## Velocity Probe on Cosmology Tsz Yan LAM (IPMU, U of Tokyo)

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## Overview

Primordial non-Gaussianity

Dynamical mass estimate of massive clusters

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In this talk the local type primordial non-Gaussianity will be used as an illustration ( $f_{nl}$  denotes the strength of primordial non-Gaussianity).

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Dynamical mass estimate of massive clusters

Phase-space distribution as an estimate of cluster (dynamical) mass.

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Possibility to test modified gravity.

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- Halos are aspherical difficult to model.
- Use lensing technique stacking halos.
- After stacking can assume spherical symmetry.
- Use observable quantities to form the phase space distribution: line-of-sight velocity vs projected separation from halo center.

Stacking massive halos



 $N_{halo} = 1$ 

#### Stacking massive halos



 $N_{halo} = 10$ 

-5 4000 3000 -10 2000 -15 1000 VIOS 0 -20 -1000 -25 -2000 -30 -3000 -4000 10 12 6 14 8 2 0 4 r<sub>2d</sub>

Lam et al. in preparation.

 $10^{14} \le M_{halo} \le 1.5 \times 10^{14}$ 



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 $4 \times 10^{14} \le M_{halo} \le 5 \times 10^{15}$ 

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 1-halo terms: virial motions (non-linear)
 2-halo terms: spherical collapse model



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Fitting formula for v\_sigma (Evrard et al. 2008):

$$\sigma_{\rm DM}(M,z) = \sigma_{\rm DM,15} \left[ \frac{(1+z)^{3/2} M_{200b}}{10^{15} \rm M_{sun}} \right]^{\alpha}$$

 $\sigma_{\text{DM},15} = 880 \text{km/s and } \alpha = 0.355$  $M_{200b} = Mass of halo defined as 200 times background density$ 

### 2-halo term: Radial infalling



Infalling velocity predicted by spherical collapse model

BUT it does NOT work





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- Need to add dispersion to the 2-halo term.
- In halo model, the 2-halo term is contributed by particles residing inside a secondary halo.
- Virial motion inside secondary halo.





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$$\sigma_{v,\text{eff}}^{2}(r_{3d},M) = \frac{1}{\bar{w}} \int_{m_{\min}}^{m_{\max}} dm \int_{0}^{r_{\text{vir}}(m)} d^{3}y \ w(m,\vec{y};\vec{r}_{3d},M)\sigma_{v}^{2}(m,\vec{y}),$$
mass of secondary balo relative position of  $r_{3d}$  to the center of secondary balo

 $w(m, \vec{y}; \vec{r}_{3d}, M) = n(m)\rho_{\rm NFW}(|\vec{y}|)[1 + \xi_{hh}(|\vec{x}|, M, m)],$ 

 $\sigma_{\text{eff}}^2(r_{3d}, M_{\text{halo}})$ 




Look better, but still not matching numerical measurements



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- AND assume this radial velocity is described by spherical collapse model.
- Go to numerical simulations to check if it is true after stacking.





#### Halo-halo velocity: Perpendicular to the line of separation





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- Halo-halo velocity: radial + tangential
- Radial: Spherical collapse biased by linear theory prediction (red curve) + dispersion from linear theory.
- Tangential: constant mean and dispersion (from linear theory at large separation).

Semi-empirical model



 $<sup>3</sup>x10^{14} \le M_{primary} \le 5x10^{15}$ 

 $<sup>1 \</sup>times 10^8 \le M_{secondary} \le 5 \times 10^{15}$ 

# Summary

- Semi-empirical model would capture some of the features in the numerical measurements.
- However there is still room for improvement in the current. model.
- The effect of the virial motion extends to radius bigger than the virial radius (see double peaks).
- Modified gravity: Change virial theorem (hence the virial motion) + spherical collapse (halo-halo velocity).

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- Extremum sensitive to primordial non-Gaussianity (see CMB maps above)
- Abundances of rare objects to constrain primordial non-Gaussianity.
- Scale-dependent halo bias distinctive signal at large scale (small k).



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- Peculiar velocity field density field in linear theory.
- Using velocity field to probe primordial non-Gaussianity: pairwise velocity distribution, redshift space distortion.

#### Pairwise Velocity Distribution

Measure probability of the relative velocity of tracers separated by some separation r.

 $v_{\perp}$ 

 $p(v_{\parallel}, v_{\perp}; r, f_{nl}) = p_g(v_{\parallel}, v_{\perp}; r)?$ 

## Pairwise velocity distribution

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Saussian initial conditions: Multivariate Gaussian

$$p_0(\mathbf{v}; r) = \frac{1}{(2\pi)^{3/2} \sqrt{|A|}} \exp\left(-\frac{1}{2} \mathbf{v}^T A^{-1} \mathbf{v}\right) \qquad \mathbf{v} = (v_{\parallel}, v_{\perp_a}, v_{\perp_b})$$
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 $\langle v_{\parallel}^2 \rangle, \, \langle v_{\perp_a}^2 \rangle, \, \text{and} \, \langle v_{\perp_b}^2 \rangle$ 

Can separate into a product of three univariate Gaussian

$$p_0(\mathbf{v};r) = p_0(v_{\parallel};r)p_0(v_{\perp_a};r)p_0(v_{\perp_b};r)$$

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- Higher order correlations as well as cross correlations are non-vanishing.
- Multivariate Gaussian does not work.
- Main corrections are in the order of bispectrum.

### Characteristic Function

Fourier transform of the probability density function.

 $\begin{array}{lll} \text{Moment generating function:} & \mathcal{M}(\vec{\lambda};r) &= \sum_{p=0}^{\infty} \frac{\langle (\vec{\lambda}\cdot\vec{v})^p \rangle}{p!} &= \langle \exp(\vec{\lambda}\cdot\vec{v}) \rangle \\ \\ \text{Cumulant generating function:} & \mathcal{C}(\vec{\lambda};r) &= \sum_{p=2}^{\infty} \frac{\langle (\vec{\lambda}\cdot\vec{v})^p \rangle_c}{p!} &= \langle \exp(\vec{\lambda}\cdot\vec{v}) \rangle_c \\ \\ \text{And they are related by:} & \mathcal{M}(\vec{\lambda};r) &= \exp\left[\mathcal{C}(\vec{\lambda};r)\right] &= \exp\langle \exp(\vec{\lambda}\cdot\vec{v}) \rangle_c \end{array}$ 

Characteristic function:

$$\exp\left[\sum_{j}\frac{\langle v_{j}^{2}\rangle}{2}(i\lambda_{j})^{2}+\frac{\langle v_{\parallel}^{3}\rangle}{3!}(i\lambda_{\parallel})^{3}+\frac{\langle v_{\parallel}v_{\perp_{a}}^{2}\rangle}{2}(i\lambda_{\parallel})(i\lambda_{\perp_{a}})^{2}+\frac{\langle v_{\parallel}v_{\perp_{b}}^{2}\rangle}{2}(i\lambda_{\parallel})(i\lambda_{\perp_{b}})^{2}+\dots\right]$$

cross correlations Characteristic function:  $\exp\left[\sum_{\perp} \frac{\langle v_j^2 \rangle}{2} (i\lambda_j)^2 + \frac{\langle v_{\parallel}^3 \rangle}{3!} (i\lambda_{\parallel})^3 + \frac{\langle v_{\parallel} v_{\perp_a}^2 \rangle}{2} (i\lambda_{\parallel}) (i\lambda_{\perp_a})^2 + \frac{\langle v_{\parallel} v_{\perp_b}^2 \rangle}{2} (i\lambda_{\parallel}) (i\lambda_{\perp_b})^2 + \dots\right]$ 

cross correlations Characteristic function:  $\exp\left[\sum_{i}\frac{\langle v_{j}^{2}\rangle}{2}(i\lambda_{j})^{2} + \frac{\langle v_{\parallel}^{3}\rangle}{3!}(i\lambda_{\parallel})^{3} + \frac{\langle v_{\parallel}v_{\perp_{a}}^{2}\rangle}{2}(i\lambda_{\parallel})(i\lambda_{\perp_{a}})^{2} + \frac{\langle v_{\parallel}v_{\perp_{b}}^{2}\rangle}{2}(i\lambda_{\parallel})(i\lambda_{\perp_{b}})^{2} + \dots\right]$ 

Fourier's Transform to get the probability density function.

Characteristic function:  

$$\exp\left[\sum_{j} \frac{\langle v_{j}^{2} \rangle}{2} (i\lambda_{j})^{2} + \frac{\langle v_{\parallel}^{3} \rangle}{3!} (i\lambda_{\parallel})^{3} + \frac{\langle v_{\parallel} v_{\perp_{a}}^{2} \rangle}{2} (i\lambda_{\parallel}) (i\lambda_{\perp_{a}})^{2} + \frac{\langle v_{\parallel} v_{\perp_{b}}^{2} \rangle}{2} (i\lambda_{\parallel}) (i\lambda_{\perp_{b}})^{2} + \dots\right]$$

Fourier's Transform to get the probability density function.

But in general it is difficult to do the transform.

Approach: Taylor expand around Gaussian density function.

$$\mathcal{F}_0(i\boldsymbol{\lambda};r) = \exp\left[\sum_j \frac{\langle v_j^2 \rangle}{2} (i\lambda_j)^2\right]$$

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$$\mathcal{F}_0(i\boldsymbol{\lambda};r) = \exp\left[\sum_j \frac{\langle v_j^2 \rangle}{2} (i\lambda_j)^2\right]$$

The linear relative velocity density function:

 $p(v_{\parallel}, v_{\perp_a}, v_{\perp_b}; f_{nl}, r) = p_0(v_{\parallel}, v_{\perp_a}, v_{\perp_b}; r)[1 + \alpha_{300}h_{300} + \alpha_{120}(h_{120} + h_{102})],$ 

$$\alpha_{300} = \frac{1}{6} \frac{\langle v_{\parallel}^3 \rangle}{\langle v_{\parallel}^2 \rangle^{3/2}}, \qquad \alpha_{120} = \frac{1}{2} \frac{\langle v_{\parallel} v_{\perp}^2 \rangle}{\langle v_{\parallel}^2 \rangle^{1/2} \langle v_{\perp}^2 \rangle},$$
$$h_{ijk} \equiv H_i(\nu_{\parallel}) H_j(\nu_{\perp_a}) H_k(\nu_{\perp_b})$$



 $p(v_{\parallel}, v_{\perp}; r, f_{nl})/p_g(v_{\parallel}, v_{\perp}; r)$ 

Change due to cross correlation.



Lam, Nishimichi & Yoshida 2011

Linear theory predictions:

#### Marginalized over $v_{\perp}$ , ratio < 1 when $v_{\parallel} \gg 0$ for $f_{nl} > 0$ ratio > 1 when $v_{\parallel} \ll 0$ for $f_{nl} > 0$

Marginalized over  $v_{\parallel}$ , ratio = 1 for all  $v_{\perp}$  and  $f_{nl}$ 

### Evolution of velocity distribution

Two particles separated by r at z with relative velocity v have different  $r_0$  and  $v_0$  at earlier redshift  $z_{0.}$ 

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Zeldovich Approximation: Velocity unchanged.

 $r^2 = \left(r_i + rac{D_0}{\dot{D}_i} v^i_{\parallel}
ight)^2 + \left(rac{D_0}{\dot{D}_i}
ight)^2 (v^{i-2}_{\perp_a} + v^{i-2}_{\perp_b})$ 

$$egin{aligned} v_{\parallel} &= rac{\dot{D}_0}{r} \left( rac{r_i v_{\parallel}^i}{\dot{D}_i} + rac{D_0}{\dot{D}_i^2} v^{i^2} 
ight) \ &|v_{\perp}|^2 &= v_{\perp_a}^2 + v_{\perp_b}^2 = \left( rac{\dot{D}_0}{\dot{D}_i} v^i 
ight)^2 - v_{\parallel^2}^2 \end{aligned}$$





Parallel to the line of separation



Linear theory predicts no change in PDF when marginalized over v\_parallel

Perpendicular to the line of separation

### Short summary

- Primordial non-Gaussianity leaves signatures on the velocity field: we use the pairwise velocity distribution to demonstrate this effect.
- Primordial non-Gaussianity induces non-vanishing three-point. correlation functions: auto correlation of v\_parallel AND cross correlation of v\_parallel and v\_perp.
- Linear theory does not work wrong predictions for both v\_parallel and v\_perp.
- Evolution model implementing Zeldovich Approximation works reasonably well.

What's next?

Redshift space distortion to probe primordial non-Gaussianity. (see Lam, Desjacques & Sheth 2010 or Schmidt 2010 for linear theory prediction)

Peculiar velocity of biased tracers – massive clusters



Halo-Halo: Parallel to the line of separation



Halo-Halo: Perpendicular to the line of separation