

Radiation backgrounds from the first sources and the redshifted 21 cm line

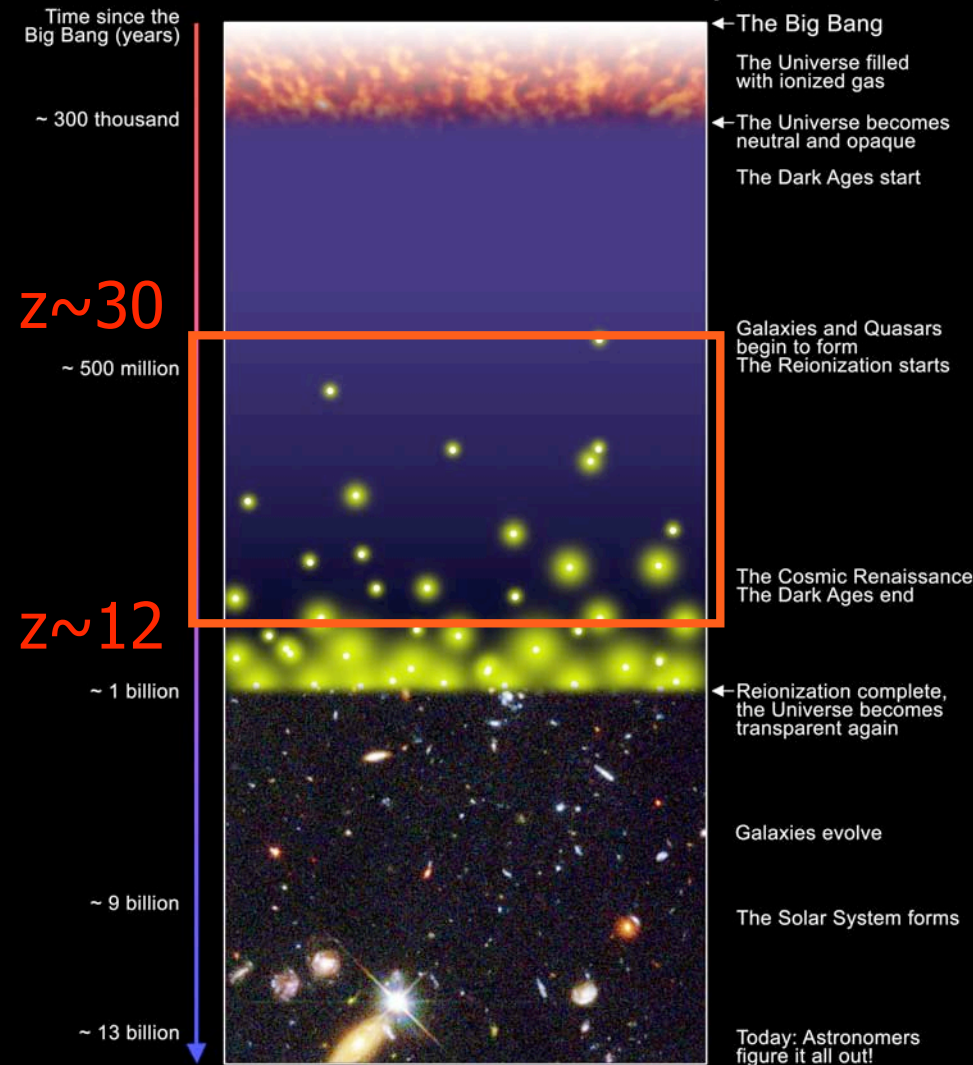
Jonathan Pritchard
(Caltech)

Collaborators: Steve Furlanetto (Yale)

Advisor: Marc Kamionkowski (Caltech)

What is the Reionization Era?

A Schematic Outline of the Cosmic History

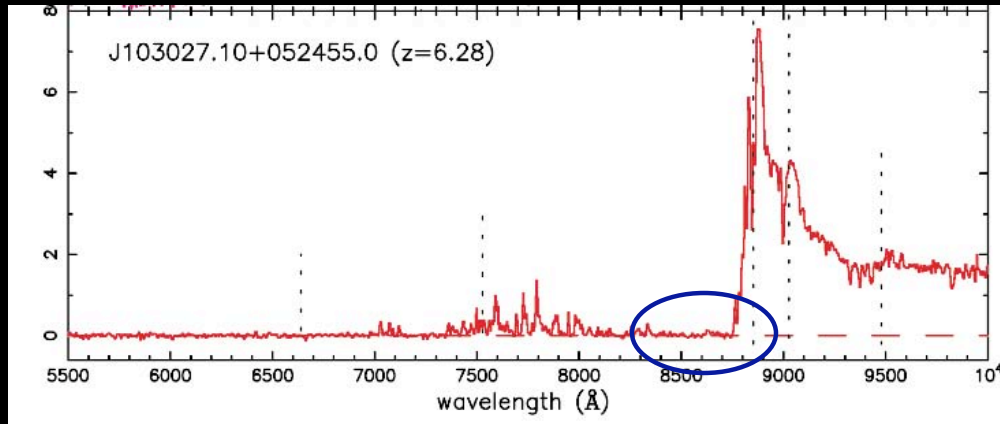


S.G. Djorgovski et al. & Digital Media Center, Caltech

- Evolution of 21 cm signal
- Detecting the first stars through 21 cm fluctuations ($\text{Ly}\alpha$)
- Inhomogeneous X-ray heating and gas temperature fluctuations (X-ray)
- Observational prospects

Ionization history

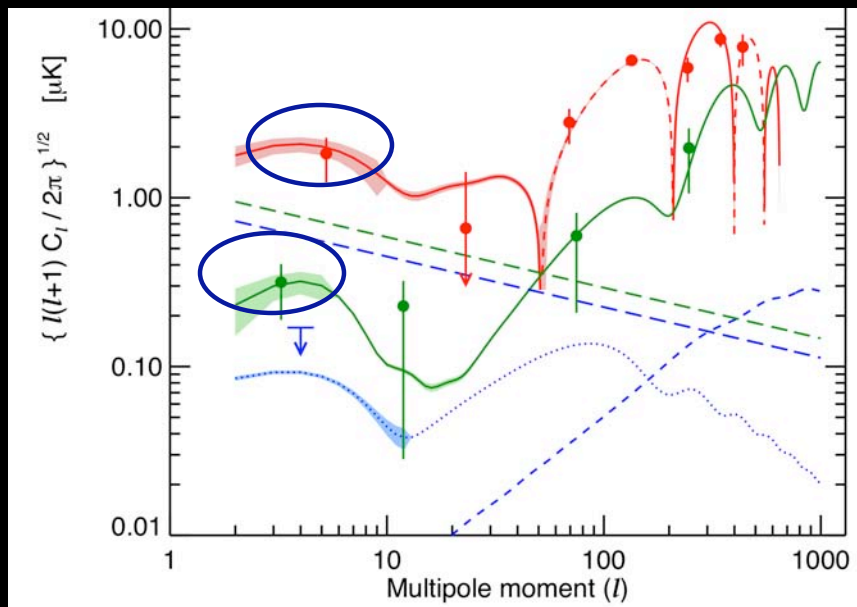
- Gunn-Peterson Trough



Becker et al. 2005

- Universe ionized below $z \sim 6$, some neutral HI at higher z
- black is black

- WMAP3 measurement of $\tau \sim 0.09$ (down from $\tau \sim 0.17$)

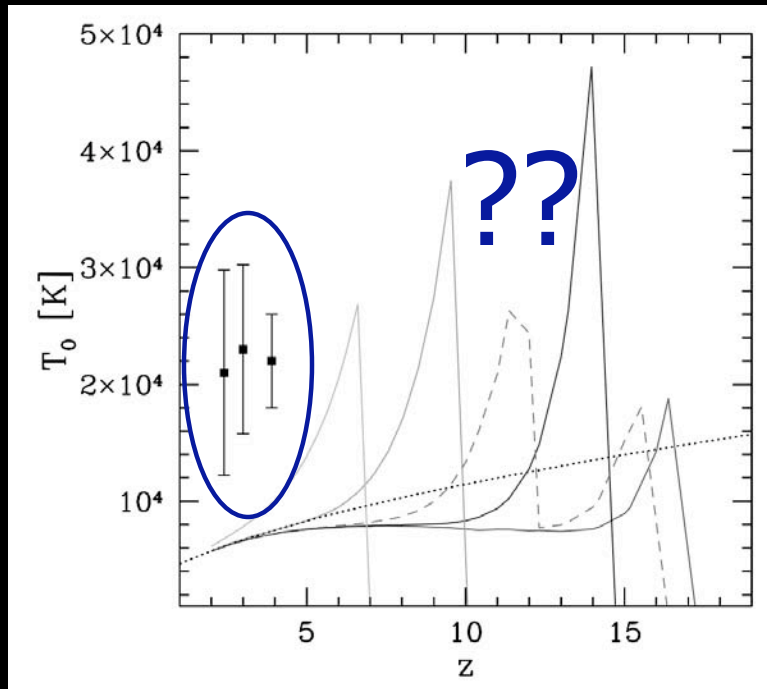


Page et al. 2006

- Integral constraint on ionization history
- Better TE measurements + EE observations

Thermal history

- Ly α forest



Zaldarriaga, Hui, & Tegmark 2001
Hui & Haiman 2003

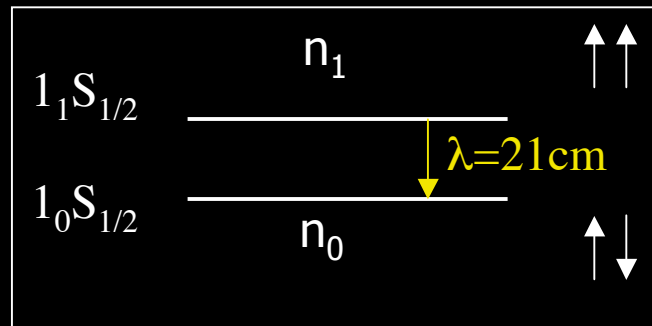
- IGM retains short term memory of reionization - suggests $z_R < 10$
- Photoionization heating erases memory of thermal history before reionization

- CMB temperature

- Knowing $T_{\text{CMB}} = 2.726$ K and assuming thermal coupling by Compton scattering followed by adiabatic expansion allows informed guess of high z temperature evolution

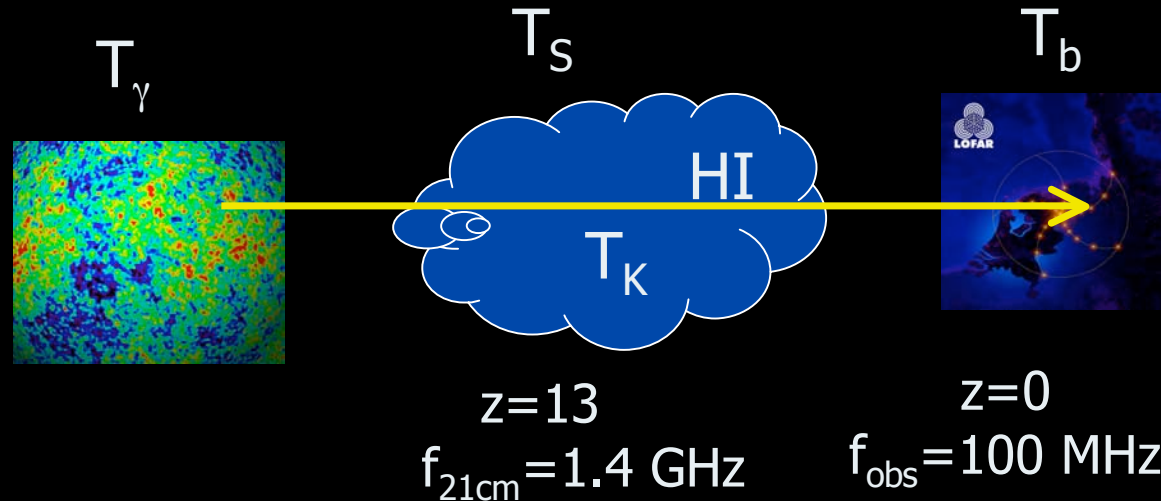
21 cm basics

- HI hyperfine structure



$$n_1/n_0 = 3 \exp(-h\nu_{21\text{cm}}/kT_s)$$

- Use CMB backlight to probe 21cm transition



- 3D mapping of HI possible - angles + frequency

- 21 cm brightness temperature

$$T_b = 27 x_{\text{HI}} (1 + \delta_b) \left(\frac{T_s - T_\gamma}{T_s} \right) \left(\frac{1+z}{10} \right)^{1/2} \text{ mK}$$

- 21 cm spin temperature

$$T_S^{-1} = \frac{T_\gamma^{-1} + x_\alpha T_\alpha^{-1} + x_c T_K^{-1}}{1 + x_\alpha + x_c}$$

Coupling mechanisms:

Radiative transitions (CMB)

Collisions

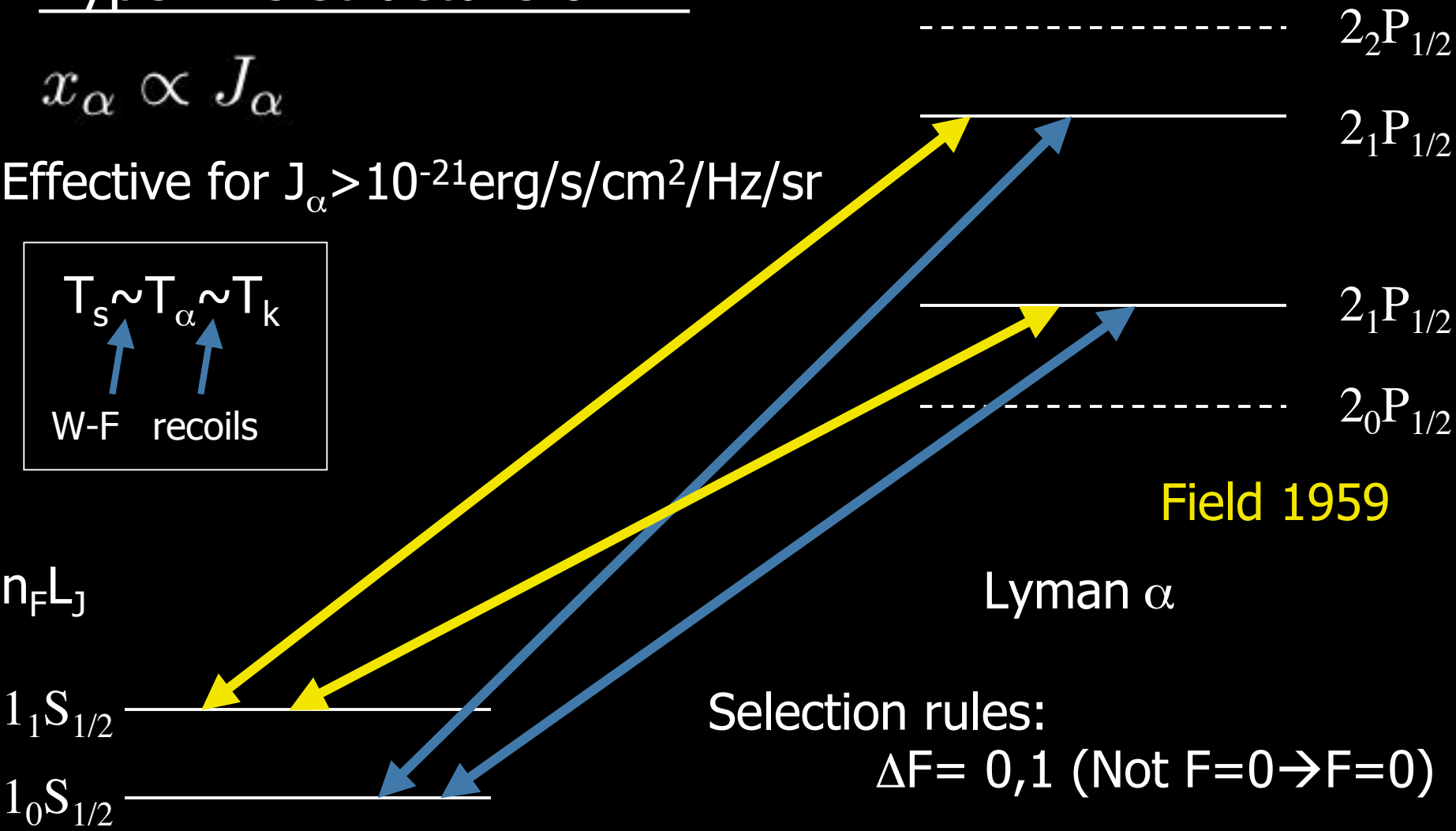
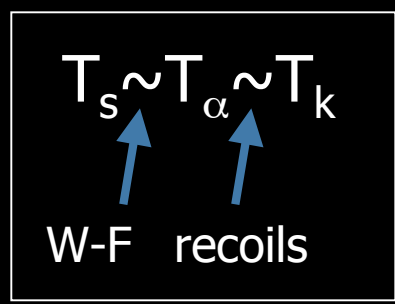
Wouthuysen-Field

Wouthysen-Field effect

Hyperfine structure of HI

$$x_{\alpha} \propto J_{\alpha}$$

Effective for $J_{\alpha} > 10^{-21} \text{erg/s/cm}^2/\text{Hz/sr}$



Field 1959

Lyman α

Selection rules:
 $\Delta F = 0, 1$ (Not $F=0 \rightarrow F=0$)

Experimental efforts

LOFAR: Netherlands
Freq: 120-240 MHz
Baselines: 100m-
100km



MWA: Australia
Freq: 80-300 MHz
Baselines: 10m-
1.5km



PAST/21CMA: China
Freq: 70-200 MHz



PAPER: UCB ?

$(f_{21\text{cm}} = 1.4 \text{ GHz})$

SKA: S Africa/Australia
Freq: 60 MHz-35 GHz
Baselines: 20m-
3000km



Foregrounds

- Many foregrounds
 - Galactic synchrotron (especially polarized component)
 - Radio Frequency Interference (RFI)
e.g. radio, cell phones, digital radio
 - Radio recombination lines
 - Radio point sources
- Foregrounds dwarf signal:
foregrounds ~ 1000 s K vs 10s mK signal
- Strong frequency dependence $T_{\text{sky}} \propto \nu^{-2.6}$
- Foreground removal exploits smoothness in frequency and spatial symmetries

Global history

Furlanetto 2006

$$T_b = T_b(x_{\text{HI}}, T_K, J_\alpha, n_H)$$

$$\frac{dT_K}{dt} = \text{Adiabatic expansion} + \text{X-ray heating} + \text{Compton heating} \quad \text{Heating}$$

$$\frac{dx_i}{dt} = \text{UV ionization} + \text{recombination} \quad \text{HII regions}$$

$$\frac{dx_e}{dt} = \text{X-ray ionization} + \text{recombination} \quad \text{IGM}$$

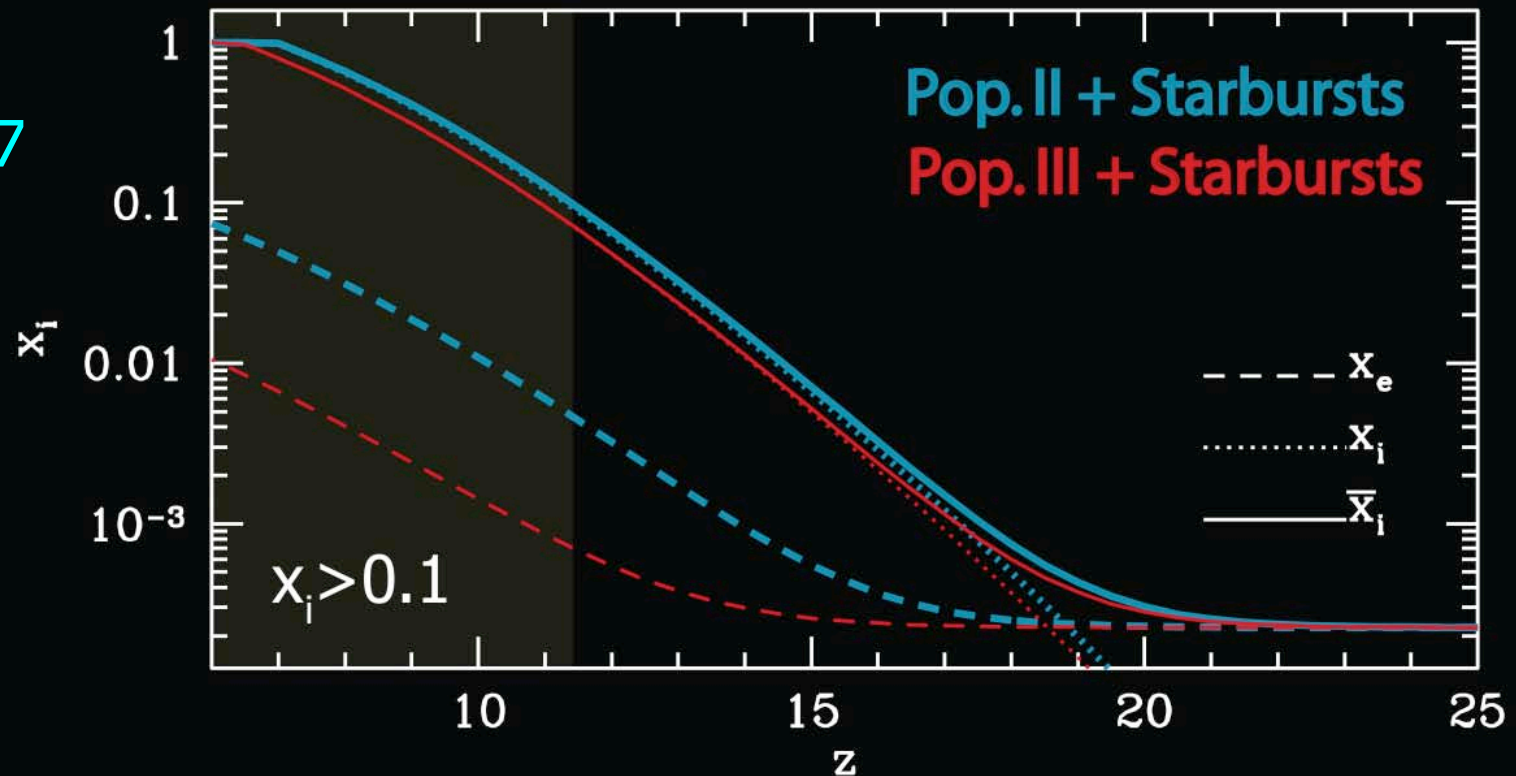
$$J_\alpha = J_C + J_{I,\star} + J_{I,X} \quad \text{Ly}\alpha \text{ flux}$$

continuum injected

- Sources: Pop. II & Pop. III stars (UV+Ly α)
Starburst galaxies, SNR, mini-quasar (X-ray)
- Source luminosity tracks star formation rate
- Many model uncertainties

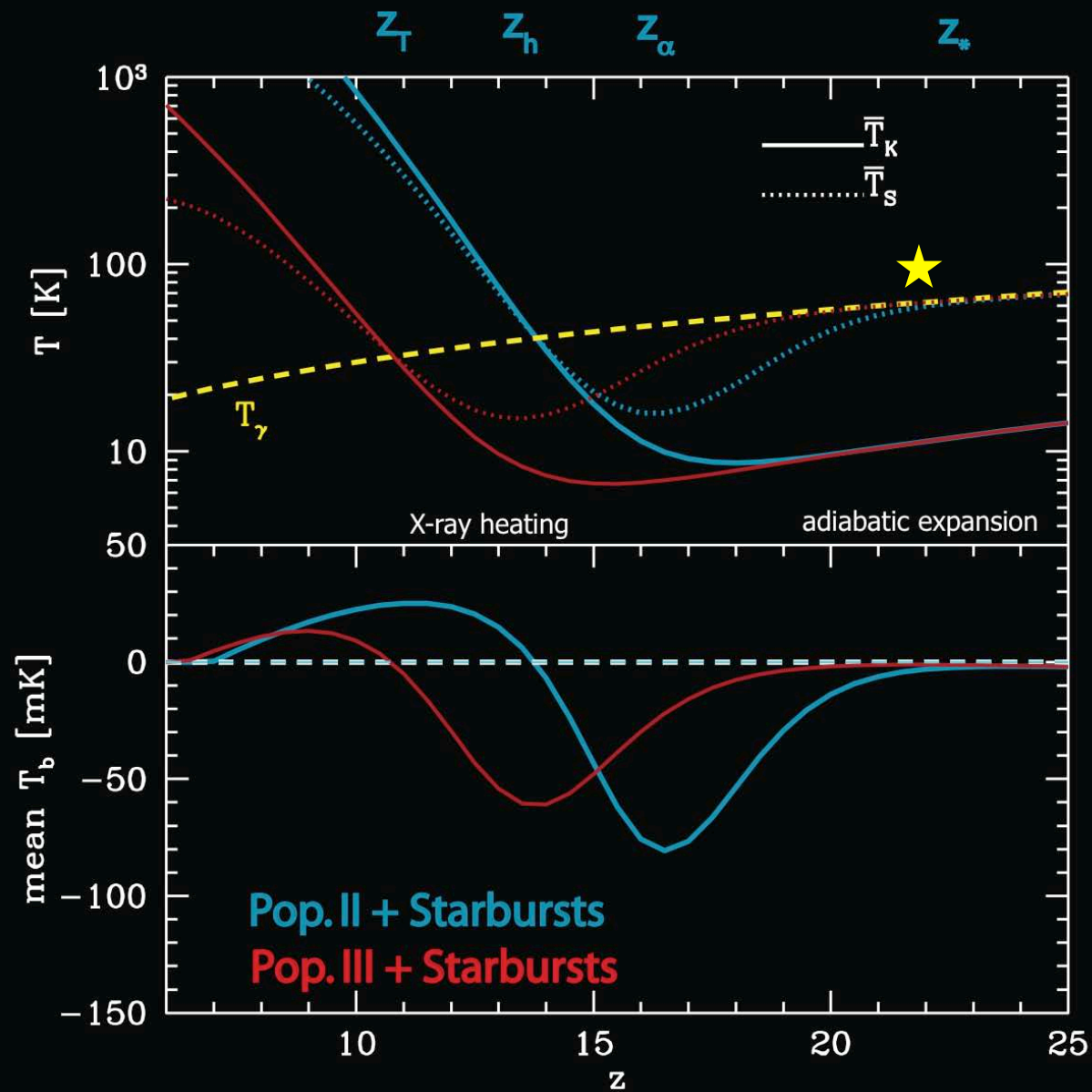
Ionization history

$z_R \sim 7$
 $\tau \sim 0.07$



- Models differ by factor ~ 10 in X-ray/Ly α per ionizing photon
- Reionization well underway at $z < 12$

Thermal history

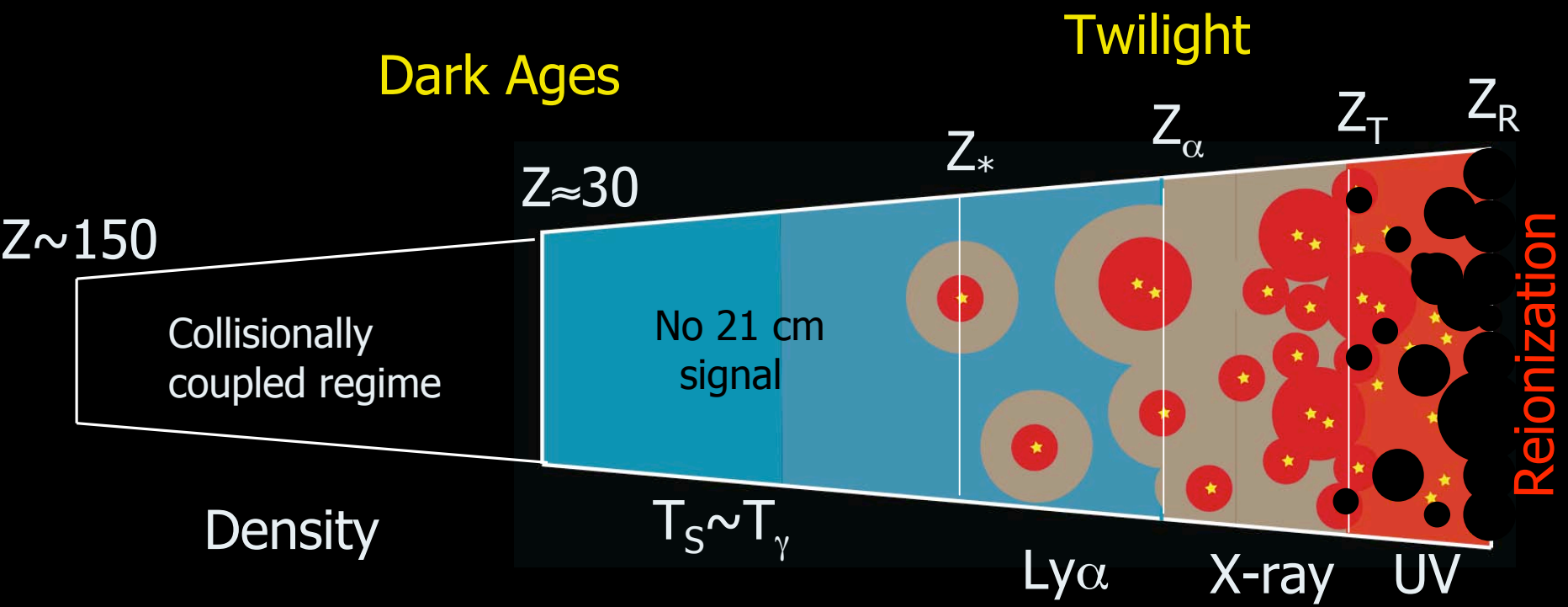


21 cm fluctuations

Brightness temperature Baryon Density Neutral fraction Gas Temperature W-F Coupling Velocity gradient

$$\delta T_b = \beta \delta + \beta_x \delta x_{HI} + \beta_T \delta T_k + \beta_\alpha \delta_\alpha - \delta_{\partial v}$$

Cosmology Reionization X-ray sources Ly α sources Cosmology



Angular separation?

Baryon
Density

Neutral
fraction

Gas
Temperature

W-F
Coupling

Velocity
gradient

$$\delta_{T_b} = \beta\delta + \beta_x\delta_{x_{HI}} + \beta_T\delta_{T_k} + \beta_\alpha\delta_\alpha - \delta_{\partial v}$$

- In linear theory, peculiar velocities correlate with overdensities

$$\delta_{d_r v_r}(k) = -\mu^2 \delta$$

Bharadwaj & Ali 2004

- Anisotropy of velocity gradient term allows angular separation

$$P_{T_b}(\mathbf{k}) = \mu^4 P_{\mu^4} + \mu^2 P_{\mu^2} + P_{\mu^0}$$

Barkana & Loeb 2005

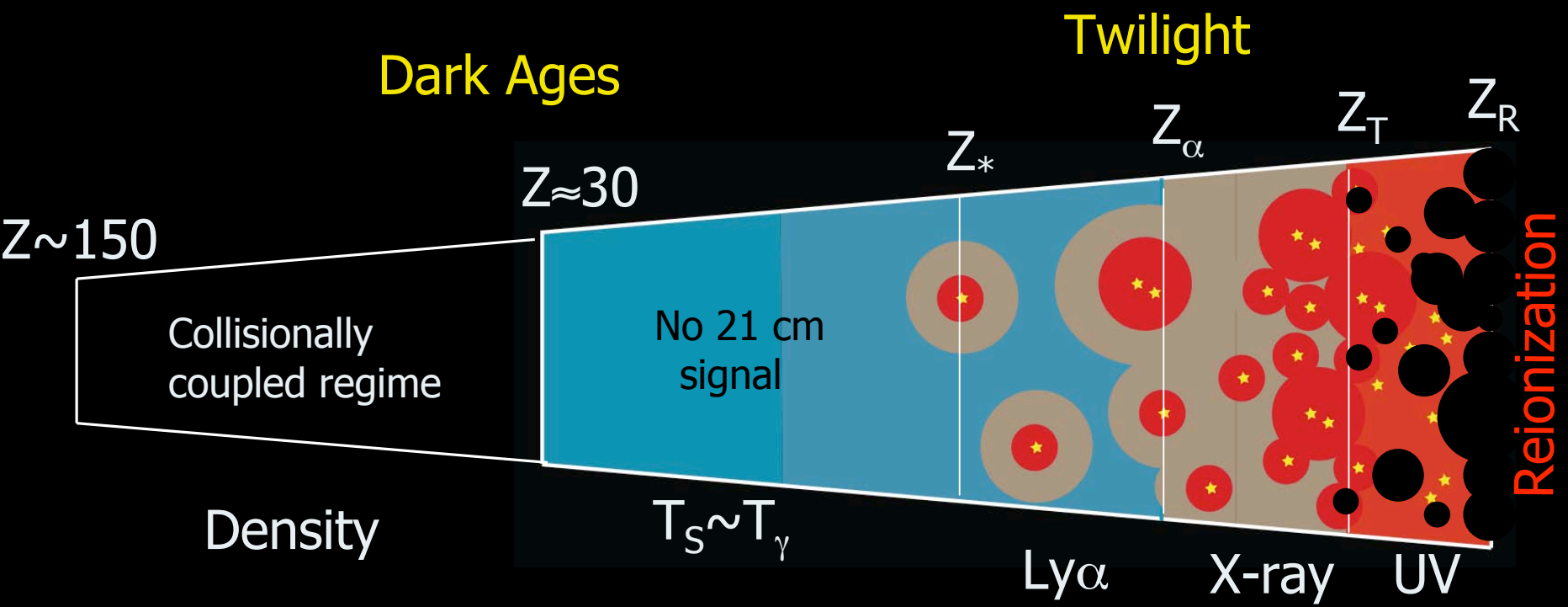
- Initial observations will average over angle to improve S/N

21 cm fluctuations

Brightness temperature Baryon Density Neutral fraction Gas Temperature W-F Coupling Velocity gradient

$$\delta T_b = \beta \delta + \beta_x \delta x_{HI} + \beta_T \delta T_k + \beta_\alpha \delta_\alpha - \delta_{\partial v}$$

Cosmology Reionization X-ray sources Ly α sources Cosmology



21 cm fluctuations: Ly α

Density	Neutral fraction	Gas Temperature	W-F Coupling	Velocity gradient
$\delta T_b = \beta \delta +$	$\beta_x \delta x_{HI}$	$+ \beta_T \delta T_k$	$+ \beta_\alpha \delta_\alpha$	$- \delta \partial v$
	IGM still mostly neutral	-negligible heating of IGM -tracks density	Ly α flux varies	

- Ly α fluctuations unimportant after coupling saturates ($x_\alpha \gg 1$)

$$\beta_\alpha \approx \frac{1}{1 + x_\alpha}$$

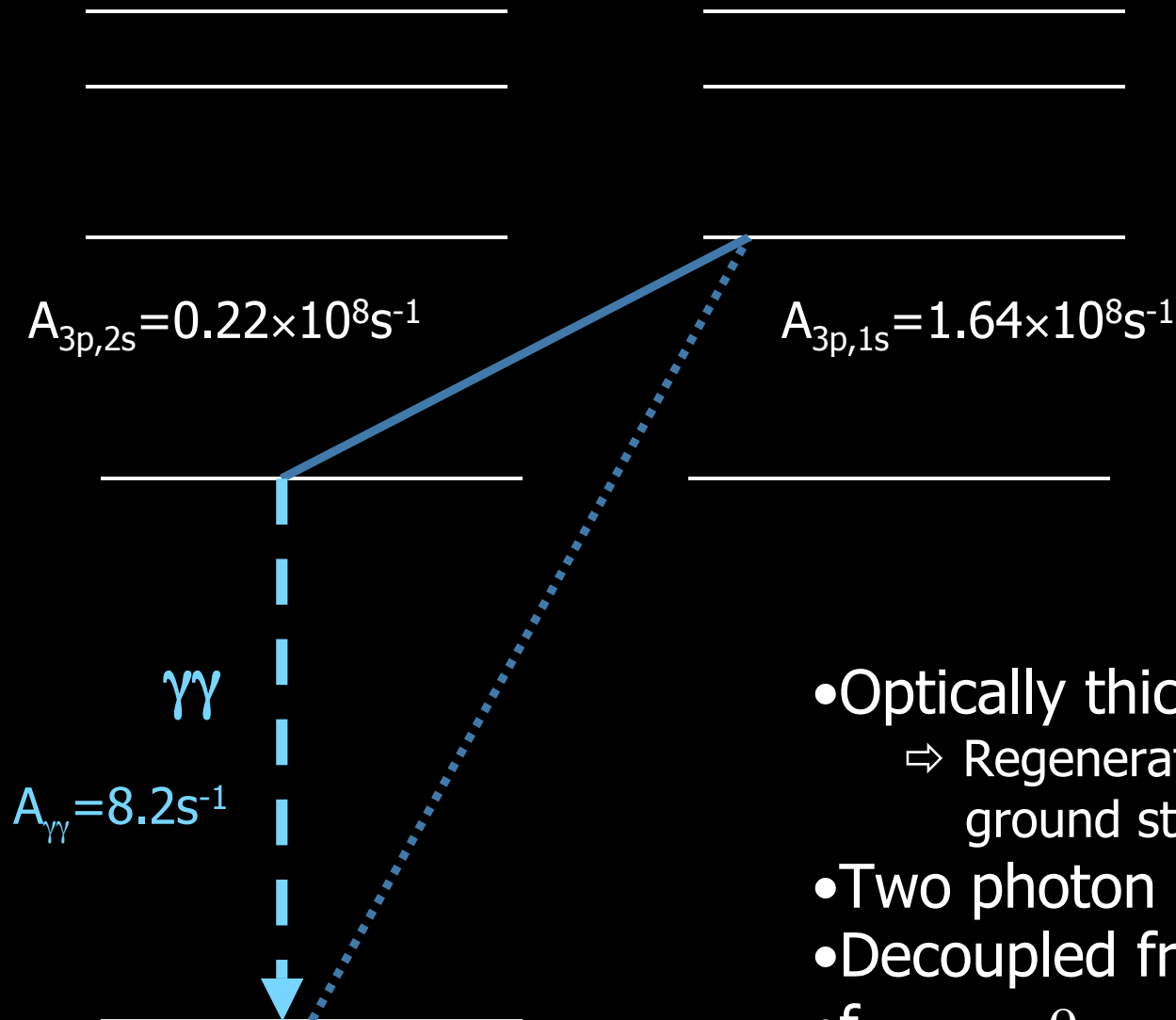
$$x_\alpha \propto J_\alpha$$

- Three contributions to Ly α flux:
 1. Stellar photons redshifting into Ly α resonance
 2. Stellar photons redshifting into higher Lyman resonances
 3. X-ray photoelectron excitation of HI

Higher Lyman series

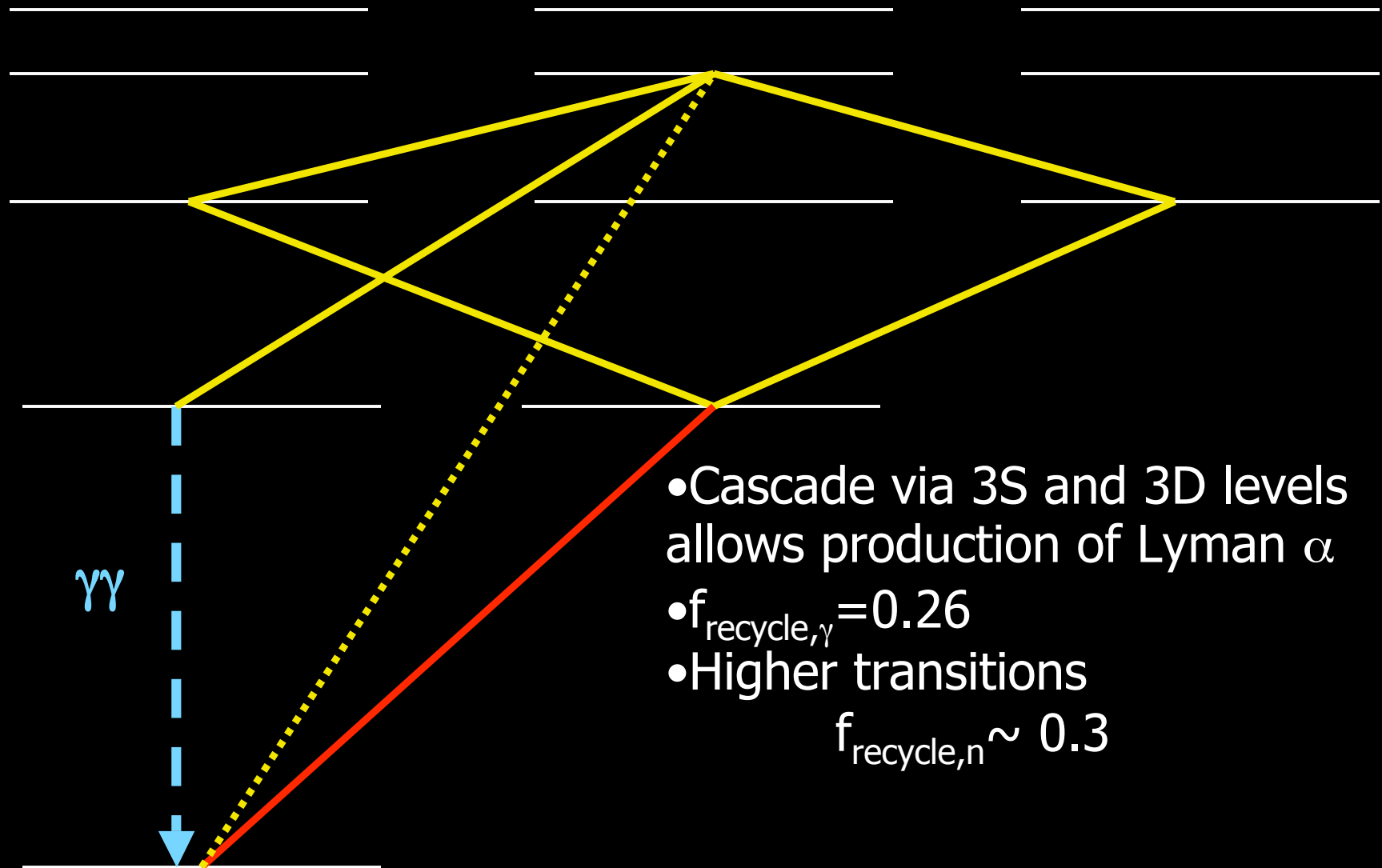
- Two possible contributions from Ly n photons
 - Direct pumping: Analogy of the W-F effect
 - Cascade: Excited state decays through cascade to generate Ly α
- Direct pumping is suppressed by the possibility of conversion into lower energy photons
 - Ly α scatters $\sim 10^6$ times before redshifting through resonance
 - Ly n scatters $\sim 1/P_{\text{abs}} \sim 10$ times before converting
 - ⇒ Direct pumping is not significant
- Cascades end through generation of Ly α or through a two photon decay
 - Use basic atomic physics to calculate fraction recycled into Ly α
 - Discuss this process in the next few slides...

Lyman β



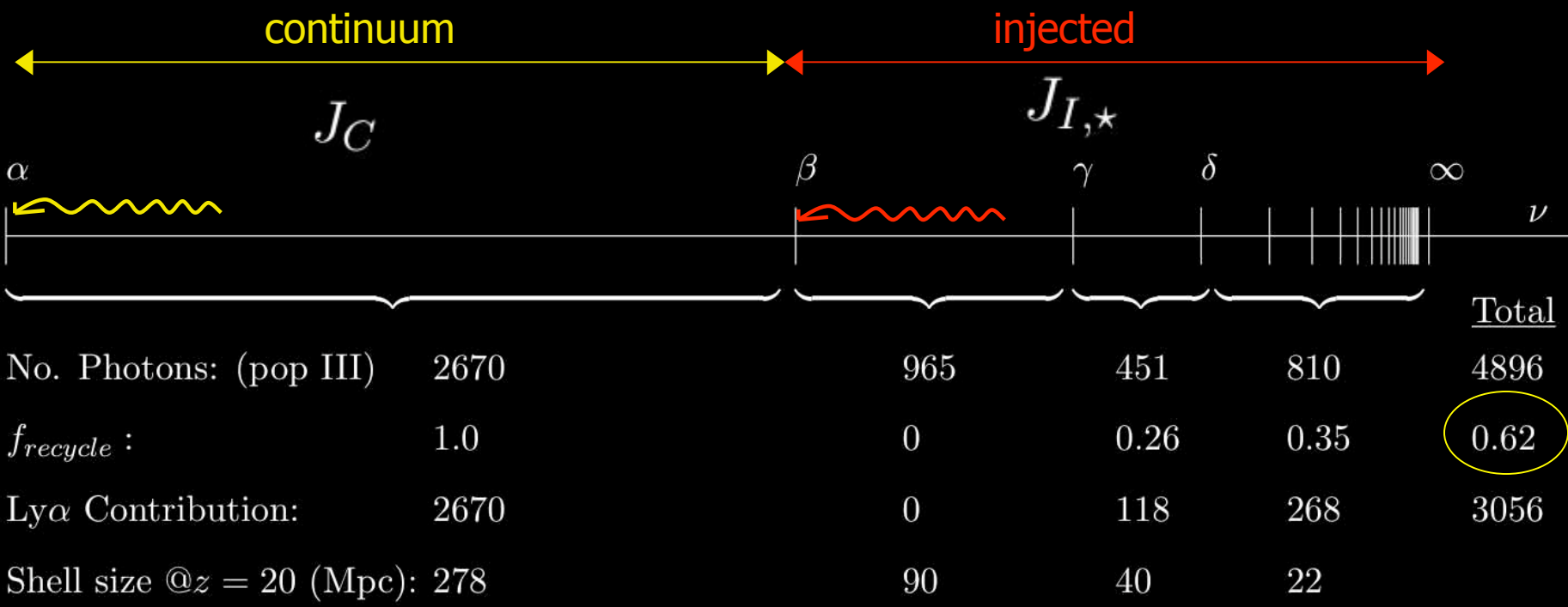
- Optically thick to Lyman series
 \Rightarrow Regenerate direct transitions to ground state
- Two photon decay from $2S$ state
- Decoupled from Lyman α
- $f_{\text{recycle},\beta} = 0$

Lyman γ



- Cascade via 3S and 3D levels allows production of Lyman α
- $f_{\text{recycle},\gamma} = 0.26$
- Higher transitions
 $f_{\text{recycle},n} \sim 0.3$

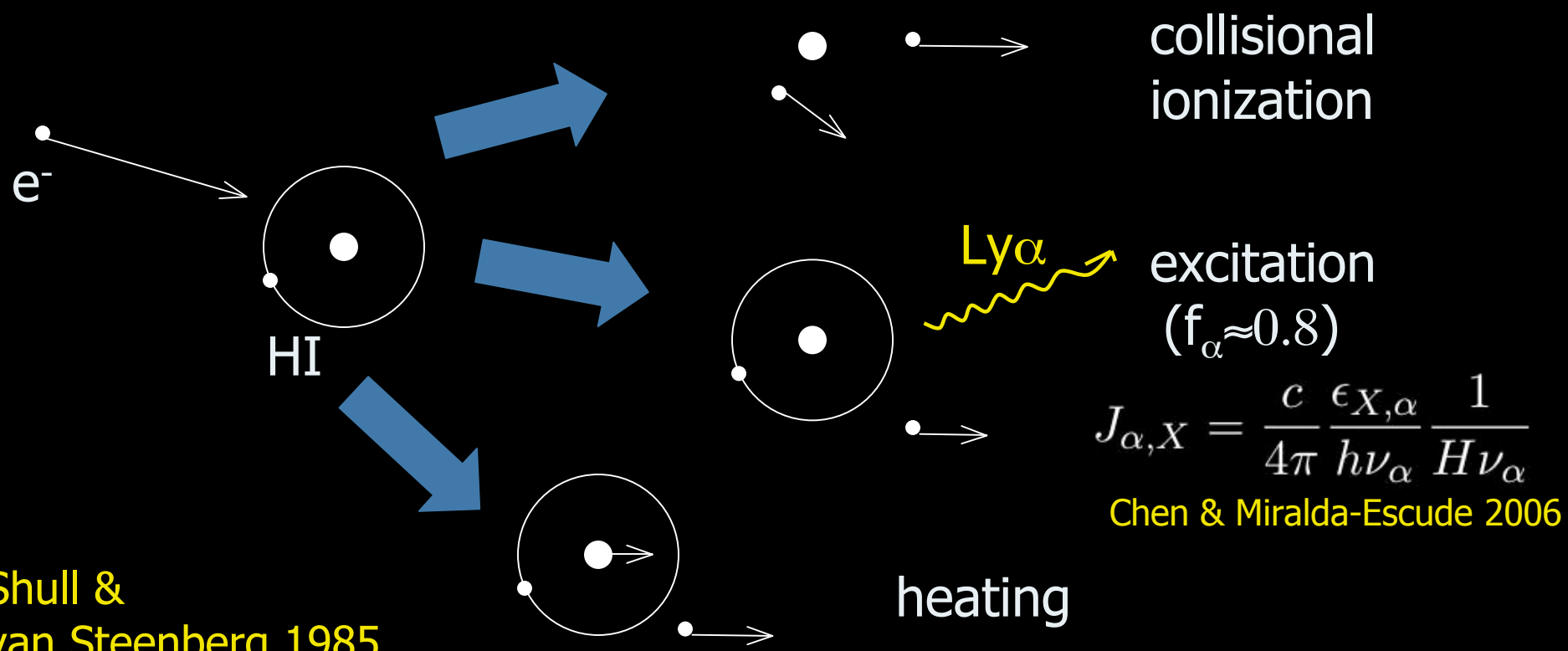
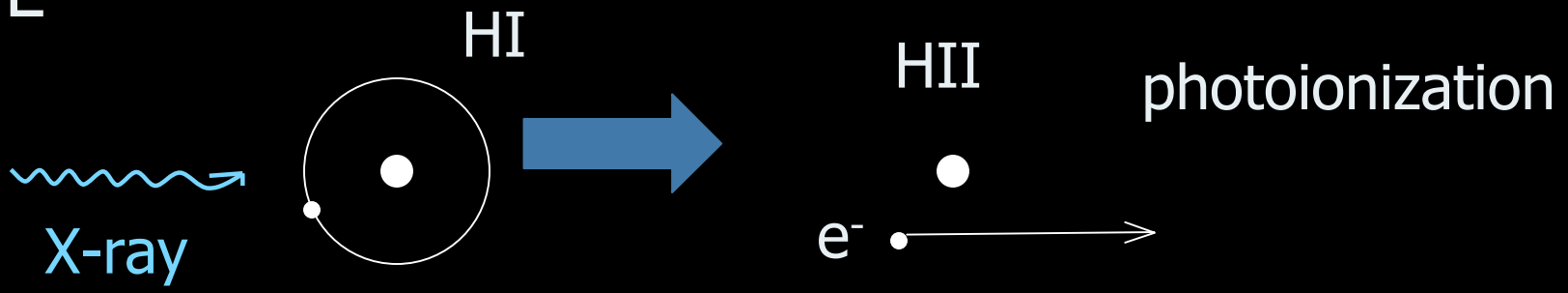
•Stellar contribution



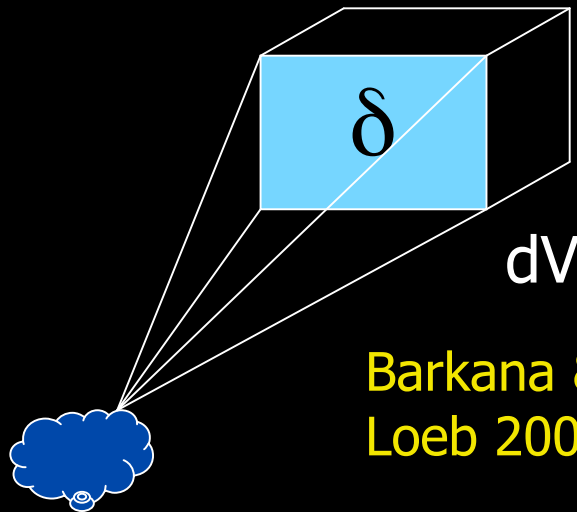
• also a contribution from X-rays...

X-rays and Ly α production

$$\sigma_{\text{pi}} \propto E^{-3}$$



Fluctuations from the first stars



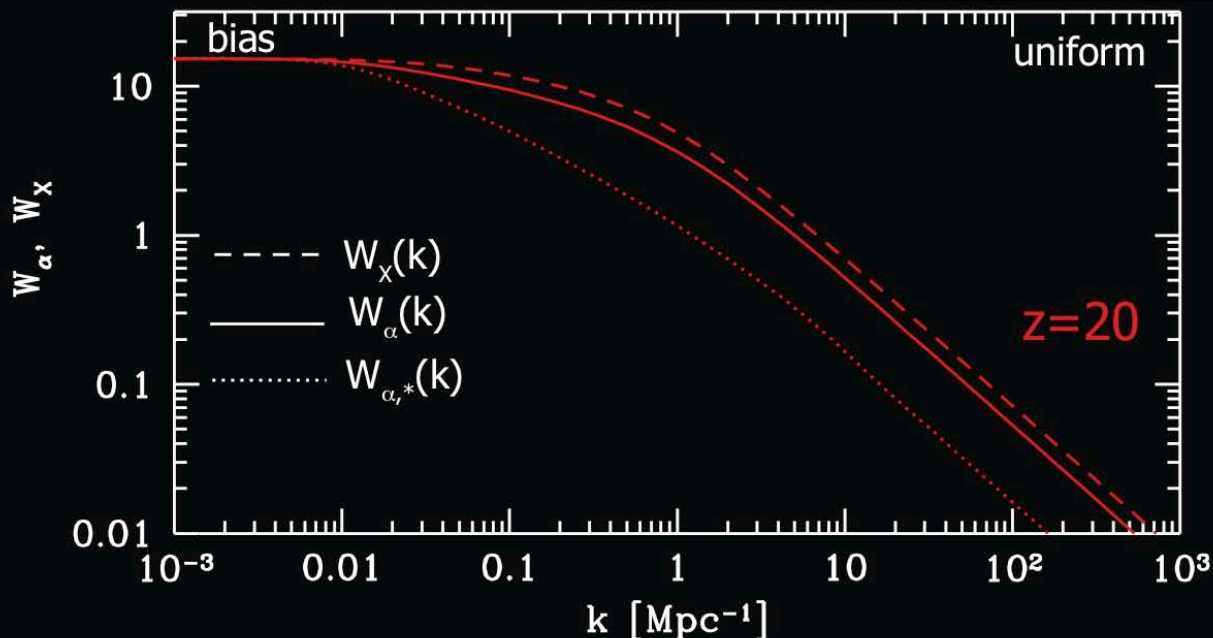
Barkana &
Loeb 2005

- Fluctuations in flux from source clustering, $1/r^2$ law, optical depth,...
- Relate $\text{Ly}\alpha$ fluctuations to overdensities

$$\delta_{x_\alpha}(\mathbf{k}) = W(k)\delta(\mathbf{k})$$

- $W(k)$ is a weighted average

$$W_\alpha = \sum_i W_{\alpha,i} (J_{\alpha,i} / J_\alpha)$$



Determining the first sources

δ_α dominates

Sources

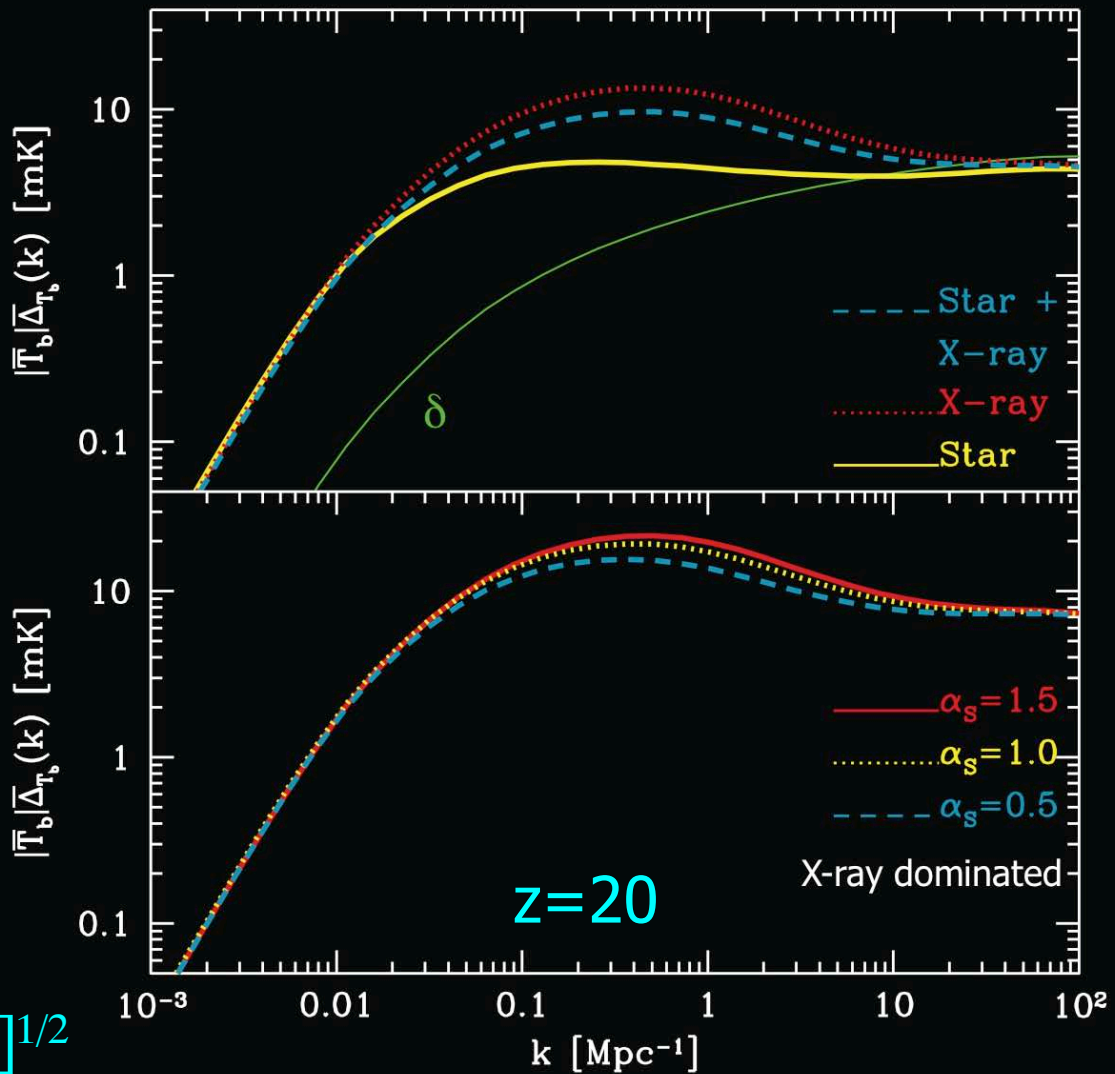
$J_{\alpha,*}$ vs $J_{\alpha,X}$

Spectra

α_S

$\Delta = [k^3 P(k) / 2\pi]^{1/2}$

bias source properties density



Chuzhoy,
Alvarez,
& Shapiro
2006

Pritchard &
Furlanetto
2006

21cm fluctuations: T_K

Density	Neutral fraction	Gas Temperature	W-F Coupling	Velocity gradient
$\delta T_b = \beta \delta +$	$\beta_x \delta x_{HI}$	<div style="border: 1px solid red; padding: 2px;">$+ \beta_T \delta T_k$</div>	$+ \beta_\alpha \delta \alpha$	$- \delta \partial v$
	IGM still mostly neutral	density + x-rays	coupling saturated	

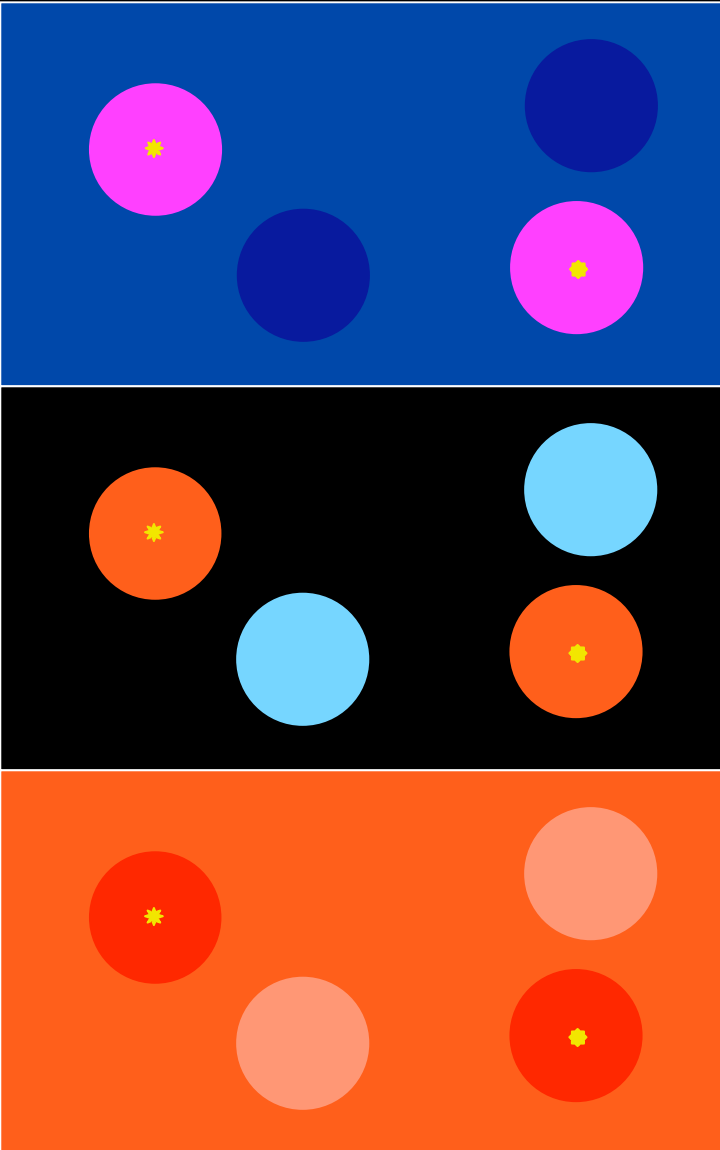
- In contrast to the other coefficients β_T can be negative

$$\beta_T \approx \frac{T_\gamma}{T_K - T_\gamma}$$

- Sign of β_T constrains IGM temperature

Temperature fluctuations

$$T_B = \tau \left(\frac{T_s - T_\gamma}{1 + z} \right)$$



$$T_S \sim T_K < T_\gamma$$

$T_b < 0$ (absorption)

Hotter region = weaker absorption

$$\beta_T < 0$$

$$T_S \sim T_K \sim T_\gamma$$

$$T_b \sim 0$$

21cm signal dominated by
temperature fluctuations

$$T_S \sim T_K > T_\gamma$$

$T_b > 0$ (emission)

Hotter region = stronger emission

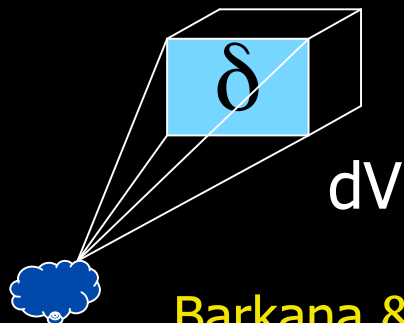
$$\beta_T > 0$$

X-ray heating

- X-rays provide dominant heating source in early universe (shocks possibly important very early on)
- X-ray heating often assumed to be uniform as X-rays have long mean free path

$$\lambda_X \approx 4.9 \bar{x}_{\text{HI}}^{1/3} \left(\frac{1+z}{15} \right)^{-2} \left(\frac{E}{300 \text{ eV}} \right)^3 \text{ Mpc}$$

- Simplistic, fluctuations may lead to observable 21cm signal
- X-ray flux \rightarrow heating rate \rightarrow temperature



Barkana &
Loeb 2005

$$\delta_T = g_T(k, z) \delta$$

↑
adiabatic index -1

Growth of fluctuations

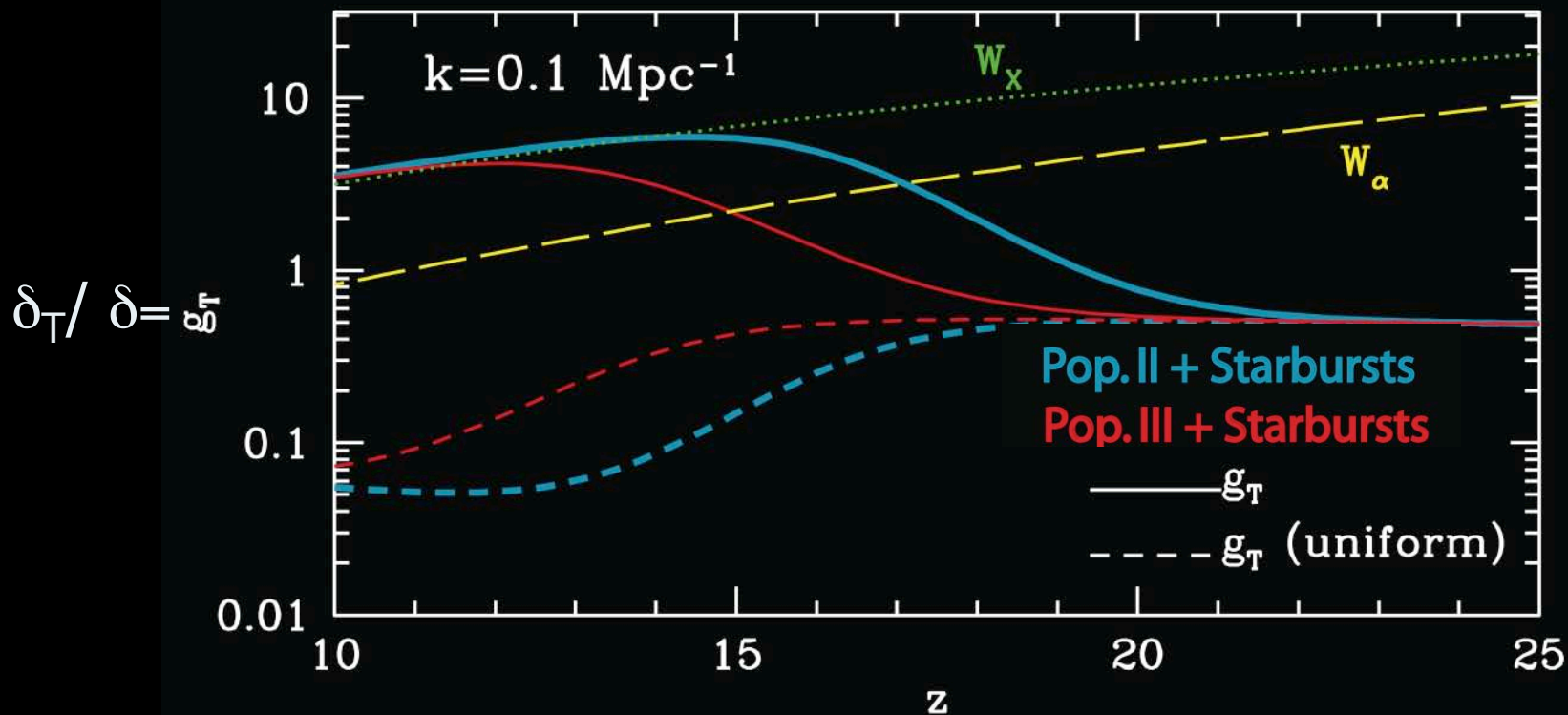
expansion X-rays Compton

$$\frac{dg_T}{dz} = \left(\frac{g_T - 2/3}{1+z} \right) - Q_X(z)[W_X(k) - g_T] - Q_C(z)g_T$$

temperature fluctuations

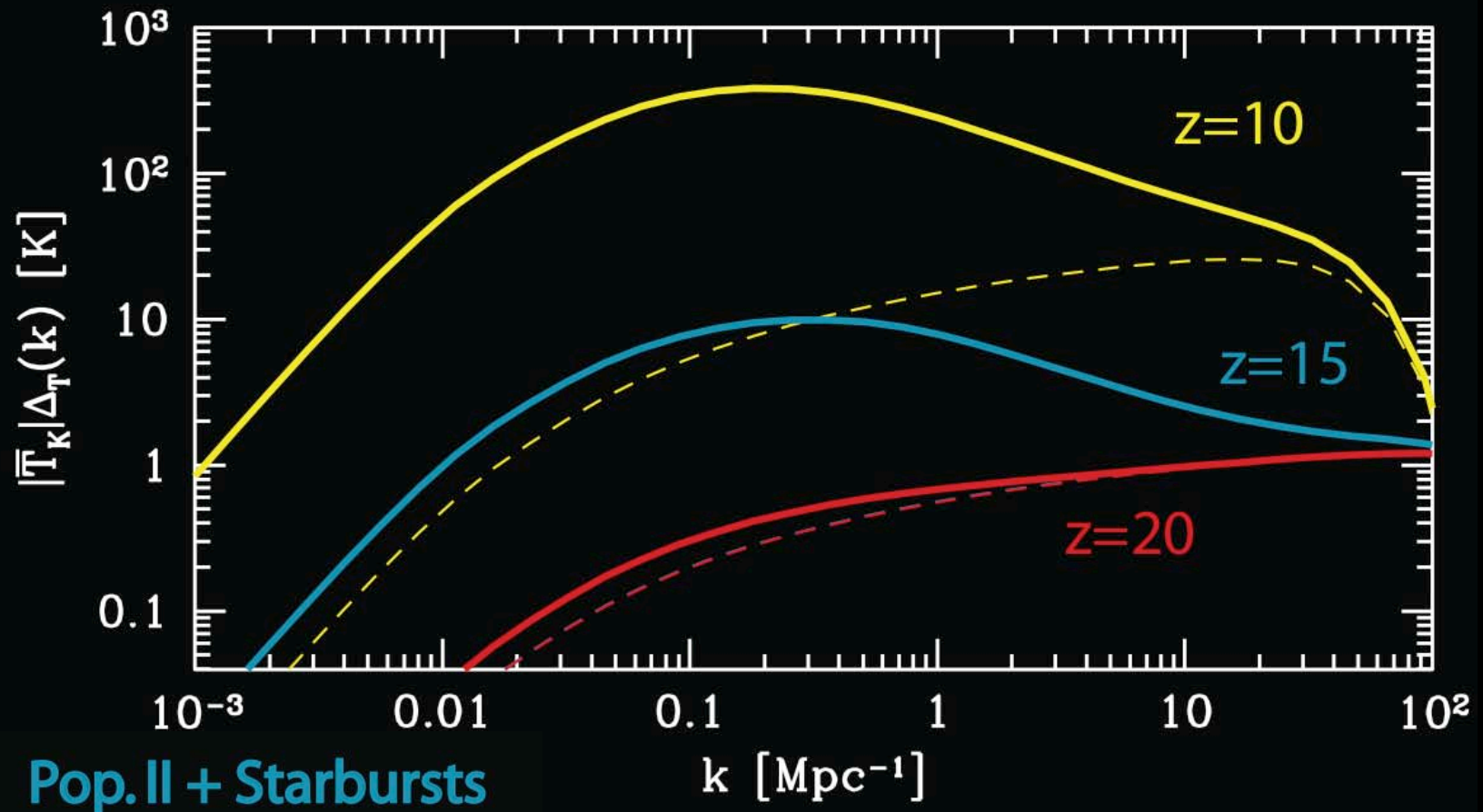
Heating fluctuations

Fractional heating per Hubble time at z



T_K fluctuations

- Fluctuations in gas temperature can be substantial
- Amplitude of fluctuations contains information about IGM thermal history



Indications of T_K

- Constrain heating transition

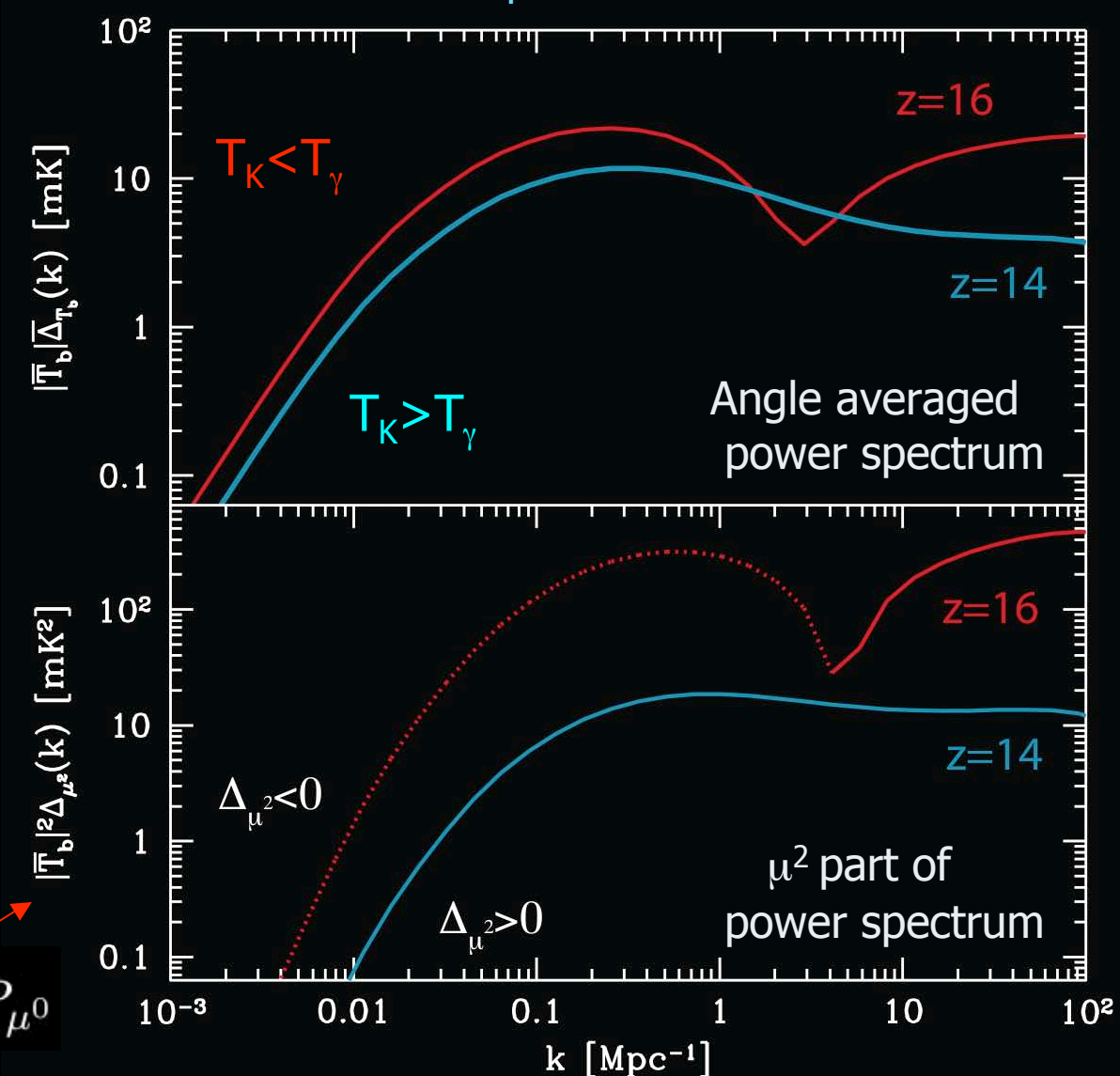
$$\delta_{T_b} \approx \delta + \beta_T \delta_{T_k}$$

$$\beta_T \approx \frac{T_\gamma}{T_K - T_\gamma}$$

- $\Delta_{\mu^2} < 0$ on large scales indicates $T_K < T_\gamma$ (but can have $P_{\delta x} < 0$)

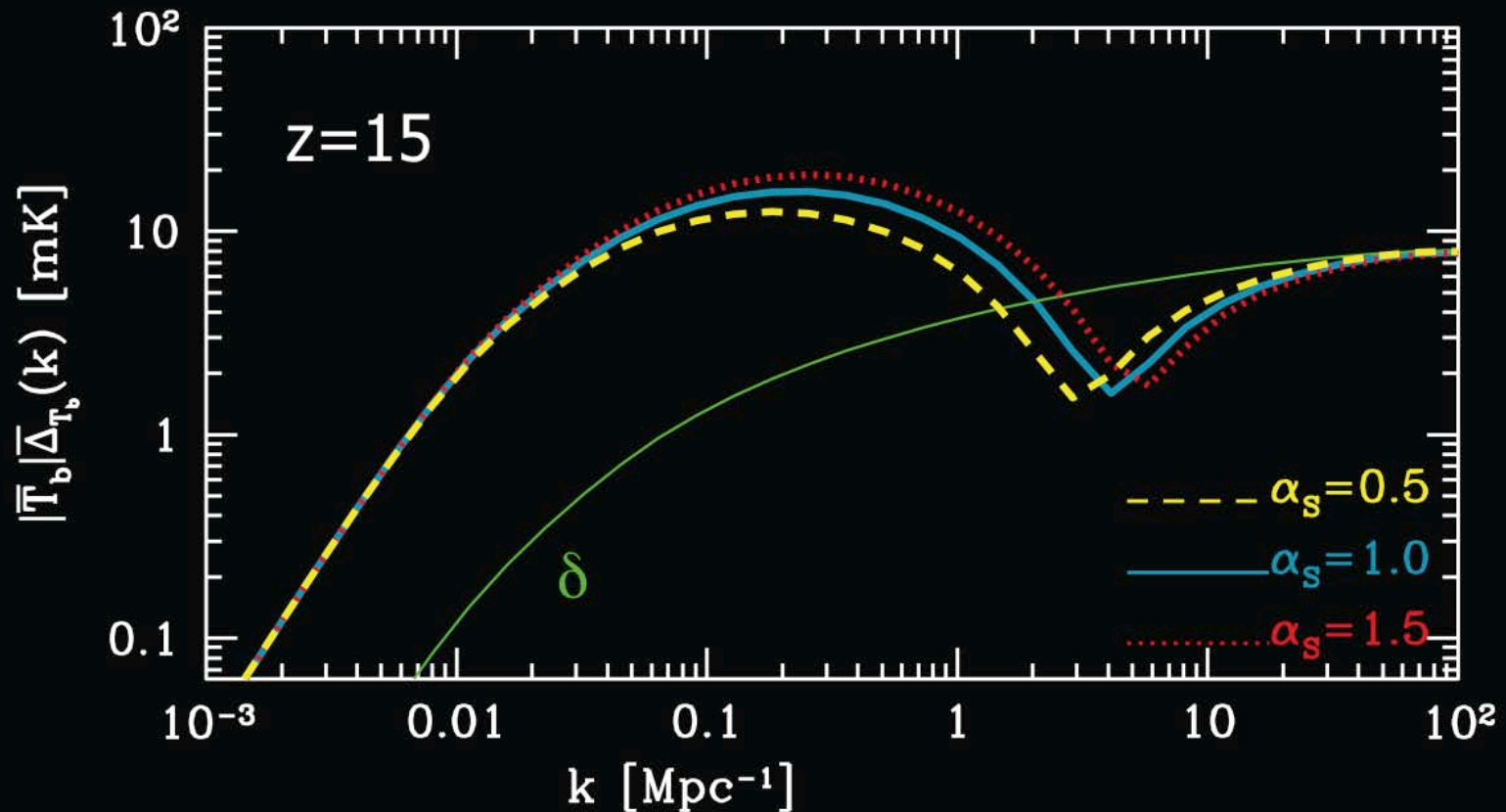
$$P_{T_b}(\mathbf{k}) = \mu^4 P_{\mu^4} + \mu^2 P_{\mu^2} + P_{\mu^0}$$

δ_T dominates



X-ray source spectra

- Sensitivity to α_s through peak amplitude and shape
- Also through position of trough
- Effect comes from fraction of soft X-rays

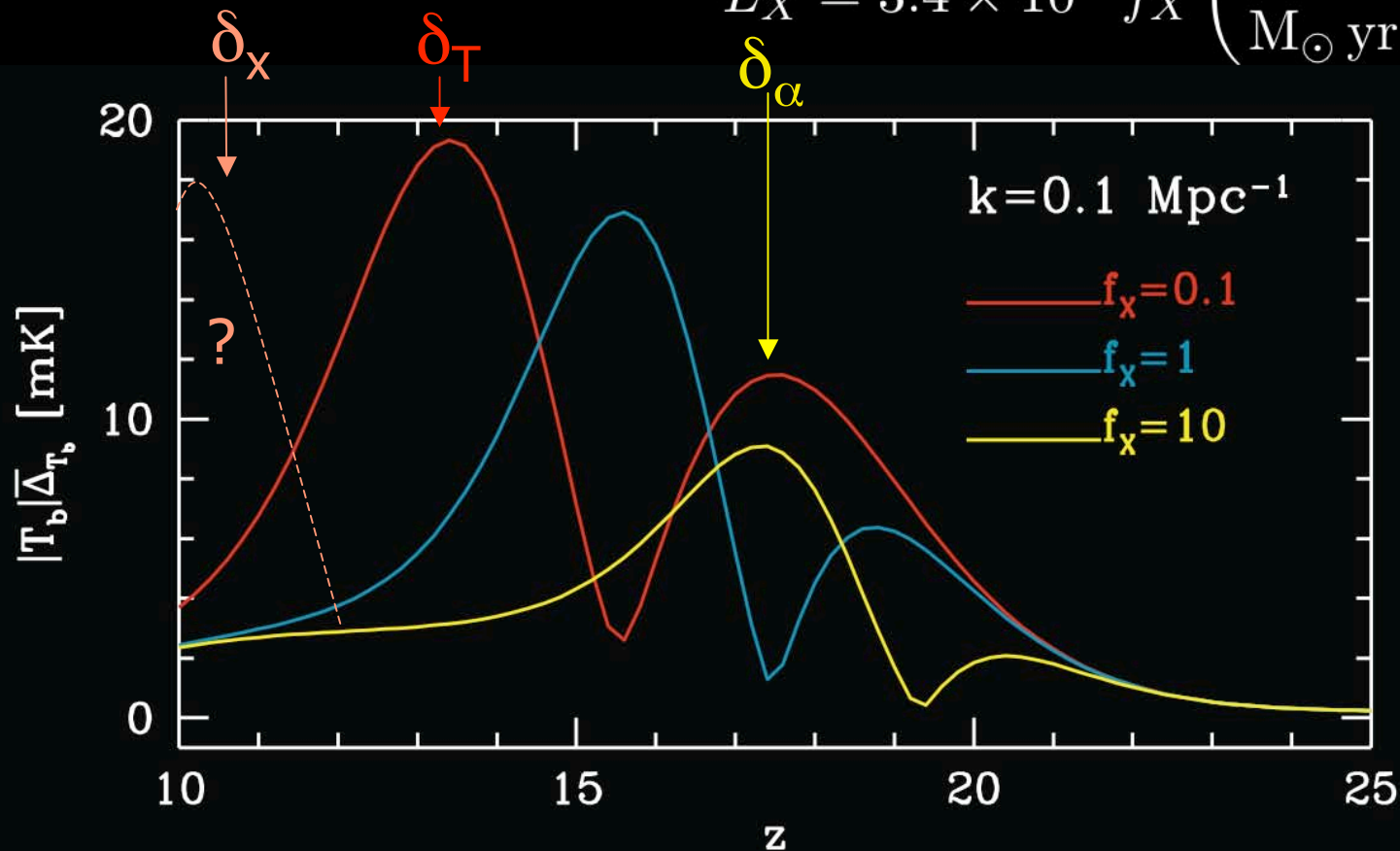


X-ray background?

- X-ray background at high z is poorly constrained

Extrapolating low- z X-ray:IR correlation gives: **Glover & Brand 2003**

$$L_X = 3.4 \times 10^{40} f_X \left(\frac{\text{SFR}}{\text{M}_\odot \text{ yr}^{-1}} \right) \text{ erg s}^{-1}$$



- 1st Experiments might see T_K fluctuations if heating late

Density

Neutral fraction

Gas Temperature

W-F Coupling

Velocity gradient

$$\delta T_b = \boxed{\beta \delta} + \cancel{\beta_\alpha \delta x_{HI}} + \cancel{\beta_T \delta T_k} + \cancel{\beta_\alpha \delta \alpha} - \boxed{\delta \partial v}$$

HII regions still small

IGM hot $T_K \gg T_\gamma$

Ly α coupling saturated

TABLE 2

ERRORS ON COSMOLOGICAL PARAMETER ESTIMATES WHEN DENSITY FLUCTUATIONS DOMINATE THE 21CM SIGNAL FOR TWO YEAR OBSERVATIONS WITH 21 CM INTERFEROMETERS AND IN COMBINATION WITH CURRENT CMB OBSERVATIONS (CCMB) AND WITH PLANCK.^a

	τ	Ω_w	w	$\Omega_m h^2$	$\Omega_b h^2$	n_s	$\delta_H \times 10^5$ ^b	α_s	Ω_ν	\bar{x}_H
	0.1	0.7	-1.0	0.14	0.022	1.0	3.91	0.0	0.0	1.0
LOFAR	-	0.09	-	0.14	0.04	0.14	3.9	-	-	-
MWA	-	0.10	-	0.13	0.03	0.13	3.9	-	-	-
MWA5000	-	0.007	-	0.011	0.003	0.03	0.31	0.012	0.008	-
SKA	-	0.005	-	0.011	0.003	0.06	0.42	0.017	0.016	-
SKA ^c	-	0.14	-	0.051	0.003	0.07	2.4	0.020	0.09	-
SKA [*]	-	0.005	-	0.009	0.002	0.04	0.26	0.011	0.009	-
MWA50K [*]	-	0.002	-	0.005	0.001	0.01	0.11	0.004	0.003	-
CCMB	0.060	0.084	-	0.017	0.0014	0.072	0.29	0.039	0.12	-
CCMB+ LOFAR	0.058	0.058	-	0.011	0.0012	0.031	0.22	0.025	0.03	0.2
CCMB+ MWA	0.057	0.058	-	0.011	0.0012	0.033	0.22	0.025	0.02	0.2
CCMB+ MWA5000	0.049	0.007	-	0.003	0.0009	0.013	0.18	0.006	0.005	0.06
CCMB + SKA	0.049	0.005	-	0.003	0.0009	0.014	0.18	0.005	0.007	0.06
Planck	0.0050	0.029	0.09	0.0023	0.00018	0.0047	0.026	0.008	0.010	-
Planck +MWA5000	0.0046	0.019	0.07	0.0011	0.00013	0.0034	0.018	0.004	0.003	0.05
Planck + SKA	0.0046	0.022	0.08	0.0009	0.00013	0.0034	0.018	0.003	0.004	0.06
Planck + SKA [*]	0.0046	0.018	0.07	0.0009	0.00013	0.0033	0.018	0.003	0.004	0.05
Planck + MWA50K [*]	0.0045	0.008	0.03	0.0004	0.00010	0.0029	0.015	0.002	0.001	0.02

<- Optimistic

improve n_s and Ω_ν most

Reionization

Density

Neutral fraction

Gas Temperature

W-F Coupling

Velocity gradient

$$\delta T_b = \beta \delta + \boxed{\beta_x \delta x_{HI}} + \cancel{\beta_T \delta T_k} + \cancel{\beta_\alpha \delta \alpha} - \delta_{\partial v}$$

HII regions large

IGM hot $T_K \gg T_\gamma$

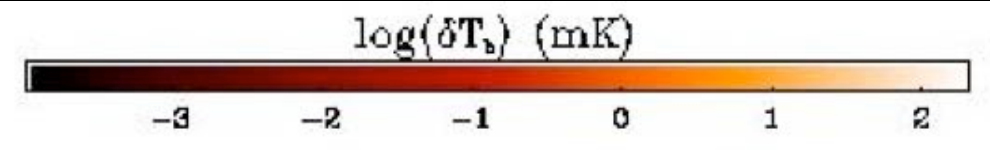
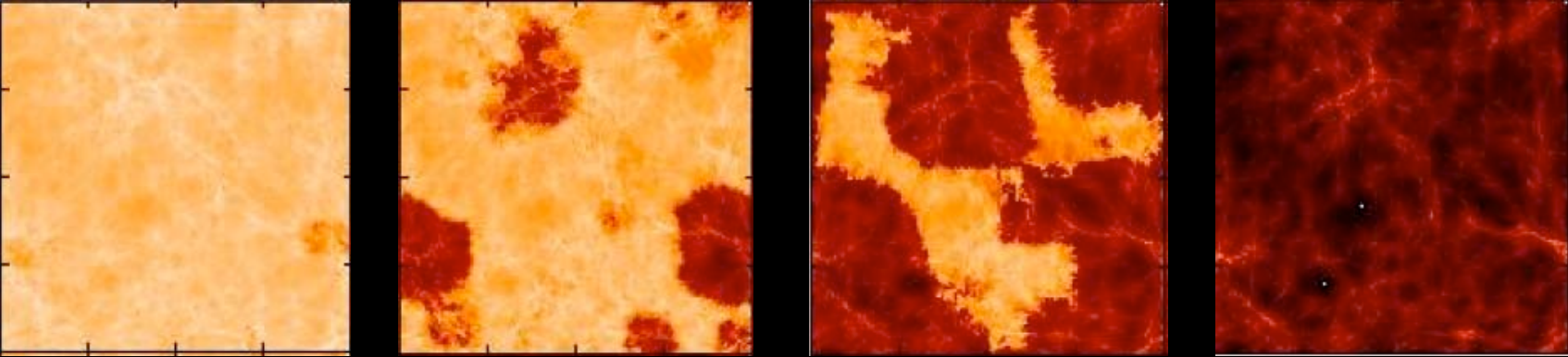
Ly α coupling saturated

Z=12.1

Z=9.2

Z=8.3

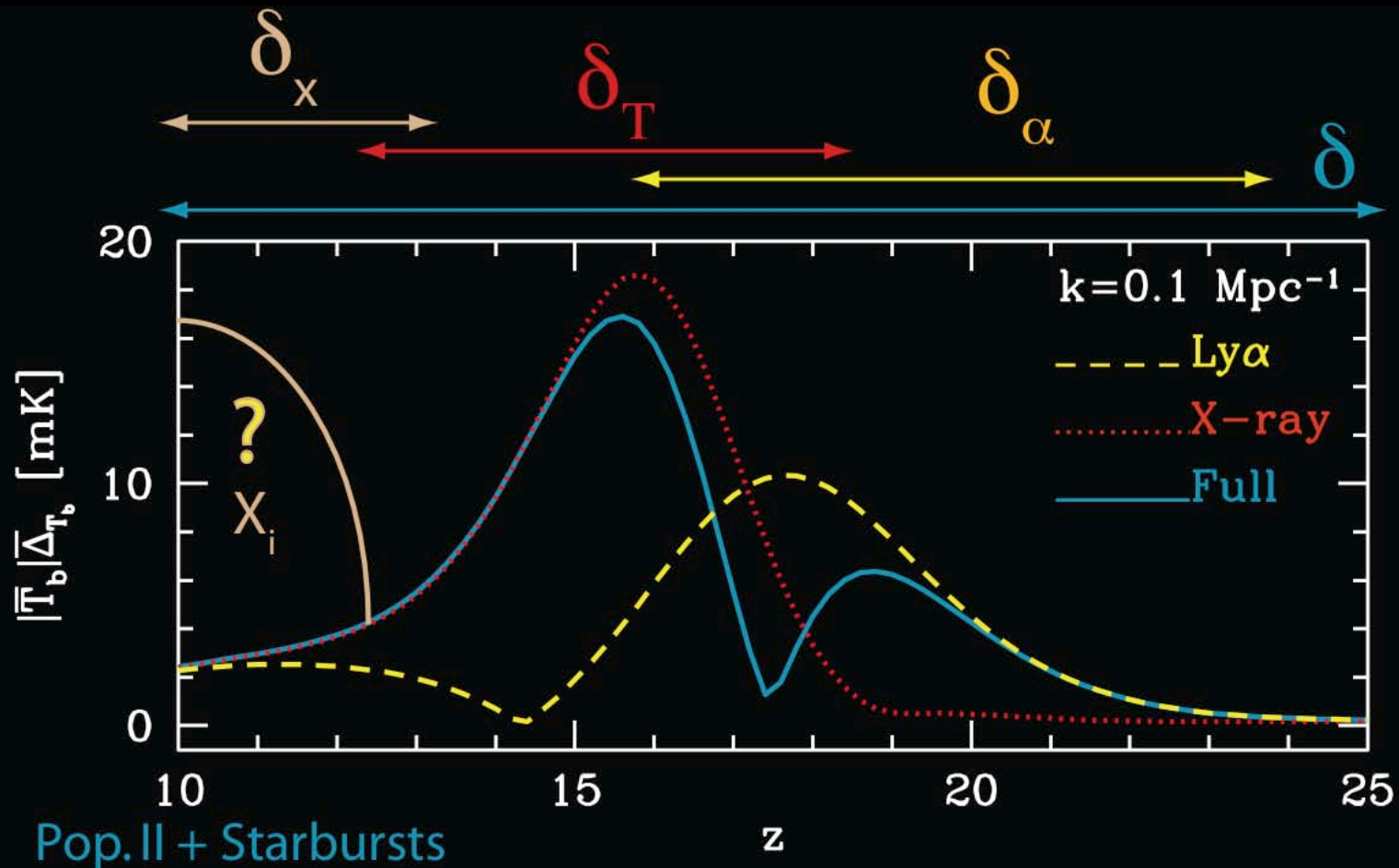
Z=7.6



Furlanetto, Sokasian, Hernquist 2003

21 cm fluctuations: z

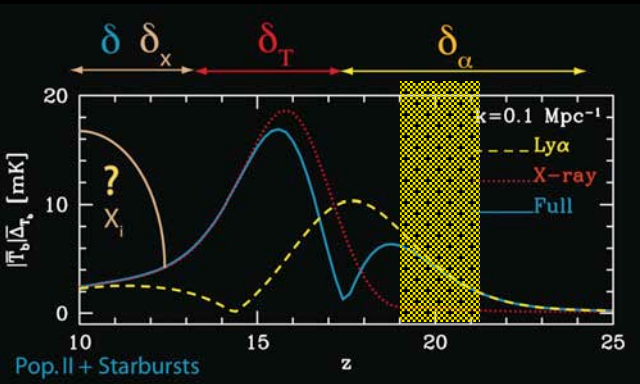
$$\delta T_b = \beta \delta + \beta_x \delta_{x_{HI}} + \beta_T \delta_{T_k} + \beta_\alpha \delta_\alpha - \delta_{\partial v}$$



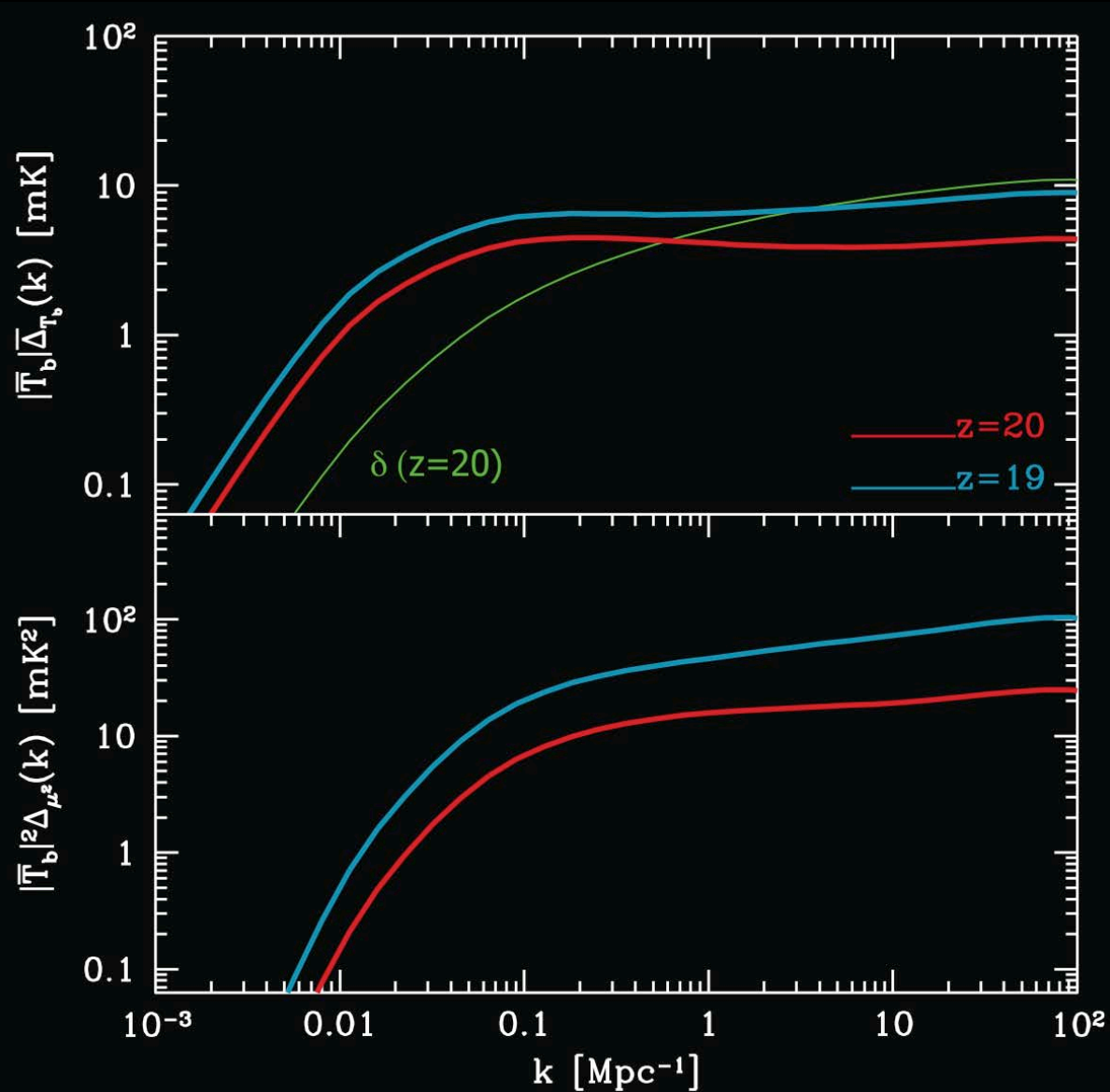
- Exact form very model dependent

Redshift slices: Ly α

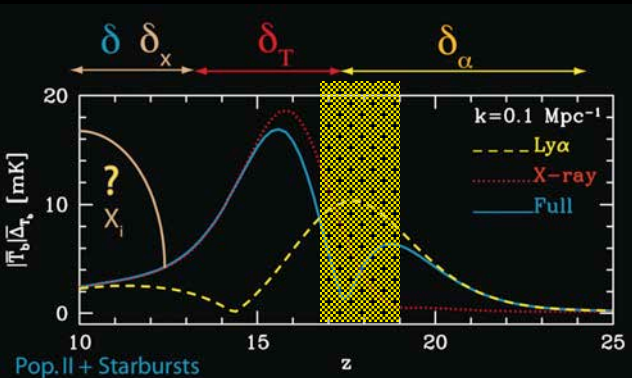
$z=19-20$



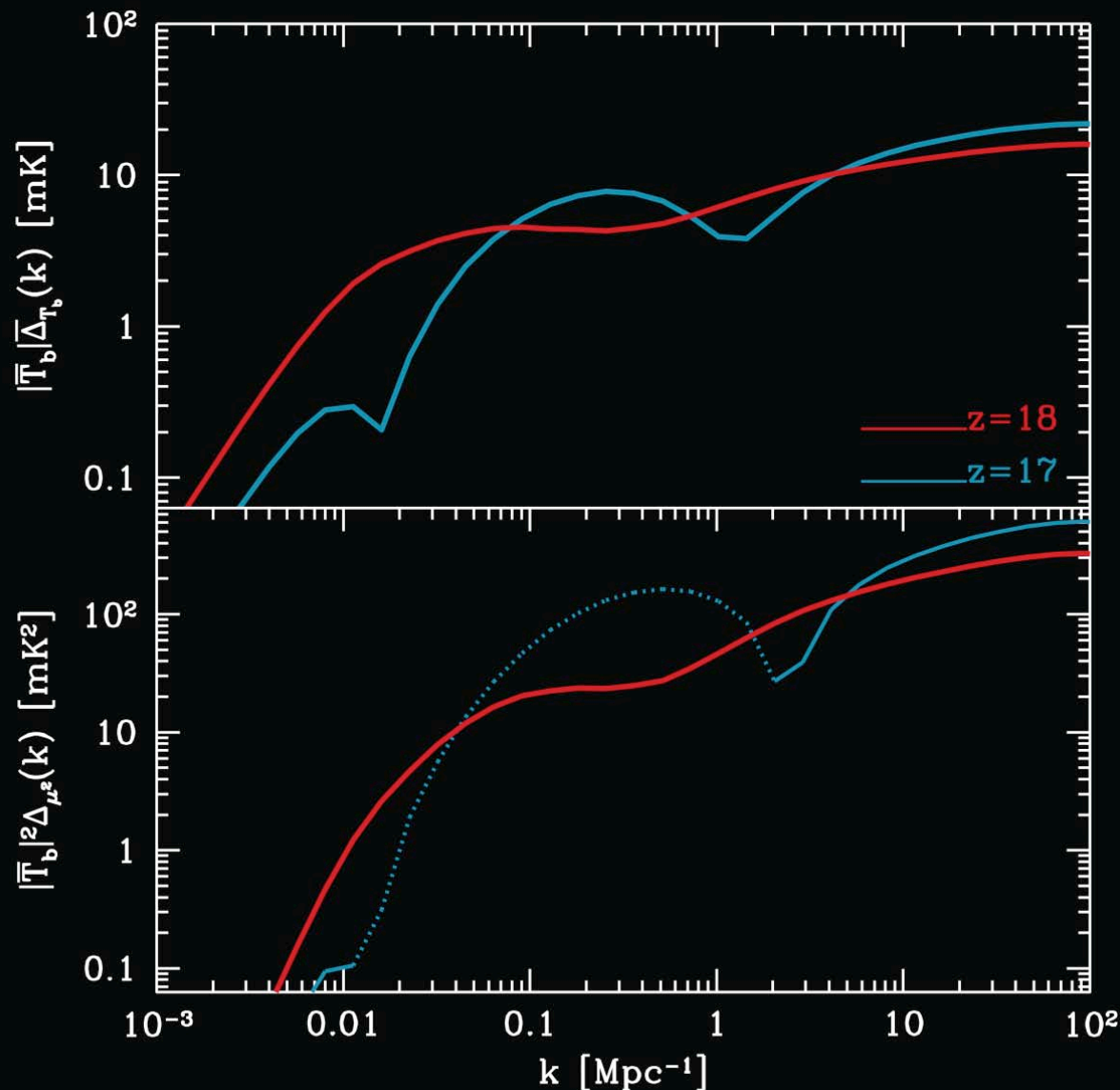
• Pure Ly α fluctuations



$z=17-18$

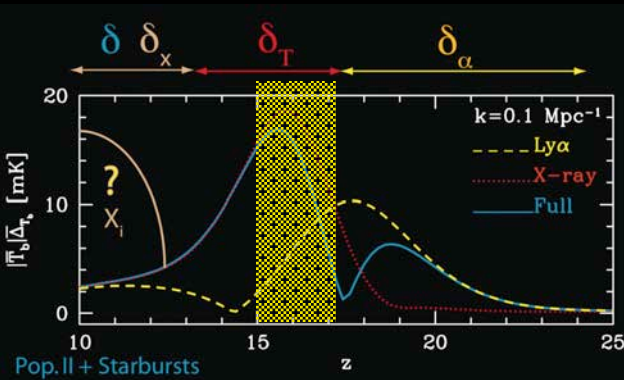


- Growing T fluctuations lead first to dip in Δ_{Tb} then to double peak structure
- Double peak requires T and $\text{Ly}\alpha$ fluctuations to have different scale dependence

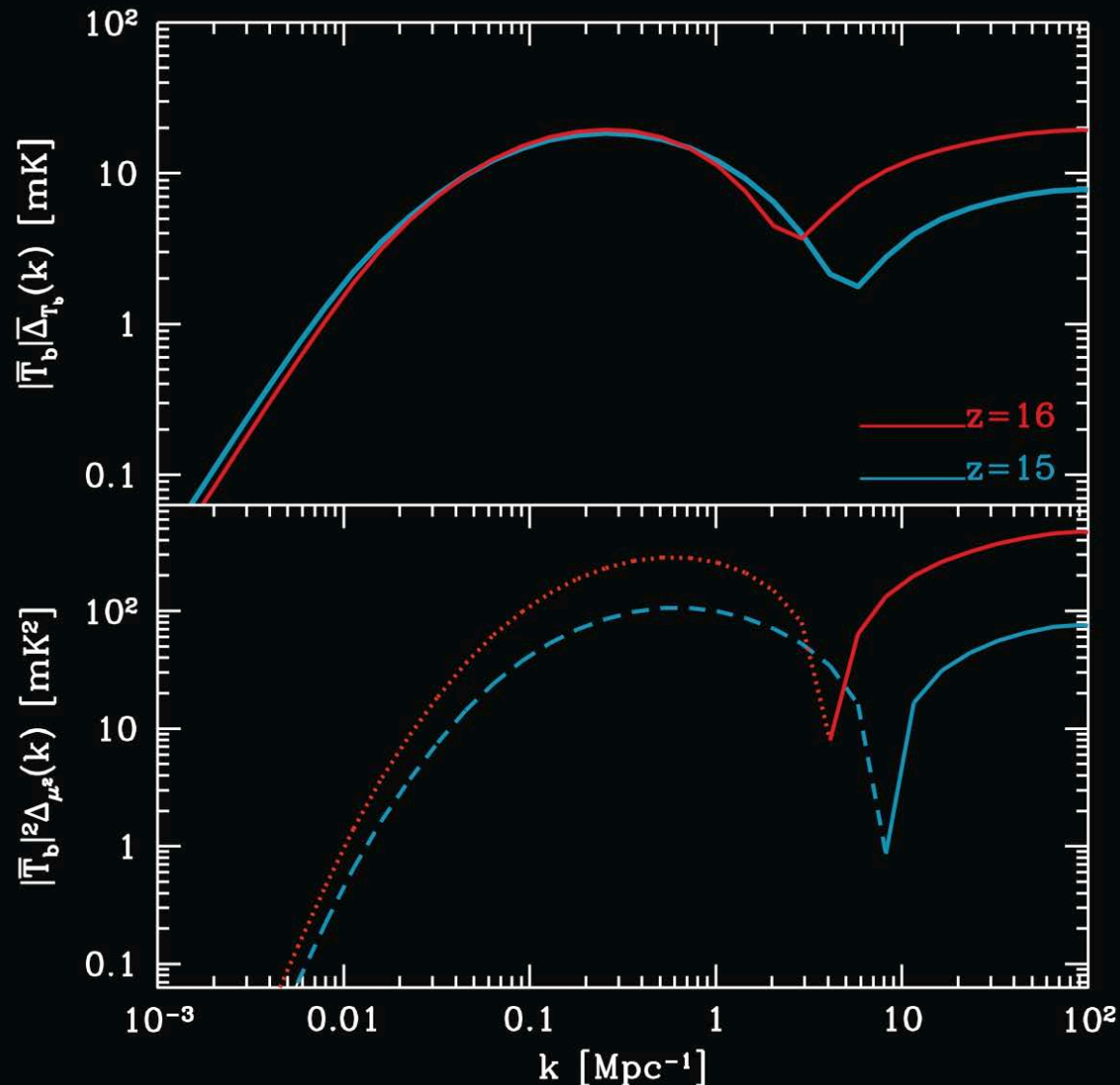


Redshift slices: T

$z=15-16$

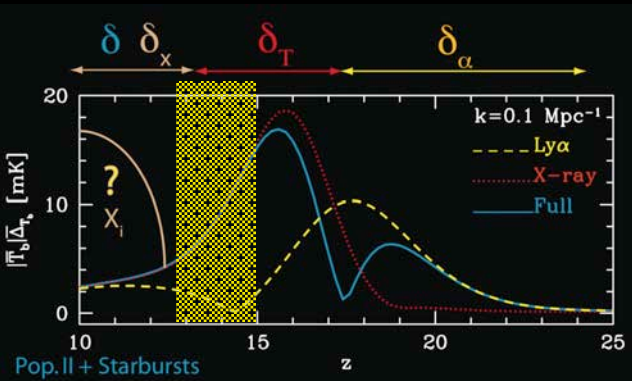


- T fluctuations dominate over $\text{Ly}\alpha$
- Clear peak-trough structure visible
- $\Delta_{\mu^2} < 0$ on large scales indicates $T_K < T_\gamma$

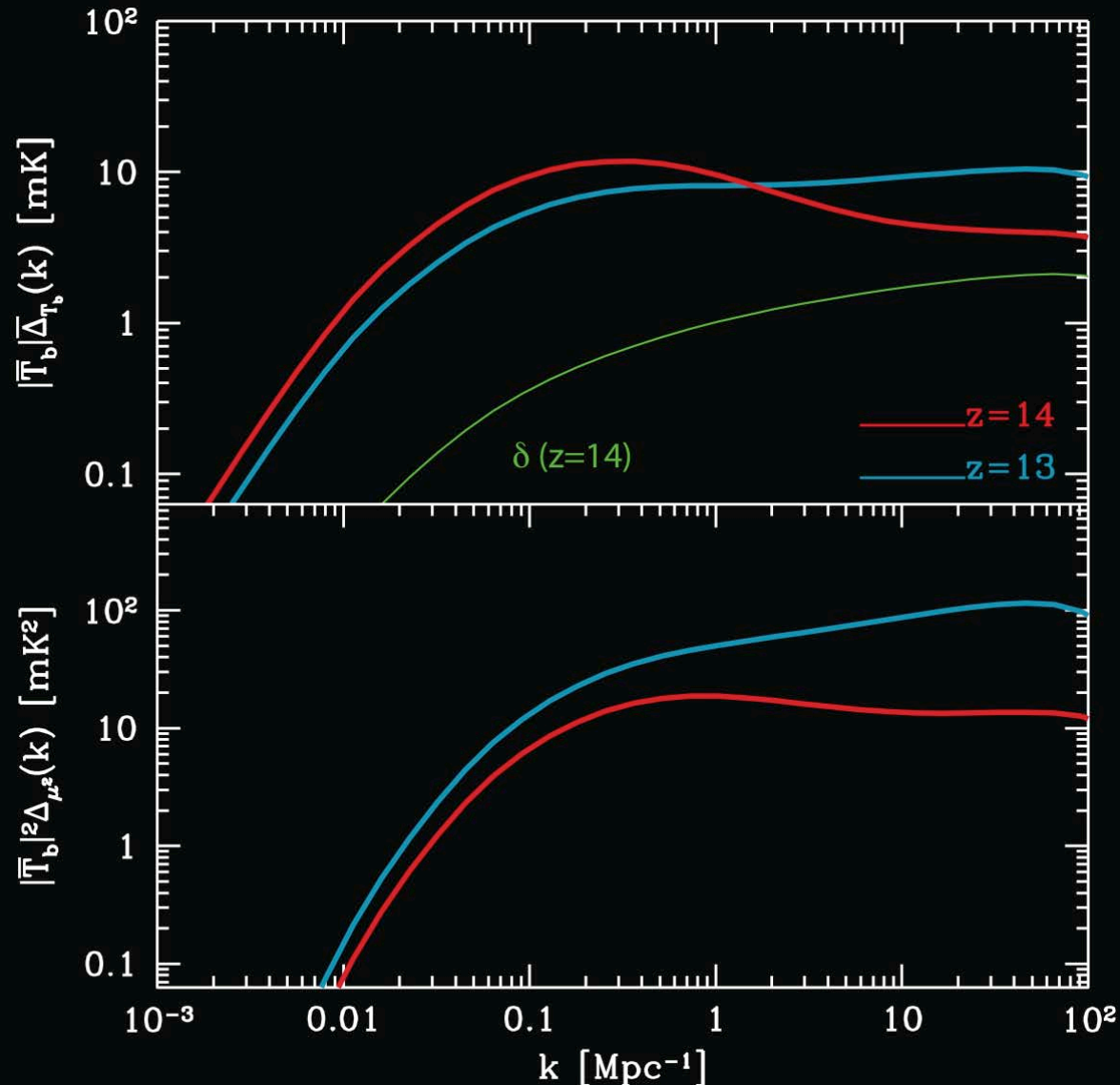


Redshift slices: T/δ

$z=13-14$

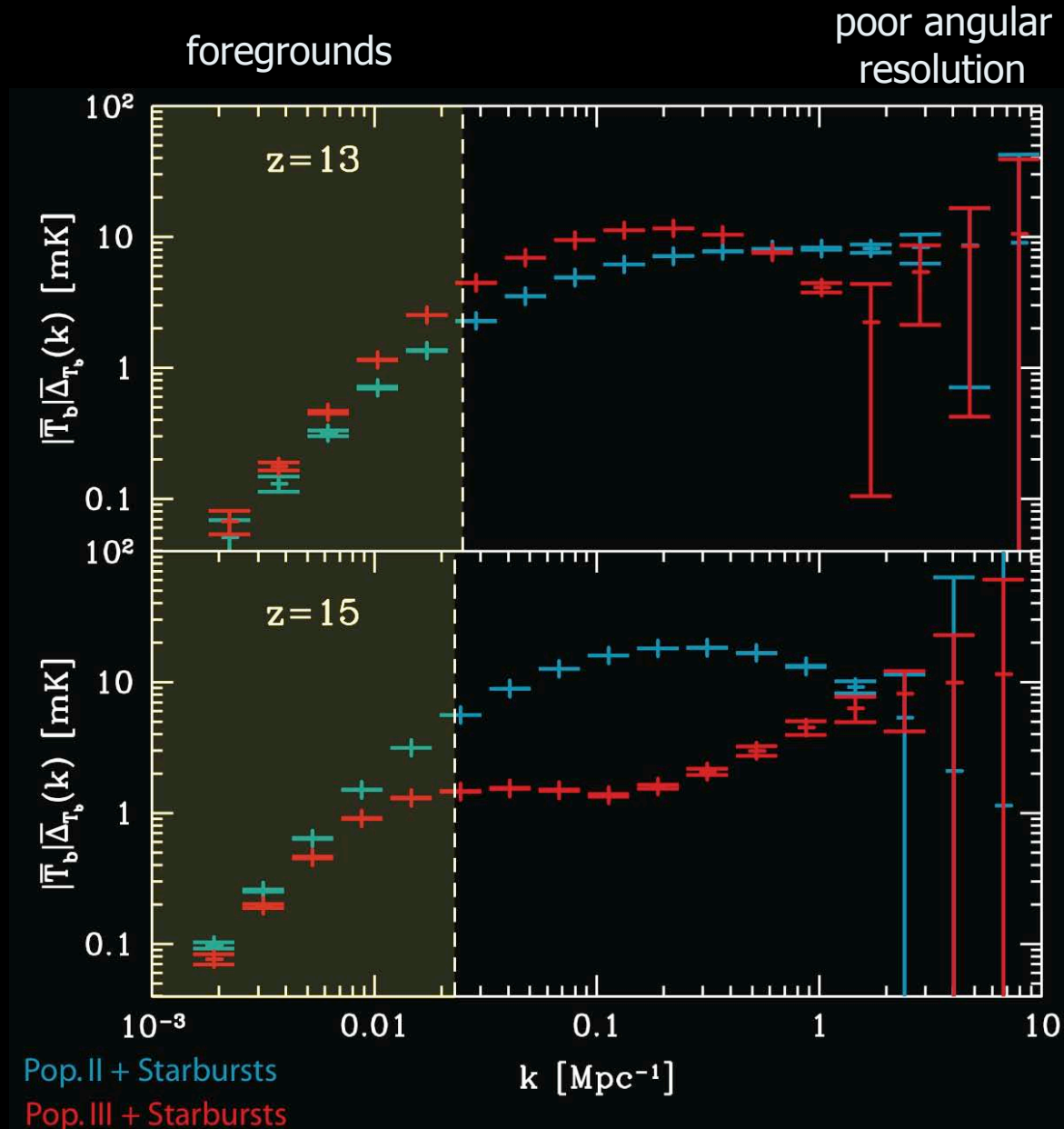


- After $T_K > T_\gamma$, the trough disappears
- As heating continues T fluctuations die out
- X_i fluctuations will start to become important at lower z



Observations

- Need SKA to probe these brightness fluctuations
- Observe scales $k=0.025\text{--}2\text{ Mpc}^{-1}$
- Easily distinguish two models
- Probably won't see trough



Conclusions

- 21 cm fluctuations potentially contain much information about the first sources
 - Bias
 - X-ray background
 - X-ray source spectrum
 - IGM temperature evolution
 - Star formation rate
- Challenge may be to separate out information
- Foregrounds pose a challenging problem at high z
- Will be important to include spin temperature fluctuations in future simulations
- Many model uncertainties - X-ray reionization?

For more details see [astro-ph/0607234](#) & [astro-ph/0508381](#)