Neutrinos, Quintessence, and Structure Formation in the Universe

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ML 1405.4855, 1602.08108 Hu, Chiang, Li, ML 1605.01412 Chiang, Li, Hu, ML 1609.01701 + in progress

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Outline

- Overview of Large-scale Structure
- Large-scale Structure Beyond Cold Dark Matter
- Solution of Structure
- The Fake Separate Universe Approach
- First Simulations Results: A proof of principle with quintessence
- Conclusion

Overview of Structure in the Universe





The Thermal History of the Universe

The Universe is expanding,

10-12	10-9	10-3	10-1	1
				"size" of universe compared to today

The Thermal History of the Universe

The Universe is expanding,

10 ⁻¹²	10 ⁻⁹	10-3	10-1	1
nd coolina.				"size" of universe compared to today
hotter 10 ¹³ K	10 ⁹ K	10,000 K	col	der 3 K
				temperature

The Thermal History of the Universe

The Universe is expanding,

	10-12	10 ⁻⁹	10-3	10-1	1
•					Size ["] of universe compared to today
and cooling.					
hott	er 10 ¹³ K	10 ⁹ K	10,000 K	cold	er 3 K
					temperature

and the matter in the Universe changes states as the universe cools



The Expansion History of the Universe



 10^{9}

 10^{7}

 10^{-5}

 10^{-4}

 10^{-3}

 10^{-1}

 10^{-2}

scale factor a

 10^{0}

The Expansion History of the Universe



And the different types of matter gravitate differently, so the expansion rate changes throughout the history of the universe



The History of Structure in the Universe

gravitational evolution of perturbations in cold dark matter density $\delta \rho / \rho$

 \sim

time

The History of Structure in the Universe

gravitational evolution of perturbations in cold dark matter density $\delta \rho / \rho$

time

Self-gravity of perturbations competes with expansion of the Universe

The History of Structure in the Universe

gravitational evolution of perturbations in cold dark matter density $\delta \rho / \rho$



The History of Structure in the Universe

gravitational evolution of perturbations in cold dark matter density $\delta \rho / \rho$



evolution of δρ/ρ depends on expansion rate of the Universe

The History of Structure in the Universe



evolution of δρ/ρ depends on expansion rate of the Universe



History of Structure in the Universe: Summary



(Kravtsov)



History of Structure in the Universe: Summary





History of Structure in the Universe: Summary



(Kravtsov)

evolution of density perturbations tells us about dark energy & types of matter present



Galaxy Distribution Mapped by SDSS



Many opportunities to learn about structure!

Euclid Large Synoptic Survey Dark Energy Survey Large-scale Telescope structure surveys Sloan Digital Sky Survey High resolution cosmic microwave background Dark Energy Spectroscopic Instrument deriments South Pole Telescope Subaru Hyper Suprime Cam and WFIRST Atacama Cosmology Prime Focus Spectrograph Telescope Planck - II ----Hobby-Eberly Telescope CMB ``Stage IV" Dark Energy EXperiment 21 CM experiments (CHIME, HIRAX) Simons Observatory SPHEREX

Many opportunities to learn about structure!



And much to learn! (inflation, dark energy, neutrino properties)

Recap: The evolution of structure depends on the expansion history of the Universe

(I) Massive Neutrinos

Large-scale Structure Beyond Cold Dark Matter (I) Massive Neutrinos



(TianNu simulation, Yu et al 2016)

Large-scale Structure Beyond Cold Dark Matter (I) Massive Neutrinos

 $\rho_{cdm} + \frac{\rho_{\nu}}{\rho_{baryon} + \rho_{\nu}}$ $\gtrsim 0.005$ $n_{\nu} \sim 10^{10} n_{baryon}$



(TianNu simulation, Yu et al 2016)

(II) Quintessence

(I) Quintessence

(scalar field dark energy)

 ρ dark energy ~ ρ matter



(Springel)

(I) Quintessence

(scalar field dark energy)

 ρ dark energy ~ ρ matter

could dark energy cluster too?



(Springel)

The Formation of (nonlinear) Structure

matter distribution $\delta_m = \delta \rho_m / \rho_m$



(Springel)

galaxy distribution $\delta_g = \delta n_g / n_g$





matter distribution $\delta_m = \delta \rho_m / \rho_m$



(Springel)

galaxy distribution $\delta_g = \delta n_g / n_g$



(Springel)

matter = cold dark matter (CDM), baryons (behaves like CDM, mostly), massive neutrinos, . . . ?!

matter distribution $\delta_m = \delta \rho_m / \rho_m$



(Springel)

galaxy distribution $\delta_g = \delta n_g / n_g$



⁽Springel)

Accurate models of the large-scale matter distribution, and the large-scale galaxy distribution are <u>crucial</u> for extracting cosmological information from cosmological datasets



when matter density fluctuations are small, can linearize the equations and gravitational evolution is easy







when matter density fluctuations become large (δρ/ρ ~1) gravity couples modes, evolution is hard! Need simulations or tricks
matter distribution $\delta_m = \delta \rho_m / \rho_m$



galaxy distribution $\delta_g = \delta n_g / n_g$



⁽Springel)

A particularly hard thing to understand is relationship between *galaxy* distribution and the *matter* distribution

matter distribution $\delta_m = \delta \rho_m / \rho_m$



galaxy distribution $\delta_g = \delta n_g / n_g$



(Springel)

Understanding the relationship between the galaxies and the matter is even harder when there is more than one type of matter around! (e.g. our universe! which, at least has neutrino dark matter and possibly dynamical dark energy or quintessence)

matter distribution $\delta_m = \delta \rho_m / \rho_m$



galaxy distribution $\delta_g = \delta n_g / n_g$

(Springel)

An intermediate step

matter distribution $\delta_m = \delta \rho_m / \rho_m$



galaxy distribution $\delta_g = \delta n_g / n_g$

galaxies live in halos (Springel)

An intermediate step

matter distribution $\delta_m = \delta \rho_m / \rho_m$



(Springel)

galaxy distribution $\delta_g = \delta n_g / n_g$



(Springel)

on large scales, halos and galaxy abundances vary with matter density, but the fractional over/under-densities are not identical $n_g(x) = \overline{n}_g(1 + \delta_g(x)), \ \rho_m(x) = \overline{\rho_m}(1 + \delta_m(x))$ $\delta_g \approx \frac{dlnn_g}{d\delta_m} \delta_m \approx b \delta_m$

matter distribution $\delta_m = \delta \rho_m / \rho_m$



(Springel)

galaxy distribution $\delta_g = \delta n_g / n_g$



(Snringel)

 $\boldsymbol{\delta}_{g} \approx \frac{d \ln n_{g}}{d \delta_{m}} \boldsymbol{\delta}_{m} \approx$ **b** δ_m

matter distribution $\delta_m = \delta \rho_m / \rho_m$



(Springel)

galaxy distribution $\delta_g = \delta n_g / n_g$



 $\boldsymbol{\delta}_{g} \approx \frac{dlnn_{g}}{d\delta_{m}} \boldsymbol{\delta}_{m}$ $\boldsymbol{b} \boldsymbol{\delta}_{m}$ ≈

the bias

matter distribution $\delta_m = \delta \rho_m / \rho_m$



(Springel)

galaxy distribution $\delta_g = \delta n_g / n_g$



(Snringel)

A parameter that depends on mass, luminosity or other properties of the tracer

$$\boldsymbol{\delta}_{g} \approx \frac{dlnn_{g}}{d\delta_{m}} \boldsymbol{\delta}_{m} \approx \boldsymbol{b} \boldsymbol{\delta}_{m}$$

matter distribution $\delta_m = \delta \rho_m / \rho_m$



(Springel)

galaxy distribution $\delta_g = \delta n_g / n_g$



*From now on use halo bias and galaxy and galaxy bias interchangeably , even though the truth is more complicated!



matter distribution $\delta_m = \delta \rho_m / \rho_m$



(Springel)

galaxy distribution $\delta_g = \delta n_g / n_g$



the bias $\delta_{g} \approx \frac{dlnn_{g}}{d\delta_{m}} \delta_{m} \approx b \delta_{m}$

Separate Universe Approach (To nonlinear evolution and bias)

The bias is a measure of the response of the number of halos to a long wavelength density fluctuation

$$\boldsymbol{\delta}_{h} \approx \frac{d \ln n_{h}}{d \delta_{m}} \boldsymbol{\delta}_{m} \approx \boldsymbol{b} \boldsymbol{\delta}_{m}$$

 $n_h(x) = \overline{n}_h(1 + \delta_h(x))$

 $\delta_m > 0$

$$n_h(x) = \overline{n}_h$$

$$\delta_{\rm m} = 0$$

The bias is a measure of the response of the number of halos to a long wavelength density fluctuation

$$\boldsymbol{\delta}_{h} \approx \frac{d \ln n_{h}}{d \delta_{m}} \boldsymbol{\delta}_{m} \approx \boldsymbol{b} \boldsymbol{\delta}_{m}$$

Usually measure by correlating halo fluctuation δ_h(x) with matter fluctuation δ_m(x)

 $n_h(x) = \overline{n}_h(1 + \delta_h(x))$

δ_m > 0

b $\approx \langle \delta_{h}(x_{1})\delta_{m}(x_{2}) \rangle / \langle \delta_{m}(x_{1})\delta_{m}(x_{2}) \rangle$

The bias is a measure of the response of the number of halos to a long wavelength density fluctuation

$$\boldsymbol{\delta}_{h} \approx \frac{d \ln n_{h}}{d \delta_{m}} \boldsymbol{\delta}_{m} \approx \boldsymbol{b} \boldsymbol{\delta}_{m}$$

But from this perspective, can also

$$n_{h}(x) = \overline{n}_{h}(1 + \delta_{h}(x))$$

$$\delta_{m} \ge 0$$

$$n_{h}(x) = n_{h}(1 + \delta_{h}(x))$$
$$\delta_{m} < 0$$

The bias is a measure of the response of the number of halos to a long wavelength density fluctuation

$$\boldsymbol{\delta}_{h} \approx \frac{d \ln n_{h}}{d \delta_{m}} \boldsymbol{\delta}_{m} \approx \boldsymbol{b} \boldsymbol{\delta}_{m}$$

But from this perspective, can also

$$n_h(x) = \overline{n}_h(1 + \delta_h(x))$$

 $\delta_m > 0$

 $n_h(x) = \overline{n}_h(1 + \delta_h(x))$



$$\mathbf{b} = \frac{1}{n_{h}} \frac{n_{h}(\delta_{m} > 0) - n_{h}(\delta_{m} < 0)}{2 \delta_{m}}$$

Similarly, the ``squeezed" *bispectrum* is a measure of the response of the small-scale power spectrum to a long-wavelength mode

Similarly, the ``squeezed" *bispectrum* is a measure of the response of the small-scale power spectrum to a long-wavelength mode

$$P(k, x) \approx P(k)_{average} + \frac{\partial P(k)}{\partial \delta_{m}} \quad \delta_{m}(x)$$

$$B(k, -k - k_{L}, k_{L}) \approx \frac{\partial P(k)}{\partial \delta_{m}} P(k_{L})$$

$$P(k) = P(k)_{average}$$

$$\delta_{m} = 0$$

$$P(k) = P(k)_{average}(1 + \delta P)$$

$$\delta_{m} > 0$$

The Separate Universe Approach formalizes this average density region Our universe

$$\Leftrightarrow$$

 $Ω_m$, $Ω_Λ$, $Ω_κ$, h, . . .

 $\delta_{\rm m} = 0$

Sirko 2005; Gnedin & Kravtsov 2011 Baldauf, Seljak, Senatore, Zaldarriaga 2011, 2015 Li, Hu, Takada 2014, 2016 Chiang, Wagner, Schmidt, Komatsu 2014a, (+perm) 2014b

The Separate Universe Approach formalizes this average density region Our universe

$$\iff$$

 $Ω_m$, $Ω_Λ$, $Ω_κ$, h, . . .

 $\delta_m = 0$

Large overdense region



Sirko 2005; Gnedin & Kravtsov 2011 Baldauf, Seljak, Senatore, Zaldarriaga 2011, 2015 Li, Hu, Takada 2014, 2016 Chiang, Wagner, Schmidt, Komatsu 2014a, (+perm) 2014b

Separate, Closed Universe

 $\Omega_{mW}, \Omega_{\Lambda W}, \Omega_{\kappa W}, h_{W, \ldots}$

(i.e. universe w/different cosmological parameters)

The Separate Universe Approach formalizes this Separate, Closed Universe

Large overdense region



 $\Omega_{mW}, \Omega_{\Lambda W}, \Omega_{\kappa W}, h_{W, \ldots}$

 $\Omega_{m}, \Omega_{\Lambda}, \Omega_{\kappa}, h, \dots \qquad \Omega_{mW}, \Omega_{\Lambda W}, \Omega_{\kappa W}, h_{W, \dots}$

To study coupling between δ_m and small scale modes, or halo abundance, just run calculations with the new cosmological parameters!

Sirko 2005; Gnedin & Kravtsov 2011 Baldauf, Seljak, Senatore, Zaldarriaga 2011, 2015 Li, Hu, Takada 2014, 2016 Chiang, Wagner, Schmidt, Komatsu 2014a, (+perm) 2014b

The Separate Universe Approach formalizes this Separate, Closed Universe

Large overdense region



 $\Omega_{mW}, \Omega_{\Lambda W}, \Omega_{\kappa W}, h_{W, \ldots}$

So far, restricted to CDM and Λ CDM so that there is <u>only one type</u> of energy fluctuation δ_m

Sirko 2005; Gnedin & Kravtsov 2011 Baldauf, Seljak, Senatore, Zaldarriaga 2011, 2015 Li, Hu, Takada 2014, 2016 Chiang, Wagner, Schmidt, Komatsu 2014a, (+perm) 2014b

What if?

Large overdense region



 $\delta_{neutrino} > 0$

??



What if?

Large overdense region



 $\delta_{neutrino} > 0$



does this look like a separate, curved universe with

??

 Ω_{mW} , $\Omega_{neutrinoW}$, $\Omega_{\kappa W}$, $h_{W, \ldots}$

or Ω_mw, Ω_Qw, Ω_κw, h_{w,...}

??

What if?

Large overdense region



 $\delta_{neutrino} > 0$



In particular, another fluid may have nongravitational interactions or other behavior that prevents the energies from evolving like they would in a separate universe

does this look like a separate, curved universe with

<u>?</u>?

 Ω_{mW} , $\Omega_{neutrinoW}$, $\Omega_{\kappa W}$, $h_{W, \ldots}$

or

 $\Omega_{mW},\ \Omega_{QW},\ \Omega_{\kappa W},\ h_{W,\ \ldots}$



In particular, another fluid may have non-gravitational interactions or other behavior that prevents it from evolving like a separate universe

Example I

quick thought experiment: initially coherent matter and neutrino perturbations



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Example I

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Example I

quick thought experiment: initially coherent matter and neutrino perturbations



 a_W does <u>not</u> look like scale factor for another universe with matter, neutrinos

Example I

quick thought experiment: initially coherent matter and neutrino perturbations

On the other hand, a very large scale perturbation

time



This scale-dependent growth is the effect that gives main cosmological constraints on neutrino mass



Hu, Eisenstein, Tegmark 1998

massive neutrinos reduce the typical amplitude of density perturbations



-> less gravitational lensing than a universe where all matter is gravitationally clustered



massive neutrinos reduce the typical amplitude of density perturbations



Bond, Efstathiou, Silk 1980 Hu, Eisenstein, Tegmark 1998

Future: *O*_{MV} ≈ 0.02 eV ??

3 σ detection of Normal Hierarchy (0.06eV)

"Stage IV CMB"



Abazajian et al 2013

CMB S4, LSST (Large Synoptic Survey Telescope), DESI (Dark Energy Spectroscopic Instrument), WFIRST, Euclid (ESA mission), SPHEREx . . .

Example II

quick thought experiment: initially coherent matter and quintessence perturbations **below** quintessence Jeans scale



time

Example II

quick thought experiment: initially coherent matter and quintessence perturbations **below** quintessence Jeans scale



 $\rho_m(a_W) \sim a_W^{-3} = \rho_Q(a_W)$ pressure supported, dilutes faster relative to ρ_m time

a_W does <u>not</u> look like scale factor for another universe with matter, quintessence behaving the same way

Hu, Chiang, Li, ML 1605.01412

Example II

quick thought experiment: initially coherent matter and quintessence perturbations **above** quintessence Jeans scale



Hu, Chiang, Li, ML 1605.01412

Separate Universe

It turns out that even in the funny, sub-Jeans cases one can still construct a "fake separate universe"



Separate Universe with additional weird energy densities

 Ω_{mW} , $\Omega_{neutrinoW}$, $\Omega_{\kappa W}$, Ω_{sw} , $h_{W, \ldots}$

 $\Omega_{mW}, \Omega_{QW}, \Omega_{KW}, \Omega_{SW}, h_{W, \ldots}$



Hu, Chiang, Li, ML 1605.01412
Separate Universe

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Separate Universe with additional weird energy densities

 Ω_{mW} , $\Omega_{neutrinoW}$, $\Omega_{\kappa W}$, Ω_{sw} , $h_{W, \ldots}$

 $\Omega_{mW}, \Omega_{QW}, \Omega_{\kappa W}, \Omega_{SW}, h_{W, \ldots}$

But the separate universe construction is still well-defined if we know evolution of δ_{cdm} , $\delta_{neutrino}$, $\delta_{Quintessence}$

Hu, Chiang, Li, ML 1605.01412

Separate Universe

It turns out that even in the funny, sub-Jeans cases one can still construct a "fake separate universe"



But the separate universe construction is still well-defined if we know evolution of $\delta_{cdm}(t)$, $\delta_{neutrino}(t)$, $\delta_{Quintessence}(t)$



define local expansion history $a_{W}(t), H_{W}(t)$

Hu, Chiang, Li, ML 1605.01412









Solve the nonlinear evolution of spherical over density (spherical cow halo) in two regions





Solve the nonlinear evolution of spherical over density (spherical cow halo) in two regions

Simplest prediction for number of halos that can collapse by time t

 $n(M,t \mid \delta_m(t))$



Solve the nonlinear evolution of spherical over density (spherical cow halo) in two regions

Simplest prediction for number of halos that can collapse by time t

n(M,t | ठ_m(t))

Determine response bias in each region

$$\mathbf{p} = \frac{1}{n_{h}} \frac{n_{h}(\delta_{m} > 0) - n_{h}(\delta_{m} < 0)}{2 \delta_{m}}$$

Halo bias - simplest model In a universe with massive neutrinos:

b(k < kfree-streaming)</td>





Halo bias - simplest model In a universe with massive neutrinos: b(k < kfree-streaming) ≠ b(k > kfree-steaming)



Scale-dependent change in the halo bias:

$$\frac{\Delta b}{b} \approx f_{\nu} + \frac{b-1}{b} (\# f_{\nu}) \qquad \text{where} \qquad f_{\nu} = \frac{\rho_{\nu}}{\rho_{cdm} + \rho_{\nu}}$$

Scale-dependent change in the halo bias:

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The fraction of energy in neutrinos may be tiny ($f_v \ge 0.5\%$)

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Because the feature in the halo bias can be used to measure neutrino mass

Scale-dependent change in the halo bias:

$$\frac{\Delta b}{b} \approx f_{\nu} + \frac{b-1}{b} (\# f_{\nu}) \qquad \text{where} \qquad f_{\nu} = \frac{\rho_{\nu}}{\rho_{cdm} + \rho_{\nu}}$$

The fraction of energy in neutrinos may be tiny ($f_v \ge 0.5\%$) Why care about such a small change to the halo bias?

- Because the feature in the halo bias can be used to measure neutrino mass
- Because this may be a serious systematic for measurements of m_{ν} from galaxy clustering

First Simulations Application: Quintessence

First Simulations Application: Quintessence

matter and quintessence perturbations below the quintessence Jeans scale



matter and quintessence perturbations above the quintessence Jeans scale



First Simulations Application: Quintessence

matter and quintessence perturbations below the quintessence Jeans scale



matter and quintessence perturbations below the quintessence Jeans scale



Super and sub-Jeans scale perturbations <u>clearly</u> map to inequivalent Separate Universes

Small-scale observables in these universes should be different too

Super and sub-Jeans scale perturbations <u>clearly</u> map to inequivalent Separate Universes



Small-scale observables in these universes should be different too



Amplitude of Power Spectrum is different at Final time

Ratio of small-scale matter power spectrum responses in two separate universes



Amplitude of Power Spectrum is different at Final time



Abundance of Halos, and so halo bias b, is different at final time



The difference between super and sub Jeans scale would correspond to a scale-dependent difference in non-linear quantities



e.g. **b** $\approx \langle \delta_{h}(x_{1})\delta_{m}(x_{2}) \rangle / \langle \delta_{m}(x_{1})\delta_{m}(x_{2}) \rangle$ for $|x_{1} - x_{2}| \langle Jeans scale$ **b** $\approx \langle \delta_{h}(x_{1})\delta_{m}(x_{2}) \rangle / \langle \delta_{m}(x_{1})\delta_{m}(x_{2}) \rangle$ for $|x_{1} - x_{2}| \rangle$ Jeans scale with **b** \neq **b**





Our results are consistent with

 $\Delta b_L/b_L = const$

and with spherical cow model and are inconsistent with constant Lagrangian bias w.r.t initial Pmm(k)



Conclusions

Nonlinear structure formation is complicated! But, can lead to new phenomena that may provide new insights into neutrinos, quintessence, and beyond!

Conclusions

- Nonlinear structure formation is complicated! But, can lead to new phenomena that may provide new insights into neutrinos, quintessence, and beyond!
- The separate universe can be extended to situations with non-gravitational forces and a Jeans scale
- This method provides a trick for being able to simulate (a limited set of important observables) in cosmologies with multiple fluids and non-gravitational forces
- The presence of a Jeans scale can lead to *new* observables (scale dependent bias, scale-dependent squeezed bispectrum)
- The distinction between super/sub-Jeans observables can be understood as a difference in the local expansion history in the two regimes a local model of halo bias can not capture this