

Shape of Halo Bispectrum from Non-Gaussian Initial Conditions in Cosmological N-body Simulations

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Local-type NG and Large Scale Structure

local-type primordial Non-Gaussianity

$$\Phi(\mathbf{x}) = \phi(\mathbf{x}) + f_{\text{NL}}[\phi^2(\mathbf{x}) - \langle \phi^2(\mathbf{x}) \rangle]$$

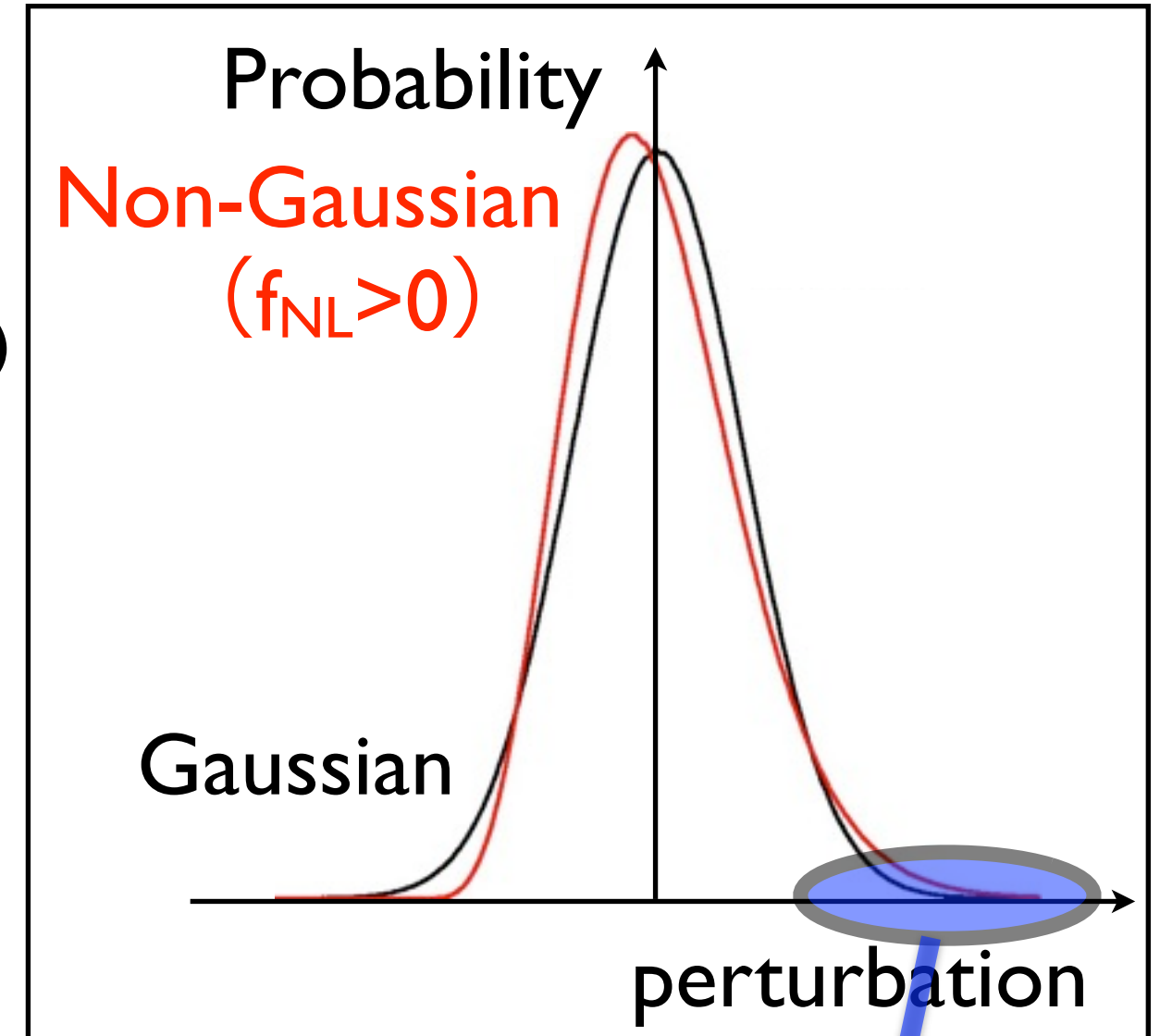
Gaussian $\nearrow$ $-10 < f_{\text{NL}} < 74$ (95% CL) WMAP7(Komatsu+10)

We observe ``galaxy”,
not ``matter” density fluctuations

biasing changes things dramatically!

- halo mass function
- halo power spectrum
- halo bispectrum

: **Understanding of the halo/galaxy
biasing is the key!**



change is relatively
large at the high
density tail!!
→ halos (galaxies)

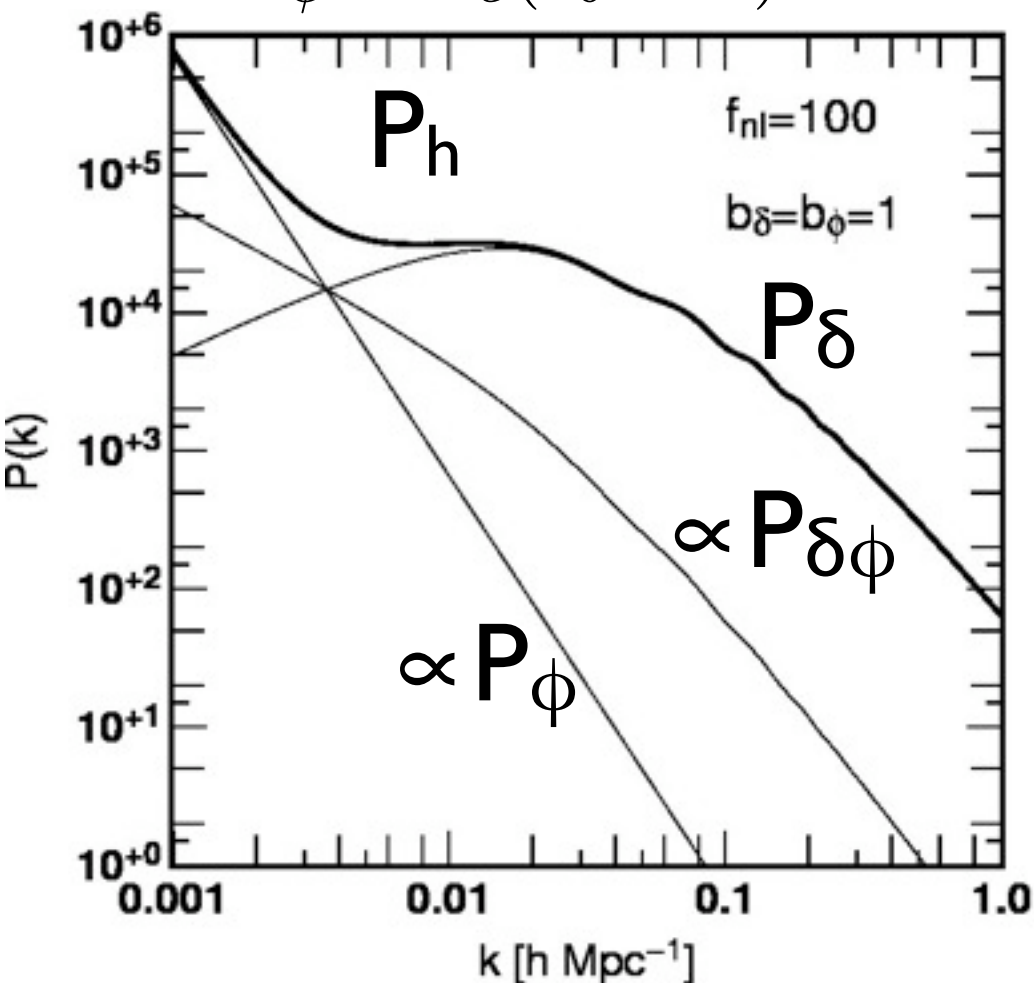
Scale-dependent bias

Theory

mass density curvature

$$\delta_{\vec{x}}^h = b_{\delta} \delta_{\vec{x}} + 2f_{\text{nl}} b_{\phi} \phi_{\vec{x}}$$

$$b_{\phi} = \delta_c (b_{\delta} - 1)$$



peak-background split
and/or peak(local) bias

Dalal+07
Slosar+08
Matarrese&Verde08
Afshordi&Tolley08
McDonald08
Taruya+09
Giannantonio&Porciani10
and more ...

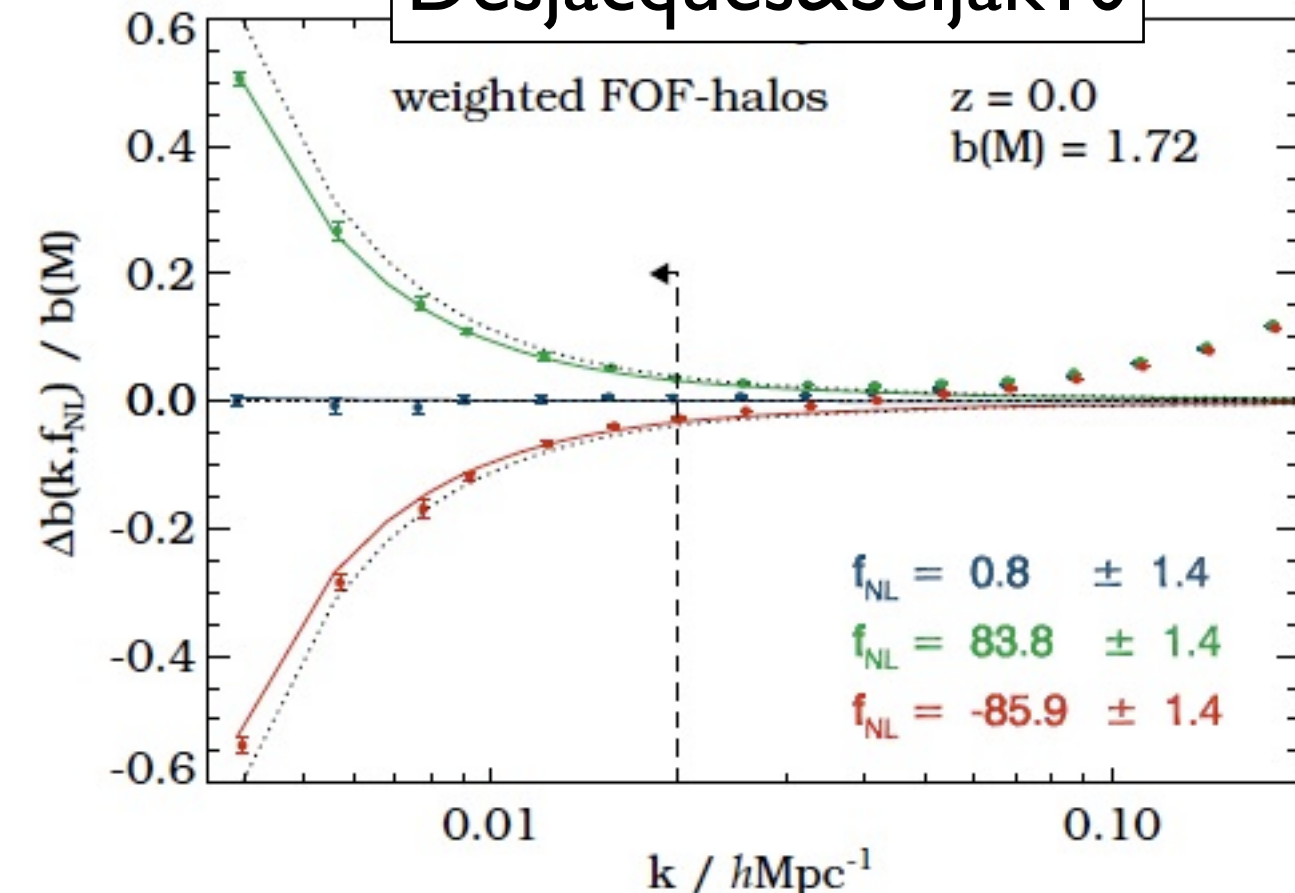
Slosar+08

$$-29 < f_{\text{nl}} < 69 \text{ (QSOs+more)}$$

calibration by simulations

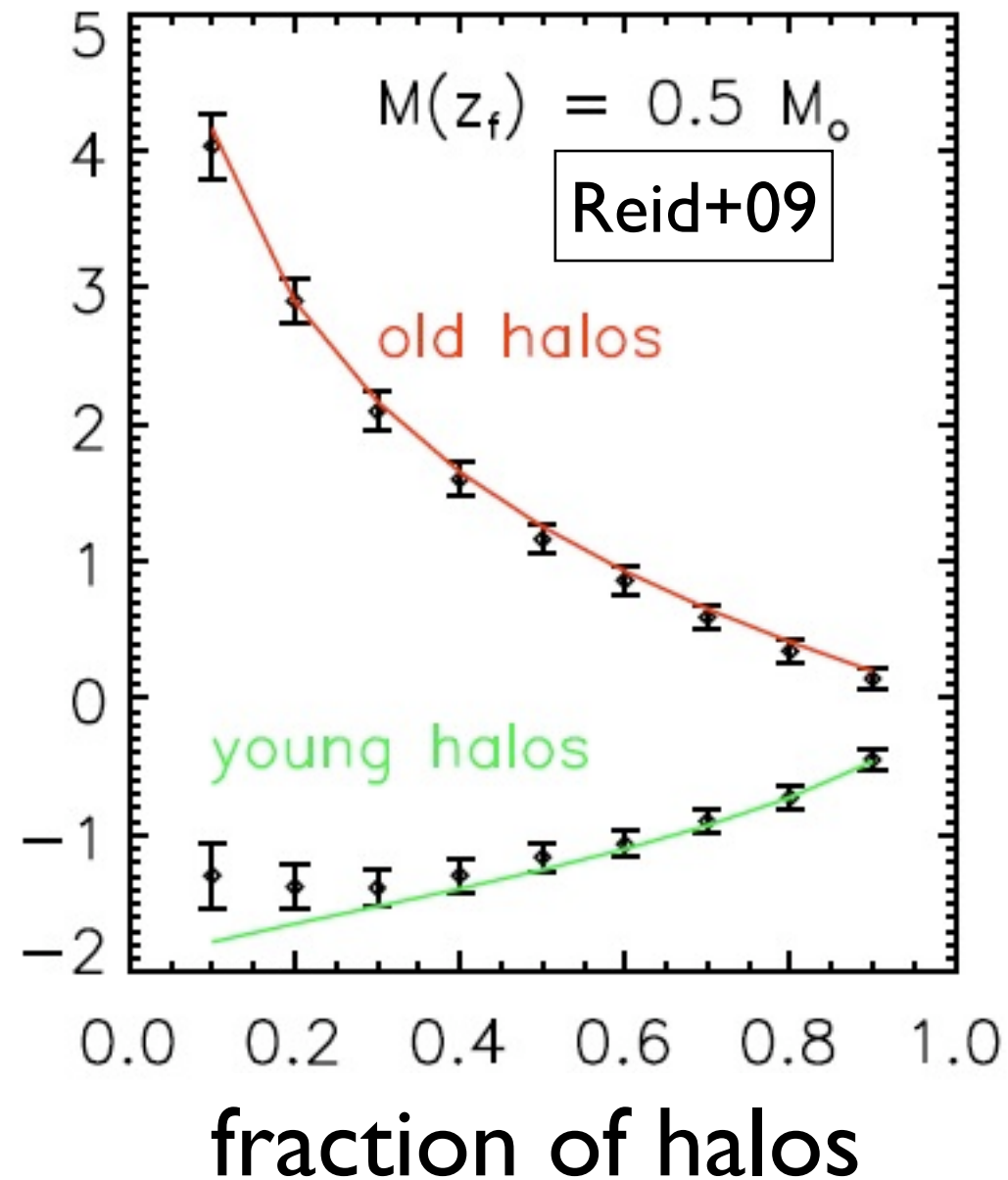
Desjacques+09
Grossi+09
Pillepich+10

Desjacques&Seljak10

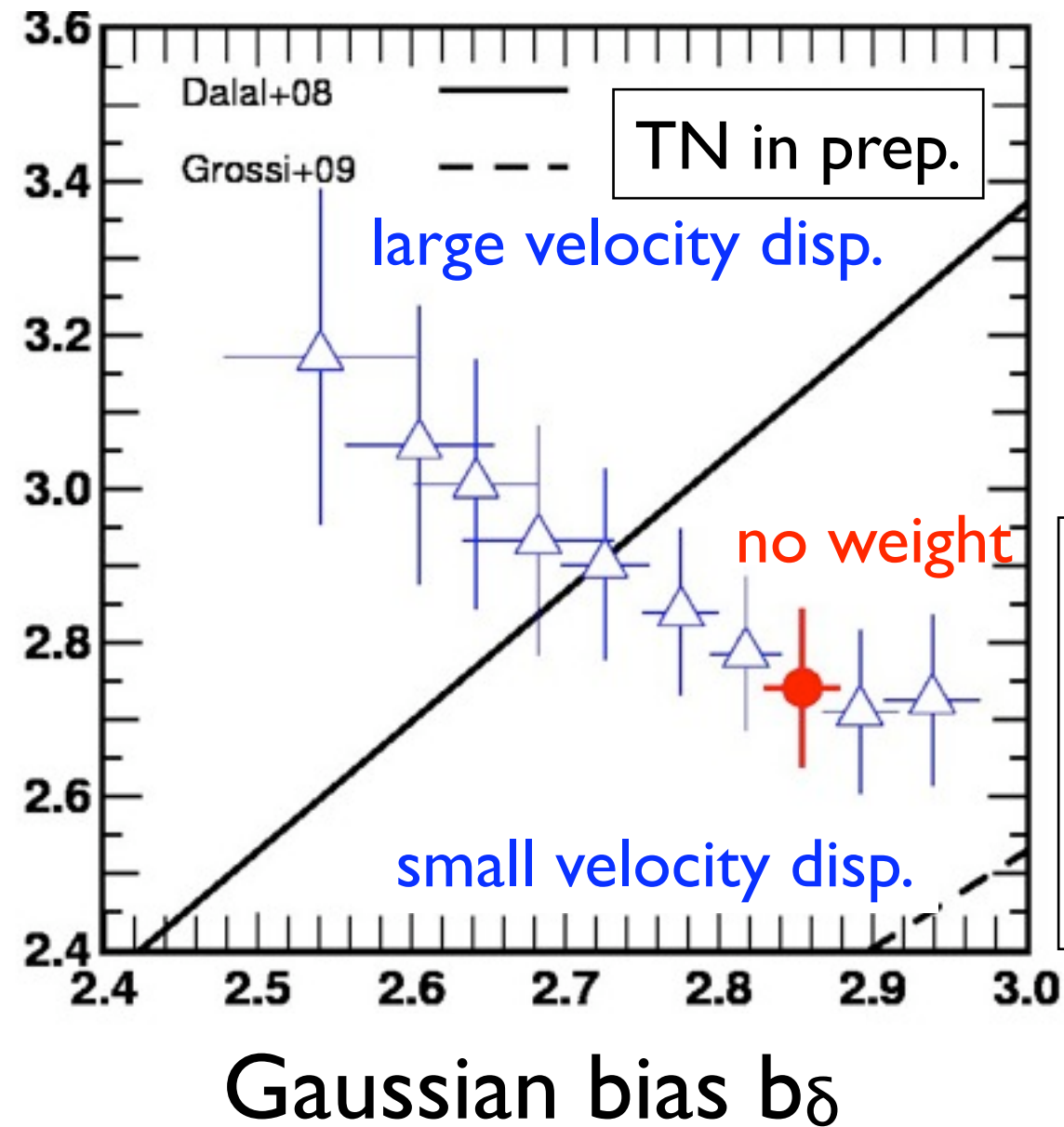


Halo assembly bias

non-Gaussian correction



non-Gaussian bias b_ϕ



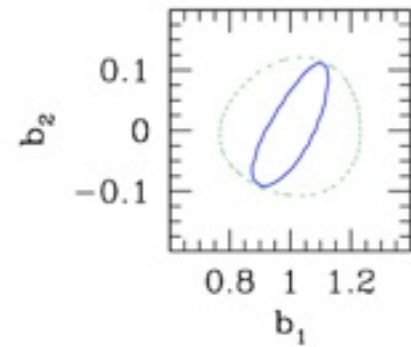
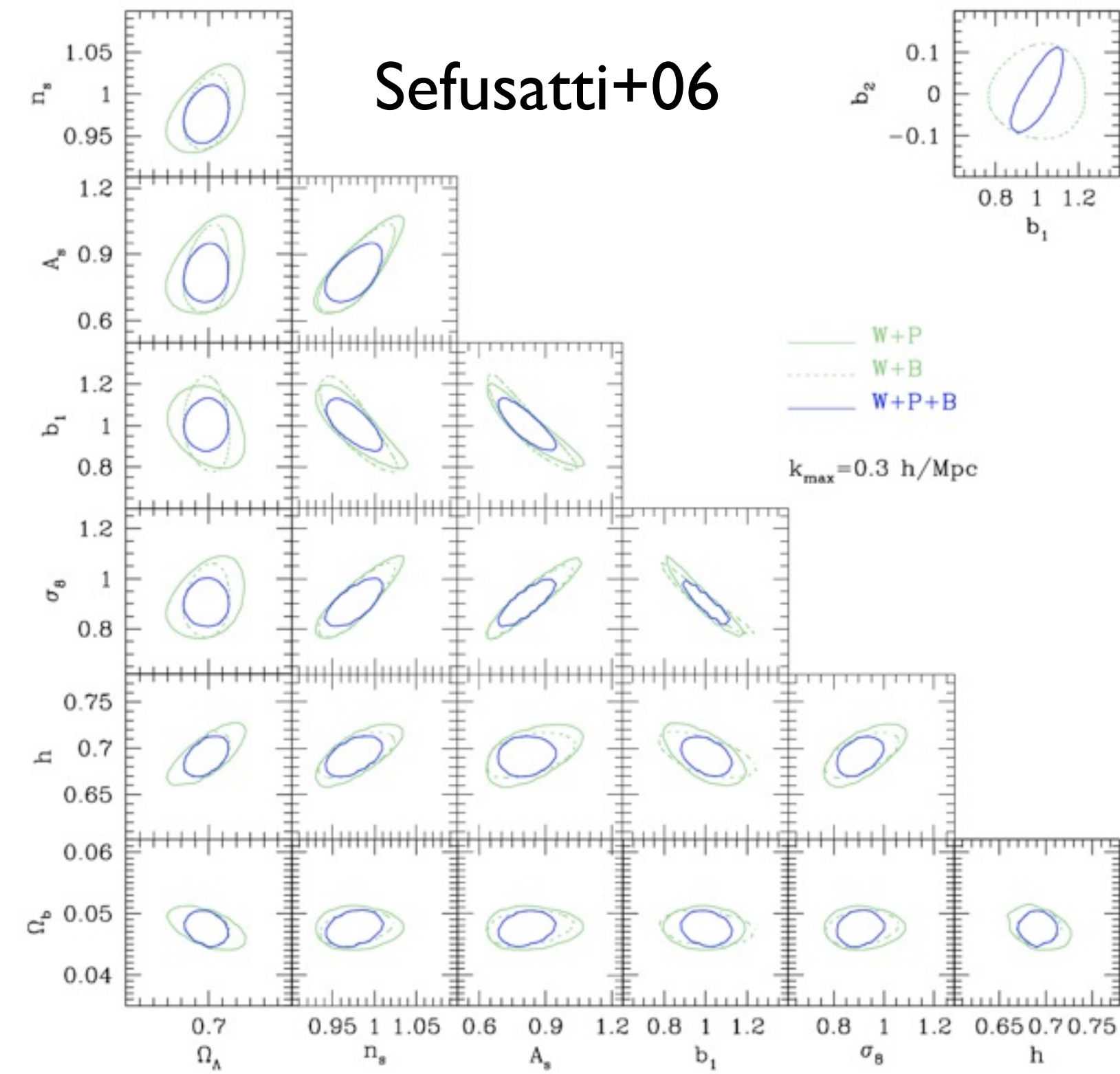
$$\delta_{\vec{x}}^h = b_\delta \delta_{\vec{x}} + 2f_{nl} b_\phi \phi_{\vec{x}}$$

usually $b_\phi = \delta_c(b_\delta - 1)$

assembly bias can change the amplitude of non-G correction for given Gaussian bias

Bispectrum?

Sefusatti+06



adds information

contains primordial non-G
information by definition

B_h vs. B_m

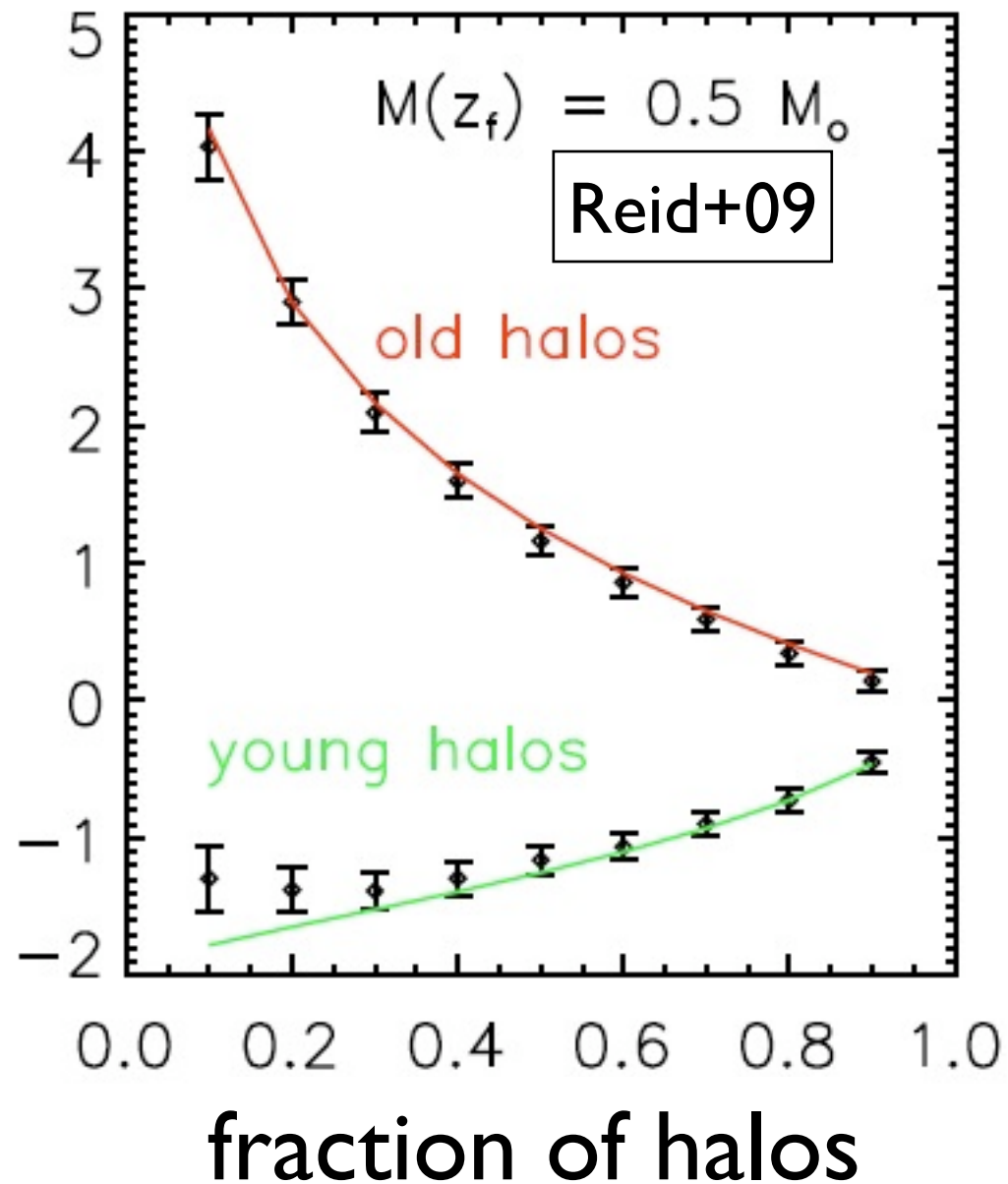
“scale-dependent bias” should exist

$$\begin{aligned}
 & B_h(\mathbf{k}_1, \mathbf{k}_2, \mathbf{k}_3) \quad \text{local bias} \quad \text{Jeong\&Komatsu09} \\
 & = b_1^3 \left[B_R(\mathbf{k}_1, \mathbf{k}_2, \mathbf{k}_3) + \frac{b_2}{b_1} \{ P_R(k_1) P_R(k_2) + (2 \text{ cyclic}) \} \right. \\
 & \quad \left. + \frac{1}{2} \frac{b_2}{b_1} \int \frac{d^3 q}{(2\pi)^3} T_R(\mathbf{q}, \mathbf{k}_1 - \mathbf{q}, \mathbf{k}_2, \mathbf{k}_3) + (2 \text{ cyclic}) \right]. (13)
 \end{aligned}$$

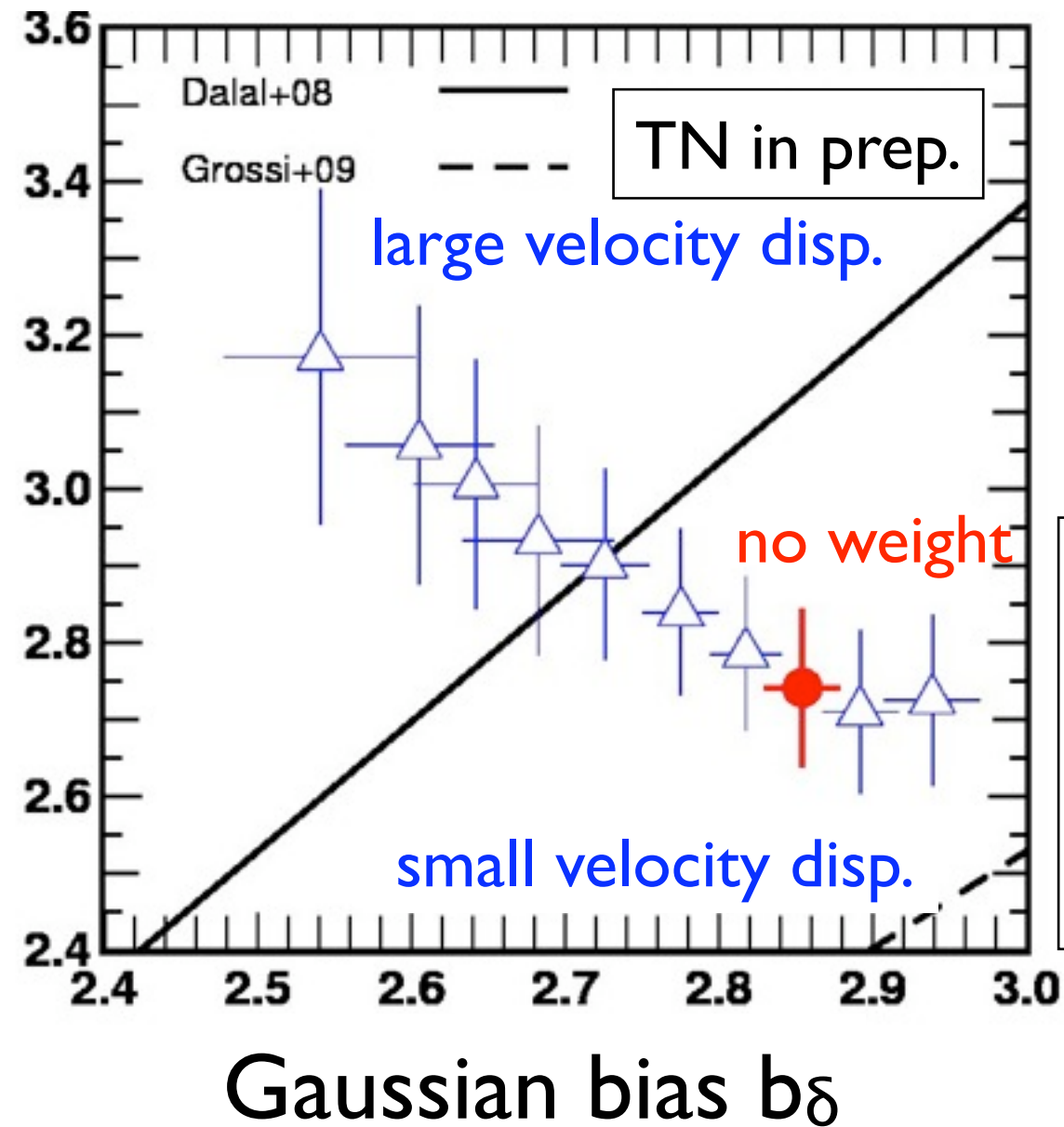
but not investigated numerically

Halo assembly bias

non-Gaussian correction



non-Gaussian bias b_ϕ



$$\delta_{\vec{x}}^h = b_\delta \delta_{\vec{x}} + 2f_{nl} b_\phi \phi_{\vec{x}}$$

usually $b_\phi = \delta_c(b_\delta - 1)$

assembly bias can change the amplitude of non-G correction for given Gaussian bias

Bispectrum can break the degeneracy: $\Delta P \propto f_{nl} b_\phi$, $\Delta B \propto f_{nl}^2 b_\phi$

N-body simulations with f_{NL}

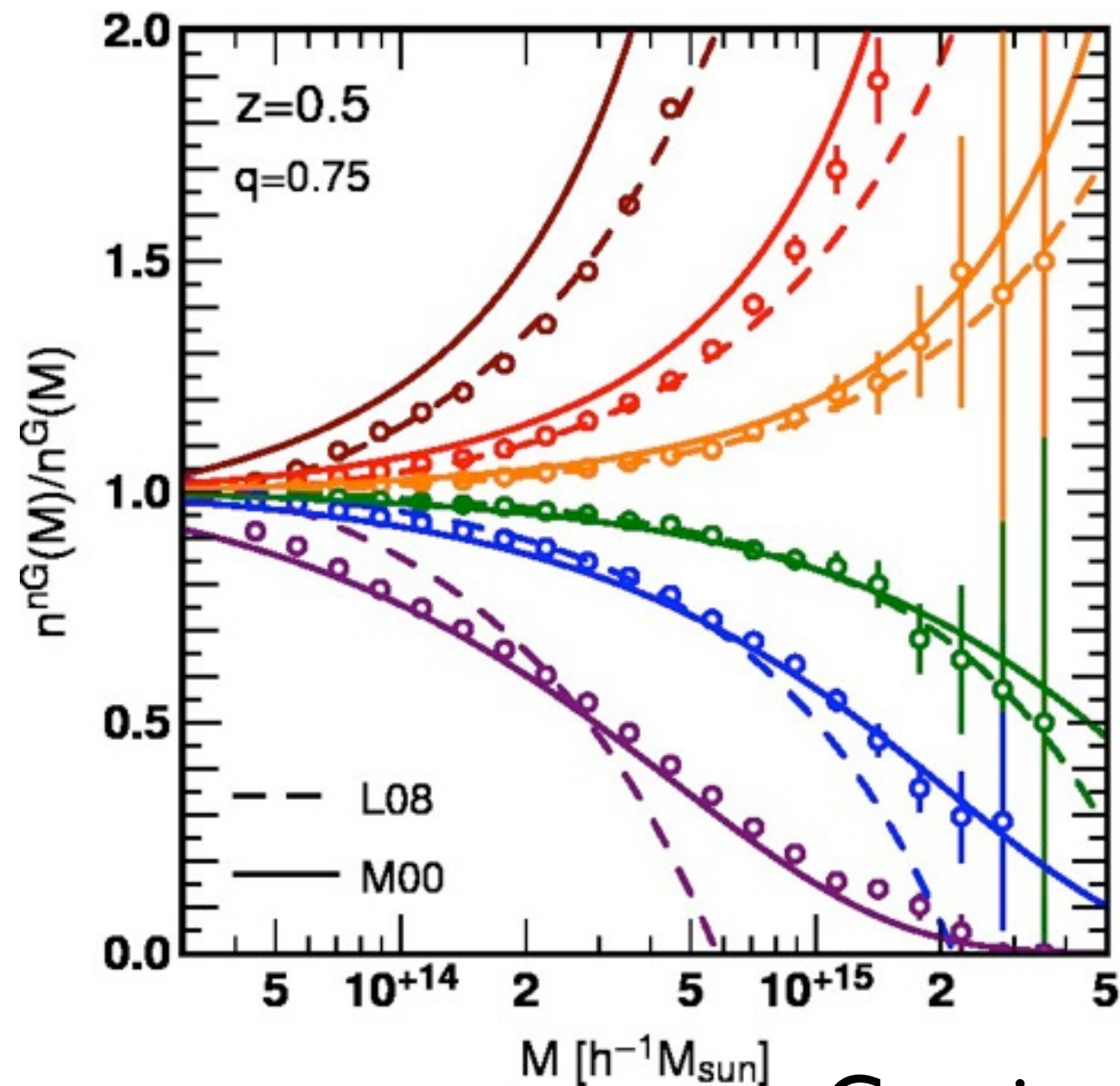
Author	f_{NL}	volume/run [$h^{-3}\text{Gpc}^3$]	run/model	particle/run	mass/particle [$h^{-1}M_{\text{sun}}$]
Kang+07	-580,134	0.027	1 (2)	128^3	$\sim 10^{12}$
Grossi+07	$0, \pm 100, \pm 500, \pm 1000$	0.13	1	800^3	2.0×10^{10}
Dalal+08	$0, \pm 5, \pm 50, \pm 500$ (± 5000)	0.51	1	512^3	2.5×10^{11}
Pillepich+10	$0, \pm 27, \pm 80, 250, 500, 750$	1.7	1 (2)	1024^3	1.2×10^{11}
Desjacques+09	$0, \pm 100$ ($\pm 10, \pm 30$)	4.1	5 (2)	1024^3	3.0×10^{11}
Grossi+09	$0, \pm 100, \pm 200$	1.7	1	960^3	1.4×10^{11}
This work	$0, \pm 100, \pm 300, \pm 1000$	8.0	20	512^3	4.6×10^{12}

examine the clustering of massive haloes at large scale

We need a sufficient number of modes to see bispectrum!

Our halo catalog

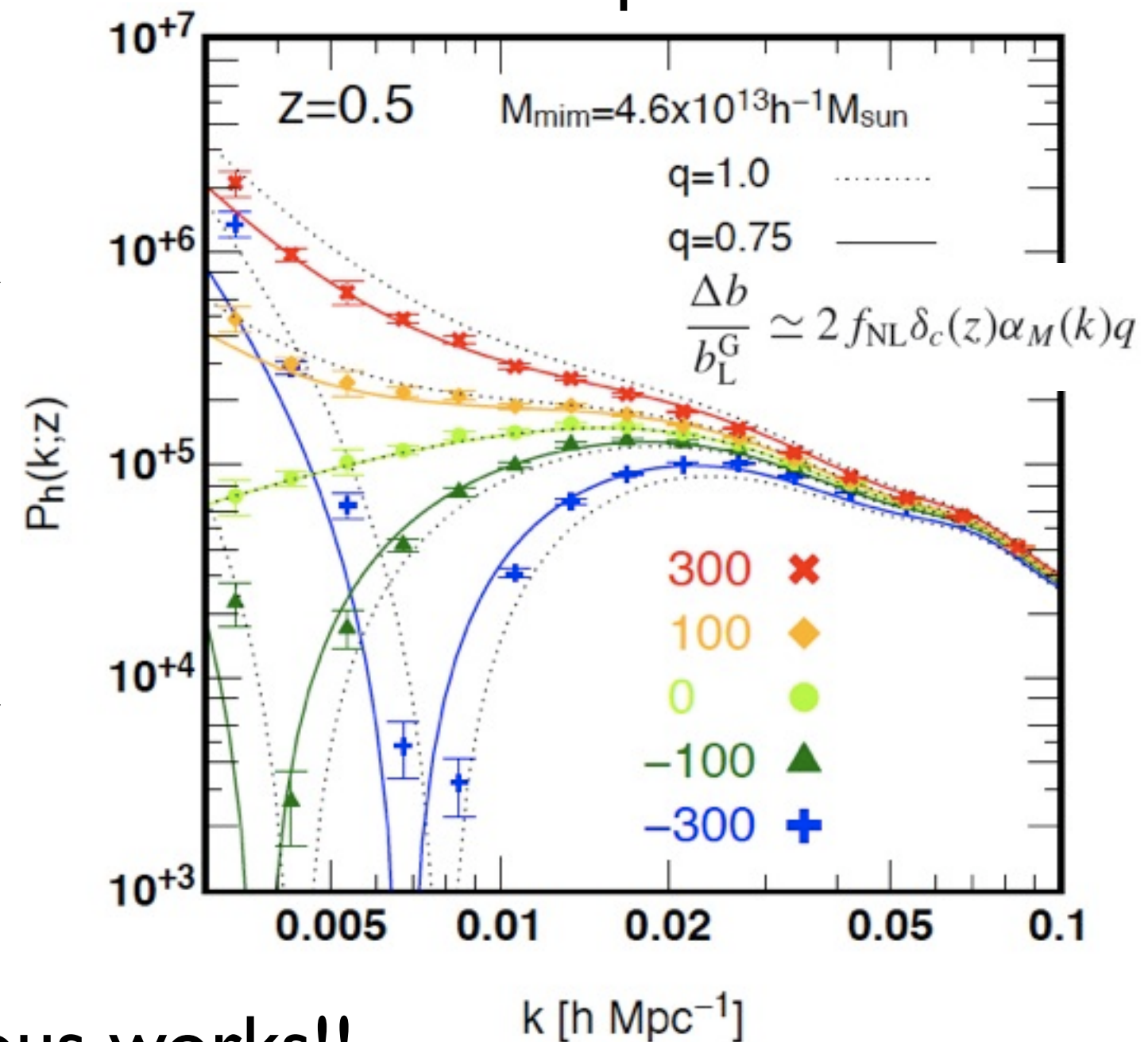
Mass function



$f_{\text{NL}} > 0$

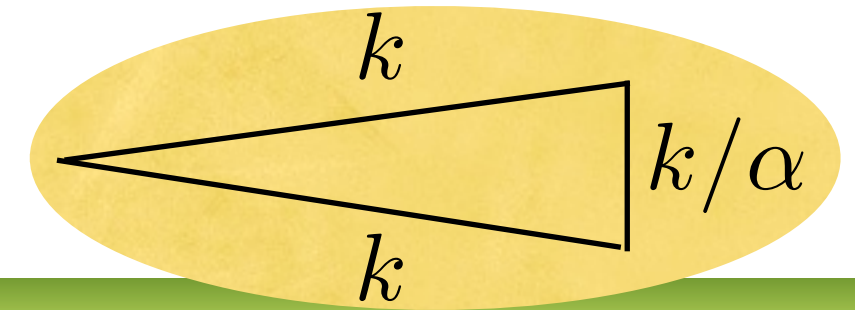
$f_{\text{NL}} < 0$

Power spectrum

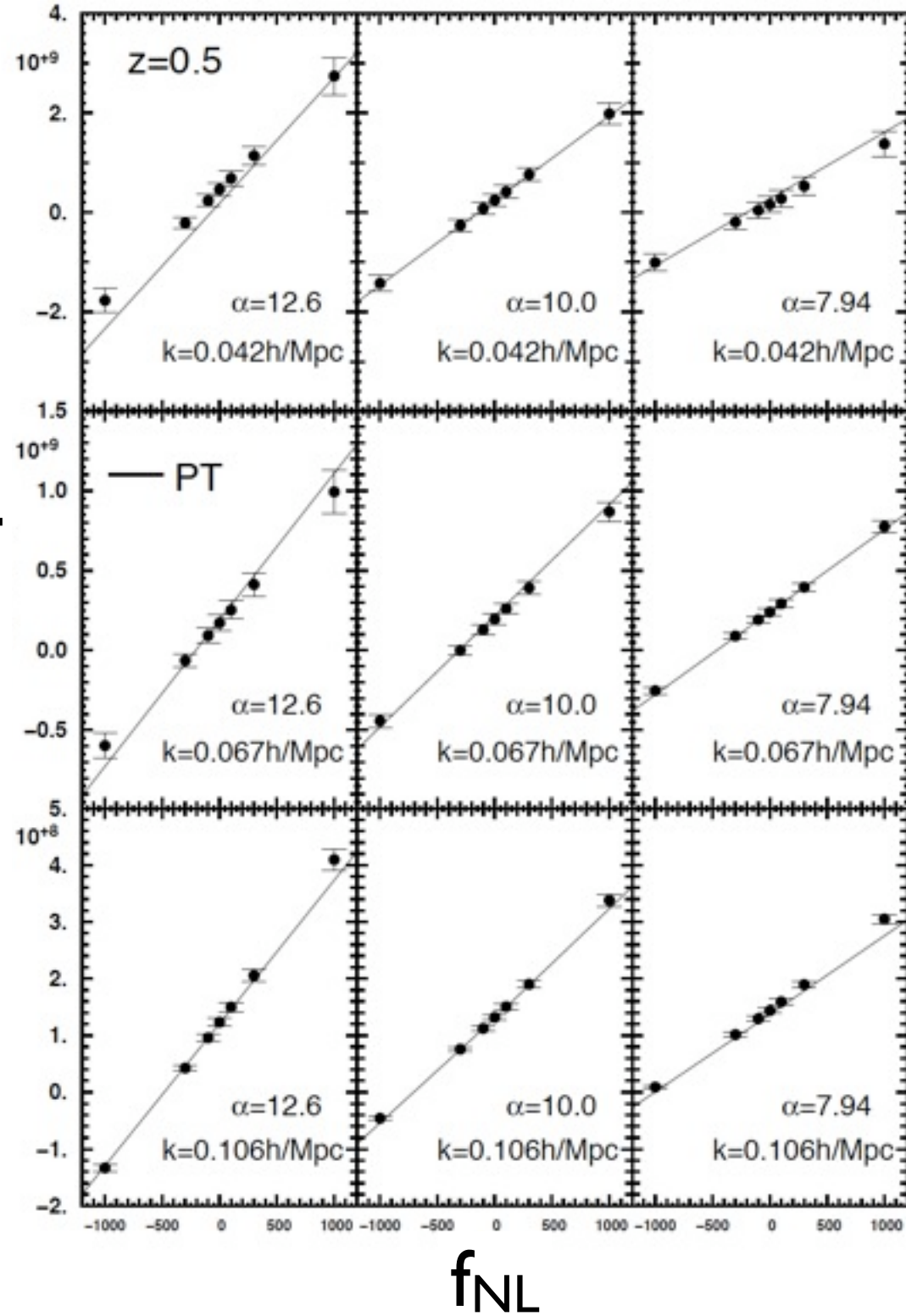


Consistent with previous works!!

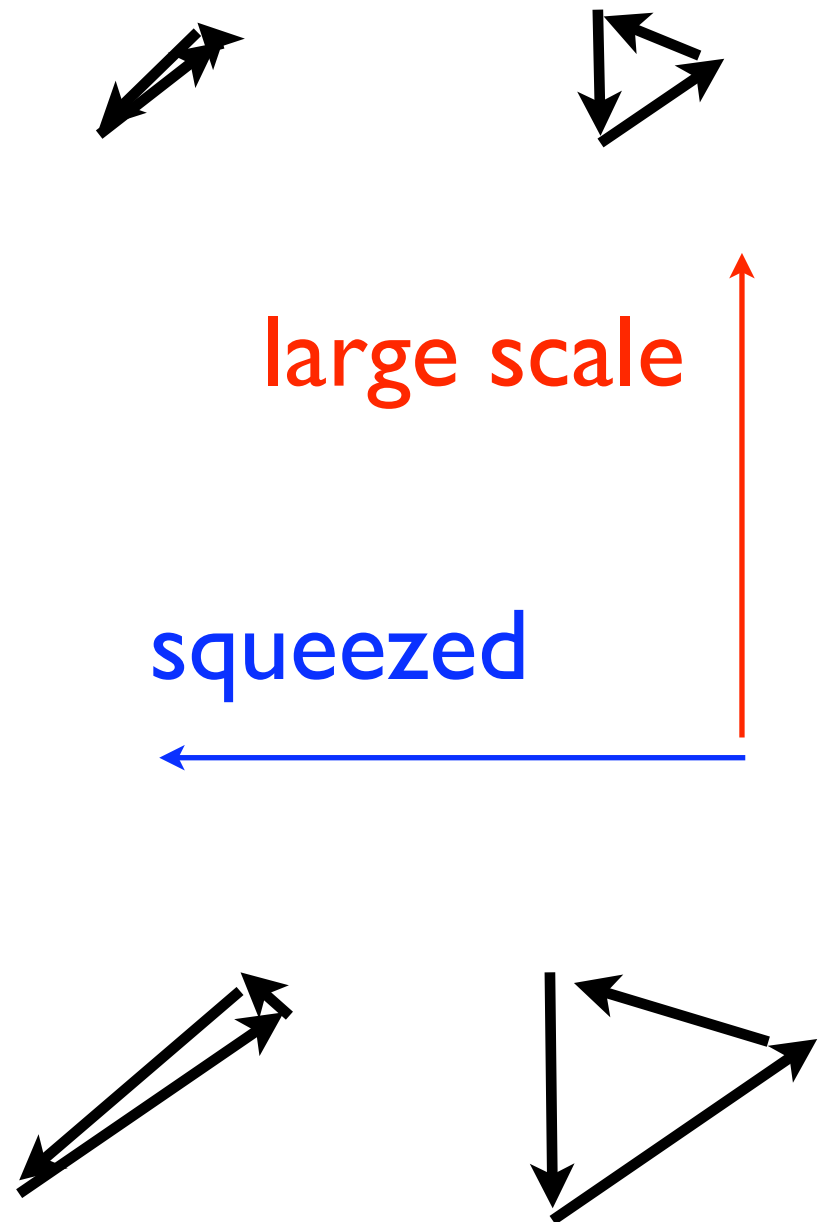
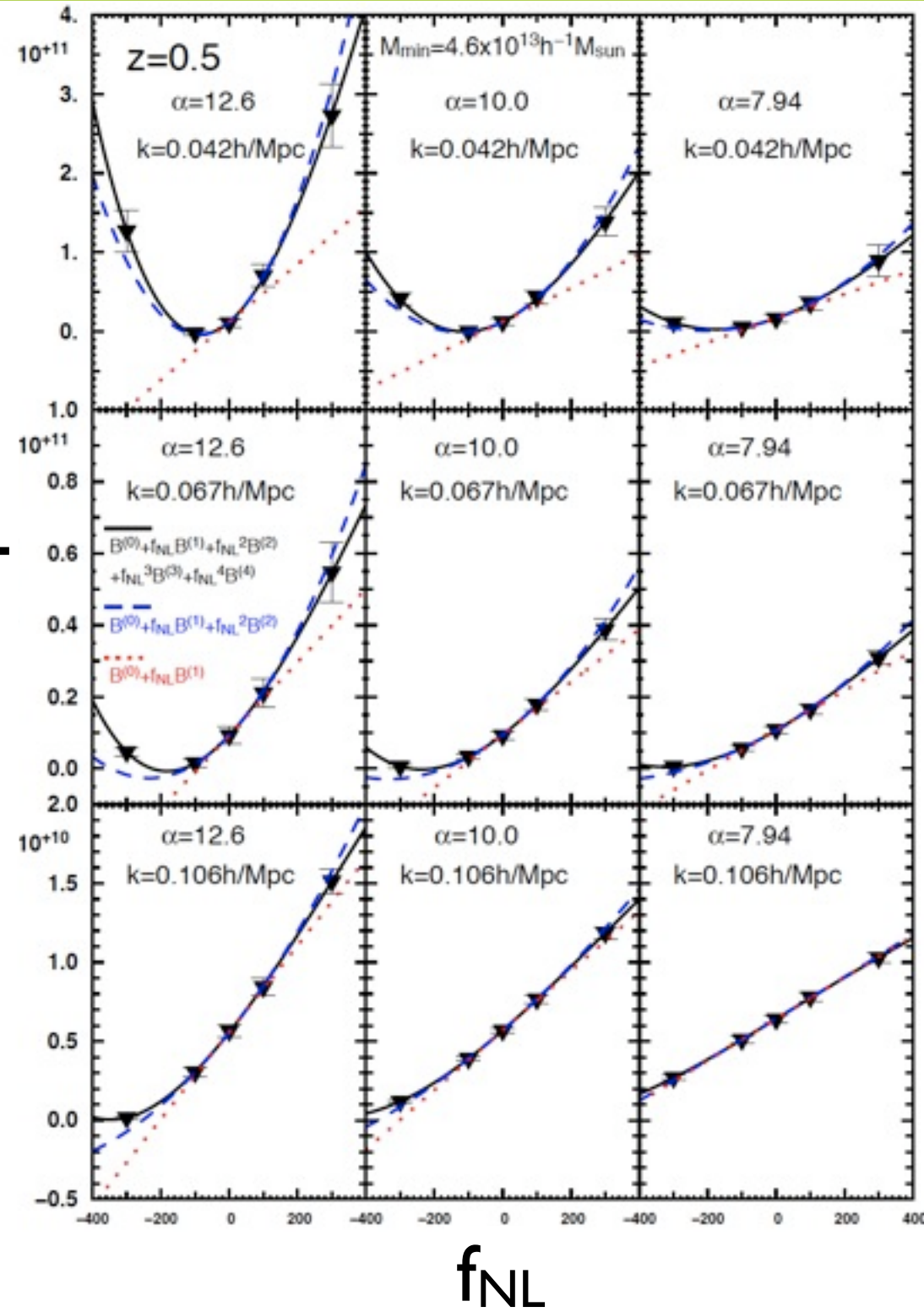
Result



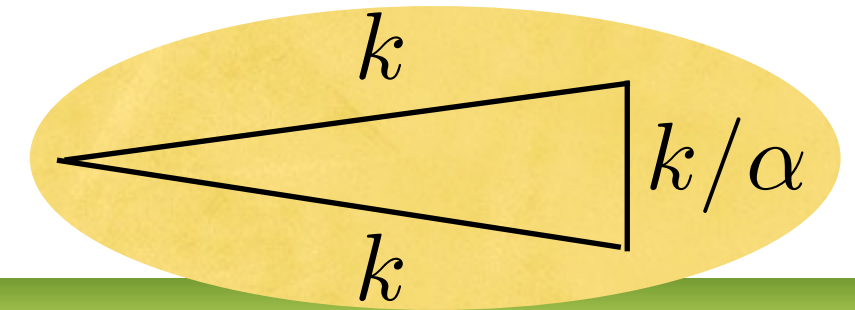
Matter bispectrum



Halo bispectrum



Result



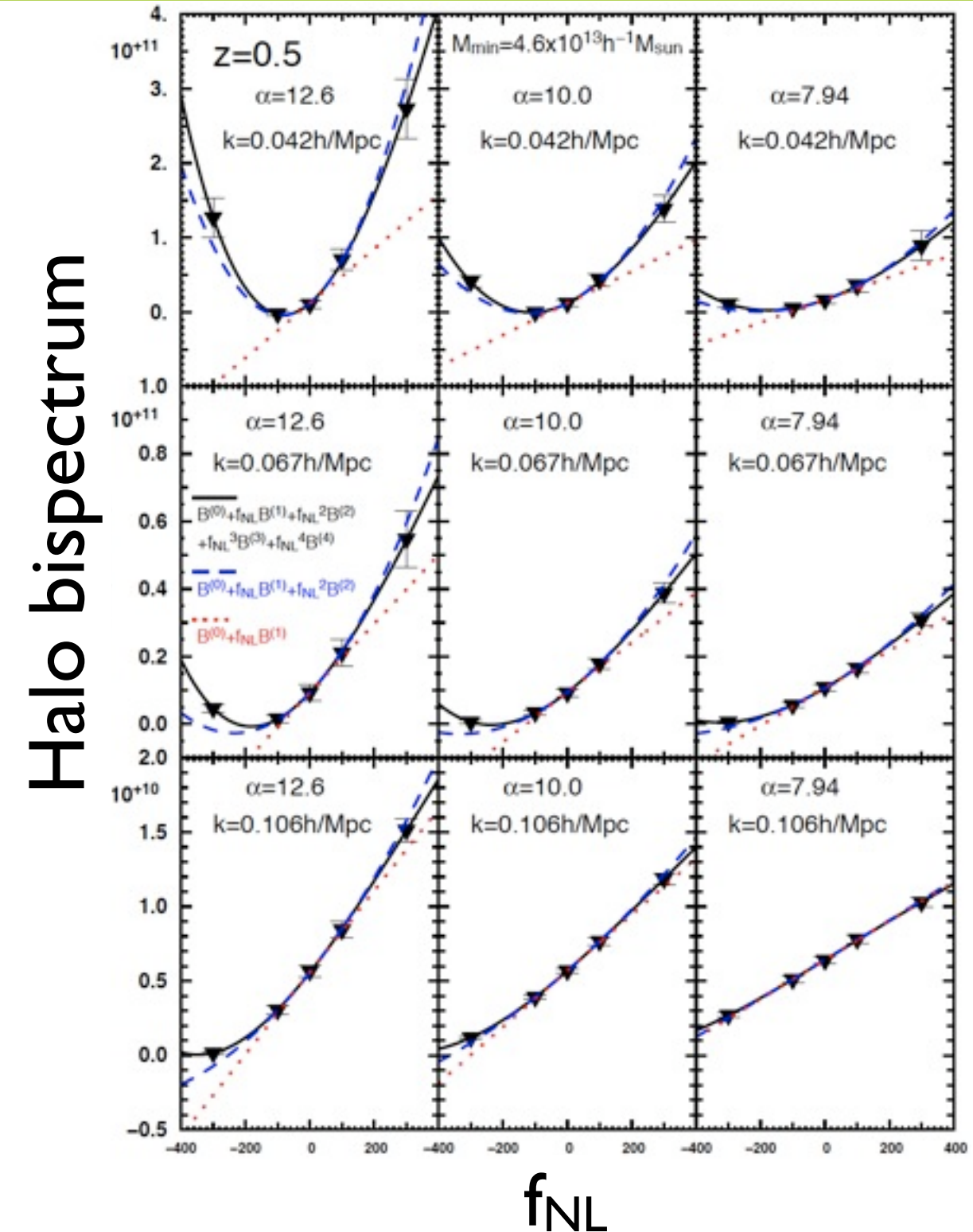
Halo bispectra at $z=0.5$ for isosceles triangles

Fitted by 4th-order polynomial:

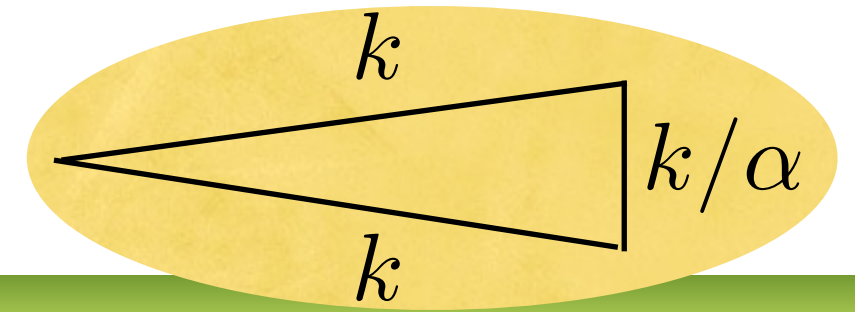
$$B^h(k, \alpha; f_{\text{NL}}) = B_0(k, \alpha) + f_{\text{NL}} B_1(k, \alpha) + f_{\text{NL}}^2 B_2(k, \alpha) + f_{\text{NL}}^3 B_3(k, \alpha) + f_{\text{NL}}^4 B_4(k, \alpha)$$

For reasonable values of f_{NL} ,

- B_0, B_1 & B_2 are enough to describe N-body data
- The new term (B_2) is important for squeezed triangles (small k & large α)



Comparison with theory



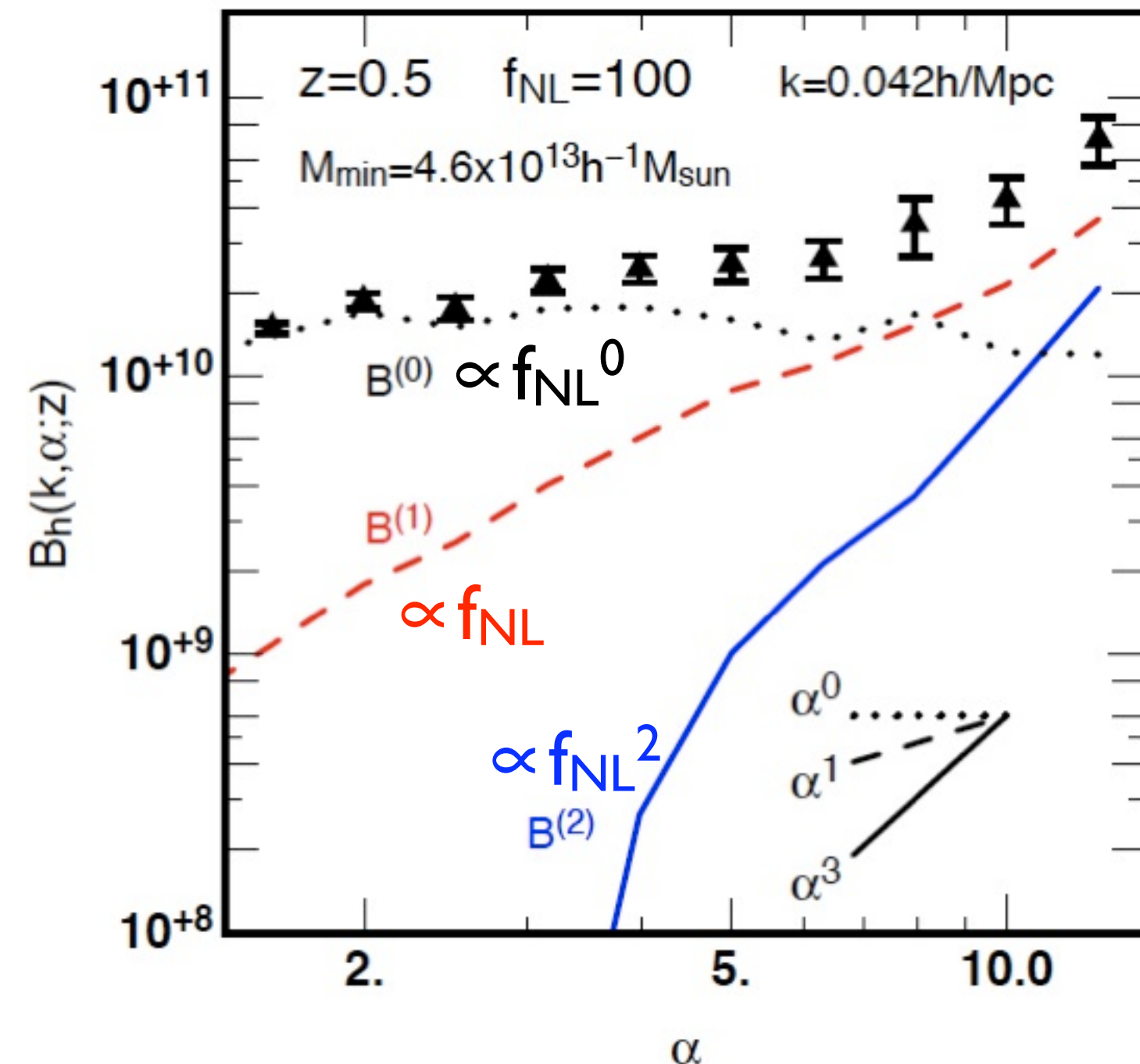
Assuming local bias model

$$\begin{aligned}
 & B_h(\mathbf{k}_1, \mathbf{k}_2, \mathbf{k}_3) \quad \text{Jeong, Komatsu09} \\
 &= b_1^3 \left[B_R(\mathbf{k}_1, \mathbf{k}_2, \mathbf{k}_3) + \frac{b_2}{b_1} \{ P_R(k_1) P_R(k_2) + (2 \text{ cyclic}) \} \right. \\
 & \quad \left. + \frac{1}{2} \frac{b_2}{b_1} \int \frac{d^3 q}{(2\pi)^3} T_R(\mathbf{q}, \mathbf{k}_1 - \mathbf{q}, \mathbf{k}_2, \mathbf{k}_3) + (2 \text{ cyclic}) \right]. \quad (13)
 \end{aligned}$$

Asymptotic behavior of halo bispectrum at “squeezed limit” ($\alpha \rightarrow \infty$)

$$\propto f_{\text{NL}}^2 \alpha^3 k^{-2}$$

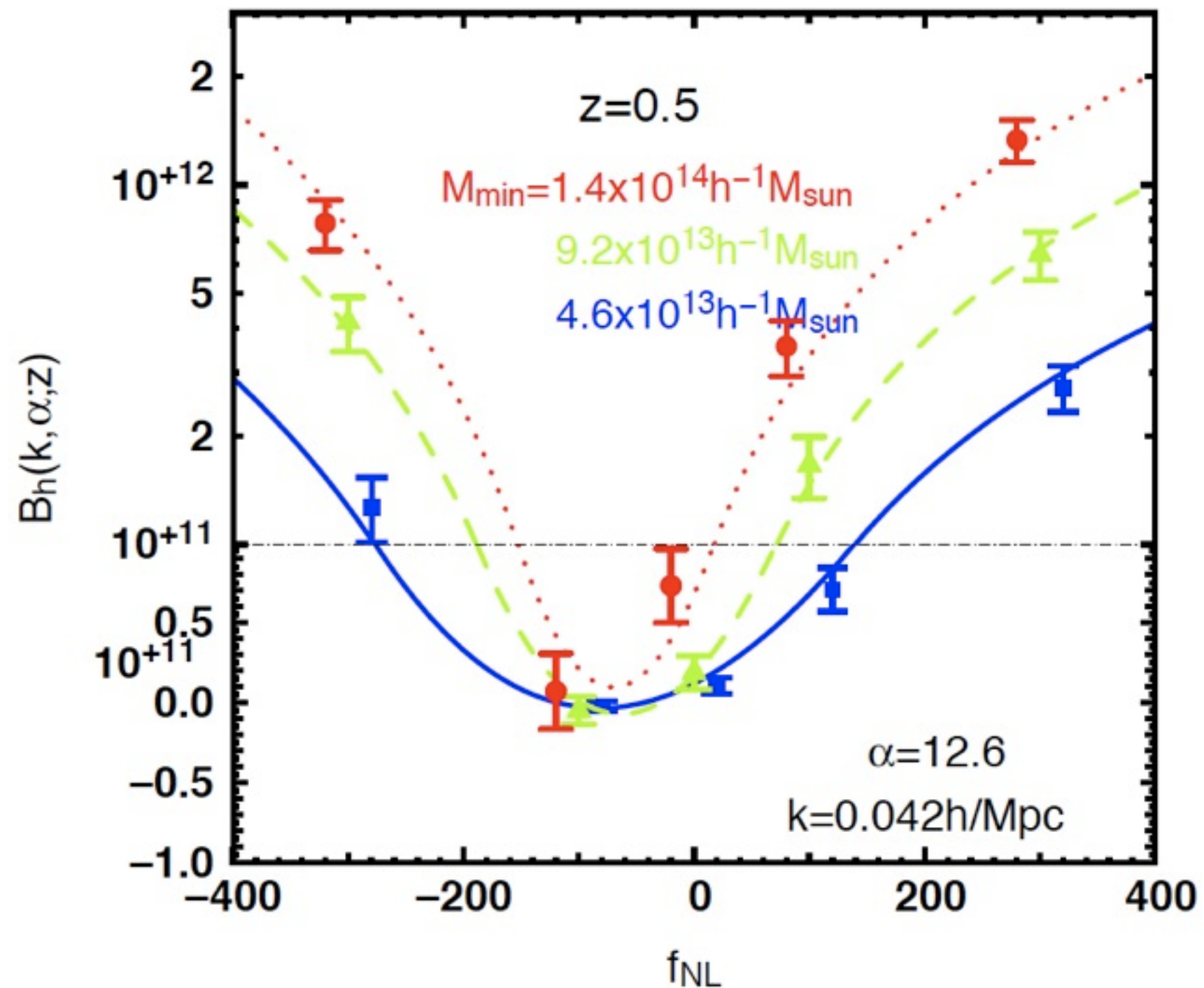
c.f. matter bispectrum: $\propto f_{\text{NL}}^1 \alpha^1 k^0$



Dependence on the halo mass

Halo bispectrum for different minimum halo masses

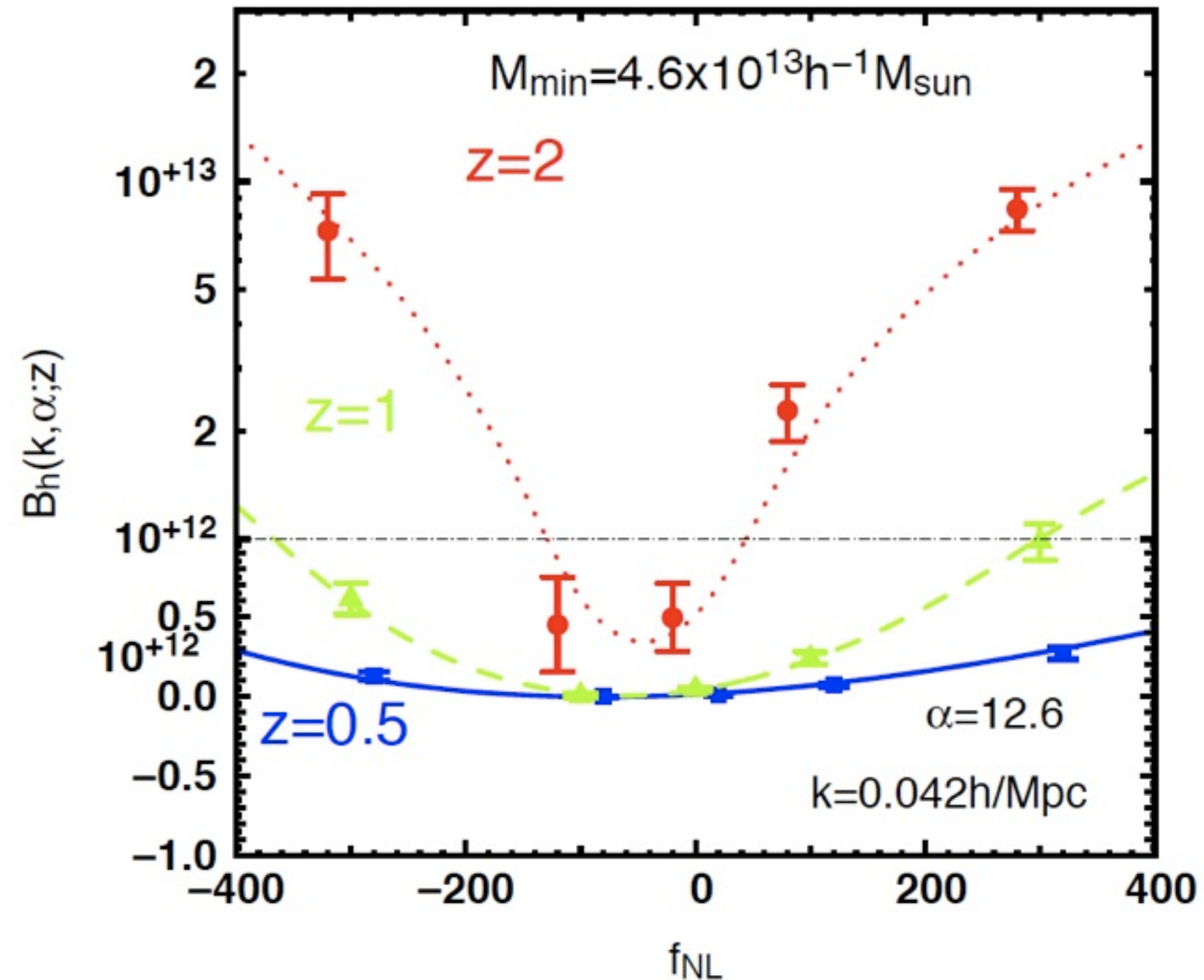
f_{NL}^2 term is more important for more massive haloes



Dependence on redshift

halo bispectrum at
different epochs

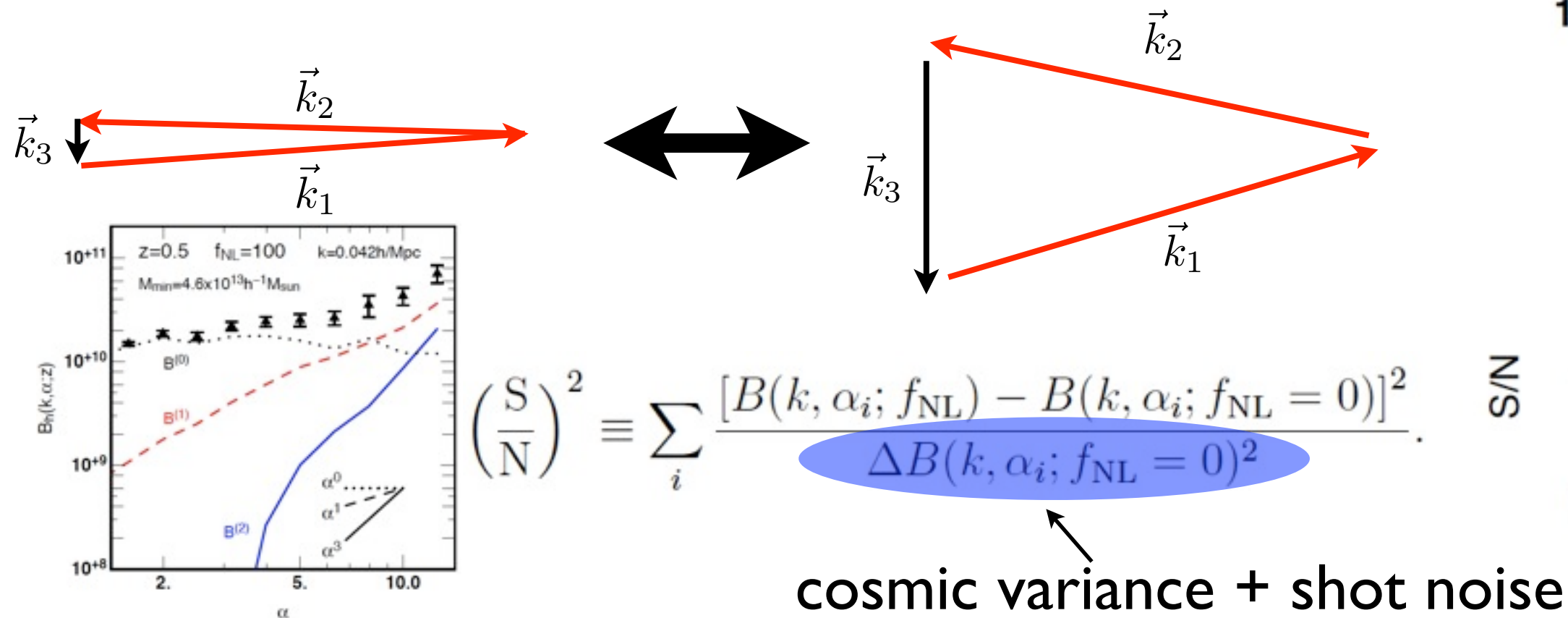
f_{NL}^2 term is more
important at higher
redshift



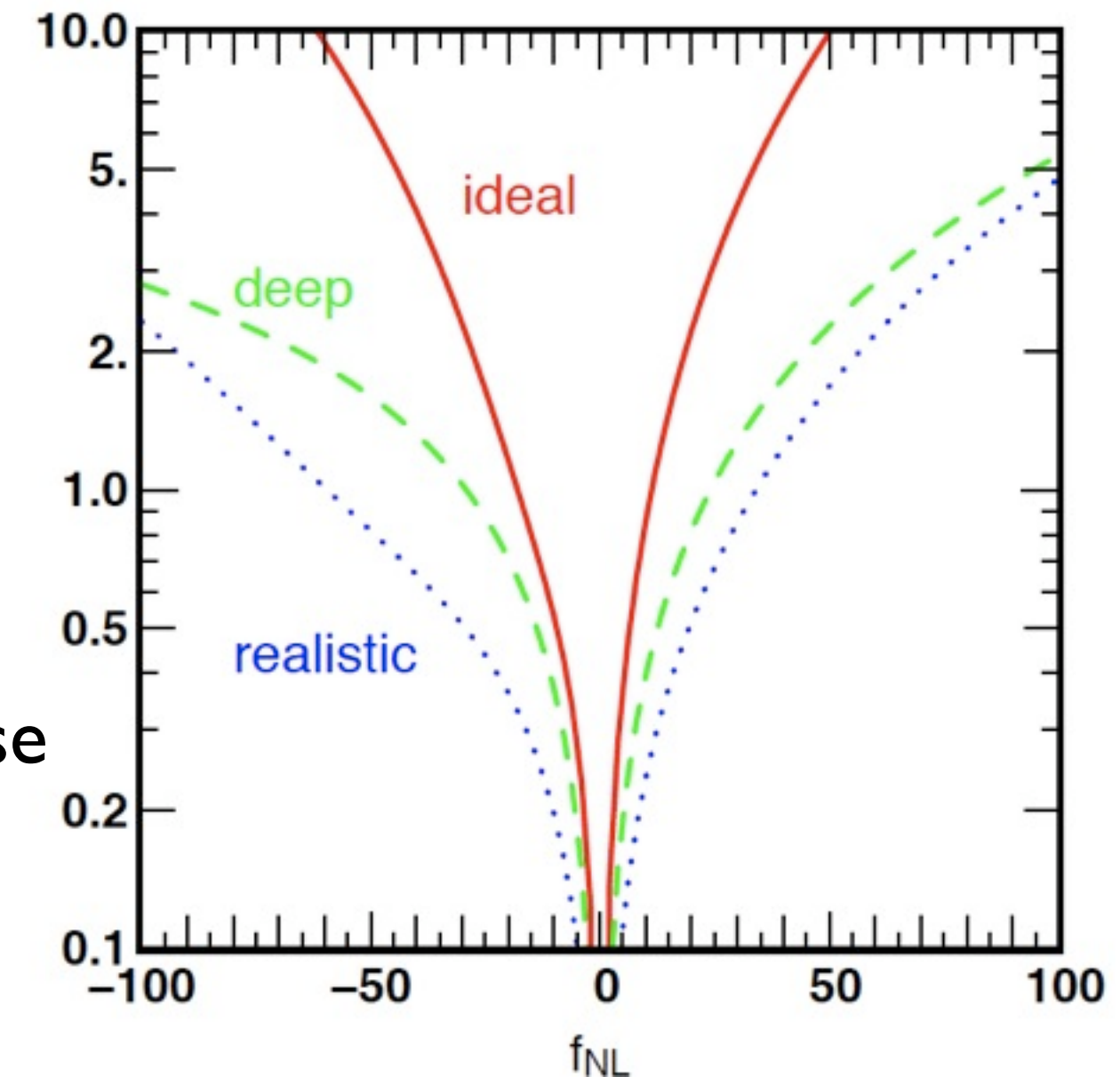
Detectability

detectability when using only limited configurations of $k_1=k_2=0.042h/\text{Mpc}$

survey	z	$V [h^{-3}\text{Gpc}^3]$	$n_g [h^3\text{Mpc}^{-3}]$	$M_{\text{min}} [h^{-1}M_{\odot}]$	b_1
IDEAL	1.0	100	1×10^{-3}	2.8×10^{12}	1.9
REALISTIC	1.0	10	5×10^{-4}	5.0×10^{12}	2.2
DEEP	2.0	3	3×10^{-4}	3.4×10^{12}	3.3



$|f_{\text{NL}}| < 10$ from ultimate future surveys!
 $(|f_{\text{NL}}| < 1 \text{ if we use full information!})$

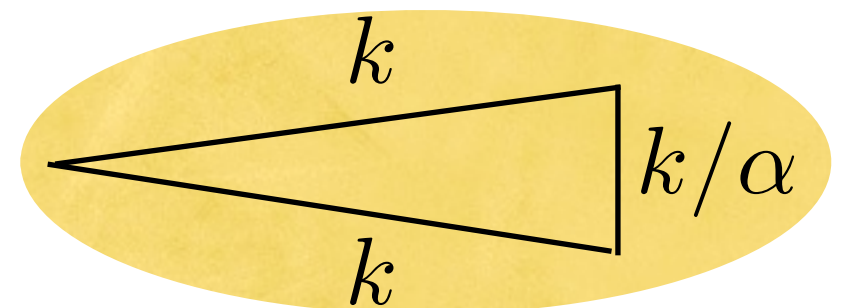


Summary

We have examined the effects of local-type primordial non-Gaussianity on the halo bispectrum:

- Big difference between the matter and halo bispectra
- f_{NL}^2 term is important for squeezed configurations at large scale
- Especially for massive haloes at high z
- The trend can be explained by analytical model assuming local bias

new term $\propto f_{\text{NL}}^2 \alpha^3 k^{-2}$

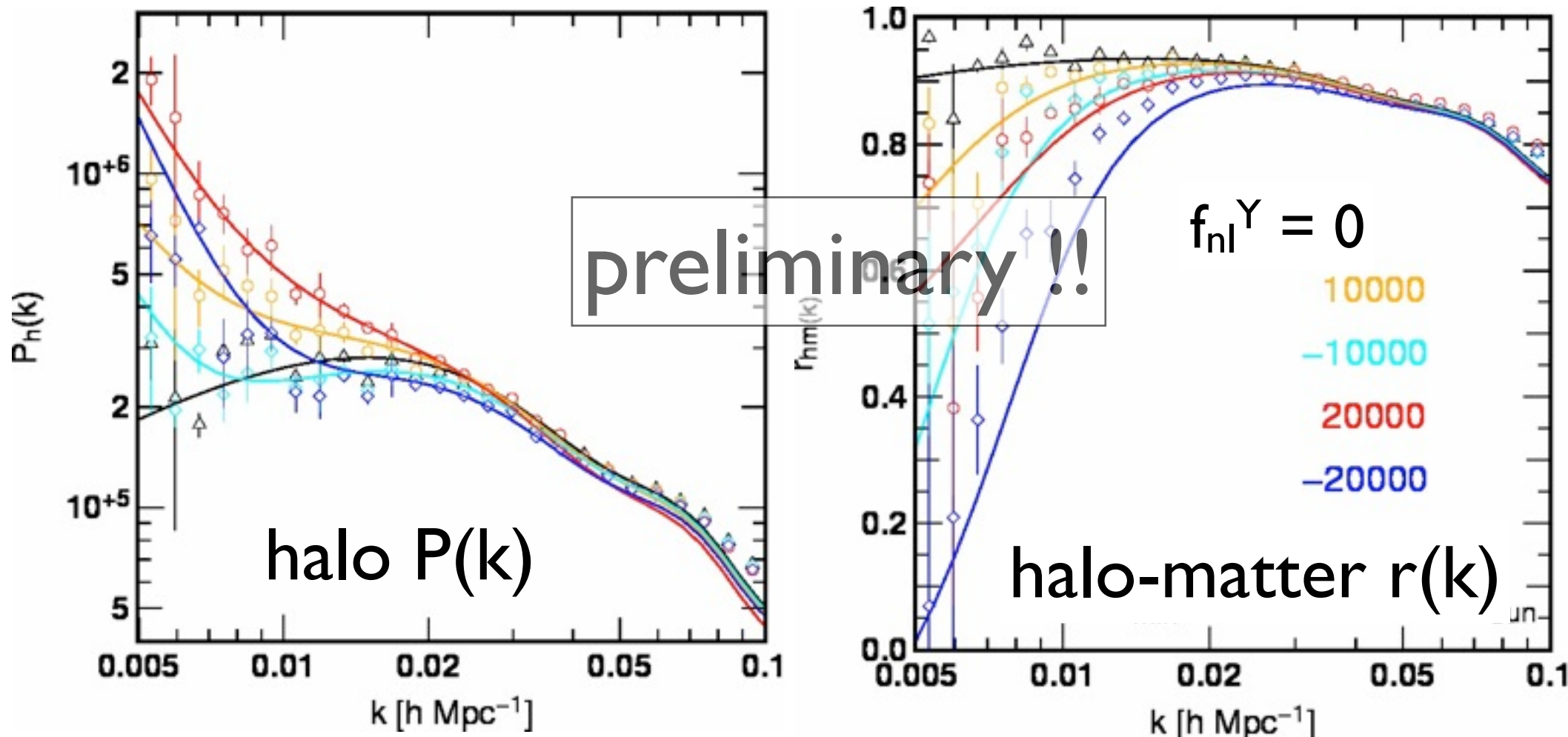


On going project

2 field non-Gaussian models

$$\delta(k) = M_X(k) X(k) + M_Y(k) Y(k)$$

Gaussian $\nearrow$ non-Gaussian $\nearrow$ $\left\{ \begin{array}{l} \eta + f_{nl}^Y (\eta^2 - \langle \eta^2 \rangle) \\ \eta^2 - \langle \eta^2 \rangle \end{array} \right. \times \left\{ \begin{array}{l} \zeta \\ s \end{array} \right. \times \text{correlation}$



c.f., inflaton+curvaton model

Tseliakhovich+I0: theory

Smith&LoVerde I0: simulation