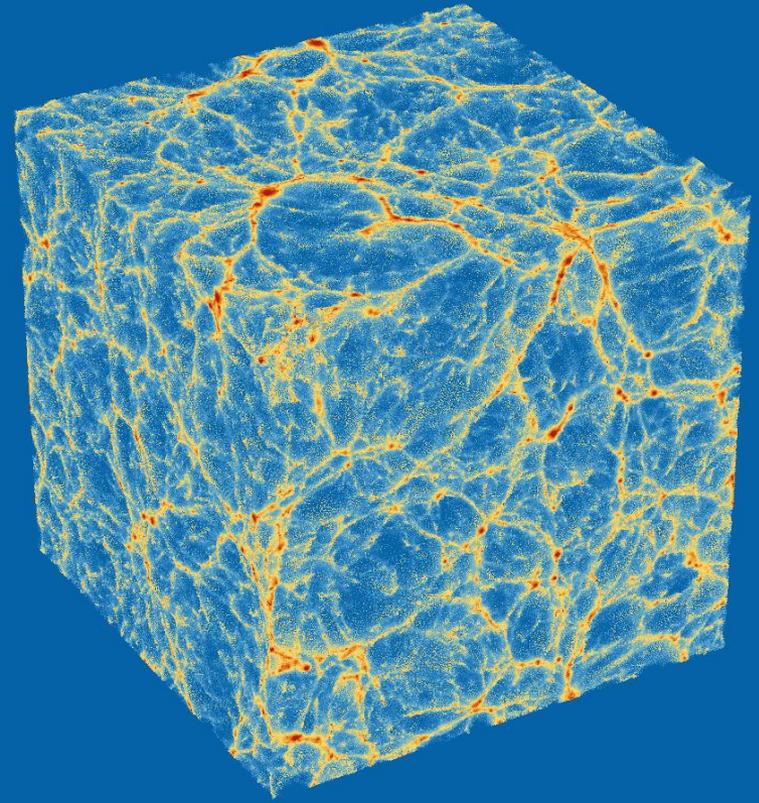
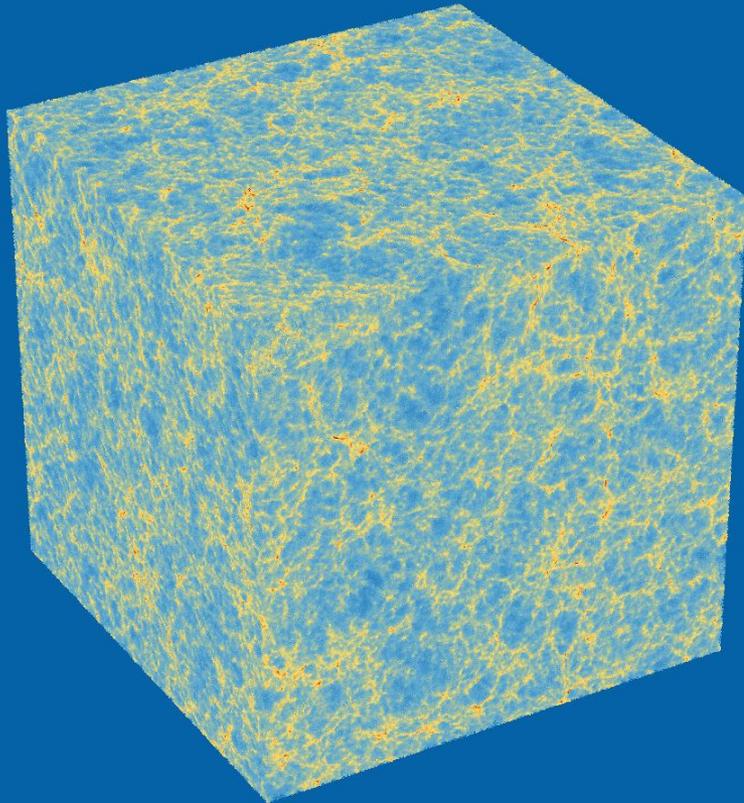


# Cosmological simulations of galaxy formation

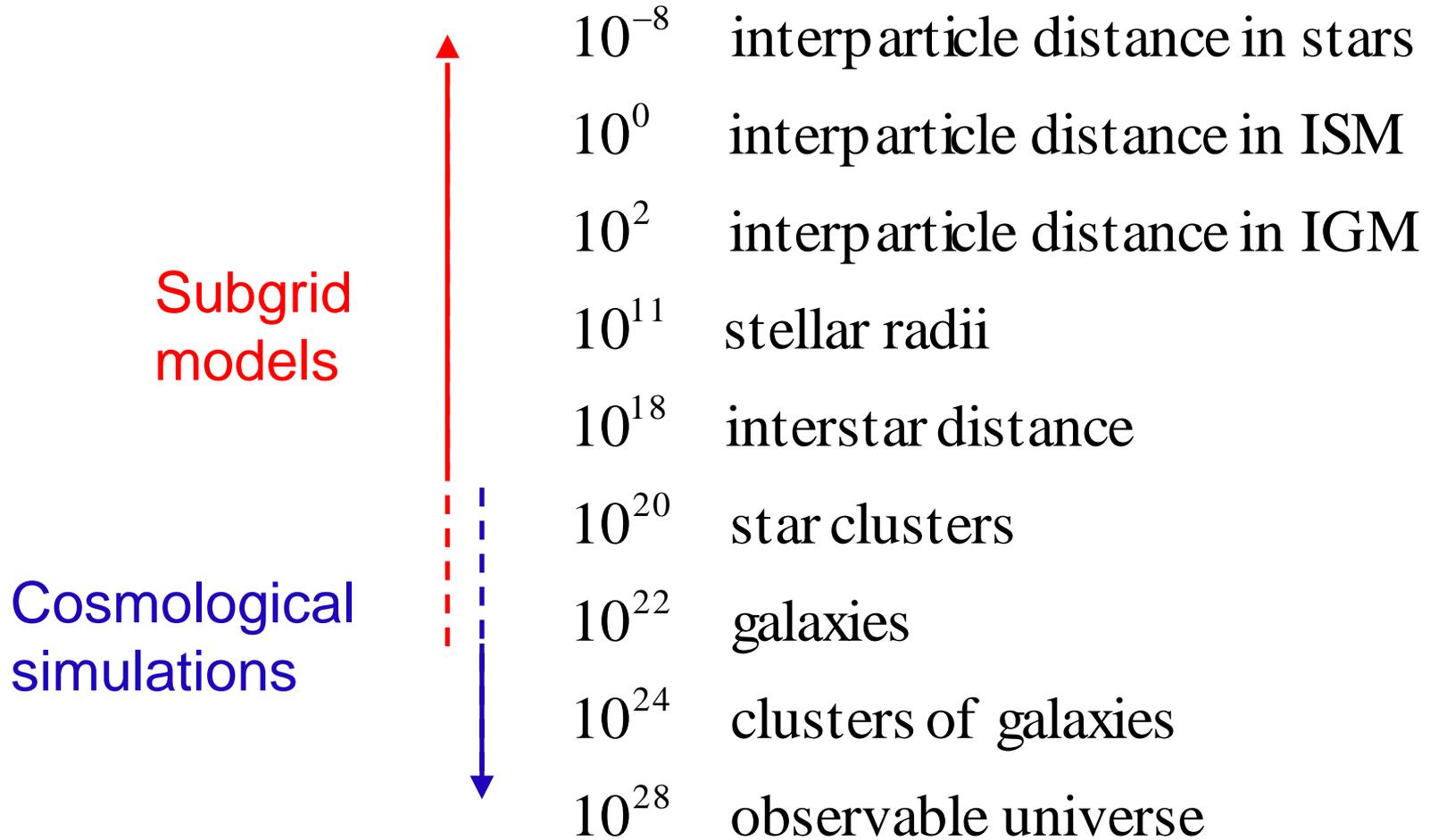
Joop Schaye (Leiden)  
(Yope Shea)



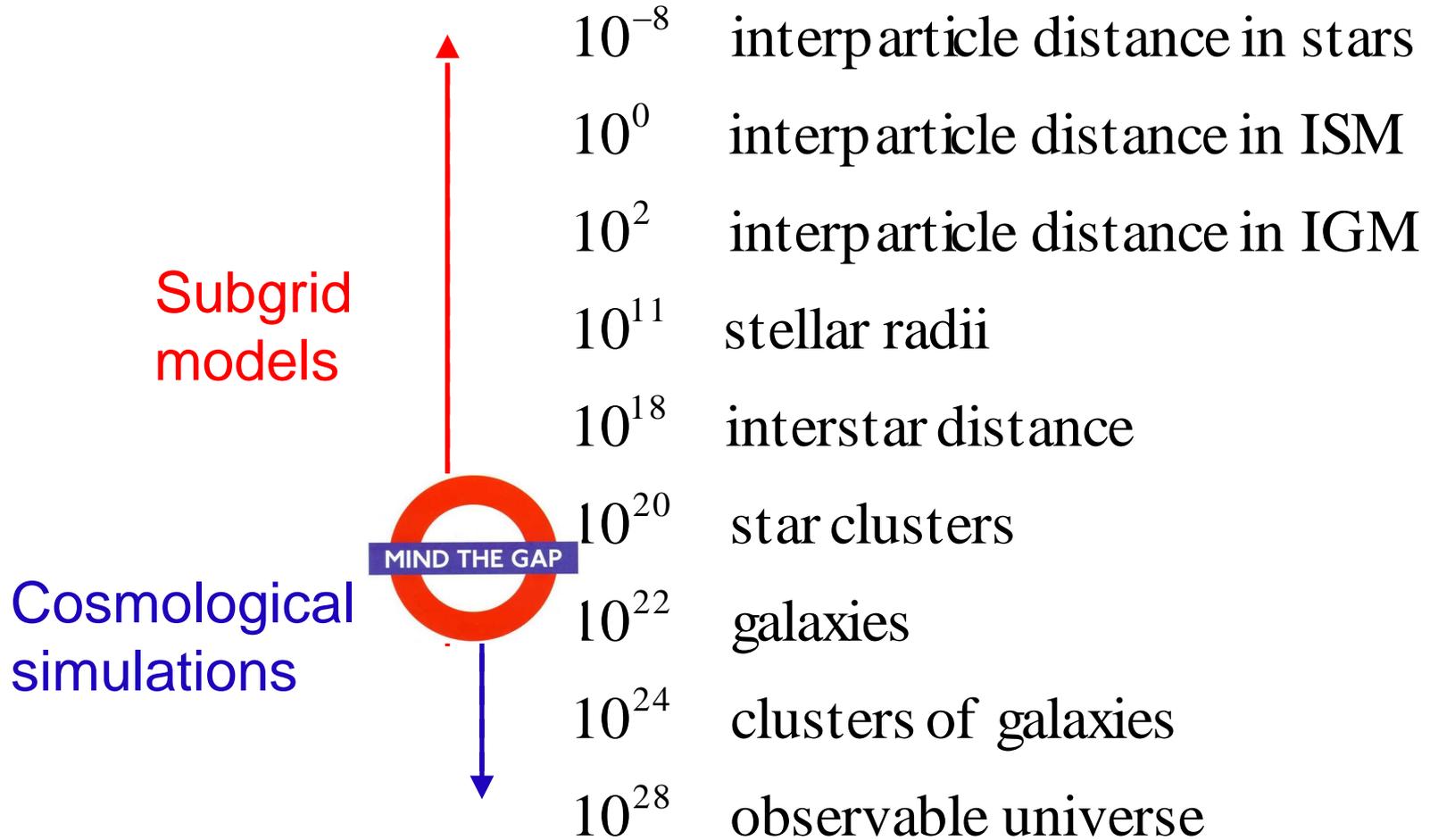
# Cosmological hydro simulations

- Evolution from  $z \gg 100$  to  $z \sim 10$  of a representative part of the universe
- Expansion solved analytically and scaled out
- Initial conditions from the CMB & LSS
- Boundary conditions: periodic
- Components: cold dark matter, gas, stars, radiation (optically thin)
- Discretization: time, mass (SPH) or length (AMR)
- Gravity and hydro solvers (and MHD, RT, ...)
- **Sub-grid modules are a crucial part of the game**

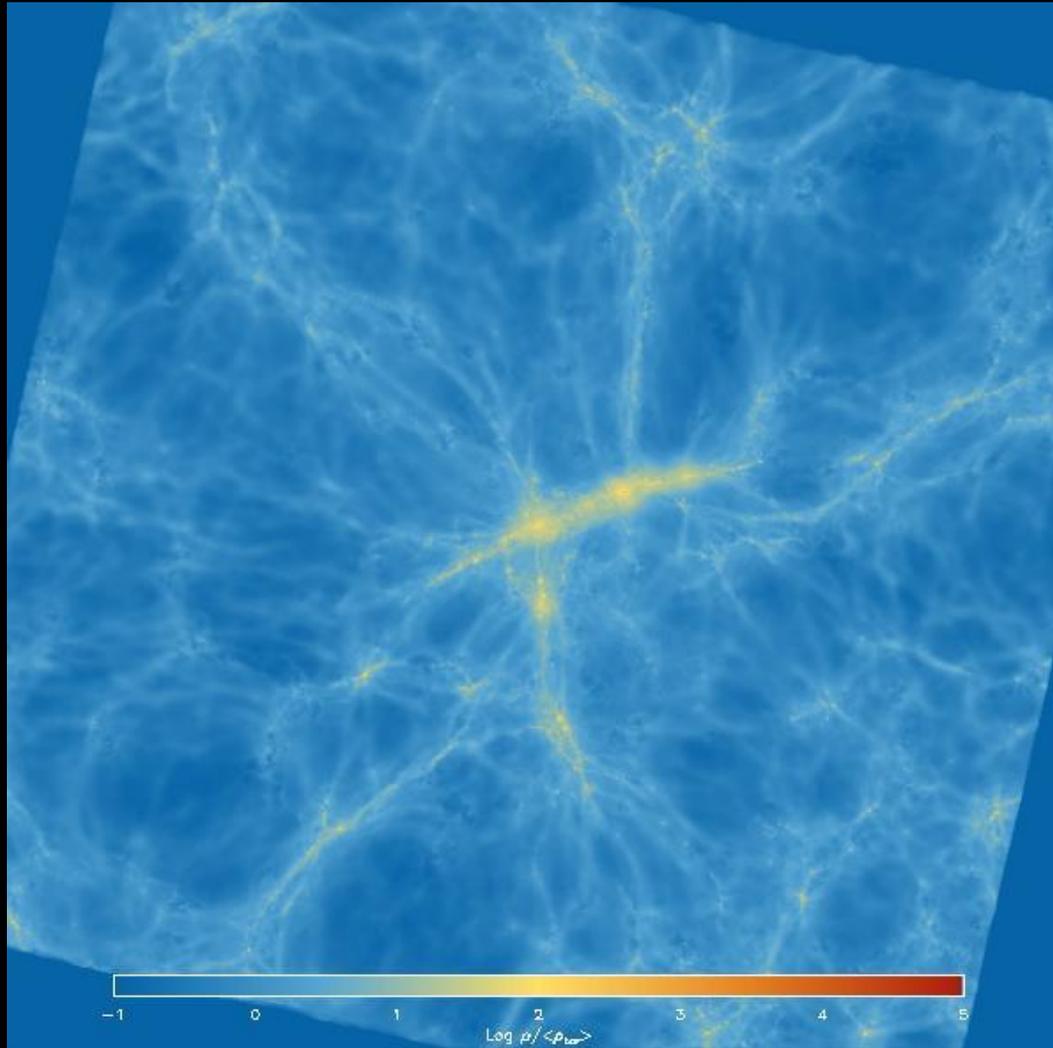
# Length Scales (cm)



# Length Scales (cm)



# Zooming into a massive galaxy at $z=2$ : Gas density



Depth: 2 Mpc/h

Log M = 12.6

Log M\* = 11.5

Simulation:

OWLS REF

L025

N512

← 25 Mpc/h →

# Where to put the gap?

Transition from warm ( $T \sim 10^4$  K) to cold, molecular ( $T \ll 10^4$  K) ISM expected at  $\Sigma_{\text{H}} \sim 10 M_{\odot} \text{pc}^{-2}$  ( $n_{\text{H}} \sim 10^{-2} - 10^{-1} \text{cm}^{-3}$  in warm phase).

- Determined by (dust) column needed to shield UV
- Associated with sharp reduction in Jeans scale  $\rightarrow$  star formation
- Threshold decreases with metallicity and increases with UV

(JS 04, Gnedin+ 09, Krumholz+ 09, 11, Gnedin & Kravtsov 11, Feldmann+ 11, Glover & Clark 12, Clark & Glover 13, ...)

Well-posed challenge:

**Resolve the Jeans scales in the warm ISM**

# Basic resolution requirements

- Convergence requires resolving the Jeans scales:

$$M_J \approx 1 \times 10^7 h^{-1} M_\odot f_g^{3/2} \left( \frac{n_H}{10^{-1} \text{ cm}^{-3}} \right)^{-1/2} \left( \frac{T}{10^4 \text{ K}} \right)^{3/2}$$

$$L_J \approx 1.5 h^{-1} \text{ kpc} f_g^{1/2} \left( \frac{n_H}{10^{-1} \text{ cm}^{-3}} \right)^{-1/2} \left( \frac{T}{10^4 \text{ K}} \right)^{1/2}$$

- Resolving the warm phase requires:
  - Particle mass  $\ll 10^7 M_\odot$
  - Spatial resolution  $\ll 1 \text{ kpc}$
- Resolving gas with  $n_H \sim 10^1 \text{ cm}^{-3}$  and  $T \sim 10^2 \text{ K}$  requires :
  - particle mass  $\ll 10^3 M_\odot$
  - spatial resolution  $\ll 10 \text{ pc}$
  - Radiative transfer
  - Complex chemistry



The cold phase is  
still too demanding for  
cosmological simulations



But we are about to cross the gap in simulations of individual galaxies!



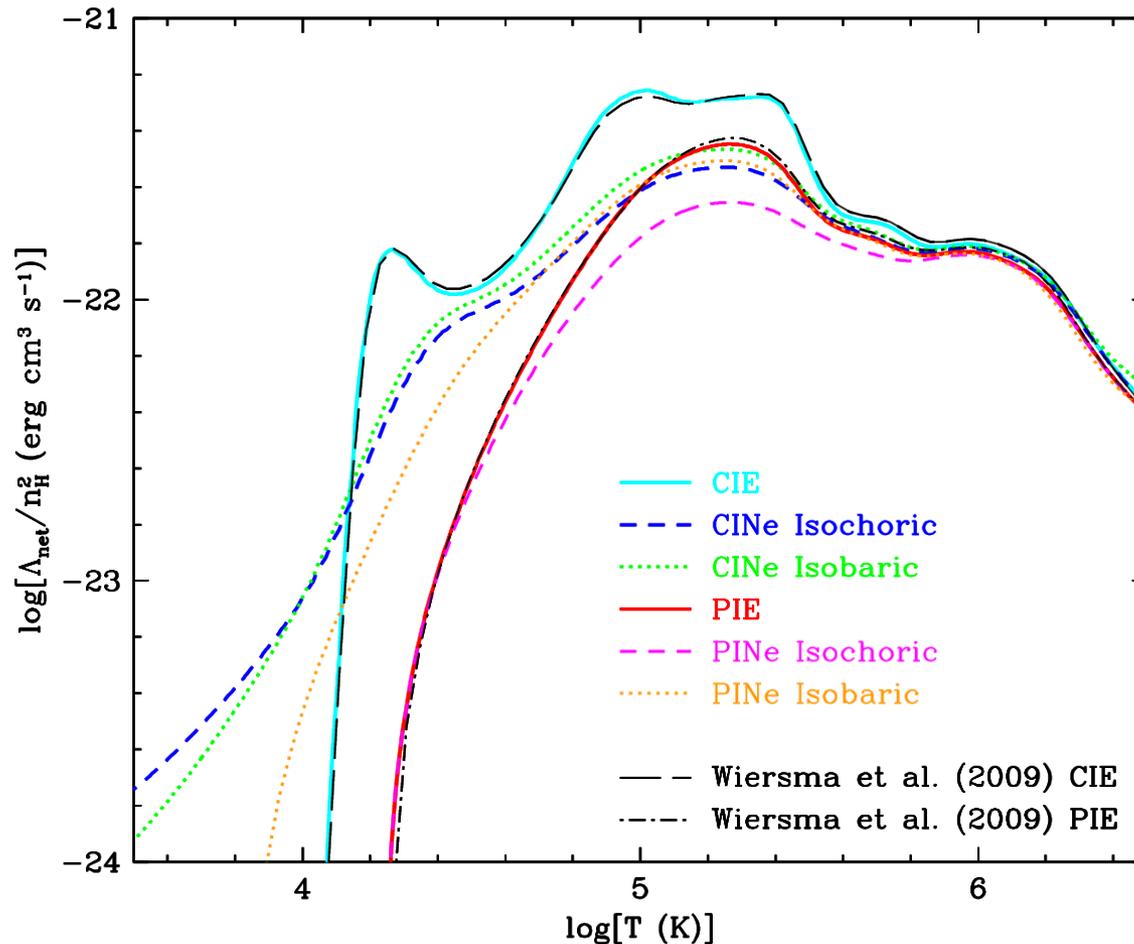
# Subgrid models for cosmological hydro simulations

- Radiative cooling/heating
- Star formation
- Chemodynamics/stellar evolution
- Galactic winds driven by feedback from SF
- Black holes and AGN feedback
- Less conventional things. E.g.:
  - Turbulence (incl. mixing)
  - Cosmic rays
  - Dust

# Radiative cooling

- Standard assumptions:
  - H & He in photo-ionisation equilibrium (optically thin, UV background only)
  - Metals in collisional ionisation equilibrium (though many studies still assume primordial abundances!)
- Recent developments (e.g. Wiersma, JS & Smith '09; Shen+ '10; Vogelsberger+ '13, Aumer+ '13):
  - Metals also in photo-ionisation equilibrium
  - Relative abundance variations
- Cutting edge/future:
  - Non-equilibrium ionization
  - Radiative transfer
  - Local radiation sources
  - Molecules
  - Dust

# Cooling: effect of non-equil. and photo-ionisation

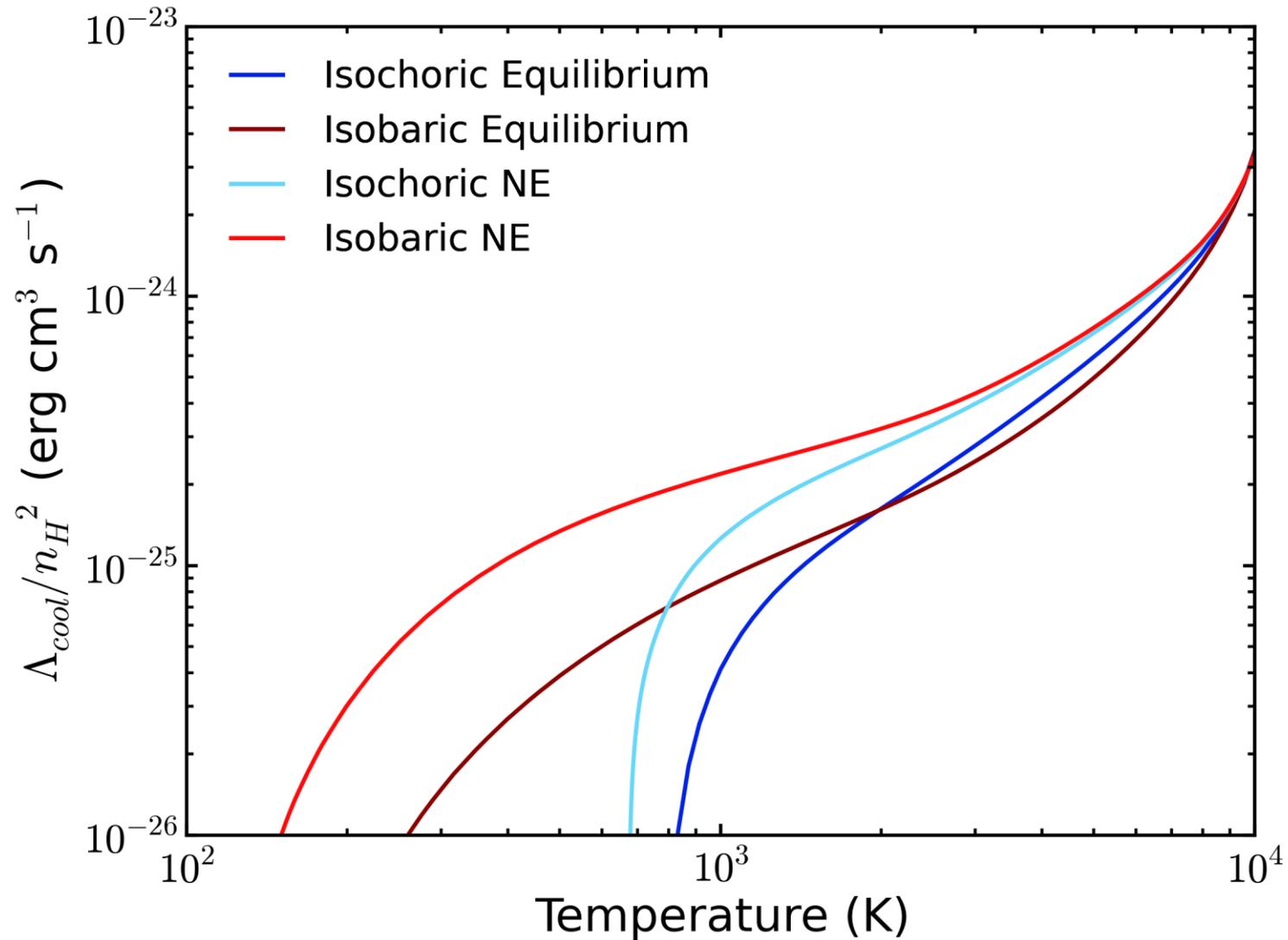


$n_{\text{H}}=10^{-4} \text{ cm}^{-3}$ ,  $z=1$ ,  $Z=Z_{\odot}$

Oppenheimer & JS (2013a)



# Cooling: effect of non-equil. at $T < 10^4$ K

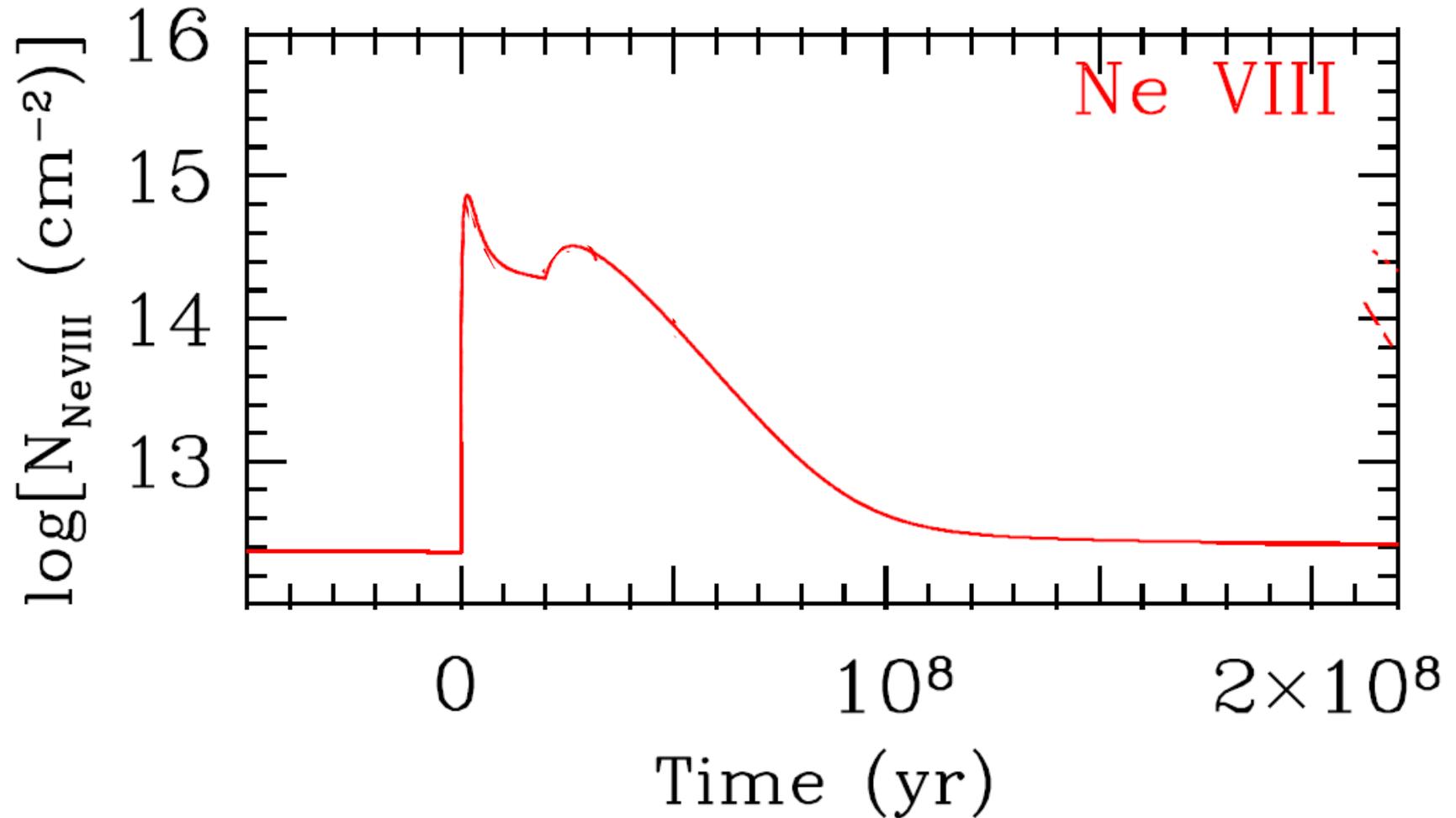


$n_H = 1 \text{ cm}^{-3}$ ,  $Z = Z_\odot$

Richings, JS & Oppenheimer (in prep)



# AGN proximity zone fossils



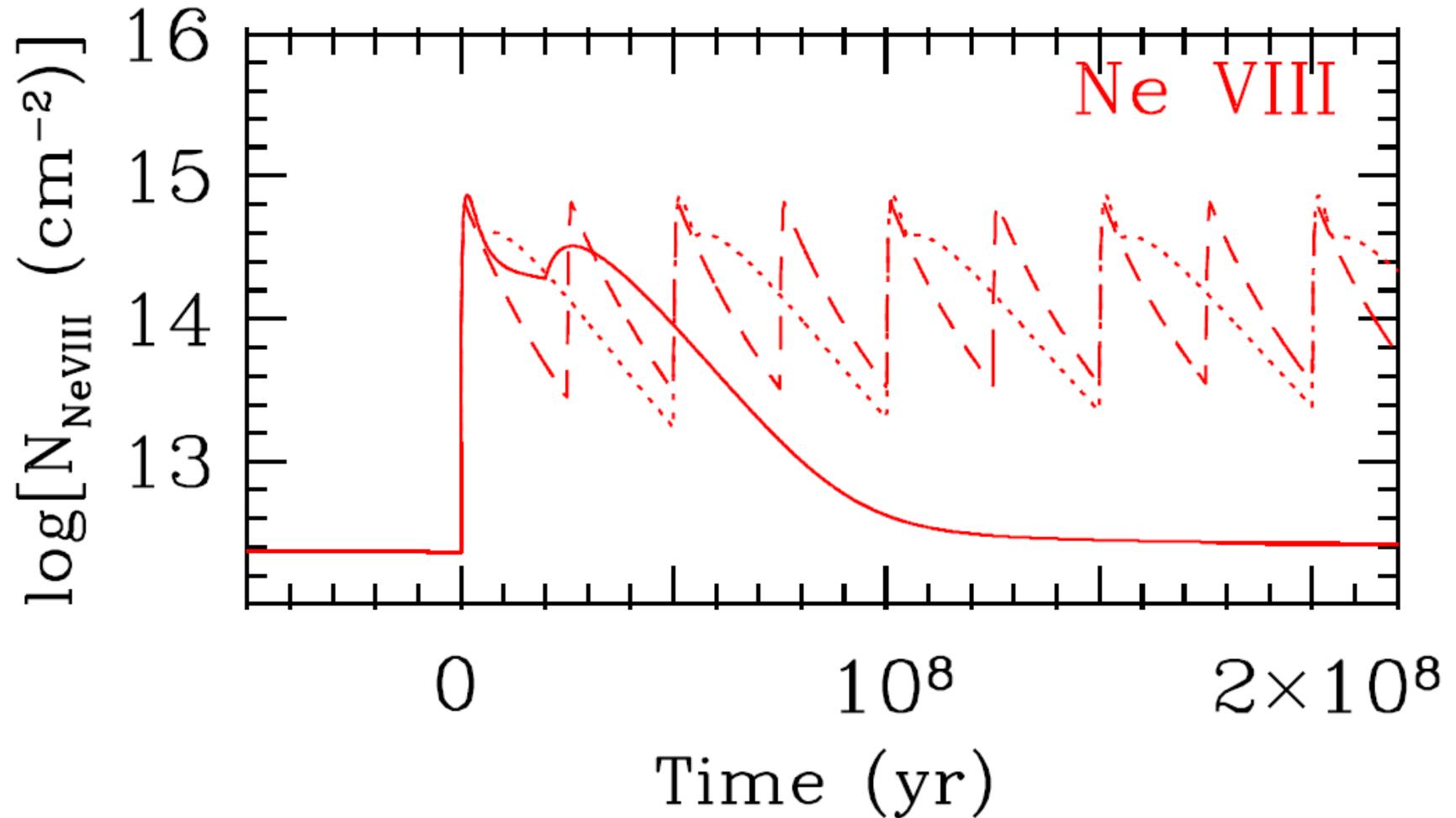
—  $\tau_{\text{AGN}} = 20$  Myr once

$n_{\text{H}} = 10^{-4} \text{ cm}^{-3}$ ,  $T = 10^4 \text{ K}$ ,  $z=1$ ,  $Z=Z_{\odot}$

Oppenheimer & JS (2013b)



# AGN proximity zone fossils



—  $\tau_{\text{AGN}} = 20$  Myr once  
.....  $\tau_{\text{AGN}} = 5$  Myr every 50 Myr  
- - -  $\tau_{\text{AGN}} = 1$  Myr every 25 Myr

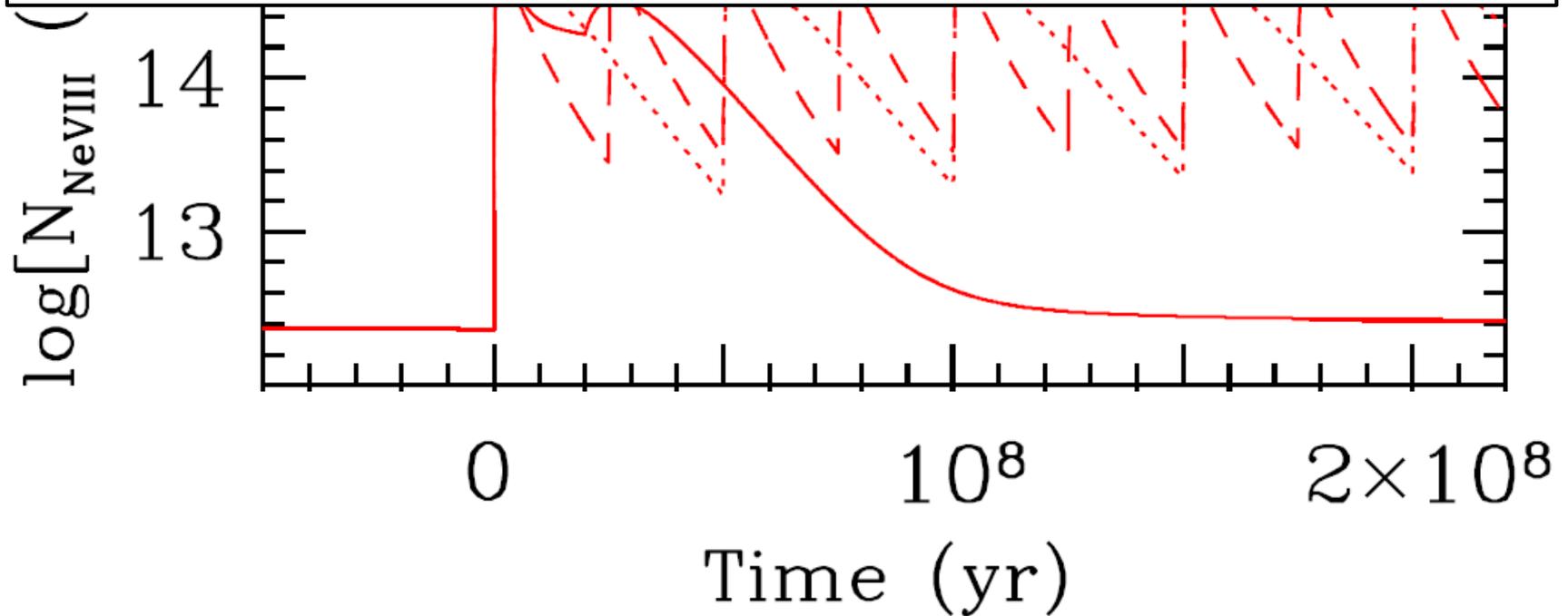
$n_{\text{H}} = 10^{-4} \text{ cm}^{-3}$ ,  $T = 10^4 \text{ K}$ ,  $z=1$ ,  $Z=Z_{\odot}$

Oppenheimer & JS (2013b)



# AGN proximity zone fossils

Most intergalactic metals may reside in out-of-equilibrium AGN fossil zones!



—  $\tau_{\text{AGN}} = 20$  Myr once  
.....  $\tau_{\text{AGN}} = 5$  Myr every 50 Myr  
- - -  $\tau_{\text{AGN}} = 1$  Myr every 25 Myr

$n_{\text{H}} = 10^{-4} \text{ cm}^{-3}$ ,  $T = 10^4 \text{ K}$ ,  $z=1$ ,  $Z=Z_{\odot}$

Oppenheimer & JS (2013b)



# Galactic winds driven by SF

- Winds may be:
  - Energy-driven
  - Momentum-driven
  - Both
- Sources of energy/momentum:
  - Supernovae
  - Radiation pressure:
    - On dust
    - From photo-ionisation
    - From trapping of Ly $\alpha$
  - Stellar winds
  - Cosmic rays
  - Combination of the above

# Galactic winds driven by SF: WARNINGS

- Efficient feedback is required to match observations
- Feedback is often inefficient due to the numerical implementation...
- ... but inefficient feedback is sometimes interpreted as a need for different physical processes
- Nearly all implementations are extremely crude (e.g. radiation pressure w/o radiative transfer)
- At the current resolution, the different feedback processes are hardly distinguishable
- Many hydro simulations use tricks that make them more like SAMs than you may think. E.g.:
  - Wind velocity depends on halo mass or dark matter velocity dispersion (e.g. Okamoto+, Davé/Oppenheimer+, Viel+, Vogelsberger+)
  - Temporarily turn off hydro for winds (e.g. Springel/Hernquist, Davé/Oppenheimer, Viel, Vogelsberger)
  - Temporarily turn off radiative cooling (nearly everyone else)

# Implementing FB: recognized problems

Simplest recipe: star particles inject thermal energy into surroundings (e.g. Katz+ '96)

Recognized problems:

- Much of the mass in the ISM is in the cold phase ( $T \ll 10^4 \text{ K}$ )
- Simulations do not model cold phase
  - intercloud density too high
    - cooling rate too high
    - feedback too inefficient
  - SF insufficiently clustered
    - feedback too inefficient

# Driving winds: brute force solution

- Allow for a cold phase
- Increase SF threshold (only sensible for cold phase)
- Still require subgrid recipe, but on smaller scale

Problem: need very high resolution

- Can only model a small number of galaxies (zoomed simulations)
- Need to pick initial conditions (e.g. merger history)

(e.g. Ceverino & Klypin '07, Hopkins+ '12, Ceverino+ '13)

# Driving winds: subgrid recipes

- **Multiphase particles** (e.g. Marri & White '93, Scannapieco, Murante, Aumer/White)
- **Suppress cooling by hand** (e.g. Gerritsen '97, Thacker, Stinson/Brook/Gibson/Governato/Maccio/Mayer/Wadsley/...)
- **Inject momentum (i.e. kinetic feedback)** (e.g. Navarro & White '93, Springel/Hernquist, Davé/Oppenheimer, Teyssier, OWLS/GIMIC, Vogelsberger, ...)
  - Most relevant advantage: can decrease initial mass loading
- **Temporarily decouple winds from the hydrodynamics** (e.g. Springel/Hernquist '03, Davé/Oppenheimer, Viel, Vogelsberger, ...)
- **Multiple feedback processes** (e.g. Hopkins+, Stinson+ '13, ...)

# Implementing FB: less recognized problems

- Reality: SNe (or BHs) inject lots of energy in very little mass
  - High temperatures
  - Long cooling times
  - Efficient feedback
- Simulations: inject energy in large gas mass
  - Low heating temperatures
  - Short cooling times
  - Inefficient feedback

e.g. Kay+ '03; Booth & JS '09; Creasey+ '11; Dalla Vecchia & JS '12

# Implementing FB: less recognized problems

- The SNII of an SSP of mass  $m_*$  can heat a mass  $m_{g,heat}$  by  
$$\Delta T = 4 \times 10^7 \text{ K} (m_*/m_{g,heat})$$
- Reality:  $m_* \gg m_{g,heat}$  initially
  - $\Delta T \gg 10^8 \text{ K}$
  - $t_c \gg 10^8 \text{ yr} (n_H/1 \text{ cm}^{-3})^{-1}$
- In simulations:  $m_* \sim 0.01 - 0.1 m_{g,heat}$ 
  - $\Delta T \sim 10^6 \text{ K}$
  - $t_c \sim 10^5 \text{ yr} (n_H/1 \text{ cm}^{-3})^{-1}$
  - overcooling
- Note that in SPH simulations  $(m_*/m_{g,heat})$  is independent of resolution!



# Implementing thermal FB: requirements

- FB only efficient if heated resolution elements expand faster than they cool radiatively:

$$t_c \gg t_s = h/c_s$$

where  $h$  is the spatial resolution

- Adiabatic expansion does not change  $t_c / t_s$  (assuming Bremsstrahlung)
- Required  $T$  depends on density and resolution

$$\frac{t_c}{t_s} = 2.8 \times 10^2 \left( \frac{n_H}{1 \text{ cm}^{-3}} \right)^{-1} \left( \frac{T}{10^{7.5} \text{ K}} \right) \left( \frac{h}{100 \text{ pc}} \right)^{-1}$$

$$\frac{t_c}{t_s} \simeq 98 \left( \frac{n_H}{1 \text{ cm}^{-3}} \right)^{-2/3} \left( \frac{T}{10^{7.5} \text{ K}} \right) \left( \frac{\langle m \rangle}{7 \times 10^4 M_\odot} \right)^{-1/3}$$

Dalla Vecchia & JS (2012)



# Implementing efficient thermal FB:

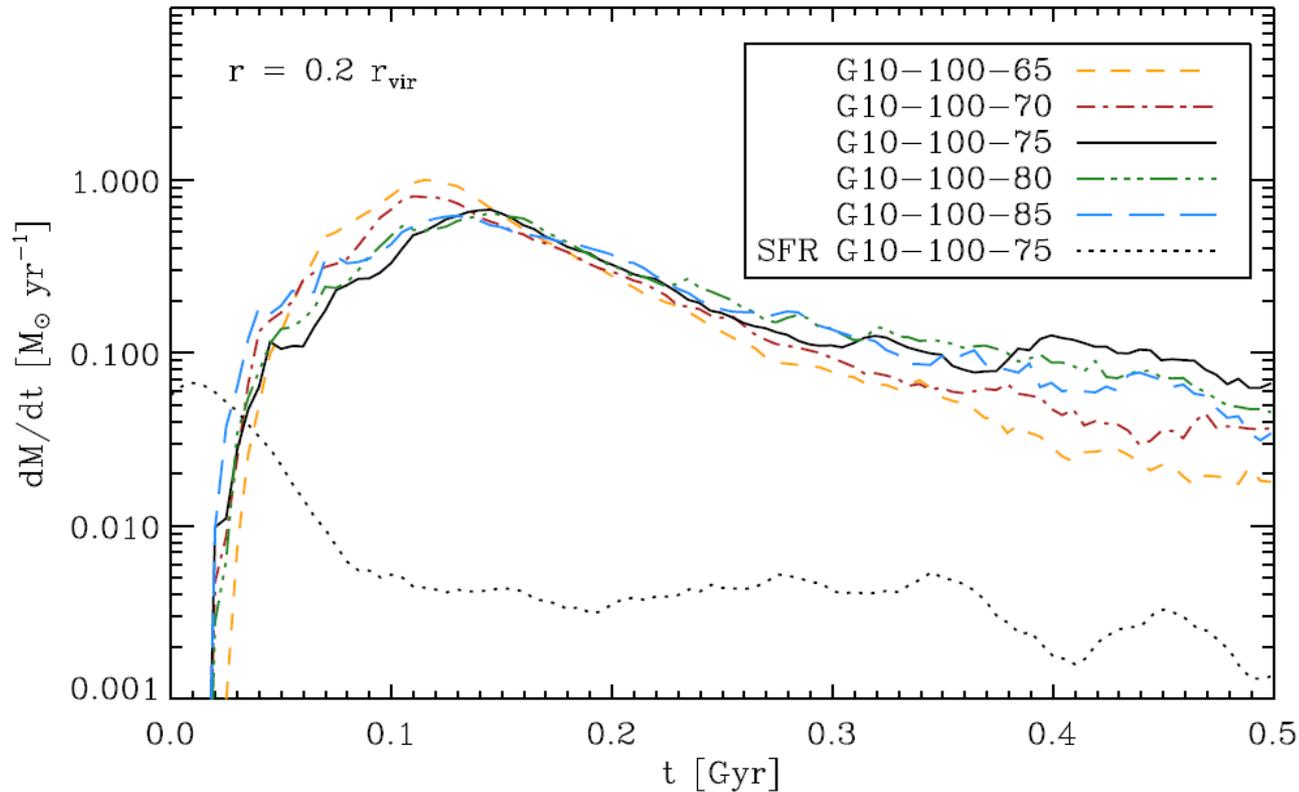
- $\Delta T$  determined by resolution
- Stochastic FB
  - given  $\Delta T$ , fraction of available energy that is injected,  $f_{th}$ , determines heating probability
- $f_{th}$  not predicted, unresolved thermal losses need to be calibrated

Dalla Vecchia & JS (2012)



# Mass outflow rate: $10^{10} M_{\odot}$ halo

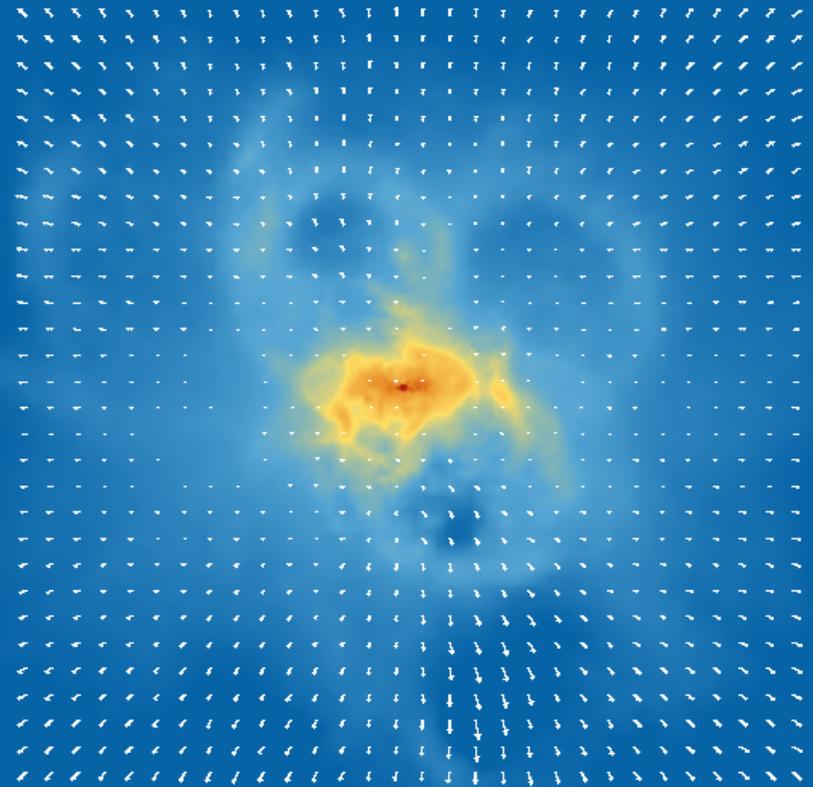
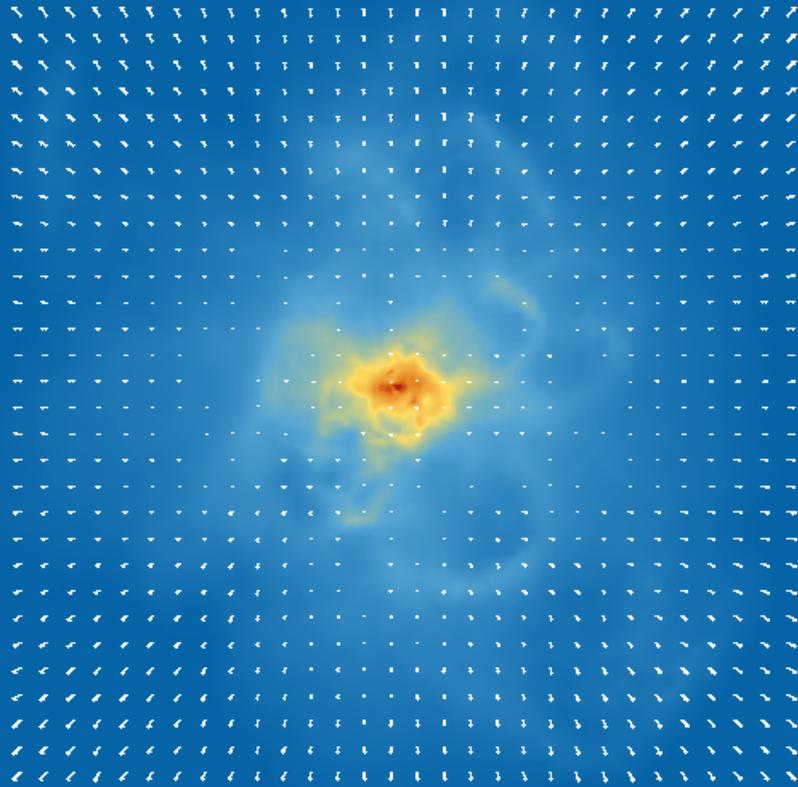
- Particle mass  $7 \times 10^2 M_{\odot} \rightarrow \frac{t_c}{t_s} \cong 5 \times 10^2 \left( \frac{n_H}{1 \text{ cm}^{-3}} \right)^{-2/3} \left( \frac{\Delta T}{10^{7.5} \text{ K}} \right)$
- Max  $n_H \sim 10^2 \text{ cm}^{-3} \rightarrow$  insensitive to  $\Delta T$  for  $\Delta T \geq 10^{6.5} \text{ K}$



# $10^{10} M_{\odot}$ halo, edge-on, gas density

$\Delta T = 10^{6.5} \text{ K}$

$\Delta T = 10^{7.5} \text{ K}$

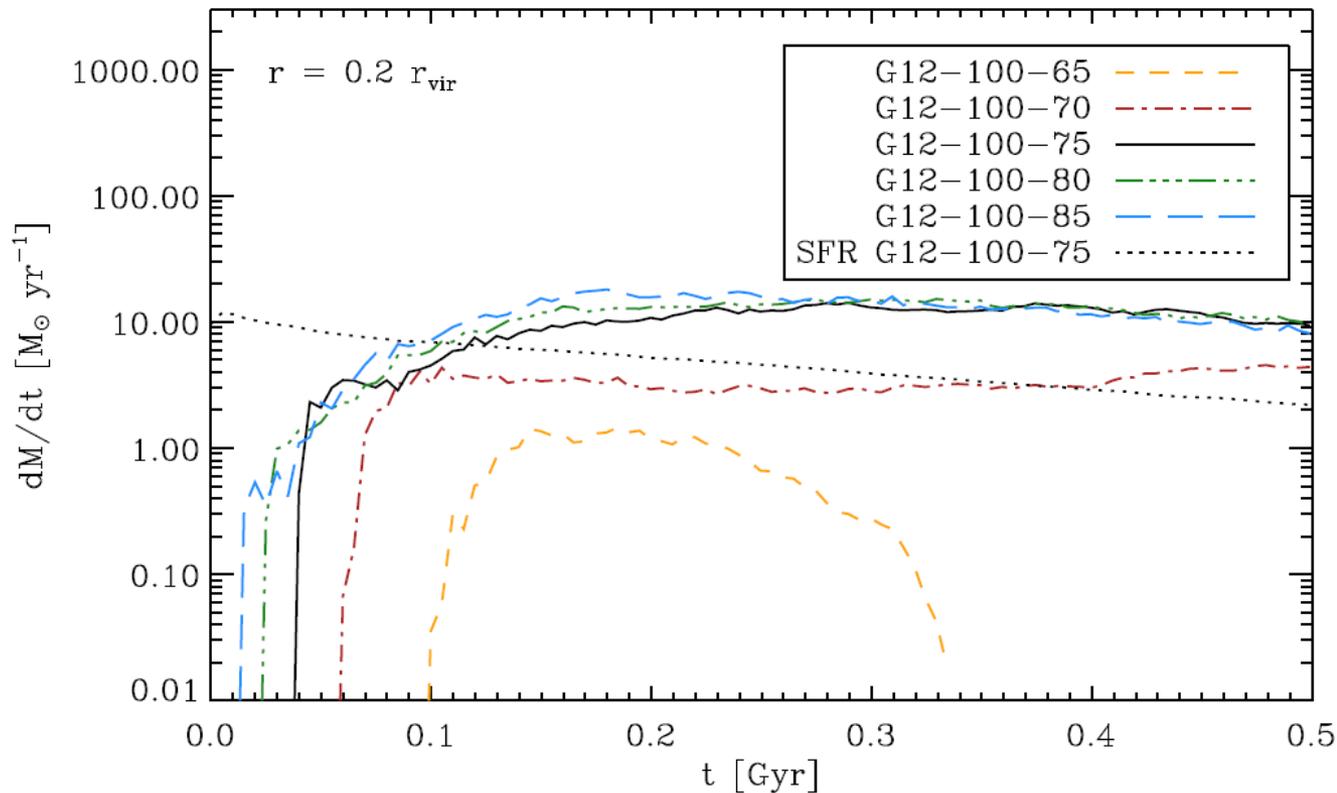


← 17.5 kpc/h →

Dalla Vecchia & JS (2012)

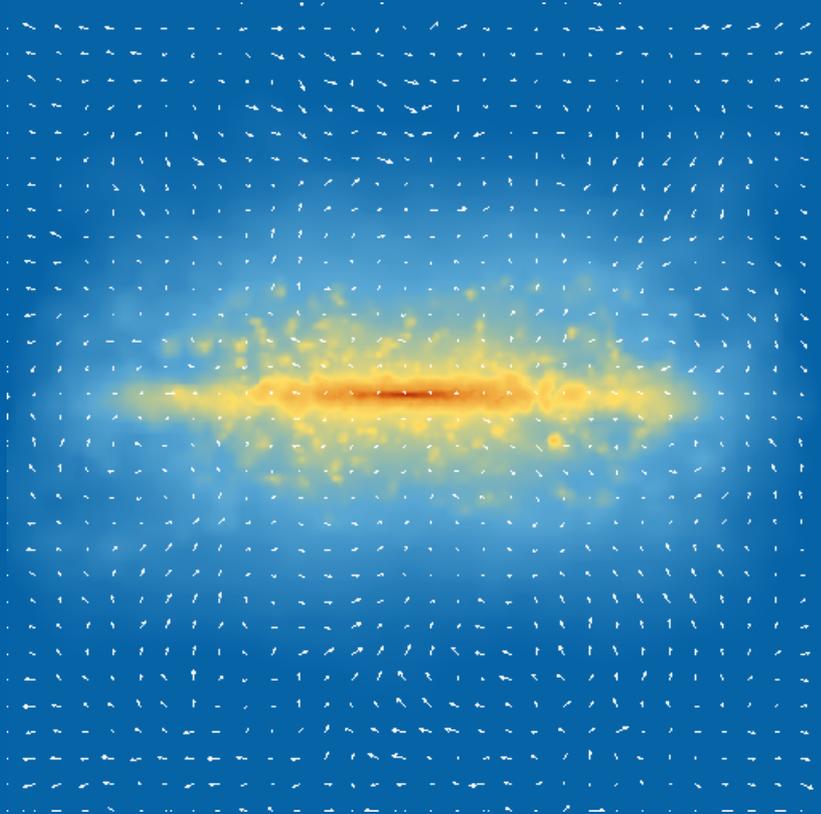
# Mass outflow rate: $10^{12} M_{\odot}$ halo

- Particle mass  $7 \times 10^4 M_{\odot} \rightarrow \frac{t_c}{t_s} \cong 1 \times 10^2 \left( \frac{n_H}{1 \text{ cm}^{-3}} \right)^{-2/3} \left( \frac{\Delta T}{10^{7.5} \text{ K}} \right)$
- Max  $n_H \sim 10^3 \text{ cm}^{-3} \rightarrow$  insensitive to  $\Delta T$  for  $\Delta T \geq 10^{7.5} \text{ K}$

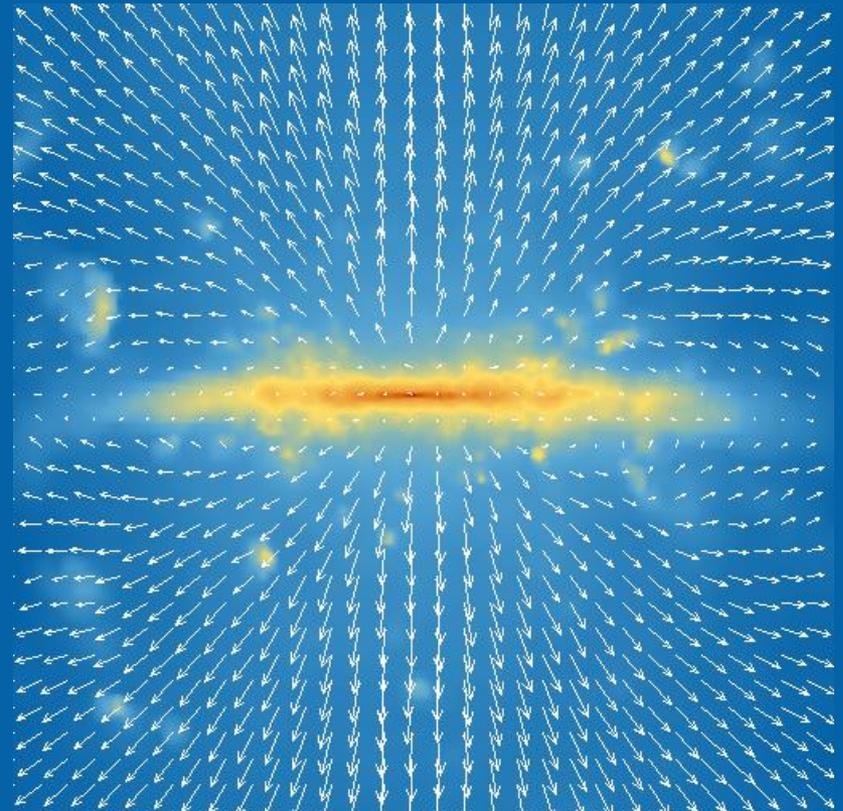


# $10^{12} M_{\odot}$ halo, edge-on, gas density

$\Delta T = 10^{6.5} \text{ K}$



$\Delta T = 10^{7.5} \text{ K}$



← 45 kpc/h →

Dalla Vecchia & JS (2012)

# Self-regulated galaxy formation

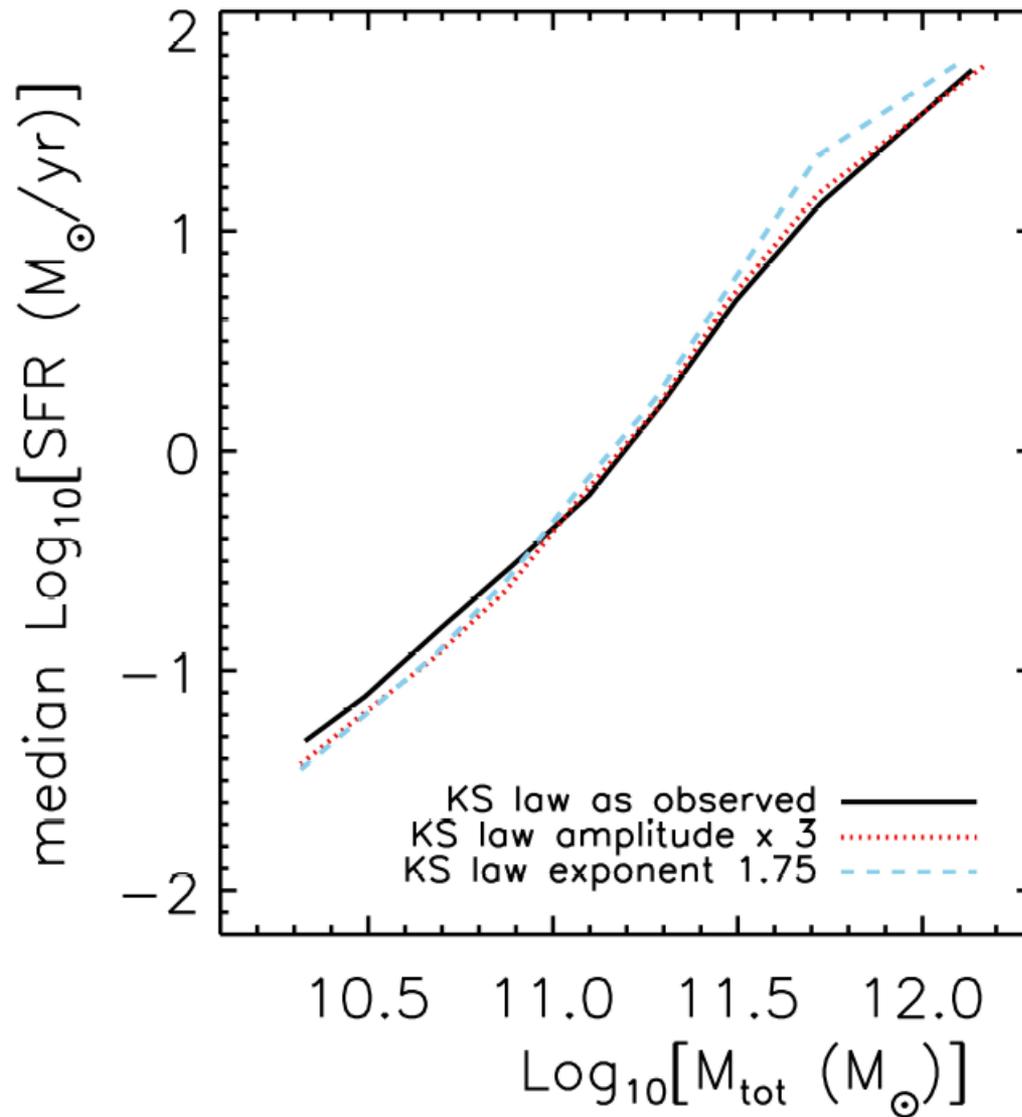
- Feedback too weak compared to accretion
  - Gas density increases
  - Star formation /BH growth rate increases
  - Feedback increases
- Feedback too strong compared to accretion
  - Gas density decreases
  - Star formation/BH growth rate decreases
  - Feedback decreases

# Consequences of self-regulated GF

Outflow rate is determined by inflow rate.  
Hence, it is independent of:

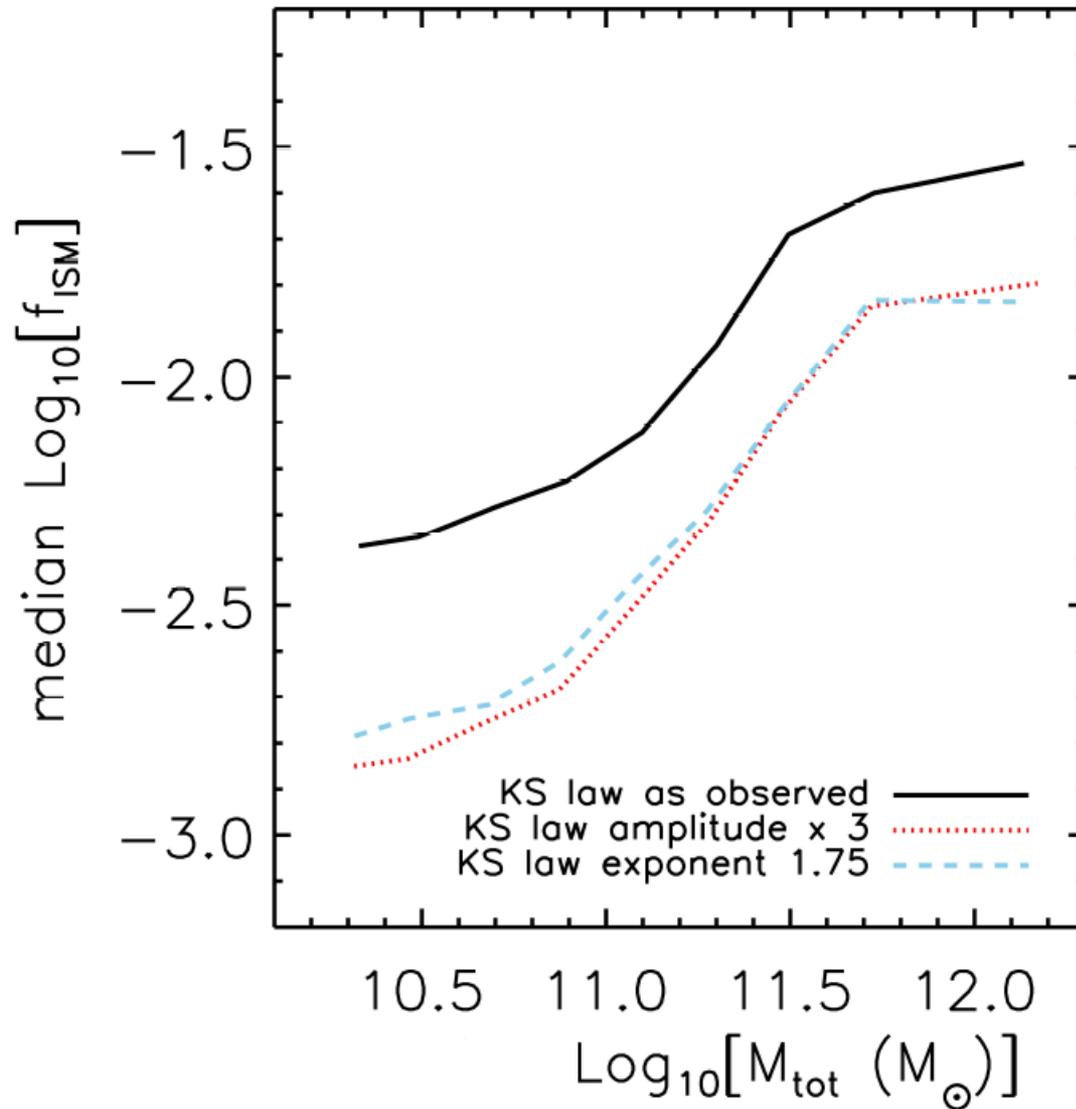
- SF law
  - SFR independent of SF law
  - SF law determines the amount of gas that is involved in SF ("gas fraction")

# Varying the SF law: SFR(M)



Haas, JS, et al. (2013a)

# Varying the SF law: Gas fraction



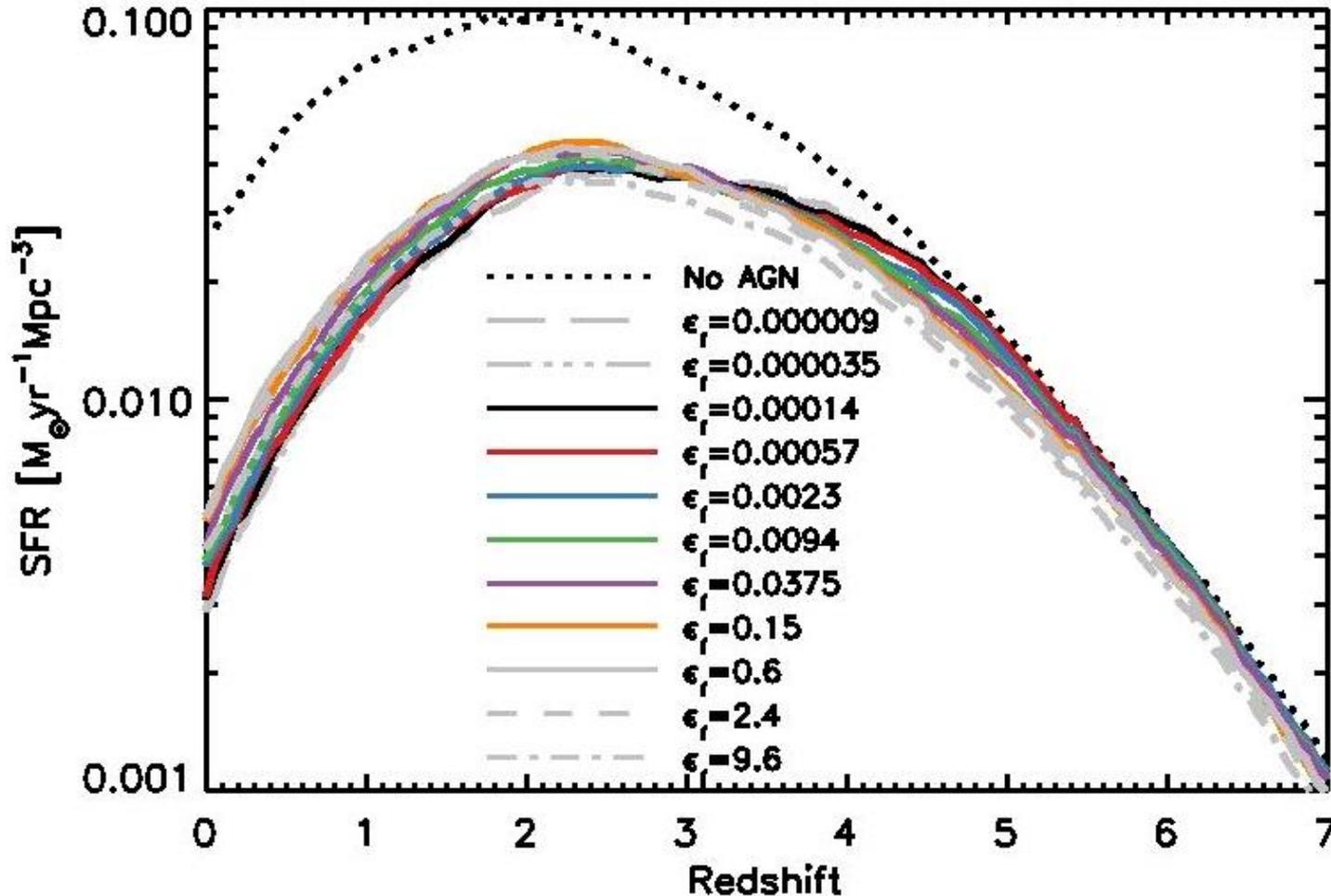
Haas, JS, et al. (2013a)

# Consequences of self-regulated GF

Outflow rate is determined by inflow rate.  
Hence, it is independent of:

- SF law
  - SFR independent of SF law
    - SF law determines the amount of gas that is involved in SF ("gas fraction")
- SF feedback efficiency
  - SFR, and hence  $M_*$ , inversely proportional to efficiency of SF feedback
- AGN feedback efficiency
  - BH accretion rate, and hence  $M_{\text{BH}}$ , inversely proportional to efficiency of AGN feedback
  - SFR independent of AGN feedback efficiency

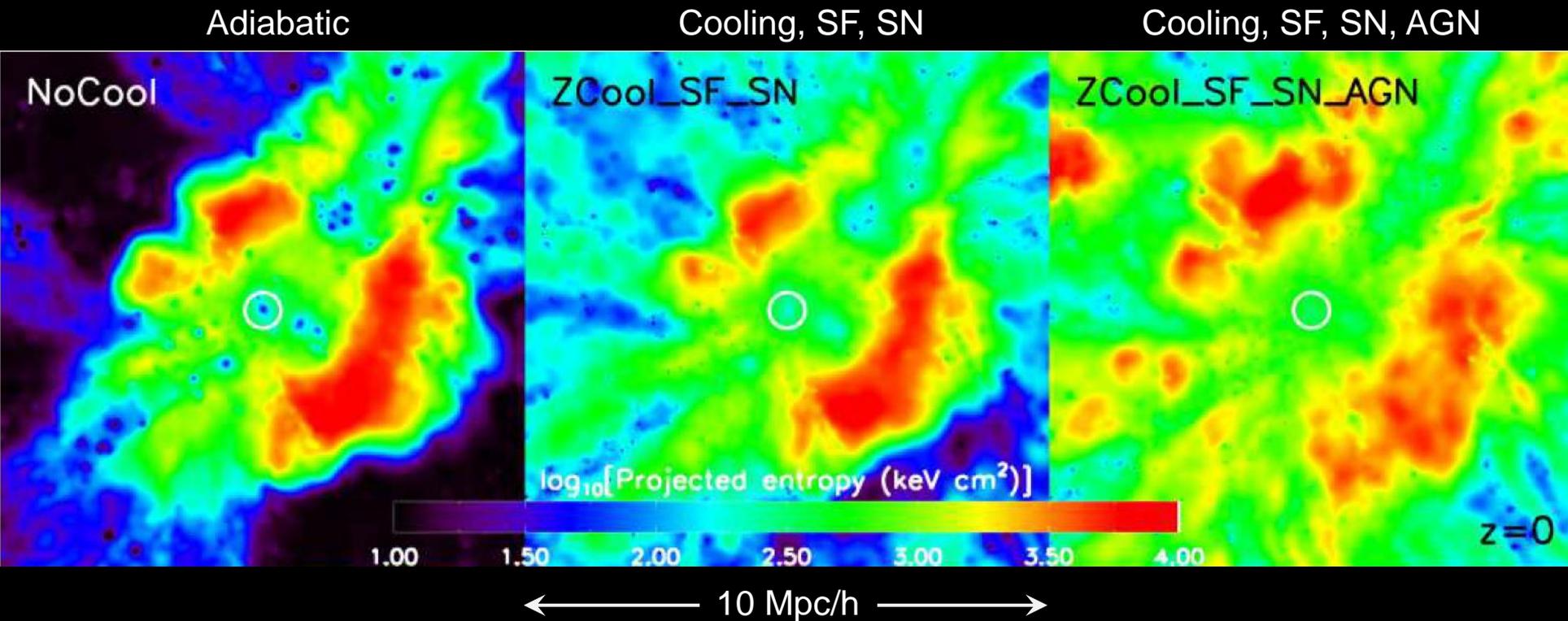
# Varying the efficiency of AGN feedback



Self-regulation on scale  
of DM haloes! (Booth &  
JS 2010, 2011)

Booth & JS (2009)

# The effect of baryons on the distribution of matter



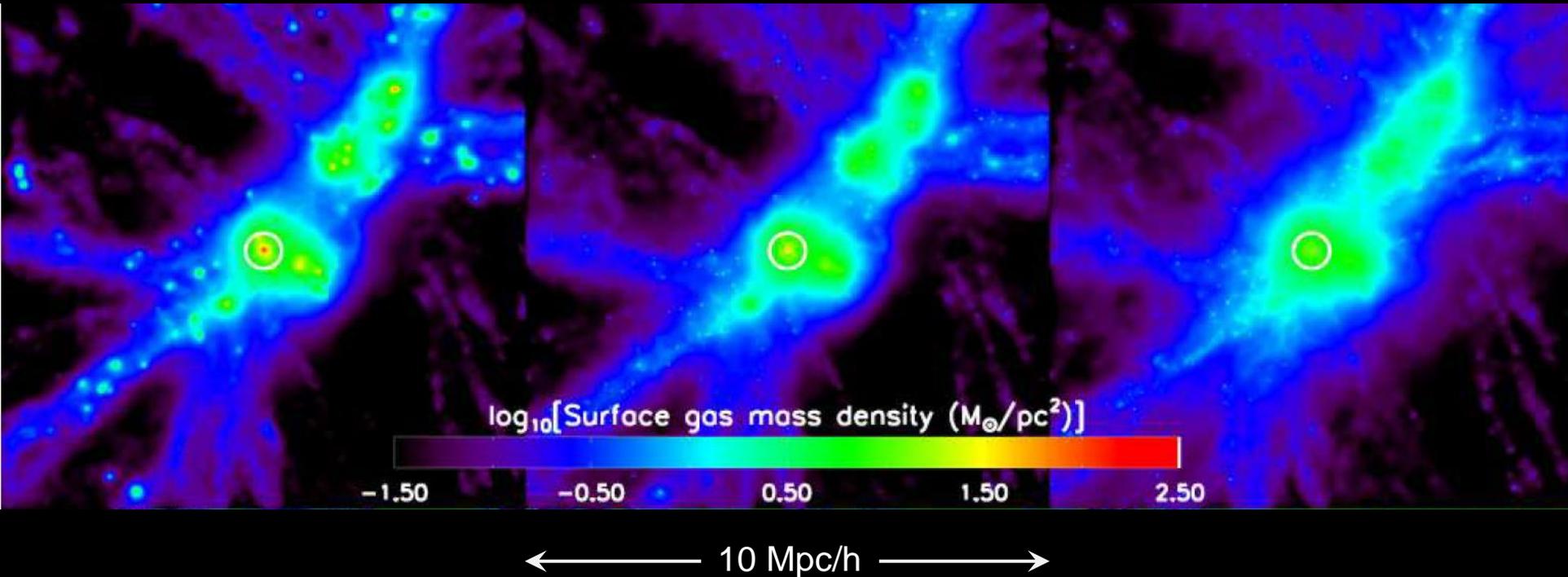
McCarthy, JS+ (2011)

# The effect of baryons on the distribution of matter

Adiabatic

Cooling, SF, SN

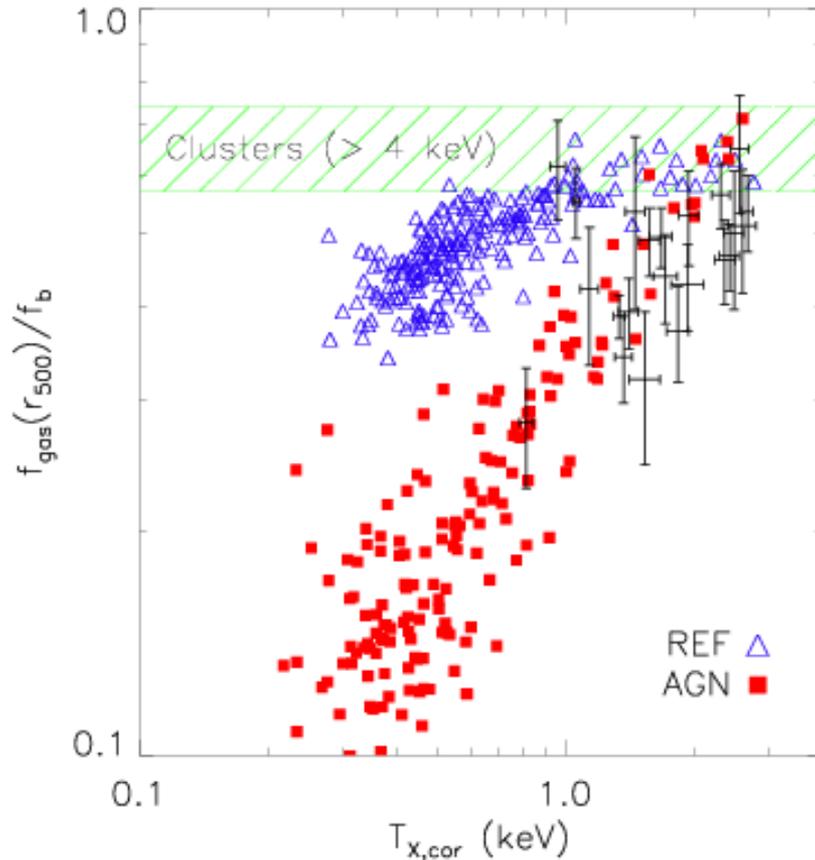
Cooling, SF, SN, AGN



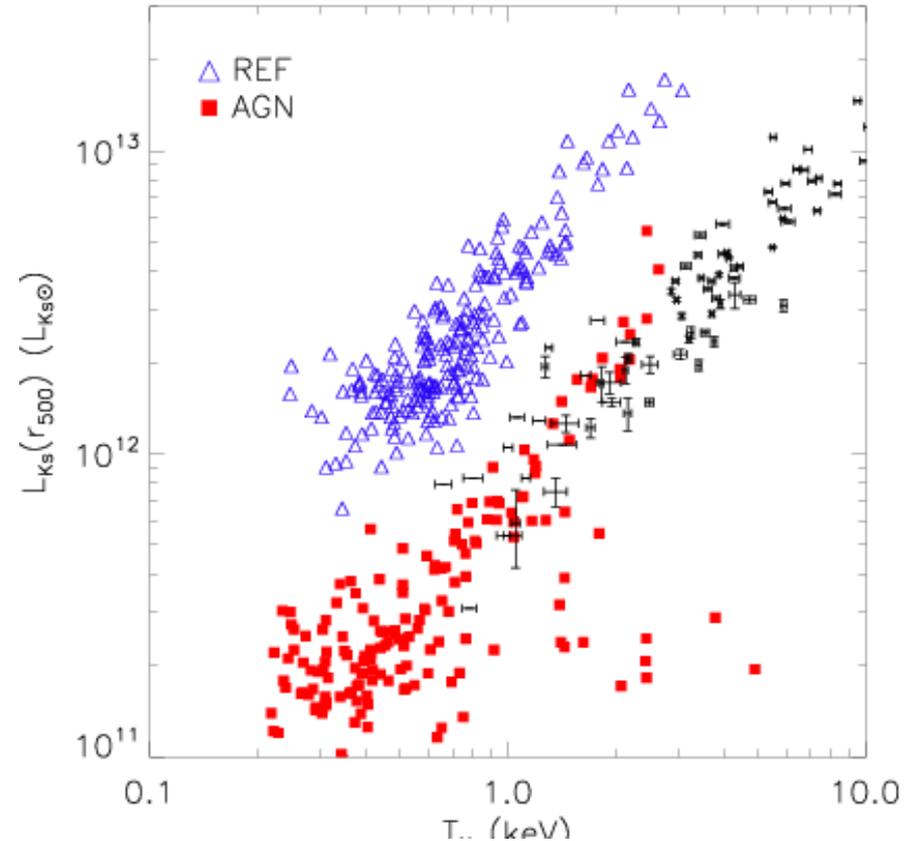
McCarthy, JS+ (2011)

# Group gas and stellar contents

Gas fraction vs T



K-band luminosity vs T

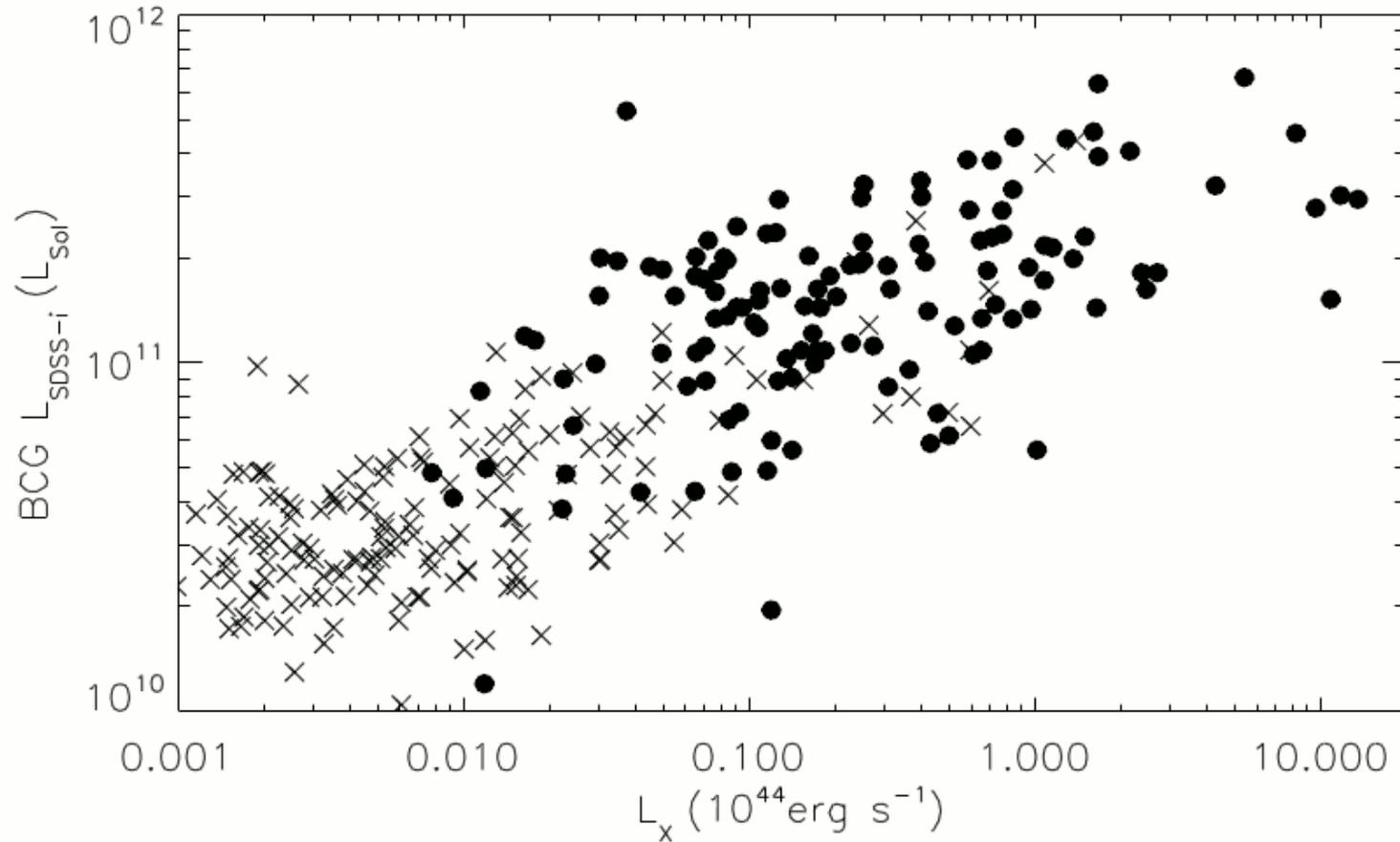


McCarthy, JS+ (2010)

Observations: Lin & Mohr 2004, Horner 2001,  
Rasmussen & Ponman (2009)



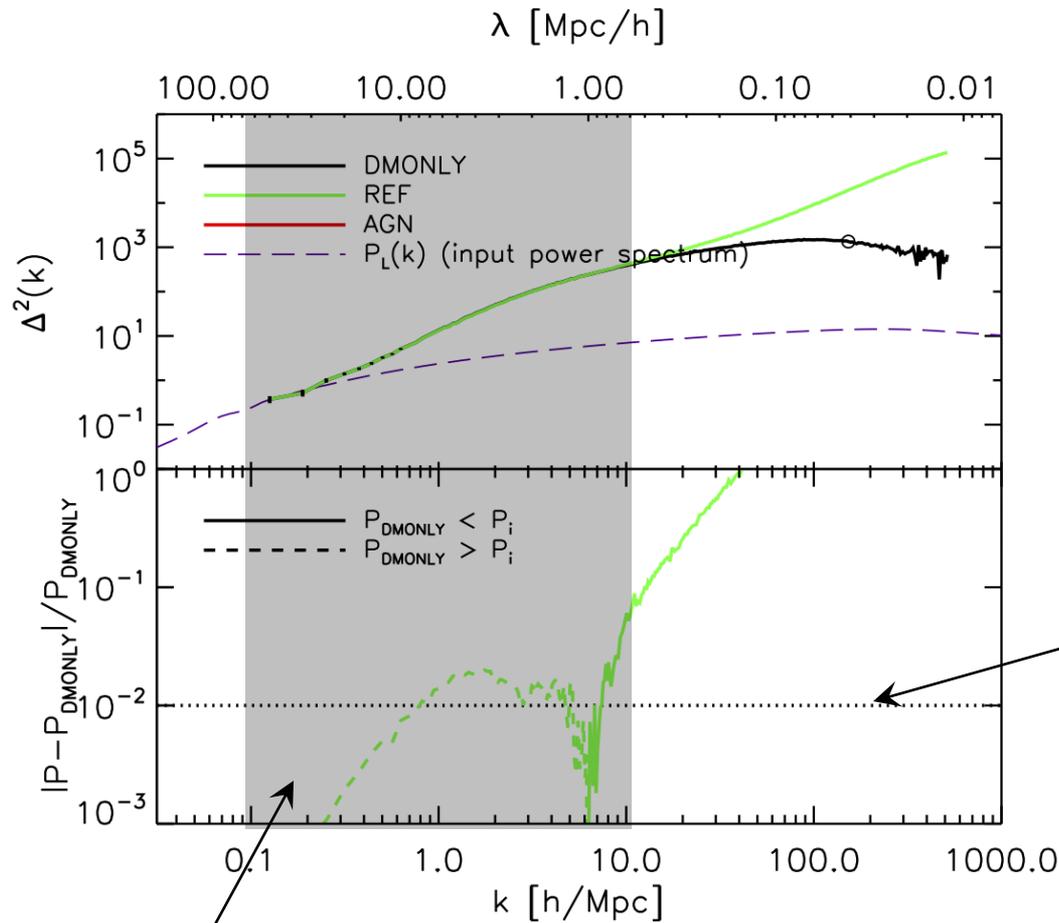
# Optical vs X-ray luminosity



Stott+ (2012)



# Baryons and the matter power spectrum



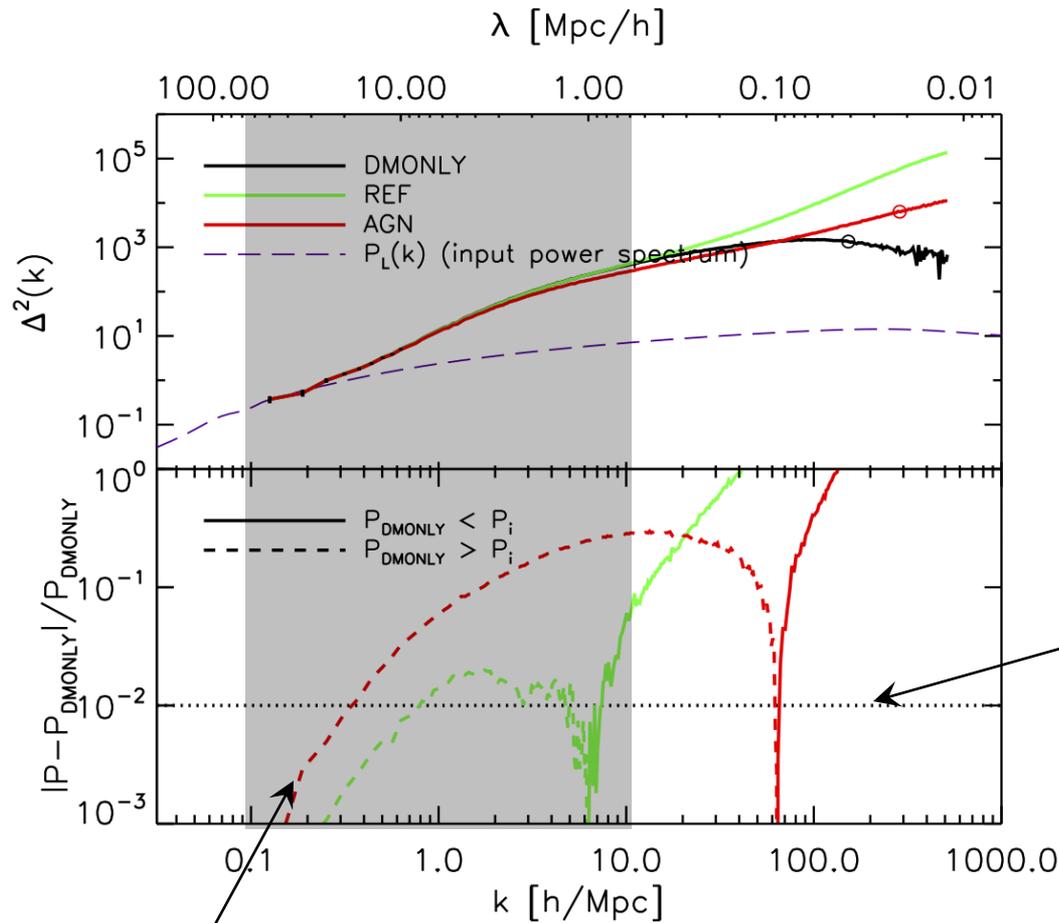
1% difference wrt dark matter only

Range of interest for cosmic shear

Van Daalen, JS+ '11



# Baryons and the matter power spectrum



The feedback required to solve the overcooling problem suppresses power on large scales

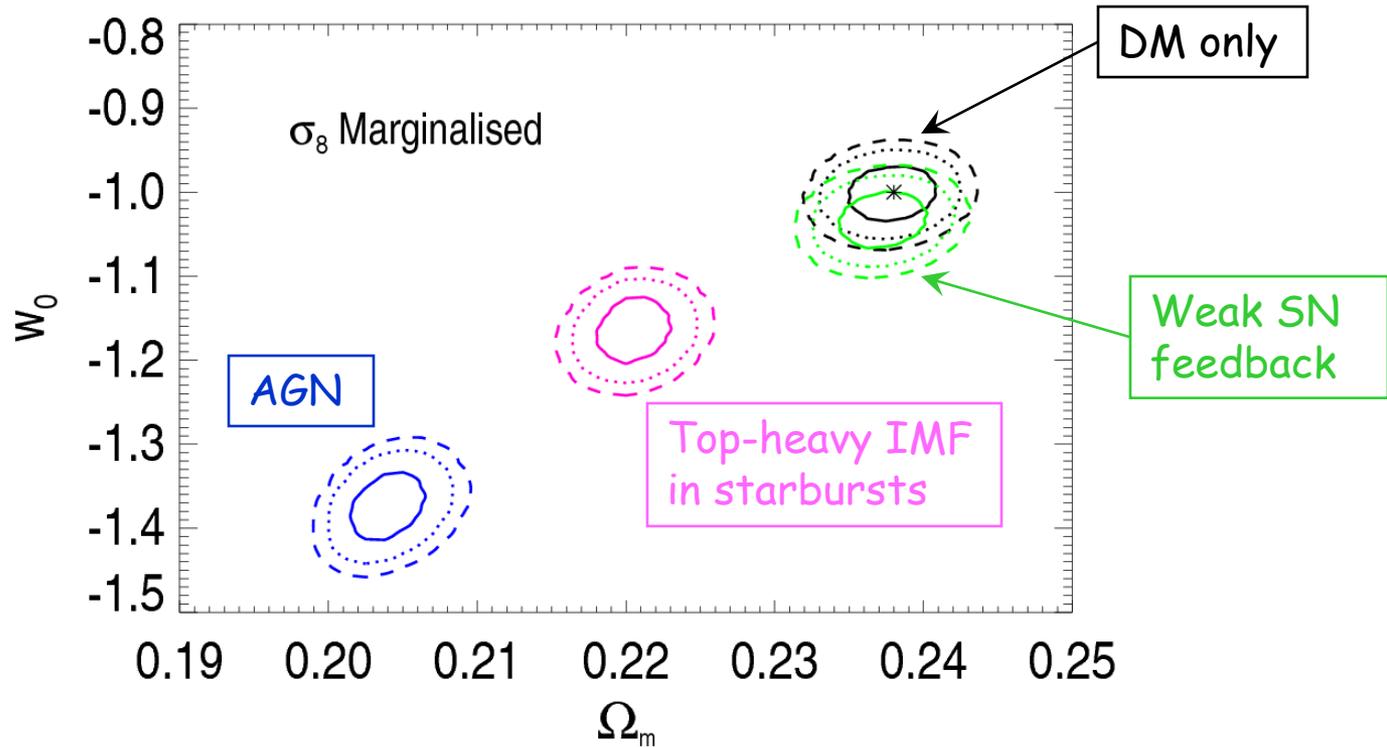
1% difference wrt dark matter only

Range of interest for cosmic shear

Van Daalen, JS+ '11



# Biases due to galaxy formation for a Euclid-like weak lensing survey

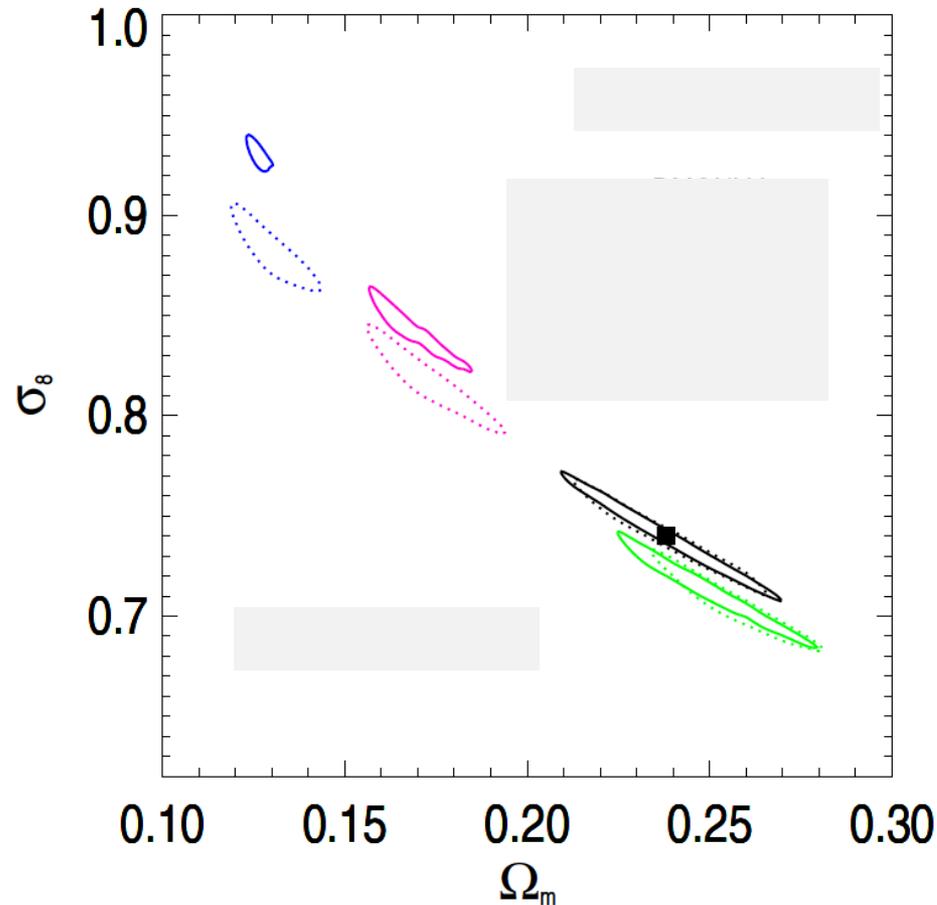


Galaxy formation provides a challenge (target?)  
for weak lensing

Semboloni+ (2011)



# Two and three point statistics



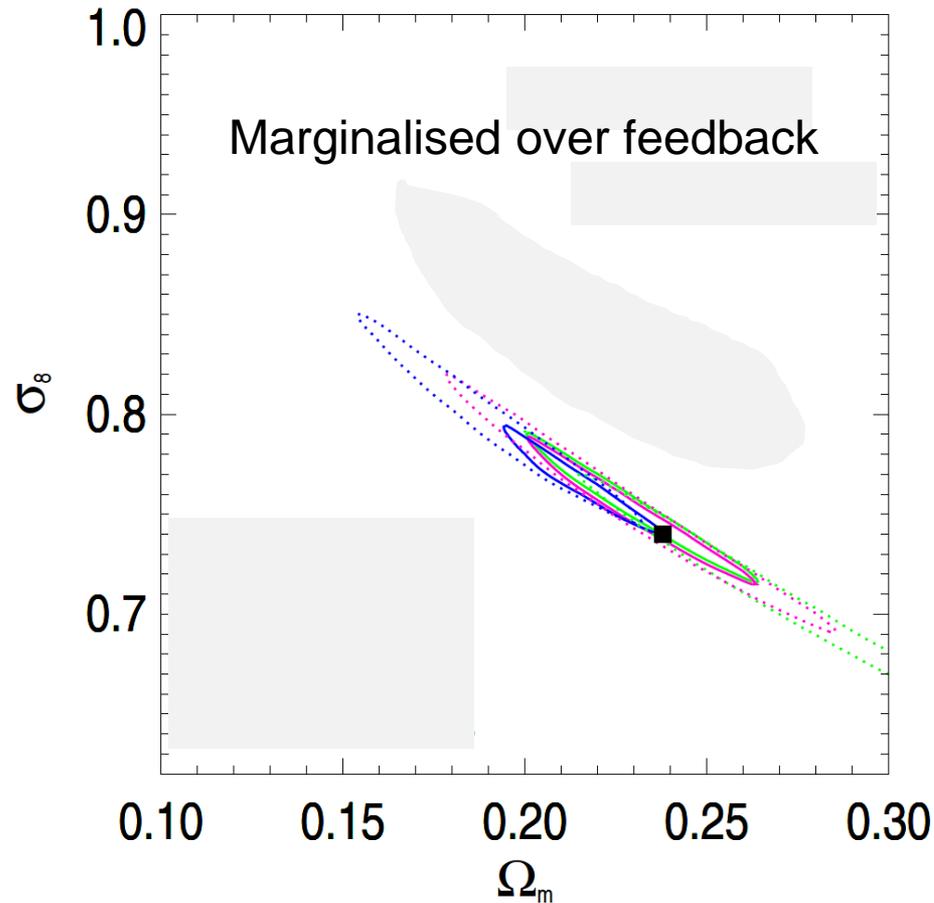
- DMONLY
- AGN
- REF
- DBLIMFV1618

Semboloni, Hoekstra, JS '13

Euclid  
 $w_0$  marginalised



# Two and three point statistics



- DMONLY
- AGN
- REF
- DBLIMFV1618

Semboloni, Hoekstra, JS '13

Euclid  
 $w_0$  marginalised



# Cosmic shear

- Baryonic effects can be dramatic (bias for Euclid ~10s of per cent if untreated)
- Power spectrum affected up to very large scales ( $k > 0.3 \text{ h/Mpc}$ )
- Impact on bispectrum even higher
- Effects dominated by gas ejection
  - Underestimated by models suffering from overcooling
  - Calibration should use gas rather than star fractions
- Modified halo models can capture most of the effects, but currently use input from simulations

# What next for galaxy simulations?

- Resolve cold ISM down to  $10^2$  K in individual (dwarf) galaxy
- Non-equilibrium chemistry including molecules
- Radiation-hydrodynamics

# What next for cosmological simulations?

# EAGLE:

## Evolution and Assembly of GaLaxies and their Environments



- 100 Mpc volume
- $2 \times 1504^3$  particles
- To  $z = 0$
- Resolves warm ISM
- Subgrid recipes dependent only on local hydro quantities
- Feedback efficiency calibrated to match mass function



EAGLE

