac{\ell}{\chi}; z(\chi) \right) \\ & \\ \boldsymbol{\xi}_+(\theta) &= \frac{1}{2\pi} \int_0^\infty d\ell \; \ell P_\gamma(\ell) J_0(\ell\theta) \\ & \\ \boldsymbol{\xi}_+(\theta) &= \frac{1}{2\pi} \int_0^\infty d\ell \; \ell P_\gamma(\ell) J_0(\ell\theta) \\ & \\ \boldsymbol{\xi}_+(\theta) &= \frac{1}{2\pi} \int_0^\infty d\ell \; \ell P_\gamma(\ell) J_0(\ell\theta) \\ & \\ \boldsymbol{\xi}_+(\theta) &= \frac{1}{2\pi} \int_0^\infty d\ell \; \ell P_\gamma(\ell) J_0(\ell\theta) \\ & \\ \boldsymbol{\xi}_+(\theta) &= \frac{1}{2\pi} \int_0^\infty d\ell \; \ell P_\gamma(\ell) J_0(\ell\theta) \\ & \\ \boldsymbol{\xi}_+(\theta) &= \frac{1}{2\pi} \int_0^\infty d\ell \; \ell P_\gamma(\ell) J_0(\ell\theta) \\ & \\ \boldsymbol{\xi}_+(\theta) &= \frac{1}{2\pi} \int_0^\infty d\ell \; \ell P_\gamma(\ell) J_0(\ell\theta) \\ & \\ \boldsymbol{\xi}_+(\theta) &= \frac{1}{2\pi} \int_0^\infty d\ell \; \ell P_\gamma(\ell) J_0(\ell\theta) \\ & \\ \boldsymbol{\xi}_+(\theta) &= \frac{1}{2\pi} \int_0^\infty d\ell \; \ell P_\gamma(\ell) J_0(\ell\theta) \\ & \\ \boldsymbol{\xi}_+(\theta) &= \frac{1}{2\pi} \int_0^\infty d\ell \; \ell P_\gamma(\ell) J_0(\ell\theta) \\ & \\ \boldsymbol{\xi}_+(\theta) &= \frac{1}{2\pi} \int_0^\infty d\ell \; \ell P_\gamma(\ell) J_0(\ell\theta) \\ & \\ \boldsymbol{\xi}_+(\theta) &= \frac{1}{2\pi} \int_0^\infty d\ell \; \ell P_\gamma(\ell) J_0(\ell\theta) \\ & \\ \boldsymbol{\xi}_+(\theta) &= \frac{1}{2\pi} \int_0^\infty d\ell \; \ell P_\gamma(\ell) J_0(\ell\theta) \\ & \\ \boldsymbol{\xi}_+(\theta) &= \frac{1}{2\pi} \int_0^\infty d\ell \; \ell P_\gamma(\ell) J_0(\ell\theta) \\ & \\ \boldsymbol{\xi}_+(\theta) &= \frac{1}{2\pi} \int_0^\infty d\ell \; \ell P_\gamma(\ell) J_0(\ell\theta) \\ & \\ \boldsymbol{\xi}_+(\theta) &= \frac{1}{2\pi} \int_0^\infty d\ell \; \ell P_\gamma(\ell) J_0(\ell\theta) \\ & \\ \boldsymbol{\xi}_+(\theta) &= \frac{1}{2\pi} \int_0^\infty d\ell \; \ell P_\gamma(\ell) J_0(\ell\theta) \\ & \\ \boldsymbol{\xi}_+(\theta) &= \frac{1}{2\pi} \int_0^\infty d\ell \; \ell P_\gamma(\ell) J_0(\ell\theta) \\ & \\ \boldsymbol{\xi}_+(\theta) &= \frac{1}{2\pi} \int_0^\infty d\ell \; \ell P_\gamma(\ell) J_0(\ell\theta) \\ & \\ \boldsymbol{\xi}_+(\theta) &= \frac{1}{2\pi} \int_0^\infty d\ell \; \ell P_\gamma(\ell\theta) J_0(\ell\theta) \\ & \\ \boldsymbol{\xi}_+(\theta) &= \frac{1}{2\pi} \int_0^\infty d\ell \; \ell P_\gamma(\ell\theta) J_0(\ell\theta) \\ & \\ \boldsymbol{\xi}_+(\theta) &= \frac{1}{2\pi} \int_0^\infty d\ell \; \ell P_\gamma(\ell\theta) J_0(\ell\theta) \\ & \\ \boldsymbol{\xi}_+(\theta) &= \frac{1}{2\pi} \int_0^\infty d\ell \; \ell P_\gamma(\ell\theta) J_0(\ell\theta) \\ & \\ \boldsymbol{\xi}_+(\theta) &= \frac{1}{2\pi} \int_0^\infty d\ell \; \ell P_\gamma(\ell\theta) J_0(\ell\theta) \\ & \\ \boldsymbol{\xi}_+(\theta) &= \frac{1}{2\pi} \int_0^\infty d\ell \; \ell P_\gamma(\ell\theta) J_0(\ell\theta) \\ & \\ \boldsymbol{\xi}_+(\theta) &= \frac{1}{2\pi} \int_0^\infty d\ell \; \ell P_\gamma(\ell\theta) J_0(\ell\theta) \\ & \\ \boldsymbol{\xi}_+(\theta) &= \frac{1}{2\pi} \int_0^\infty d\ell \; \ell P_\gamma(\ell\theta) J_0(\ell\theta) \\ & \\ \boldsymbol{\xi}_+(\theta) &= \frac{1}{2\pi} \int_0^\infty d\ell \; \ell P_\gamma(\ell\theta) J_0(\ell\theta) \\ & \\ \boldsymbol{\xi}_+(\theta) &= \frac{1}{2\pi} \int_0^\infty d\ell \; \ell P_\gamma(\ell\theta) J_0(\ell\theta) \\ & \\ \boldsymbol{\xi}_+(\theta) &= \frac{1}{2\pi} \int_0^\infty d\ell \; \ell P_\gamma(\ell\theta) J_0(\ell\theta) \\ & \\ \boldsymbol{\xi}_+(\theta) &= \frac{1}{2\pi} \int_0^\infty d\ell \; \ell P_\gamma(\ell\theta) J_0(\ell\theta) \\ & \\ \boldsymbol{\xi}_+(\theta) &= \frac{1}{2\pi} \int_$$

KL Decomposition:

Now that we have the theoretical correlation matrix ξ , we can compute the KL basis via an eigenvalue decomposition: $\xi = \Psi \Lambda \Psi^{\dagger}$

These eigenvectors Ψ are an orthogonal basis, and give the *optimal* low-rank reconstruction of a shear vector.

$$oldsymbol{\gamma} = oldsymbol{\Psi} oldsymbol{a} \ oldsymbol{a} = oldsymbol{\Psi}^\dagger oldsymbol{\gamma}$$

$$\hat{\boldsymbol{\gamma}}^{(n)} \equiv \sum_{i=1}^{n < N} a_i \boldsymbol{\Psi}_i$$

KL Decomposition:

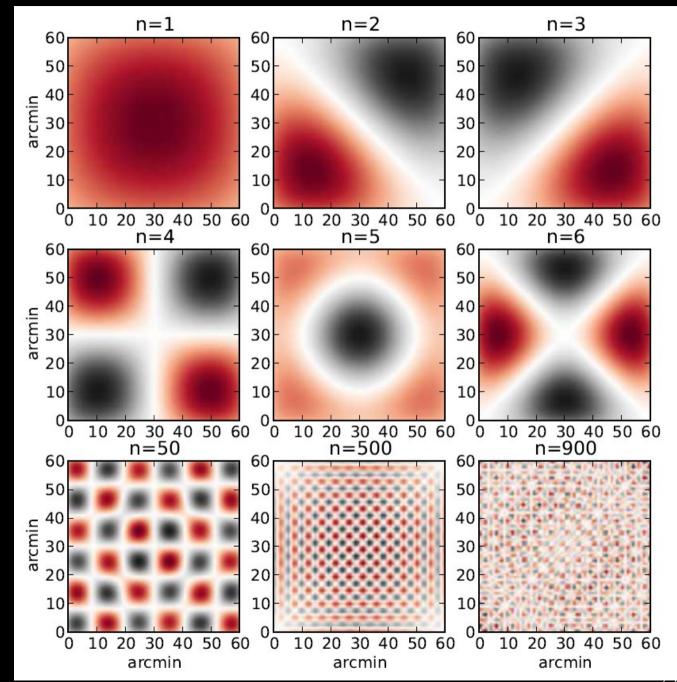
We can use the unmasked region to constrain the coefficients

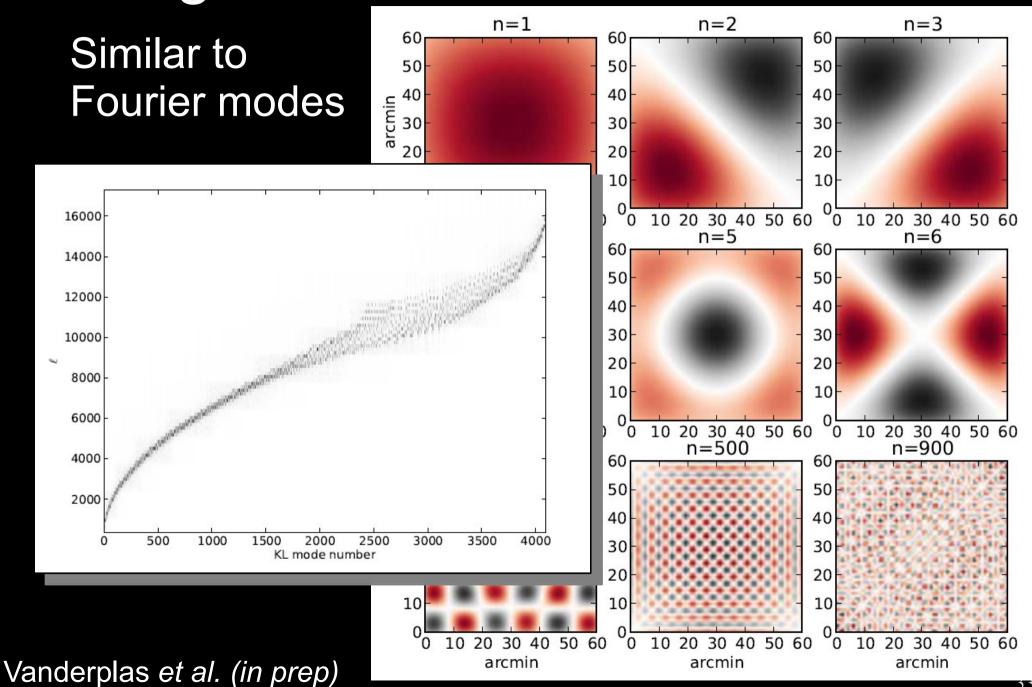
$$\hat{\boldsymbol{a}}_{(n)} = \boldsymbol{M}_{(n)}^{-1} \boldsymbol{\Psi}_{(n)}^{\dagger} \boldsymbol{W} \boldsymbol{\mathcal{N}}_{\gamma}^{-1/2} \boldsymbol{\gamma}^{o}$$

$$oldsymbol{M}_{(n)} \equiv oldsymbol{\Psi}^\dagger_{(n)} oldsymbol{W} oldsymbol{\Psi}_{(n)}$$

These coefficients can then be used to estimate the masked-out shear, and additionally filter noise from the entire field

$$oldsymbol{\hat{\gamma}}^{(n)} = oldsymbol{\mathcal{N}}_{\gamma}^{1/2} oldsymbol{\Psi}_{(n)} oldsymbol{\hat{a}}_{(n)}$$





Eigenvalues encode Signal-to-Noise

10¹

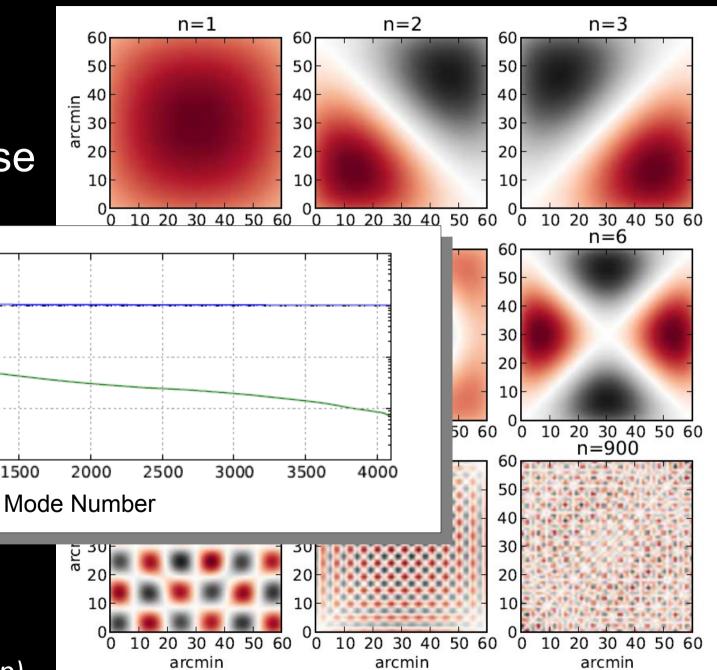
10°

10⁻¹

10-2

10⁻³

Signal / Noise



Vanderplas et al. (in prep)

500

1000

1500

Eigenvalues encode Signal-to-Noise

10¹

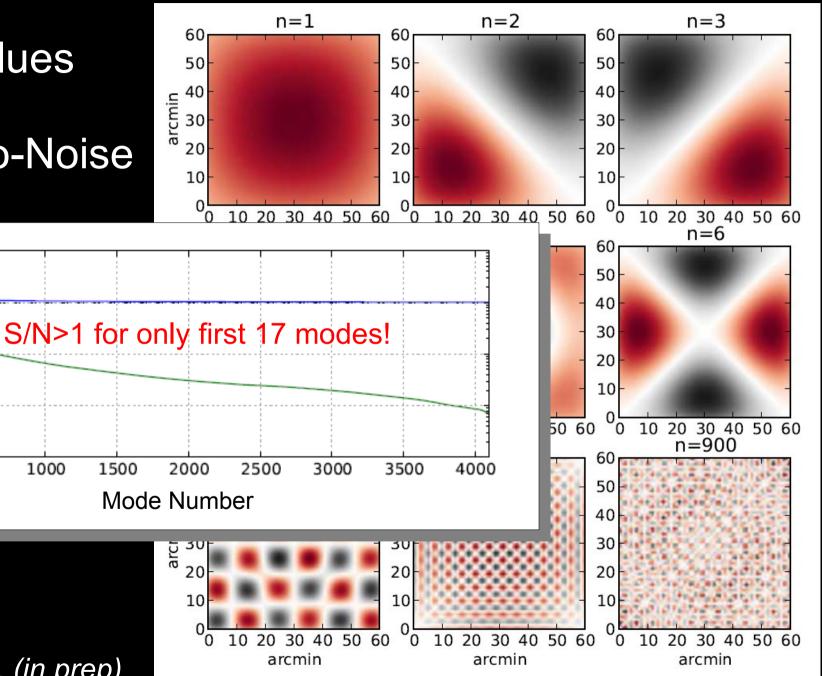
10°

10-1

10-2

10⁻³

Signal / Noise



500

1000

1500

Eigenvalues encode Signal-to-Noise

10¹

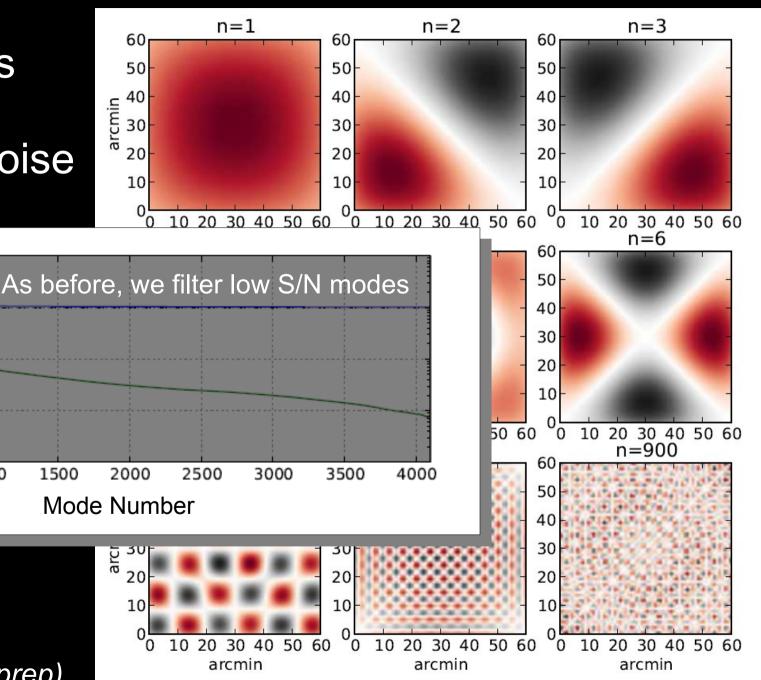
10°

10⁻¹

10-2

10⁻³

Signal / Noise



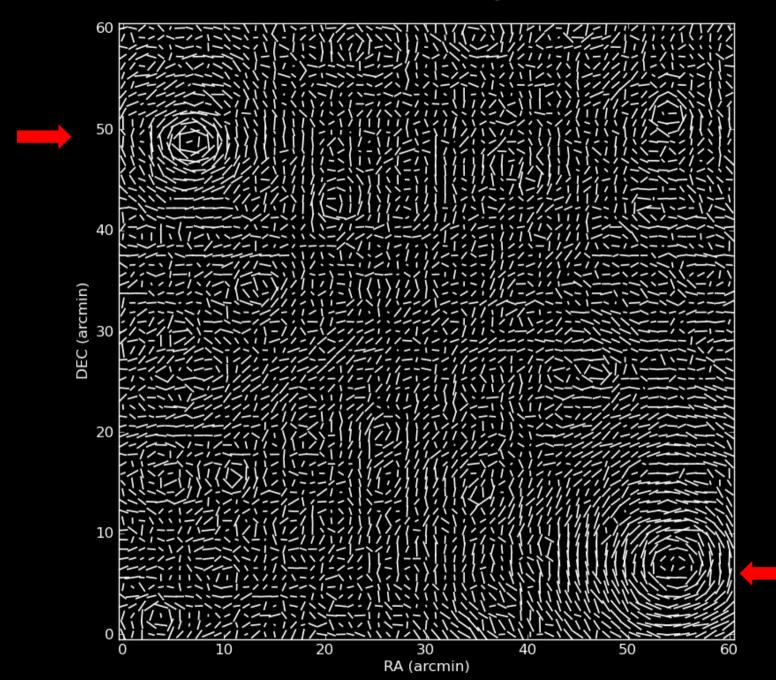
500

1000

1500

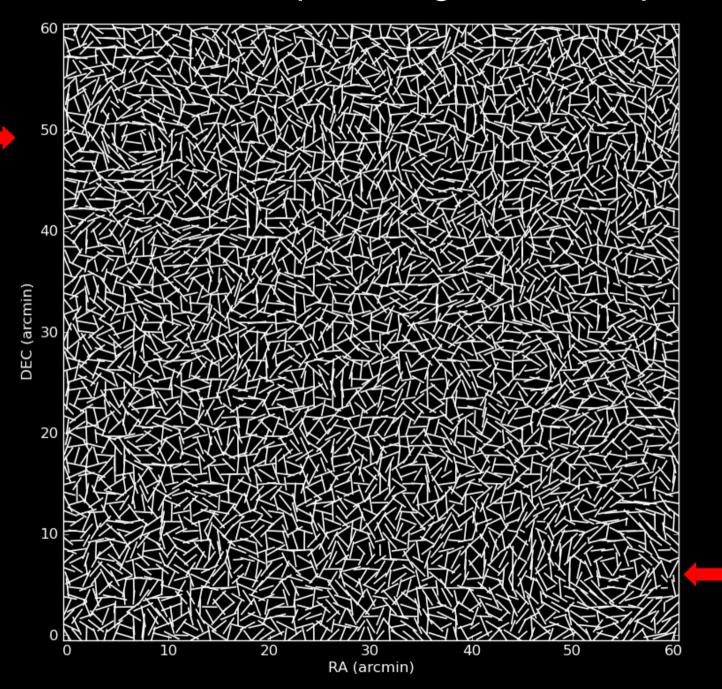
Noiseless Shear

1 square degree, 64x64 pixels



Noisy Shear

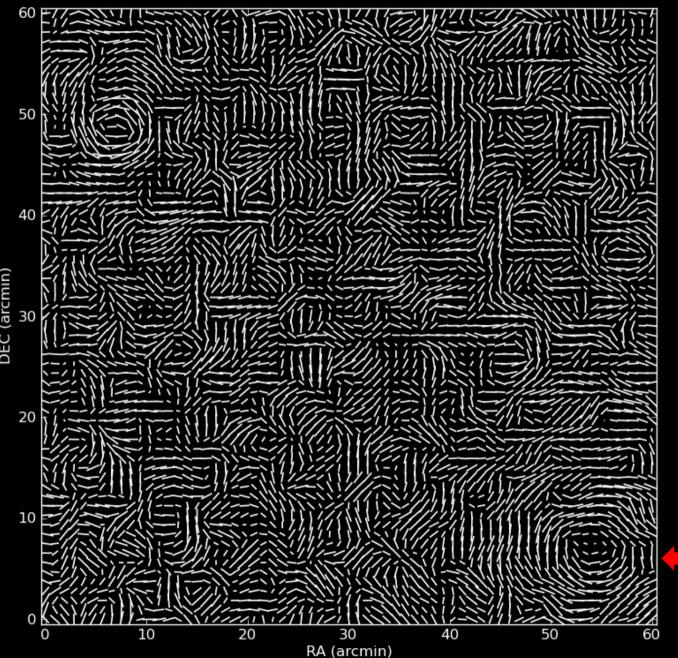
1 square degree, 64x64 pixels



900/4096 KL modes,

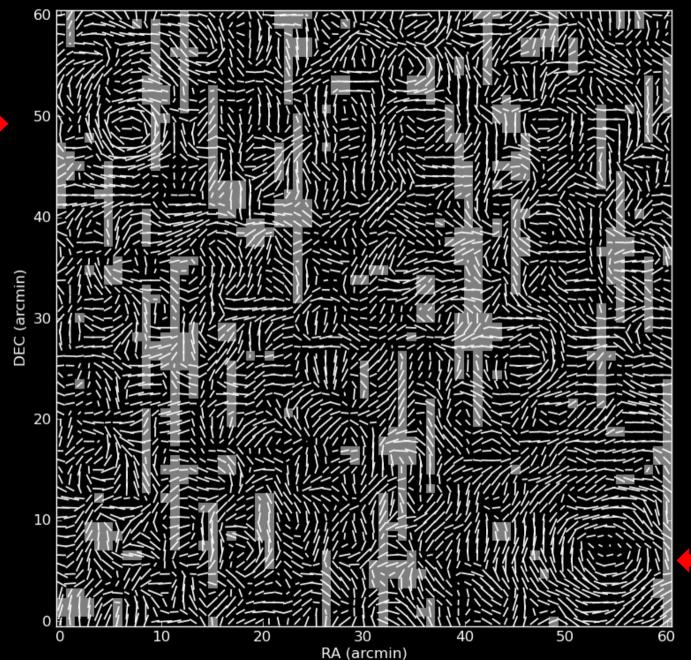
1 square degree, 64x64 pixels

unmasked



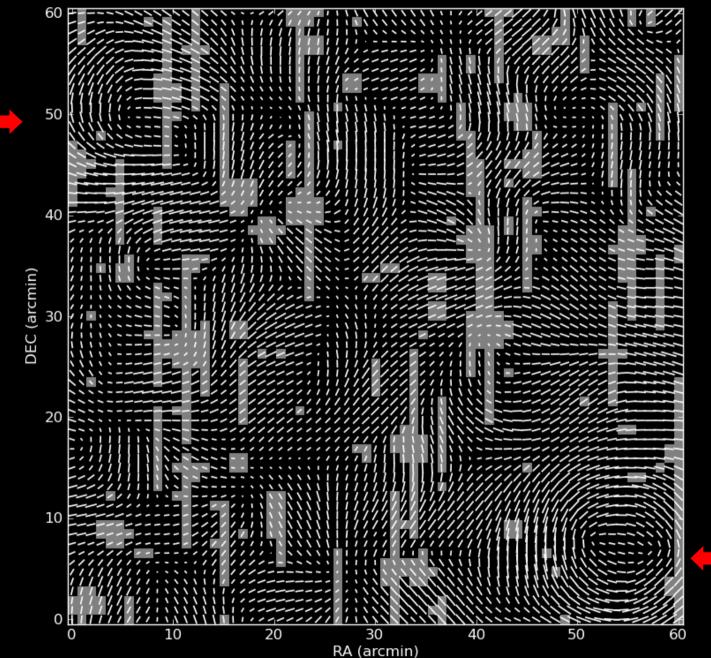
900/4096 KL modes, 20% mask

1 square degree, 64x64 pixels



100/4096 KL modes, 20% mask

1 square degree, 64x64 pixels

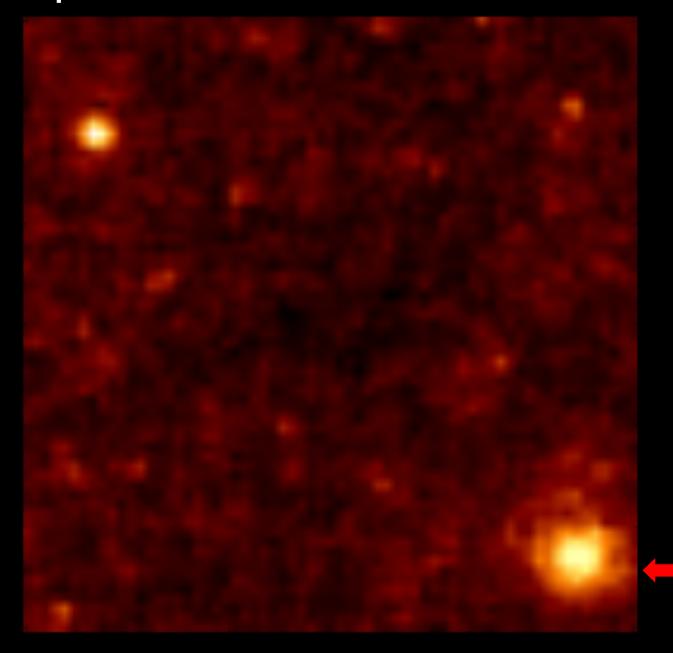


2000/4096 KL modes, ¹ square degree, 64x64 pixels 20% mask

RA (arcmin)

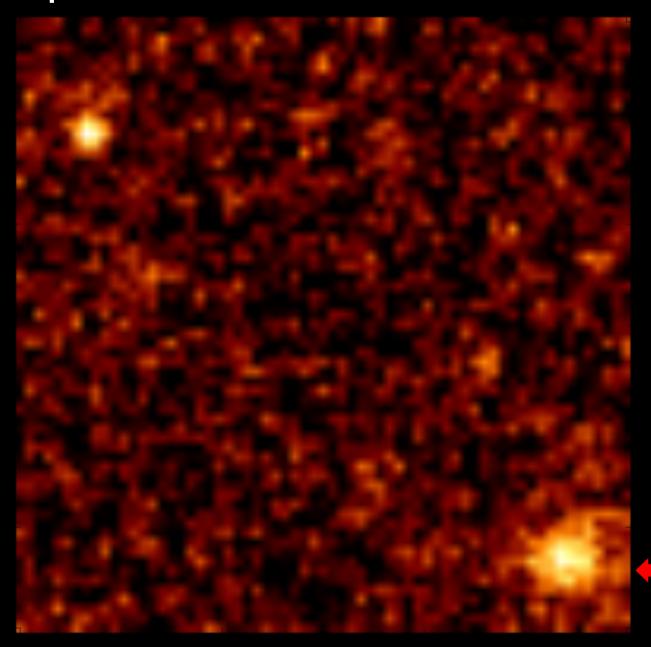
1 square degree, 64x64 pixels

noiseless



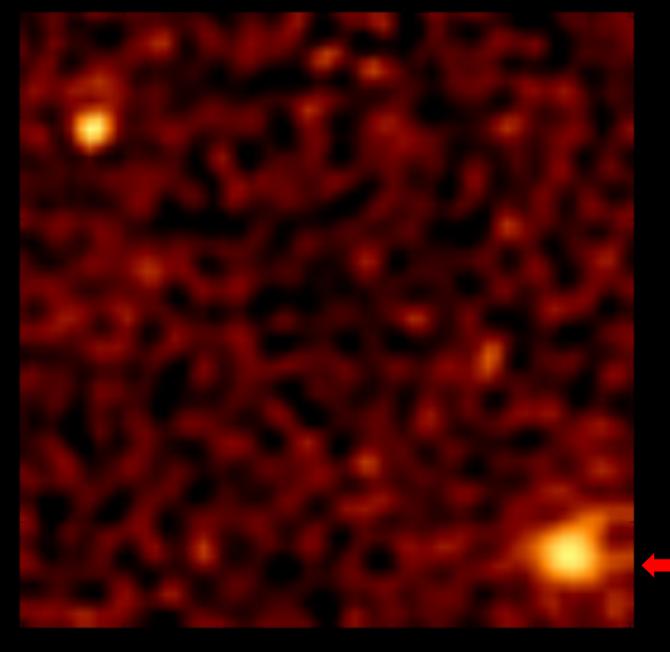
1 square degree, 64x64 pixels

noisy



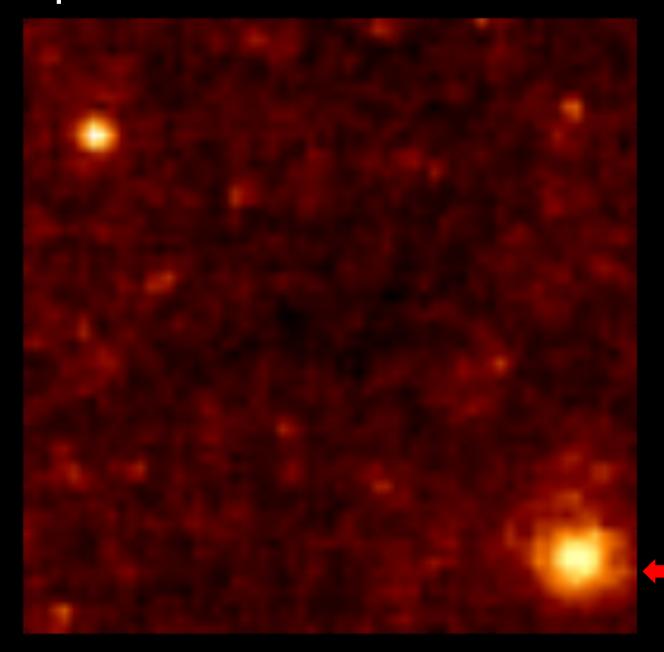
1 square degree, 64x64 pixels

900 modes

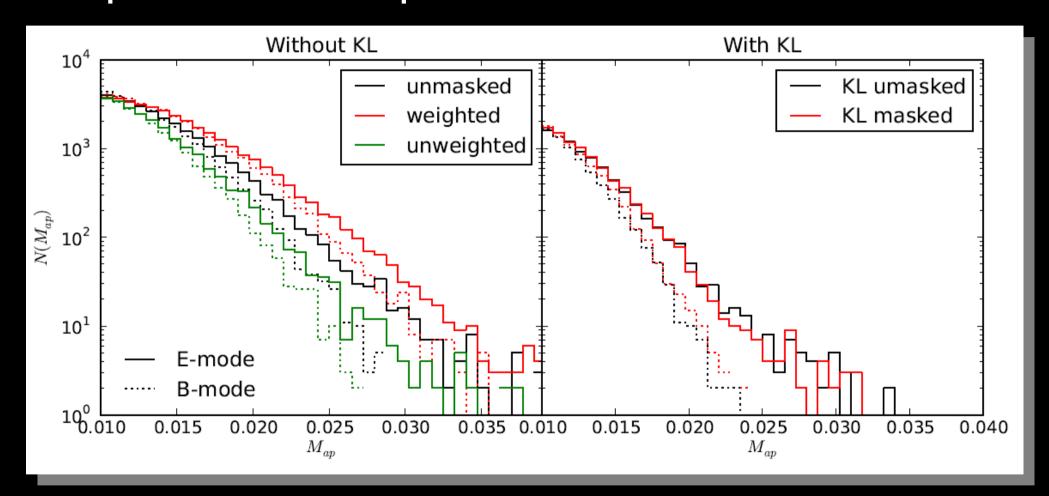


1 square degree, 64x64 pixels

noiseless



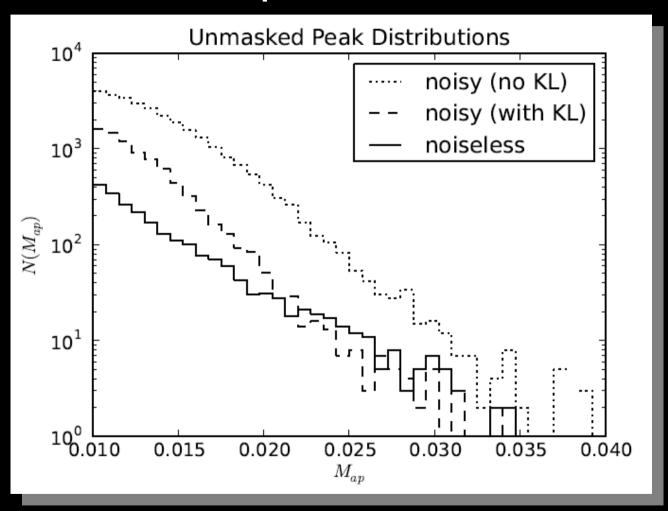
Quantitative comparison: Aperture mass peak distribution



KL recovers the unmasked peak distribution

Quantitative comparison:

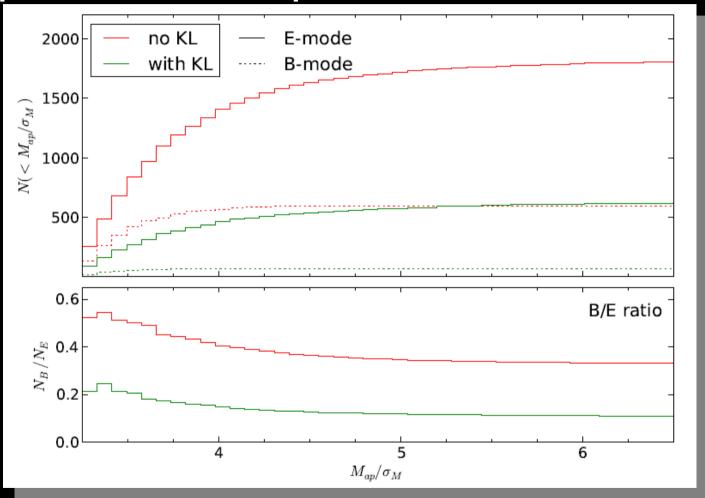
Aperture mass peak distribution



Addition of noise adds a factor of ~3 more peaks. KL filtering reduces this number.

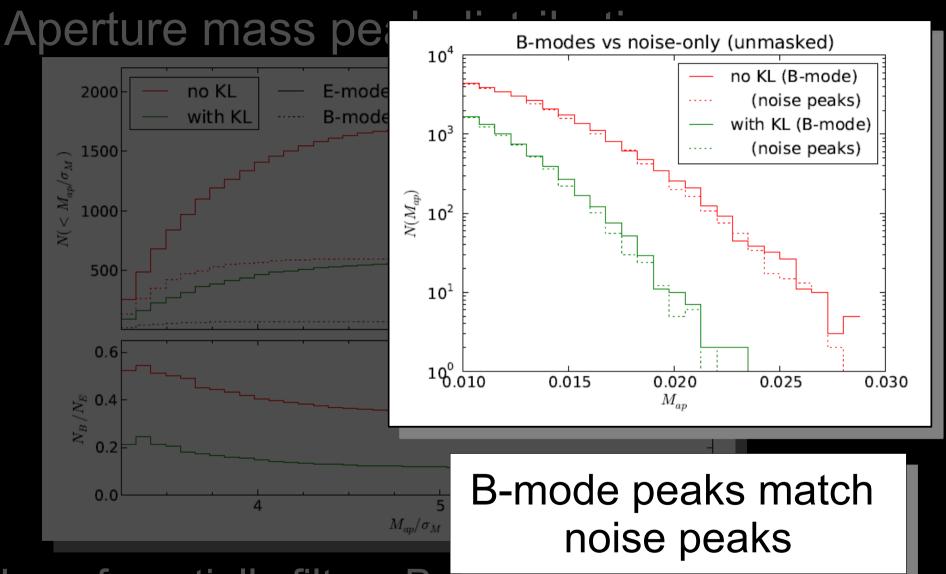
Quantitative comparison:

Aperture mass peak distribution



KL preferentially filters B-modes, leading to a factor of 3 relative reduction

Quantitative comparison:



KL preferentially filters B-rhoues, reading to a factor of 3 relative reduction

Conclusions

Using statistics of shear peaks can evade some pitfalls of cluster cosmology.

The KL method robustly interpolates between masked regions of the shear field.

The KL method reduces the statistical error in the peak function: measured by total number of peaks, and B/E ratio.

This suggests that the KL method could improve cosmological constraints from shear peaks: further study is needed to quantify this.

Very Briefly...

Modified Gravity and Dwarf Galaxies

(preliminary results of work over the last ~2 weeks with Bhuvnesh Jain)

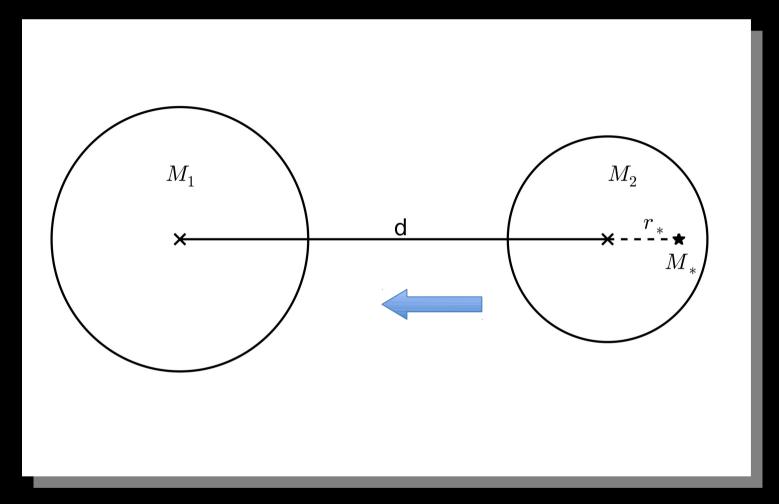
Many modified gravity theories [e.g. f(R)] involve scalar fields that provide an attractive, fifth-force

BUT... GR must be restored in the Milky Way - via "natural" mechanisms that work for massive/dense objects.

(Khoury & Weltman 2004; Vainshtein 1972)

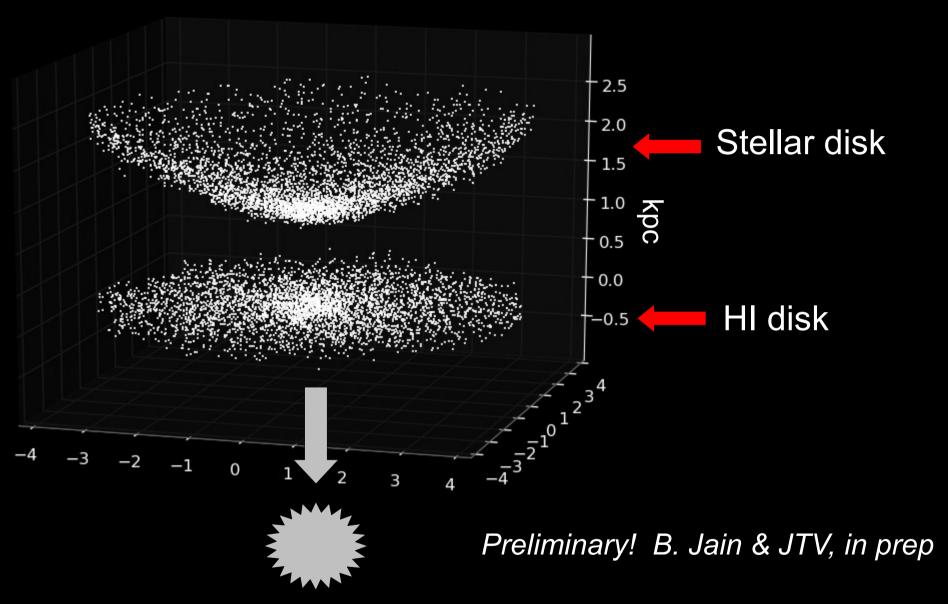
So small galaxies or the outer regions of big galaxy/cluster halos may show deviations from GR. (Kesden & Kamionkowski 2006; Hui et al 2009; Chang & Hui 2011; Davis et al 2011)

The Situation: Colliding dwarf galaxies:



Preliminary! B. Jain & JTV, in prep

- Unscreened HI disk tracks dark matter
- Self-screened stellar disk can be offset and distorted

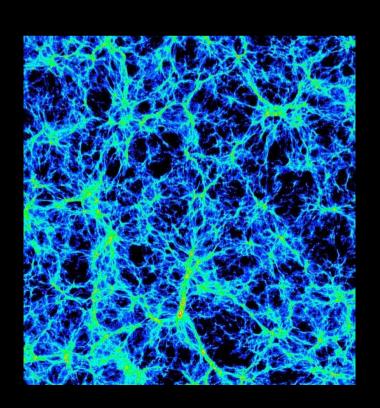


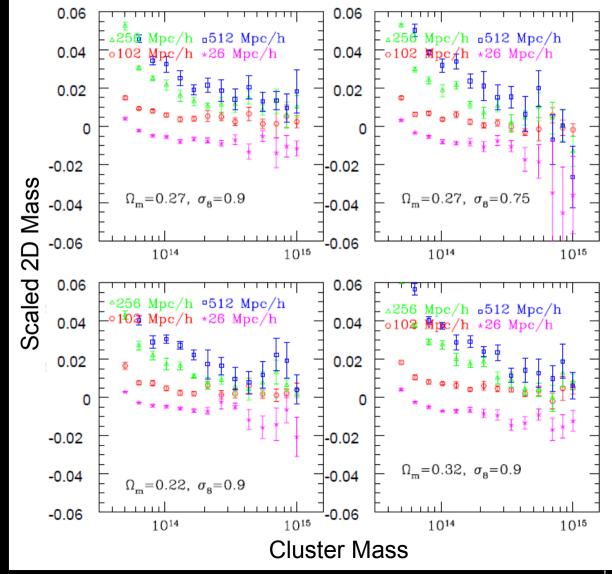
- Unscreened HI disk tracks dark matter
- Self-screened stellar disk can be offset and distorted

Followup: Marian et al. 2010

Slab Thickness: 206 Mpc/h 512 Mpc/h 102 Mpc/h 26 Mpc/h

"Correlated Projections" can affect the projected mass

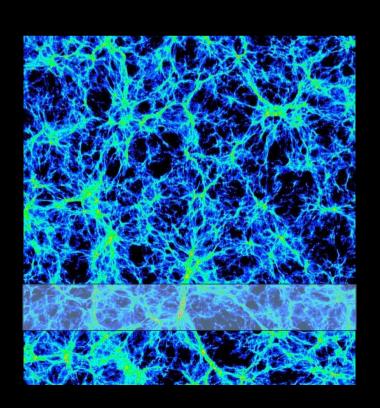


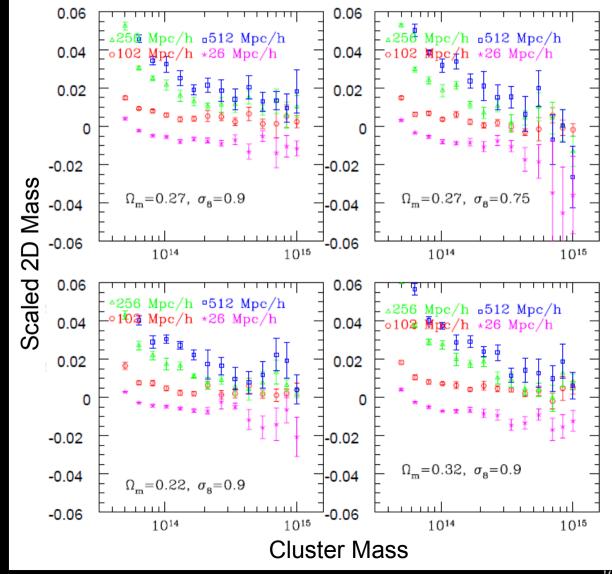


Followup: Marian et al. 2010

Slab Thickness: 206 Mpc/h 512 Mpc/h 102 Mpc/h 26 Mpc/h

"Correlated Projections" can affect the projected mass

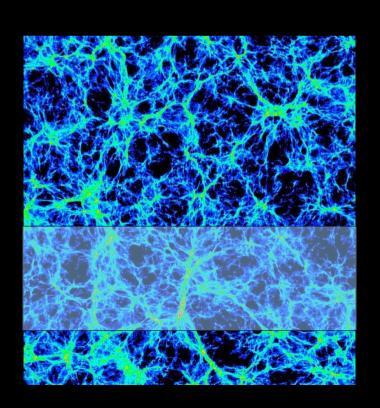


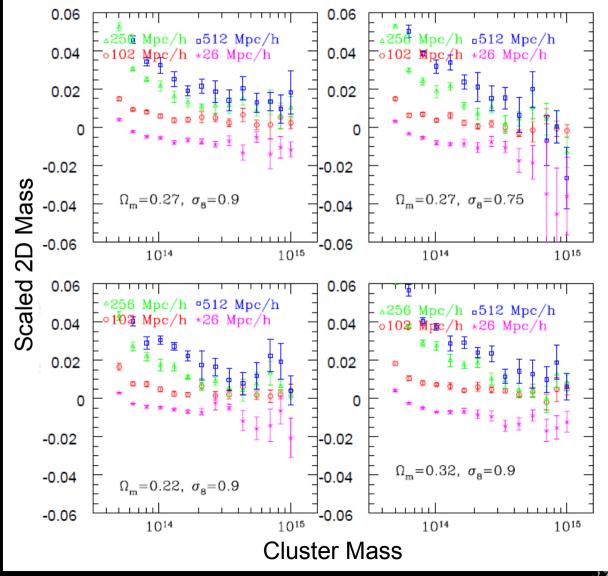


Followup: Marian et al. 2010

Slab Thickness: 206 Mpc/h 512 Mpc/h 102 Mpc/h 26 Mpc/h

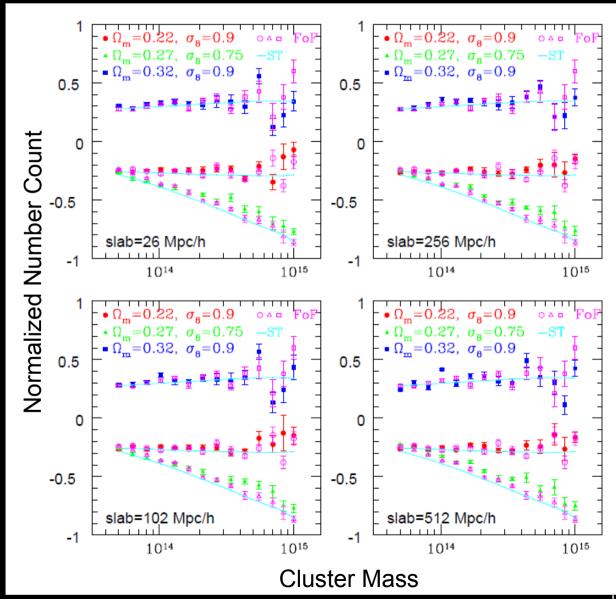
"Correlated Projections" can affect the projected mass



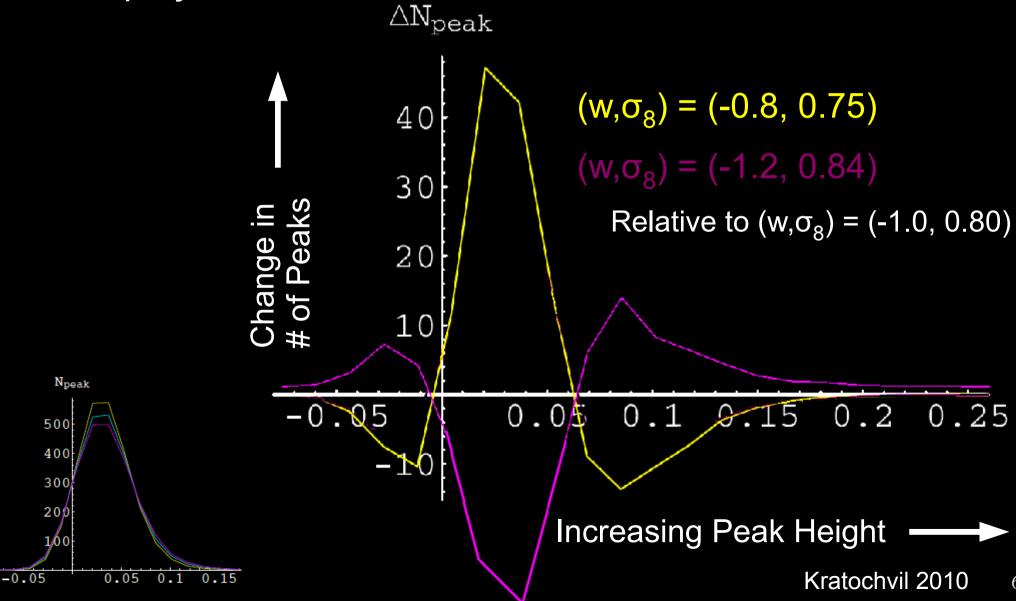


Followup: Marian et al. 2010

But cosmological scaling is unaffected



Is there information in these correlated projections?

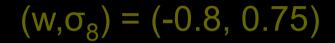


Projections boost the signal of small peaks, and these carry information ΔN_{peak}



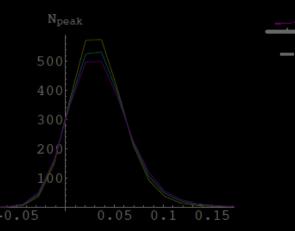
Large Peaks scale with cosmology as expected:

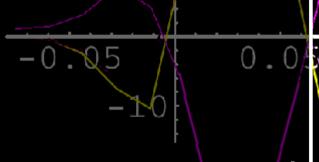
More negative w leads to later dark energy turn-on, and so more large peaks

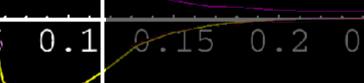


$$(w,\sigma_8) = (-1.2, 0.84)$$

Relative to $(w, \sigma_8) = (-1.0, 0.80)$

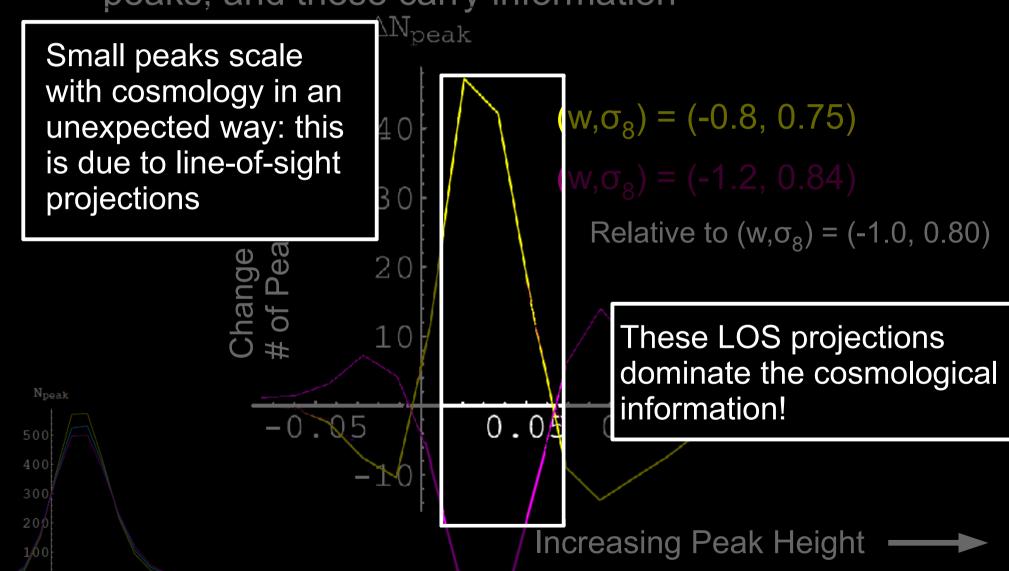


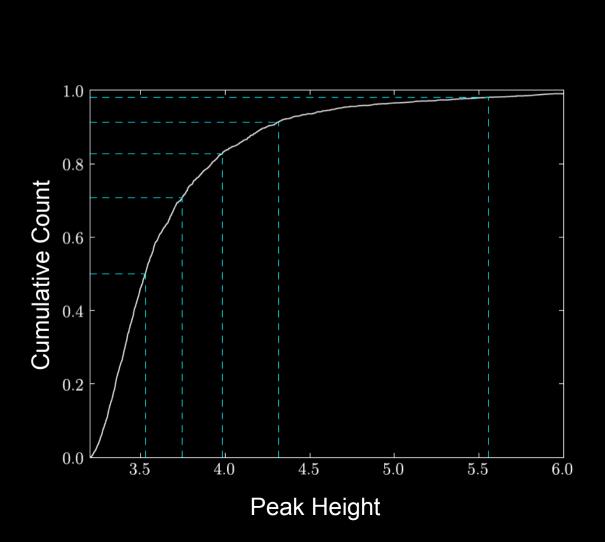


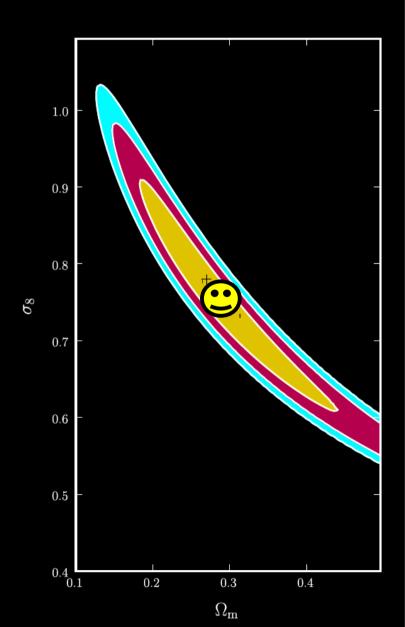


Ir creasing Peak Height

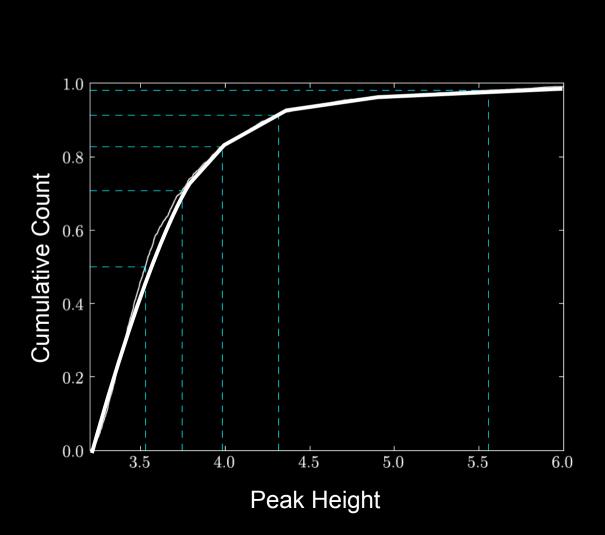
Projections boost the signal of small peaks, and these carry information

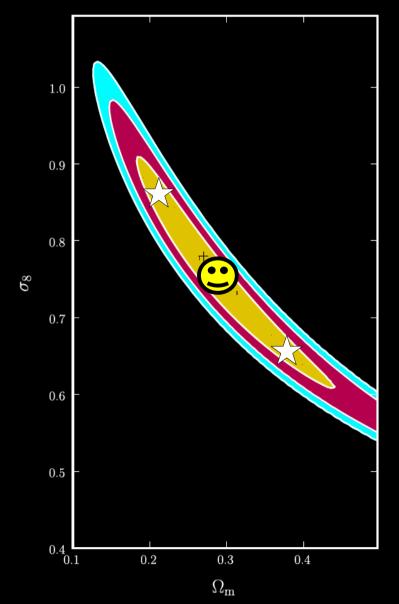






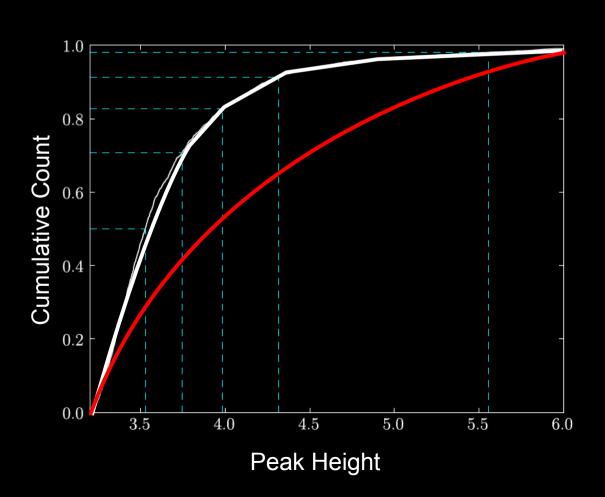
Increased normalization + less matter Decreased normalization + more matter

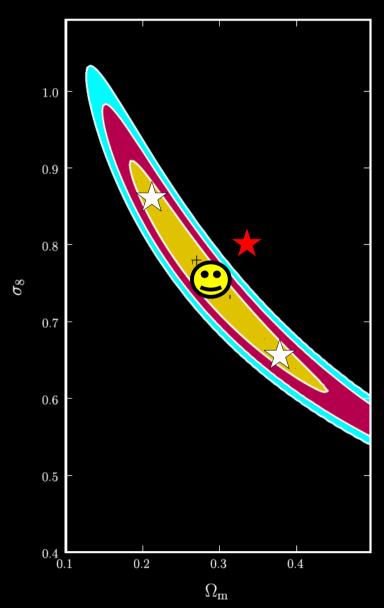




Increased normalization + less matter Decreased normalization + more matter

Too many large peaks





Increased normalization + less matter
Decreased normalization + more matter

Too many large peaks
Too many small peaks

