

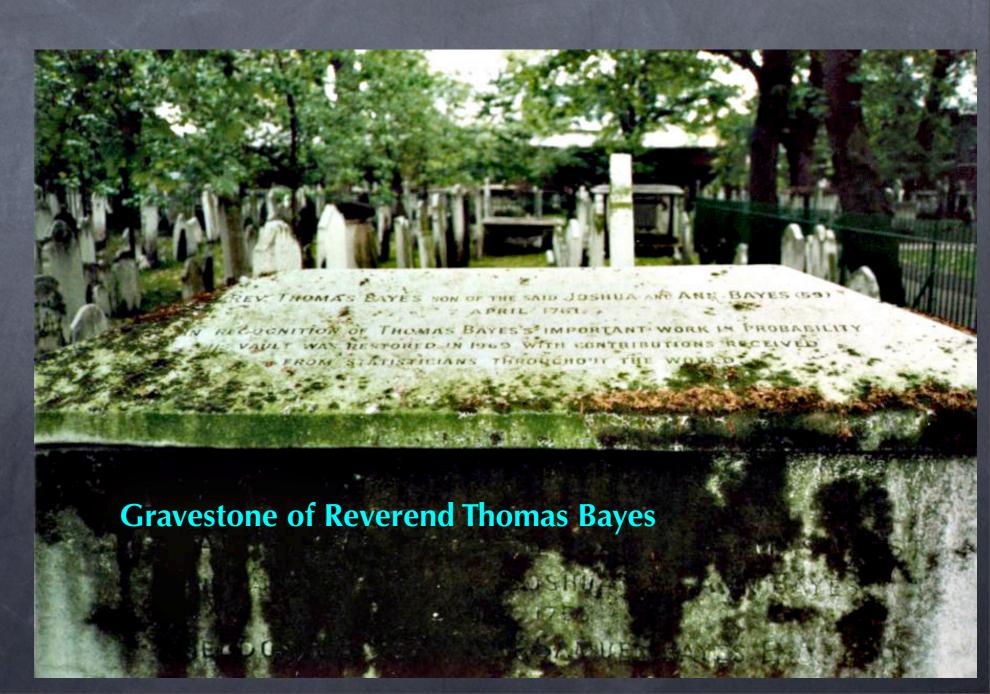
University of Hawaii Institute for Astronomy



Model Selection and Dark Energy

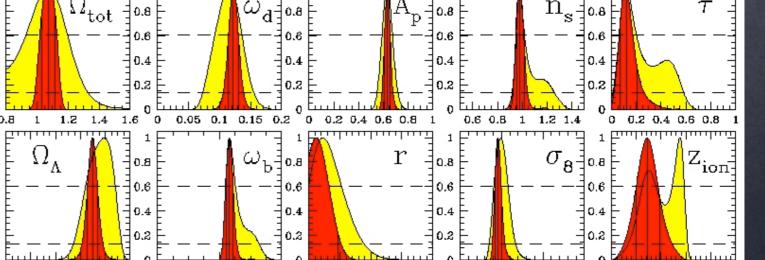
Andrew Liddle February 2007

Microsoft-free presentation



Choose model:
Set of parameter stimation
Set of parameters to be varied
Prior ranges for those parameters
Compute likelihood function
Obtain posterior parameter distribution
Interpret

Choose model:
Set of parameters to be varied
Prior ranges for those parameters
Compute likelihood function
Obtain posterior parameter distribution
Interpret



Bayesian model selection

Choose model:
Set of parameters to be varied
Prior ranges for those parameters
Compute likelihood function
Obtain posterior parameter distribution
Interpret

Bayesian model selection

Choose model M₁:
Set of parameters to be varied
Prior ranges for those parameters
Compute likelihood function
Obtain posterior parameter distribution

Bayesian model selection

Choose model M₁:
Set of parameters to be varied
Prior ranges for those parameters
Compute likelihood function
Obtain posterior parameter distribution

Choose model M₂:
Set of parameters to be varied
Prior ranges for those parameters
Compute likelihood function
Obtain posterior parameter distribution

Bayesian model selection

Choose model M₁:
Set of parameters to be varied
Prior ranges for those parameters
Compute likelihood function
Obtain posterior parameter distribution

Choose model M₂:
Set of parameters to be varied
Prior ranges for those parameters
Compute likelihood function
Obtain posterior parameter distribution

Bayesian model selection

Choose model M₁:
Set of parameters to be varied
Prior ranges for those parameters
Compute likelihood function
Obtain posterior parameter distribution

Choose model M₂:
Set of parameters to be varied
Prior ranges for those parameters
Compute likelihood function
Obtain posterior parameter distribution

Assign model probability $P(M_1)$

Assign model probability $P(M_2)$.

Bayesian model selection

Choose model M₁:
Set of parameters to be varied
Prior ranges for those parameters
Compute likelihood function
Obtain posterior parameter distribution

Choose model M₂:
Set of parameters to be varied
Prior ranges for those parameters
Compute likelihood function
Obtain posterior parameter distribution

Assign model probability $P(M_1)$

Assign model probability P(M₂)

Compute model likelihoods, known as the Bayesian evidence
Update prior model probabilities to posterior ones
[option: multi-model inference by bayesian model averaging]
Interpret

Model Selection

Model selection is the study of sets of parameters. It is a higher level of inference than parameter estimation.

In many contexts we don't actually know which parameters are the relevant ones. A particular situation is deciding if a new parameter, describing some new physical effect seen in data, is actually required.

Model Selection

Model selection is the study of sets of parameters. It is a higher level of inference than parameter estimation.

In many contexts we don't actually know which parameters are the relevant ones. A particular situation is deciding if a new parameter, describing some new physical effect seen in data, is actually required.

A suitable baseline cosmological model to consider is the simplest one giving an adequate fit to current data. It is a spatially-flat adiabatic Λ CDM model with five fundamental parameters and two phenomenological ones.

Model Selection

Model selection is the study of sets of parameters. It is a higher level of inference than parameter estimation.

In many contexts we don't actually know which parameters are the relevant ones. A particular situation is deciding if a new parameter, describing some new physical effect seen in data, is actually required.

A suitable baseline cosmological model to consider is the simplest one giving an adequate fit to current data. It is a spatially-flat adiabatic

\CDM model with five fundamental parameters and two

phenomenological ones.

$\Omega_{ m m}$	matter density
$\Omega_{ m b}$	baryon density
$\Omega_{\mathbf{r}}$	radiation density
h	hubble parameter
A	adiabatic density perturbation amplitude
τ	reionization optical depth
b	bias parameter (or parameters)

There are many, many ways in which this base cosmological model can be extended.

Table 2. Candidate parameters: those which might be relevant for cosmological observations, but for which there is presently no convincing evidence requiring them. They are listed so as to take the value zero in the base cosmological model. Those above the line are parameters of the background homogeneous cosmology, and those below describe the perturbations. Of the latter set, the first six refer to adiabatic perturbations, the next three to tensor perturbations, and the remainder to isocurvature perturbations.

Ω_k	spatial curvature
$N_{\nu} - 3.04$	effective number of neutrino species (CMBFAST definition)
$m_{ u_i}$	neutrino mass for species 'i'
	[or more complex neutrino properties]
$m_{ m dm}$	(warm) dark matter mass
w+1	dark energy equation of state
dw/dz	redshift dependence of w
	[or more complex parametrization of dark energy evolution]
$c_{\rm S}^2 - 1$	effects of dark energy sound speed
$1/r_{ m top}$	topological identification scale
	[or more complex parametrization of non-trivial topology]
dlpha/dz	redshift dependence of the fine structure constant
dG/dz	redshift dependence of the gravitational constant
n-1	scalar spectral index
$dn/d \ln k$	running of the scalar spectral index
$k_{ m cut}$	large-scale cut-off in the spectrum
$A_{ m feature}$	amplitude of spectral feature (peak, dip or step)
$k_{ m feature}$	and its scale
	[or adiabatic power spectrum amplitude parametrized in N bins]
$f_{ m NL}$	quadratic contribution to primordial non-gaussianity
	[or more complex parametrization of non-gaussianity]
r	tensor-to-scalar ratio
$r + 8n_{\mathrm{T}}$	violation of the inflationary consistency equation
$dn_{ m T}/d\ln k$	running of the tensor spectral index
\mathcal{P}_S	CDM isocurvature perturbation
n_S	and its spectral index
$\mathcal{P}_{S\mathcal{R}}$	and its correlation with adiabatic perturbations
$n_{SR}-n_S$	and the spectral index of that correlation
	[or more complicated multi-component isocurvature perturbation]
$G\mu$	cosmic string component of perturbations

How do we compare different cosmological models (i.e. different choices of fundamental parameters)?

Can we say which model is best?

How do we compare different cosmological models (i.e. different choices of fundamental parameters)? Can we say which model is best?

Problem 1: if we add extra parameters, typically the maximum likelihood will increase, even if the new parameter actually has no physical relevance.

How do we compare different cosmological models (i.e. different choices of fundamental parameters)? Can we say which model is best?

Problem 1: if we add extra parameters, typically the maximum likelihood will increase, even if the new parameter actually has no physical relevance.

Problem 2: as we add extra parameters, the uncertainties on existing parameters increase, and eventually we learn nothing useful about anything.

How do we compare different cosmological models (i.e. different choices of fundamental parameters)? Can we say which model is best?

Problem 1: if we add extra parameters, typically the maximum likelihood will increase, even if the new parameter actually has no physical relevance.

Problem 2: as we add extra parameters, the uncertainties on existing parameters increase, and eventually we learn nothing useful about anything.

We need a way of penalizing use of extra parameters - an implementation of Ockham's razor.

Liddle, MNRAS, astro-ph/0401198 astro-ph/0701113

Liddle, MNRAS, astro-ph/0401198 astro-ph/0701113

Akaike information criterion

(Akaike 1974)

Bayesian information criterion (Schwarz 1978)

Bayesian evidence

(Jeffreys 1961 etc)

- Akaike information criterion $AIC = -2 \ln \mathcal{L}_{max} + 2k$
- Bayesian information criterion

Bayesian evidence

Liddle, MNRAS, astro-ph/0401198 astro-ph/0701113

(Akaike 1974)

k = number of parameters

(Schwarz 1978)

(Jeffreys 1961 etc)

- Akaike information criterion $AIC = -2 \ln \mathcal{L}_{max} + 2k$
- Bayesian information criterion $BIC = -2 \ln \mathcal{L}_{max} + k \ln N$
- Bayesian evidence

Liddle, MNRAS, astro-ph/0401198 astro-ph/0701113

(Akaike 1974)

k = number of parameters

(Schwarz 1978)

N = number of datapoints

(Jeffreys 1961 etc)

Akaike information criterion

$$AIC = -2 \ln \mathcal{L}_{max} + 2k$$

Bayesian information criterion

$$BIC = -2\ln \mathcal{L}_{max} + k\ln N$$

Bayesian evidence

$$E = \int d\theta \, \mathcal{L}(\theta) \, \operatorname{pr}(\theta)$$

Liddle, MNRAS, astro-ph/0401198 astro-ph/0701113

(Akaike 1974)

k = number of parameters

(Schwarz 1978)

N = number of datapoints

(Jeffreys 1961 etc)

 θ = parameter vector, pr = prior

Liddle, MNRAS, astro-ph/0401198 astro-ph/0701113

Akaike information criterion

$$AIC = -2 \ln \mathcal{L}_{max} + 2k$$

Bayesian information criterion

$$BIC = -2\ln \mathcal{L}_{max} + k \ln N$$

Bayesian evidence

$$E = \int d\theta \, \mathcal{L}(\theta) \, \mathrm{pr}(\theta)$$

(Akaike 1974)

k = number of parameters

(Schwarz 1978)

N = number of datapoints

(Jeffreys 1961 etc)

 θ = parameter vector, pr = prior

The preferred model is the one which minimizes the information criterion, or maximizes the evidence.

Liddle, MNRAS, astro-ph/0401198 astro-ph/0701113

Akaike information criterion

$$AIC = -2 \ln \mathcal{L}_{max} + 2k$$

k = number of parameters

Bayesian information criterion

$$BIC = -2\ln \mathcal{L}_{max} + k \ln N$$

N = number of datapoints

Bayesian evidence

(Schwarz 1978)

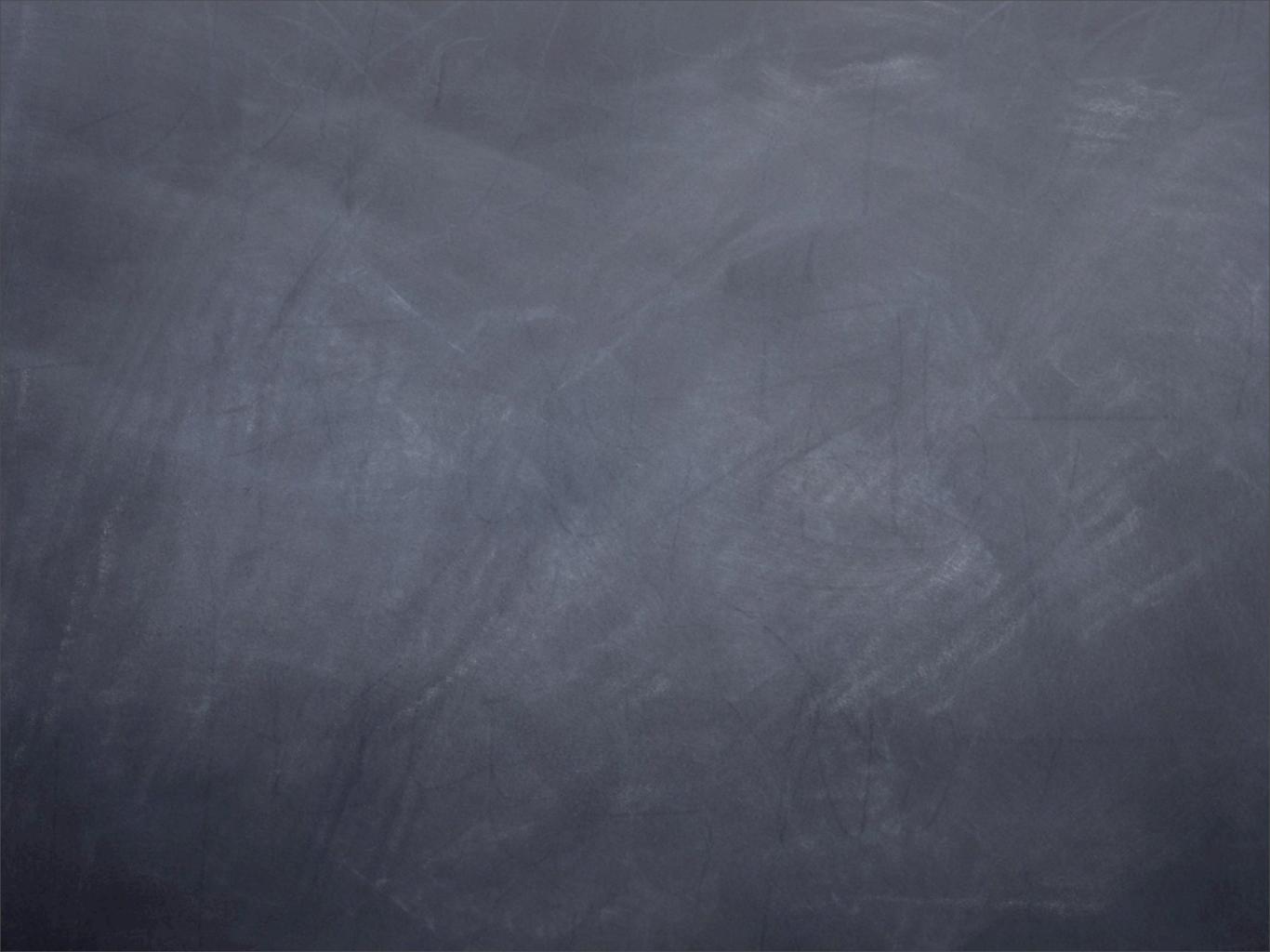
(Akaike 1974)

$$E = \int d\theta \, \mathcal{L}(\theta) \, \operatorname{pr}(\theta)$$

 θ = parameter vector, pr = prior

The preferred model is the one which minimizes the information criterion, or maximizes the evidence.

NB: the ratio of evidences between two models is also known as the Bayes factor.



The Bayesian evidence is the most powerful of these. It is a full implementation of Bayesian inference, and literally gives the probability of the data given the model (note, not the probability of particular parameter values). If multiplied by the prior model probability it gives the posterior model probability. However it can be hard to calculate, being a highly-peaked multi-dimensional integral.

The Bayesian evidence is the most powerful of these. It is a full implementation of Bayesian inference, and literally gives the probability of the data given the model (note, not the probability of particular parameter values). If multiplied by the prior model probability it gives the posterior model probability. However it can be hard to calculate, being a highly-peaked multi-dimensional integral.

The Bayesian Information Citerion was derived using Bayesian statistics. It gives a crude approximation to the Bayesian evidence. While it can give guidance, the assumptions for its validity are questionable in cosmological applications (eg parameter degeneracies).

The Bayesian evidence is the most powerful of these. It is a full implementation of Bayesian inference, and literally gives the probability of the data given the model (note, not the probability of particular parameter values). If multiplied by the prior model probability it gives the posterior model probability. However it can be hard to calculate, being a highly-peaked multi-dimensional integral.

The Bayesian Information Cite in was derived using Bayesian statistics. It gives a crude approximation to the Bayesian evidence. While it can give guidance, the assumptions for its validity are questionable in cosmological applications (eg parameter degeneracies).

The Akaike Information Criterion was derived using information theory techniques. It gives an approximate minimization of the so-called Kullback-Leibler information entropy, which is a measure of the difference between two probability distributions. It is however 'dimensionally inconsistent'.

Model selection techniques are essential when considering whether or not new data requires the addition of new parameters to describe it.

 Statistical fluke: By definition important only if people do their error analysis wrongly.

- Statistical fluke: By definition important only if people do their error analysis wrongly.
- Publication bias: Only positive results get published, enhancing their apparent statistical significance (recognised as a major problem in clinical trials).

- Statistical fluke: By definition important only if people do their error analysis wrongly.
- Publication bias: Only positive results get published, enhancing their apparent statistical significance (recognised as a major problem in clinical trials).
- Inappropriate "a posteriori" reasoning: choosing "interesting" features from the data and assessing their significance via Monte Carlo analyses.

- Statistical fluke: By definition important only if people do their error analysis wrongly.
- Publication bias: Only positive results get published, enhancing their apparent statistical significance (recognised as a major problem in clinical trials).
- Inappropriate "a posteriori" reasoning: choosing "interesting" features from the data and assessing their significance via Monte Carlo analyses.
- Neglect of model dimensionality: using parameter estimation rather than model selection.

Interpretational scale

For guidance, people usually appeal to the Jeffreys' scale.

Interpretational scale

For guidance, people usually appeal to the Jeffreys' scale.

Jeffreys' Scale:

 $\Delta \ln E < 1$

Not worth more than a bare mention

 $1 < \Delta \ln E < 2.5$

Substantial evidence

 $2.5 < \Delta \ln E < 5$

Strong to very strong evidence

 $5 < \Delta \ln E$

Decisive evidence

Interpretational scale

For guidance, people usually appeal to the Jeffreys' scale.

Jeffreys' Scale:

 $\Delta \ln E < 1$

Not worth more than a bare mention

 $1 < \Delta \ln E < 2.5$

Substantial evidence

 $2.5 < \Delta \ln E < 5$

Strong to very strong evidence

 $5 < \Delta \ln E$

Decisive evidence

The most useful divisions are 2.5 (odds ratio of 12:1) and 5 (odds ratio of 150:1).

The parameter placed in the most interesting position is the scalar spectral index, which for the first time is apparently measured as significantly less than one.

The parameter placed in the most interesting position is the scalar spectral index, which for the first time is apparently measured as significantly less than one.

$$n_S = 0.951^{+0.015}_{-0.019}$$

The parameter placed in the most interesting position is the scalar spectral index, which for the first time is apparently measured as significantly less than one.

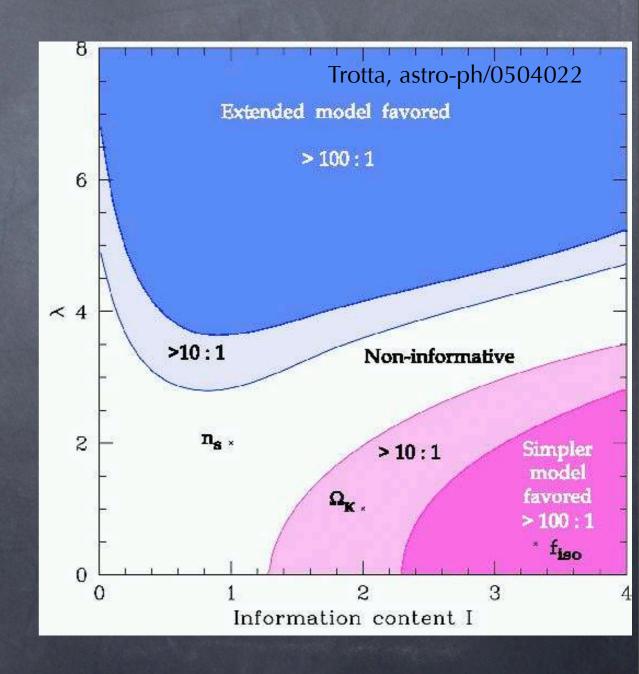
$$n_S = 0.951^{+0.015}_{-0.019}$$

Does a result that is more than 3-sigma need a model selection analysis? Definitely yes - Lindley's paradox operates most strongly in the range 2 to 4 sigma.

The parameter placed in the most interesting position is the scalar spectral index, which for the first time is apparently measured as significantly less than one.

$$n_S = 0.951^{+0.015}_{-0.019}$$

Does a result that is more than 3-sigma need a model selection analysis? Definitely yes - Lindley's paradox operates most strongly in the range 2 to 4 sigma.

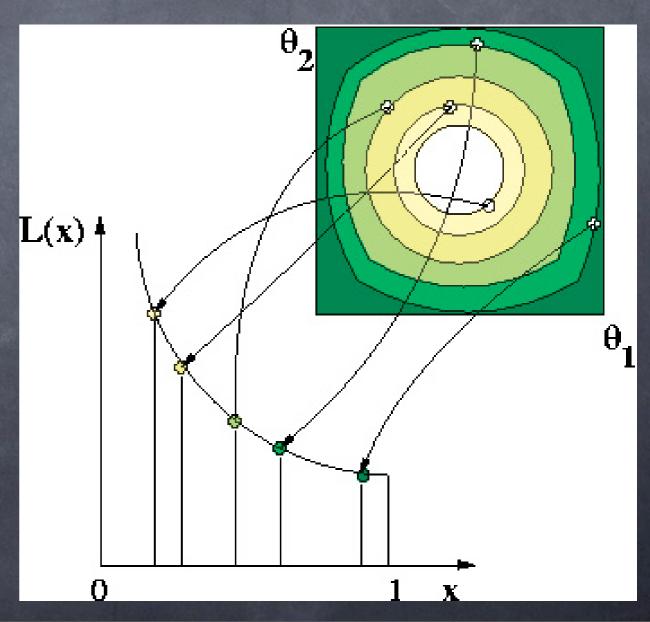


Nested Sampling: CosmoNest

Mukherjee, Parkinson and Liddle, ApJL, astro-ph/0508461 Parkinson, Mukherjee and Liddle, PRD, astro-ph/0605003 www.cosmonest.org

Computing the evidence accurately had been only just within reach of supercomputers. We have recently implemented Skilling's Nested

Sampling algorithm for cosmology.



Nested Sampling: CosmoNest

Mukherjee, Parkinson and Liddle, ApJL, astro-ph/0508461 Parkinson, Mukherjee and Liddle, PRD, astro-ph/0605003 www.cosmonest.org

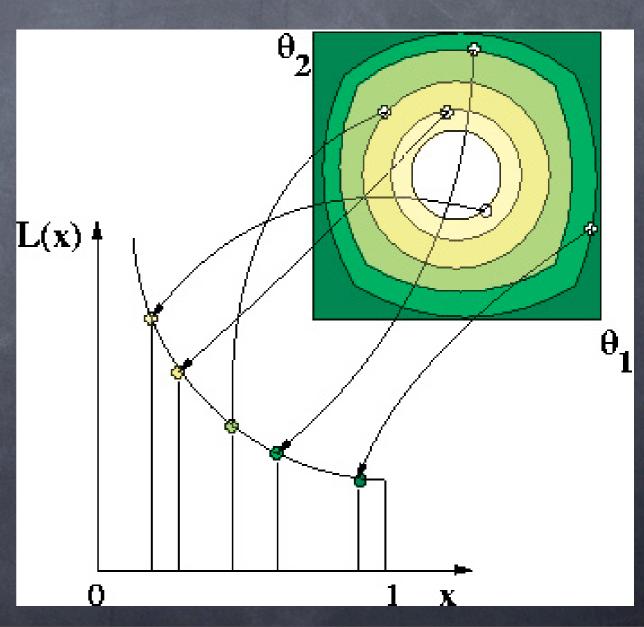
Computing the evidence accurately had been only just within reach of supercomputers. We have recently implemented Skilling's Nested

Sampling algorithm for cosmology.

Skilling (2004) rewrote the evidence as

$$E = \int \mathcal{L}(\theta) \operatorname{pr}(\theta) d\theta = \int_0^1 \mathcal{L}(X) dX$$

where X is the fractional prior mass.



Nested Sampling: CosmoNest

Mukherjee, Parkinson and Liddle, ApJL, astro-ph/0508461 Parkinson, Mukherjee and Liddle, PRD, astro-ph/0605003 www.cosmonest.org

Computing the evidence accurately had been only just within reach of supercomputers. We have recently implemented Skilling's Nested

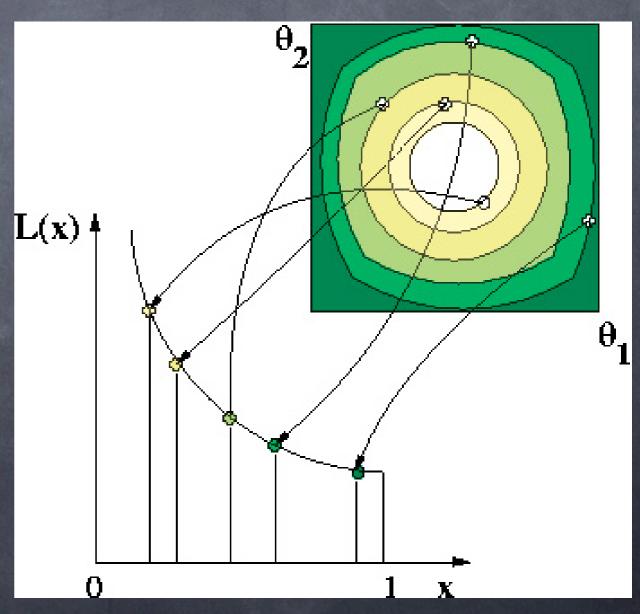
Sampling algorithm for cosmology.

Skilling (2004) rewrote the evidence as

$$E = \int \mathcal{L}(\theta) \operatorname{pr}(\theta) d\theta = \int_0^1 \mathcal{L}(X) dX$$

where X is the fractional prior mass.

This can then be evaluated using Monte Carlo samples to trace the variation of likelihood with prior mass, peeling away thin nested isosurfaces of equal likelihood.

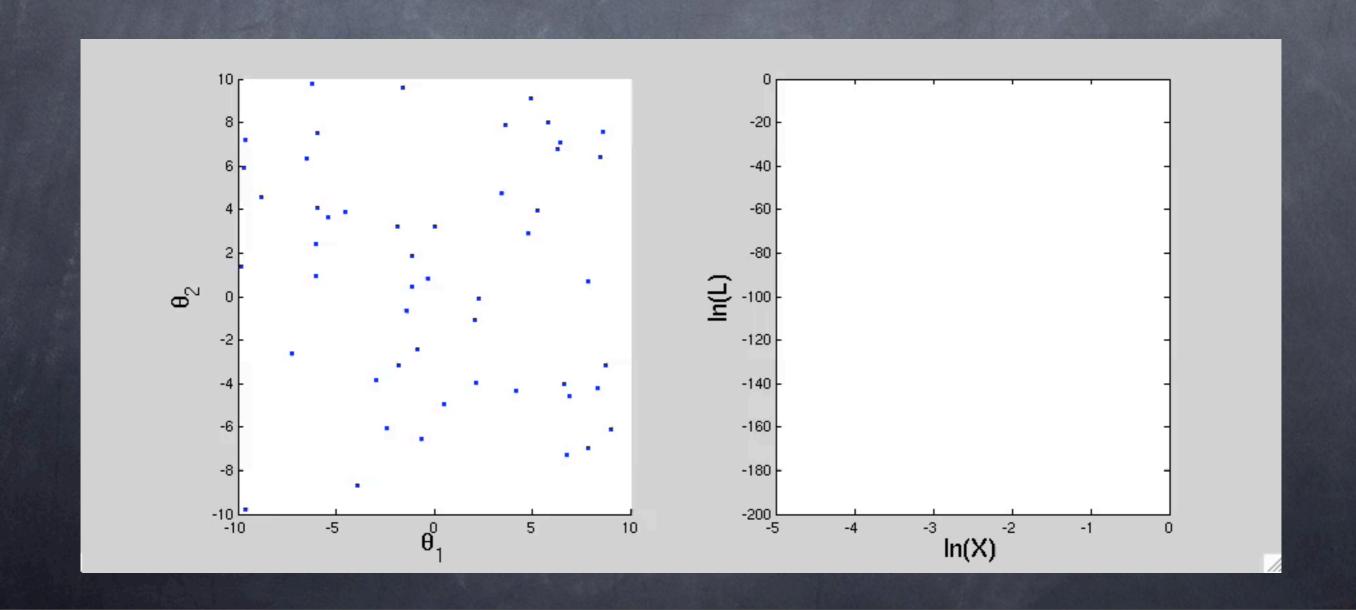


Nested Sampling

The method `walks' a set of points (eg 300) into the high-likelihood region using replacement. The main difficulty in implementing the algorithm successfully is in efficiently generating replacement points which are uniformly sampled from the remaining prior volume.

Nested Sampling

The method `walks' a set of points (eg 300) into the high-likelihood region using replacement. The main difficulty in implementing the algorithm successfully is in efficiently generating replacement points which are uniformly sampled from the remaining prior volume.



A model selection example: spectral index from WMAP3

Parkinson, Mukherjee and Liddle, astro-ph/0605003

WMAP3 has been interpretted as giving ruling out the Harrison-Zel'dovich $n_S = 1$ spectrum and hence favouring inflation, e.g. $n_S = 0.951^{+0.015}_{-0.019}$. But this ignores model dimensionality. Using CosmoNest we find

A model selection example: spectral index from WMAP3

Parkinson, Mukherjee and Liddle, astro-ph/0605003

WMAP3 has been interpretted as giving ruling out the Harrison-Zel'dovich $n_S = 1$ spectrum and hence favouring inflation, e.g. $n_S = 0.951^{+0.015}_{-0.019}$. But this ignores model dimensionality. Using CosmoNest we find

Datasets	Model	In E	
WMAP only	HZ varying n	0.0 0.34 ± 0.26	
WMAP+all	HZ varying n n and r (uniform on r) n and r (log on r)	0.0 1.99 ± 0.26 -1.45 ± 0.45 1.90 ± 0.24	

Datasets	Model	In E	
WMAP only	HZ varying n	0.0 0.34 ± 0.26	
WMAP+all	HZ varying n n and r (uniform on r) n and r (log on r)	0.0 1.99 ± 0.26 -1.45 ± 0.45 1.90 ± 0.24	



Datasets	Model	In E
WMAP only	HZ varying n	0.0 0.34 ± 0.26
WMAP+all	HZ varying n n and r (uniform on r) n and r (log on r)	0.0 1.99 ± 0.26 -1.45 ± 0.45 1.90 ± 0.24

1. WMAP alone cannot distinguish between HZ and a varying spectral index.



Datasets	Model	In E
WMAP only	HZ varying n	0.0 0.34 ± 0.26
WMAP+all	HZ varying n n and r (uniform on r) n and r (log on r)	0.0 1.99 ± 0.26 -1.45 ± 0.45 1.90 ± 0.24

- 1. WMAP alone cannot distinguish between HZ and a varying spectral index.
- 2. Adding other datasets starts to prefer varying n, but only at odds of about 8:1.



Datasets	Model	In E	
WMAP only	HZ varying n	0.0 0.34 ± 0.26	
WMAP+all	HZ varying n n and r (uniform on r) n and r (log on r)	0.0 1.99 ± 0.26 -1.45 ± 0.45 1.90 ± 0.24	}

- 1. WMAP alone cannot distinguish between HZ and a varying spectral index.
- 2. Adding other datasets starts to prefer varying n, but only at odds of about 8:1.
- 3. However inflation predicts we should include both n and r, which is actually disfavoured as compared to HZ...



Datasets	Model	In E		
WMAP only	HZ varying n	0.0 0.34 ± 0.26		
WMAP+all	HZ varying n n and r (uniform on r) n and r (log on r)	0.0 1.99 ± 0.26 -1.45 ± 0.45 1.90 ± 0.24		

- 1. WMAP alone cannot distinguish between HZ and a varying spectral index.
- 2. Adding other datasets starts to prefer varying n, but only at odds of about 8:1.
- 3. However inflation predicts we should include both n and r, which is actually disfavoured as compared to HZ...
- 4. ... unless you use a logarithmic prior for r, which puts you back close to the r=0 case.



Model selection for survey comparison/design

As well as applying to present data, a powerful tool is forecasts of the model selection capabilities of upcoming experiments, eg dark energy surveys.

Model selection for survey comparison/design

As well as applying to present data, a powerful tool is forecasts of the model selection capabilities of upcoming experiments, eg dark energy surveys.

■ Fisher matrix approach:

simulate data for a fiducial model (eg LambdaCDM); estimate expected parameter uncertainties about that model; interpret that if the true model is outside the contours, LambdaCDM is excluded.

Model selection for survey comparison/design

As well as applying to present data, a powerful tool is forecasts of the model selection capabilities of upcoming experiments, eg dark energy surveys.

Fisher matrix approach:

simulate data for a fiducial model (eg LambdaCDM); estimate expected parameter uncertainties about that model; interpret that if the true model is outside the contours, LambdaCDM is excluded.

Bayes factor approach:

simulate data at each point in parameter plane; compute Bayes factor (ie evidence ratio) of full model versus eg LambdaCDM at each point.

Upcoming experiments are usually motivated not by their ability to constrain parameters, but by their ability to discover new physical effects, requiring new parameters (e.g. dark energy evolution).

- Upcoming experiments are usually motivated not by their ability to constrain parameters, but by their ability to discover new physical effects, requiring new parameters (e.g. dark energy evolution).
- Usually interpretted as giving an experiment's ability to rule out LambdaCDM in favour of a dark energy model whose data is however not that simulated.

- Upcoming experiments are usually motivated not by their ability to constrain parameters, but by their ability to discover new physical effects, requiring new parameters (e.g. dark energy evolution).
- Usually interpretted as giving an experiment's ability to rule out LambdaCDM in favour of a dark energy model whose data is however not that simulated.
- The criterion for ruling out LambdaCDM is exactly the same as that used to rule out any other value in the plane, e.g. w=-0.99. Special status of LambdaCDM is not recognised.

- Upcoming experiments are usually motivated not by their ability to constrain parameters, but by their ability to discover new physical effects, requiring new parameters (e.g. dark energy evolution).
- Usually interpretted as giving an experiment's ability to rule out LambdaCDM in favour of a dark energy model whose data is however not that simulated.
- The criterion for ruling out LambdaCDM is exactly the same as that used to rule out any other value in the plane, e.g. w=-0.99. Special status of LambdaCDM is not recognised.
- Fisher matrix approach often assumes a gaussian likelihood.

Dark energy forecasting

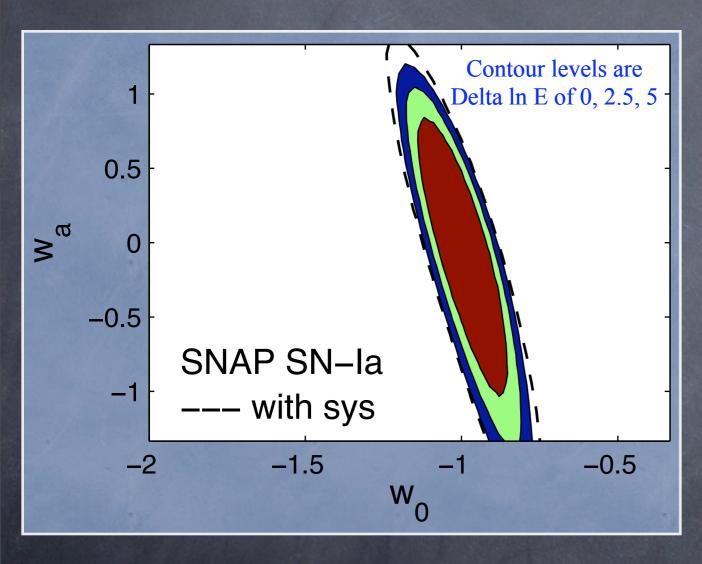
Mukherjee, Parkinson, Corasaniti, Liddle, Kunz, MNRAS, astro-ph/0512484

$$w = w_0 + (1 - a)w_a$$

Dark energy forecasting

Mukherjee, Parkinson, Corasaniti, Liddle, Kunz, MNRAS, astro-ph/0512484

$$w = w_0 + (1 - a)w_a$$

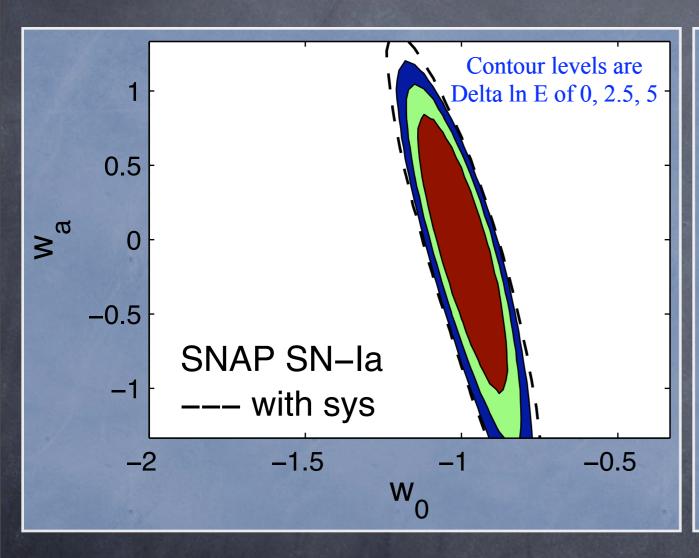


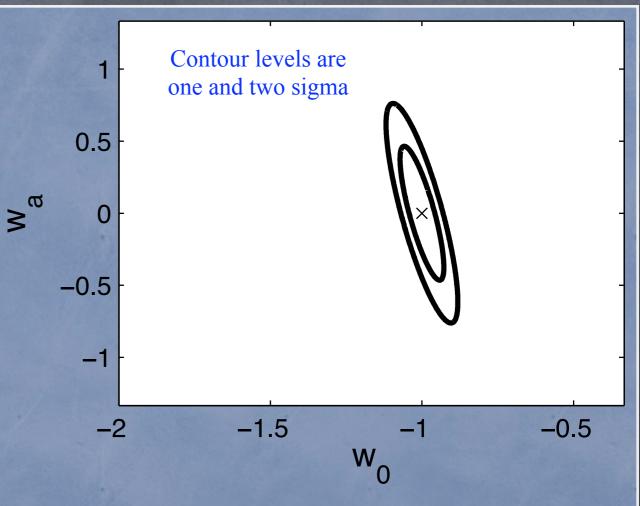
Projected Bayes factor plot against LambdaCDM, SNAP supernovae only

Dark energy forecasting

Mukherjee, Parkinson, Corasaniti, Liddle, Kunz, MNRAS, astro-ph/0512484

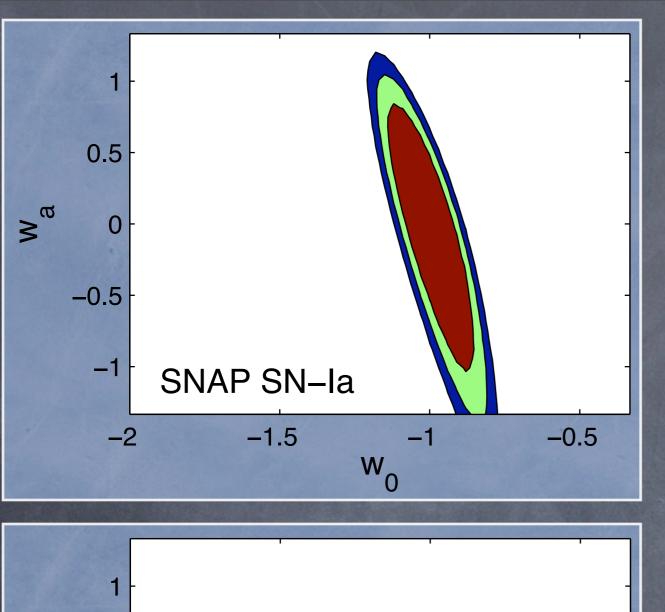
$$w = w_0 + (1 - a)w_a$$

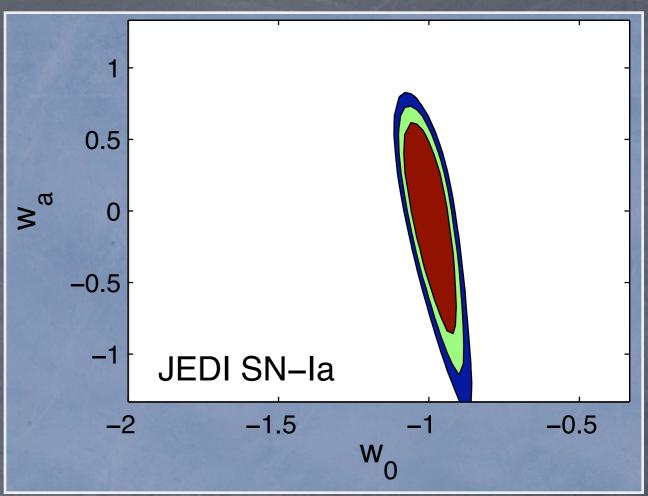


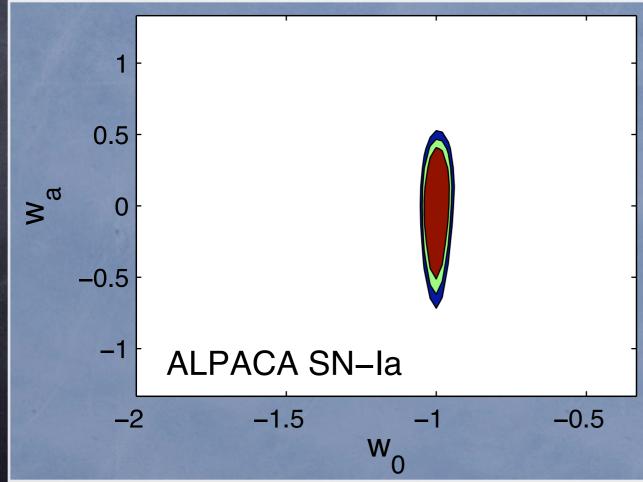


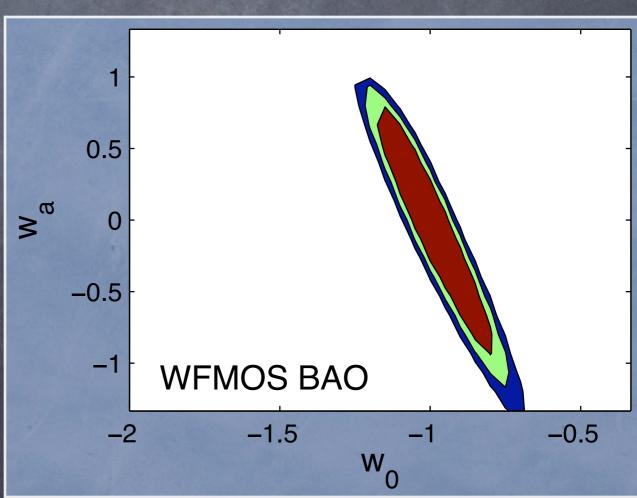
Projected Bayes factor plot against LambdaCDM, SNAP supernovae only

Projected Fisher matrix uncertainties about LambdaCDM









Liddle, Mukherjee, Parkinson, and Wang, astro-ph/0610126

Liddle, Mukherjee, Parkinson, and Wang, astro-ph/0610126

CMB shift+BAO(SDSS)+SN

LambdaCDM

Constant W

W₀-W_a

data used	Model			
WMAP+SDSS+	$\Delta \ln E$	Н	$\chi^2_{ m min}$	parameter constraints
				Model I: Λ
Riess04	0.0	5.7	30.5	$\Omega_{\rm m} = 0.26 \pm 0.03, \ H_0 = 65.5 \pm 1.0$
Astier05	0.0	6.5	94.5	$\Omega_{\rm m} = 0.25 \pm 0.03, \ H_0 = 70.3 \pm 1.0$
				Model II: constant w , flat prior $-1 \le w \le -0.33$
Riess04	-0.1 ± 0.1	6.4	28.6	$\Omega_{\rm m} = 0.27 \pm 0.04, \ H_0 = 64.0 \pm 1.4, \ w < -0.81, -0.70^a$
Astier05	-1.3 ± 0.1	8.0	93.3	$\Omega_{\rm m} = 0.24 \pm 0.03, \ H_0 = 69.8 \pm 1.0, \ w < -0.90, -0.83^a$
	Model III: constant w , flat prior $-2 \le w \le -0.33$			
Riess04	-1.0 ± 0.1	7.3	28.6	$\Omega_{\rm m} = 0.27 \pm 0.04, \ H_0 = 64.0 \pm 1.5, \ w = -0.87 \pm 0.1$
Astier05	-1.8 ± 0.1	8.2	93.3	$\Omega_{\rm m} = 0.25 \pm 0.03, H_0 = 70.0 \pm 1.0, w = -0.96 \pm 0.08$
		M	odel I	V: $w_0 - w_a$, flat prior $-2 \le w_0 \le -0.33$, $-1.33 \le w_a \le 1.33$
Riess04	-1.1 ± 0.1	7.2	28.5	$\Omega_{\rm m} = 0.27 \pm 0.04, \ H_0 = 64.1 \pm 1.5, \ w_0 = -0.83 \pm 0.20, \ w_a =^b$
Astier05	-2.0 ± 0.1	8.2	93.3	$\Omega_{\rm m} = 0.25 \pm 0.03, \ H_0 = 70.0 \pm 1.0, \ w_0 = -0.97 \pm 0.18, \ w_a =^b$
				Model V: w_0-w_a , $-1 \le w(a) \le 1$ for $0 \le z \le 2$
Riess04	-2.4 ± 0.1	9.1	28.5	$\Omega_{\rm m} = 0.28 \pm 0.04, \ H_0 = 63.6 \pm 1.3, \ w_0 < -0.78, -0.60^a, \ w_a = -0.07 \pm 0.34$
Astier05	-4.1 ± 0.1	11.1	93.3	$\Omega_{\rm m} = 0.24 \pm 0.03, \ H_0 = 69.5 \pm 1.0, \ w_0 < -0.90, -0.80^a, \ w_a = 0.12 \pm 0.22$

Liddle, Mukherjee, Parkinson, and Wang, astro-ph/0610126

CMB shift+BAO(SDSS)+SN

LambdaCDM

Constant W

W₀-W_a

data used		Model		
WMAP+SDSS+	$\Delta \ln E$	$H = \chi^2_{\min}$ parameter constraints		
		Model I: Λ		
Riess04	0.0	5.7 30.5 $\Omega_{\rm m} = 0.26 \pm 0.03, H_0 = 65.5 \pm 1.0$		
Astier05	0.0	6.5 94.5 $\Omega_{\rm m} = 0.25 \pm 0.03, H_0 = 70.3 \pm 1.0$		
		Model II: constant w , flat prior $-1 \le w \le -0.33$		
Riess04	-0.1 ± 0.1	6.4 28.6 $\Omega_{\rm m} = 0.27 \pm 0.04, H_0 = 64.0 \pm 1.4, w < -0.81, -0.70^a$		
Astier05	-1.3 ± 0.1	8.0 93.3 $\Omega_{\rm m} = 0.24 \pm 0.03, H_0 = 69.8 \pm 1.0, w < -0.90, -0.83^a$		
		Model III: constant w , flat prior $-2 \le w \le -0.33$		
Riess04	-1.0 ± 0.1	7.3 28.6 $\Omega_{\rm m} = 0.27 \pm 0.04, H_0 = 64.0 \pm 1.5, w = -0.87 \pm 0.1$		
Astier05	-1.8 ± 0.1	8.2 93.3 $\Omega_{\rm m} = 0.25 \pm 0.03, H_0 = 70.0 \pm 1.0, w = -0.96 \pm 0.08$		
		Model IV: w_0-w_a , flat prior $-2 \le w_0 \le -0.33$, $-1.33 \le w_a \le 1.33$		
Riess04	-1.1 ± 0.1	7.2 28.5 $\Omega_{\rm m} = 0.27 \pm 0.04, H_0 = 64.1 \pm 1.5, w_0 = -0.83 \pm 0.20, w_a =b$		
Astier05	-2.0 ± 0.1	8.2 93.3 $\Omega_{\rm m} = 0.25 \pm 0.03, H_0 = 70.0 \pm 1.0, w_0 = -0.97 \pm 0.18, w_a =b$		
		Model V: w_0-w_a , $-1 \le w(a) \le 1$ for $0 \le z \le 2$		
Riess04	-2.4 ± 0.1	9.1 $\Omega_{\rm m} = 0.28 \pm 0.04, H_0 = 63.6 \pm 1.3, w_0 < -0.78, -0.60^a, w_a = -0.07 \pm 0.34$		
Astier05	-4.1 ± 0.1	11.1 93.3 $\Omega_{\rm m} = 0.24 \pm 0.03, H_0 = 69.5 \pm 1.0, w_0 < -0.90, -0.80^a, w_a = 0.12 \pm 0.22$		

Liddle, Mukherjee, Parkinson, and Wang, astro-ph/0610126

CMB shift+BAO(SDSS)+SN

LambdaCDM

Constant W

W₀-W_a

data used		Model		
WMAP+SDSS+	$\Delta \ln E$	$H = \chi^2_{\min}$ parameter constraints		
		Model I: Λ		
Riess04	0.0	5.7 30.5 $\Omega_{\rm m} = 0.26 \pm 0.03, H_0 = 65.5 \pm 1.0$		
Astier05	0.0	6.5 94.5 $\Omega_{\rm m} = 0.25 \pm 0.03, H_0 = 70.3 \pm 1.0$		
		Model II: constant w , flat prior $-1 \le w \le -0.33$		
Riess04	-0.1 ± 0.1	6.4 28.6 $\Omega_{\rm m} = 0.27 \pm 0.04, H_0 = 64.0 \pm 1.4, w < -0.81, -0.70^a$		
Astier05	-1.3 ± 0.1	8.0 93.3 $\Omega_{\rm m} = 0.24 \pm 0.03, H_0 = 69.8 \pm 1.0, w < -0.90, -0.83^a$		
		Model III: constant w , flat prior $-2 \le w \le -0.33$		
Riess04	-1.0 ± 0.1	7.3 28.6 $\Omega_{\rm m} = 0.27 \pm 0.04, H_0 = 64.0 \pm 1.5, w = -0.87 \pm 0.1$		
Astier05	-1.8 ± 0.1	8.2 93.3 $\Omega_{\rm m} = 0.25 \pm 0.03, H_0 = 70.0 \pm 1.0, w = -0.96 \pm 0.08$		
		Model IV: w_0-w_a , flat prior $-2 \le w_0 \le -0.33$, $-1.33 \le w_a \le 1.33$		
Riess04	-1.1 ± 0.1	7.2 28.5 $\Omega_{\rm m} = 0.27 \pm 0.04, H_0 = 64.1 \pm 1.5, w_0 = -0.83 \pm 0.20, w_a =b$		
Astier05	-2.0 ± 0.1	8.2 93.3 $\Omega_{\rm m} = 0.25 \pm 0.03, H_0 = 70.0 \pm 1.0, w_0 = -0.97 \pm 0.18, w_a =b$		
		Model V: w_0-w_a , $-1 \le w(a) \le 1$ for $0 \le z \le 2$		
Riess04	-2.4 ± 0.1	9.1 $\Omega_{\rm m} = 0.28 \pm 0.04, H_0 = 63.6 \pm 1.3, w_0 < -0.78, -0.60^a, w_a = -0.07 \pm 0.34$		
Astier05	-4.1 ± 0.1	11.1 93.3 $\Omega_{\rm m} = 0.24 \pm 0.03, H_0 = 69.5 \pm 1.0, w_0 < -0.90, -0.80^a, w_a = 0.12 \pm 0.22$		

Conclusion: LambdaCDM currently favoured but all models still alive

Future forecasts informed by current data

Liddle, Mukherjee, Parkinson, and Wang, astro-ph/0610126

Bayesian philosophy: continual updating of probabilities as new data comes in.

⇒ Use current probabilities to forecast future experiment outcomes

Liddle, Mukherjee, Parkinson, and Wang, astro-ph/0610126.

Bayesian philosophy: continual updating of probabilities as new data comes in.

- ⇒ Use current probabilities to forecast future experiment outcomes
 - If LambdaCDM is right, are upcoming experiments (eg DES, WFMOS, SNAP) good enough to favour it decisively?

Liddle, Mukherjee, Parkinson, and Wang, astro-ph/0610126

Bayesian philosophy: continual updating of probabilities as new data comes in.

- ⇒ Use current probabilities to forecast future experiment outcomes
 - If LambdaCDM is right, are upcoming experiments (eg DES, WFMOS, SNAP) good enough to favour it decisively?
 - What is the probability that upcoming experiments will robustly detect dark energy evolution?

Liddle, Mukherjee, Parkinson, and Wang, astro-ph/0610126

Bayesian philosophy: continual updating of probabilities as new data comes in.

- ⇒ Use current probabilities to forecast future experiment outcomes
 - If LambdaCDM is right, are upcoming experiments (eg DES, WFMOS, SNAP) good enough to favour it decisively?
 - What is the probability that upcoming experiments will robustly detect dark energy evolution?
 - If future experiments are still inconclusive, how tight will be the limits they can impose on dark energy properties?

Liddle, Mukherjee, Parkinson, and Wang, astro-ph/0610126

- If LambdaCDM is right, are upcoming experiments (eg DES, WFMOS, SNAP) good enough to favour it decisively?
- What is the probability that upcoming experiments will robustly detect dark energy evolution?
- If future experiments are still inconclusive, how tight will be the limits they can impose on dark energy properties?

Liddle, Mukherjee, Parkinson, and Wang, astro-ph/0610126

Under particular prior assumptions we made (the effect of whose variation is readily tested), the answers are ...

- If LambdaCDM is right, are upcoming experiments (eg DES, WFMOS, SNAP) good enough to favour it decisively?
- What is the probability that upcoming experiments will robustly detect dark energy evolution?
- If future experiments are still inconclusive, how tight will be the limits they can impose on dark energy properties?

Liddle, Mukherjee, Parkinson, and Wang, astro-ph/0610126

Under particular prior assumptions we made (the effect of whose variation is readily tested), the answers are ...

- If LambdaCDM is right, are upcoming experiments (eg DES, WFMOS, SNAP) good enough to favour it decisively?
- What is the probability that upcoming experiments will robustly detect dark energy evolution?
- If future experiments are still inconclusive, how tight will be the limits they can impose on dark energy properties?

YES

Liddle, Mukherjee, Parkinson, and Wang, astro-ph/0610126

Under particular prior assumptions we made (the effect of whose variation is readily tested), the answers are ...

If LambdaCDM is right, are upcoming experiments (eg DES, WFMOS, SNAP) good enough to favour it decisively?

YES

What is the probability that upcoming experiments will robustly detect dark energy evolution?

About 25%

If future experiments are still inconclusive, how tight will be the limits they can impose on dark energy properties?

Liddle, Mukherjee, Parkinson, and Wang, astro-ph/0610126

Under particular prior assumptions we made (the effect of whose variation is readily tested), the answers are ...

If LambdaCDM is right, are upcoming experiments (eg DES, WFMOS, SNAP) good enough to favour it decisively?

YES

What is the probability that upcoming experiments will robustly detect dark energy evolution?

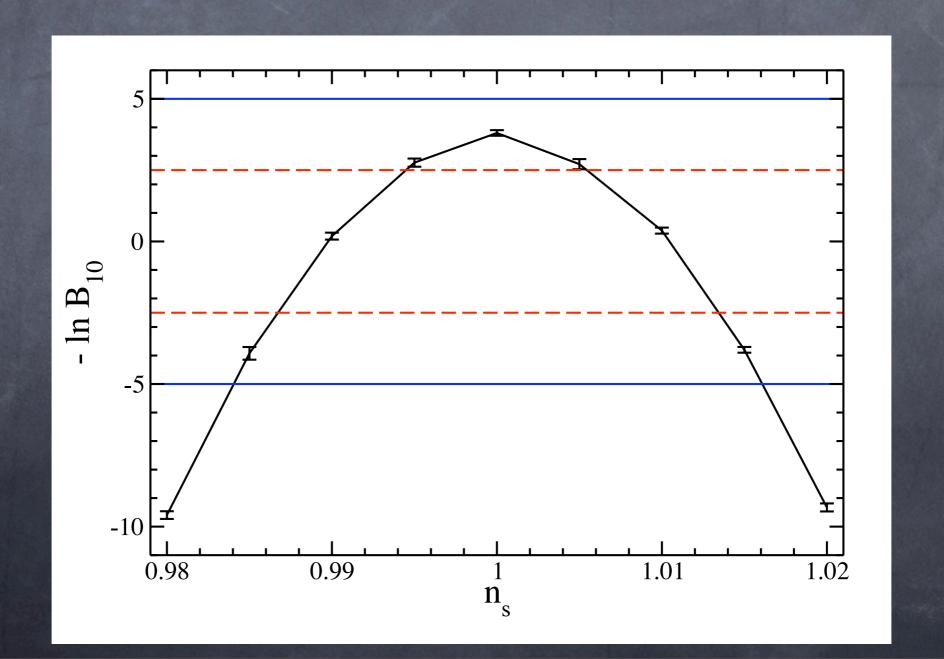
About 25%

If future experiments are still inconclusive, how tight will be the limits they can impose on dark energy properties?

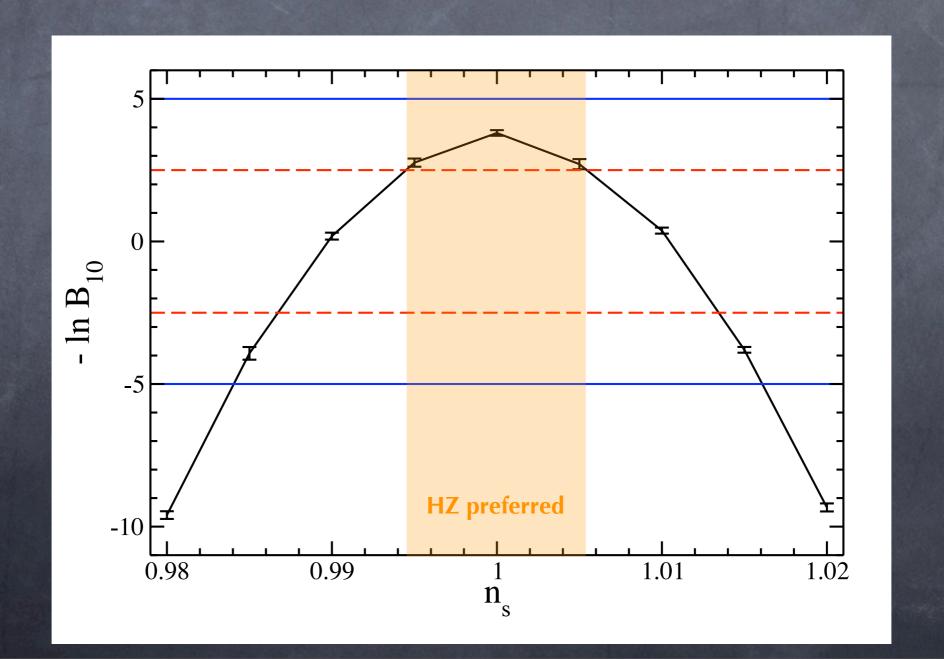
Tighter than you expect!

Pahud, Liddle, Mukherjee, and Parkinson, PRD, astro-ph/0605004

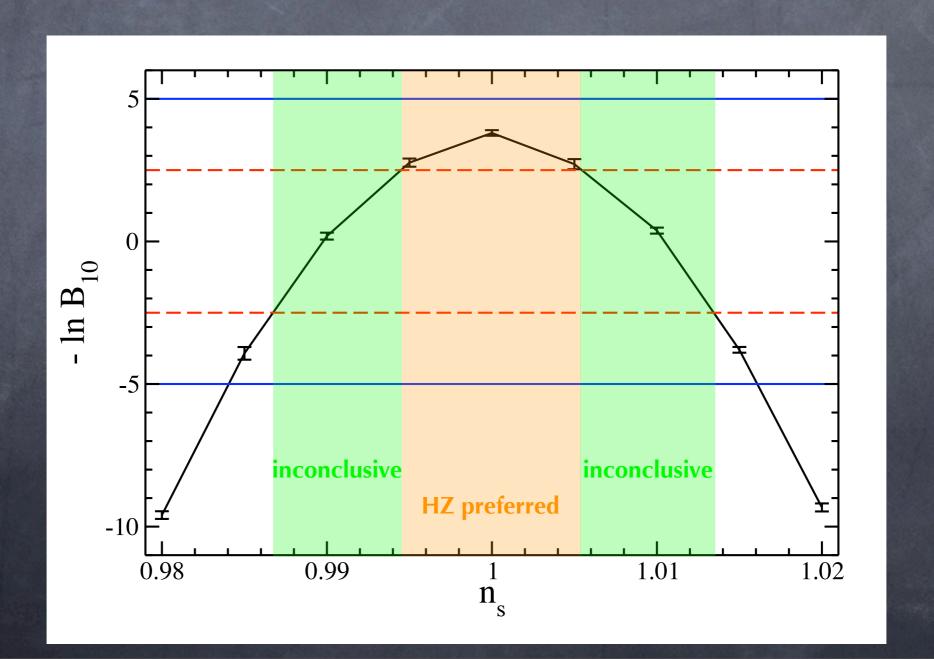
Pahud, Liddle, Mukherjee, and Parkinson, PRD, astro-ph/0605004



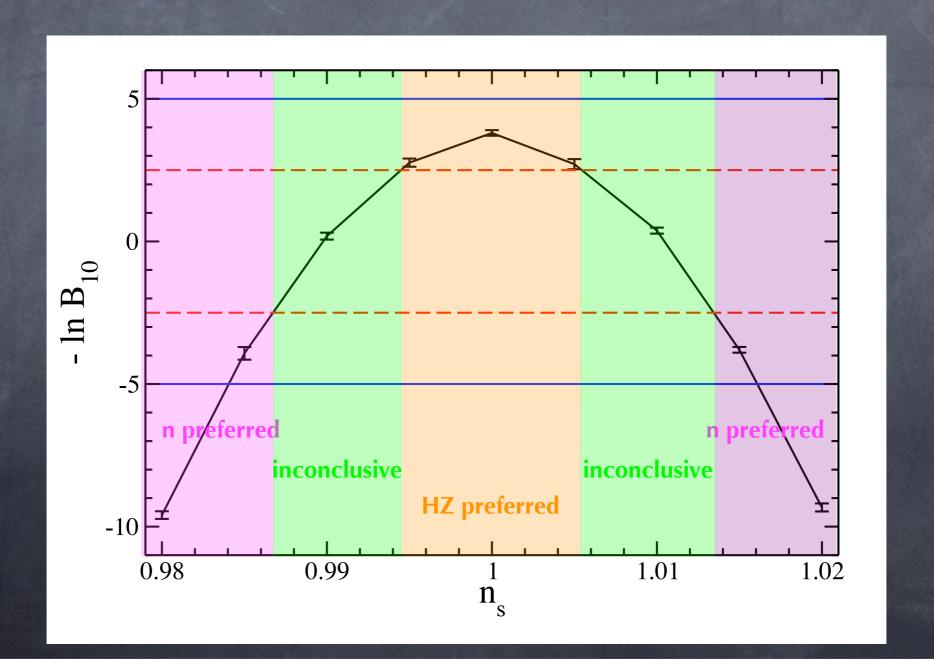
Pahud, Liddle, Mukherjee, and Parkinson, PRD, astro-ph/0605004



Pahud, Liddle, Mukherjee, and Parkinson, PRD, astro-ph/0605004



Pahud, Liddle, Mukherjee, and Parkinson, PRD, astro-ph/0605004



A rigorous approach to defining the Standard Cosmological Model requires Model Selection techniques. Such techniques can positively support simpler models, and set more stringent conditions for inclusion of new parameters.

- A rigorous approach to defining the Standard Cosmological Model requires Model Selection techniques. Such techniques can positively support simpler models, and set more stringent conditions for inclusion of new parameters.
- The Bayesian evidence is the most powerful available tool. It is challenging to compute but nested sampling makes it feasible.

- A rigorous approach to defining the Standard Cosmological Model requires Model Selection techniques. Such techniques can positively support simpler models, and set more stringent conditions for inclusion of new parameters.
- The Bayesian evidence is the most powerful available tool. It is challenging to compute but nested sampling makes it feasible.
- An application to adiabatic models shows current data are comparably well explained by the Harrison-Zel'dovich model and a varying spectral index model, with slight preference for the latter.

- A rigorous approach to defining the Standard Cosmological Model requires Model Selection techniques. Such techniques can positively support simpler models, and set more stringent conditions for inclusion of new parameters.
- The Bayesian evidence is the most powerful available tool. It is challenging to compute but nested sampling makes it feasible.
- An application to adiabatic models shows current data are comparably well explained by the Harrison-Zel'dovich model and a varying spectral index model, with slight preference for the latter.
- Model selection forecasting is a powerful new tool for experimental design and comparison, and is readily applied to dark energy experiments.