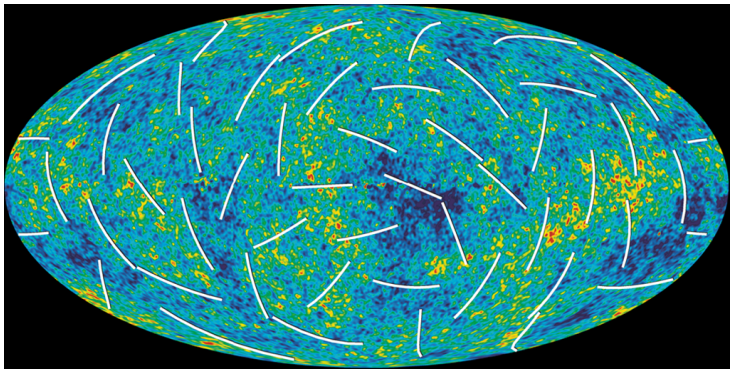


CMB polarization: past, present, and future



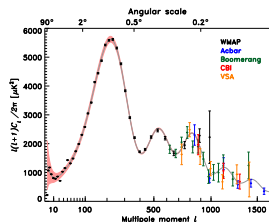
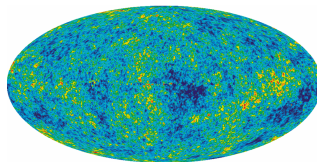
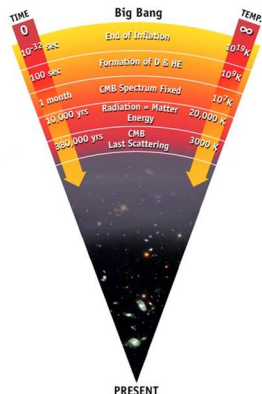
Kendrick Smith
University of Chicago
February 2007, Berkeley

Outline

- I Overview; what can be learned from CMB polarization?
- II CAPMAP (2003-2007)
- III QUIET (2007-)
- IV E-B separation

What have we learned from the CMB temperature?

Primary CMB temperature signal: snapshot of acoustic oscillations at recombination ($z \sim 1100$).



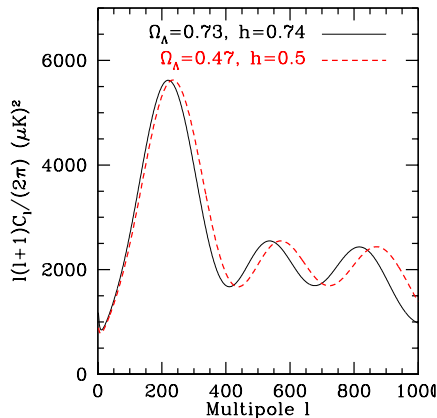
What have we learned from the CMB temperature?

1. Angular scale of acoustic peaks:

$$\ell_a = \pi \frac{D_*}{s_*}$$

← Angular diameter distance to recombination

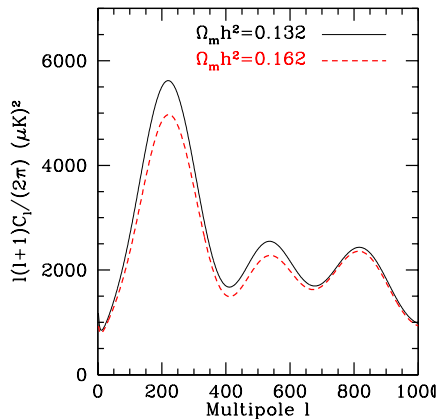
← Distance sound travels before recombination



What have we learned from the CMB temperature?

2. Radiation-matter ratio:

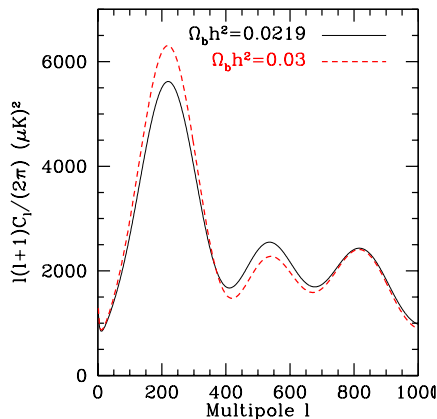
$$r_* = \left(\frac{\rho_r}{\rho_m} \right)_{a_*} \propto (\Omega_m h^2)^{-1}$$



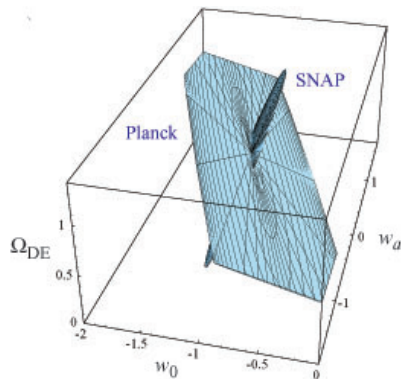
What have we learned from the CMB temperature?

3. Photon-baryon ratio:

$$R_* = \left(\frac{3\rho_b}{4\rho_\gamma} \right)_{a_*} \propto (\Omega_b h^2)$$



Importance of the CMB prior



Hu, Huterer & Smith, ApJL (2006)

CMB constraints on

$$\{\Omega_m h^2, \Omega_b h^2, D_*\}$$

are important for every flavor of cosmological data!

(Supernova example shown here)

In addition, C_ℓ^{TT} contains information about the shape of the initial power spectrum (tilt n_s , running α , etc.)

CMB polarization: E-B decomposition

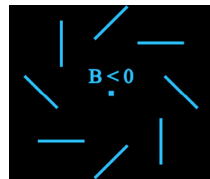
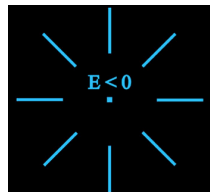
CMB polarization can be decomposed into...

E-modes: $\Pi_{ab} = (\nabla_a \nabla_b - \frac{1}{2} \delta_{ab} \nabla^2) \phi$.

- ▶ Dominant component: generated by first-order perturbations at recombination and reionization.

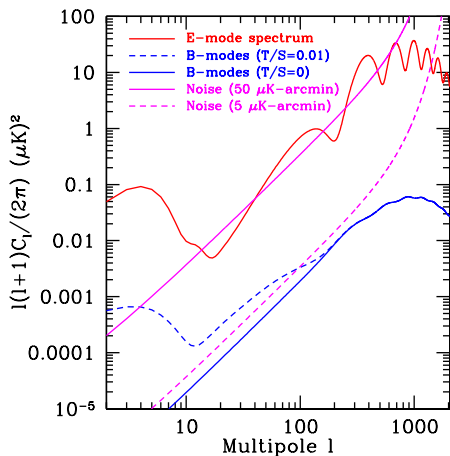
B-modes: $\Pi_{ab} = (\frac{1}{2} \epsilon_a^c \nabla_b \nabla_c + \frac{1}{2} \epsilon_b^c \nabla_a \nabla_c) \phi$.

- ▶ Generated by second-order effects (mainly gravitational lensing of the larger E-mode signal).
- ▶ Also generated by **gravitational waves** from inflation.



This is the spin-2 analogue of the gradient/curl decomposition for a vector field.

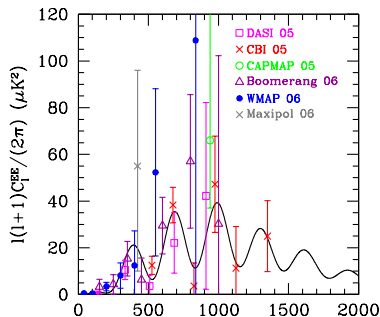
CMB polarization: E-B decomposition



- ▶ EE and BB power spectra shown, with noise power spectra for comparison ($\theta_{\text{FWHM}} = 10$ arcmin).
- ▶ Current experiments ($\sim 50 \mu K\text{-arcmin}$) have detected EE to 10σ , with upper limits on BB.
- ▶ Future ground-based experiments ($\sim 5 \mu K\text{-arcmin}$) should make precision measurements of EE and detect B-modes.

CMB polarization: acoustic peaks (C_ℓ^{EE})

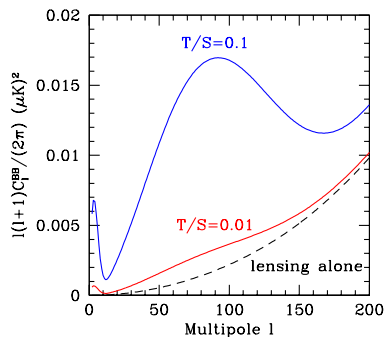
In principle, measuring acoustic peaks in C_ℓ^{EE} can improve errors on $\{\Omega_m h^2, \Omega_b h^2, D_*\}$ by a factor of $\sim \sqrt{2}$. However, for signal-to-noise reasons, C_ℓ^{EE} probably best regarded as predicted by $C_\ell^{TT} \dots$



Exception: polarization is complementary to temperature when estimating the primordial power spectrum $P^{\zeta\zeta}(k)$ (Hu & Okamoto 2003).

Exception: isocurvature modes (Bucher, Moodley & Turok 2000)

CMB polarization: gravity waves (C_ℓ^{BB} at low ℓ)



Low B-mode multipoles ultimately have more sensitivity to gravity waves in the early universe (parameterized by tensor-to-scalar ratio T/S) than any other type of data.

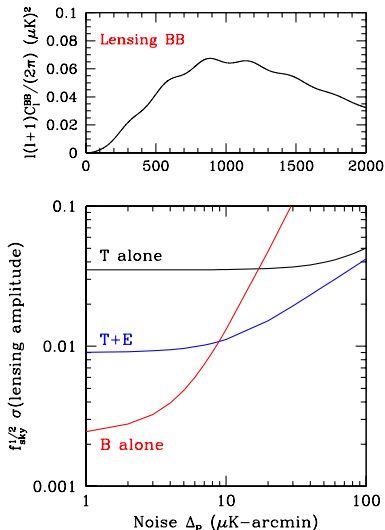
Well-motivated in some inflationary models; observation would rule out others (e.g. ekpyrotic).

CMB polarization: gravitational lensing

“Guaranteed” B-mode signal from gravitational lensing of the larger E-mode signal.

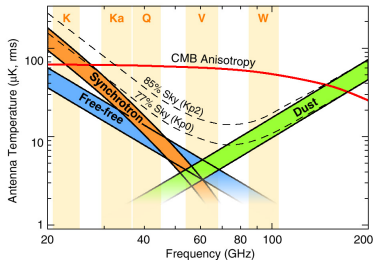
Probes new parameters, e.g. overall lensing amplitude depends on $\sum m_\nu$ by 50% per eV (Eisenstein & Hu 1997)

Gravitational lensing appears in temperature and polarization, but polarization is more sensitive.



Smith, Hu & Kaplinghat, PRD (2005)

Foregrounds

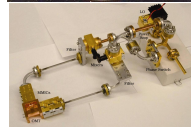


Good frequency coverage is essential for controlling foregrounds

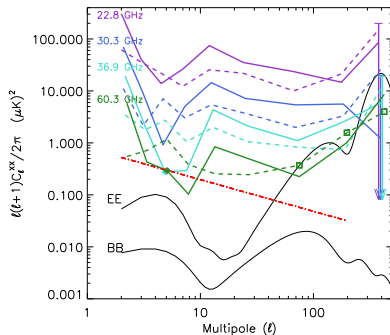
Bennett et al (2003)

Bolometers ($\nu \geq 100$ GHz)
BICEP, Clover, EBEX, Spider

Coherent detectors ($\nu \leq 90$ GHz)
QUIET



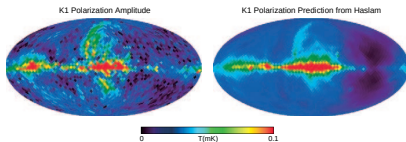
Foregrounds



Taken over large regions of sky, foregrounds are comparable to or larger than expected CMB polarization signals

However...

To date, polarized foregrounds have not been a significant contaminant for ground-based experiments sampling “clean” patches of sky.



Page et al (2006)

Part II: CAPMAP

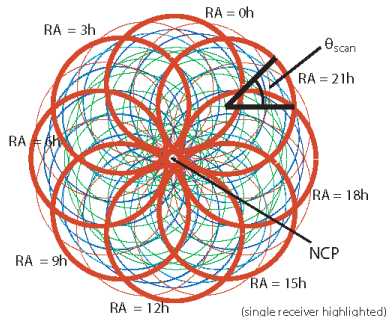


Coherent polarimeters + 7m telescope
(Crawford Hill, NJ)

First (2003, 4 detectors) observing
season resulted in 2σ detection of EE
([Barkats et al 2005](#))

Second (2005, 16 detectors) observing
season: [soon!](#)

CAPMAP: focal plane



Jeff McMahon

12 W-band (90 GHz) detectors + 4 Q-band (40 GHz) detectors
Scan small patch near NCP

Total sky area: 7.3 deg² (90 GHz), 9.2 deg² (40 GHz)

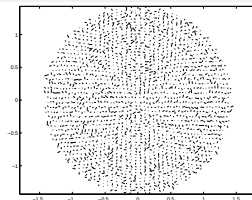
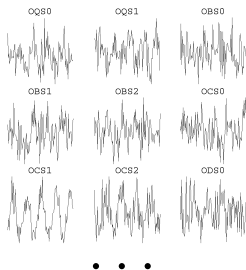
CAPMAP: summary

	ν	N_{det}	ΔT	Area	Noise	θ_{FWHM}
CAPMAP I	90 GHz	4	433 hr	2.0 deg ²	100 μ K'	4'
CAPMAP II	40 GHz	4	900 hr	9.2 deg ²	70 μ K'	6'
	90 GHz	12	900 hr	7.3 deg ²	60 μ K'	3.5'

Two independent analysis pipelines:

1. **Pipeline 1:** represent data in real space (θ, ϕ), noise as dense covariance matrix ($N_{pix} \sim 2500$ for beam size pixels, ~ 10000 for 1/2 beam size).
2. **Pipeline 2:** treat noise as azimuthally symmetric and Fourier transform in ϕ (coordinates are now θ, m); big speedup since covariance matrices become block diagonal in m (can easily go to 1/3 beam size)

CAPMAP: analysis pipeline



$$\begin{pmatrix} * & * & \cdots \\ * & * & \cdots \\ \vdots & \vdots & \ddots \end{pmatrix}$$

Timestreams (+ pointing)

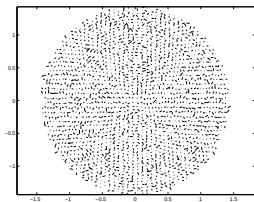
Map + noise covariance

$$N_{map}^{-1} = P N_{tod}^{-1} P^T \quad ; \quad N_{map}^{-1} m = P^T N_{tod}^{-1} t$$

Timestream noise model: white + marginalization over lowest 5 ring modes + ground-synchronous signal.

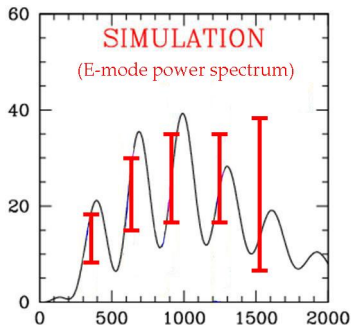
(Simulations show that this filter sufficiently removes $1/f$ noise.)

CAPMAP: analysis pipeline



$$\begin{pmatrix} * & * & \cdots \\ * & * & \cdots \\ \vdots & \vdots & \ddots \end{pmatrix}$$

Map + noise covariance

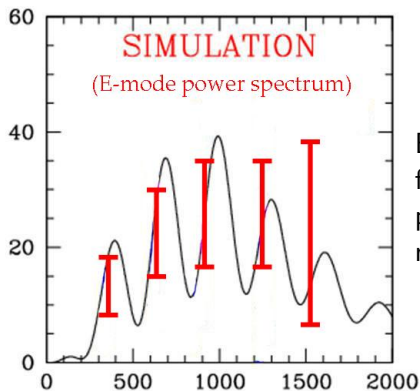


Power spectrum likelihood

Likelihood can be explored in different ways (Newton-Raphson, Markov chain, ...) but each “step” requires a dense matrix operation (Bond, Jaffe & Knox 1997)

Use S/N eigenmode compression to reduce matrix size ($N \sim 2000$ independent of pixel size) (Bond 1994)

CAPMAP: Fisher matrix from full analysis pipeline



Bandpower Fisher matrix computed from noise covariance matrix in real pipeline (equivalent to MC average of many simulations).

Pipeline also includes **null test suite**: difference two maps made from disjoint data subsets, analyze power spectrum and compare to zero. (When applied to the two frequency channels, this is a strong test for foregrounds.)

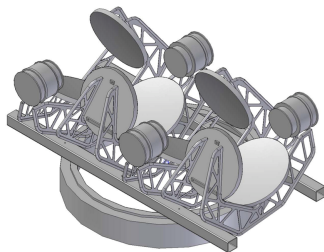
Part III: QUIET (Q/U Imaging ExperimentT)



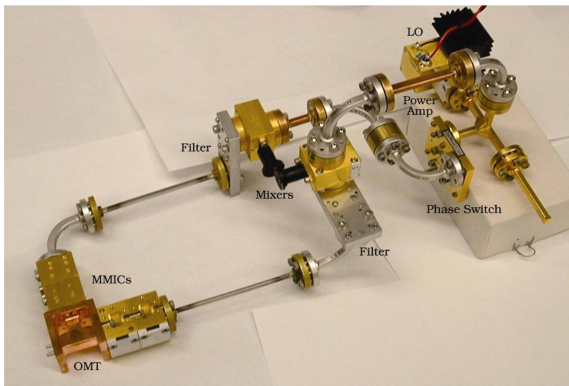
Atacama desert, Chile

Phase I (late 2007):
2m telescope, ~ 100 detectors

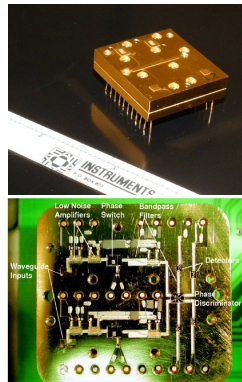
Phase II (2010):
2m telescope, ~ 1000 detectors
7m telescope, ~ 500 detectors



QUIET: "Coherent polarimeter on a chip"

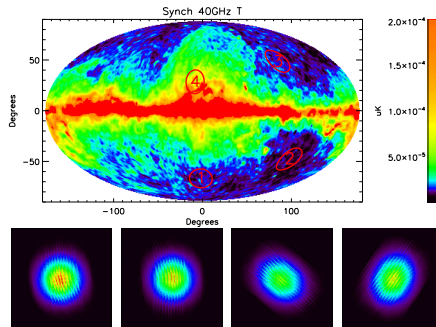


CAPMAP 90GHz polarimeter



QUIET 90GHz module

QUIET: Scan strategy



Four patches selected for low foreground contamination, distribution around SCP roughly uniform in RA.

After many repointings, get roughly isotropic coverage and good cross-linking

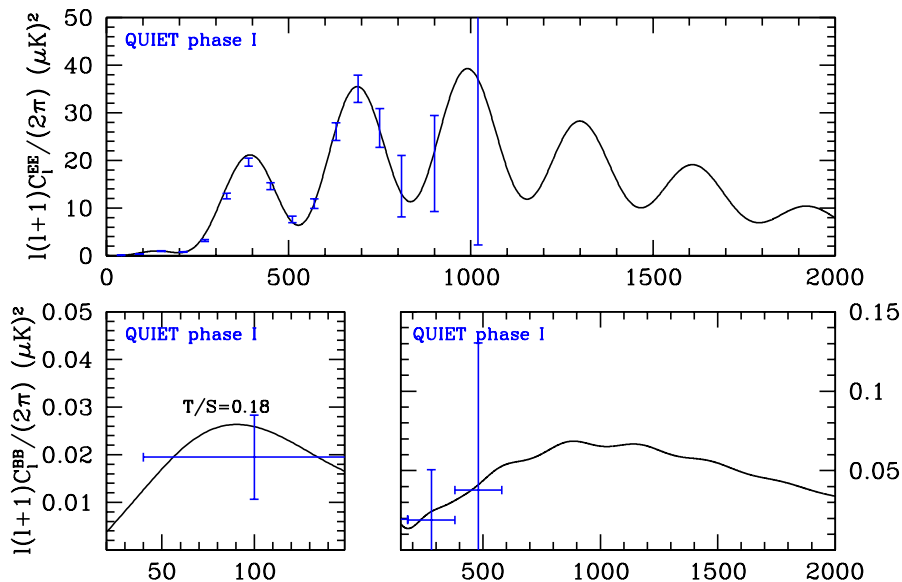
Keith Vanderlinde

QUIET: summary

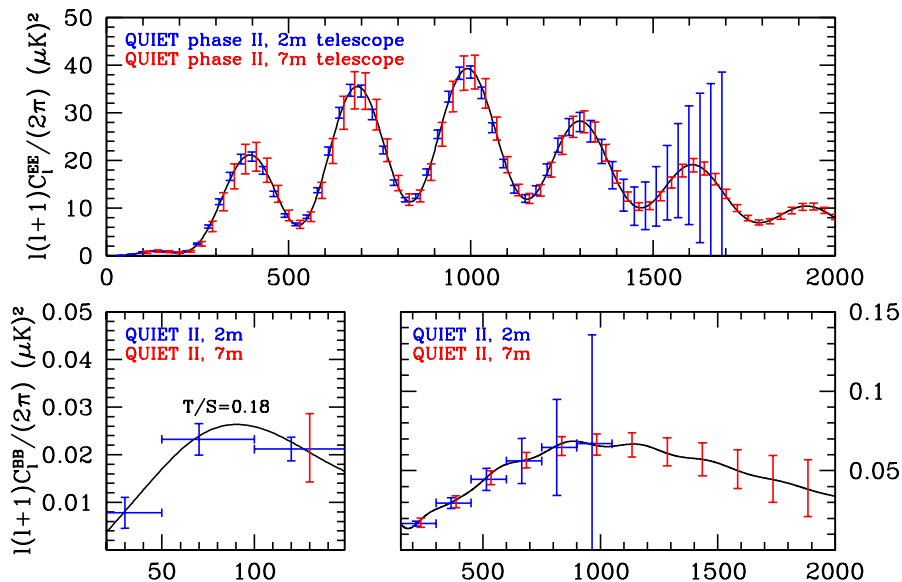
	ν	N_{det}	ΔT	Area	Noise	θ_{FWHM}
CAPMAP I	90 GHz	4	433 hr	2.0 deg ²	100 $\mu K'$	4'
CAPMAP II	40 GHz	4	900 hr	9.2 deg ²	70 $\mu K'$	6'
	90 GHz	12	900 hr	7.3 deg ²	60 $\mu K'$	3.5'
QUIET I	40 GHz	19	4000 hr	1600 deg ²	49 $\mu K'$	10'
	90 GHz	83	4000 hr	1600 deg ²	34 $\mu K'$	10'
QUIET II (2m)	40 GHz	166	4000 hr	1600 deg ²	9.6 $\mu K'$	10'
	90 GHz	714	4000 hr	1600 deg ²	7.3 $\mu K'$	10'
QUIET II (7m)	40 GHz	83	8000 hr	160 deg ²	3.2 $\mu K'$	6'
	90 GHz	357	8000 hr	160 deg ²	2.5 $\mu K'$	3.5'

Forecasting methodology: Compute dense matrix at low ℓ (including scan strategy and mode removal); use simple f_{sky} scaling at high ℓ .

QUIET: phase I forecasts (90 GHz alone)



QUIET: phase II forecasts (90 GHz alone)



QUIET phase I: analysis pipeline

Problem size:

$$N_{\text{tod}} = (2)(91)(30 \text{ Hz})(10^7 \text{ sec}) = 5.4 \times 10^{10}$$

$$N_{\text{pix}} = (3)(f_{\text{sky}})(12N_{\text{side}}^2) = 3.8 \times 10^5 \quad (N_{\text{side}} = 1024)$$

Unlike CAPMAP, fully optimal analysis seems prohibitive.

Exploring Monte Carlo based alternatives:

- ▶ Map-making: destriping
- ▶ Map-making: MASTER approach (high-pass + binning)
- ▶ Power spectrum estimation: pseudo- C_ℓ

Computational cost: probably $\sim 10^4$ CPU-hours for a full analysis

Running problem: **E-B mixing**

Part IV: E-B mixing

How are E-mode and B-mode power spectra estimated from data?

- ▶ Under the simplifying assumption of all-sky isotropic noise, estimating C_ℓ^{EE} , C_ℓ^{BB} from a noisy map Π_{ab} is straightforward:

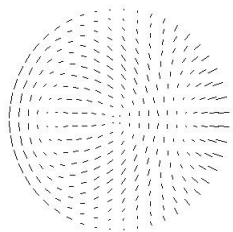
$$a_{\ell m}^E = \int d^2x \Pi^{ab}(x) Y_{(\ell m)ab}^E(x) \quad a_{\ell m}^B = \int d^2x \Pi^{ab}(x) Y_{(\ell m)ab}^B(x)$$

$$\hat{C}_\ell^{EE} = \frac{1}{2\ell+1} \sum_{m=-\ell}^{\ell} a_{\ell m}^E a_{\ell m}^{E*} \quad \hat{C}_\ell^{BB} = \frac{1}{2\ell+1} \sum_{m=-\ell}^{\ell} a_{\ell m}^B a_{\ell m}^{B*}$$

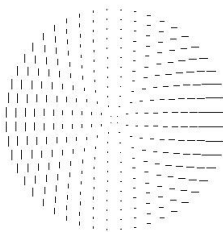
- ▶ The E-mode and B-mode spherical harmonics $Y_{\ell m}^E(x)$, $Y_{\ell m}^B(x)$ provide a complete basis for E-mode and B-mode power on the sky, and are decoupled from each other.

What is the E-B mixing problem?

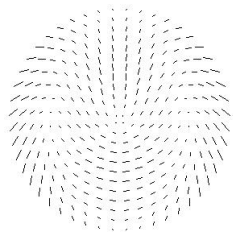
- ▶ In the presence of sky cuts or inhomogeneous noise, the E/B decomposition becomes more complicated:



pure E-mode



ambiguous mode

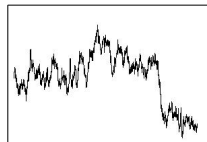


pure B-mode

- ▶ Pure E-modes and B-modes are expensive to compute directly
- ▶ **Challenge:** Find a B-mode estimator which only receives contributions from pure B-modes, and is computationally fast.

Pseudo- C_ℓ power spectrum estimation: 1-D analogy

- ▶ Start with timeseries on a finite interval:
- ▶ Take FFT, and estimate power spectra, as if the timeseries were periodic on a larger interval:



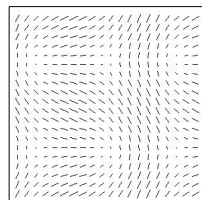
$$\begin{aligned}\tilde{f}(\omega) &= \int dt f(t) W(t) \exp(i\omega t) \\ \tilde{P}(\omega) &= |\tilde{f}(\omega)|^2\end{aligned}$$

- ▶ Final step: fix up the normalization

$$\hat{P}(\omega) = \frac{1}{K} \tilde{P}(\omega)$$

Pseudo- C_ℓ power spectrum estimation: 2-D version

- ▶ Start with CMB polarization on a finite patch:
- ▶ Take spherical transform, and estimate power spectrum, as if the polarization were defined all-sky:



$$\tilde{a}_{\ell m}^E = \int \Pi^{ab}(x) W(x) Y_{(\ell m)ab}^{E*}(x) \quad \tilde{a}_{\ell m}^B = \int \Pi^{ab}(x) W(x) Y_{(\ell m)ab}^{B*}(x)$$

$$\tilde{C}_\ell^{EE} = \frac{1}{2\ell+1} \sum_{m=-\ell}^{\ell} \tilde{a}_{\ell m}^E \tilde{a}_{\ell m}^{E*} \quad \tilde{C}_\ell^{BB} = \frac{1}{2\ell+1} \sum_{m=-\ell}^{\ell} \tilde{a}_{\ell m}^B \tilde{a}_{\ell m}^{B*}$$

- ▶ Final step: debias

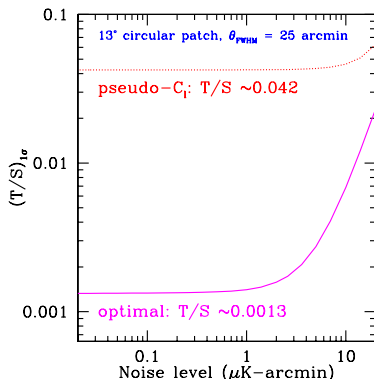
$$\begin{pmatrix} \hat{C}_\ell^{EE} \\ \hat{C}_\ell^{BB} \end{pmatrix} = \begin{pmatrix} K_{\ell\ell'}^+ & K_{\ell\ell'}^- \\ K_{\ell\ell'}^- & K_{\ell\ell'}^+ \end{pmatrix}^{-1} \begin{pmatrix} \tilde{C}_{\ell'}^{EE} - \langle \tilde{C}_{\ell'}^{EE} \rangle_{\text{noise}} \\ \tilde{C}_{\ell'}^{BB} - \langle \tilde{C}_{\ell'}^{BB} \rangle_{\text{noise}} \end{pmatrix}$$

Pseudo- C_ℓ power spectrum estimation: tradeoffs

- ▶ Pseudo- C_ℓ : suboptimal but very fast
- ▶ Practical problem: how to choose pixel weight function $W(x)$?
- ▶ For B-modes, there is a more fundamental problem: the estimator mixes E and B and therefore limits B-mode sensitivity at low noise levels.

$$\tilde{a}_{\ell m}^B = \int \Pi^{ab}(x) W(x) Y_{(\ell m)ab}^B(x)$$

Pseudo- C_ℓ power spectrum estimation: EB mixing



- Challinor & Chon (2004): For $f_{\text{sky}} \sim 0.01$, pseudo- C_ℓ limits the gravity wave signal which can be detected to $(T/S) \sim 0.05$.
- In the pseudo- C_ℓ method, $E \rightarrow B$ mixing is treated like noise: can subtract bias, but extra variance remains.

Smith (2005)

$$\begin{pmatrix} \hat{C}_\ell^{EE} \\ \hat{C}_\ell^{BB} \end{pmatrix} = \begin{pmatrix} K_{\ell\ell'}^+ & K_{\ell\ell'}^- \\ K_{\ell\ell'}^- & K_{\ell\ell'}^+ \end{pmatrix}^{-1} \begin{pmatrix} \tilde{C}_{\ell'}^{EE} - \langle \tilde{C}_{\ell'}^{EE} \rangle_{\text{noise}} \\ \tilde{C}_{\ell'}^{BB} - \langle \tilde{C}_{\ell'}^{BB} \rangle_{\text{noise}} \end{pmatrix}$$

A “pure” pseudo- C_ℓ estimator:

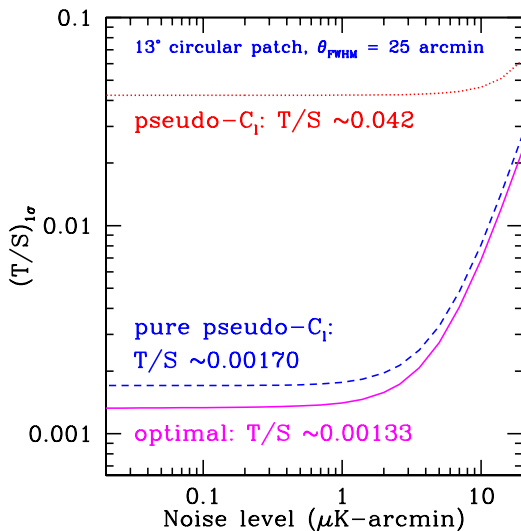
Proposal: Add higher-spin weights W_a , W_{ab} and counterterms to cancel the E-B mixing ([Smith 2005](#))

$$\begin{aligned}\tilde{a}_{\ell m}^B = & \int d^2x \Pi^{ab}(x) \left[W(x) Y_{(\ell m)ab}^{B*} \right. \\ & + \epsilon_a^c \frac{W_b(x) Y_{(\ell m)c}^{G*} + W_c(x) Y_{(\ell m)b}^{G*}}{\sqrt{(\ell-1)(\ell+2)}} \\ & \left. + \frac{\epsilon_a^c W_{bc}(x) Y_{\ell m}^*}{\sqrt{(\ell-1)\ell(\ell+1)(\ell+2)}} \right]\end{aligned}$$

Also proposed algorithmic approach for choosing $W(x)$, $W_a(x)$, $W_{ab}(x)$ ([Smith and Zaldarriaga 2006](#)).

Based on variational principle: minimize “average power” $\langle \tilde{C}_\ell \rangle$ satisfying normalization constraint $\sum_x W(x) = 1$.

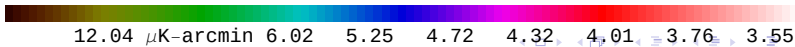
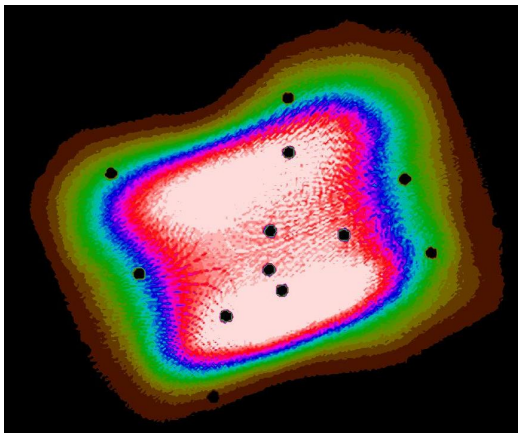
Pure pseudo- C_ℓ estimator: EB mixing



Smith (2005)

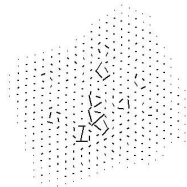
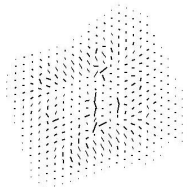
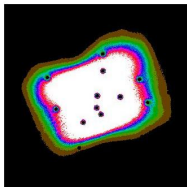
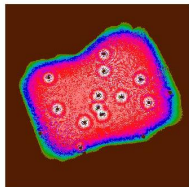
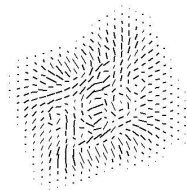
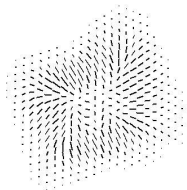
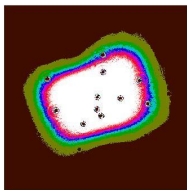
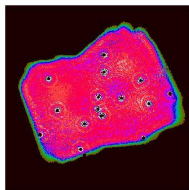
Fiducial experiment: hitcount map

- Fiducial experiment: average noise $\sim 5.75 \mu\text{K-arcmin}$, $\theta_{\text{FWHM}} = 8 \text{ arcmin}$, randomly generated point source mask, noise distribution based on preliminary EBEX simulations.

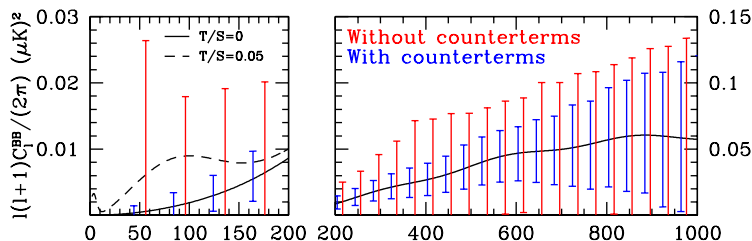


Fiducial experiment: weight functions

- ▶ Optimized weight function depends on ℓ band; shown here for $(\ell_{\min}, \ell_{\max}) = (30, 70)$ (top row) and $(510, 550)$ (bottom).
- ▶ Each “weight function” consists of four pieces (left to right): E-mode weight function $W_E(x)$, scalar piece of B-mode weight $W(x)$, and two B-mode counterterms $W_a(x)$, $W_{ab}(x)$.



Fiducial experiment: power spectrum errors



Smith & Zaldarriaga, astro-ph/0610059

- ▶ Adding counterterms significantly improves BB power spectrum errors for the fiducial experiment
- ▶ Gravity wave signal: $\sigma(T/S)$ $0.054 \rightarrow 0.0056$.
- ▶ Lensing amplitude: $\sigma(A_{\text{lens}})$ $0.258 \rightarrow 0.085$.

Concluding thoughts

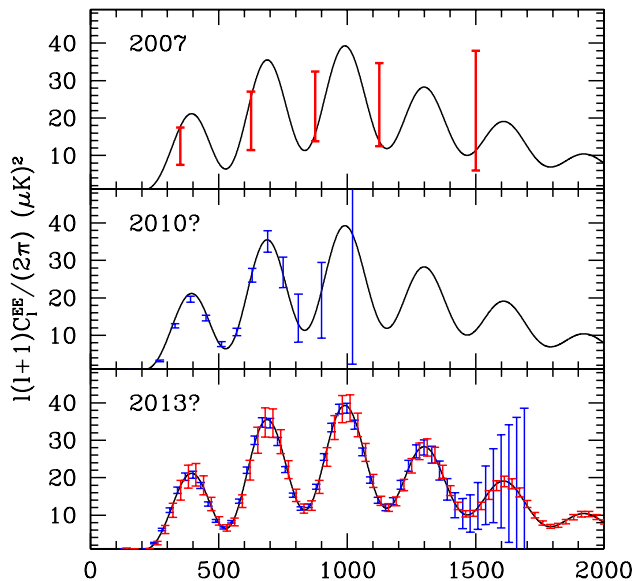
Estimator gives good E-B separation while also solving an outstanding practical problem: choosing the weight function.

Estimator performance seems convincing for “real-world” distribution of white noise, but full QUIET noise model will be more complicated ($1/f$ noise, ground synchronous modes. . .)

Related: this only solves half the EB separation problem for Monte Carlo CMB pipelines!

Map-making also mixes E and B. . .

Concluding thoughts



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