

Higher signal from Lower densities

Elena Massara

Collaborators: F. Villaescusa-Navarro (Princeton), S. Ho (CCA, Princeton, Carnegie Mellon),
N. Dalal (Perimeter), D. Spergel (CCA, Princeton)



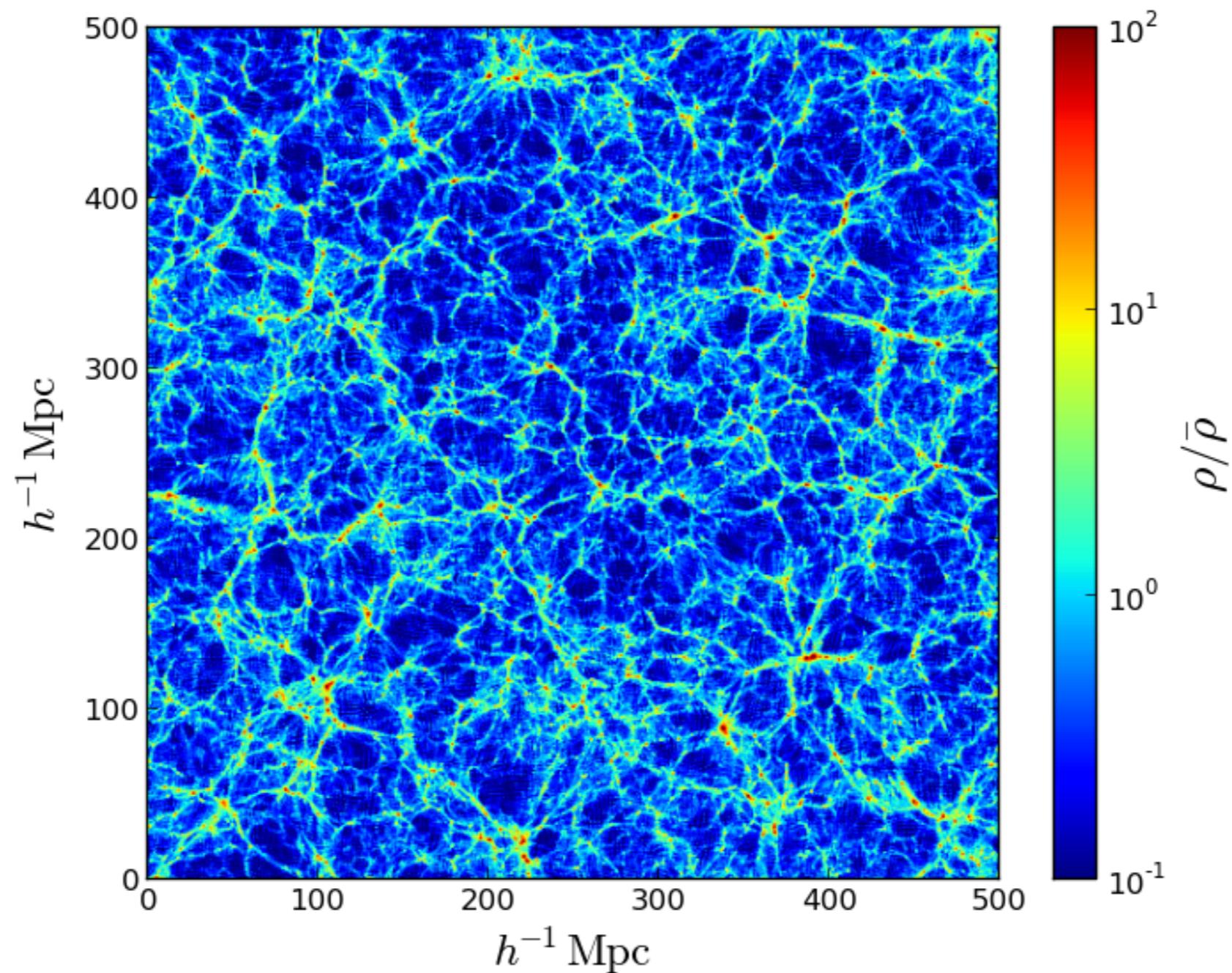
UNIVERSITY OF
WATERLOO

WATERLOO CENTRE FOR
ASTROPHYSICS

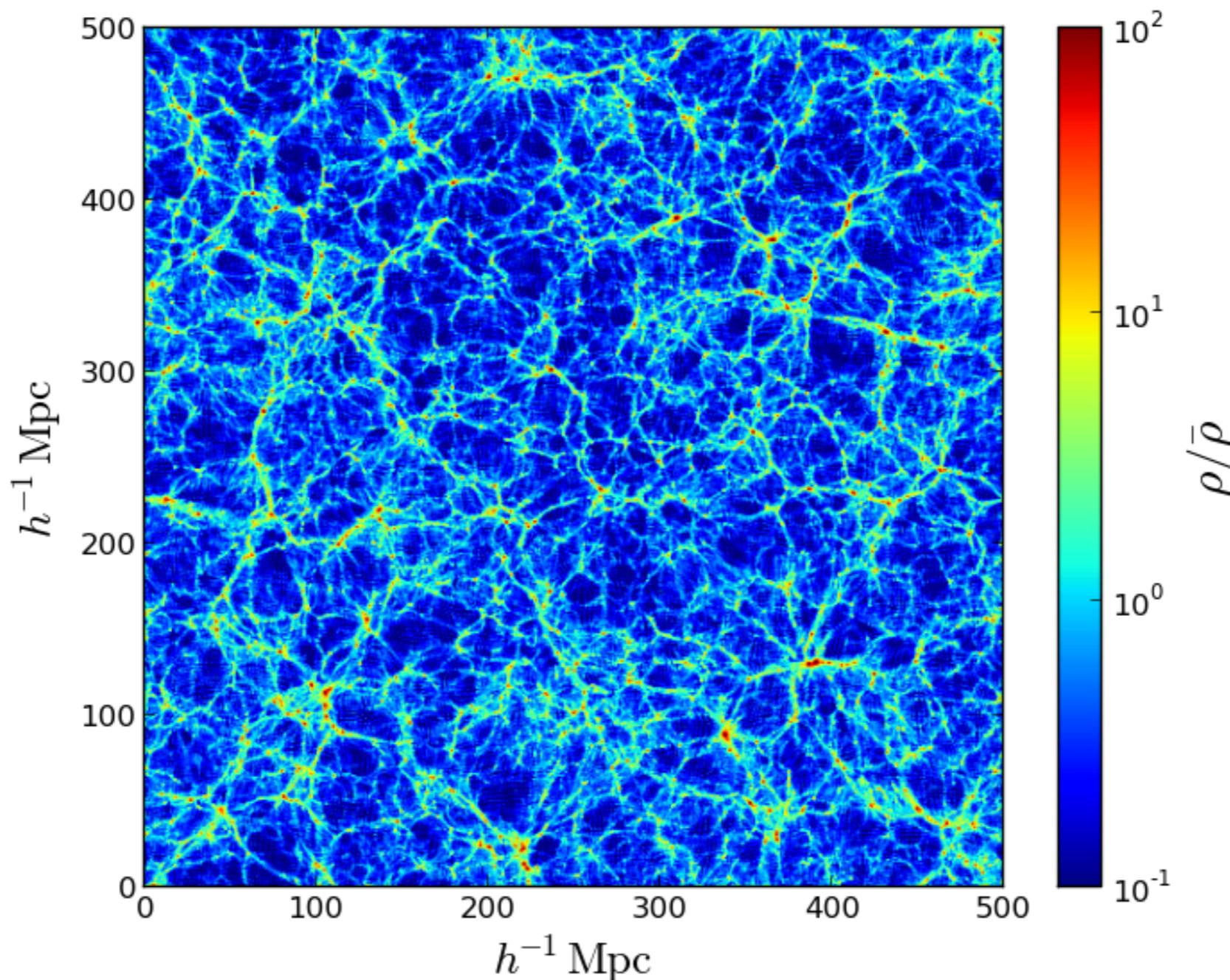
OUTLINE

- Introduction
- Reasons why low density regions are of interest when studying cosmology
- Marked power spectrum:
 - Definition
 - Fisher analysis
 - Results and their interpretation
- Conclusion

THE LARGE SCALE STRUCTURE



THE LARGE SCALE STRUCTURE



The distribution of matter in the Universe is sensitive to:

- properties of dark matter
- nature of dark energy
- neutrinos' masses
- initial condition of the Universe

OBSERVABLES

The distribution of matter in the Universe is sensitive to:

- properties of dark matter
- nature of dark energy
- neutrinos' masses
- initial condition of the Universe

$$\longrightarrow \Omega_m, \Omega_b, h, \sigma_8, n_s, M_\nu$$

OBSERVABLES

The distribution of matter in the Universe is sensitive to:

- properties of dark matter
- nature of dark energy
- neutrinos' masses
- initial condition of the Universe

$$\longrightarrow \Omega_m, \Omega_b, h, \sigma_8, n_s, M_\nu$$

What **observables** should be used to extract the
best constraints on cosmological parameters
from the large-scale structure?

OBSERVABLES

The distribution of matter in the Universe is sensitive to:

- properties of dark matter
- nature of dark energy
- neutrinos' masses
- initial condition of the Universe

$$\longrightarrow \Omega_m, \Omega_b, h, \sigma_8, n_s, M_\nu$$

What **observables** should be used to extract the **best constraints on cosmological parameters** from the large-scale structure?

On large scales/high redshift

2pt function

On small scales/low redshift

2pt-3pt-4pt.. function
voids, peaks, ...

OBSERVABLES

The distribution of matter in the Universe is sensitive to:

- properties of dark matter
- nature of dark energy
- neutrinos' masses
- initial condition of the Universe

$$\longrightarrow \Omega_m, \Omega_b, h, \sigma_8, n_s, M_\nu$$

What **observables** should be used to extract the **best constraints on cosmological parameters** from the large-scale structure?

On large scales/high redshift

2pt function

On small scales/low redshift

2pt-3pt-4pt.. function
voids, peaks, ...

NEUTRINOS

- Three species of neutrinos

NEUTRINOS

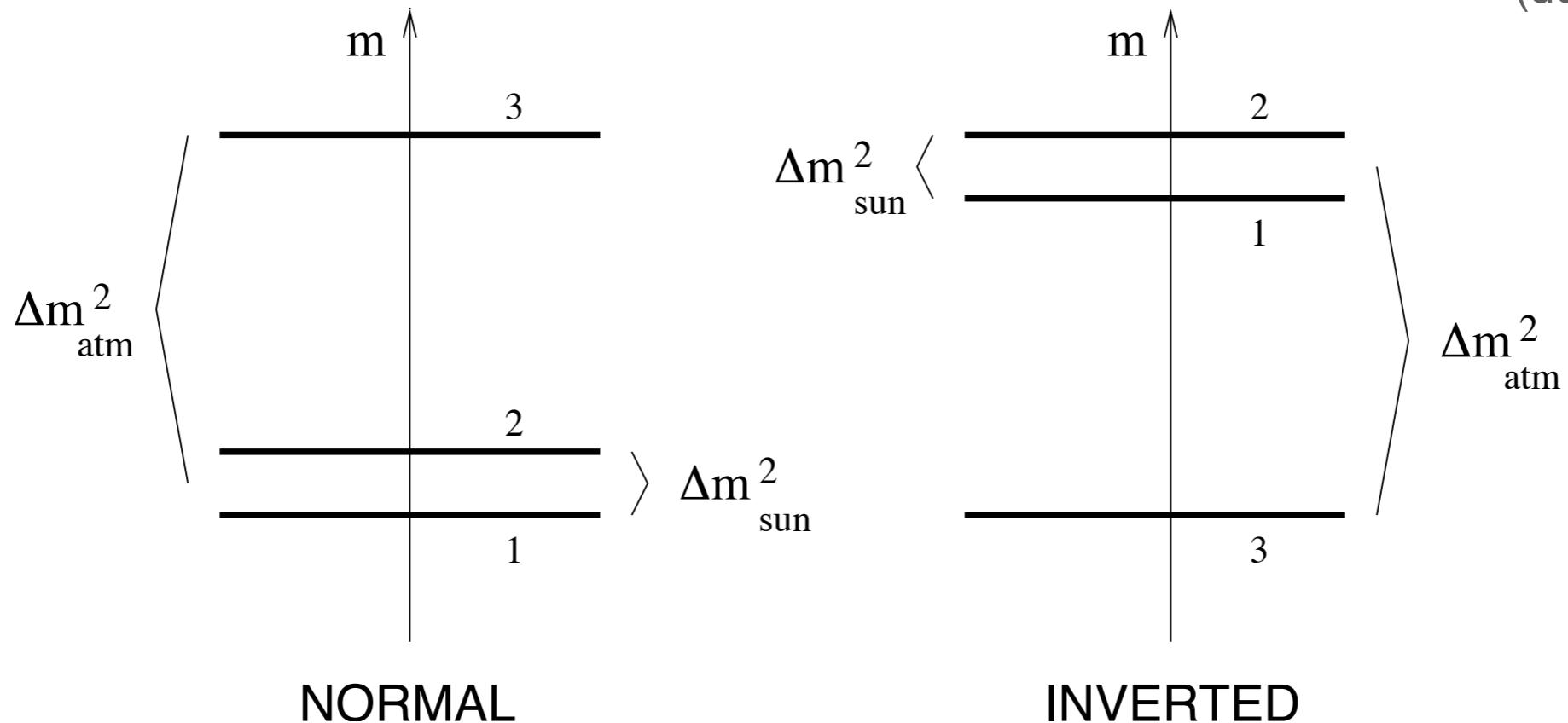
- Three species of neutrinos
- Neutrinos are massive: $\Delta m_{21}^2 = m_2^2 - m_1^2 \sim 10^{-5} eV^2$ from solar neutrinos
 $|\Delta m_{31}^2| = |m_3^2 - m_1^2| \sim 10^{-3} eV^2$ from atmospheric neutrinos

(de Salas et al 2017)

NEUTRINOS

- Three species of neutrinos
- Neutrinos are massive: $\Delta m_{21}^2 = m_2^2 - m_1^2 \sim 10^{-5} eV^2$ from solar neutrinos
 $|\Delta m_{31}^2| = |m_3^2 - m_1^2| \sim 10^{-3} eV^2$ from atmospheric neutrinos

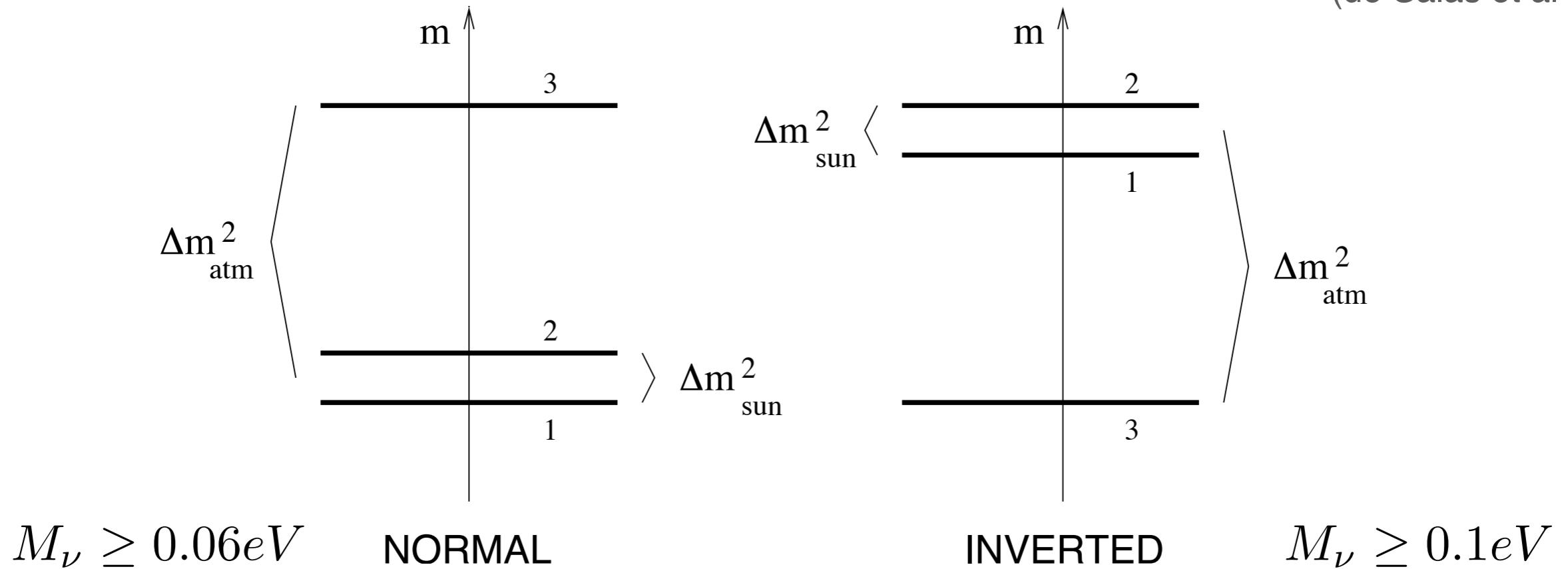
(de Salas et al 2017)



NEUTRINOS

- Three species of neutrinos
- Neutrinos are massive: $\Delta m_{21}^2 = m_2^2 - m_1^2 \sim 10^{-5} eV^2$ from solar neutrinos
 $|\Delta m_{31}^2| = |m_3^2 - m_1^2| \sim 10^{-3} eV^2$ from atmospheric neutrinos

(de Salas et al 2017)

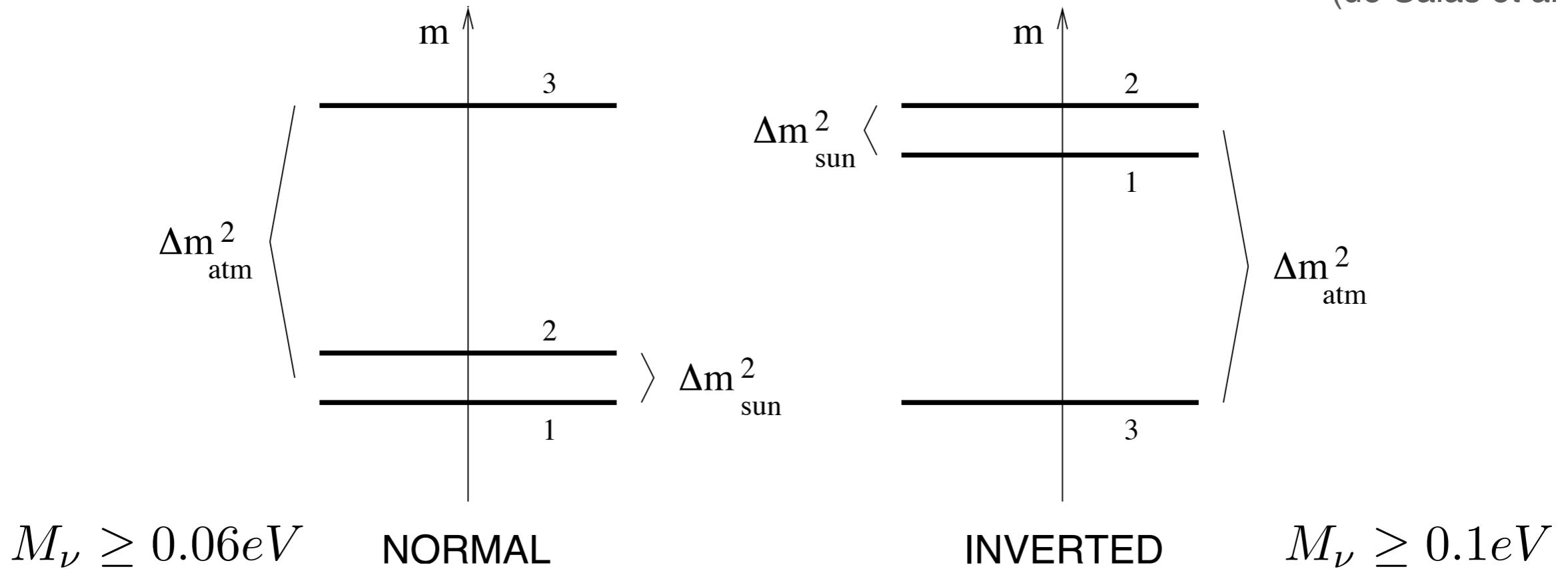


$$M_\nu = \sum m_i$$

NEUTRINOS

- Three species of neutrinos
- Neutrinos are massive: $\Delta m_{21}^2 = m_2^2 - m_1^2 \sim 10^{-5} eV^2$ from solar neutrinos
 $|\Delta m_{31}^2| = |m_3^2 - m_1^2| \sim 10^{-3} eV^2$ from atmospheric neutrinos

(de Salas et al 2017)

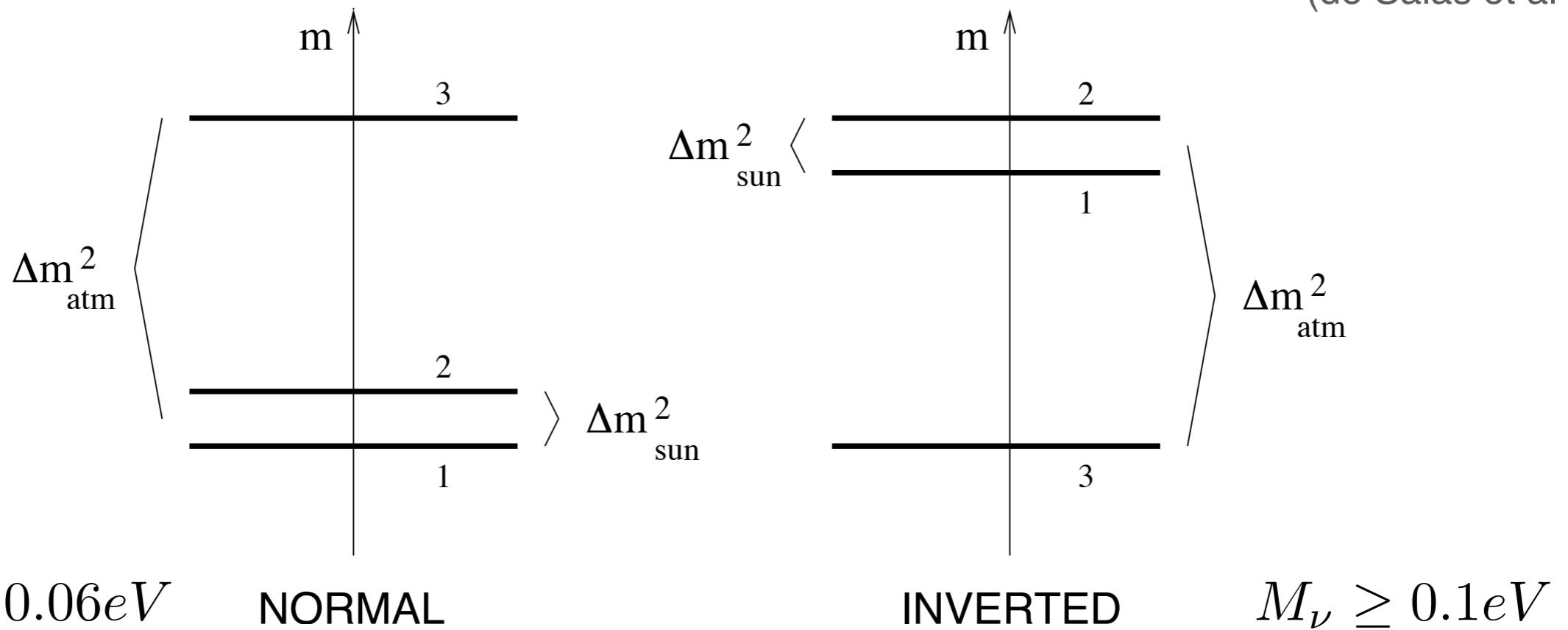


- Cosmology: $M_\nu = \sum m_i < 0.12 \text{ eV}$ (95% C.L.) (Planck 2018)

NEUTRINOS

- Three species of neutrinos
- Neutrinos are massive: $\Delta m_{21}^2 = m_2^2 - m_1^2 \sim 10^{-5} eV^2$ from solar neutrinos
 $\Delta m_{31}^2 = |m_3^2 - m_1^2| \sim 10^{-3} eV^2$ from atmospheric neutrinos

(de Salas et al 2017)



- Cosmology: $M_\nu = \sum m_i < 0.12 \text{ eV}$ (95% C.L.) (Planck 2018)

$$0.06 eV < M_\nu < 0.12 eV$$

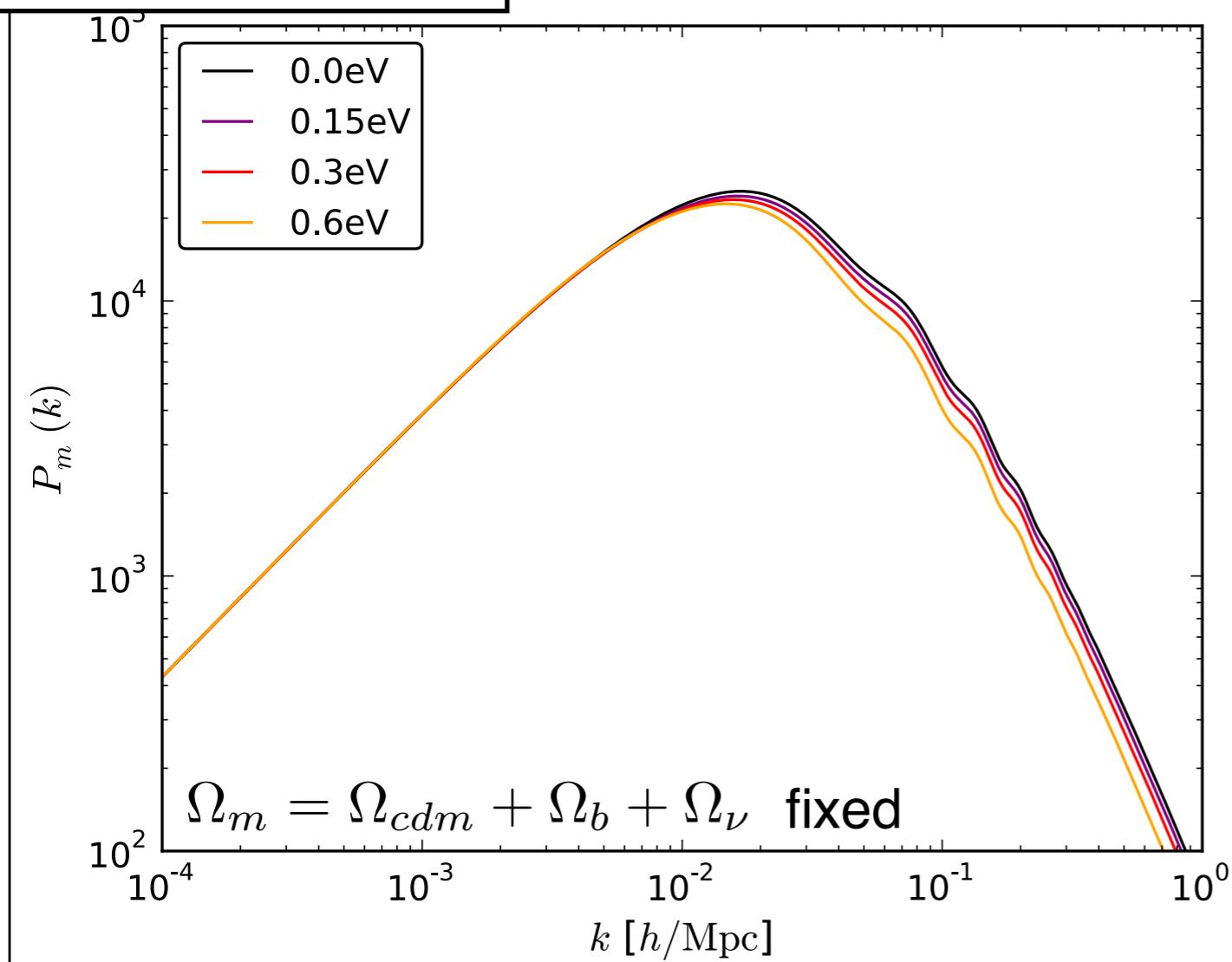
LINEAR ORDER

Neutrino masses have two effects
at linear order:

- 1) delaying the matter-radiation equality
- 2) slowing down the growth of matter perturbations

LINEAR ORDER

Linear Theory



Neutrino masses have two effects

at linear order:

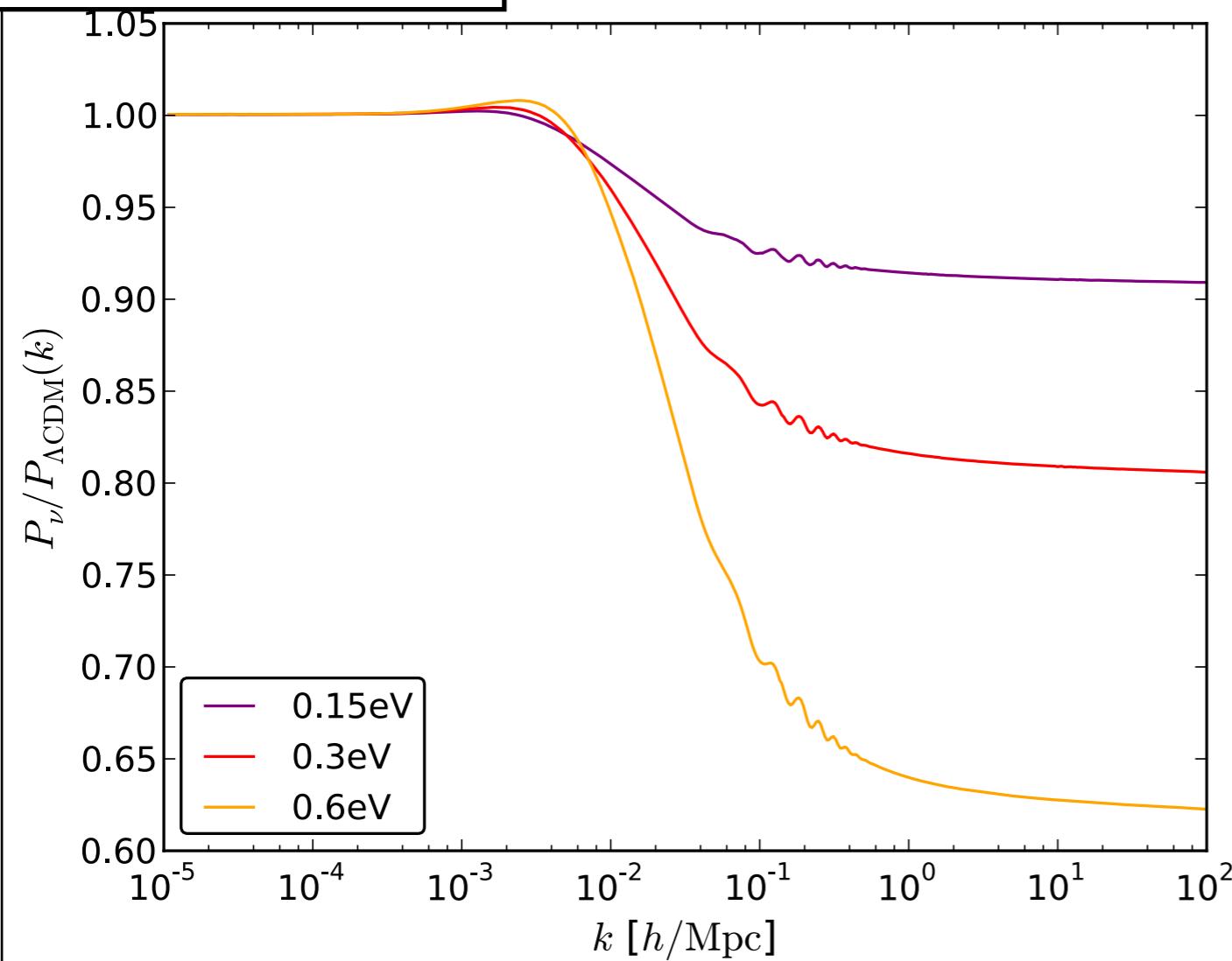
- 1) delaying the matter-radiation equality
- 2) slowing down the growth of matter perturbations



SUPPRESSION of the linear matter power spectrum at intermediate/ small scales

LINEAR ORDER

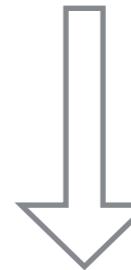
Linear Theory



Neutrino masses have two effects

at linear order:

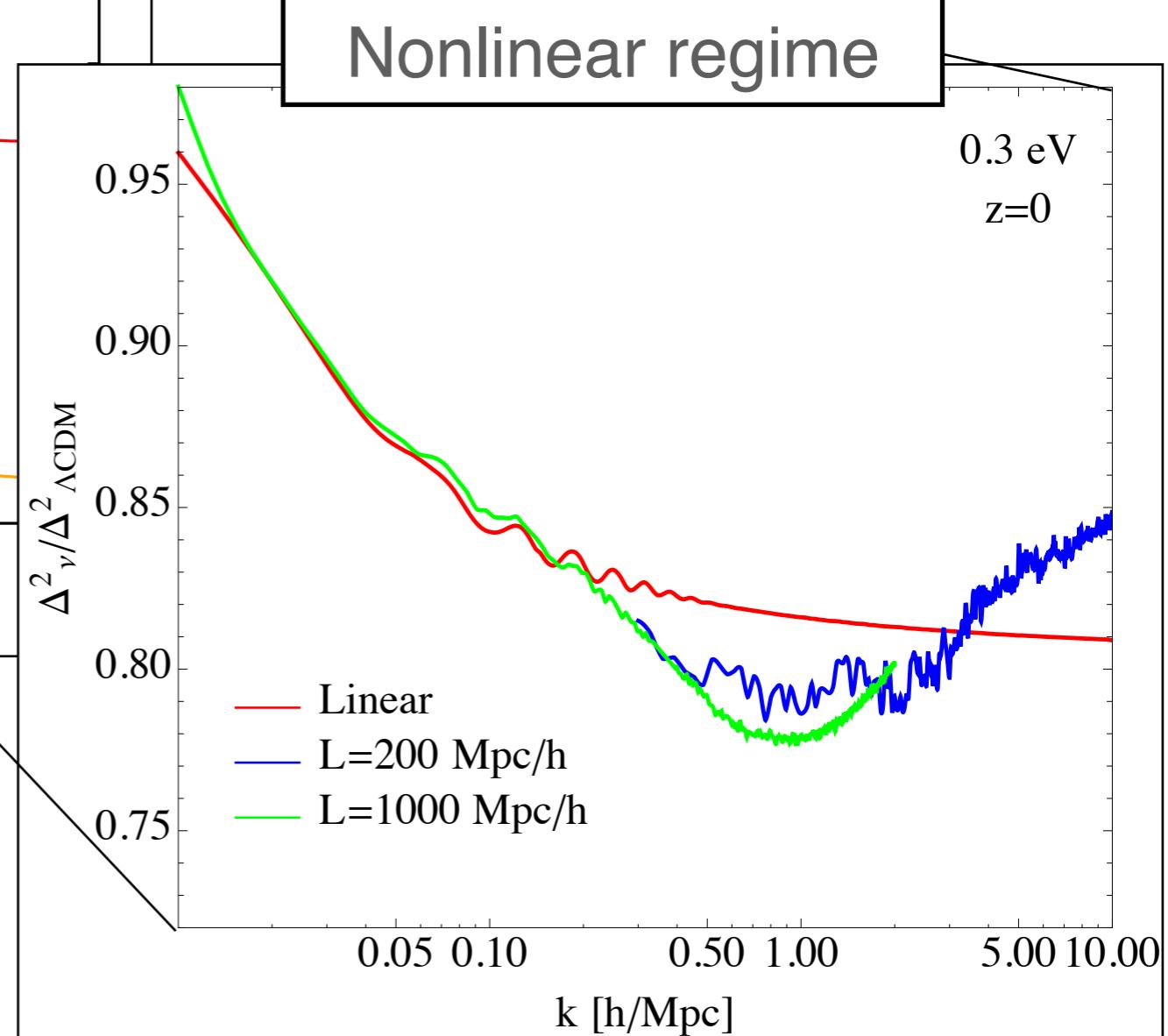
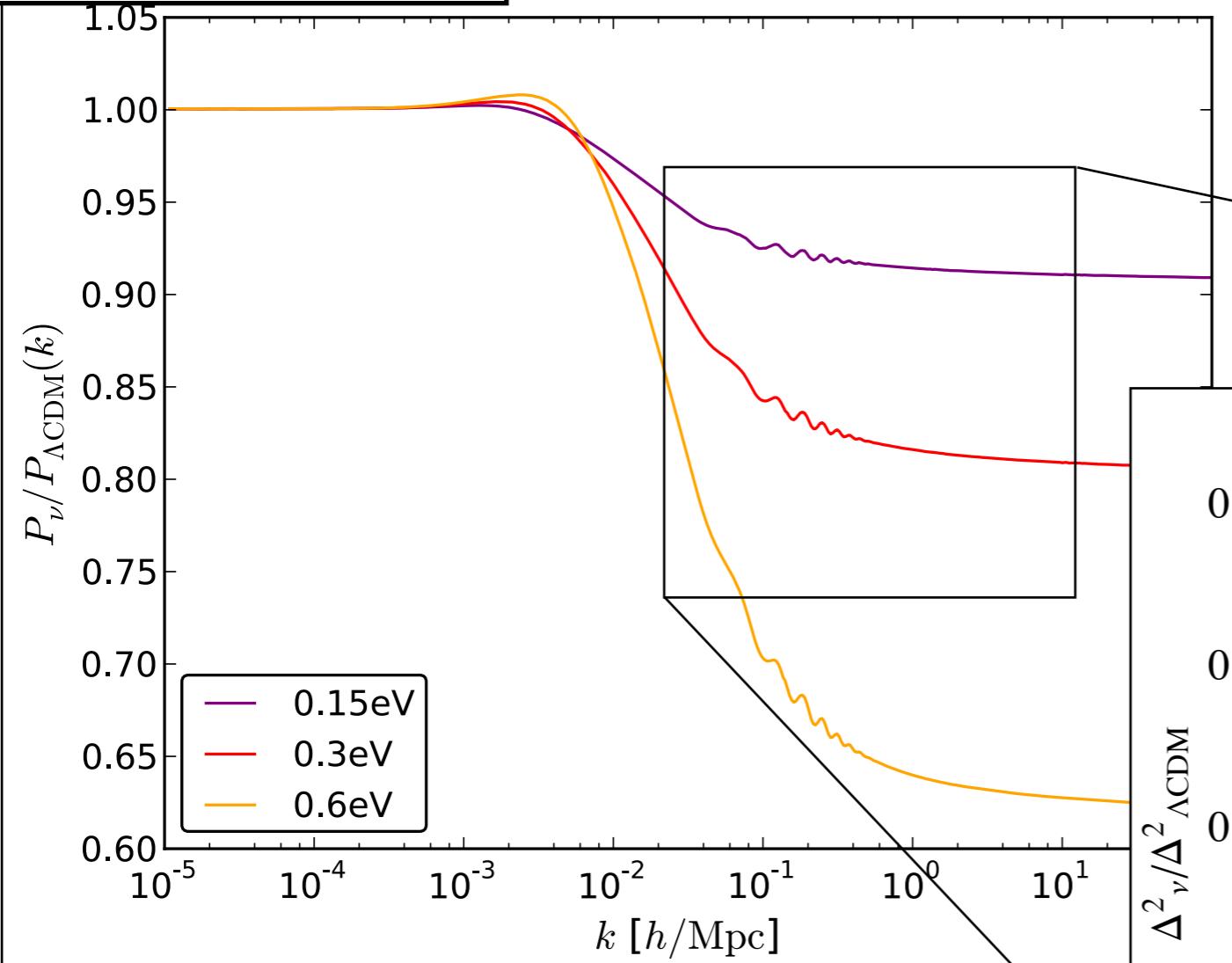
- 1) delaying the matter-radiation equality
- 2) slowing down the growth of matter perturbations



SUPPRESSION of the linear matter power spectrum at intermediate/ small scales

LINEAR & NONLINEAR

Linear Theory

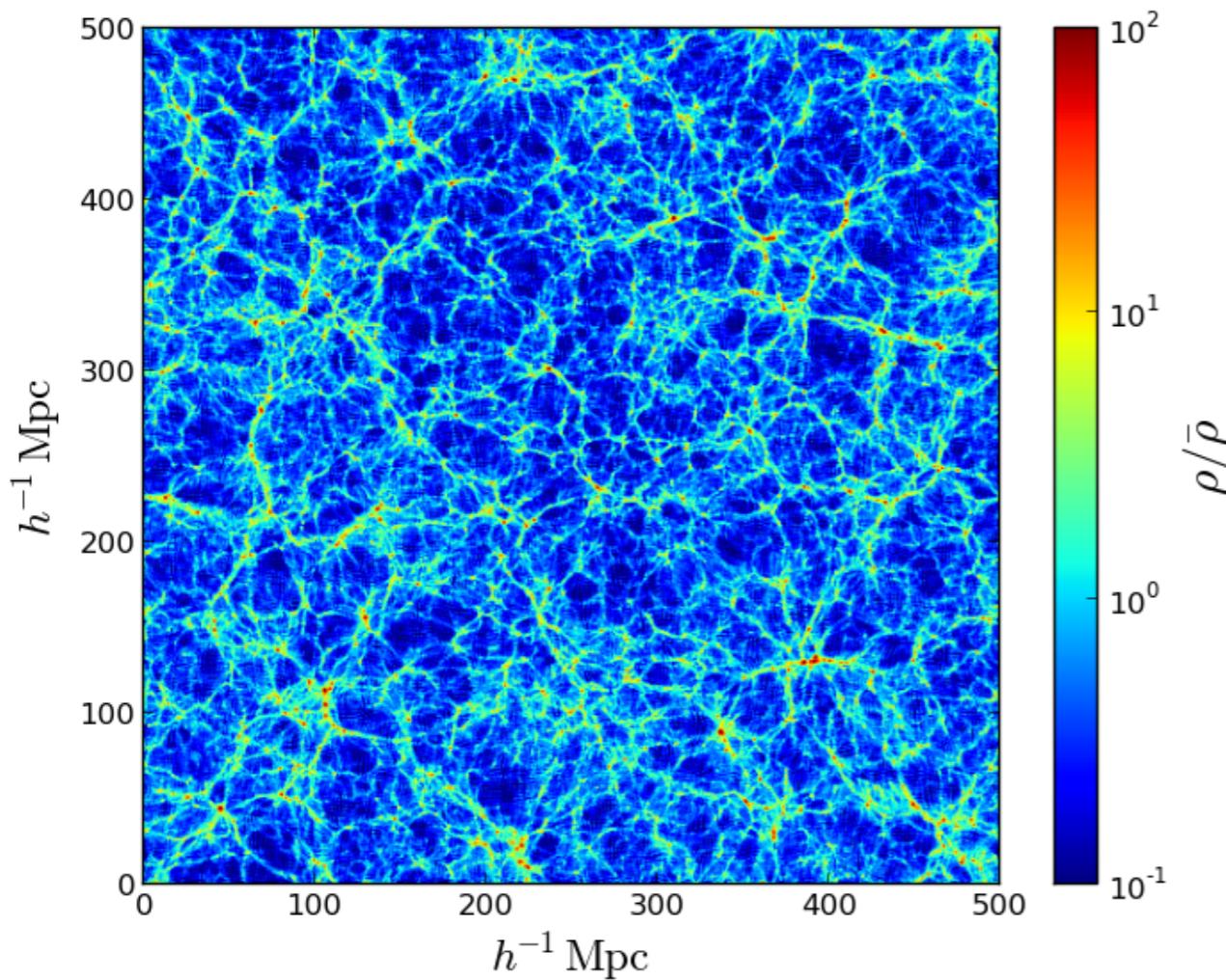


VOIDS

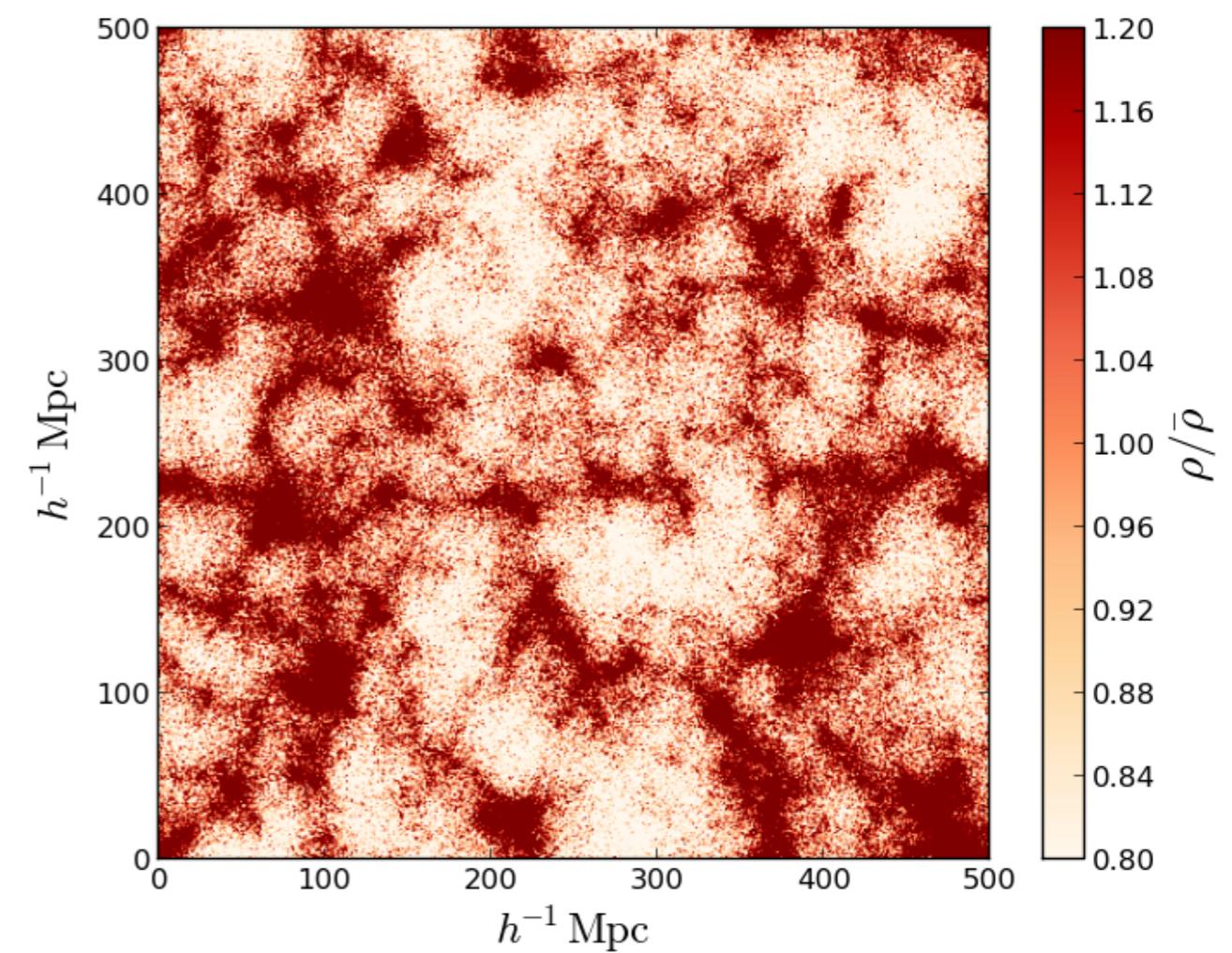
EM, Villaescusa-Navarro, Viel, Sutter, 2015

- Neutrinos: diffuse component

CDM density field



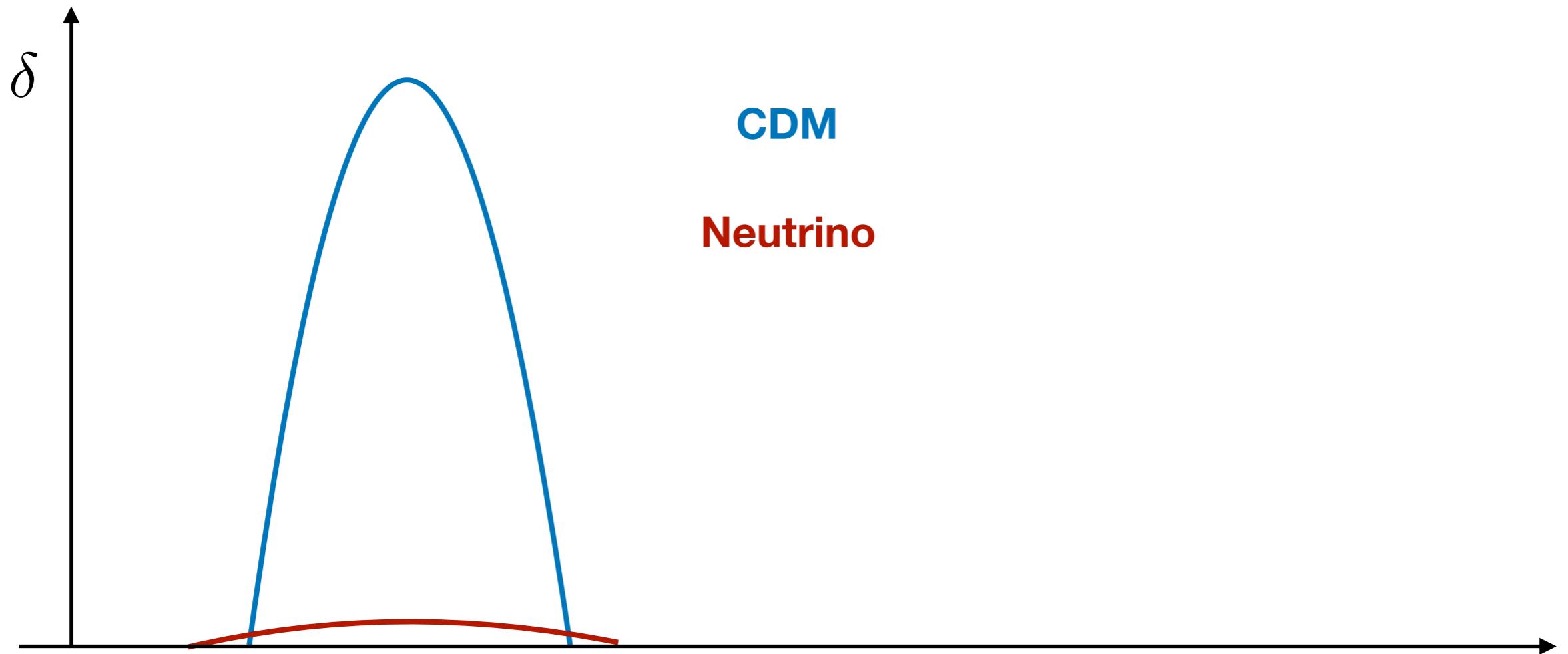
Neutrino density field



VOIDS

EM, Villaescusa-Navarro, Viel, Sutter, 2015

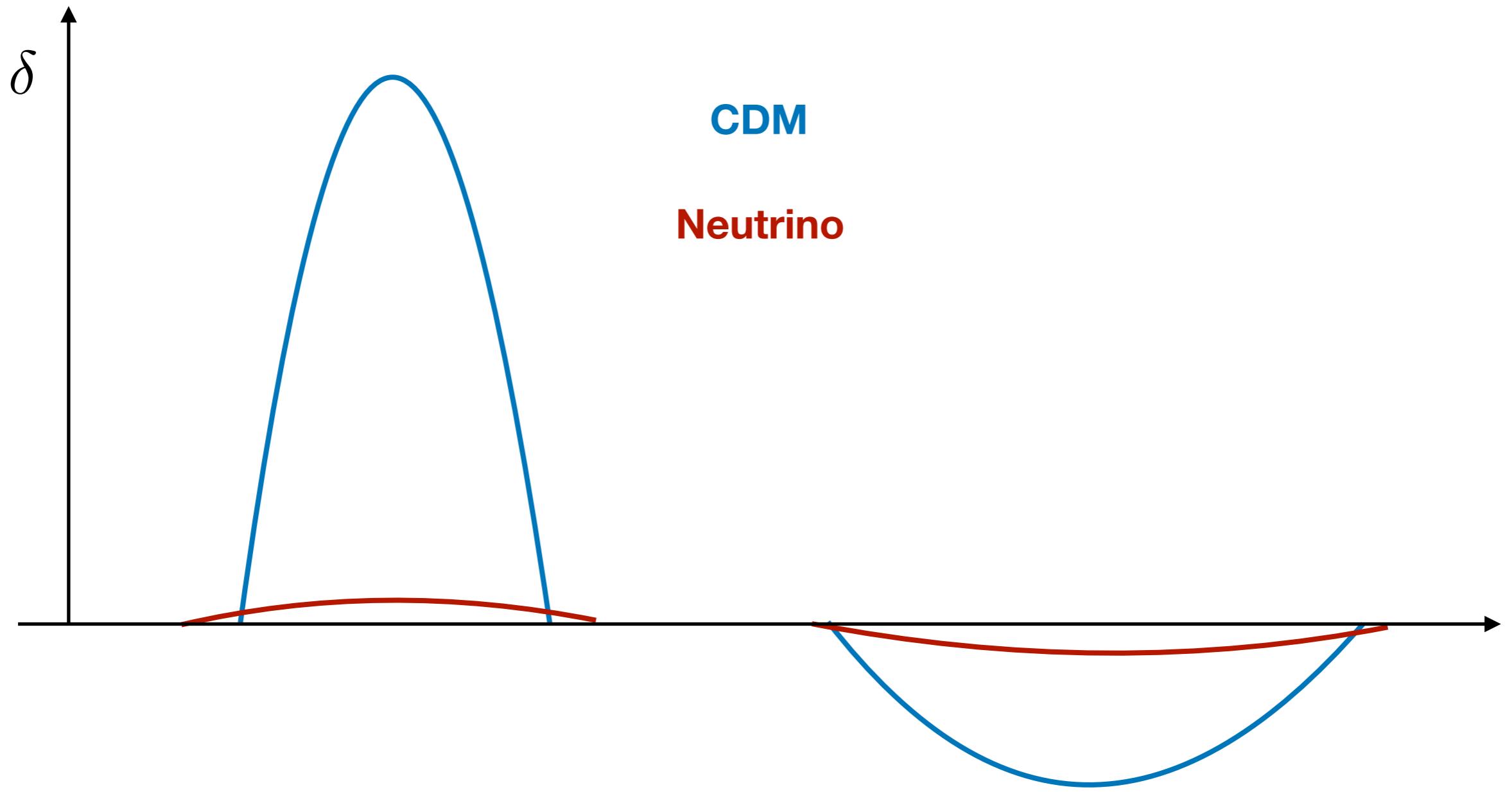
- Neutrinos: diffuse component
- They are very small fraction of the matter in halos



VOIDS

EM, Villaescusa-Navarro, Viel, Sutter, 2015

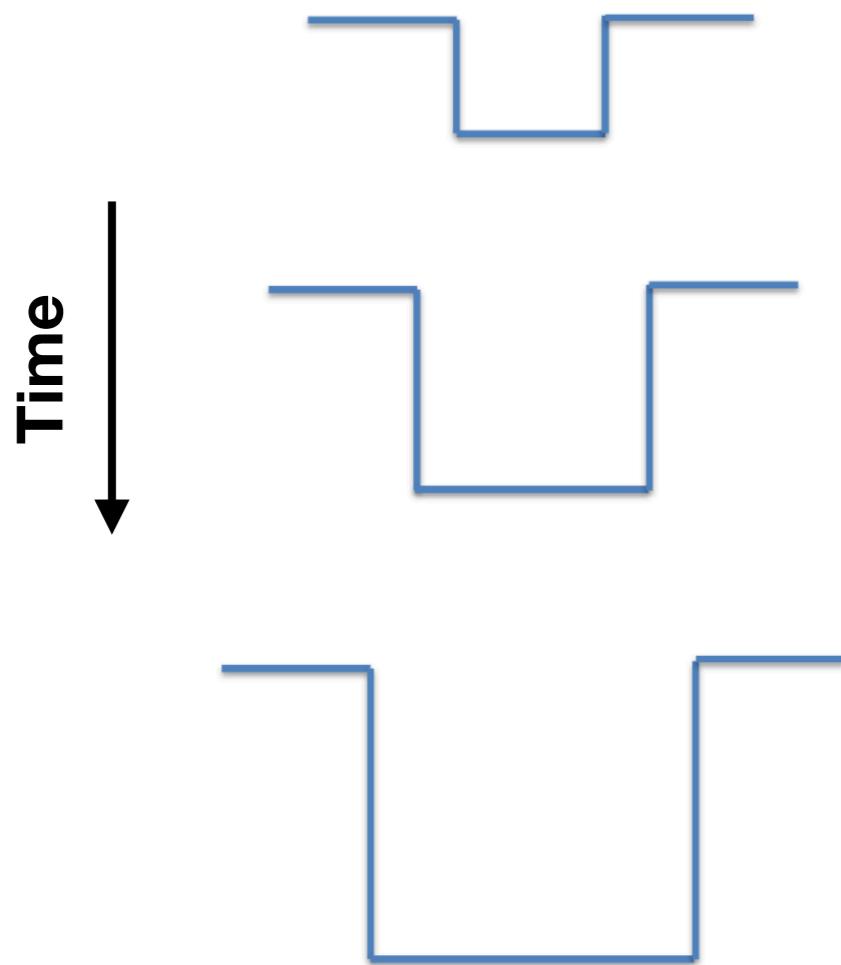
- Neutrinos: diffuse component
- They are very small fraction of the matter in halos
- Voids are under-densities in the CDM field but are filled with neutrinos, therefore in voids the ratio between neutrinos and CDM densities is highest in the Universe



VOIDS

EM, Villaescusa-Navarro, Viel, Sutter, 2015

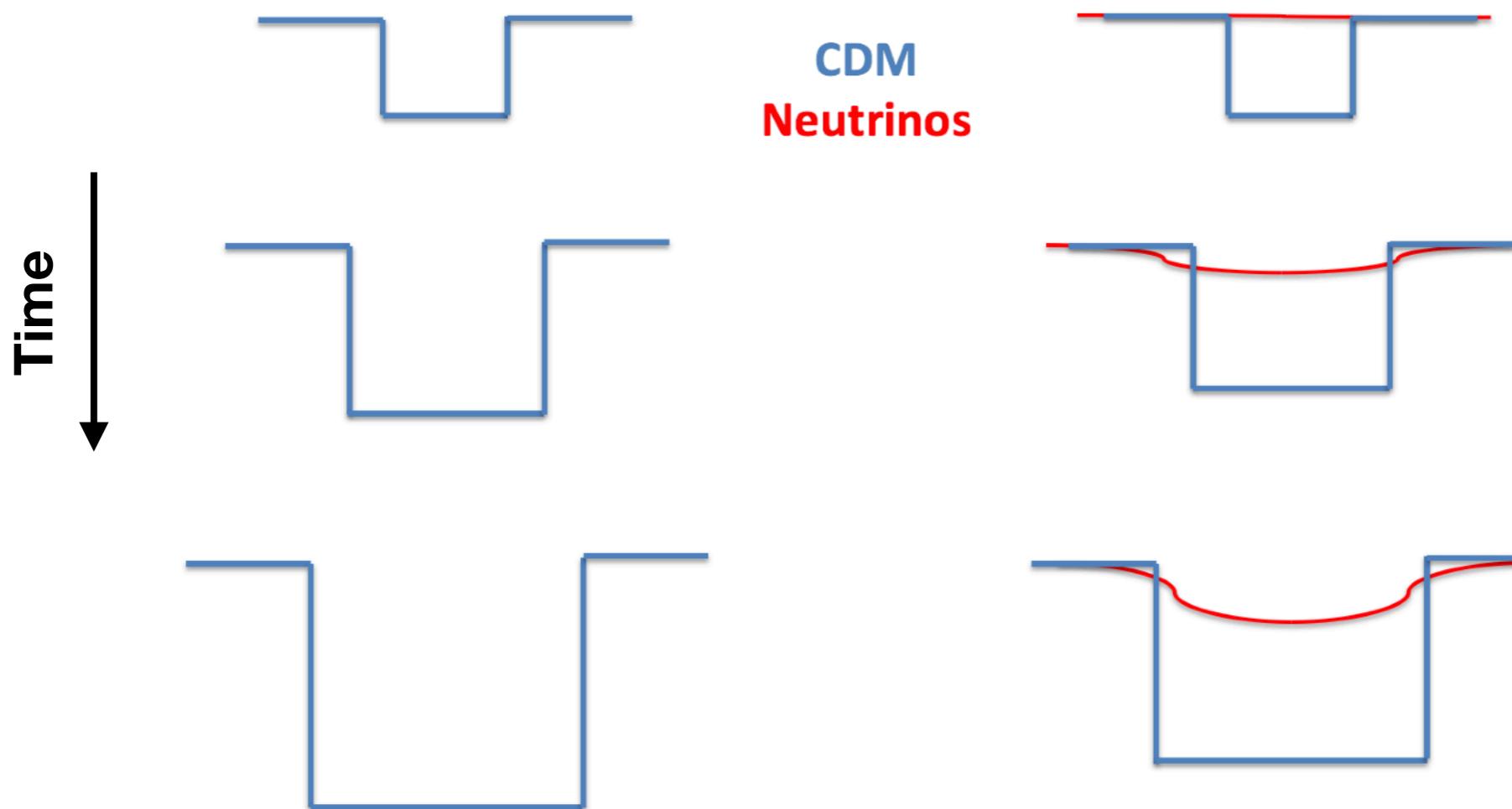
- Neutrinos: diffuse component
- They are very small fraction of the matter in halos
- Voids are under-densities in the CDM field but are filled with neutrinos, therefore in voids the ratio between neutrinos and CDM densities is highest in the Universe



VOIDS

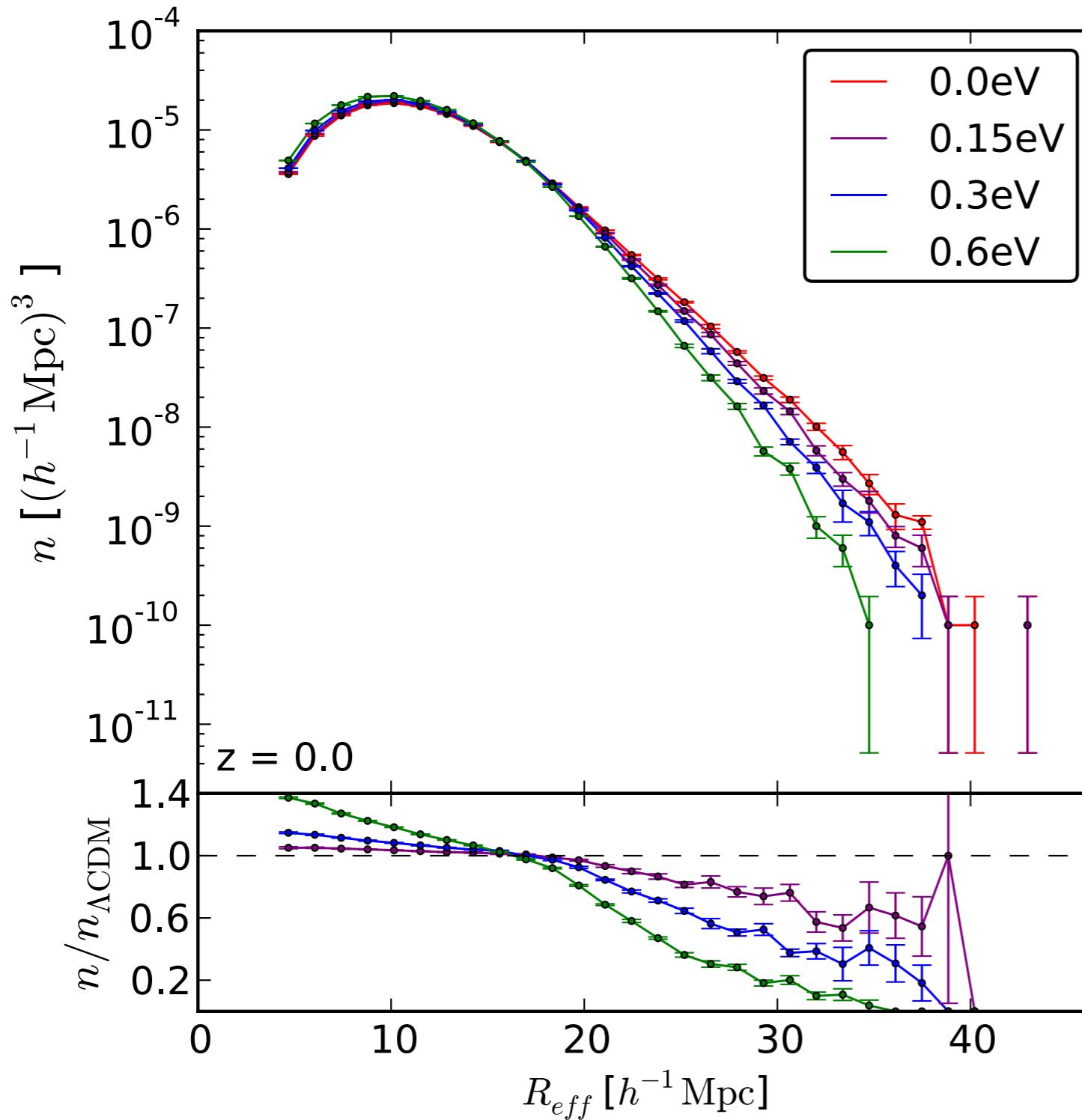
EM, Villaescusa-Navarro, Viel, Sutter, 2015

- Neutrinos: diffuse component
- They are very small fraction of the matter in halos
- Voids are under-densities in the CDM field but are filled with neutrinos, therefore in voids the ratio between neutrinos and CDM densities is highest in the Universe



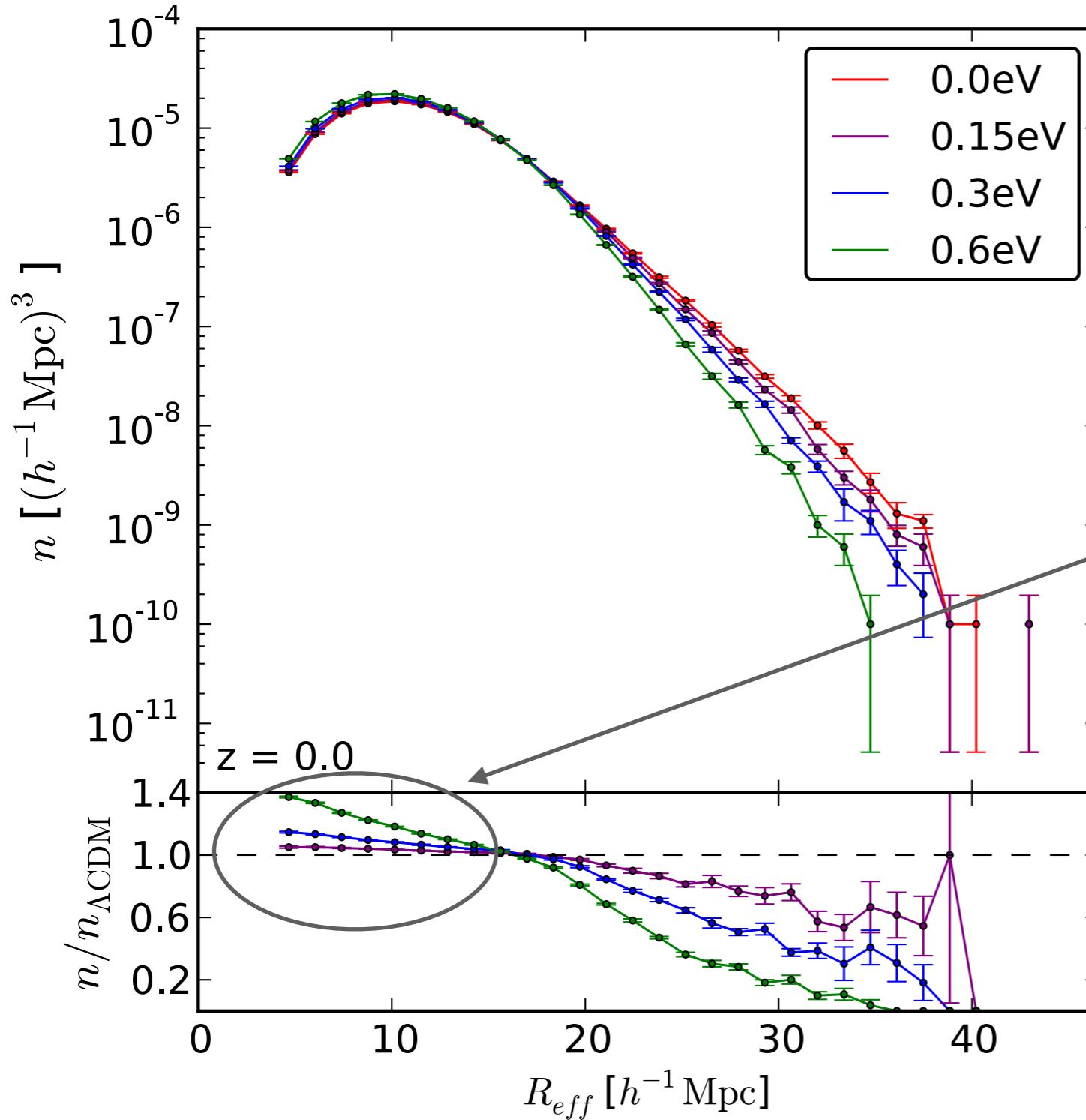
VOID SIZE FUNCTION

EM, Villaescusa-Navarro, Viel, Sutter, 2015



VOID SIZE FUNCTION

EM, Villaescusa-Navarro, Viel, Sutter, 2015

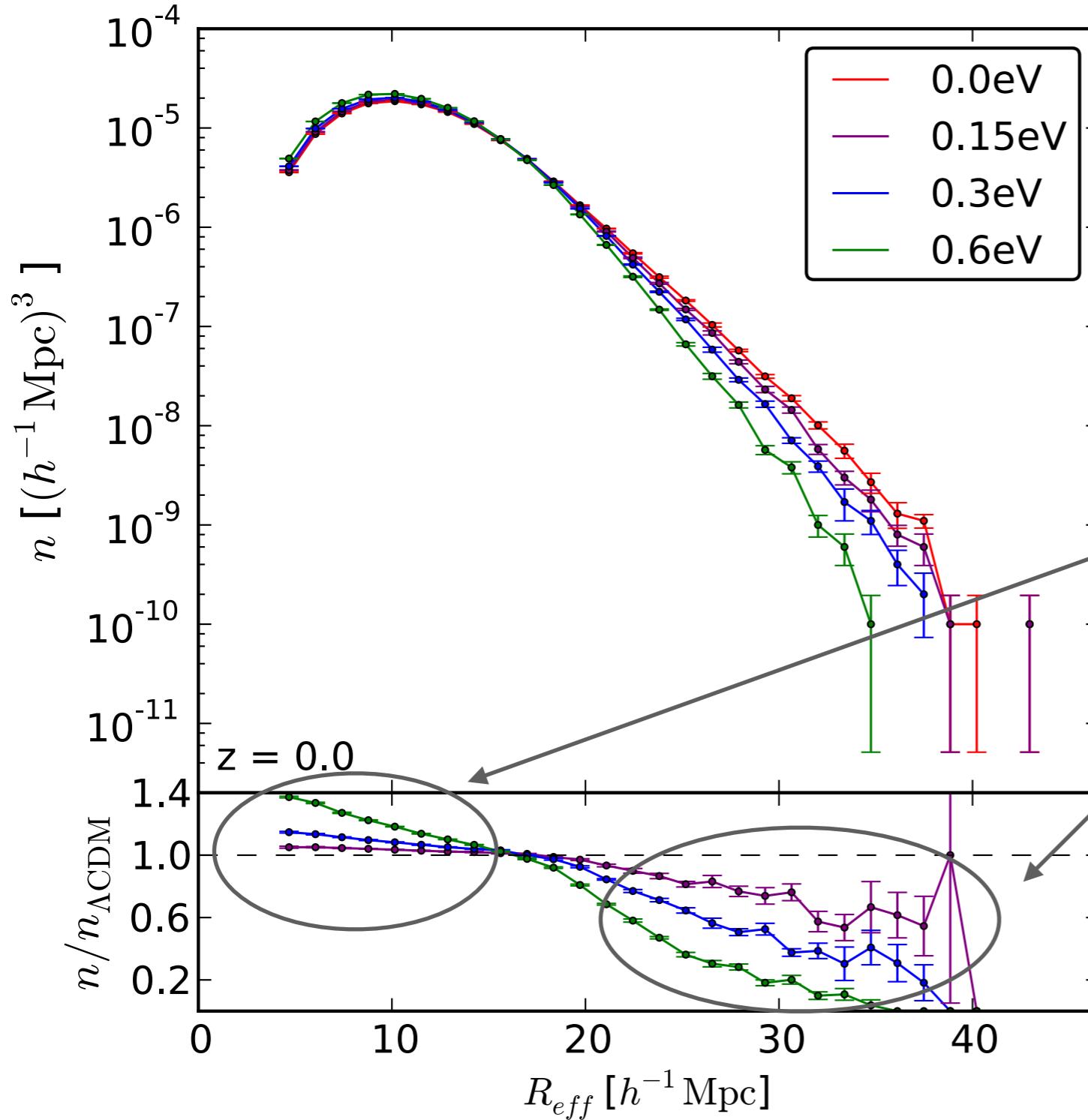


In massive neutrino cosmologies there are

- more small voids

VOID SIZE FUNCTION

EM, Villaescusa-Navarro, Viel, Sutter, 2015

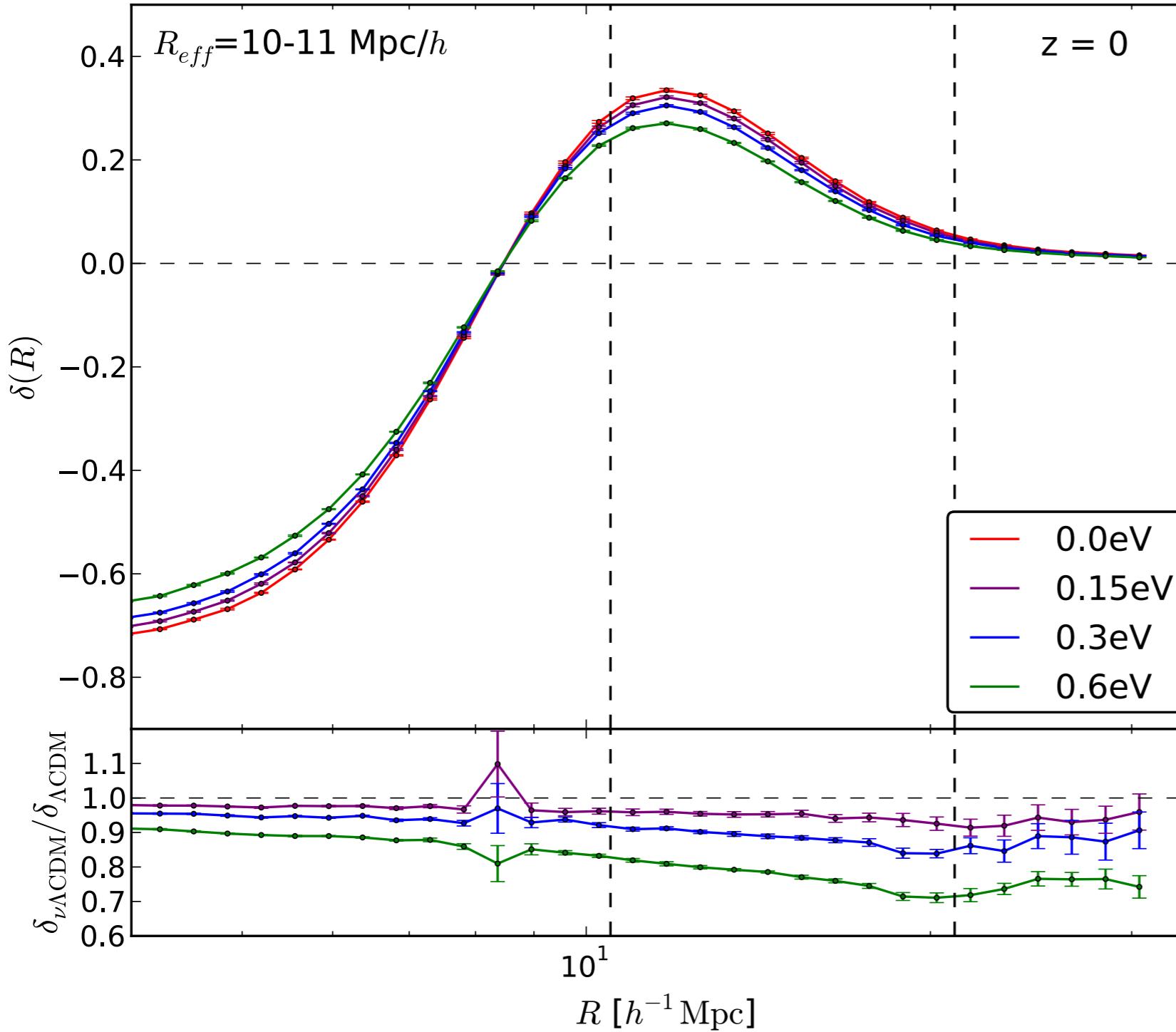


In massive neutrino cosmologies there are

- more small voids
- fewer big voids

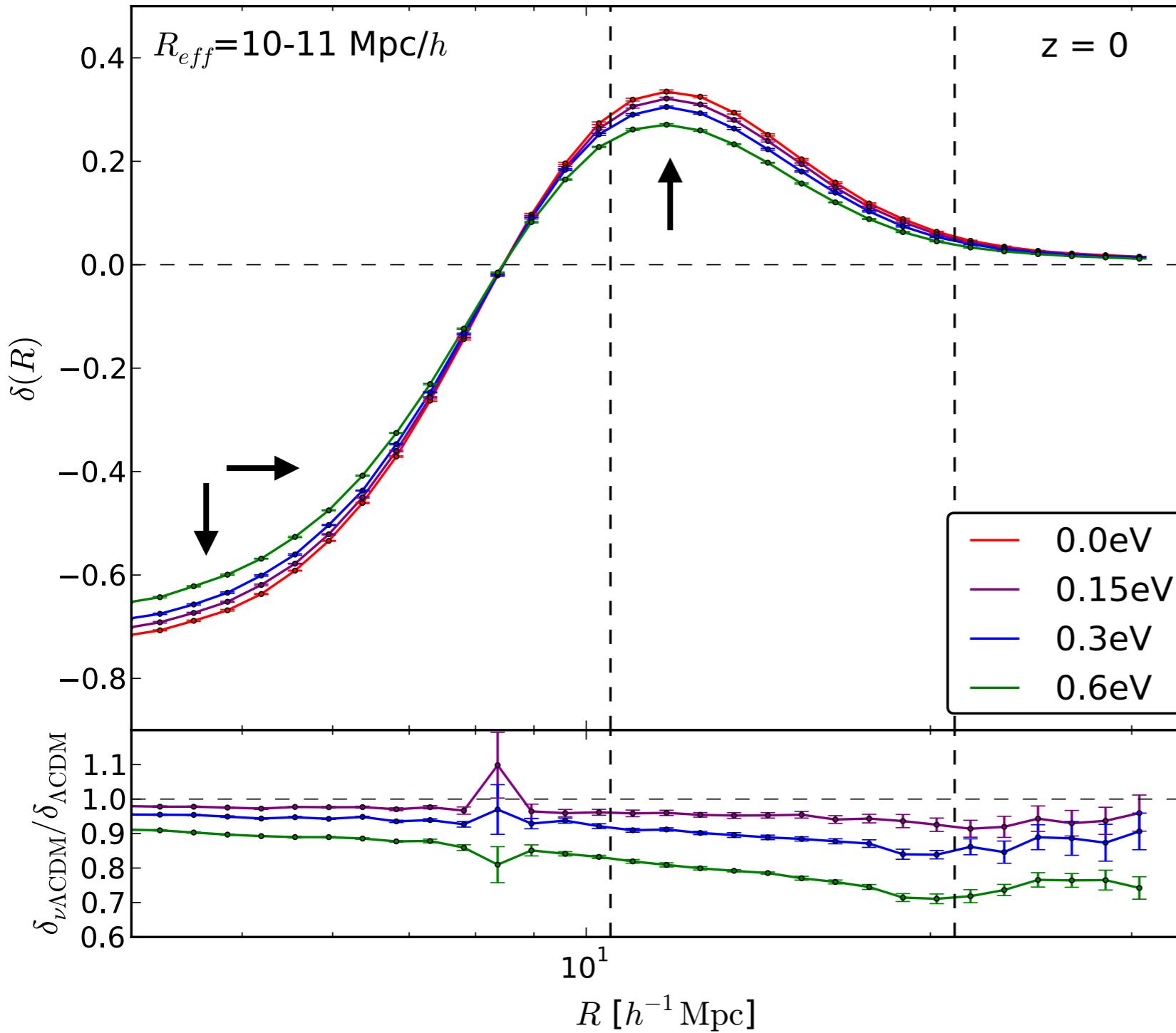
MATTER PROFILE around VOIDS

EM, Villaescusa-Navarro, Viel, Sutter, 2015



MATTER PROFILE around VOIDS

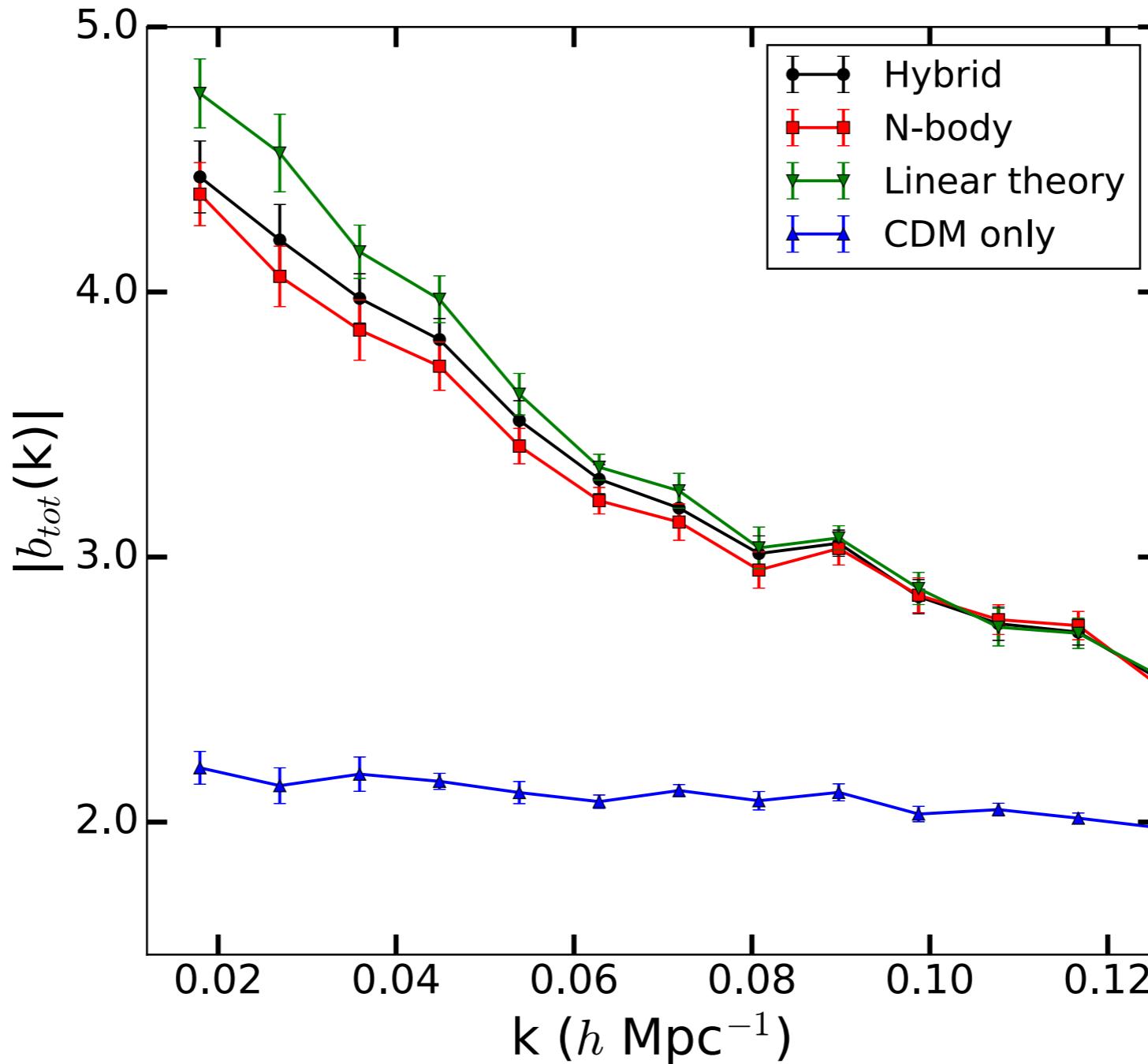
EM, Villaescusa-Navarro, Viel, Sutter, 2015



Voids in massive
neutrino cosmology are
less evolved

VOID BIAS

Banerjee and Dalal, 2016



$$\delta_{\text{void}}(k) = b_{\text{tot}}(k)\delta_{\text{tot}}(k)$$

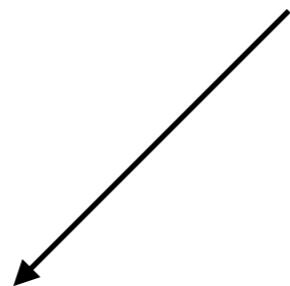
COSMOLOGY with VOIDS

Voids are good laboratories to study cosmology because

- They are sensitive to diffuse components such as
 - neutrinos
 - dark energy
- They have low densities, thus screening mechanism are inefficient and they are favored place where to study modification of gravity
- They have not undergone virialization, thus they are expected to retain most of their initial cosmological information

LOW DENSITY REGIONS

Low density regions are good probe to study cosmology

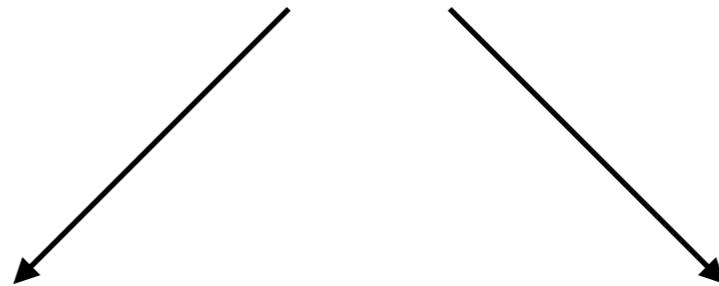


Should we build observables related
to low density regions/voids?

Void Size Function,
void-matter cross-correlation, ...

LOW DENSITY REGIONS

Low density regions are good probe to study cosmology



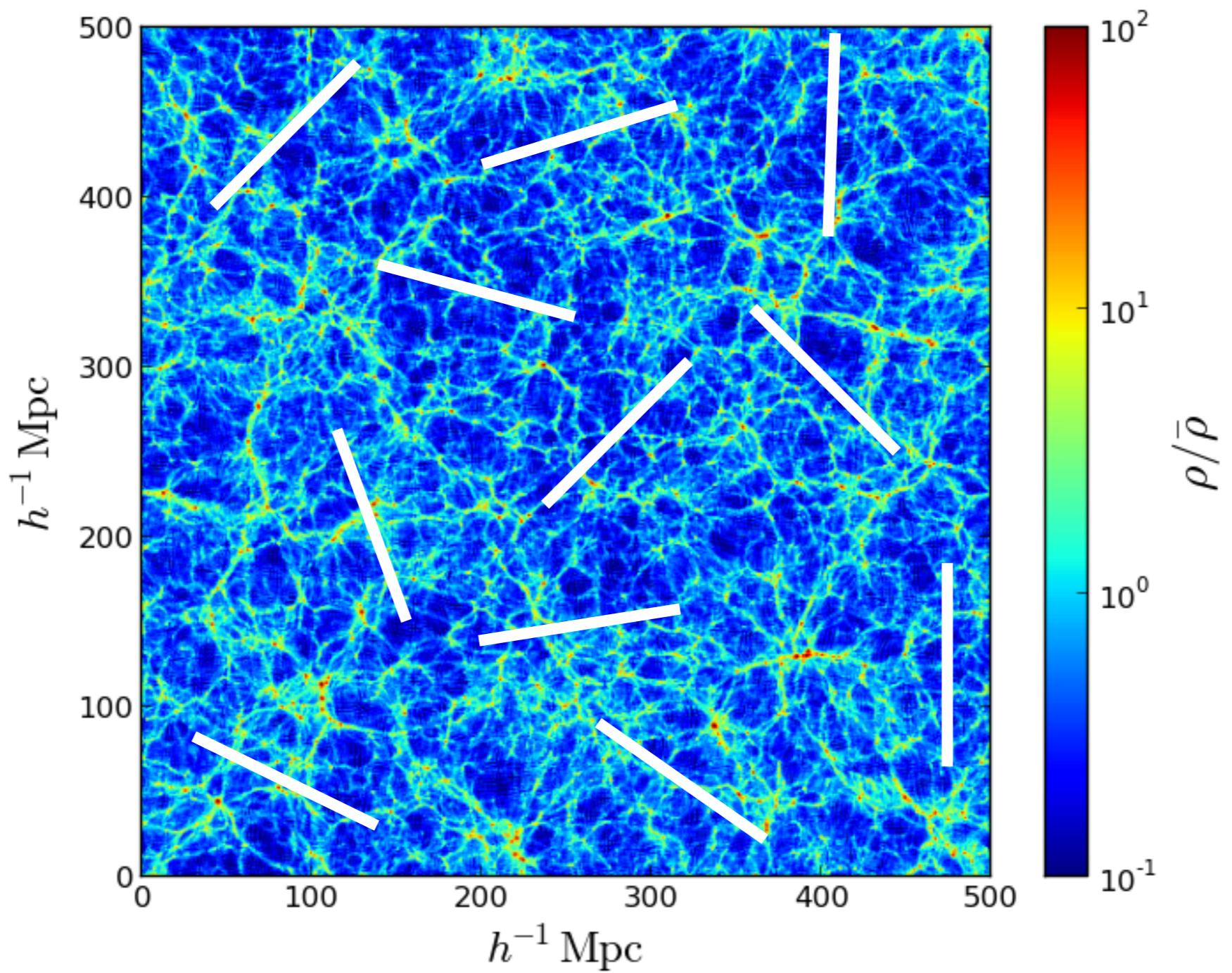
Should we build observables related
to low density regions/voids?

Void Size Function,
void-matter cross-correlation, ...

How much do correlation functions/
power spectra depend on low density
regions?

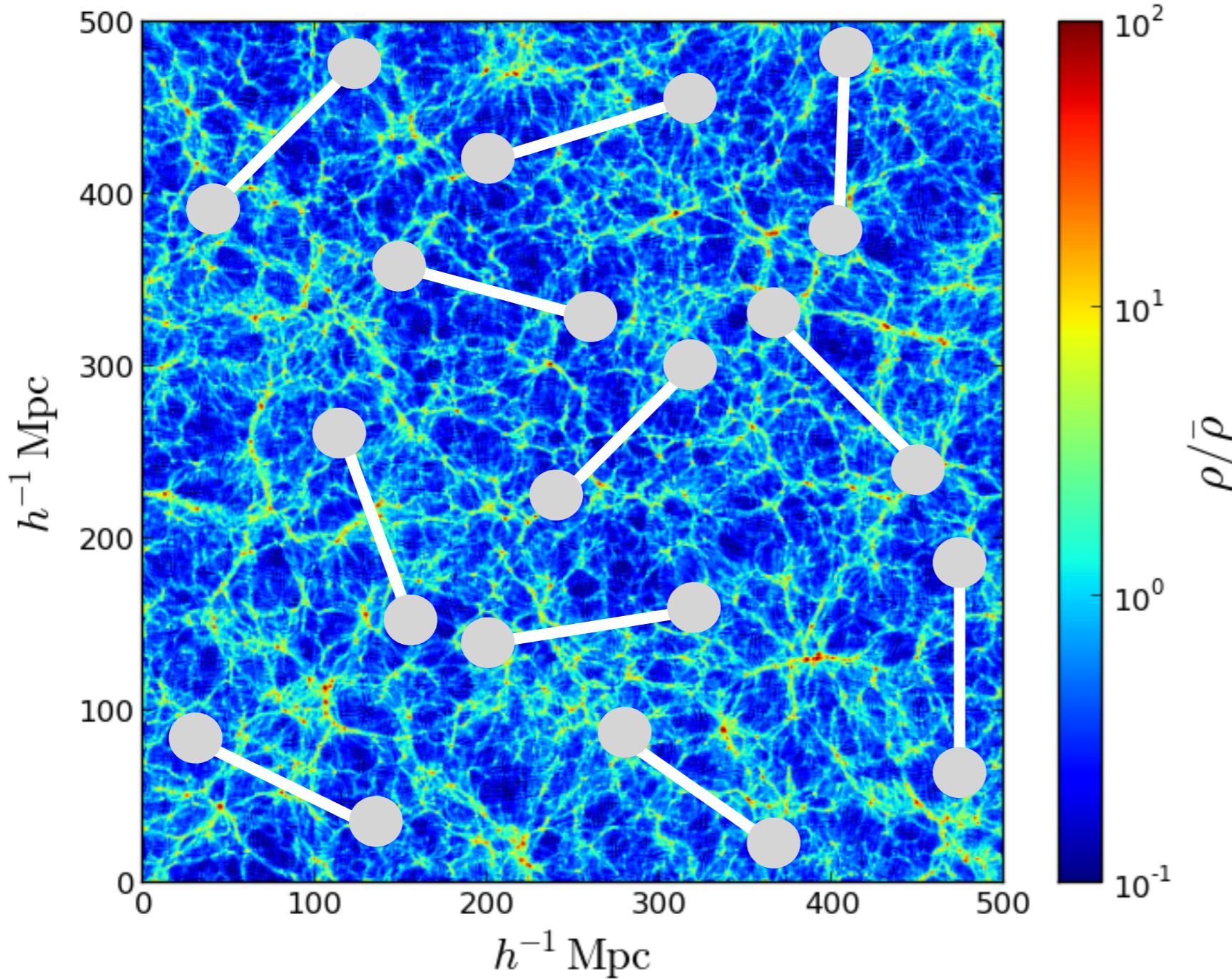
CORRELATION FUNCTION

$$1 + \xi(r) = \frac{V}{N^2} \sum_{i,j=1}^N \delta(|\vec{x}_i - \vec{x}_j| - r)$$



MARKED CORRELATION FUNCTION

$$1 + M(r, \phi) = \frac{V}{N^2} \sum_{i,j=1}^N \frac{\delta(|\vec{x}_i - \vec{x}_j| - r) m(\vec{x}_i, \phi) m(\vec{x}_j, \phi)}{\bar{m}^2}$$



We need m to up-weight
low density regions
and
down-weight high density
regions

m must depend on the local
density around each point

MARKED CORRELATION FUNCTION

$$1 + M(r, \phi) = \frac{V}{N^2} \sum_{i,j=1}^N \frac{\delta(|\vec{x}_i - \vec{x}_j| - r) m(\vec{x}_i, \phi) m(\vec{x}_j, \phi)}{\bar{m}^2}$$

$$m(\vec{x}, \phi = R, p, \delta_s) = \left[\frac{1 + \delta_s}{1 + \delta_s + \delta_R(\vec{x})} \right]^p$$

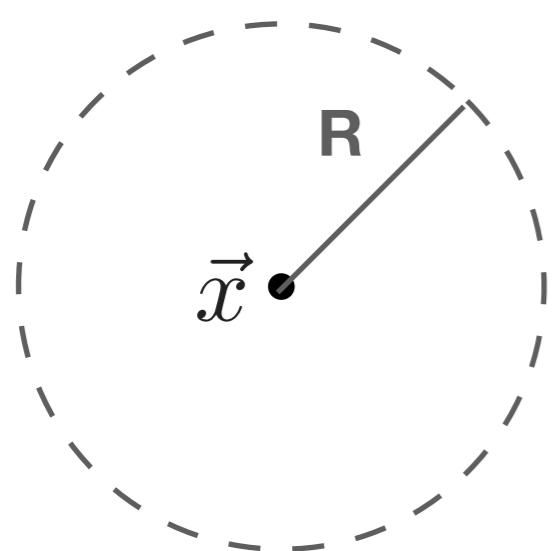
(M. White 2016)

MARKED CORRELATION FUNCTION

$$1 + M(r, \phi) = \frac{V}{N^2} \sum_{i,j=1}^N \frac{\delta(|\vec{x}_i - \vec{x}_j| - r) m(\vec{x}_i, \phi) m(\vec{x}_j, \phi)}{\bar{m}^2}$$

$$m(\vec{x}, \phi = R, p, \delta_s) = \left[\frac{1 + \delta_s}{1 + \delta_s + \delta_R(\vec{x})} \right]^p$$

(M. White 2016)



$$\delta_R(\vec{x}) = \frac{1}{V_R} \int_{V_R} d^3y \delta(\vec{y})$$

MARKED CORRELATION FUNCTION

$$1 + M(r, \phi) = \frac{V}{N^2} \sum_{i,j=1}^N \frac{\delta(|\vec{x}_i - \vec{x}_j| - r) m(\vec{x}_i, \phi) m(\vec{x}_j, \phi)}{\bar{m}^2}$$

$$m(\vec{x}, \phi = R, p, \delta_s) = \left[\frac{1 + \delta_s}{1 + \delta_s + \delta_R(\vec{x})} \right]^p$$

(M. White 2016)

$$m(\vec{x}, \phi = R, p, \delta_s \rightarrow 0) \rightarrow \left[\frac{\bar{\rho}}{\rho_R(\vec{x})} \right]^p$$

MARKED CORRELATION FUNCTION

$$1 + M(r, \phi) = \frac{V}{N^2} \sum_{i,j=1}^N \frac{\delta(|\vec{x}_i - \vec{x}_j| - r) m(\vec{x}_i, \phi) m(\vec{x}_j, \phi)}{\bar{m}^2}$$

$$m(\vec{x}, \phi = R, p, \delta_s) = \left[\frac{1 + \delta_s}{1 + \delta_s + \delta_R(\vec{x})} \right]^p$$

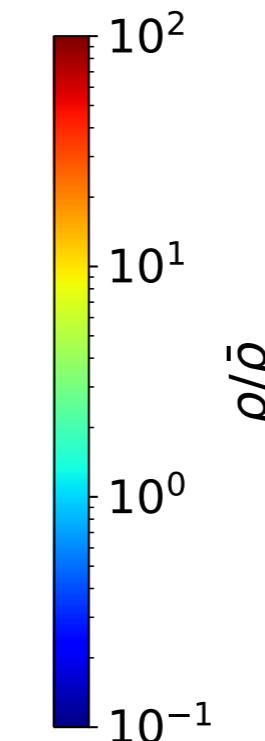
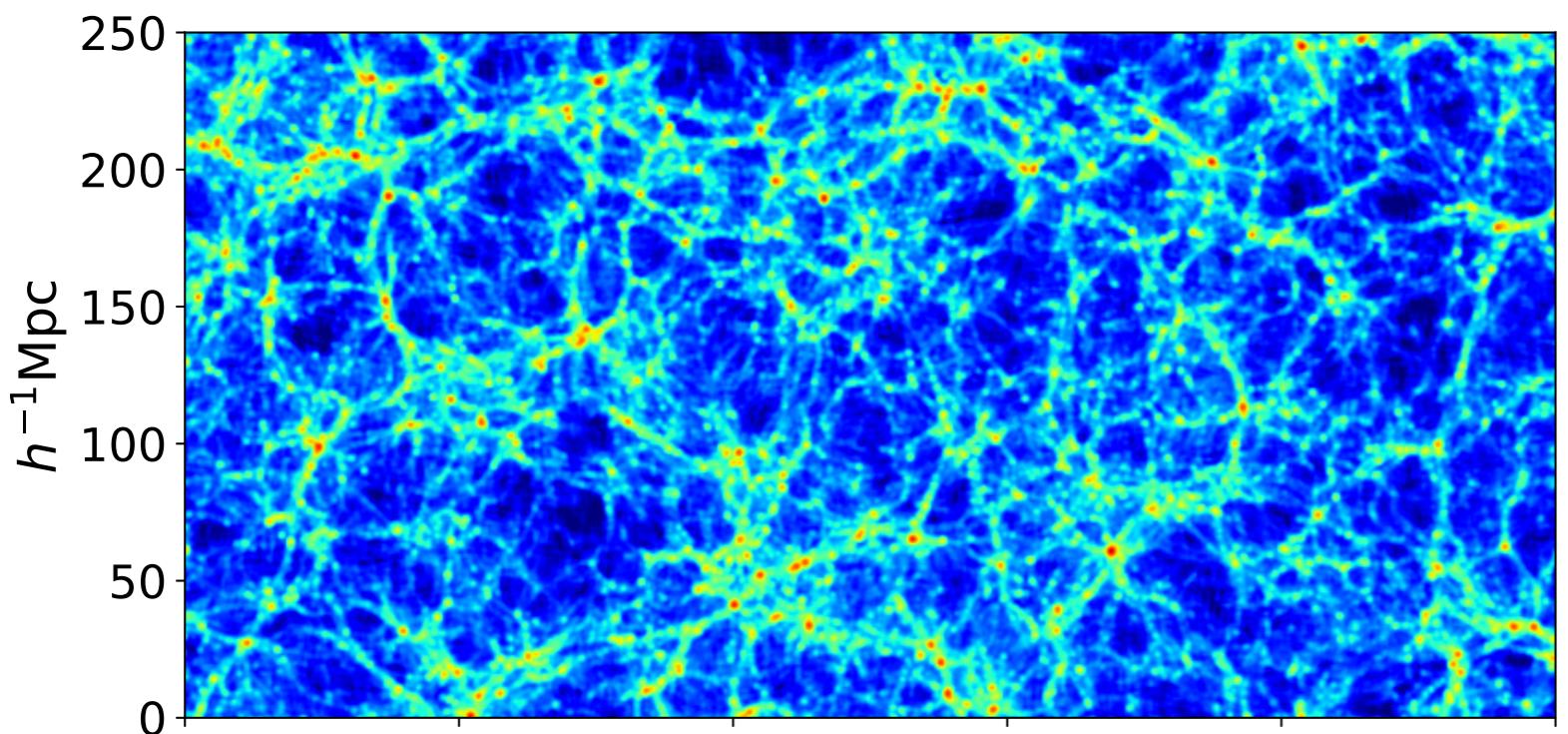
(M. White 2016)

$$m(\vec{x}, \phi = R, p, \delta_s \rightarrow 0) \rightarrow \left[\frac{\bar{\rho}}{\rho_R(\vec{x})} \right]^p$$

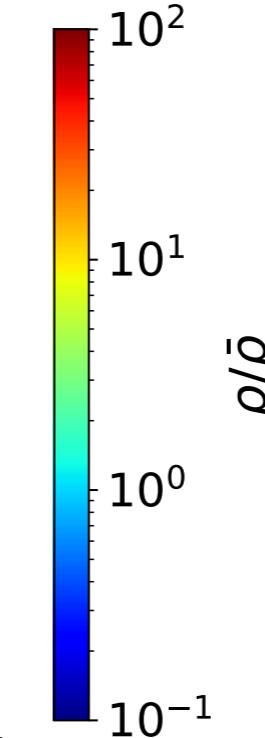
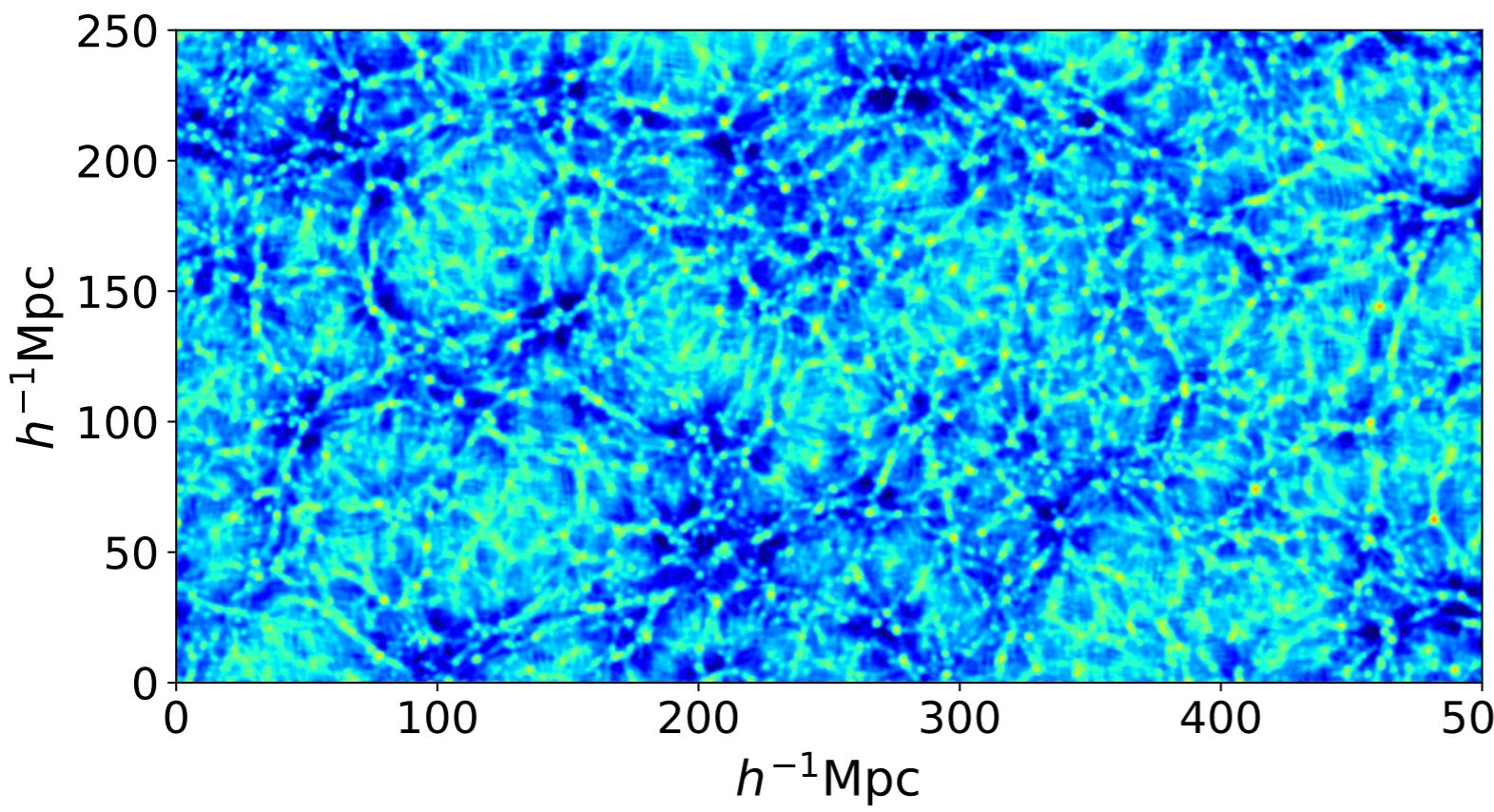
- | | |
|---------|---|
| $p > 0$ | Weight more particles/galaxies in UNDER -densities |
| $p < 0$ | Weight more particles/galaxies in OVER -densities |

MARKED DENSITY FIELD

EM et al. 2020



Density field



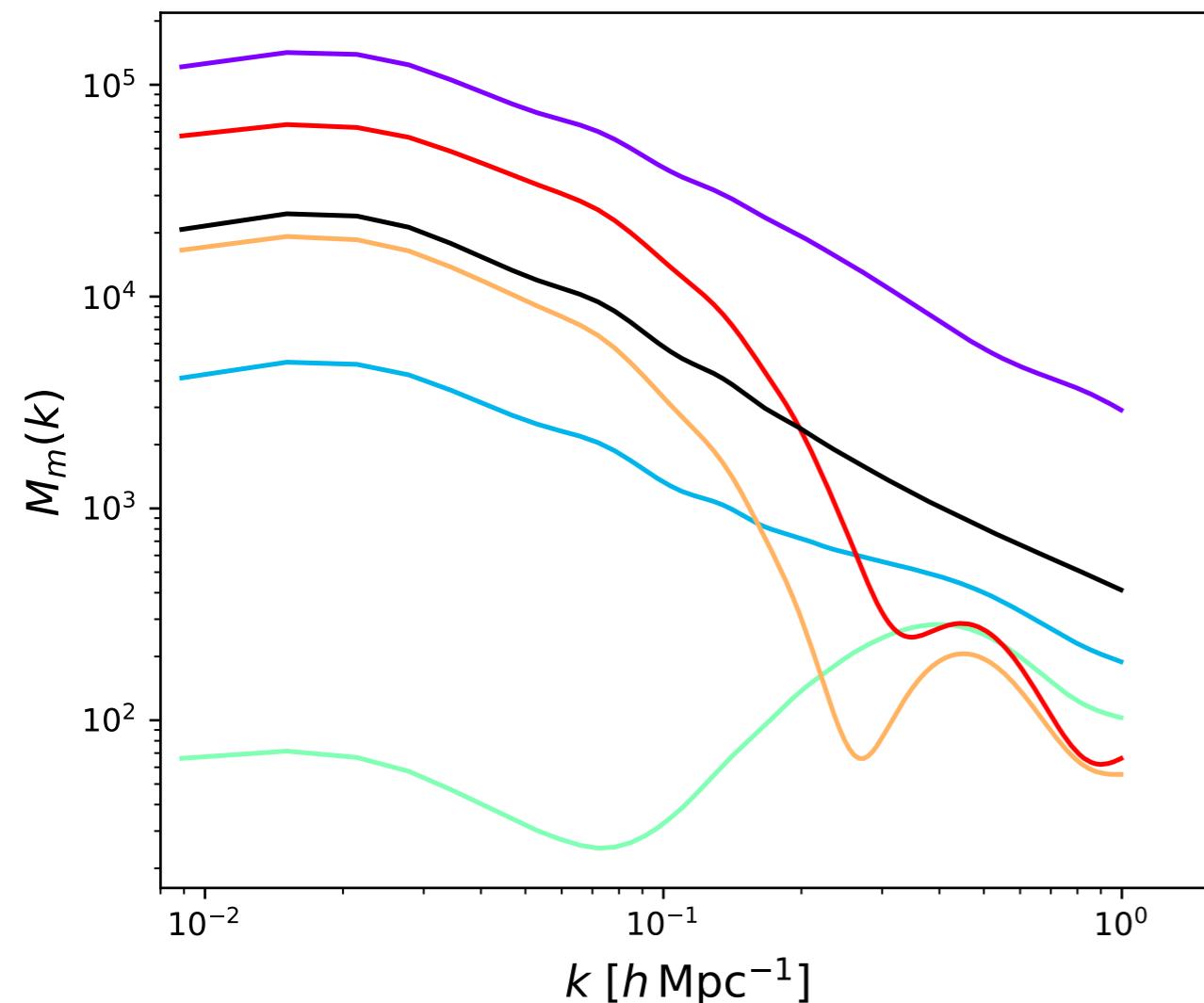
**Marked density field
with $p>0$**

MARKED POWER SPECTRUM

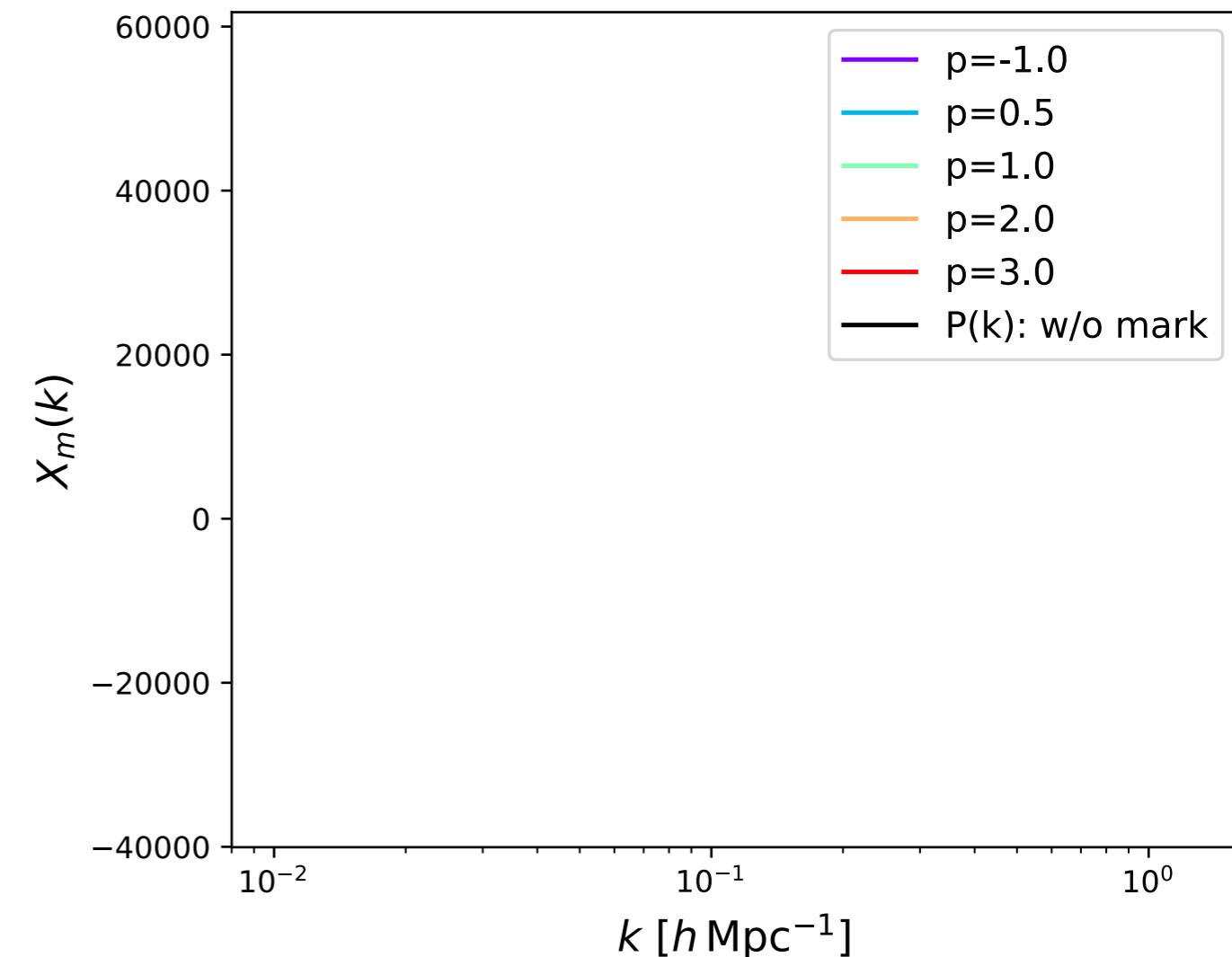
M(k) with $R = 10 \text{ Mpc}/h$, $\delta_s = 0.0$

EM et al. 2020

Marked power spectrum



Marked-standard density cross-power spectrum



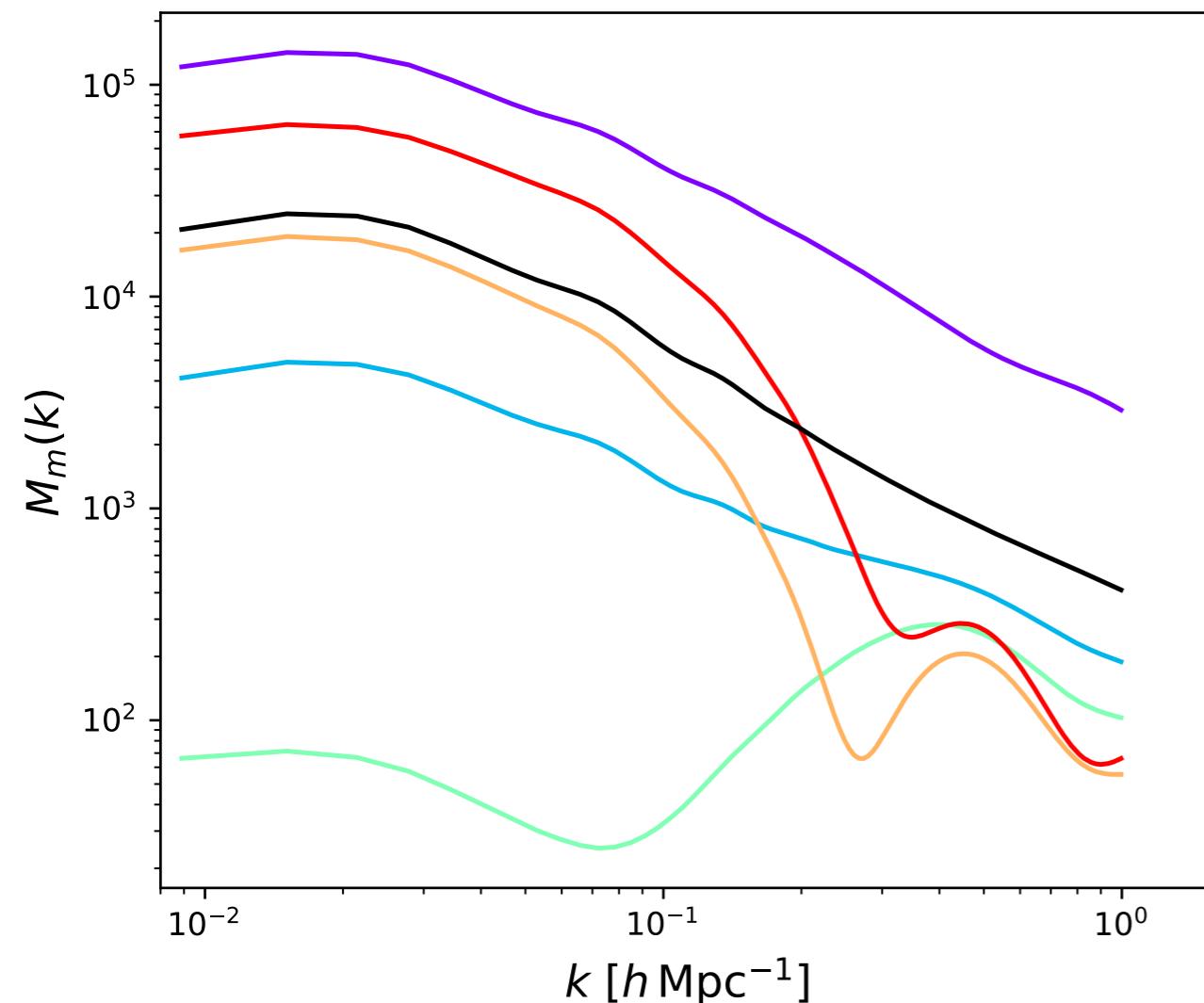
On large scales, an effective bias can describe the difference between standard and marked power spectra

MARKED POWER SPECTRUM

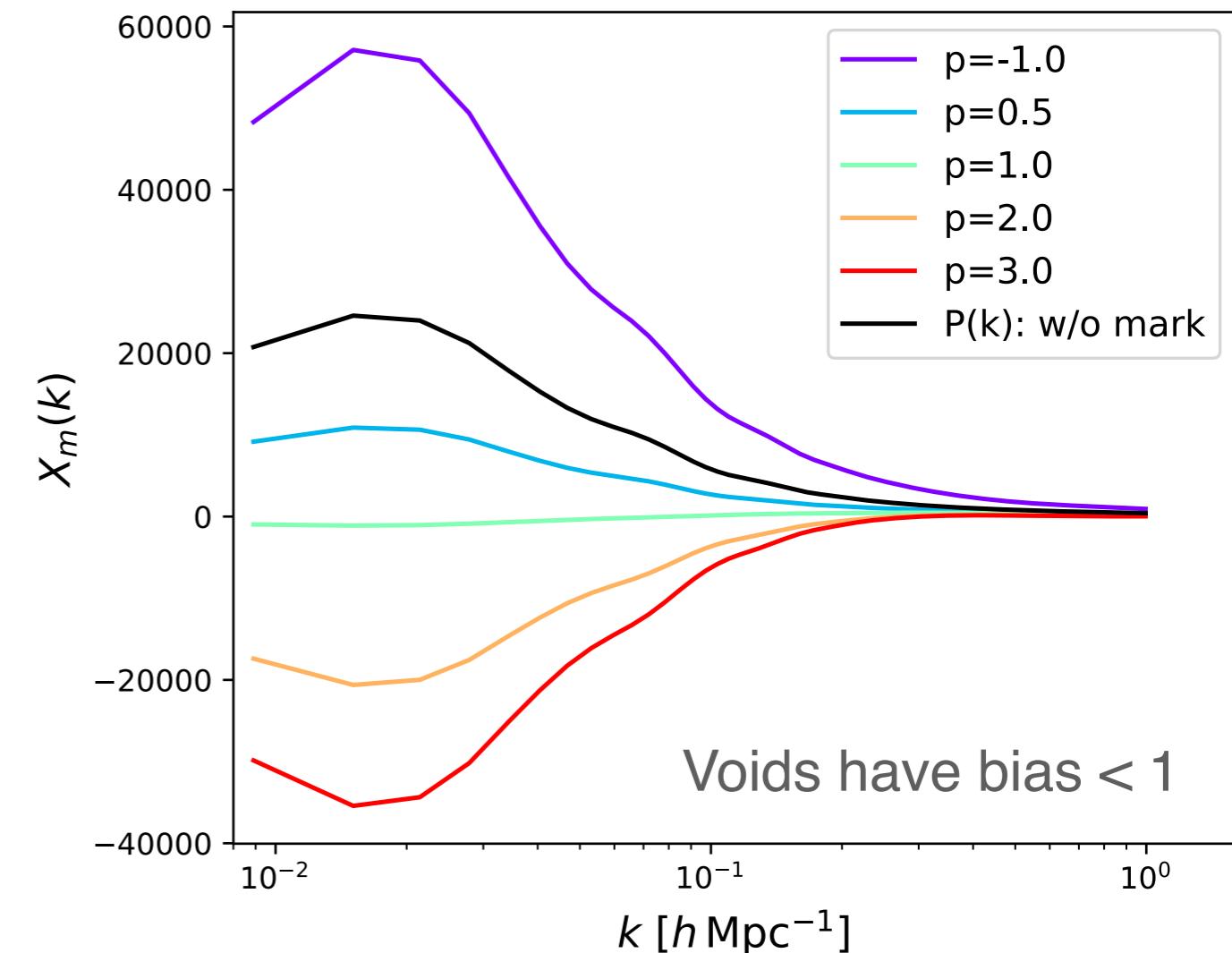
M(k) with $R = 10 \text{ Mpc}/h$, $\delta_s = 0.0$

EM et al. 2020

Marked power spectrum



Marked-standard density cross-power spectrum



On large scales, an effective bias can describe the difference between standard and marked power spectra

INFORMATION CONTENT IN MARKED POWER SPECTRA

FISHER ANALYSIS

Cosmological parameters: $\vec{\theta} = \{\Omega_m, \Omega_b, h, n_s, \sigma_8, M_\nu\}$

Data vector (observables): $\vec{d} = \{P(k_1), P(k_2), \dots, P(k_n)\}$

Error on each parameter:

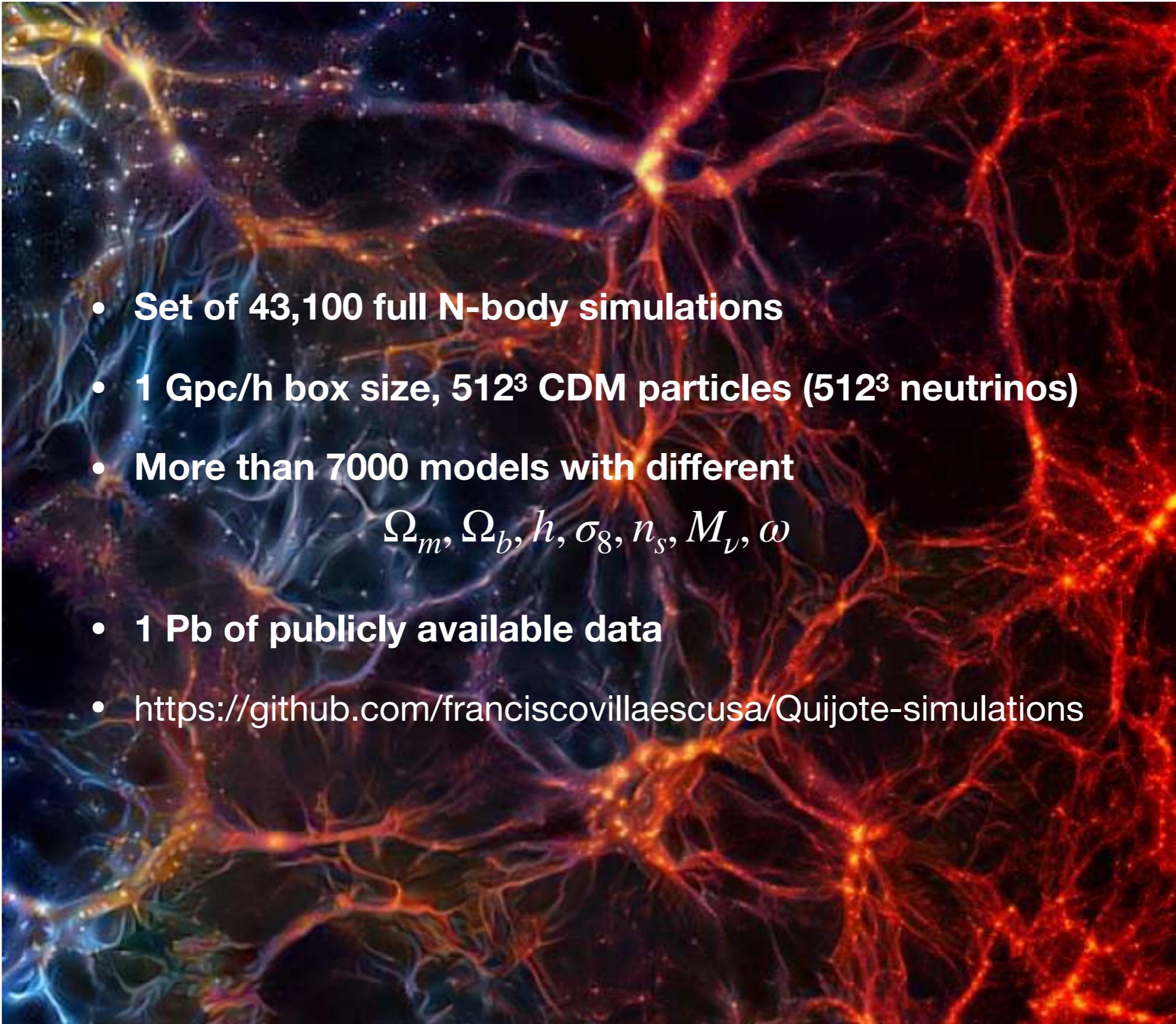
$$\sigma(\theta_\alpha) \leq \sqrt{(F^{-1})_{\alpha\alpha}}$$

Fisher matrix:

$$F_{\alpha,\beta} = \frac{\partial \vec{d}}{\partial \theta_\alpha} C^{-1} \frac{\partial \vec{d}}{\partial \theta_\beta}$$

QUIJOTE SIMULATIONS

Villaescusa-Navarro, Hanh, EM et al 2019



- Set of 43,100 full N-body simulations
- 1 Gpc/h box size, 512^3 CDM particles (512^3 neutrinos)
- More than 7000 models with different
 $\Omega_m, \Omega_b, h, \sigma_8, n_s, M_\nu, \omega$
- 1 Pb of publicly available data
- <https://github.com/franciscovillaescusa/Quijote-simulations>

QUIJOTE SIMULATIONS

Villaescusa-Navarro, Hanh, EM et al 2019

Name	Ω_m	Ω_b	h	n_s	σ_8	M_ν (eV)	w	realizations	simulations	ICs	$N_c^{1/3}$	$N_\nu^{1/3}$
Fid	<u>0.3175</u>	<u>0.049</u>	<u>0.6711</u>	<u>0.9624</u>	<u>0.834</u>	<u>0</u>	<u>-1</u>	15000 500 500 1000 100	standard standard paired fixed standard standard	2LPT Zeldovich 2LPT 2LPT 2LPT	512 512 512 256 1024	0 0 0 0 0
Ω_m^+	<u>0.3275</u>	0.049	0.6711	0.9624	0.834	0	-1	500 500	standard paired fixed	2LPT	512	0
Ω_m^-	<u>0.3075</u>	0.049	0.6711	0.9624	0.834	0	-1	500 500	standard paired fixed	2LPT	512	0
Ω_b^{++}	0.3175	<u>0.051</u>	0.6711	0.9624	0.834	0	-1	500 500	standard paired fixed	2LPT	512	0
Ω_b^+	0.3175	<u>0.050</u>	0.6711	0.9624	0.834	0	-1	500	paired fixed	2LPT	512	0
Ω_b^-	0.3175	<u>0.048</u>	0.6711	0.9624	0.834	0	-1	500	paired fixed	2LPT	512	0
Ω_b^{--}	0.3175	<u>0.047</u>	0.6711	0.9624	0.834	0	-1	500 500	standard paired fixed	2LPT	512	0
h^+	0.3175	0.049	<u>0.6911</u>	0.9624	0.834	0	-1	500 500	standard paired fixed	2LPT	512	0
h^-	0.3175	0.049	<u>0.6511</u>	0.9624	0.834	0	-1	500 500	standard paired fixed	2LPT	512	0
n_s^+	0.3175	0.049	0.6711	<u>0.9824</u>	0.834	0	-1	500 500	standard paired fixed	2LPT	512	0
n_s^+	0.3175	0.049	0.6711	<u>0.9424</u>	0.834	0	-1	500 500	standard paired fixed	2LPT	512	0
σ_8^+	0.3175	0.049	0.6711	0.9624	<u>0.849</u>	0	-1	500 500	standard paired fixed	2LPT	512	0
σ_8^-	0.3175	0.049	0.6711	0.9624	<u>0.819</u>	0	-1	500 500	standard paired fixed	2LPT	512	0
M_ν^{+++}	0.3175	0.049	0.6711	0.9624	0.834	<u>0.4</u>	-1	500 500	standard paired fixed	Zeldovich	512	512
M_ν^{++}	0.3175	0.049	0.6711	0.9624	0.834	<u>0.2</u>	-1	500 500	standard paired fixed	Zeldovich	512	512
M_ν^+	0.3175	0.049	0.6711	0.9624	0.834	<u>0.1</u>	-1	500 500	standard paired fixed	Zeldovich	512	512

QUIJOTE SIMULATIONS

Villaescusa-Navarro, Hanh, EM et al 2019

Name	Ω_m	Ω_b	h	n_s	σ_8	$M_\nu(\text{eV})$	w	realizations	simulations	ICs	$N_c^{1/3}$	$N_\nu^{1/3}$
Fid	<u>0.3175</u>	<u>0.049</u>	<u>0.6711</u>	<u>0.9624</u>	<u>0.834</u>	<u>0</u>	<u>-1</u>	15000 500 500 1000 100	standard standard paired fixed standard standard	2LPT Zeldovich 2LPT 2LPT 2LPT	512 512 512 256 1024	0 0 0 0 0
Ω_m^+	<u>0.3275</u>	0.049	0.6711	0.9624	0.834	0	-1	500 500	standard paired fixed	2LPT	512	0
Ω_m^-	<u>0.3075</u>	0.049	0.6711	0.9624	0.834	0	-1	500 500	standard paired fixed	2LPT	512	0
Ω_b^{++}	0.3175	<u>0.051</u>	0.6711	0.9624	0.834	0	-1	500 500	standard paired fixed	2LPT	512	0
Ω_b^+	0.3175	<u>0.050</u>	0.6711	0.9624	0.834	0	-1	500	paired fixed	2LPT	512	0
Ω_b^-	0.3175	<u>0.048</u>	0.6711	0.9624	0.834	0	-1	500	paired fixed	2LPT	512	0
Ω_b^{--}	0.3175	<u>0.047</u>	0.6711	0.9624	0.834	0	-1	500 500	standard paired fixed	2LPT	512	0
h^+	0.3175	0.049	<u>0.6911</u>	0.9624	0.834	0	-1	500 500	standard paired fixed	2LPT	512	0
h^-	0.3175	0.049	<u>0.6511</u>	0.9624	0.834	0	-1	500 500	standard paired fixed	2LPT	512	0
n_s^+	0.3175	0.049	0.6711	<u>0.9824</u>	0.834	0	-1	500 500	standard paired fixed	2LPT	512	0
n_s^+	0.3175	0.049	0.6711	<u>0.9424</u>	0.834	0	-1	500 500	standard paired fixed	2LPT	512	0
σ_8^+	0.3175	0.049	0.6711	0.9624	<u>0.849</u>	0	-1	500 500	standard paired fixed	2LPT	512	0
σ_8^-	0.3175	0.049	0.6711	0.9624	<u>0.819</u>	0	-1	500 500	standard paired fixed	2LPT	512	0
M_ν^{+++}	0.3175	0.049	0.6711	0.9624	0.834	<u>0.4</u>	-1	500 500	standard paired fixed	Zeldovich	512	512
M_ν^{++}	0.3175	0.049	0.6711	0.9624	0.834	<u>0.2</u>	-1	500 500	standard paired fixed	Zeldovich	512	512
M_ν^+	0.3175	0.049	0.6711	0.9624	0.834	<u>0.1</u>	-1	500 500	standard paired fixed	Zeldovich	512	512

QUIJOTE SIMULATIONS

Villaescusa-Navarro, Hanh, EM et al 2019

Name	Ω_m	Ω_b	h	n_s	σ_8	M_ν (eV)	w	realizations	simulations	ICs	$N_c^{1/3}$	$N_\nu^{1/3}$
Fid	<u>0.3175</u>	<u>0.049</u>	<u>0.6711</u>	<u>0.9624</u>	<u>0.834</u>	<u>0</u>	<u>-1</u>	15000 500 500 1000 100	standard standard paired fixed standard standard	2LPT Zeldovich 2LPT 2LPT 2LPT	512 512 512 256 1024	0 0 0 0 0
Ω_m^+	<u>0.3275</u>	0.049	0.6711	0.9624	0.834	0	-1	500 500	standard paired fixed	2LPT	512	0
Ω_m^-	<u>0.3075</u>	0.049	0.6711	0.9624	0.834	0	-1	500 500	standard paired fixed	2LPT	512	0
Ω_b^{++}	0.3175	<u>0.051</u>	0.6711	0.9624	0.834	0	-1	500 500	standard paired fixed	2LPT	512	0
Ω_b^+	0.3175	<u>0.050</u>	0.6711	0.9624	0.834	0	-1	500	paired fixed	2LPT	512	0
Ω_b^-	0.3175	<u>0.048</u>	0.6711	0.9624	0.834	0	-1	500	paired fixed	2LPT	512	0
Ω_b^{--}	0.3175	<u>0.047</u>	0.6711	0.9624	0.834	0	-1	500 500	standard paired fixed	2LPT	512	0
h^+	0.3175	0.049	<u>0.6911</u>	0.9624	0.834	0	-1	500 500	standard paired fixed	2LPT	512	0
h^-	0.3175	0.049	<u>0.6511</u>	0.9624	0.834	0	-1	500 500	standard paired fixed	2LPT	512	0
n_s^+	0.3175	0.049	0.6711	<u>0.9824</u>	0.834	0	-1	500 500	standard paired fixed	2LPT	512	0
n_s^+	0.3175	0.049	0.6711	<u>0.9424</u>	0.834	0	-1	500 500	standard paired fixed	2LPT	512	0
σ_8^+	0.3175	0.049	0.6711	0.9624	<u>0.849</u>	0	-1	500 500	standard paired fixed	2LPT	512	0
σ_8^-	0.3175	0.049	0.6711	0.9624	<u>0.819</u>	0	-1	500 500	standard paired fixed	2LPT	512	0
M_ν^{+++}	0.3175	0.049	0.6711	0.9624	0.834	<u>0.4</u>	-1	500 500	standard paired fixed	Zeldovich	512	512
M_ν^{++}	0.3175	0.049	0.6711	0.9624	0.834	<u>0.2</u>	-1	500 500	standard paired fixed	Zeldovich	512	512
M_ν^+	0.3175	0.049	0.6711	0.9624	0.834	<u>0.1</u>	-1	500 500	standard paired fixed	Zeldovich	512	512

MARKED POWER SPECTRA

EM et al. 2020

The Mark

$$m(\vec{x}, \phi = R, p, \delta_s) = \left[\frac{1 + \delta_s}{1 + \delta_s + \delta_R(\vec{x})} \right]^p$$

Considered values for the mark parameters

$$R = [5, 10, 15, 20, 30] \text{ Mpc/h}$$

$$p = [-1, 0.5, 1, 2, 3]$$

$$\delta_s = [0, 0.25, 0.5, 0.75, 1]$$

125 marked power spectra compute on the matter fields
cdm and m = cdm+neutrinos

MARKED POWER SPECTRA

EM et al. 2020

The Mark

$$m(\vec{x}, \phi = R, p, \delta_s) = \left[\frac{1 + \delta_s}{1 + \delta_s + \delta_R(\vec{x})} \right]^p$$

Considered values for the mark parameters

$$R = [5, 10, 15, 20, 30] \text{ Mpc/h}$$

$$p = [-1, 0.5, 1, 2, 3]$$

$$\delta_s = [0, 0.25, 0.5, 0.75, 1]$$



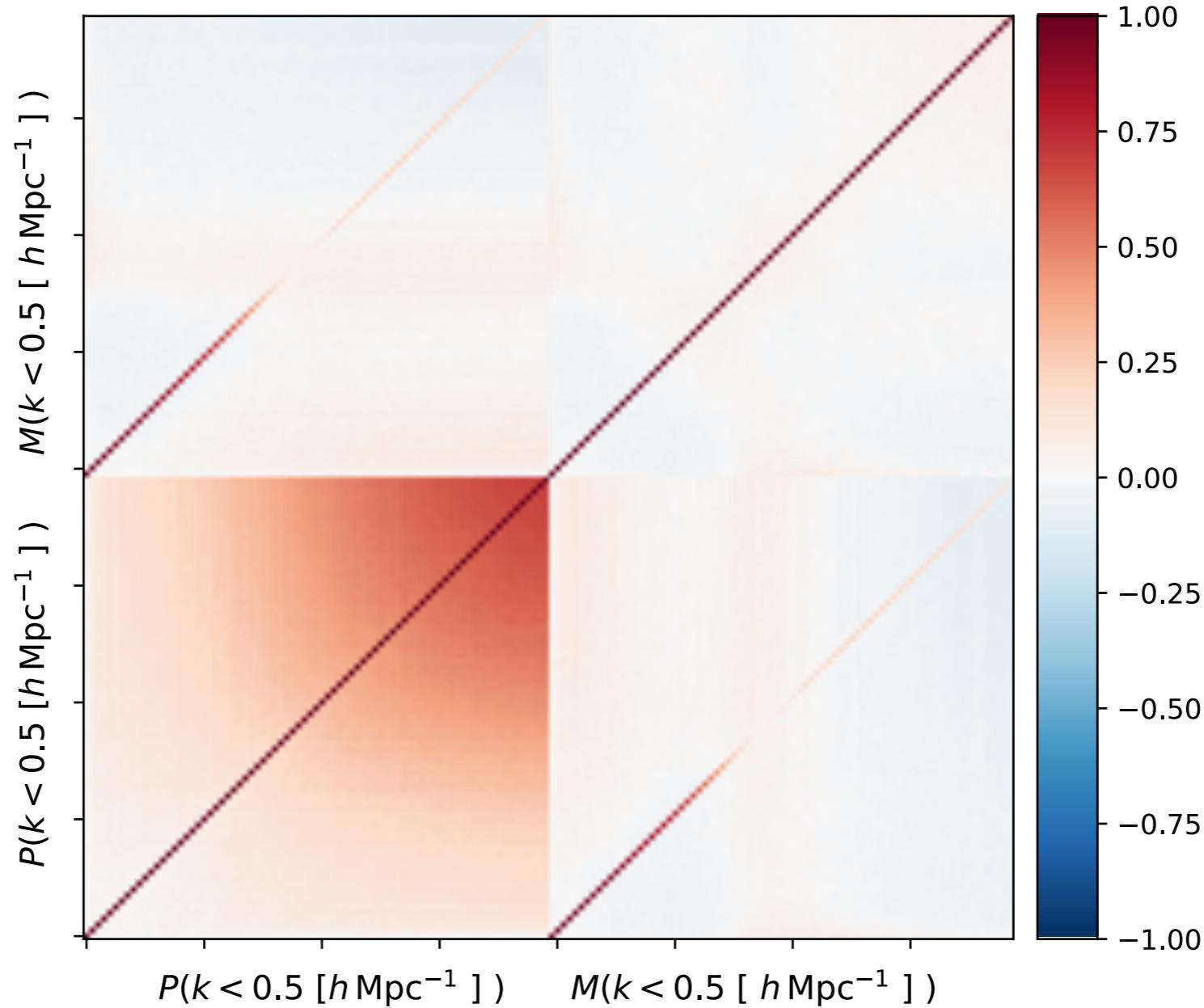
= model giving the tightest constraint on the neutrino masses

125 marked power spectra compute on the matter fields
cdm and m = cdm+neutrinos

Correlation matrix

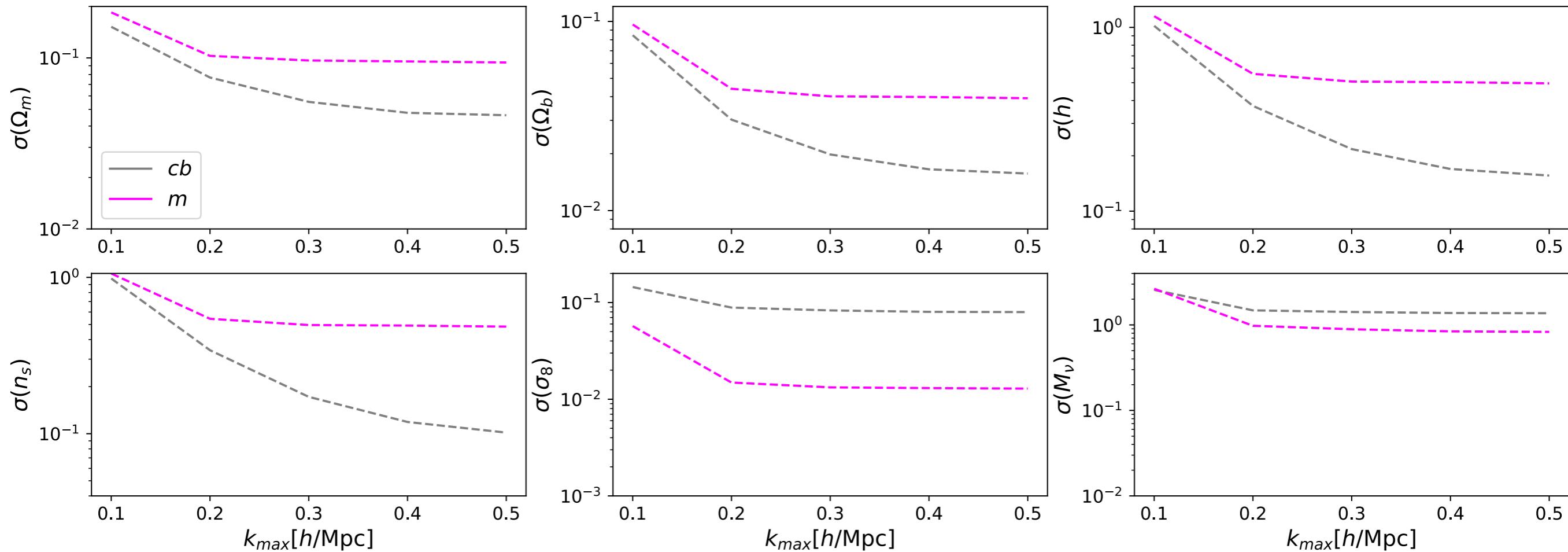
EM et al. 2020

$M(k)$ with $R = 10 \text{ Mpc}/h$, $p = 2$, $\delta_s = 0.25$



Marginalized errors

EM et al. 2020

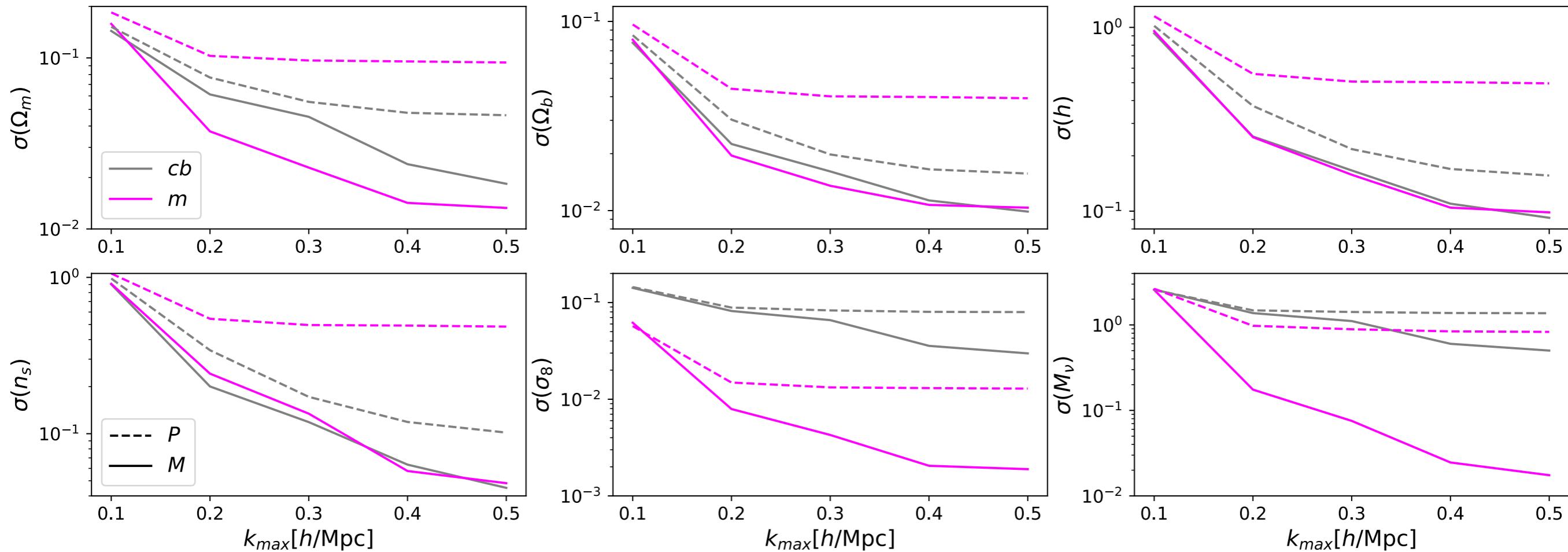


Marginalized errors for $k_{\max} = 0.5 \text{ h/Mpc}$

Parameter	P_{cb}	M_{cb}	$P_{cb} + M_{cb}$	$M_{cb} + M'_{cb}$	P_m	M_m	$P_m + M_m$	$M_m + M'_m$
Ω_m	0.046				0.094			
Ω_b	0.016				0.039			
h	0.16				0.50			
n_s	0.10				0.48			
σ_8	0.080				0.013			
M_ν	1.4				0.83			

Marginalized errors

EM et al. 2020

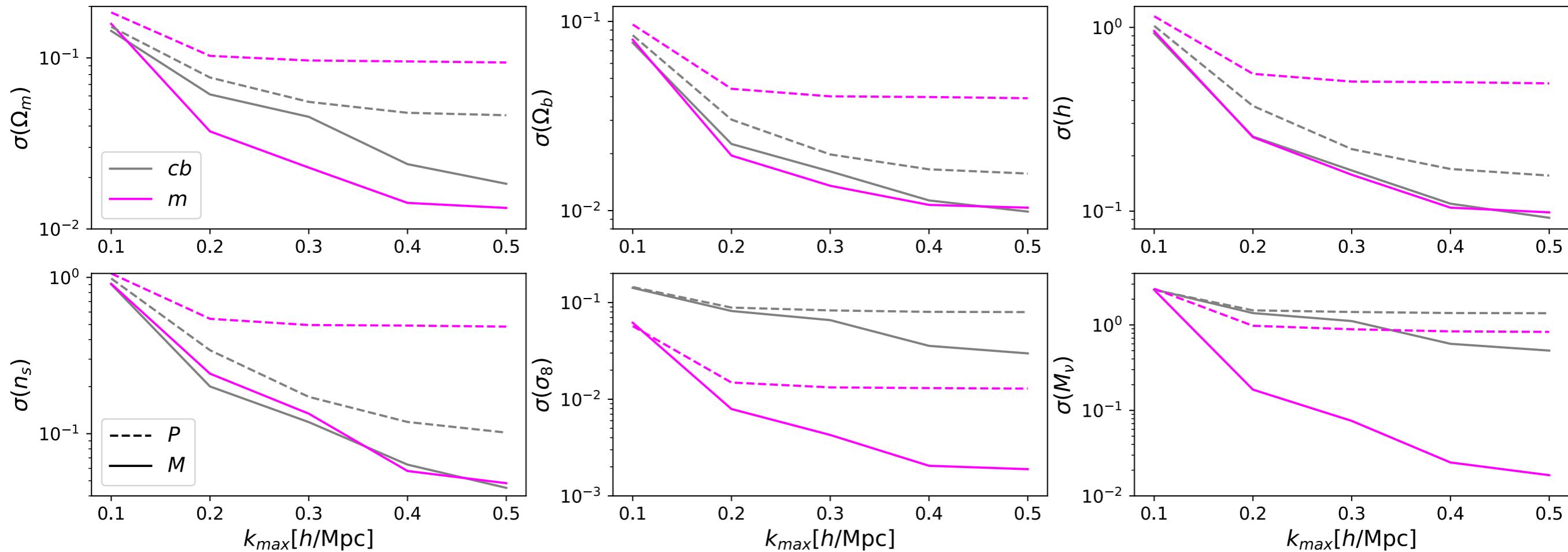


Marginalized errors for $k_{\max} = 0.5 \text{ h/Mpc}$

Parameter	P_{cb}	M_{cb}	$P_{cb} + M_{cb}$	$M_{cb} + M'_{cb}$	P_m	M_m	$P_m + M_m$	$M_m + M'_m$
Ω_m	0.046	0.018			0.094	0.013		
Ω_b	0.016	0.0099			0.039	0.010		
h	0.16	0.092			0.50	0.098		
n_s	0.10	0.045			0.48	0.048		
σ_8	0.080	0.030			0.013	0.0019		
M_ν	1.4	0.50			0.83	0.017		

Marginalized errors

EM et al. 2020



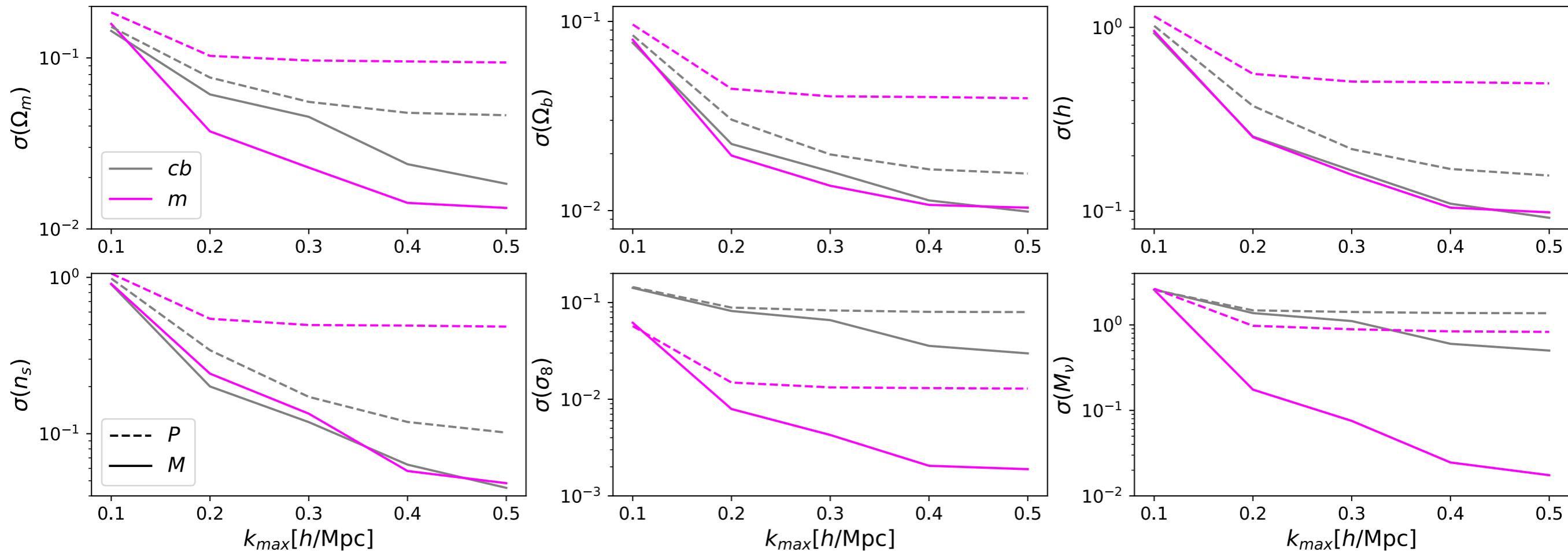
Marginalized errors for $k_{\max} = 0.5 \text{ h/Mpc}$

Parameter	P_{cb}	M_{cb}	$P_{cb} + M_{cb}$	$M_{cb} + M'_{cb}$	P_m	M_m	$P_m + M_m$	$M_m + M'_m$
Ω_m	0.046	0.018			0.094	0.013		
Ω_b	0.016	0.0099			0.039	0.010		
h	0.16	0.092			0.50	0.098		
n_s	0.10	0.045			0.48	0.048		
σ_8	0.080	0.030			0.013	0.0019		
M_ν	1.4	0.50			0.83	0.017		

3.5 σ detection of the minimum mass using 1 (Gpc/h)³ volume

Marginalized errors

EM et al. 2020



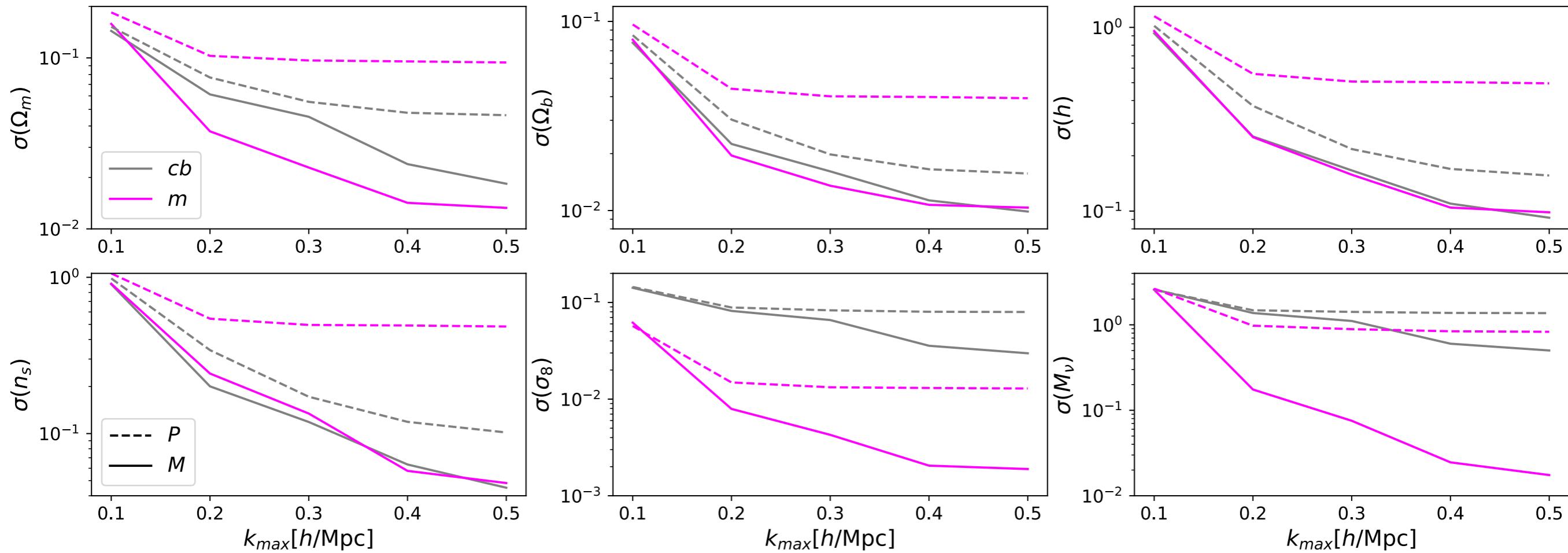
Marginalized errors for $k_{\max} = 0.5 \text{ h/Mpc}$

Parameter	P_{cb}	M_{cb}	$P_{cb} + M_{cb}$	$M_{cb} + M'_{cb}$	P_m	M_m	$P_m + M_m$	$M_m + M'_m$
Ω_m	0.046	0.018	0.017		0.094	0.013	0.012	
Ω_b	0.016	0.0099	0.0091		0.039	0.010	0.009	
h	0.16	0.092	0.083		0.50	0.098	0.082	
n_s	0.10	0.045	0.04		0.48	0.048	0.039	
σ_8	0.080	0.030	0.026		0.013	0.0019	0.0015	
M_ν	1.4	0.50	0.44		0.83	0.017	0.014	

4.2 σ detection of the minimum mass using 1 (Gpc/h)³ volume

Marginalized errors

EM et al. 2020



Marginalized errors for $k_{\max} = 0.5 \text{ h/Mpc}$

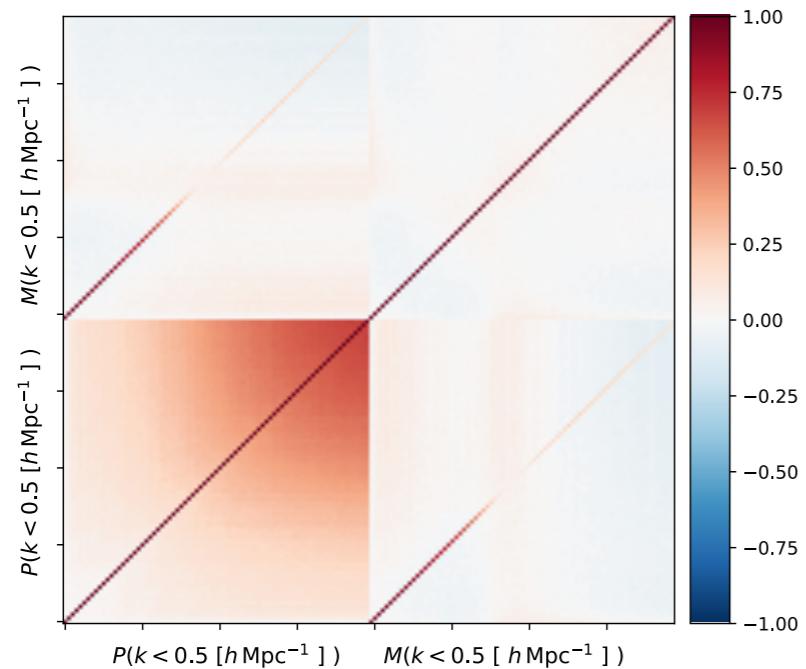
Parameter	P_{cb}	M_{cb}	$P_{cb} + M_{cb}$	$M_{cb} + M'_{cb}$	P_m	M_m	$P_m + M_m$	$M_m + M'_m$
Ω_m	0.046	0.018	0.017	0.014	0.094	0.013	0.012	0.011
Ω_b	0.016	0.0099	0.0091	0.008	0.039	0.010	0.009	0.008
h	0.16	0.092	0.083	0.068	0.50	0.098	0.082	0.069
n_s	0.10	0.045	0.04	0.029	0.48	0.048	0.039	0.028
σ_8	0.080	0.030	0.026	0.021	0.013	0.0019	0.0015	0.0015
M_ν	1.4	0.50	0.44	0.35	0.83	0.017	0.014	0.01

6 σ detection of the minimum mass using 1 (Gpc/h)³ volume

THE INFORMATION CONTENT

EM et al. 2020

- The covariance matrix of the considered marked power spectrum $M(k)$ is very diagonal

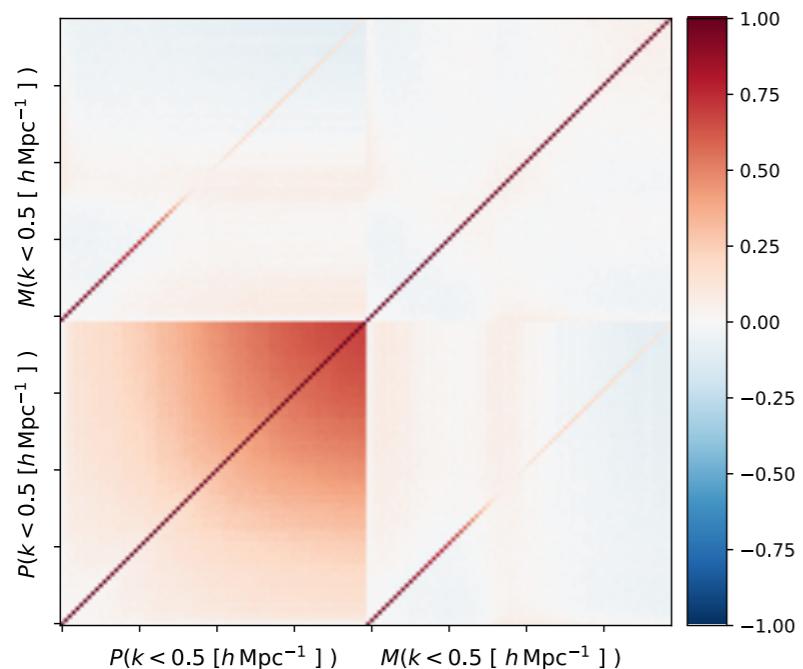


$$m(\vec{x}, \phi = R, p, \delta_s) = \left[\frac{1 + \delta_s}{1 + \delta_s + \delta_R(\vec{x})} \right]^p$$

THE INFORMATION CONTENT

EM et al. 2020

- The covariance matrix of the considered marked power spectrum $M(k)$ is very diagonal



- Marked density field is a nonlinear transformation of the density field
- Other nonlinear transformations, such as the log-transformation, have shown to make the field more Gaussian (Neyrinck et al, 2009, 2010, 2011)

$$m(\vec{x}, \phi = R, p, \delta_s) = \left[\frac{1 + \delta_s}{1 + \delta_s + \delta_R(\vec{x})} \right]^p$$

THE INFORMATION CONTENT

EM et al. 2020

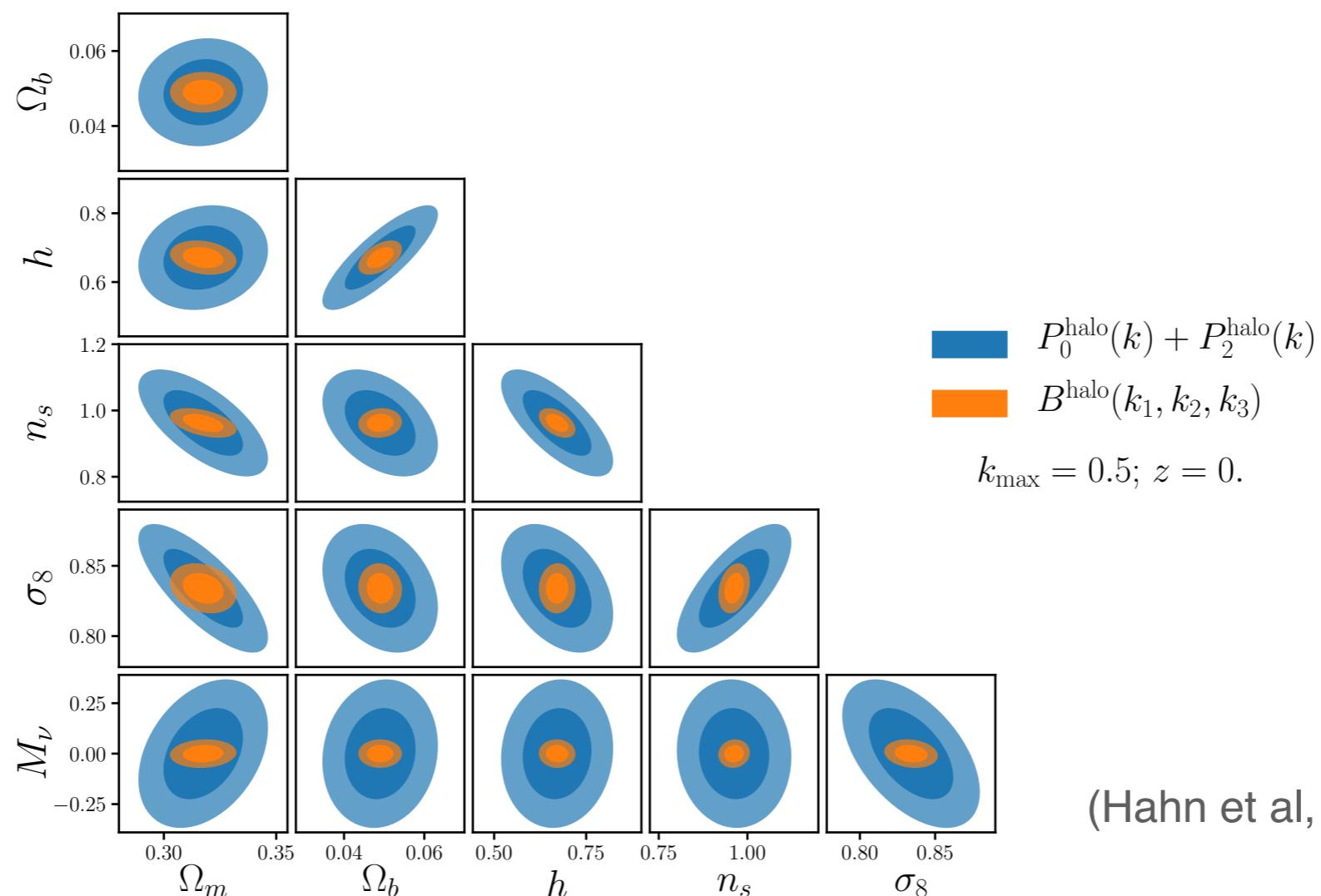
- The covariance matrix of the considered marked power spectrum $M(\mathbf{k})$ is very diagonal
- The marked power spectrum contains higher order statistics of the density field

$$m(\vec{x}, \phi = R, p, \delta_s) = \left[\frac{1 + \delta_s}{1 + \delta_s + \delta_R(\vec{x})} \right]^p$$

THE INFORMATION CONTENT

EM et al. 2020

- The covariance matrix of the considered marked power spectrum $M(k)$ is very diagonal
- The marked power spectrum contains higher order statistics of the density field



THE INFORMATION CONTENT

EM et al. 2020

- The covariance matrix of the considered marked power spectrum $M(k)$ is very diagonal
- The marked power spectrum contains higher order statistics of the density field
- The considered marked power spectrum incorporates information from voids

CONCLUSIONS

- We have studied marked power spectra computed on the MATTER density field

CONCLUSIONS

- We have studied marked power spectra computed on the MATTER density field
- Marked power spectra on the GALAXY density field will include BIAS and REDSHIFT SPACE DISTORTION, which will degrade the constraints on the sum of the neutrino masses.

CONCLUSIONS

- We have studied marked power spectra computed on the MATTER density field
- Marked power spectra on the GALAXY density field will include BIAS and REDSHIFT SPACE DISTORTION, which will degrade the constraints on the sum of the neutrino masses.
- Upcoming surveys (DESI, EUCLID, WFIRST) will have larger volumes, that will improve the constraints.

CONCLUSIONS

- We have studied marked power spectra computed on the MATTER density field
- Marked power spectra on the GALAXY density field will include BIAS and REDSHIFT SPACE DISTORTION, which will degrade the constraints on the sum of the neutrino masses.
- Upcoming surveys (DESI, EUCLID, WFIRST) will have larger volumes, that will improve the constraints.
- Mark power spectra are promising and easy-to-compute observables.

THE END



