

# Heating in the IGM and ICM

Peng Oh (UCSB)

# Heating in the IGM: hydrogen reionization

- When did it happen?
- How fast was reionization?
- What was the topology of reionization?
- What were the sources responsible?

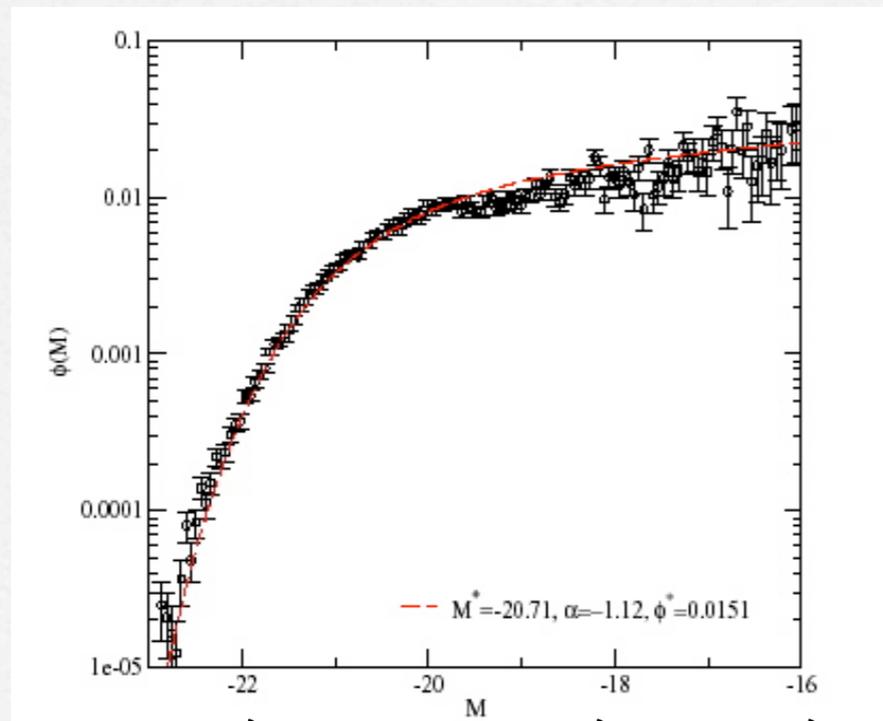
Field is driven primarily by observations

The argument from people power:  
It's the last time that baryons did anything  
together en masse...after that they all did their  
own thing.

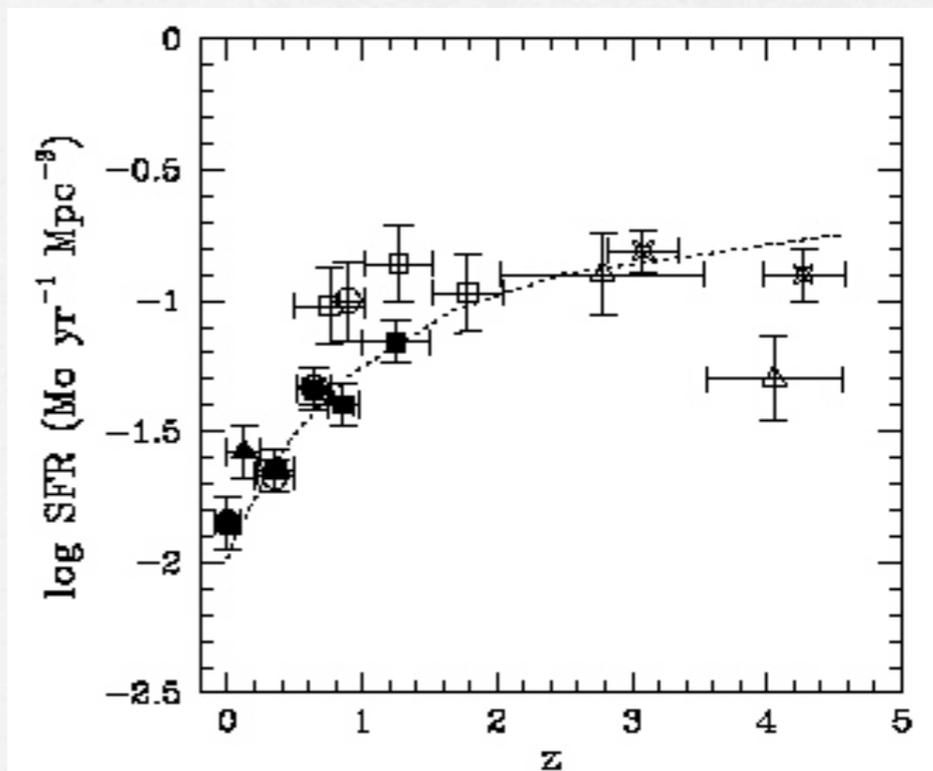


## Heating in the ICM

- How can we understand the characteristic mass scale  $L^*$  for galaxies?
- Why haven't most collapsed baryons turned to stars?



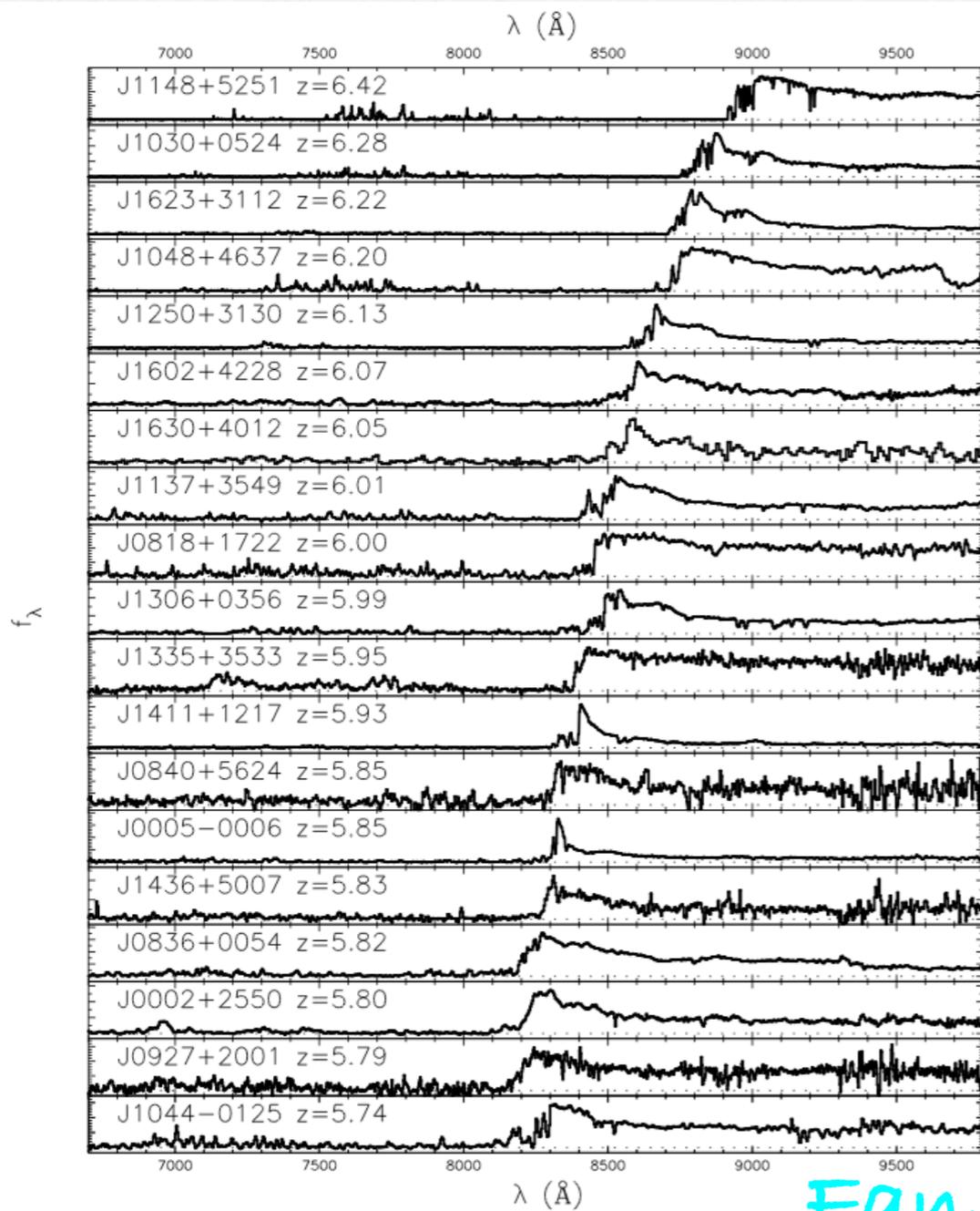
Knee in luminosity function



Downturn in comoving SFR

What do we already  
know about the  $z \sim 6$   
IGM?

# How neutral is the Universe at $z \sim 6$ ?

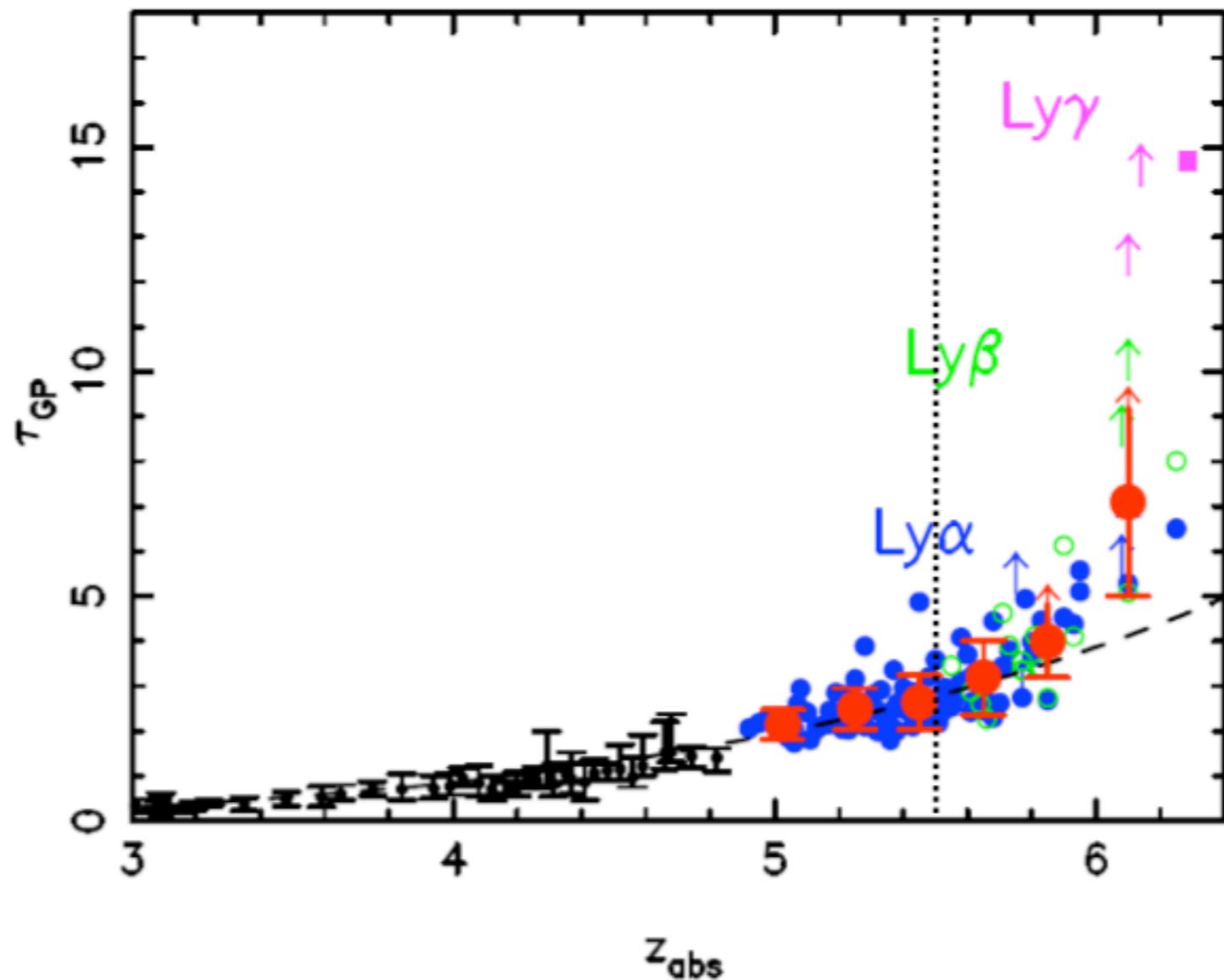


Best probe: Ly series  
absorption seen in 19  
SDSS QSOs with  
 $5.74 < z < 6.42$

saturated absorber:  
can only probe tail end  
of reionization

Fan et al 2006

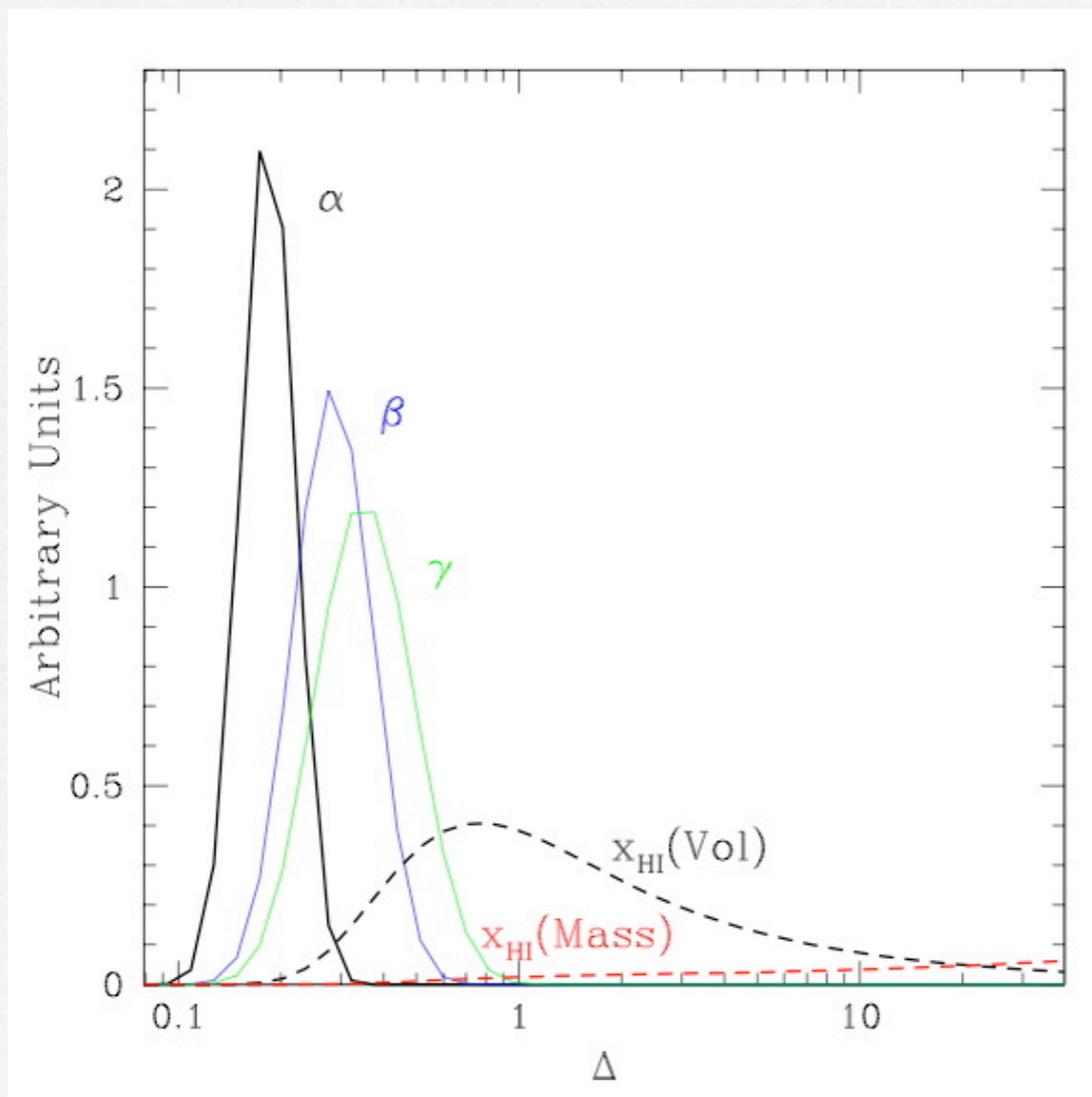
# End of reionization?



Claimed sharp increase  
in optical depth and  
scatter

Fan et al 06

...but it's really hard to infer neutral fraction...



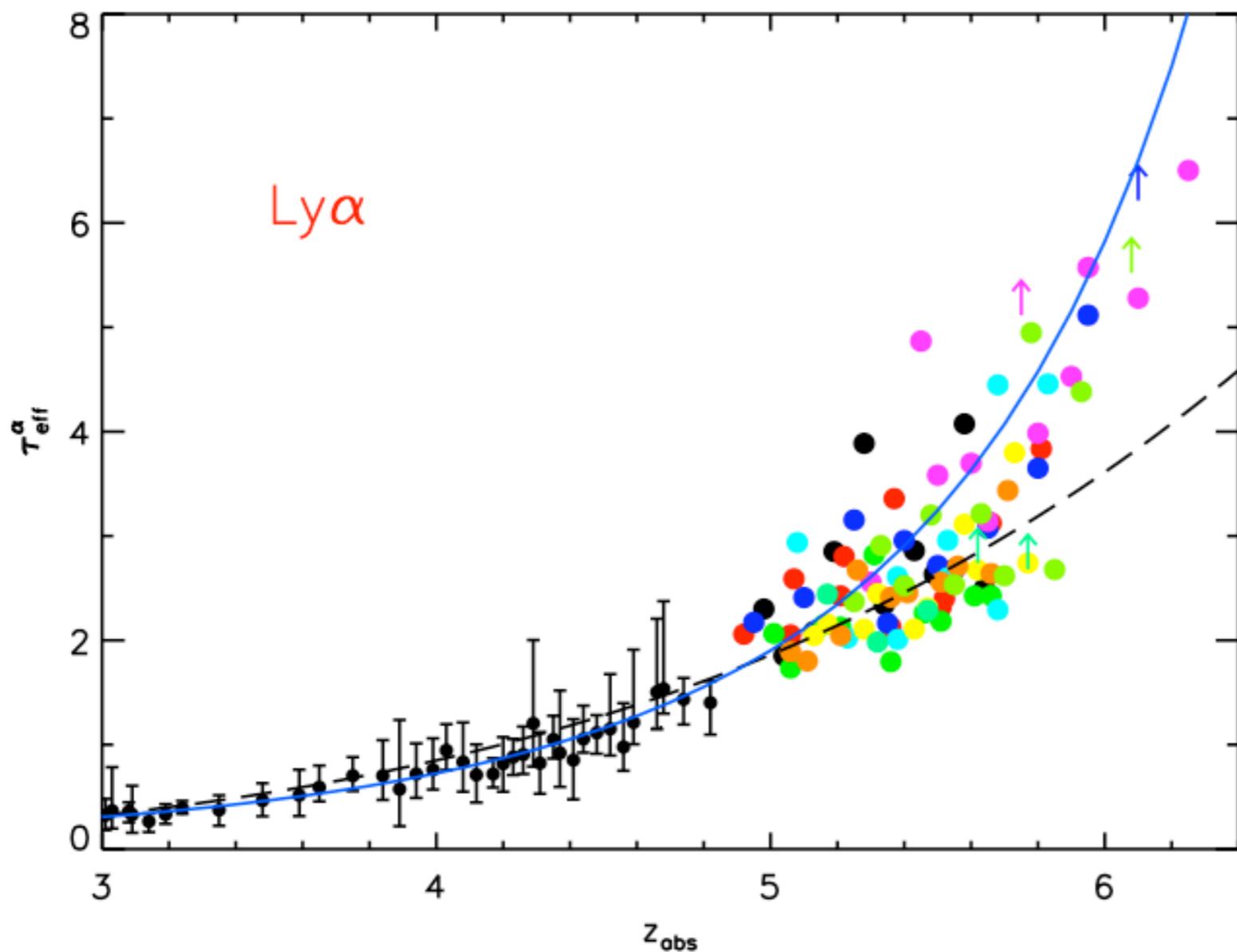
Oh & Furlanetto 2005

Transmission mainly  
due to rare voids

Most HI at higher  
overdensities

Caution: comparing  
different Lyman series  
on absolute scale is hard

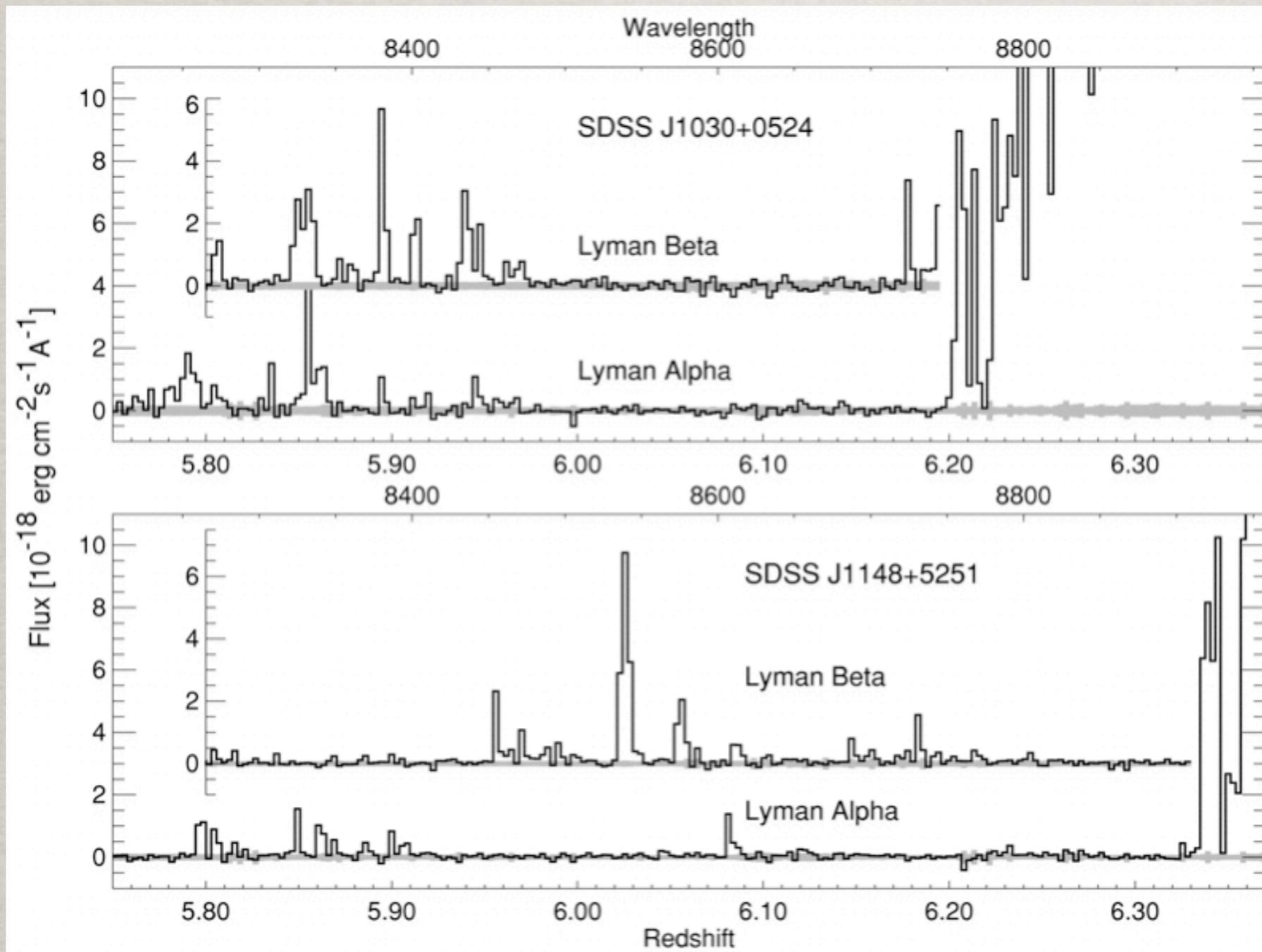
...and highly dependent on assumed density PDF



Evolution can be explained if assume lognormal PDF

Becker, Rauch & Sargent 2006

# HAVE WE SEEN PATCHY REIONIZATION?



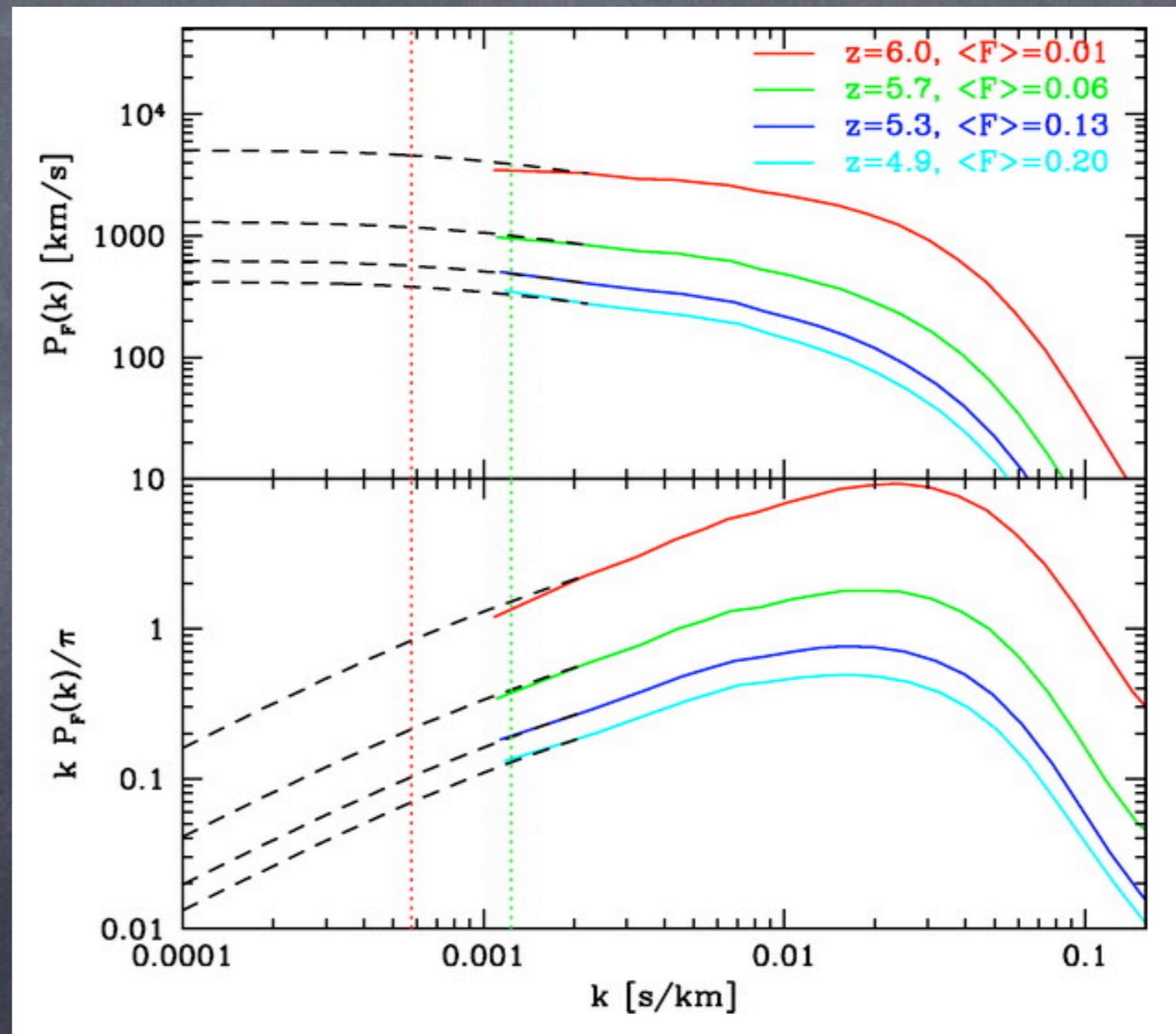
There's a lot of sight-line to sight-line scatter in QSO absorption spectra. Order unity fluctuations on LARGE (50-100 Mpc comoving) scales

# But it is VERY hard to see patchy reionization...

Lidz, Oh & Furlanetto 2006

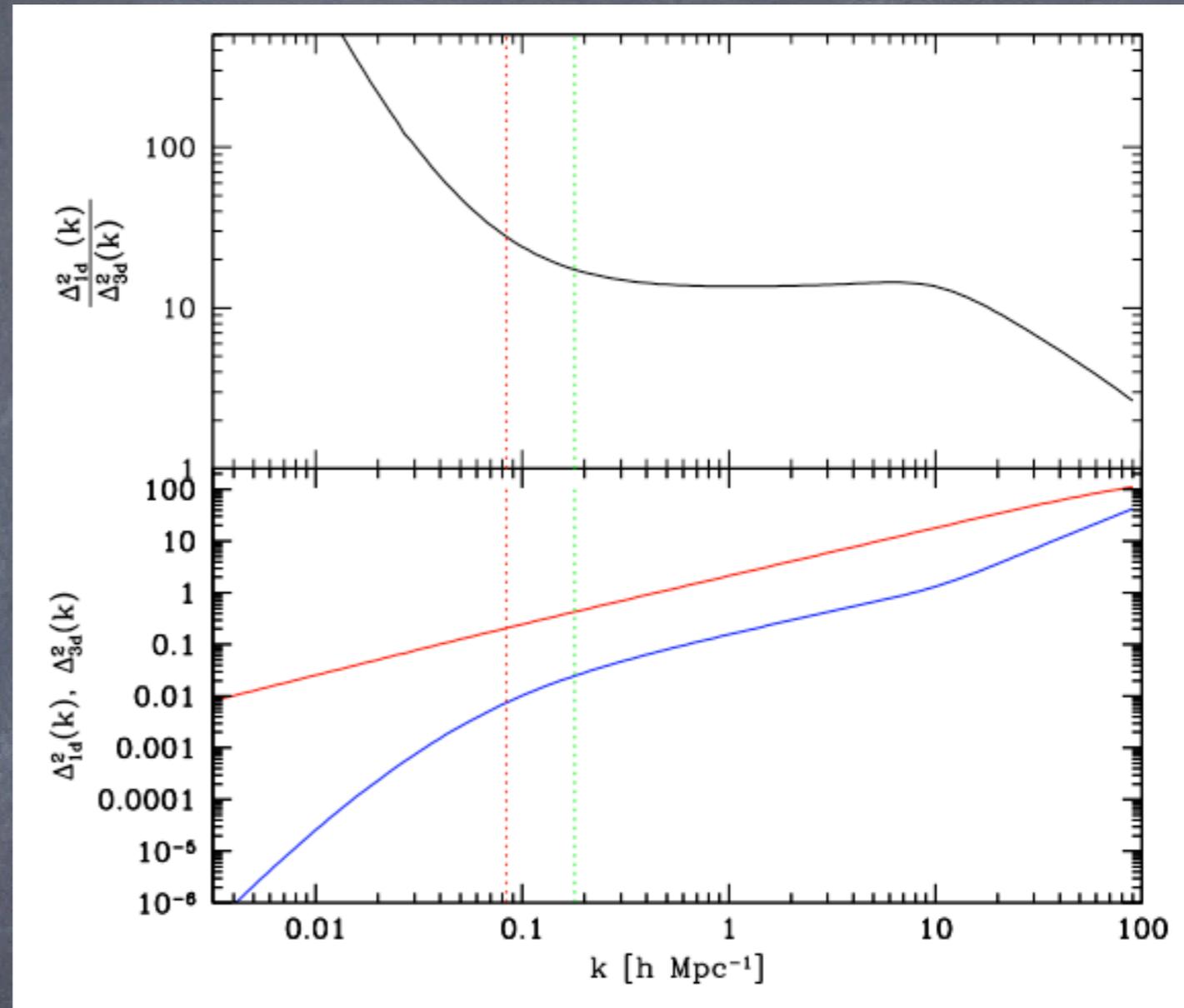
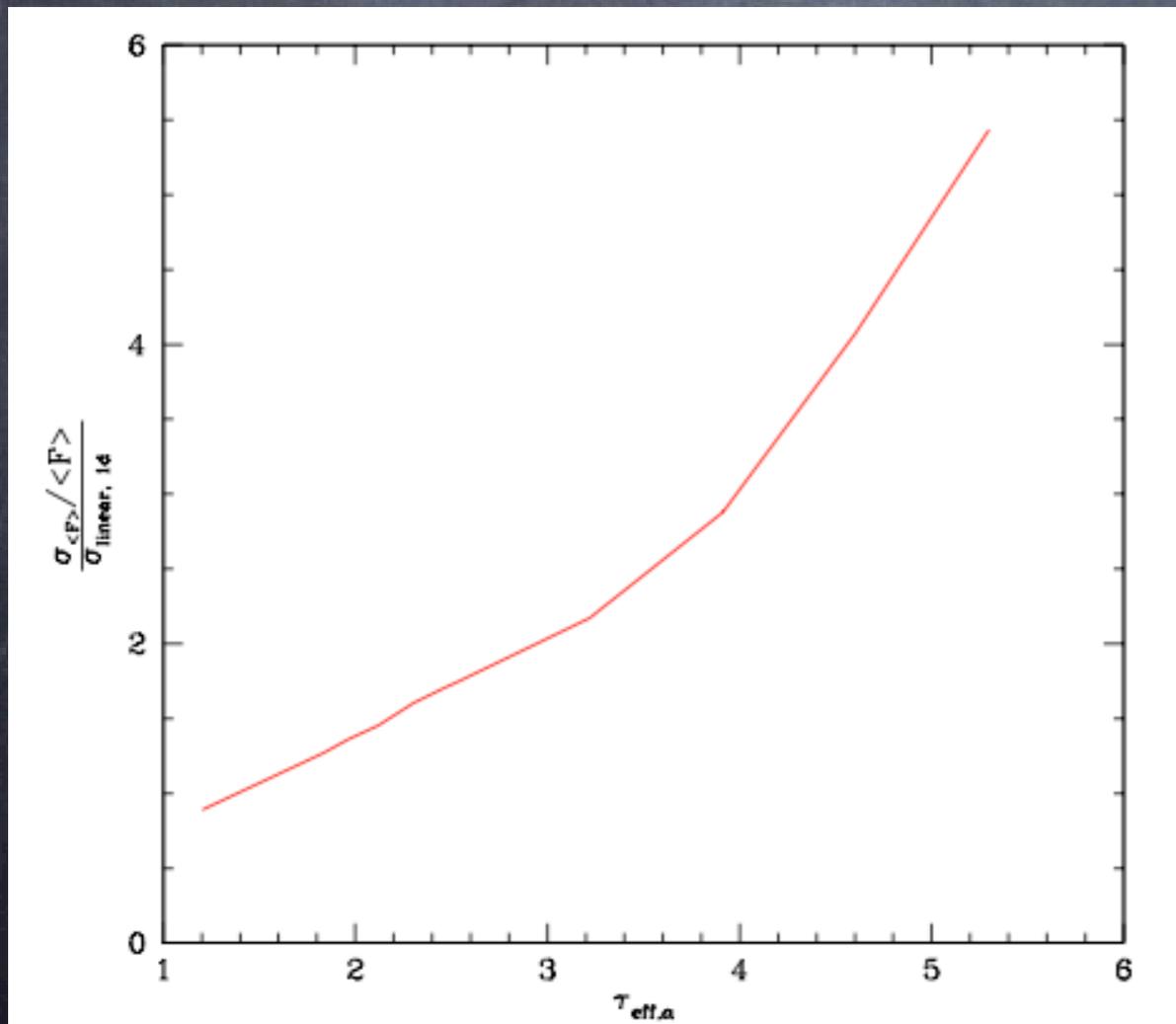
Simulate 40 Mpc  
 $h^{-1}$  box

Flux power  
spectrum declines  
slowly with scale



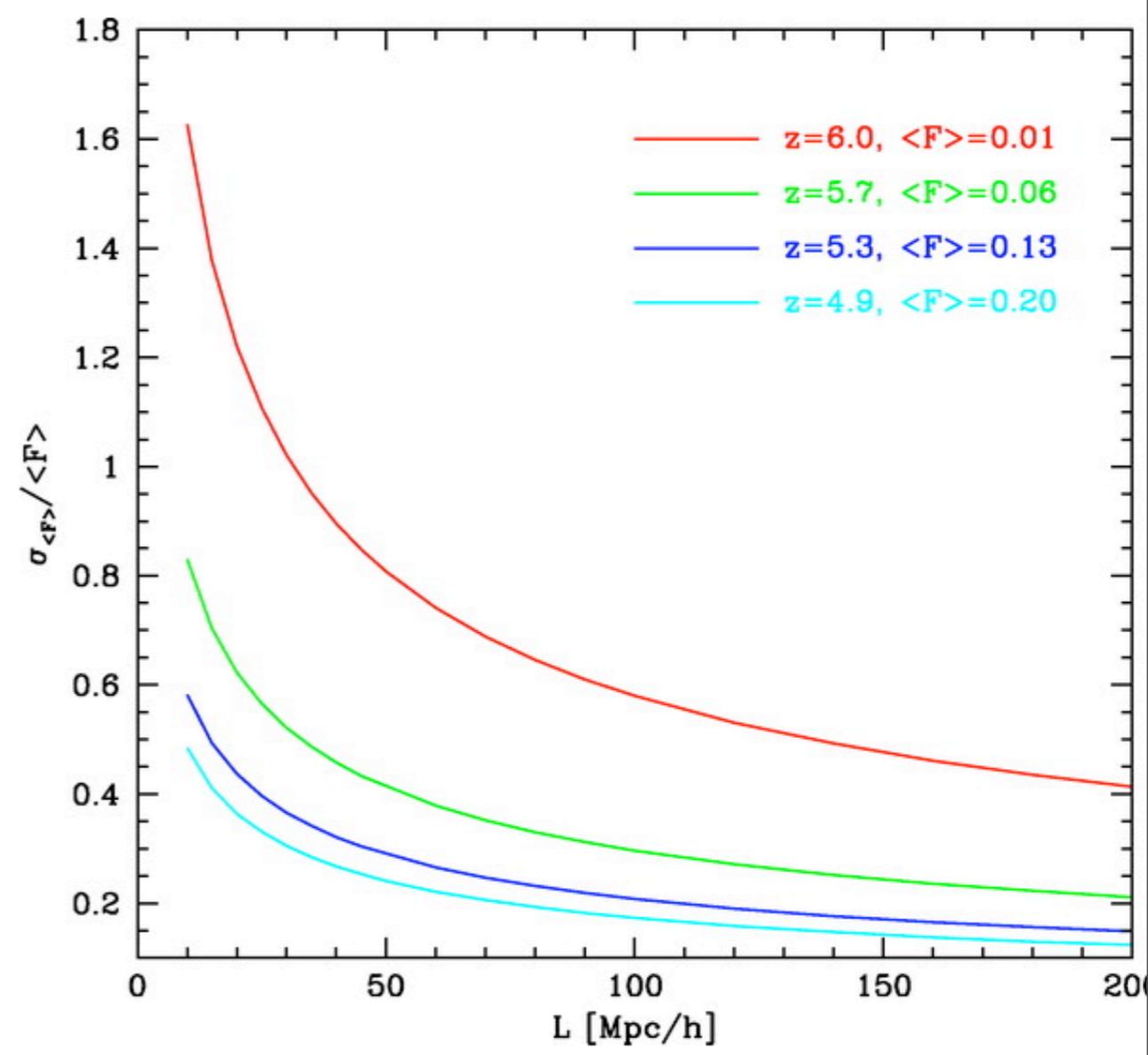
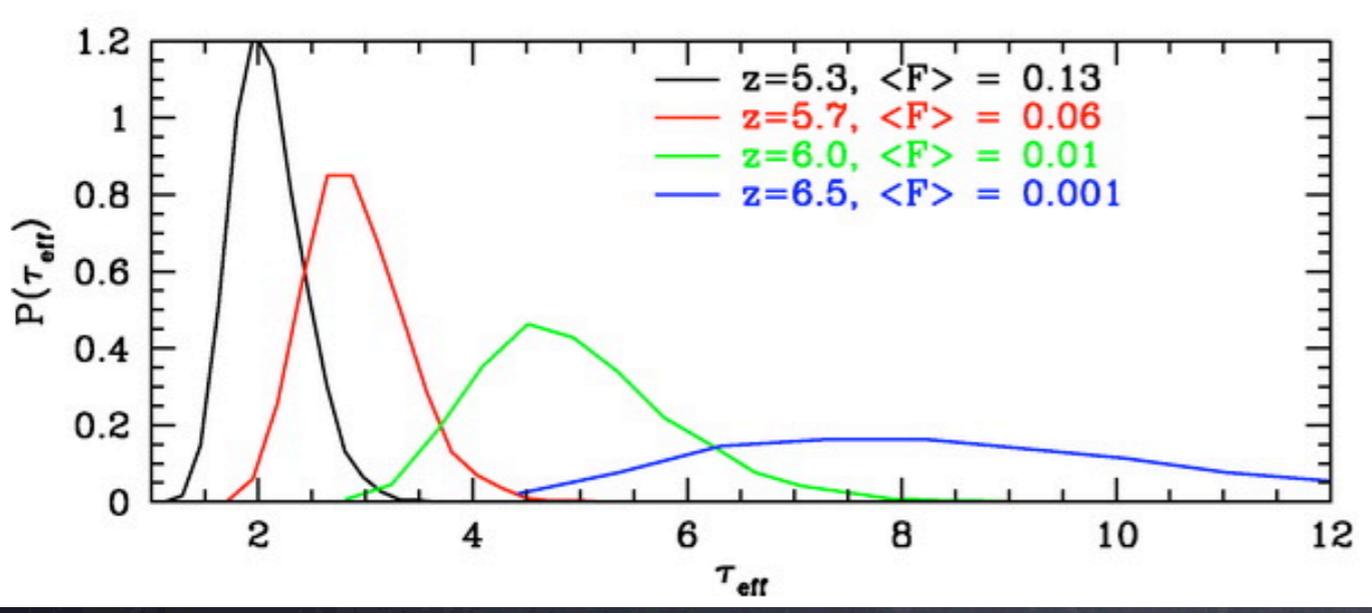
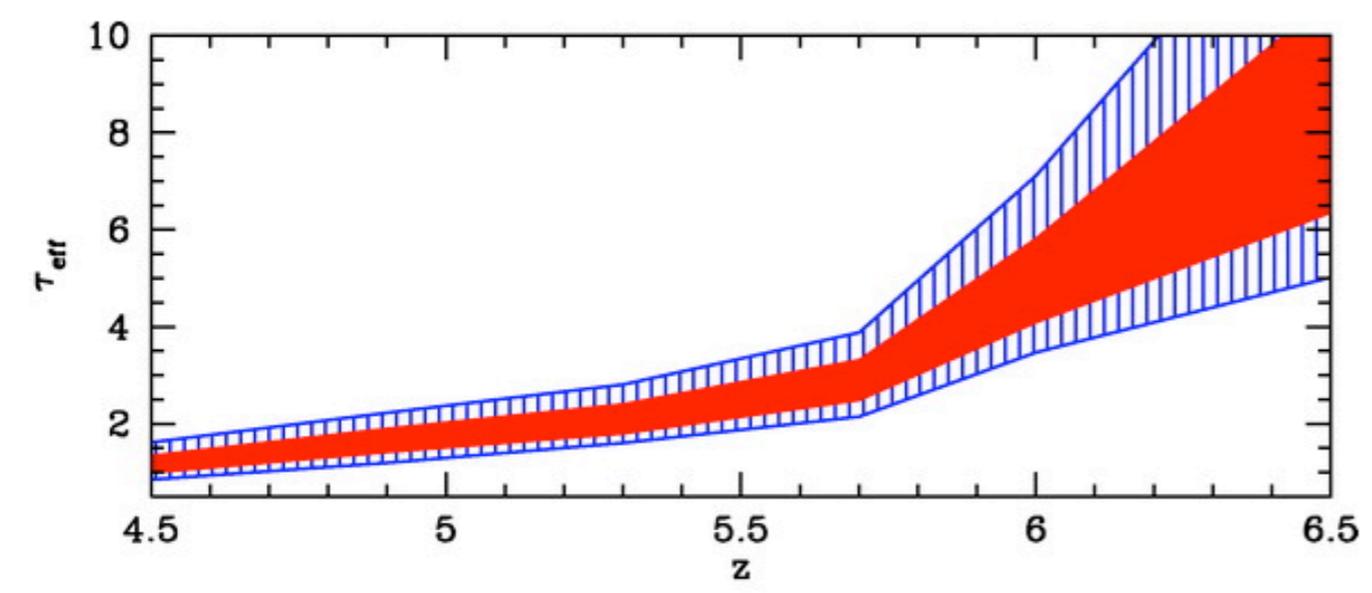
# Two effects

Aliasing boosts power on large scales



Bias increases amplitude of fluctuations

# Uniform Radiation Field is consistent w/ observed scatter...



# The future: 21 cm observations

"There are more things in heaven and earth,  
Horatio, Than are dreamt of in your philosophy."

--Hamlet (Act I, Scene V)

# Advertising: 21 cm Review Article

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Cosmology at low frequencies: The 21 cm transition and the high-redshift Universe

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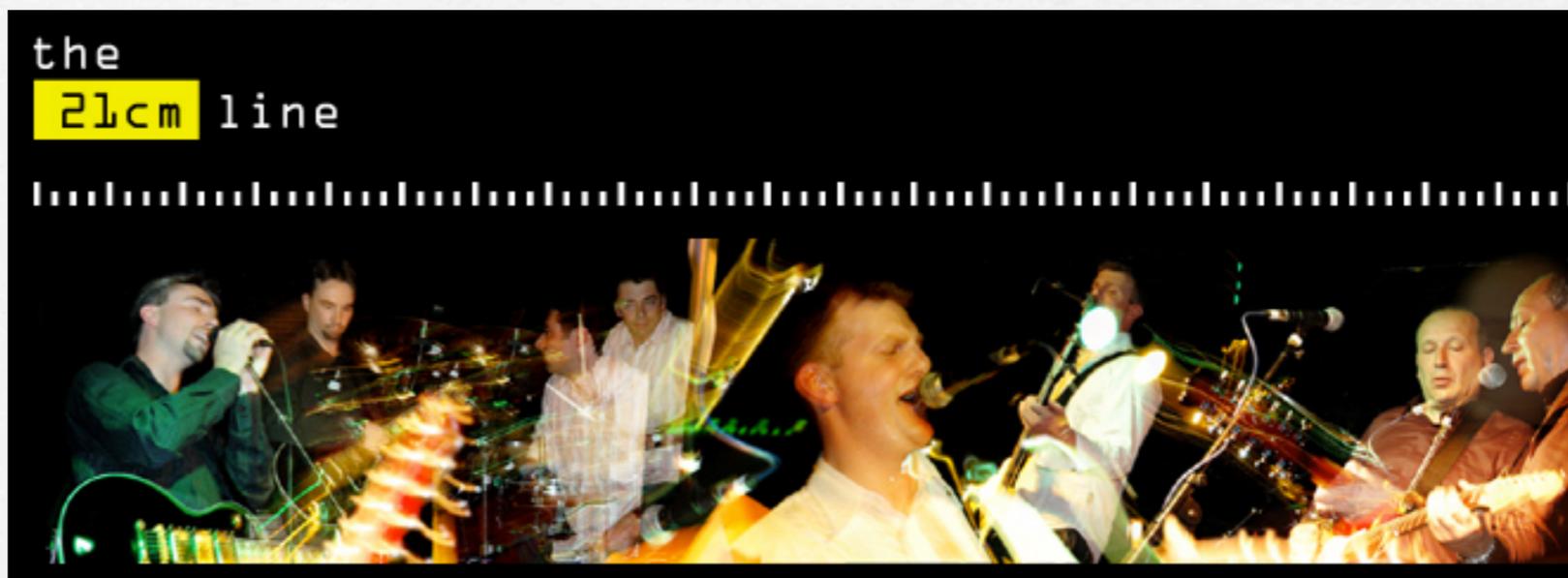
Accepted 1 August 2006

Furlanetto, Oh &  
Briggs 2006  
(FOBOG)

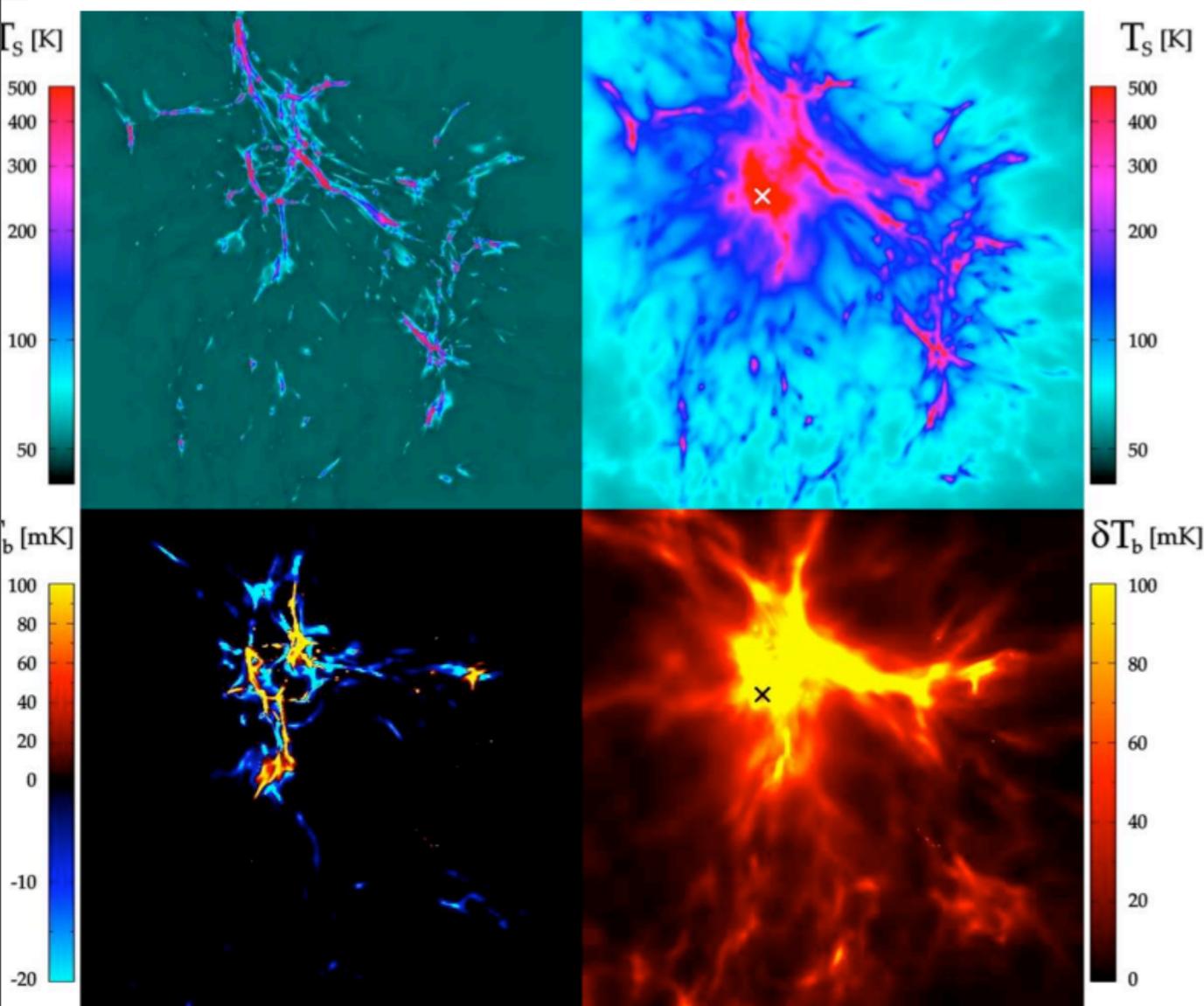
[astro-ph/0608032](https://arxiv.org/abs/astro-ph/0608032)

Top Google hit for '21cm  
transition'

But if you Google  
'21cm line' instead  
you get...



# 21 cm observations will revolutionize the field



Kuhlen & Madau 06

See 21cm emission from IGM in absorption or emission against CMB

Couple spin and kinetic temperatures by collisions or Wouthuysen-Field effect

Probe both Dark Ages and First Light

LOFAR



PAPER--USA/Australia

GMRT---India

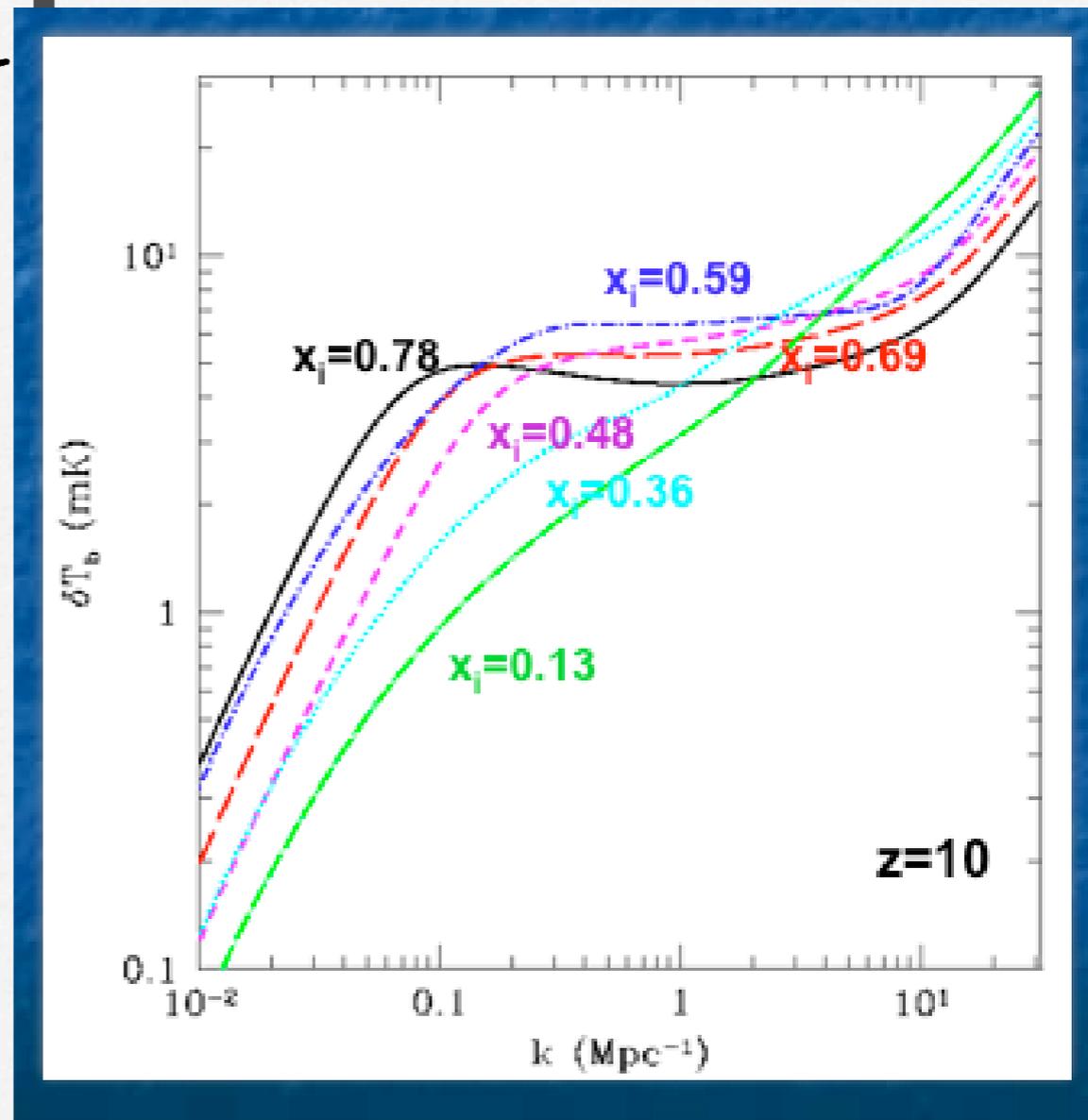
MWA---Western Australia

LOFAR--Netherlands

SKA--??

# 21cm Power Spectrum

- Language generally used for 21cm fluctuations
- Tools developed for CMB/ galaxy surveys
- Natural language for interferometer
- Good choice for Dark Ages, before ionizing sources turn on. But after that...

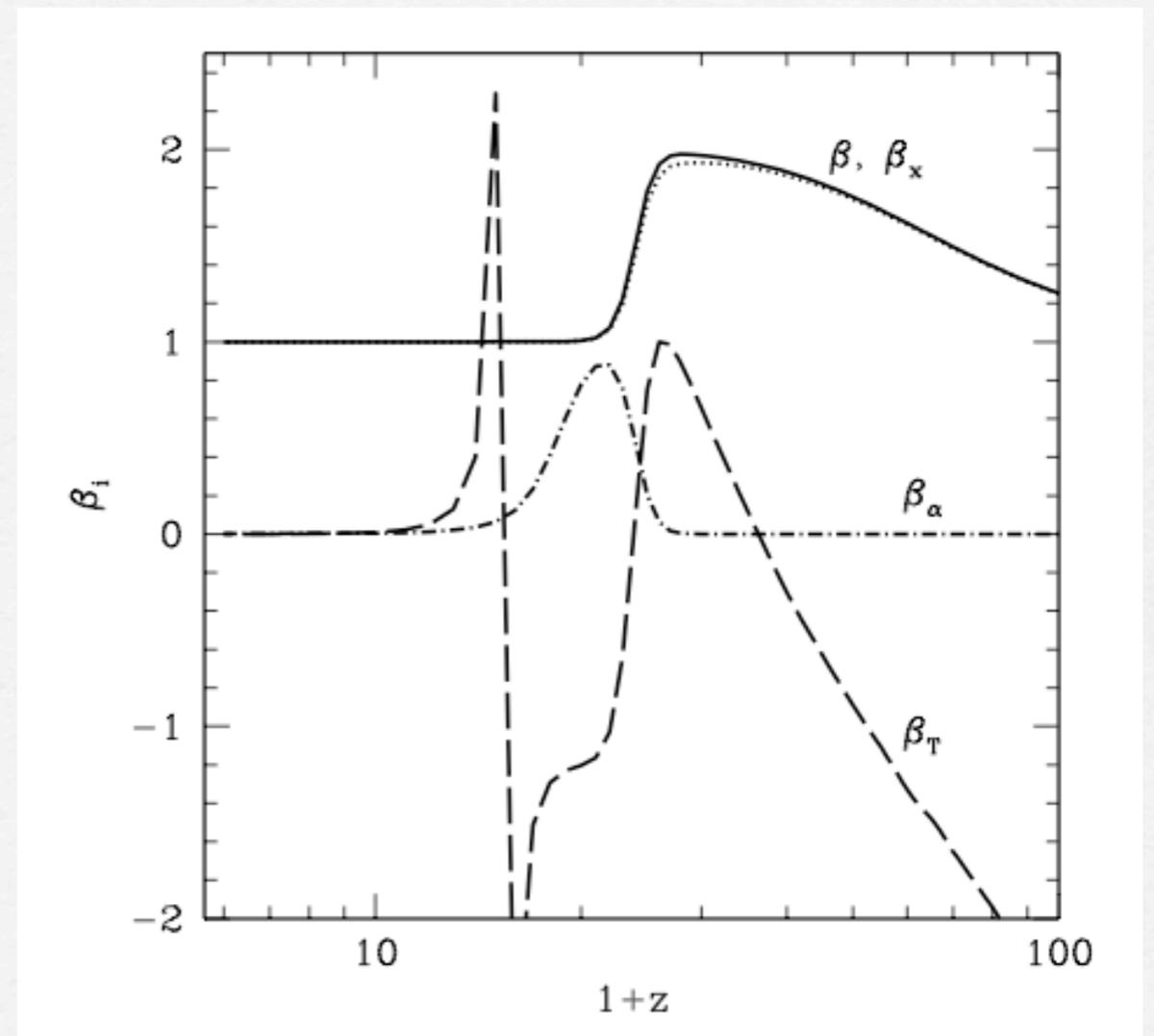


FOBOG

# ...many effects contribute

Fluctuations in...

- density (Gaussian)
- Ly-alpha flux
- ionization state
- temperature
- velocity gradients

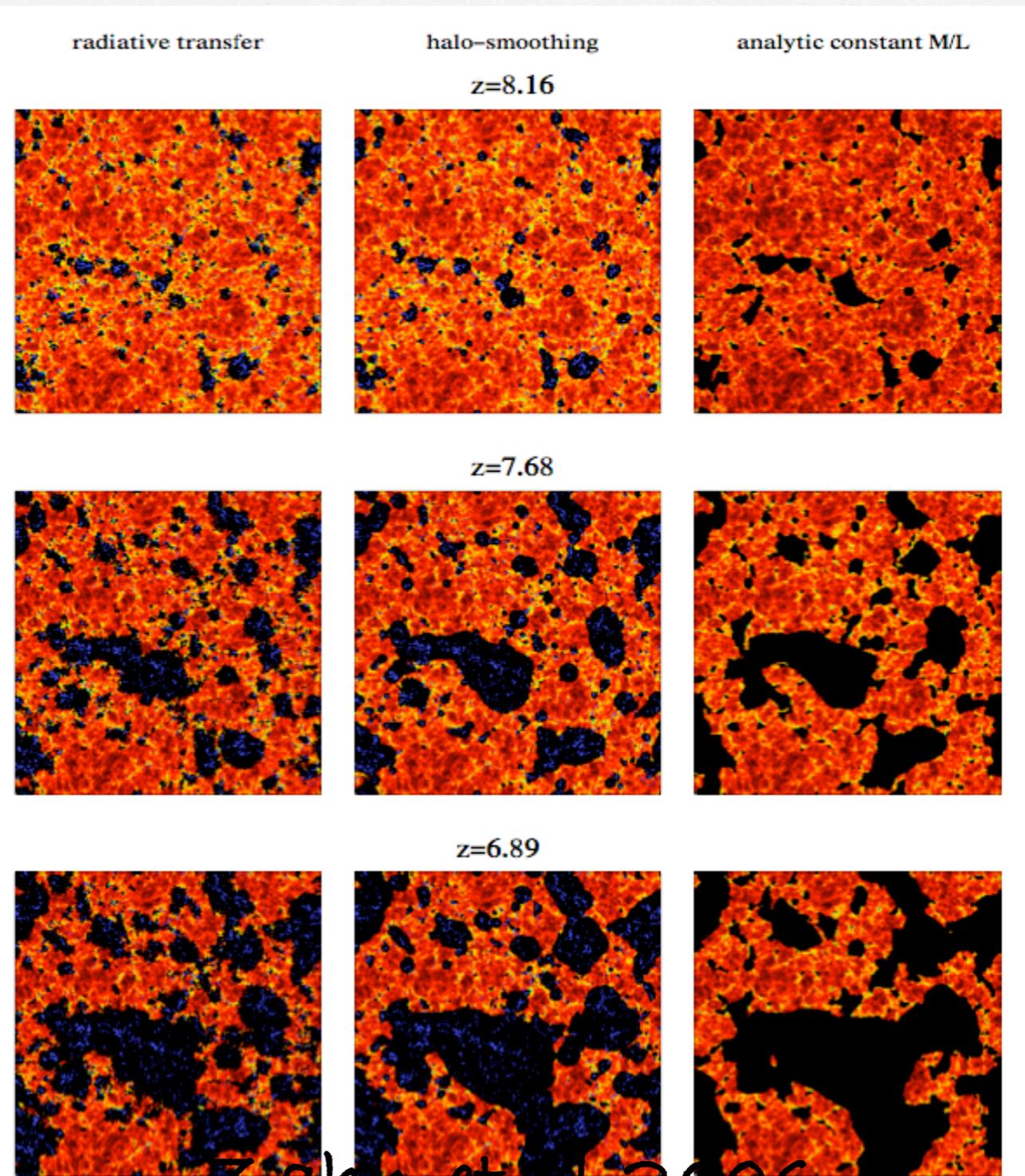


FOBOG

Many likely to be correlated

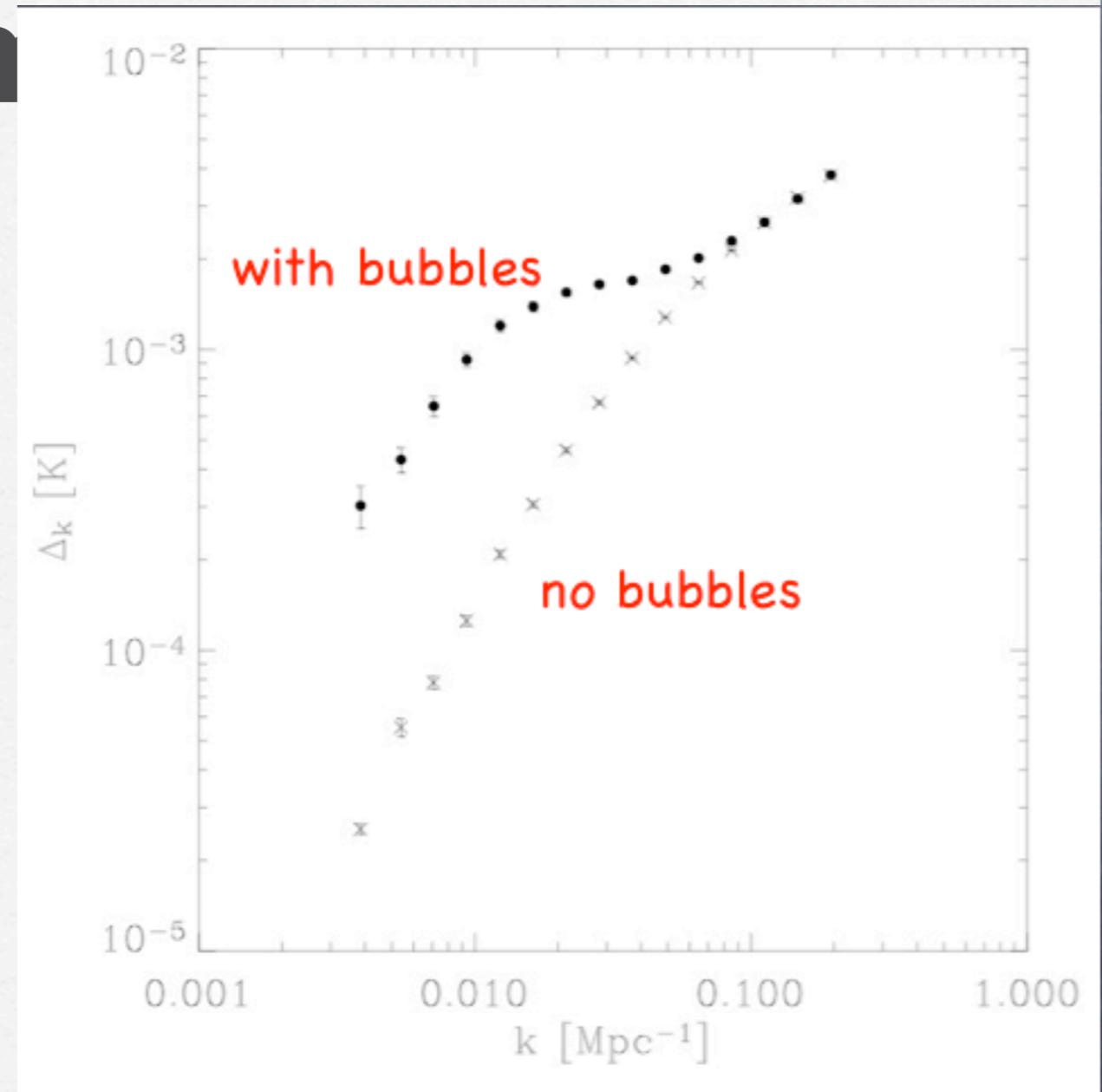
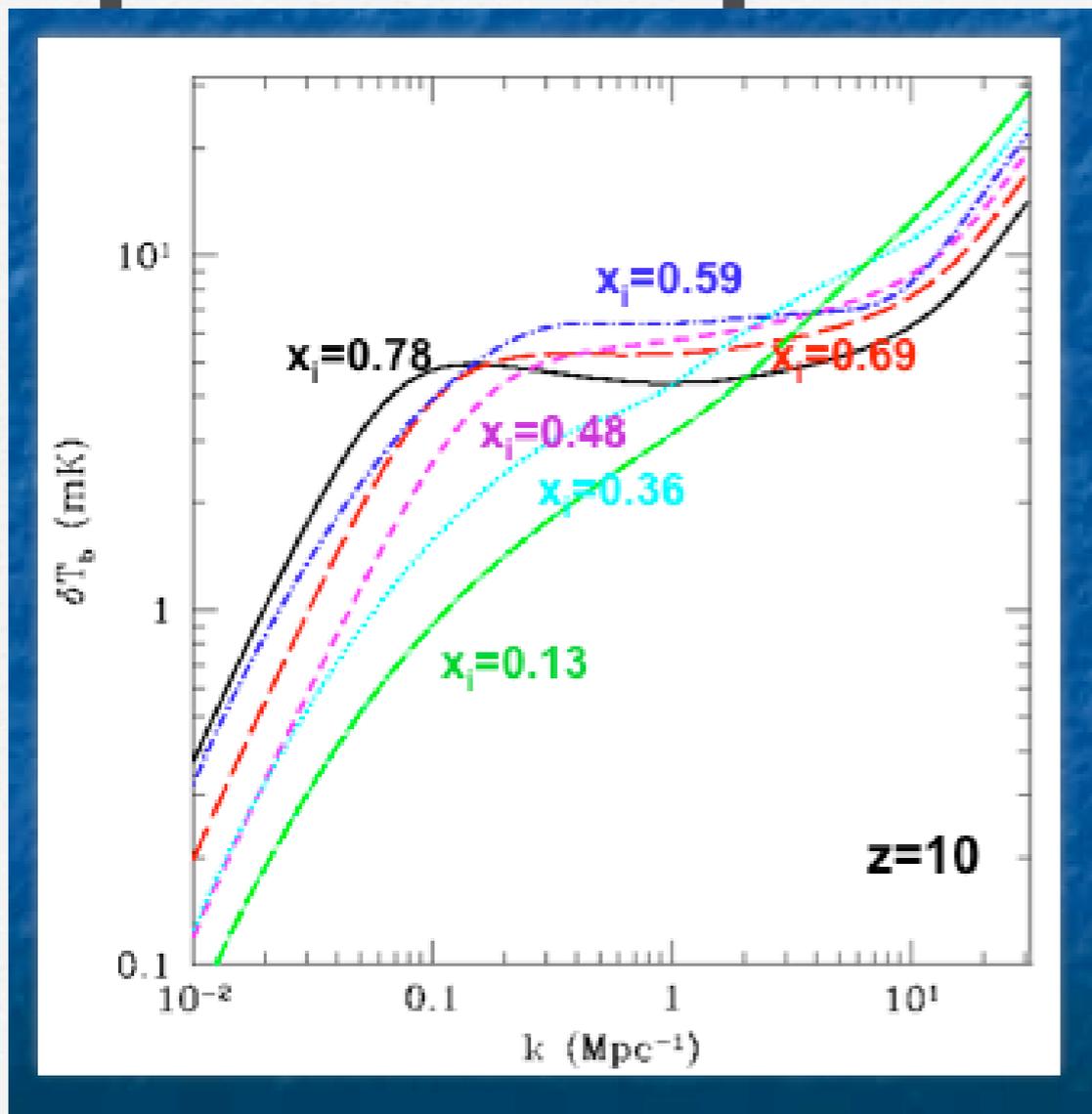
...it's a highly non-Gaussian field!

If we want to study growth and topology of reionization, we should focus on the bubbles



Zahn et al 2006

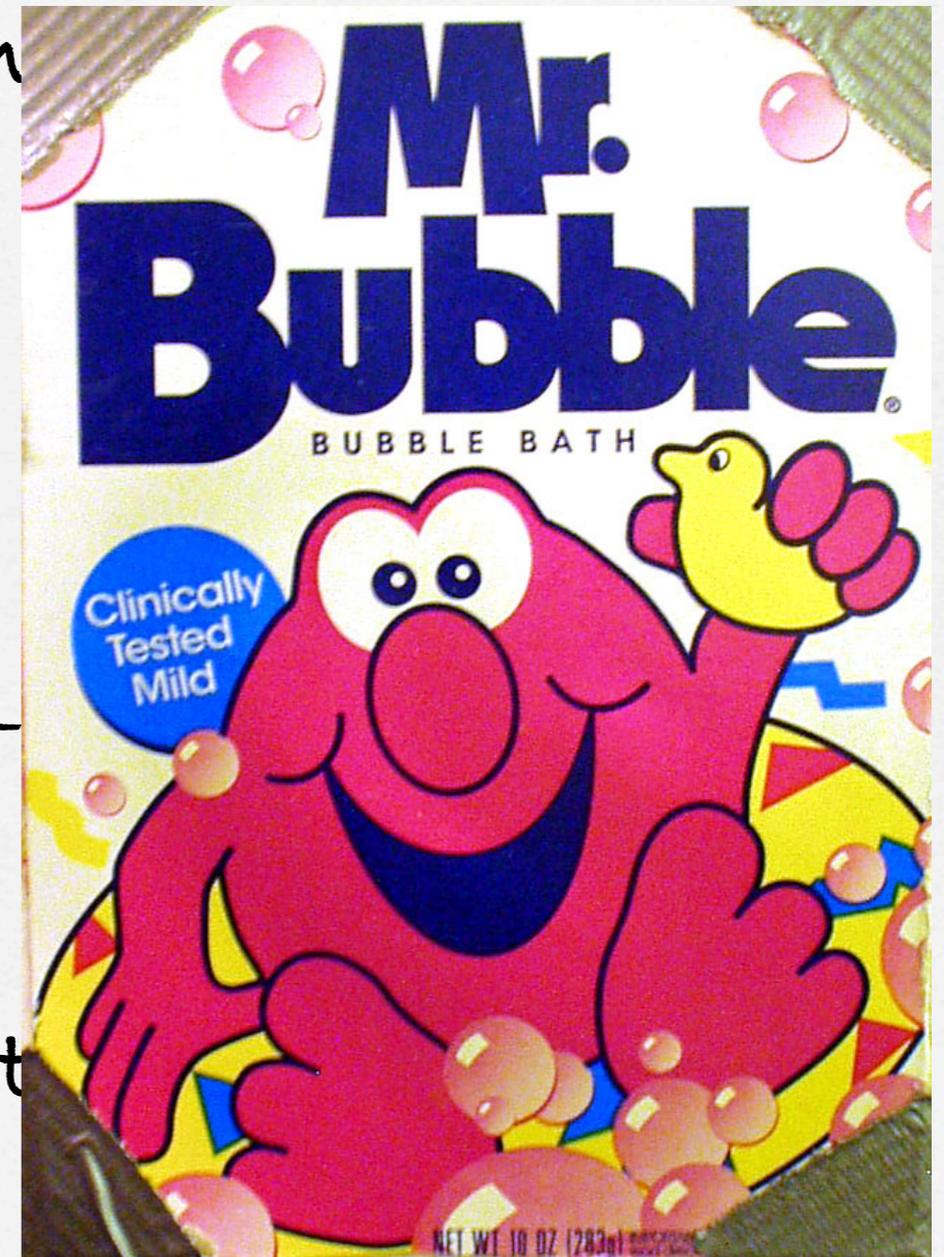
...bubbles DO strongly affect power spectrum



...but quantifying this will be model-dependent

# Bubbles are your Friend

- Probe of ionizing source population (supposed to be big)
- Directly extract HII filling factor
- Foreground calibrator:
  - Measure mean temperature  $T(z)$
  - Remove long wavelength artifact from foreground removal



# Direct Imaging

S/N high only on largest scales, need  $R \sim 20$  Mpc

Rare bright quasars (or clustered galaxies)

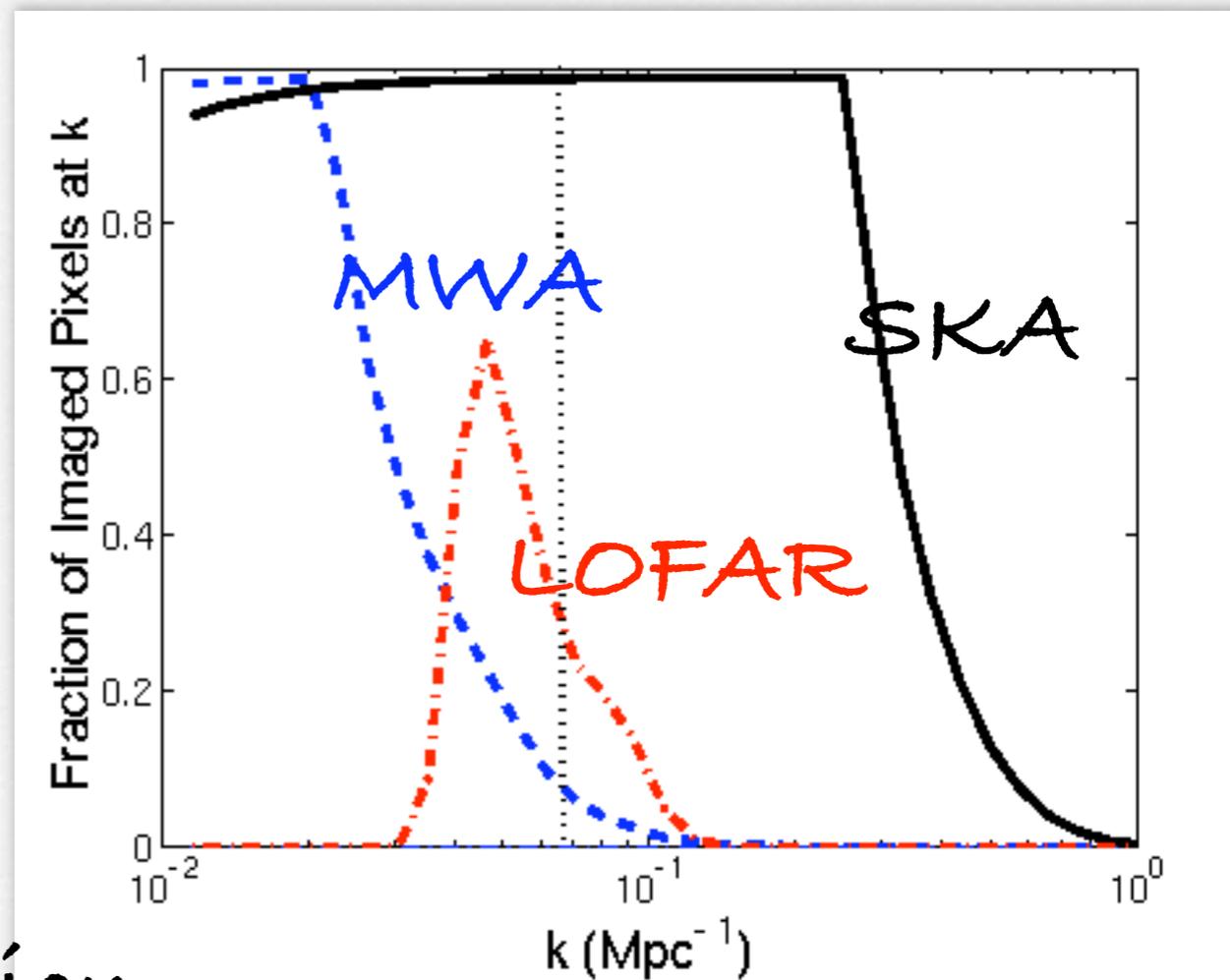
BUT: survey volume is HUGE!

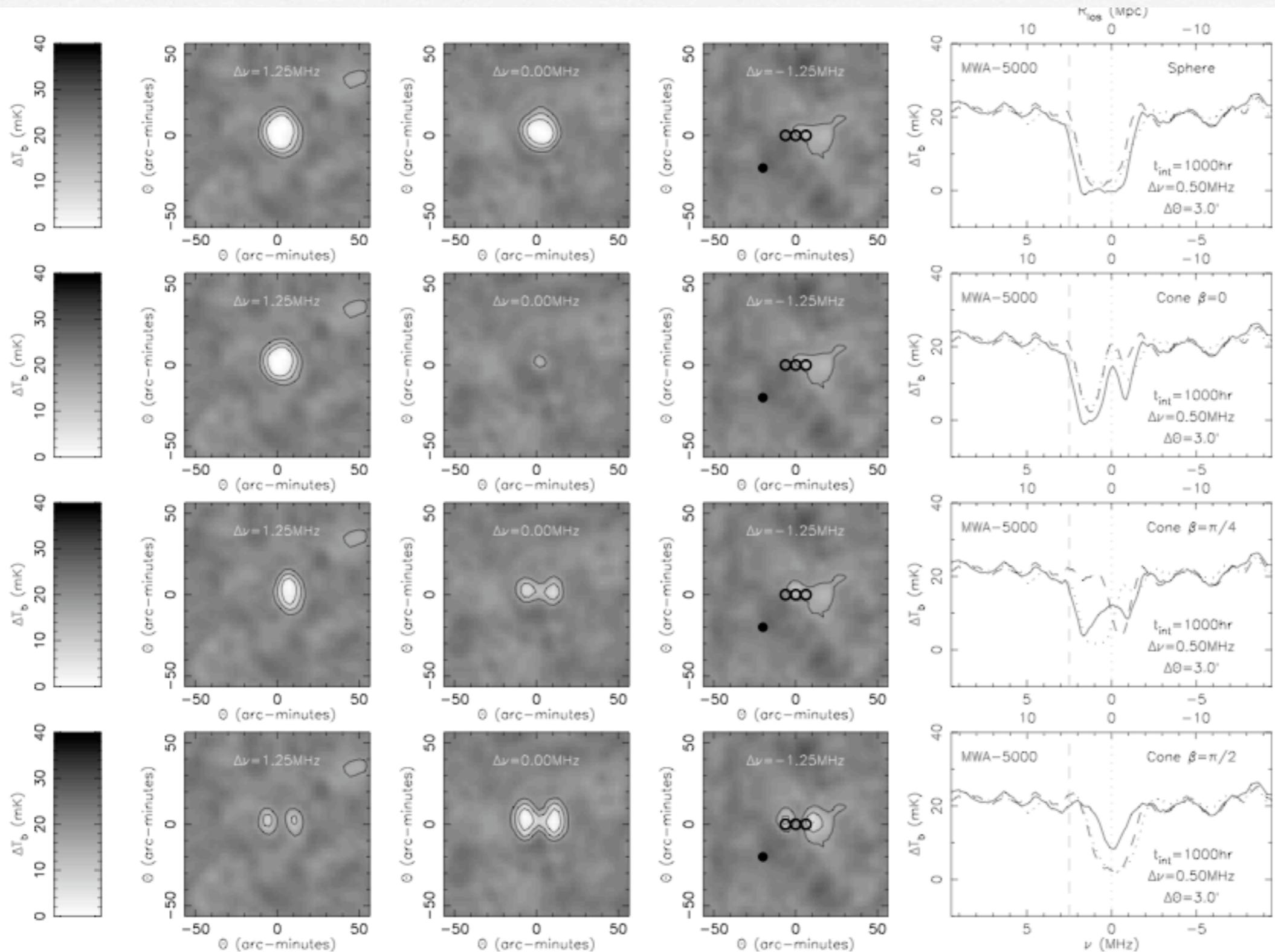
Expect 1 active/fossil HII region

in every MWA FOV with  $R > 20$  Mpc (McQuinn et al 2006)

(24, 40) Mpc at  $z=7$  (Wyithe,

Loeb & Barnes 2004)





Wyithe, Loeb & Barnes 2004

# ...what do we get?

- $\delta T_b(z)$  x-rays, fossil HII
- --
- Foreground calibrator
- Size, shape of HII region --> QSO properties
- Discover QSOs? (though mostly their fossils)
- Try to cross-correlate with galaxy population

But can we see the smaller bubbles and get  $Q_{\text{HII}}(z)$ ?

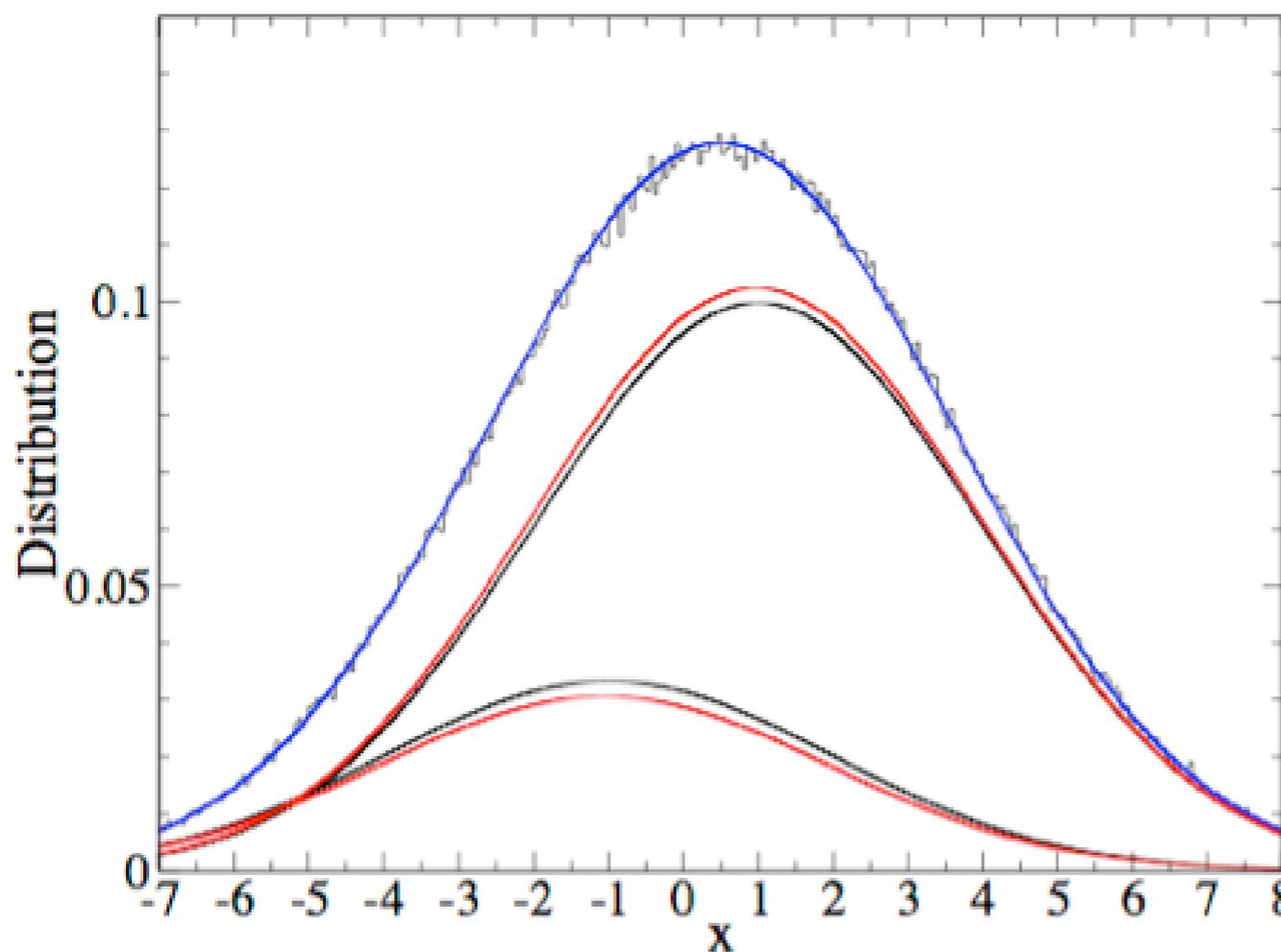
# Back to Basics: One Point Statistics

Hansen, Oh & Furlanetto  
(2007, in prep)



# One Point Statistics

Bubbles create bimodality in the PDF



Directly tells us  $Q_{HII}$   
(z)!!!

Can we pick it  
out?

## DETECTING BIMODALITY IN ASTRONOMICAL DATASETS

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### ABSTRACT

We discuss statistical techniques for detecting and quantifying bimodality in astronomical datasets. We concentrate on the KMM algorithm, which estimates the statistical significance of bimodality in such datasets and objectively partitions data into subpopulations. By simulating bimodal distributions with a range of properties we investigate the sensitivity of KMM to datasets with varying characteristics. Our results facilitate the planning of optimal observing strategies for systems where bimodality is suspected. Mixture-modeling algorithms similar to the KMM algorithm have been used in previous studies to partition the stellar population of the Milky Way into subsystems. We illustrate the broad applicability of KMM by analyzing published data on globular cluster metallicity distributions, velocity distributions of galaxies in clusters, and burst durations of gamma-ray sources. FORTRAN code for the KMM algorithm and directions for its use are available from the authors upon request.

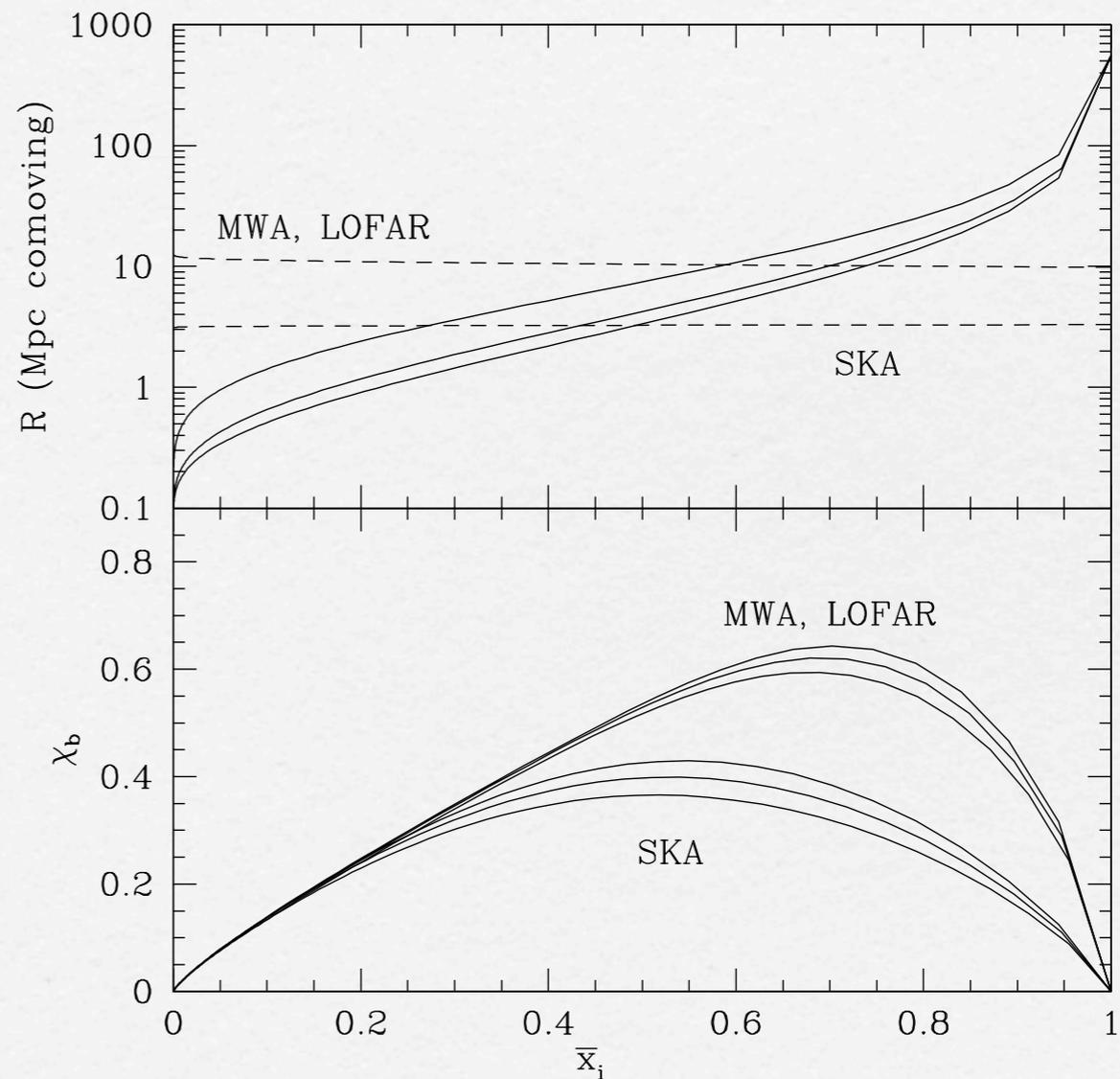
# ...partially ionized boundary pixels create complications

$$f_{bd} \approx 3 \frac{r_{pix}}{R_{bub}} Q_{\text{HII}}$$

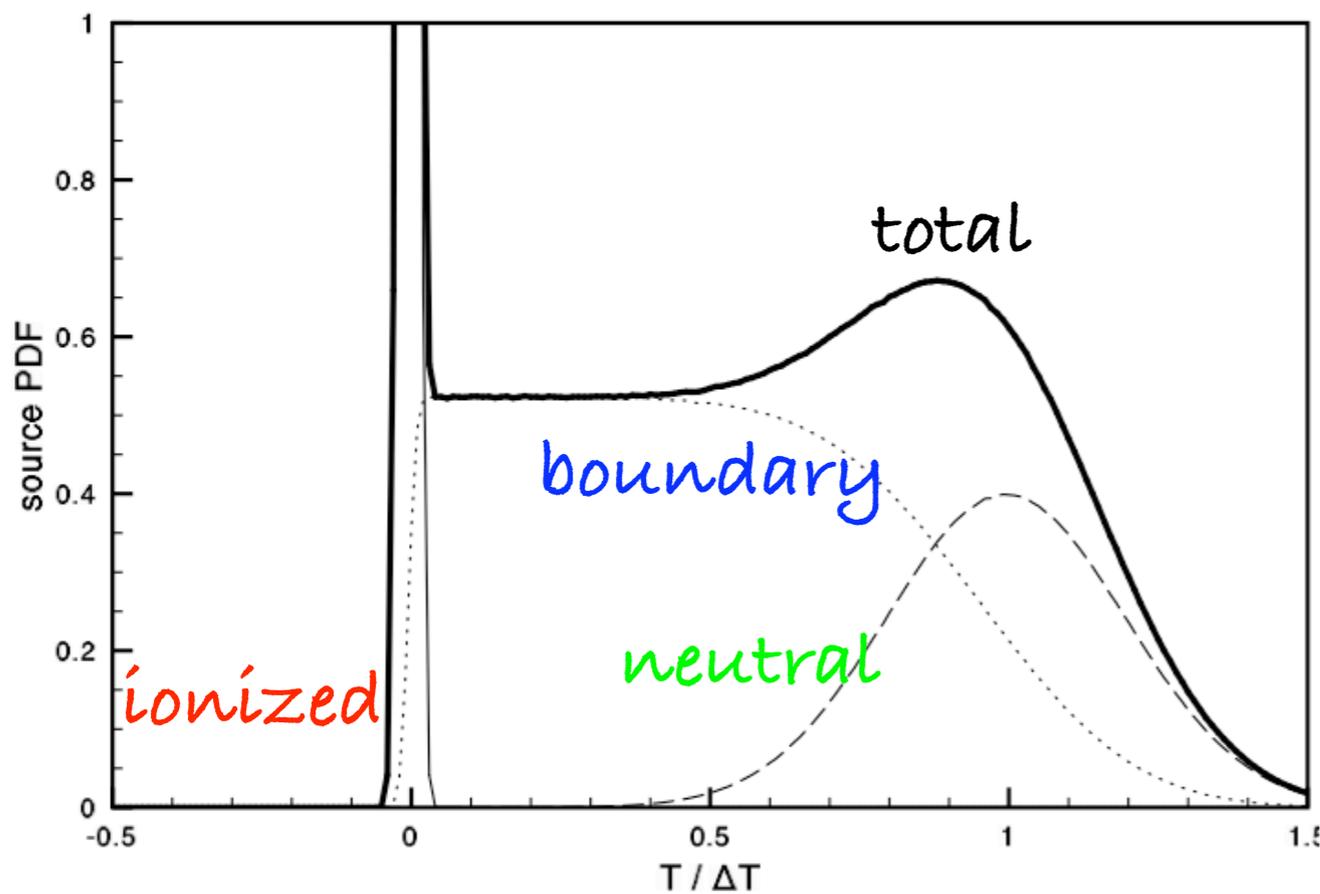
can be ~10-70% of pixels

Dependent on telescope resolution + bubble size

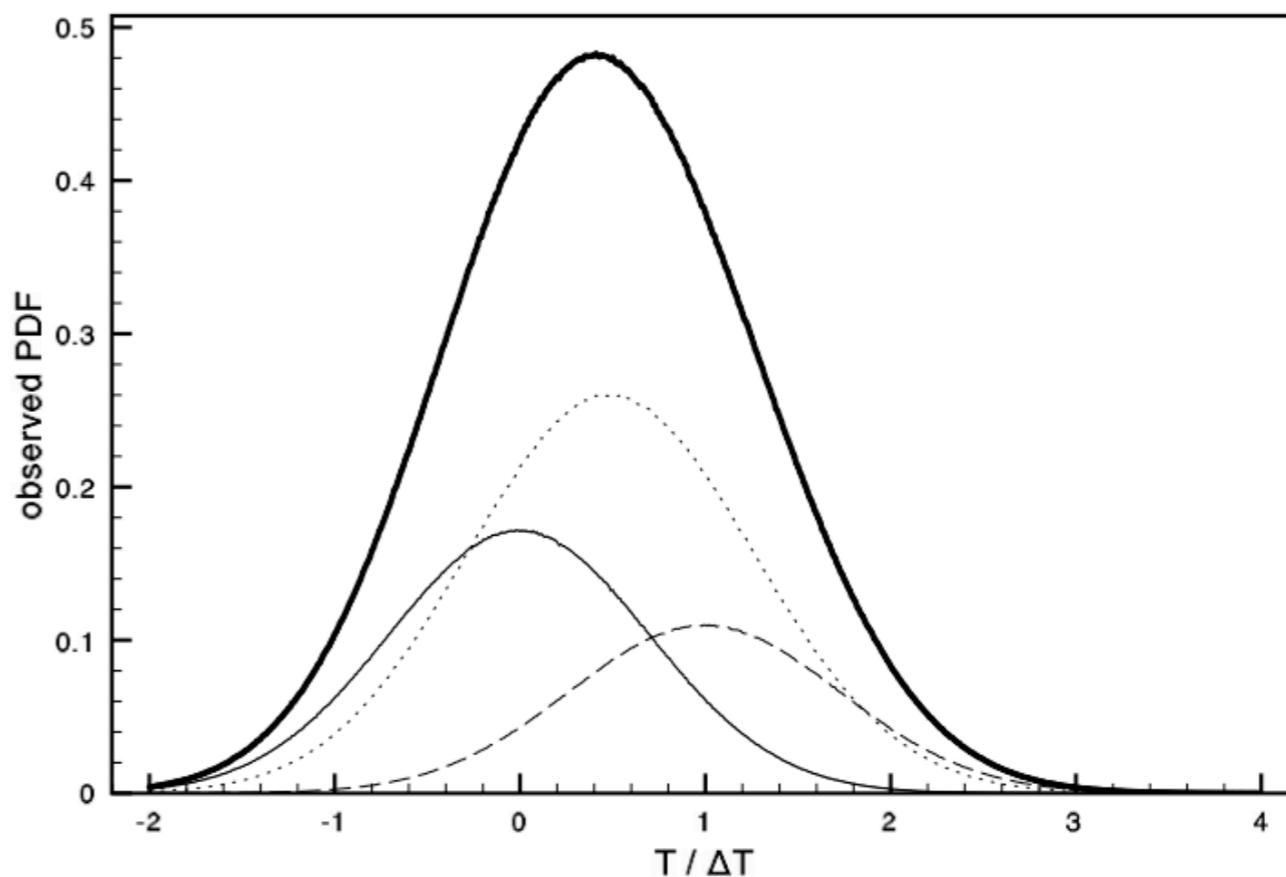
Also partially ionized pixels from x-rays, fossil regions



# ...here's a more realistic PDF



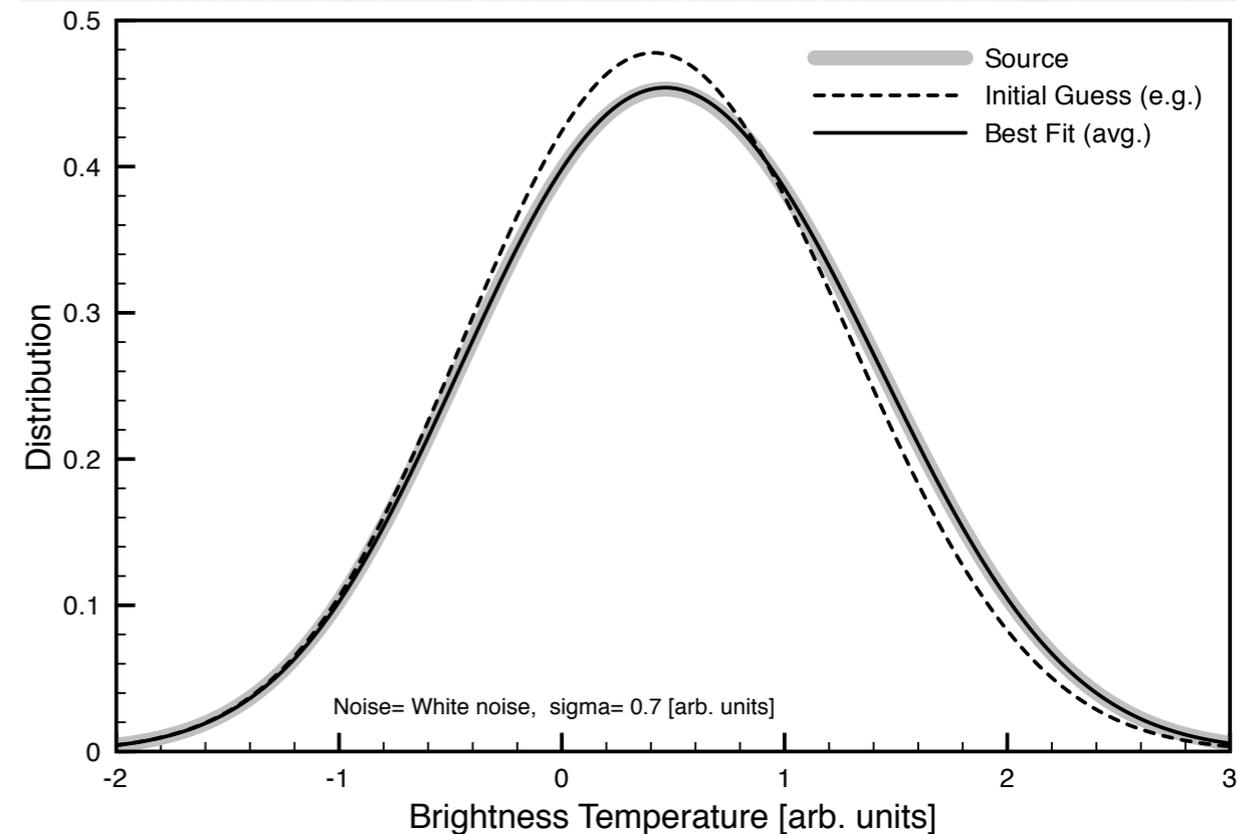
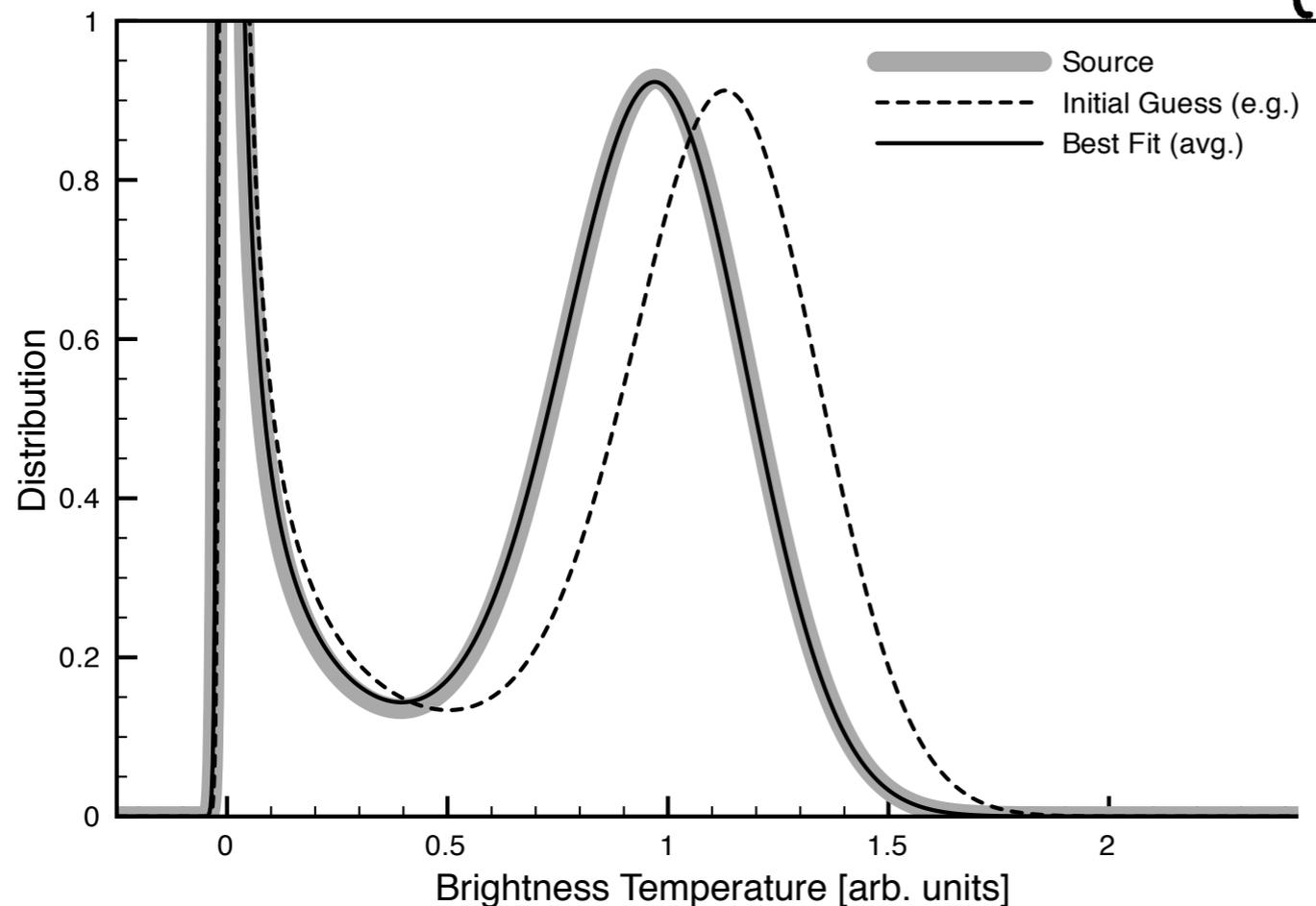
no noise



with noise

# It works!

Solve for populations via iterative Max-likelihood technique

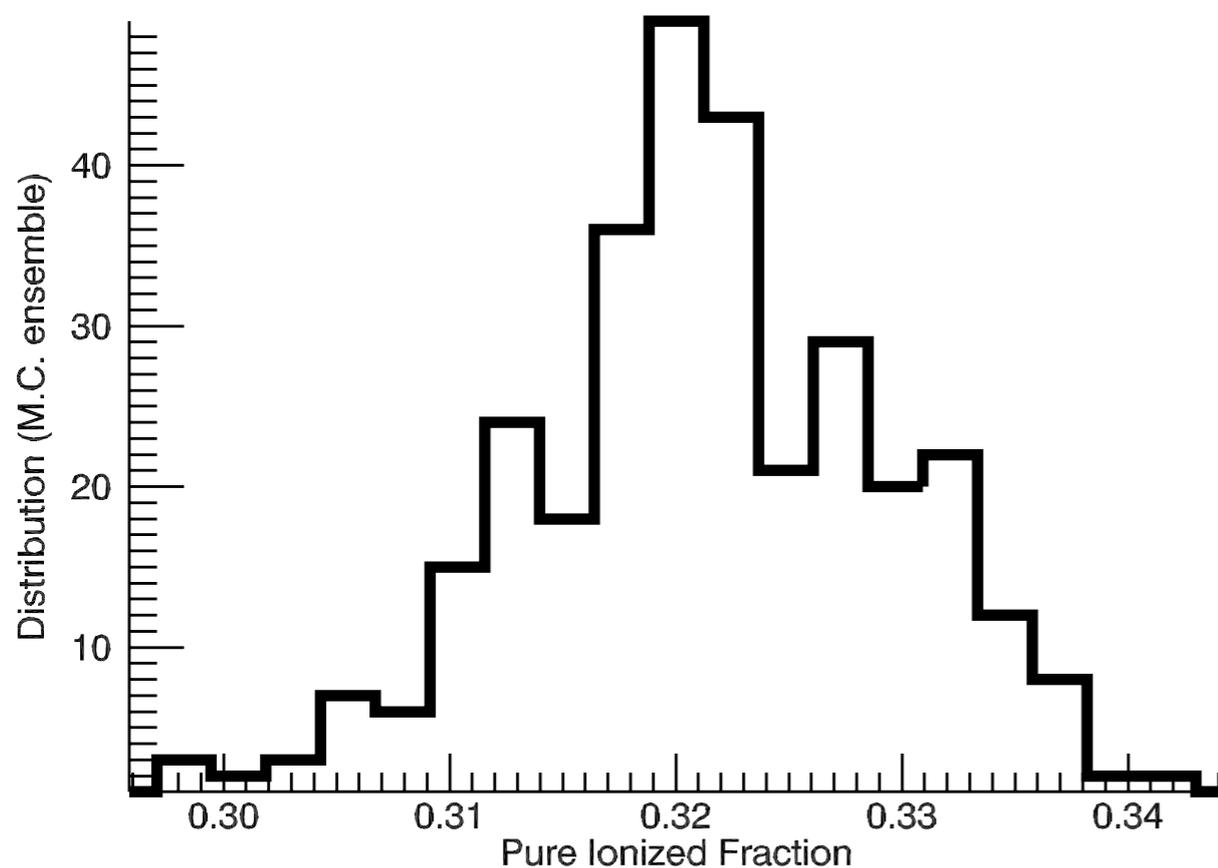


original distribution

with noise

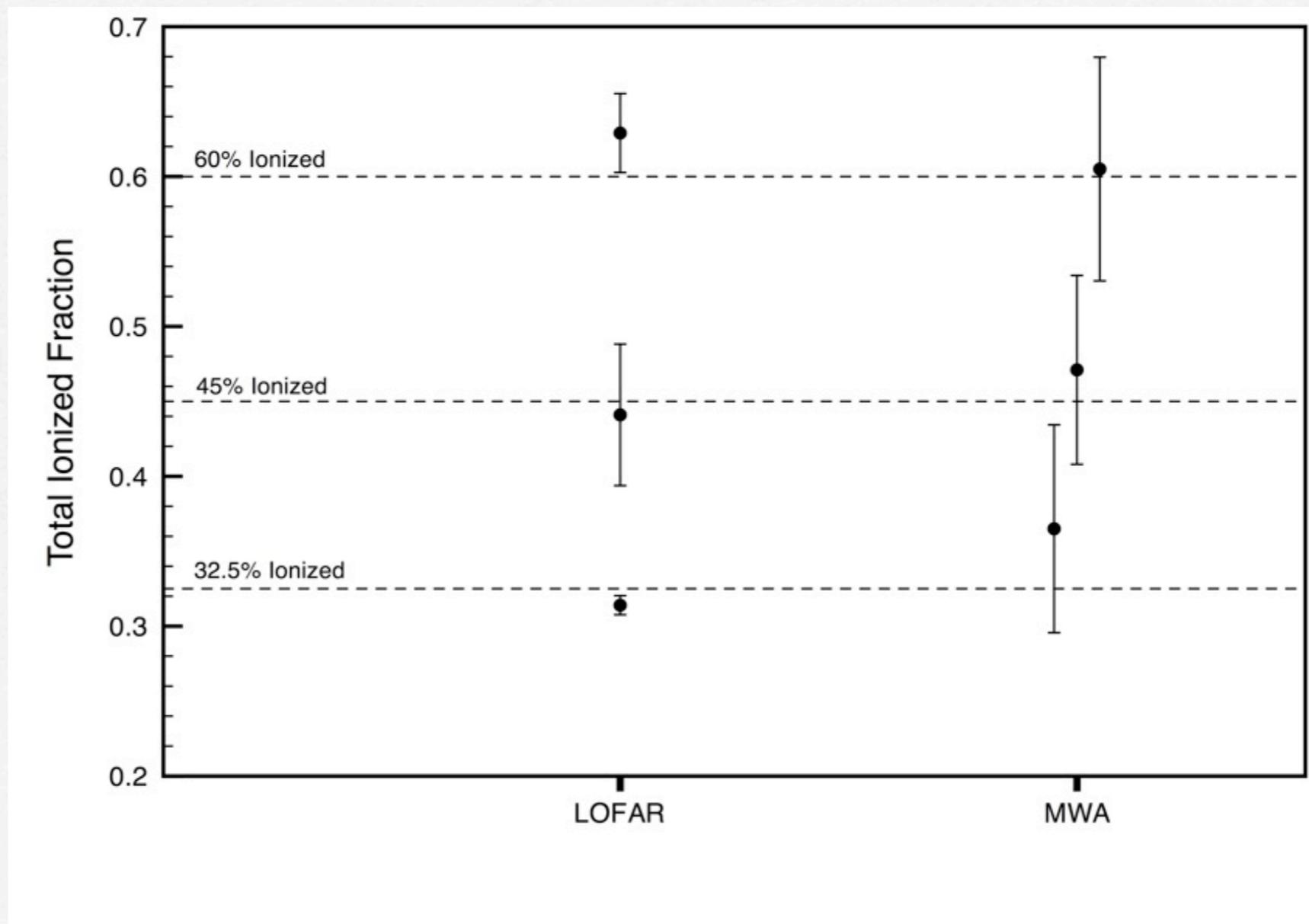
Leverage comes from having many pixels

# Monte Carlo Errors agree with Fisher Matrix estimates



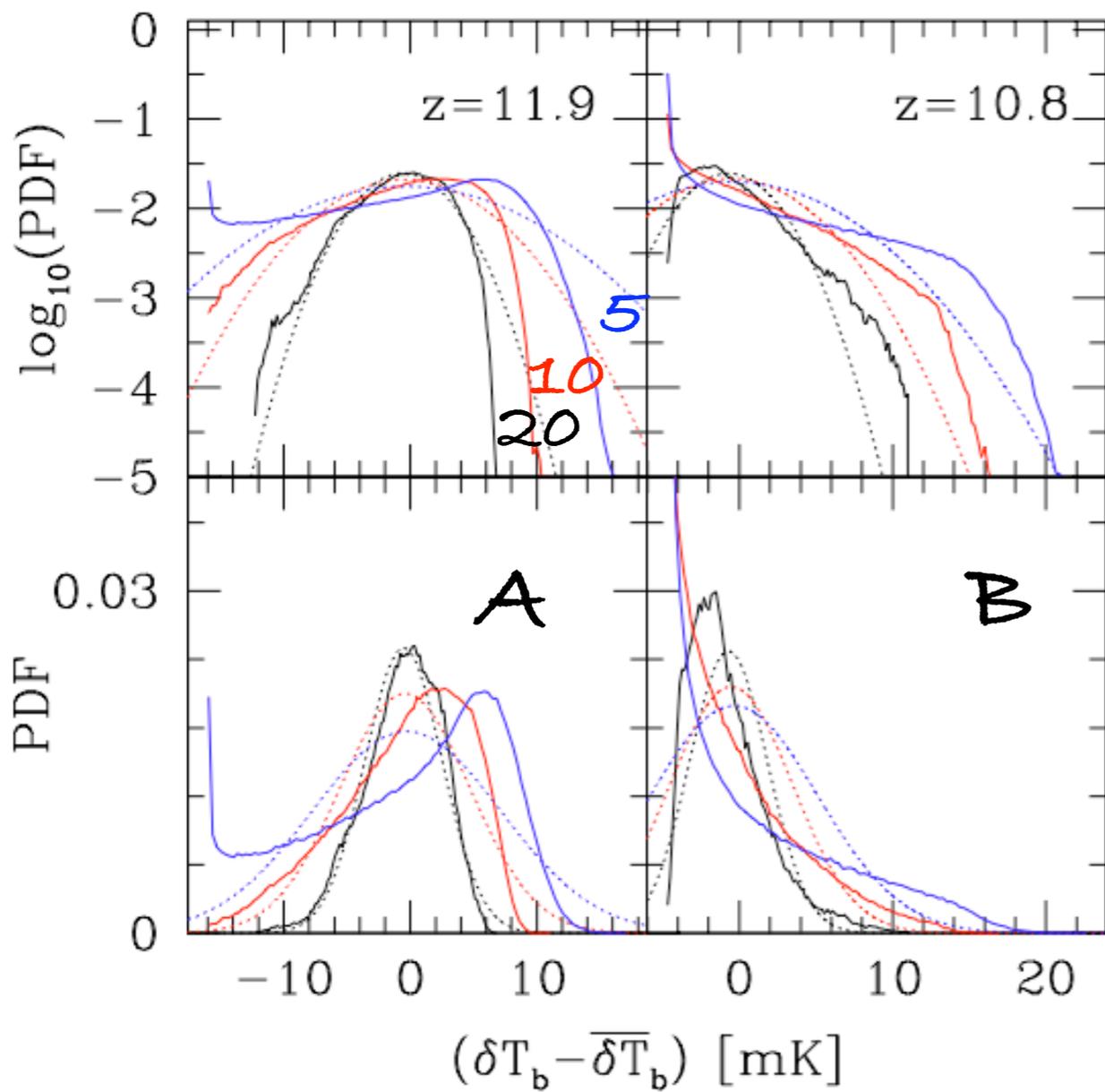
*An idealized case,  
but results are very  
encouraging...*

# Results...



Looks promising...

# In principle, PDF has info about topology too



A: Cutoff at high  $T_b$ --  
inside-out reionization

B: Tail at high  $T_b$ --  
islands of neutral gas  
inside HII regions

Note: distribution narrows  
as smoothing scale increases

# Direct Bubble Detection: The Canny Algorithm



Phelps et al (2007, in prep)

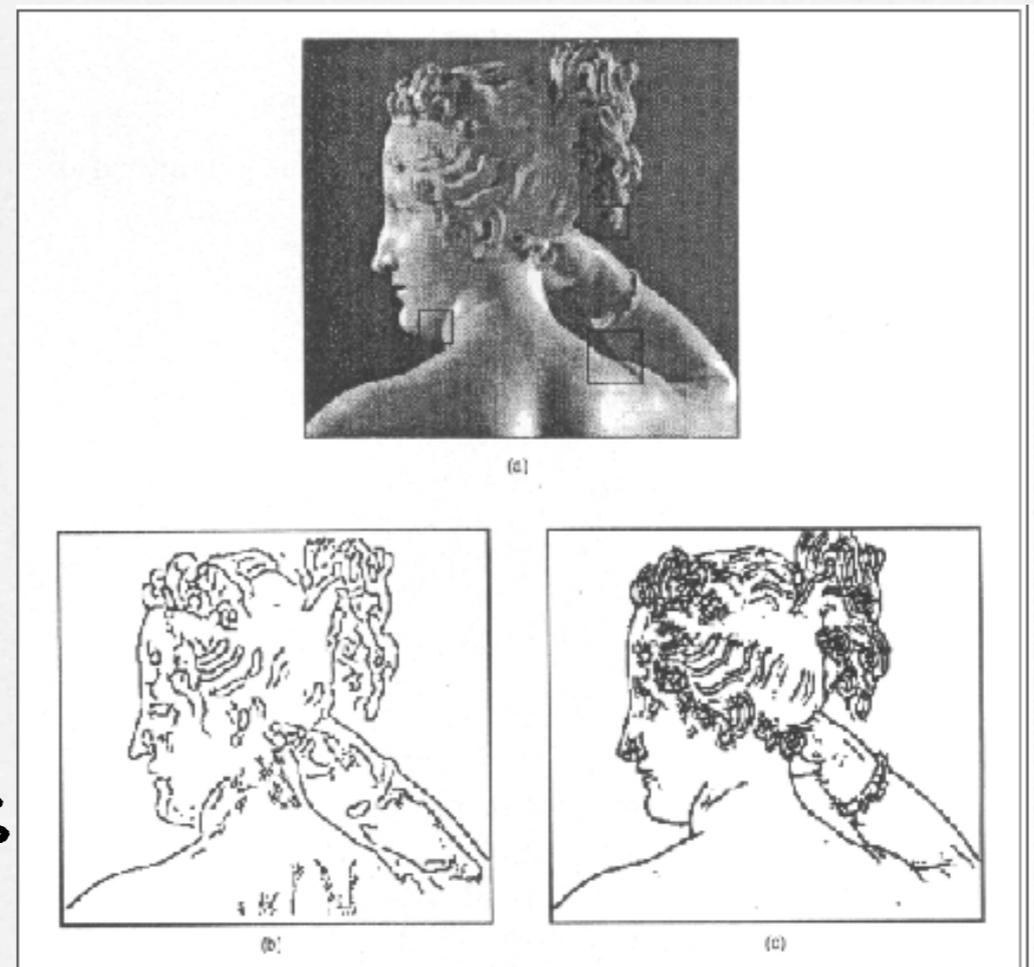
# How to detect bubbles directly?

Look for edges in noisy  
background: classic image  
processing problem

Canny Algorithm:

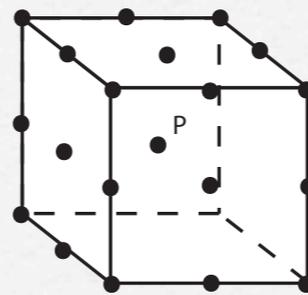
--optimal edge detector

--looks for maxima in derivatives  
of smoothed image

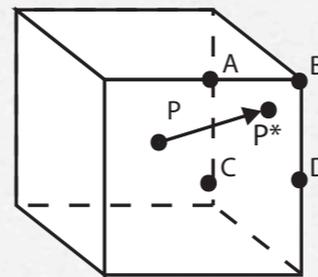


1. Apply Gaussian filter  
-- need to do this several times at different scales

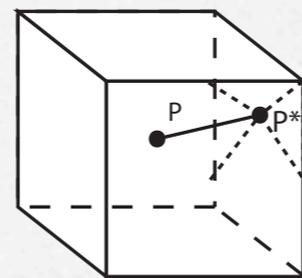
2. Find edge pixels  
-- find maximum in 3D spatial gradient



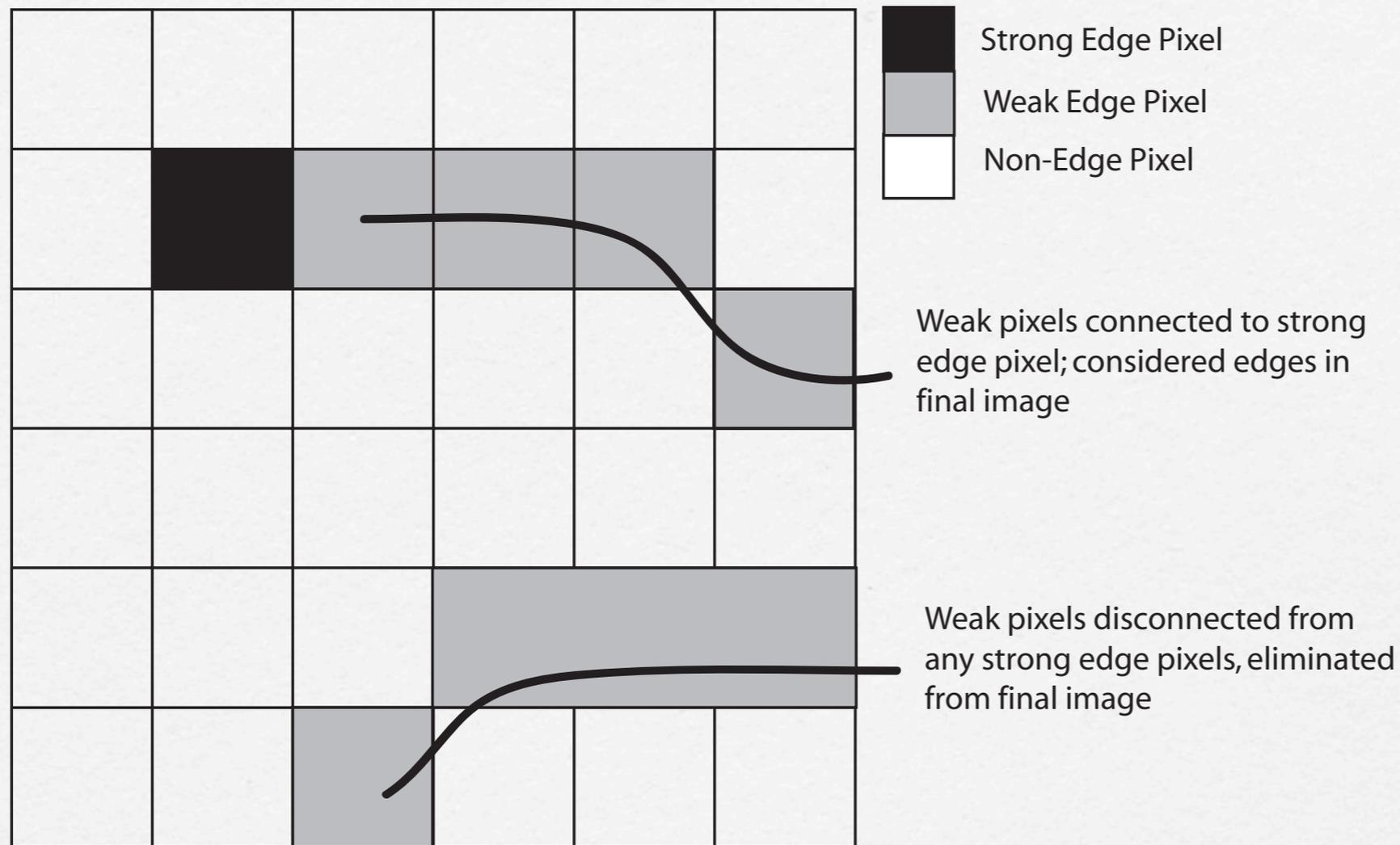
We want to determine whether pixel P (in the center of the cube) is a local maximum of the signal



By using the gradient information, we are able to determine that the gradient of P intersects the cube at point P\*, which lies closest to pixels A, B, C, and D.



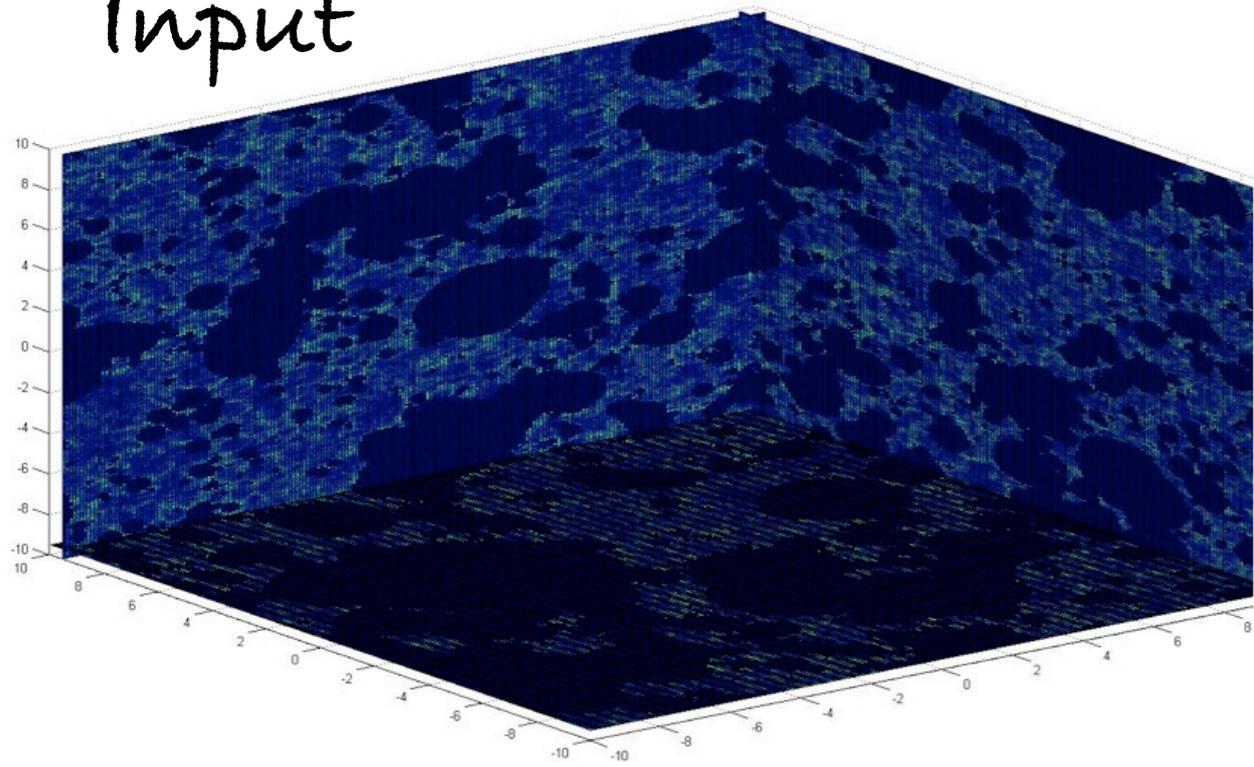
Interpolating the gradient magnitudes of the signal at pixels A, B, C, and D, we are able to determine the gradient magnitude of P\*. If this value, and the correlating value in the opposite direction, are less than the gradient magnitude at P\*, P\* is indeed a maximum.



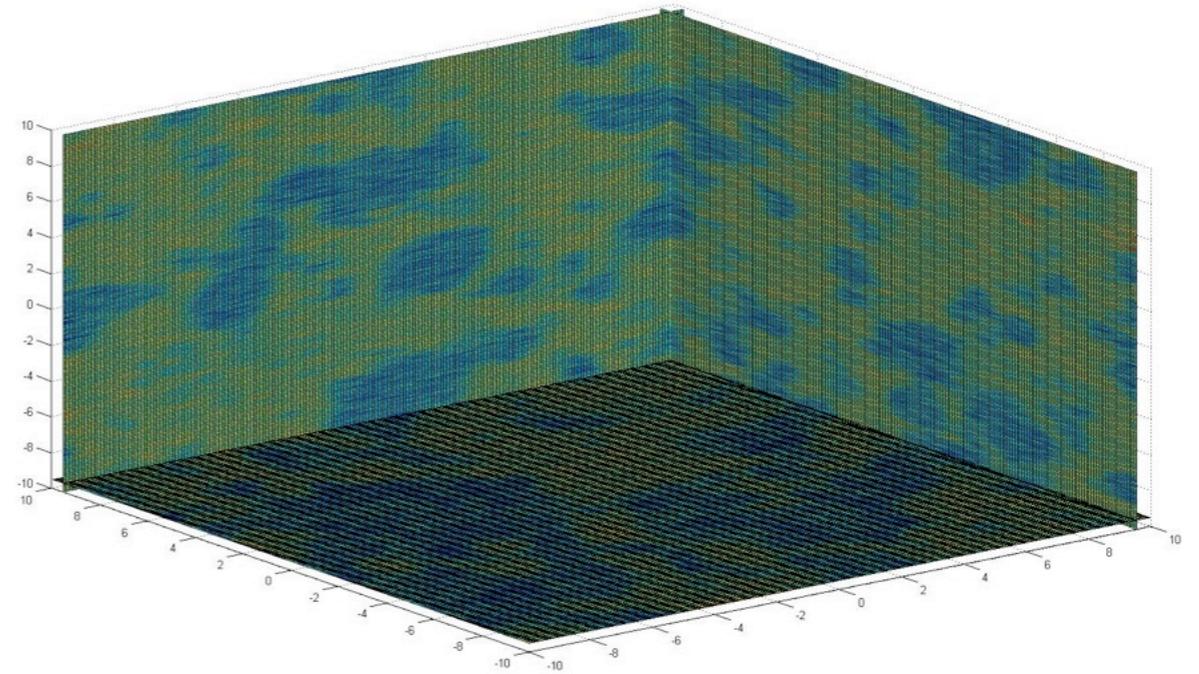
### 3. Apply thresholding with hysteresis

- strong pixel: automatically part of edge
- weak pixel: only part of edge if connected to strong pixel

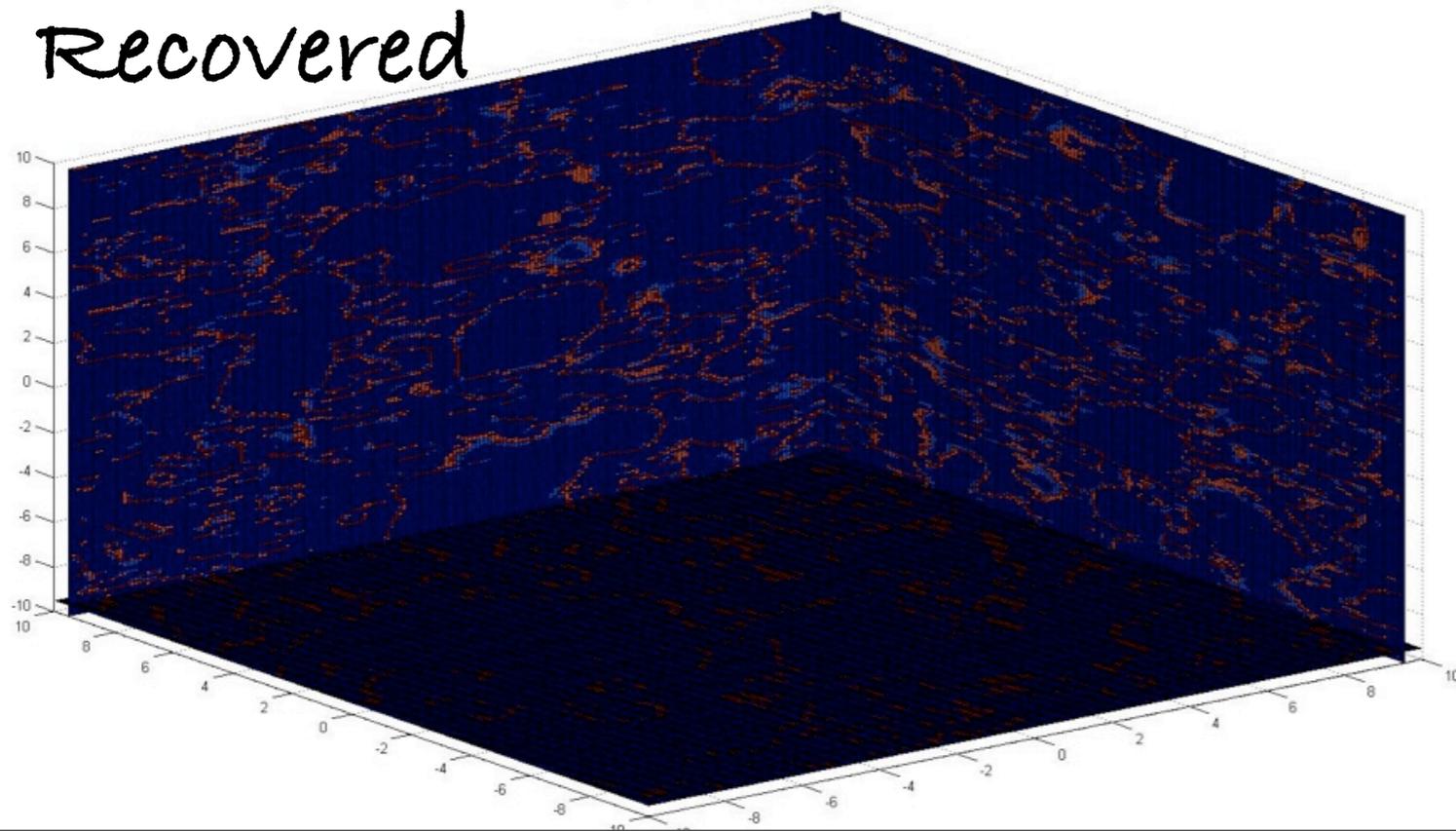
Input



with noise



Recovered



Sigma = 1.1, Highthresh = 7, Lowthresh = 2

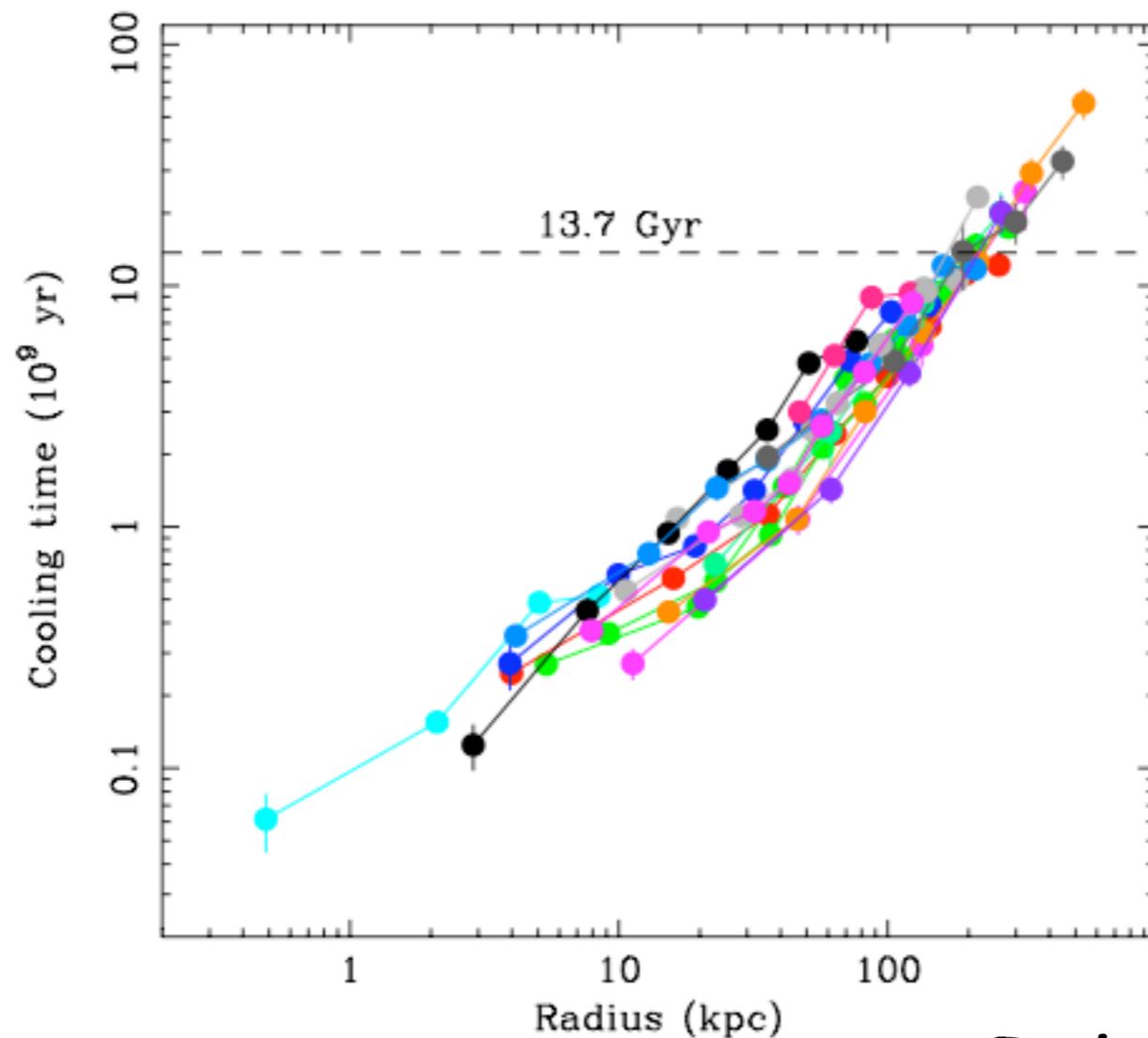
Input box from  
Mesinger & Furlanetto  
2007

# Bottom Line

- HII Bubbles are main feature (holes in 21cm emission) after first sources light up
- Much needed foreground calibrators
- Can only directly image biggest ones: sharpen detection w/ Canny algorithm
- If can detect statistically, obtain  $Q_{\text{HII}}(z)$
- More work needed!

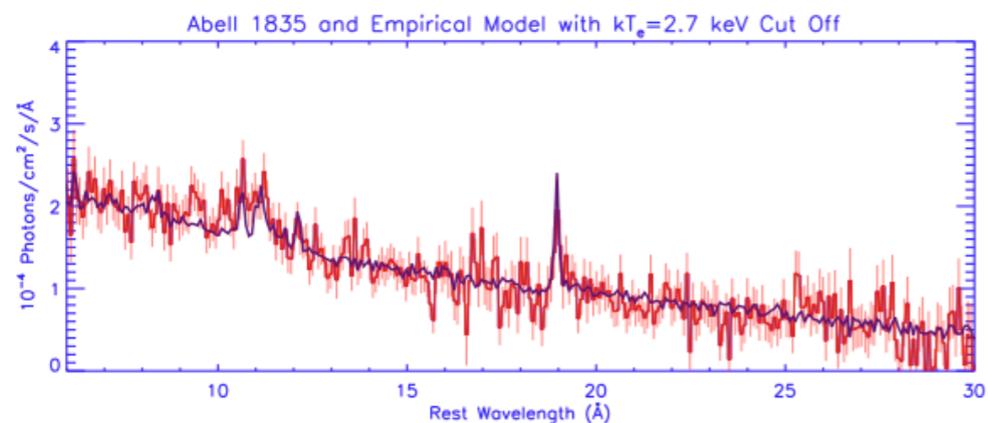
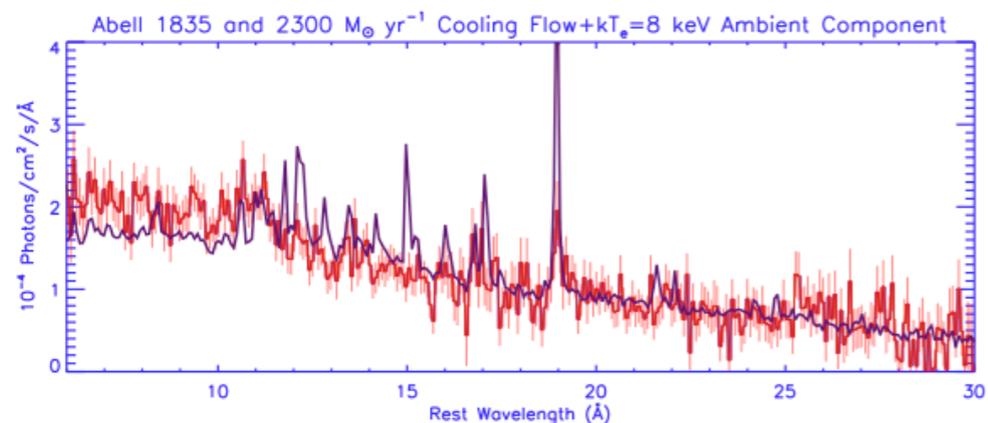
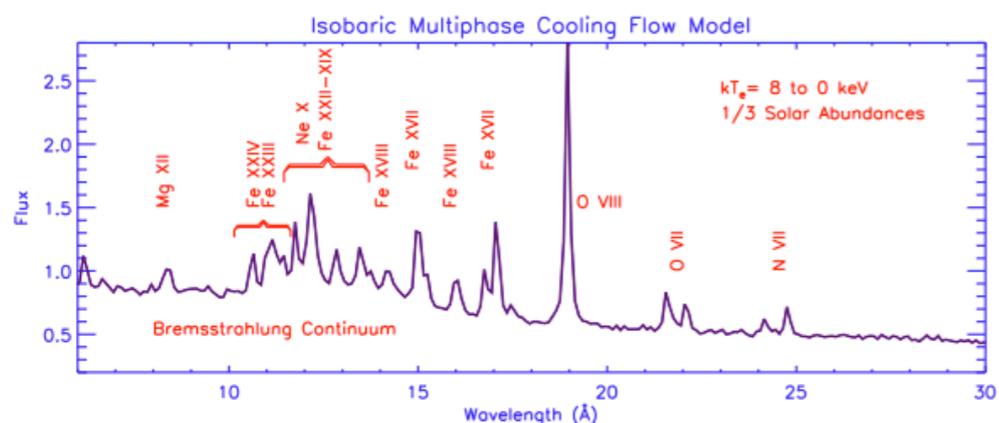
# Feedback Heating in Galaxy Clusters

# Although gas cooling times in clusters are short...



Peterson & Fabian (2006)

# Gas does not appear to cool below $\sim 1/3$ of $T_{\text{vir}}$



can only fit spectra if prevent gas from cooling below  $\sim 1/3$  of ambient temperature

universal across different cluster temperatures

Peterson et al (2001)

# What's going on?

Plasma Physics effect? (e.g.,)

- Turbulent mixing
- non-Maxellian distribution function

Feedback heating (e.g...)

- thermal conduction
- AGN mechanical heating
- cosmic ray heating (this talk)

# The Straw Man: Isobaric Cooling Flow

For a constant mass-flow rate  $\dot{M}$ , the power release is thus

$$dL = \frac{5}{2} \frac{k_B}{\mu m_H} \dot{M} dT. \quad (1)$$

We also have, from the definition of the cooling function  $\Lambda$ ,

$$\begin{aligned} dL &= n_e n_H \Lambda(T, Z) dV, \\ dL_\nu &= n_e n_H \Lambda_\nu(T, Z) dV. \end{aligned} \quad (2)$$

Hence we obtain the spectral power for a steady-state flow cooling from  $T_{\max}$  to  $T_{\min}$

$$L_\nu = \frac{5}{2} \frac{k_B}{\mu m_H} \dot{M} \int_{T_{\min}}^{T_{\max}} \frac{\Lambda_\nu(T, Z)}{\Lambda(T, Z)} dT. \quad (3)$$

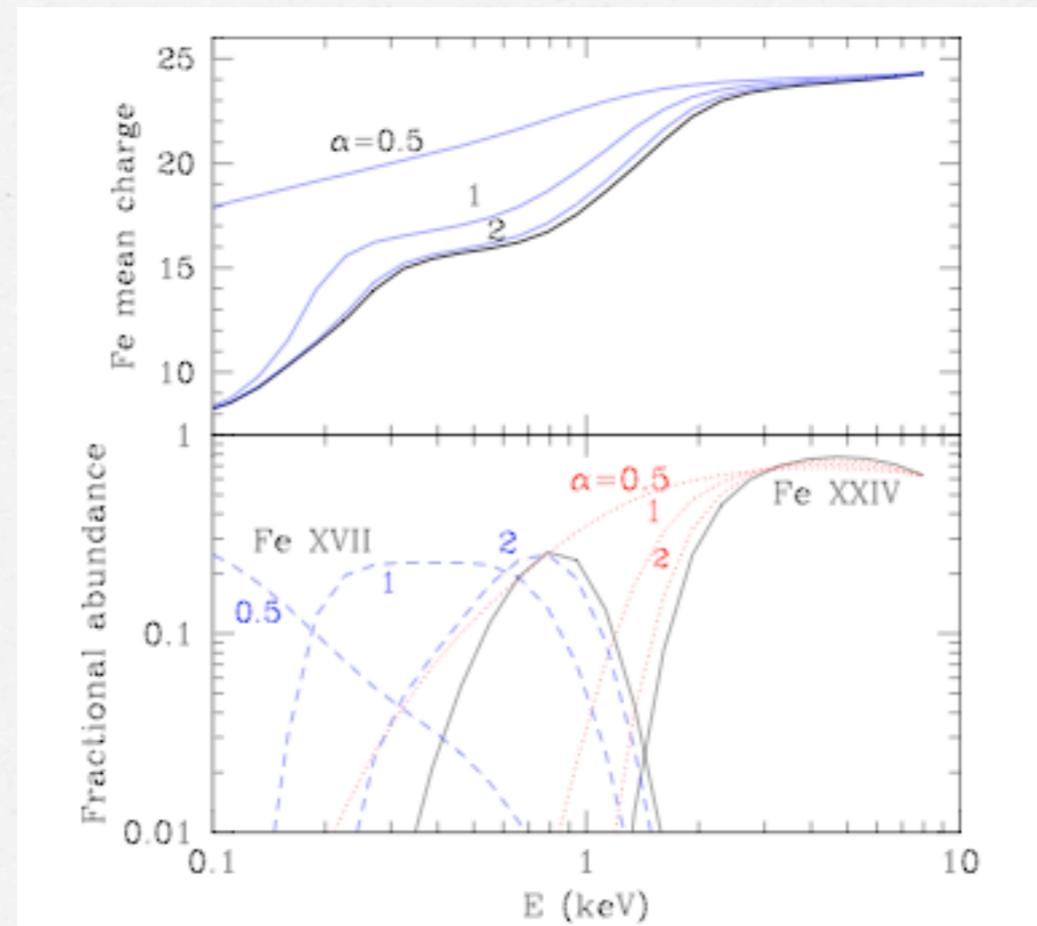
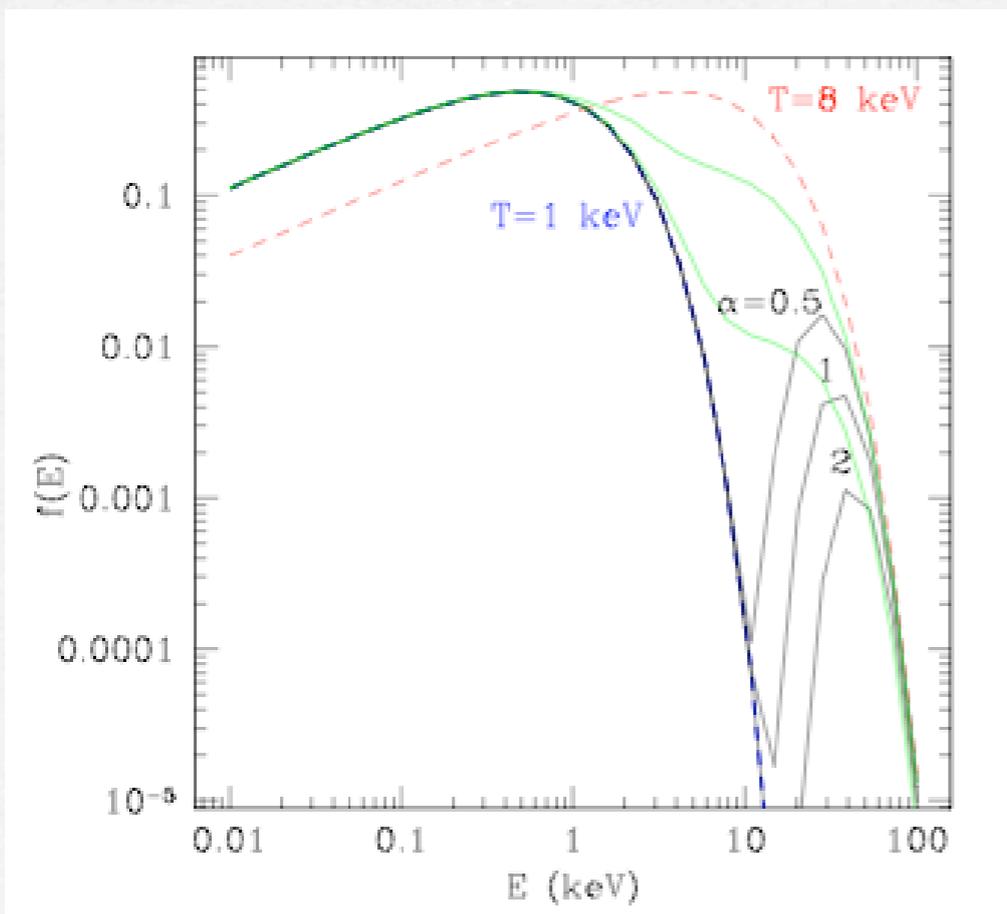
*e.g. see Fabian (1994)*

# Plasma effects I: Turbulent mixing?

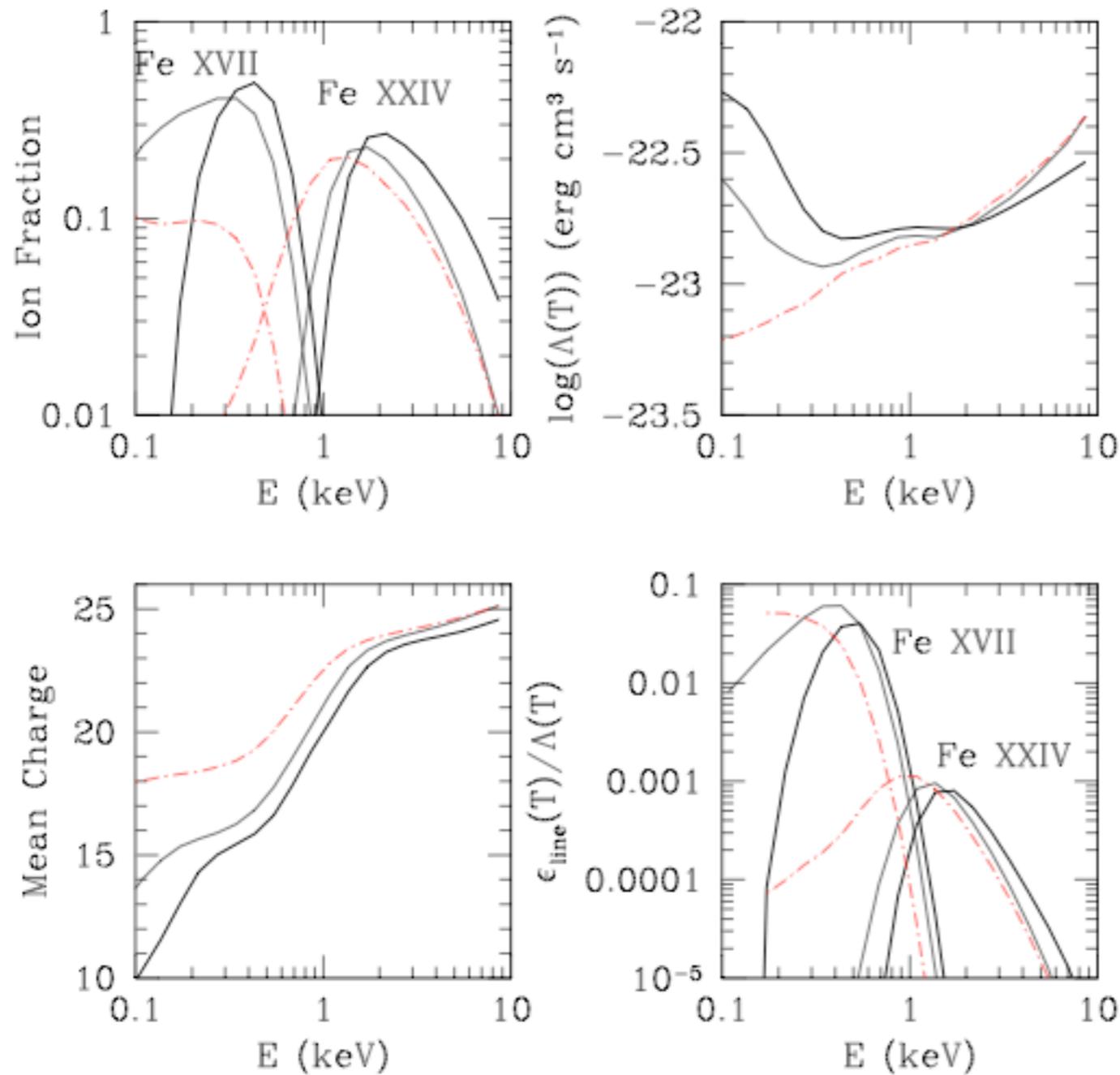
Gas cools non-radiatively by mixing hot and cold gas (Begelman & Fabian 1990)

Energy eventually escapes through optical lines

# Plasma Effects II: break collisional equilibrium?



Non-Maxwellian distribution function (Oh 2004)

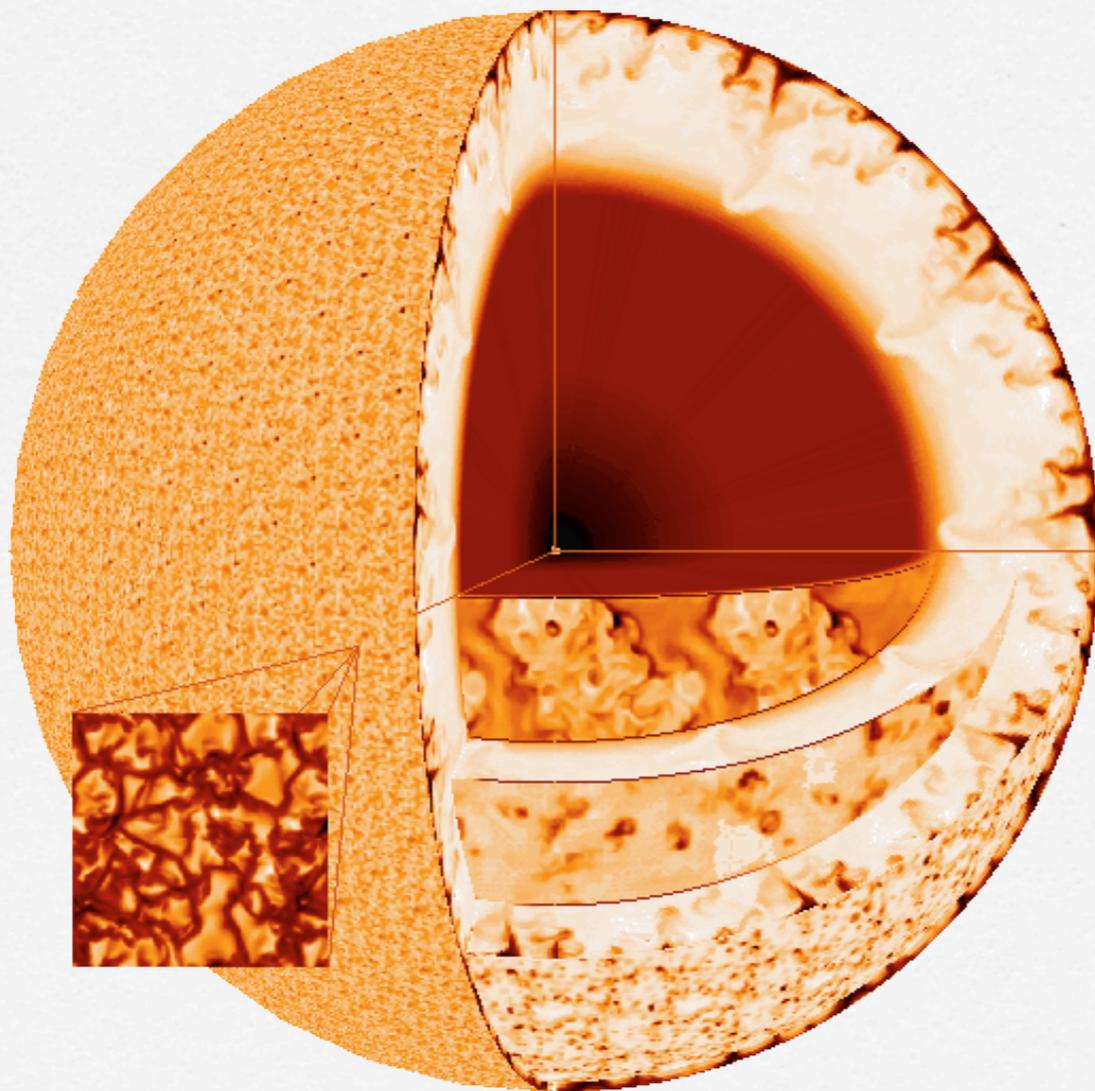


Photoionization from surrounding cluster gas or central AGN

Decreased cooling efficiency at low  $T$ : emission measure of low ionization state lines larger

Photoionization (Oh 2004)

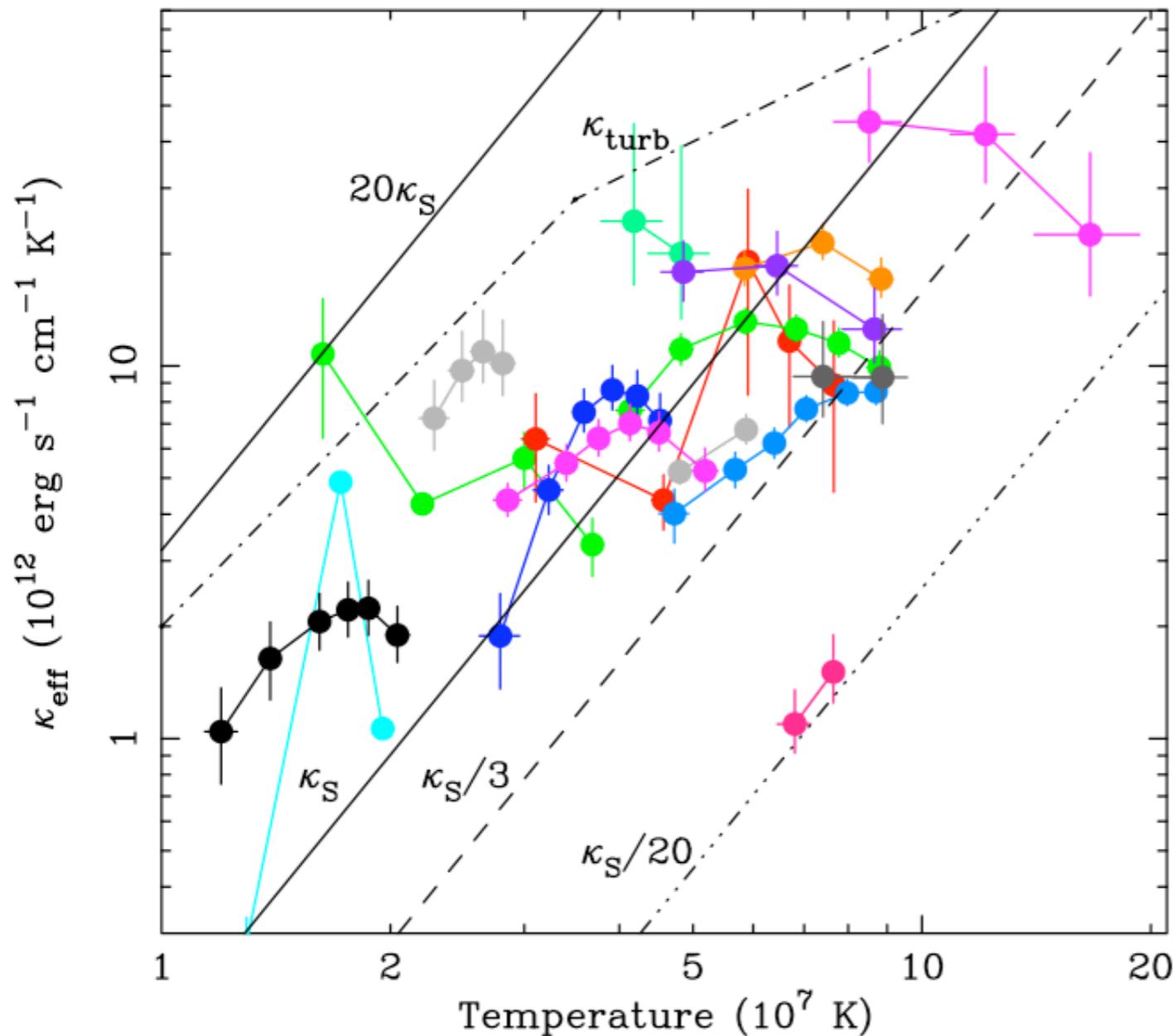
# Heating is a more popular solution



Cluster sits in quasi-thermal equilibrium: just like a star!

Also explains lack of cold gas/stars

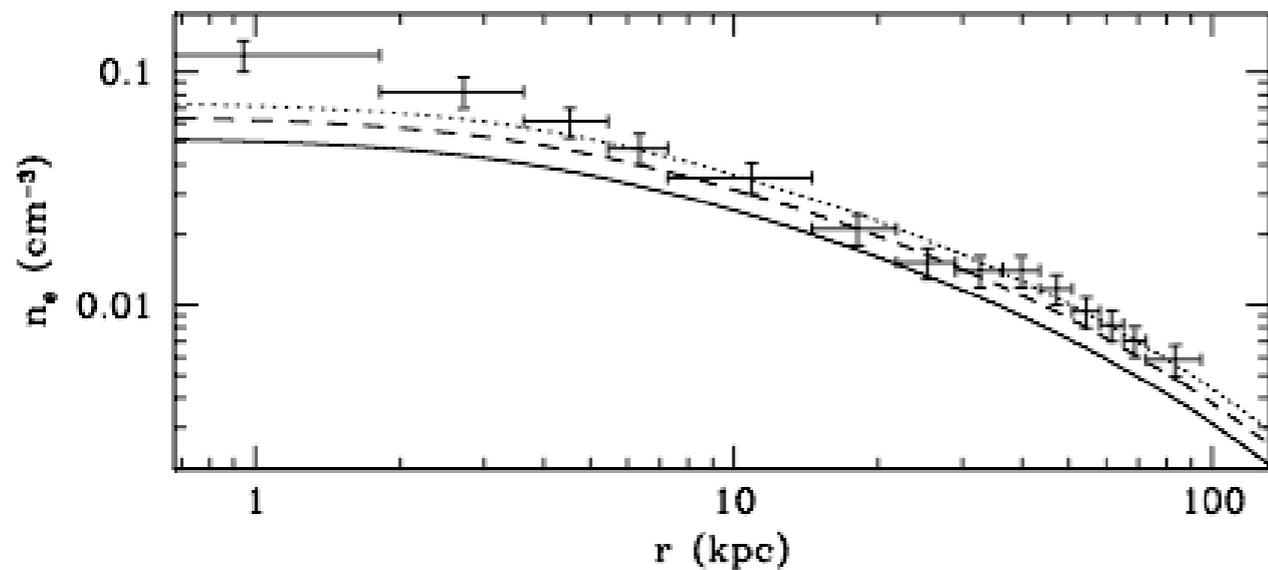
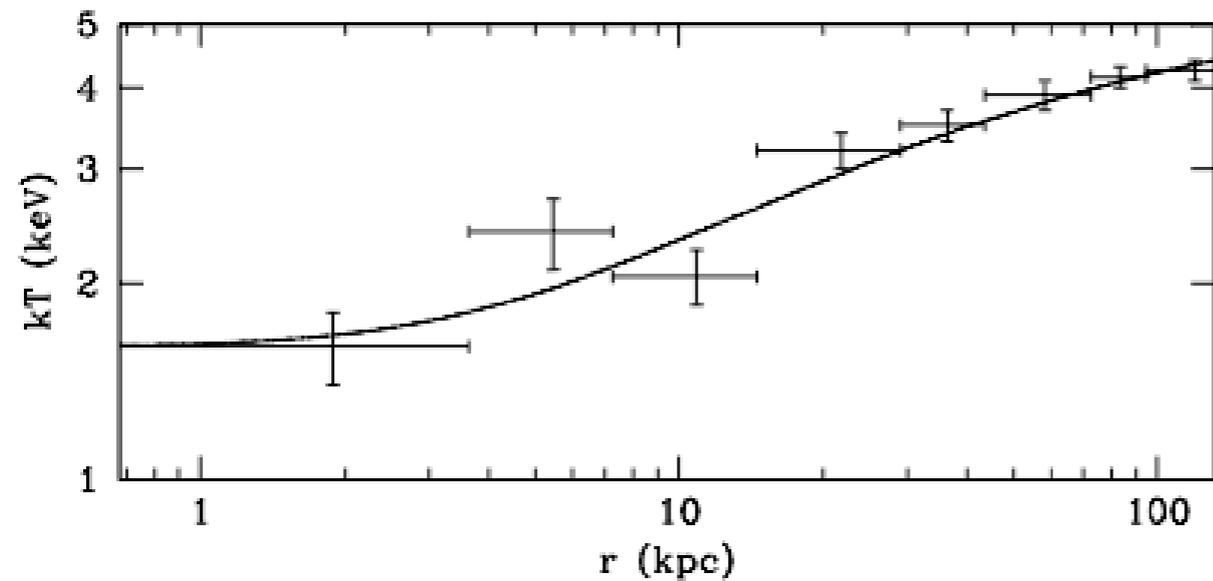
# Thermal Conduction



Conduction at  
fraction of classical  
Spitzer value close  
to what's needed.

Coincidence??

Peterson & Fabian (2006)

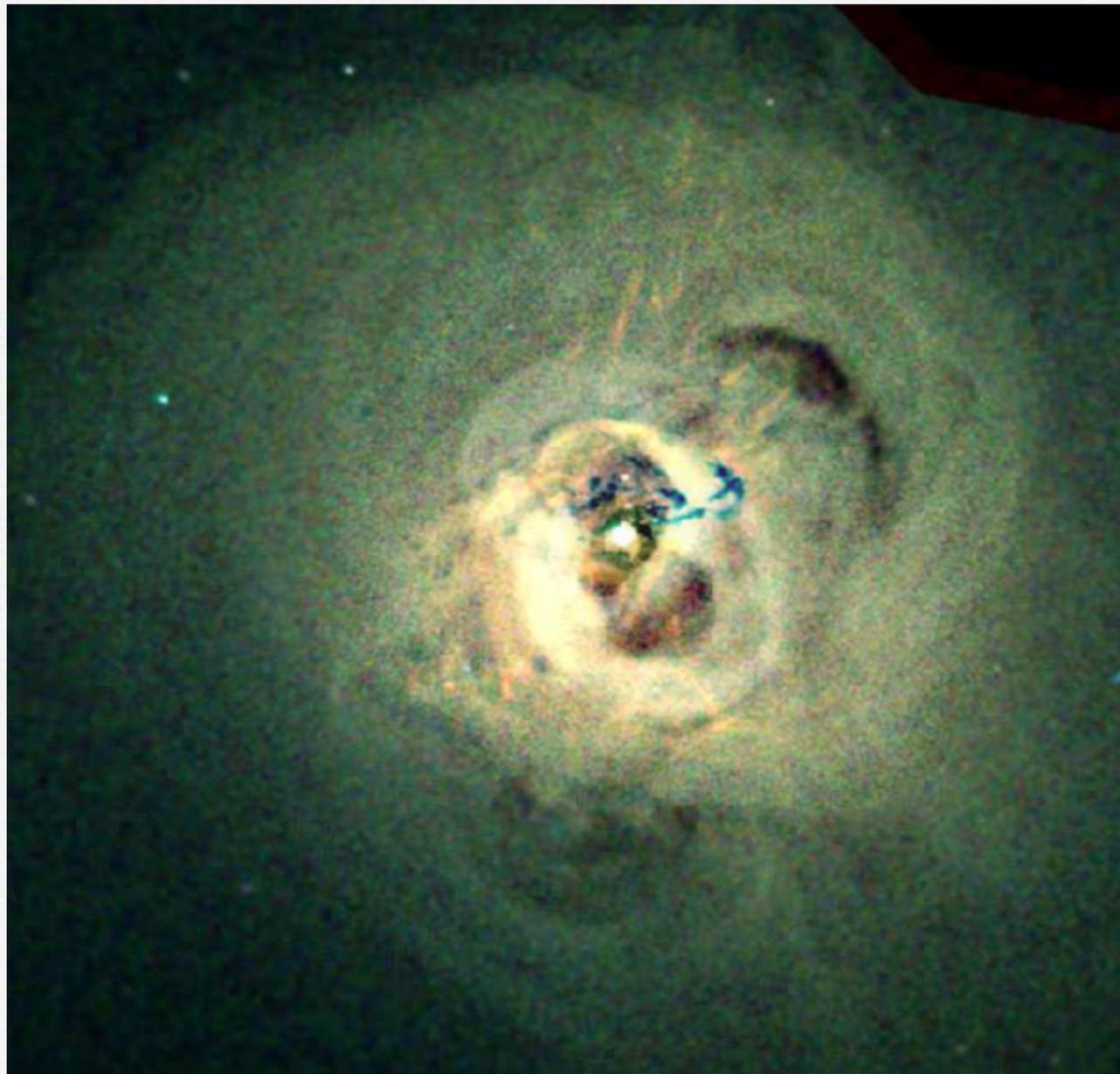


can build  
conduction-only  
models in  
hydrostatic and  
thermal equilibrium

But: suffer fine-  
tuning problems,  
tend to be globally  
unstable

Zakamska & Narayan (2002)

# AGN/radio galaxy heating



Chandra image, Perseus cluster

Bubbles observed in ICM,  
filled with hot/relativistic  
plasma

Maybe: entrain cold gas  
pdv work

This talk: **cosmic ray  
heating**

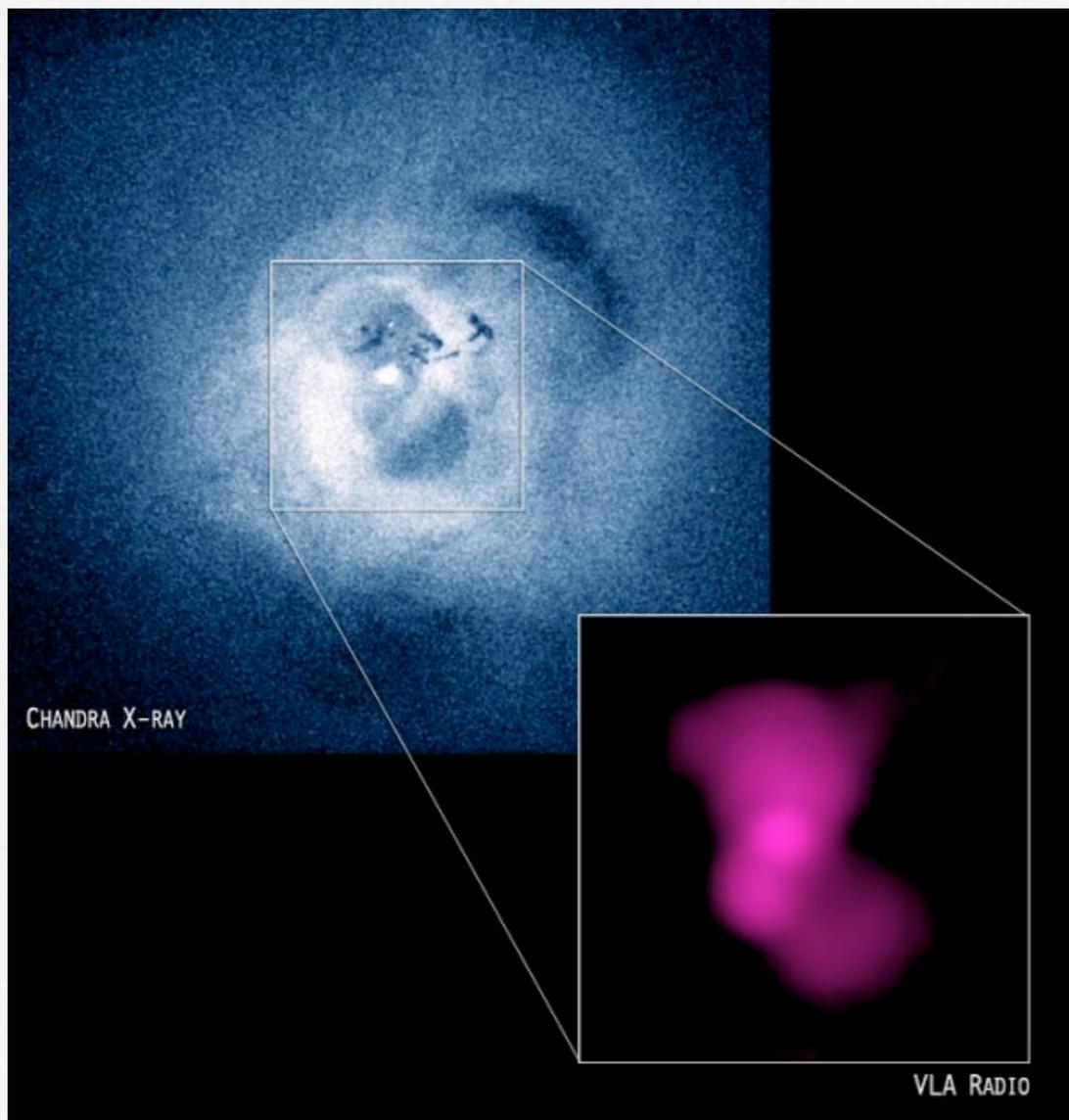
(Guo & Oh 2007)

# Why cosmic rays ?

We see radio synchrotron  
emission

Spallation products indicate  
CRs could be present (Nath, Madau &  
Silk 2005)

Many sources: jets,  
accretion shock, SN  
Provide gentle,  
distributed heating



# It's been tried before...

- Authors have considered dynamical and heating effects (via Coulomb, hadronic and Alfvén wave interactions) (Boehringer & Morfill 1988, Loewenstein et al 1991, Repaheli & Silk 1995, Colafrancesco et al 2004, Jubelgas et al 2006, Prommer et al 2006)
- None have constructed models where CRs successfully stop cooling flow

# A key problem: CR transport is slow

$$\mathbf{F}_c = \gamma_c E_c (\mathbf{u} + \mathbf{v}_A) - \mathbf{n} \kappa_c (\mathbf{n} \cdot \nabla E_c), \quad (\text{A14})$$

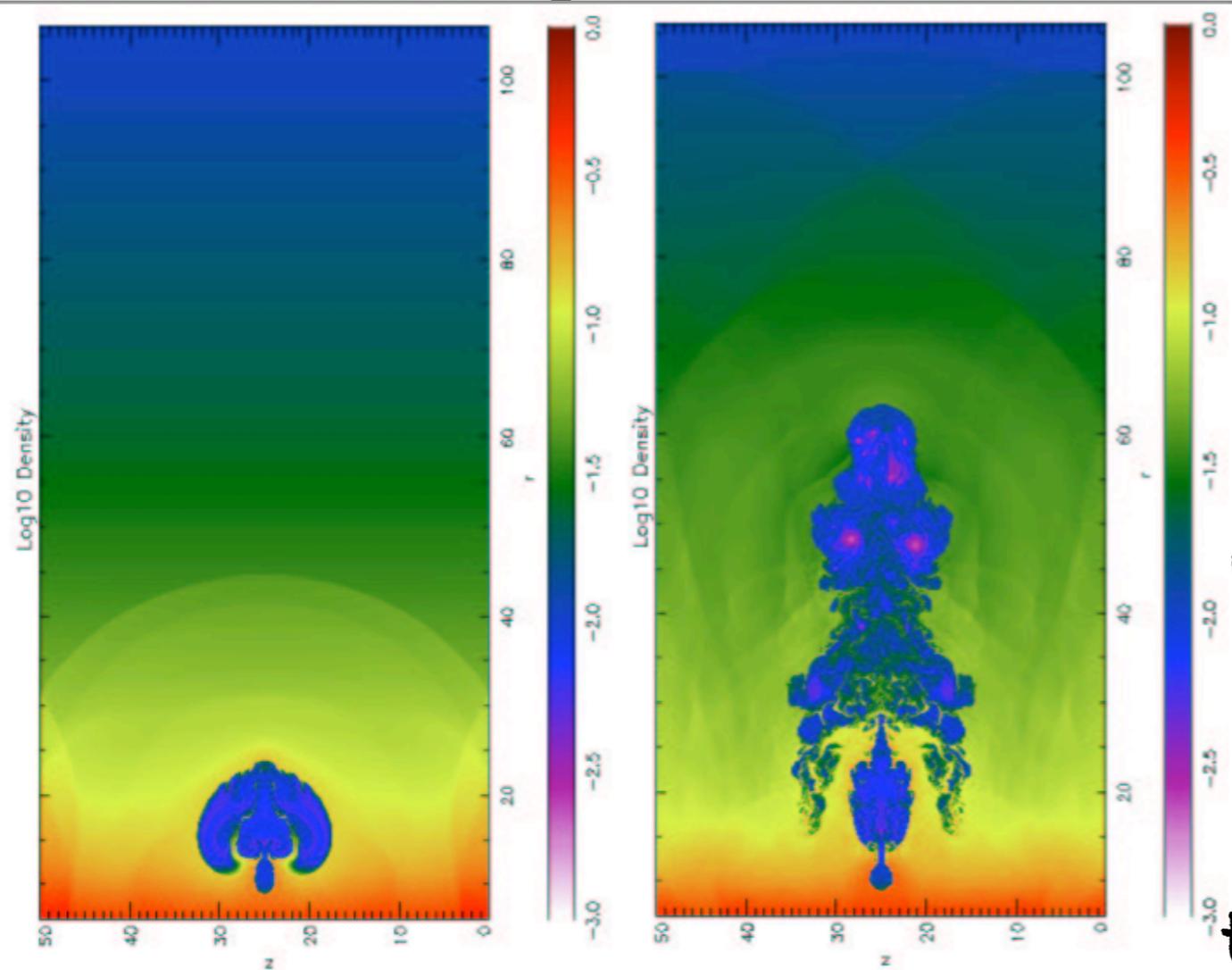
$$\frac{\partial E_c}{\partial t} = (\gamma_c - 1) (\mathbf{u} + \mathbf{v}_A) \cdot \nabla E_c - \nabla \cdot \mathbf{F}_c + \bar{Q}. \quad (\text{A15})$$

Diffusive and other CR transport timescales are long

Leads to overpressured center with insufficient heating at outskirts (though may drive turbulent convection: Chandran & collaborators)

# Our model: use bubbles to transport CRs

Bubbles disrupted by Rayleigh-Taylor & Kelvin-Helmholtz instabilities as rise



(Also: CRs diffuse out)  
Fast way of transporting CRs: rise time ~ sound crossing time

Bruggen & Kaiser (2002)

# Method

- 1D Zeus code: solve time-dependent hydrodynamic equations + CR heating & transport equations
- calculate steady steady CR spectrum, assuming Coulomb, hadronic and Alfvén-wave energy losses (latter dominates):

$$\Gamma_{wave} = v_A \frac{dP_c}{dr}$$

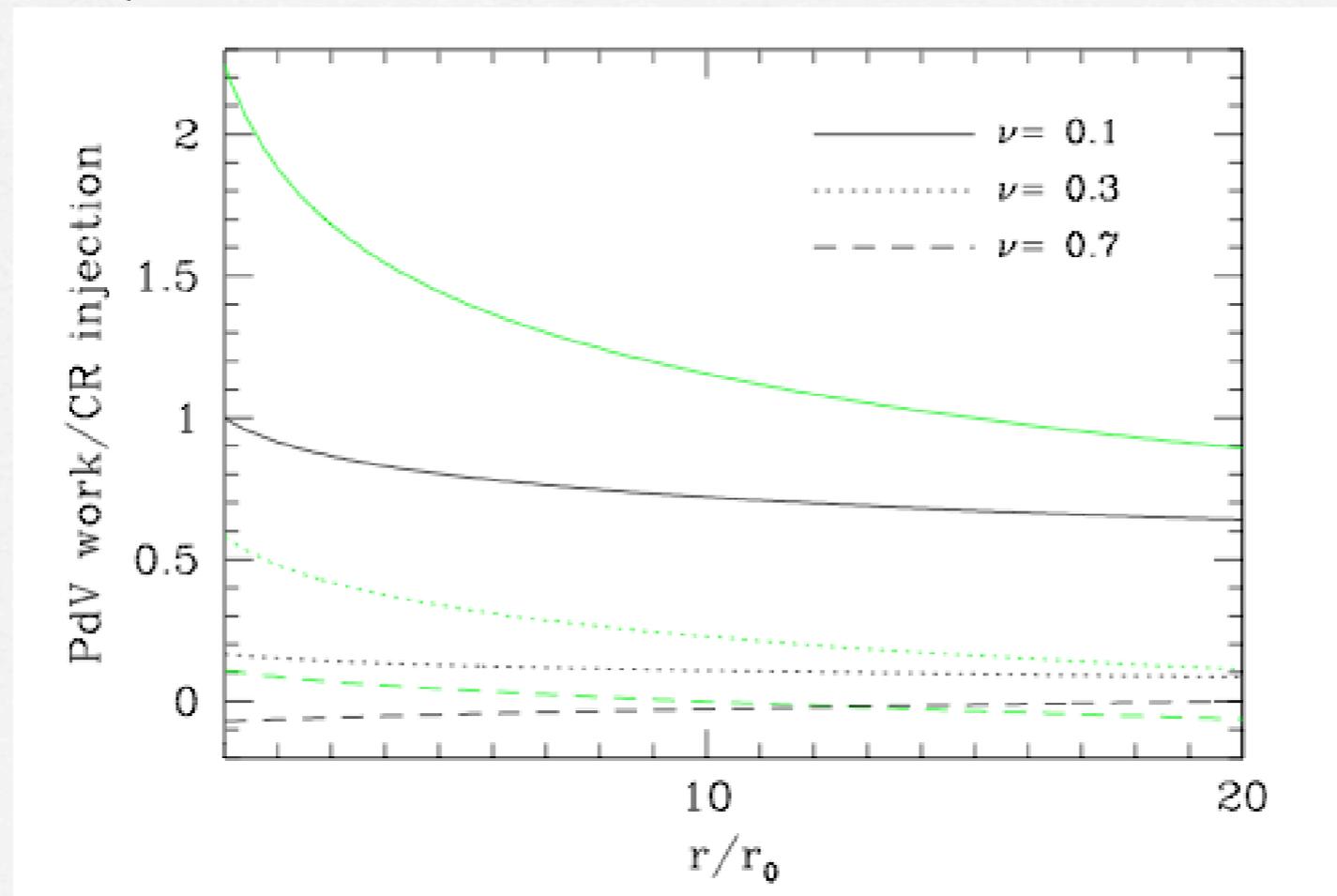
- Assume energy density in bubbles is a power-law with radius (note: CR injection rate depends on gas cooling---feedback effect)

$$L_{\text{bubble}} \sim -\epsilon \dot{M}_{\text{in}} c^2 \left( \frac{r}{r_0} \right)^{-\nu} \quad \text{for } r > r_0,$$

$$\begin{aligned} Q_c = \nabla \cdot \mathbf{F}_{\text{bubble}} &\sim -\frac{1}{4\pi r^2} \frac{\partial L_{\text{bubble}}}{\partial r} \left[ 1 - e^{-(r/r_0)^2} \right] \\ &\sim -\frac{\nu \epsilon \dot{M}_{\text{in}} c^2}{4\pi r_0^3} \left( \frac{r}{r_0} \right)^{-3-\nu} \left[ 1 - e^{-(r/r_0)^2} \right], \end{aligned} \quad (19)$$

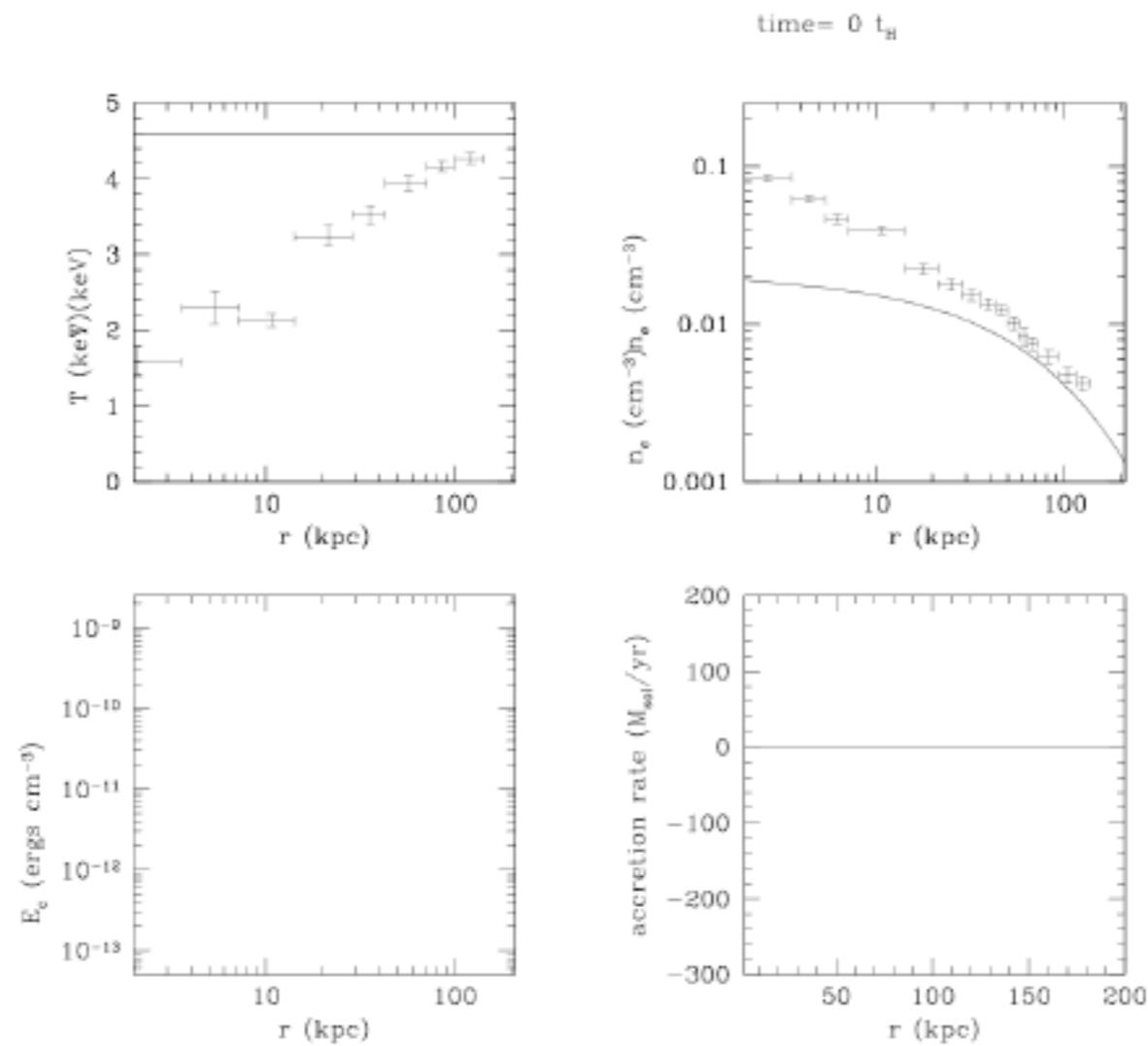
Slope is free parameter, implicitly specifies bubble disruption rate

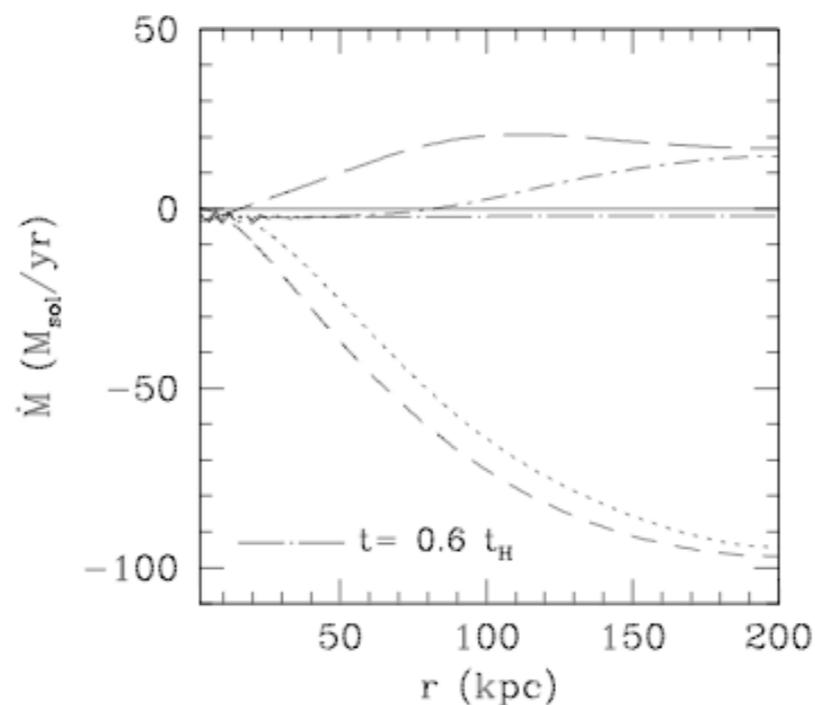
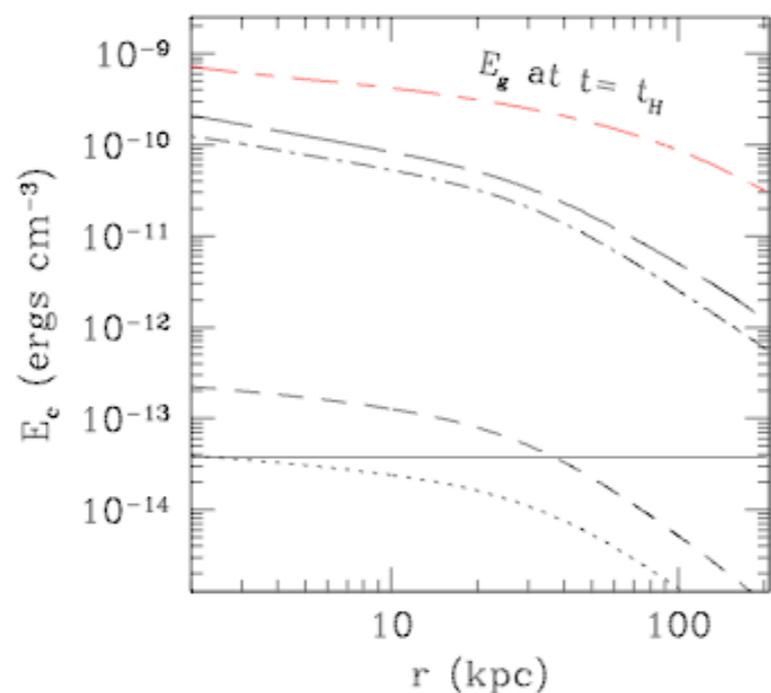
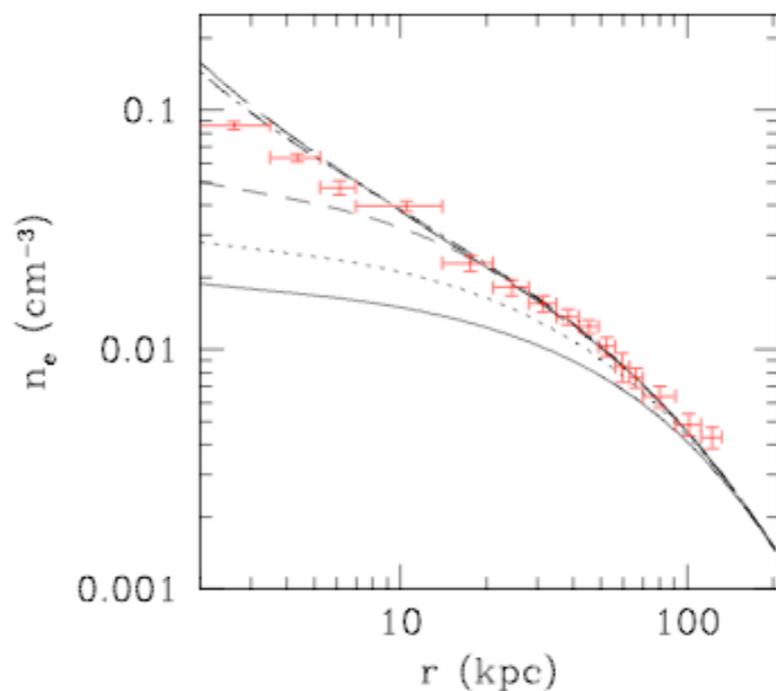
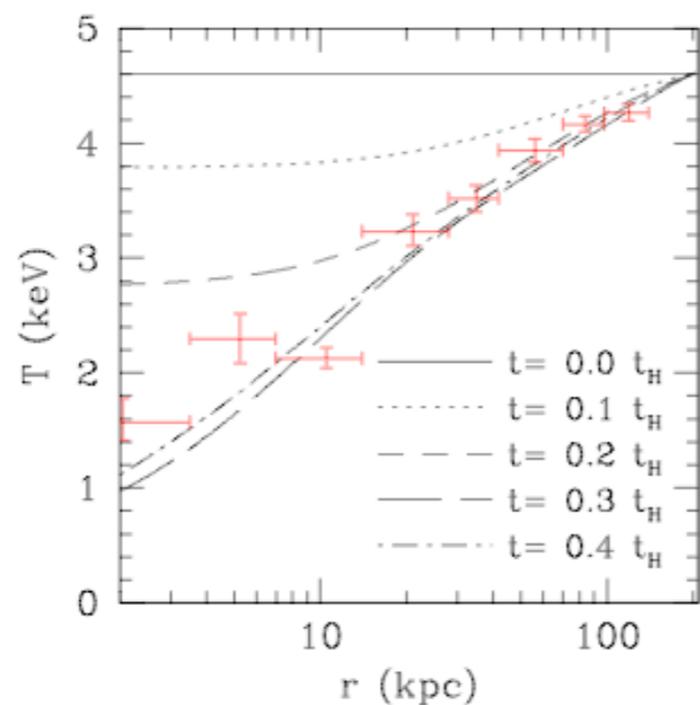
□ Amount of energy lost to pdv work is small, at most comparable to the bubble disruption rate



(pdv work can also heat ICM, we ignore it)

# Bottom line: it works!

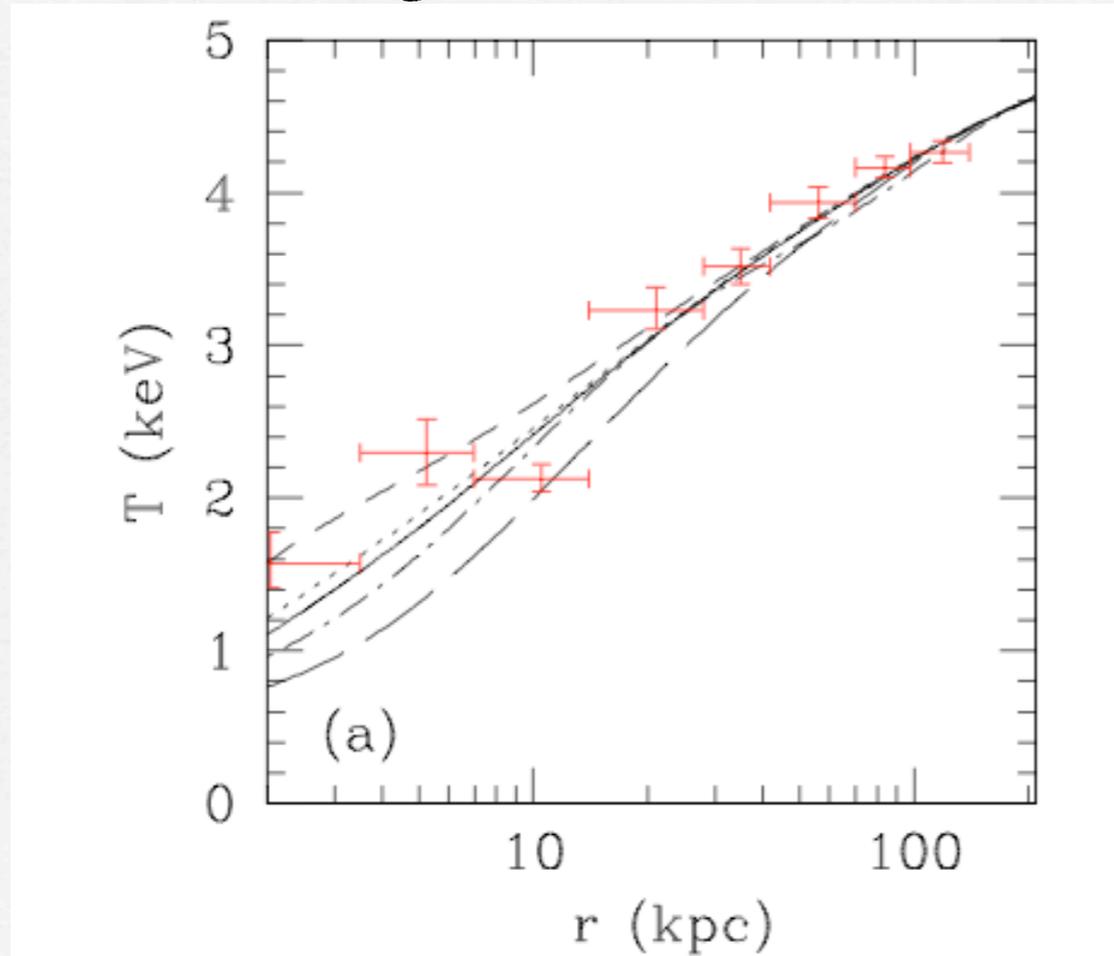




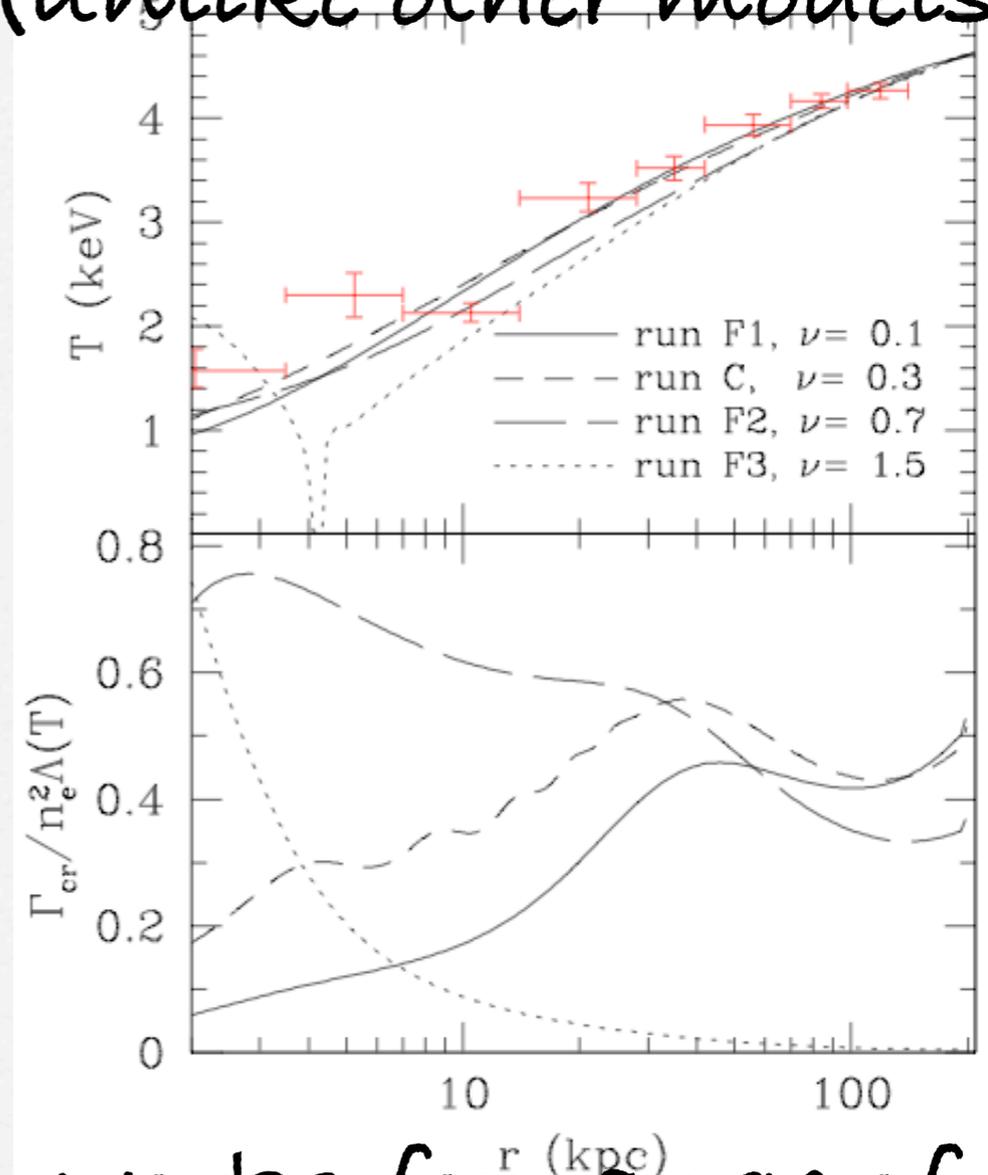
Note that CR pressure is much less than thermal pressure

# No fine tuning

works (i.e., no massive cooling flow) starting from arbitrary initial conditions (unlike other models...)



works for range of AGN +  
conduction parameters

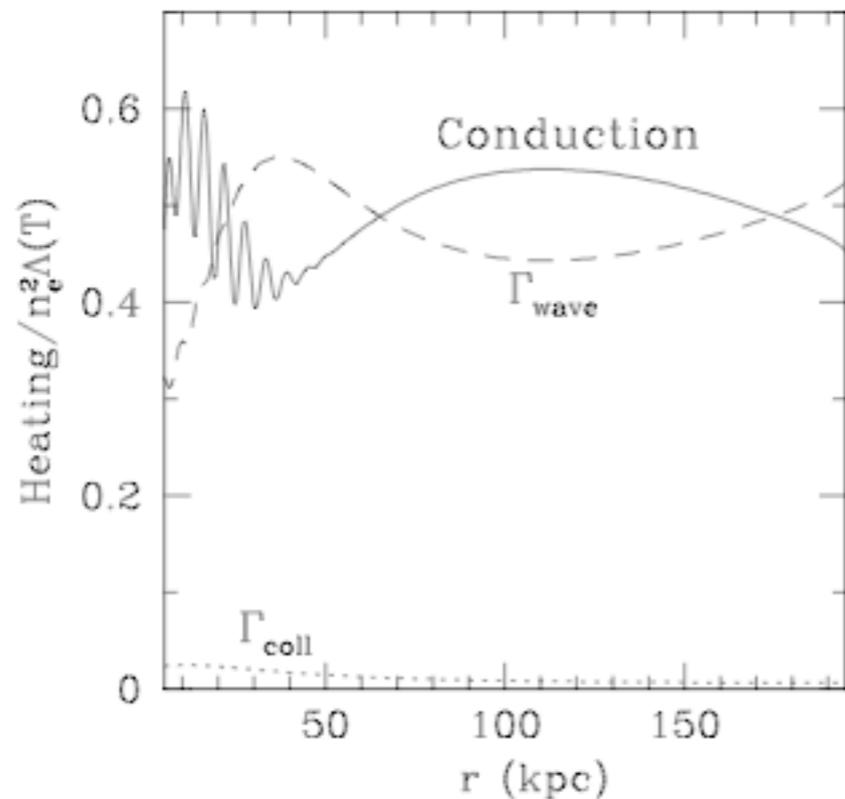
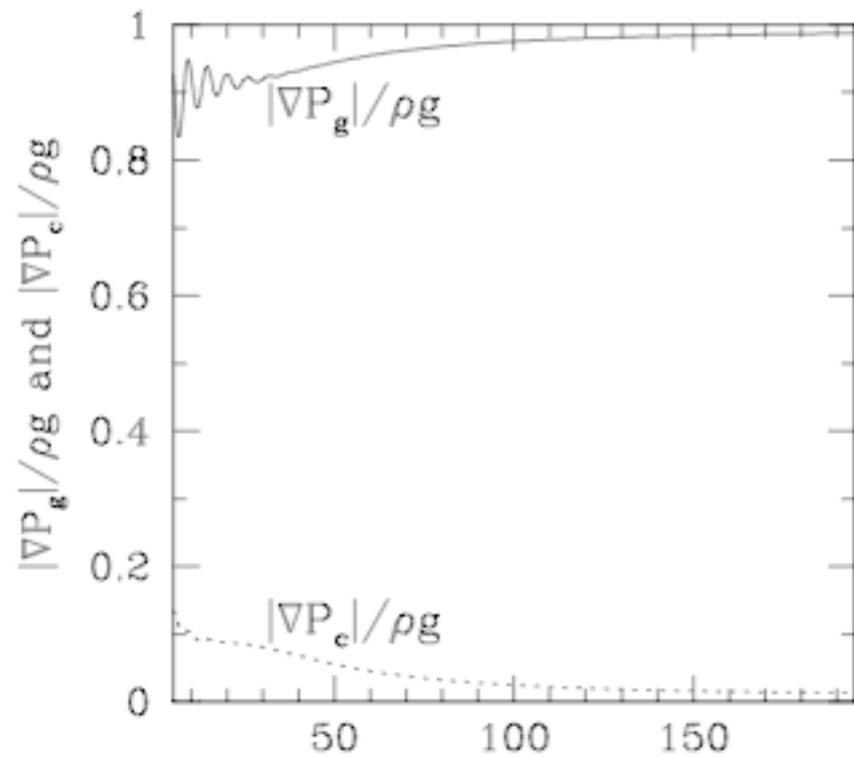


works for range of CR  
profiles

# Required CR pressure gradients

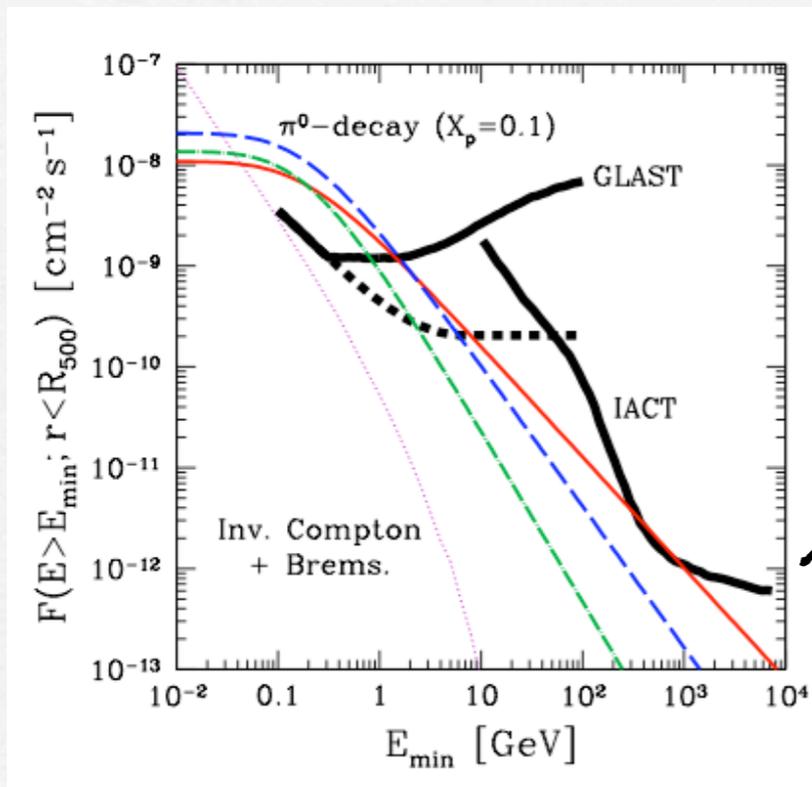
OK

Small fraction of thermal  
pressure gradient



most heating is wave heating

# Observational tests

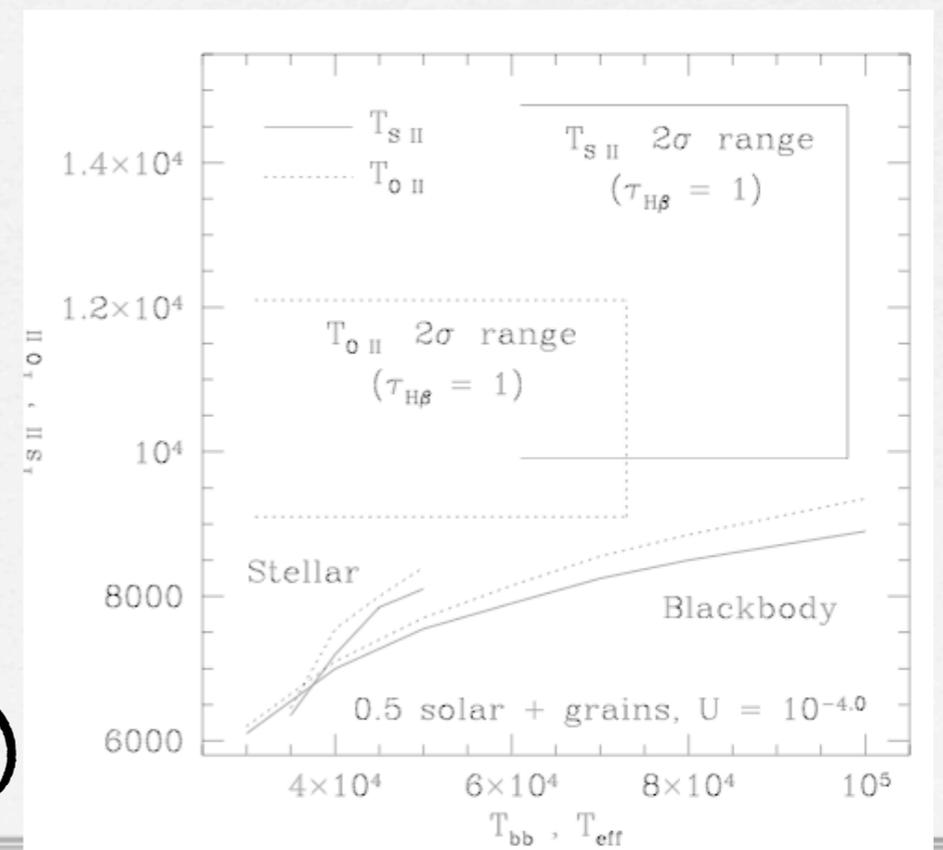


See gamma-rays from pion-decay with GLAST

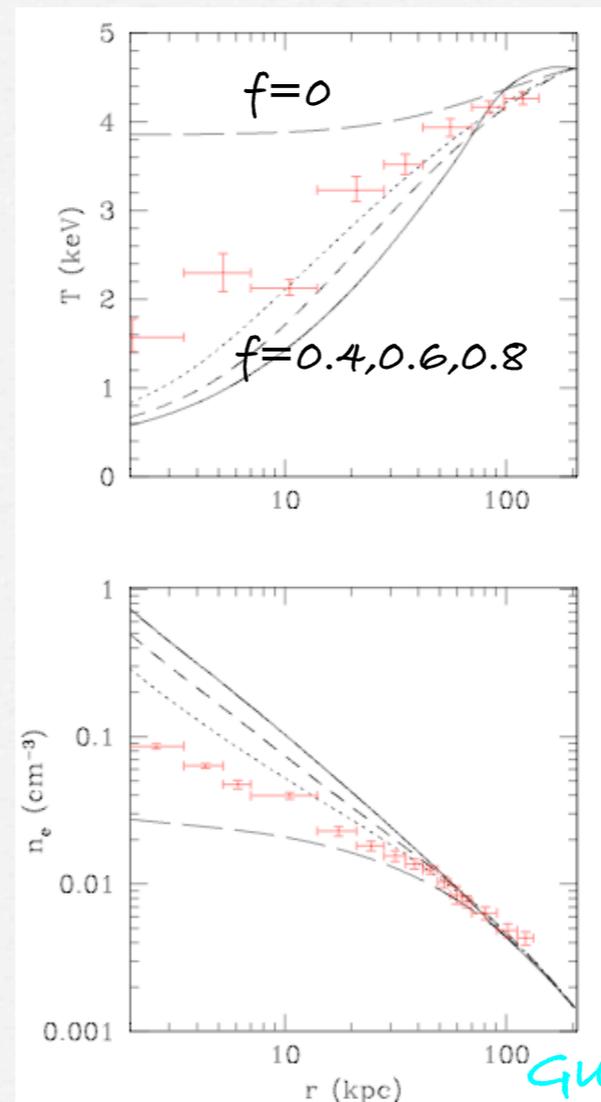
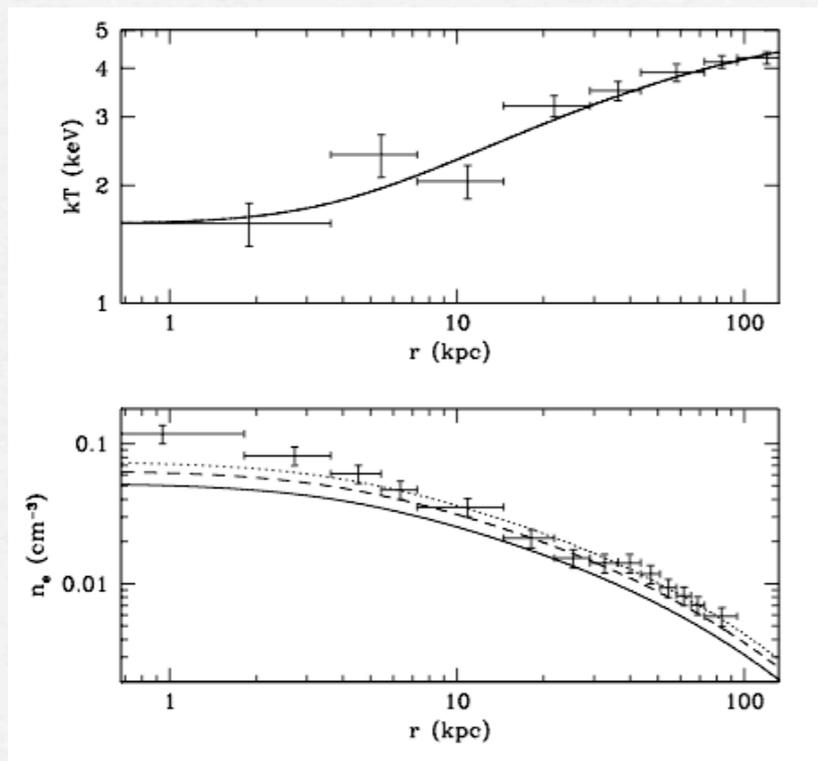
Ando & Nagai (2007)

Optical filaments: need source of anomalous heating?

Voit & Donahue (1997)



# Let's look more closely at fine-tuning issues for conduction models..



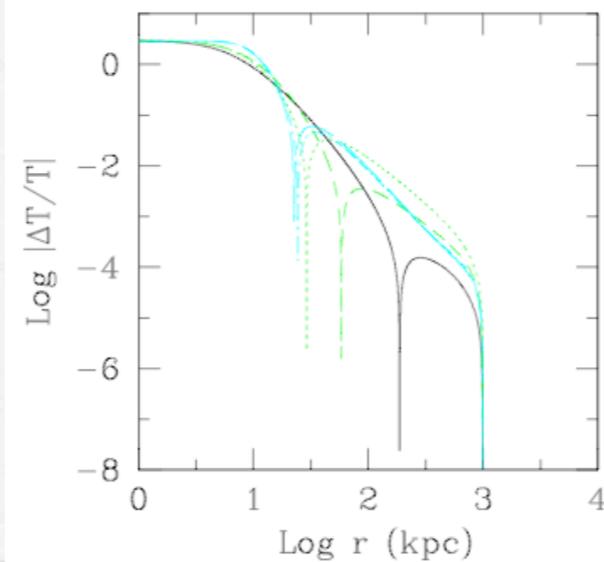
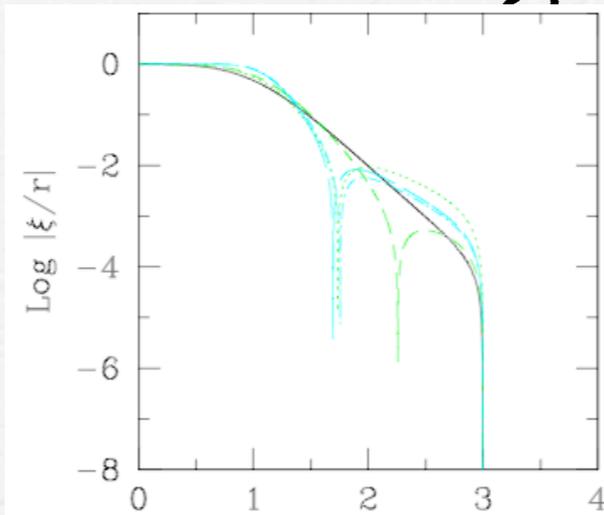
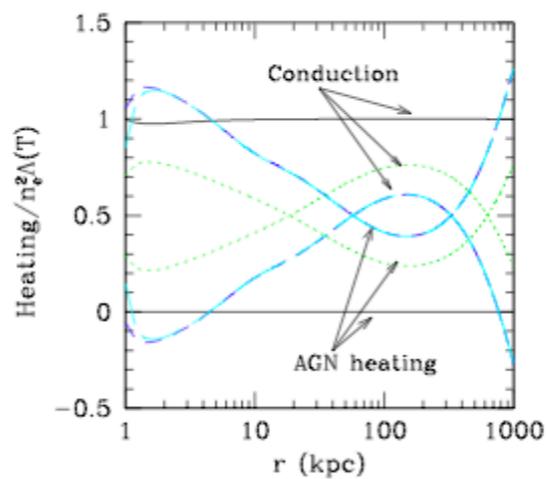
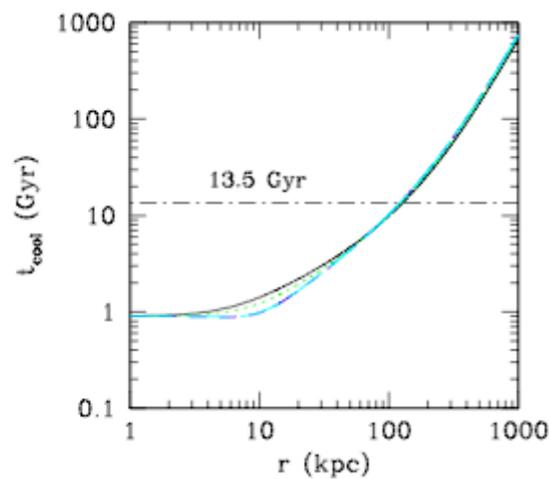
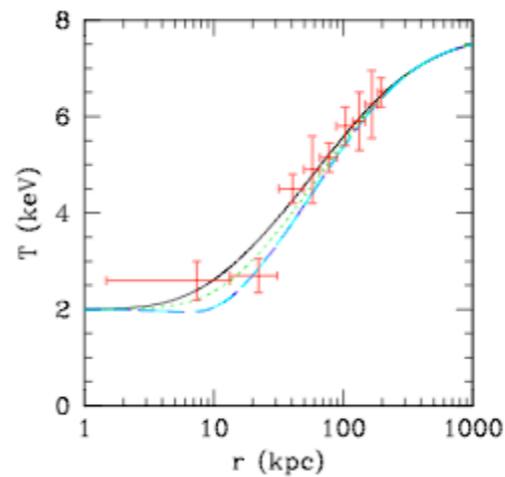
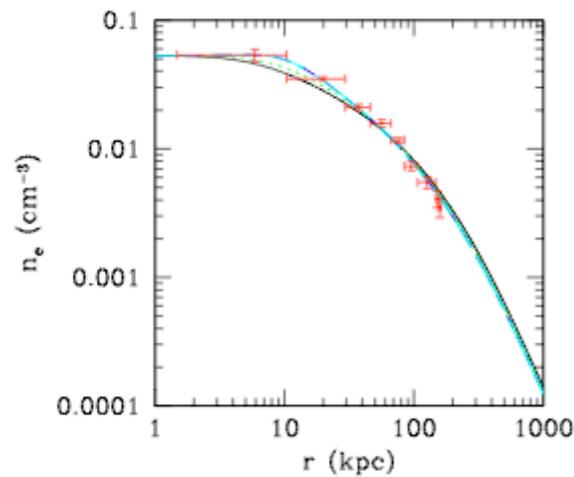
But it won't evolve toward this state in general...

can have equilibrium model which fits observations (solve eigenvalue problem)

# Global Stability Analysis

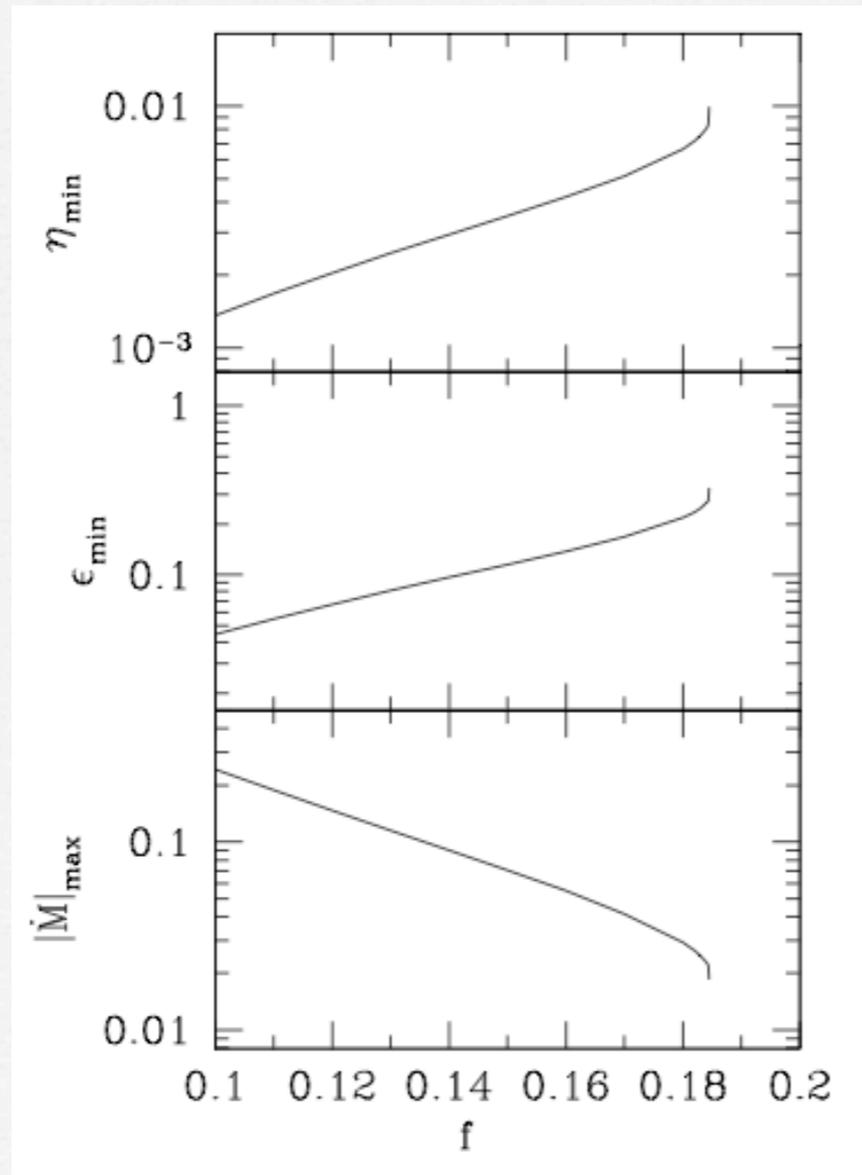
Guo, Oh & Ruszkowski 2007, in prep

Perform Lagrangian global stability analysis



Find  
unstable  
eigenmodes

# Global Unstable modes are suppressed with AGN!



(note: this never happens for conduction only model)

useful tool for analysing models w/out sims