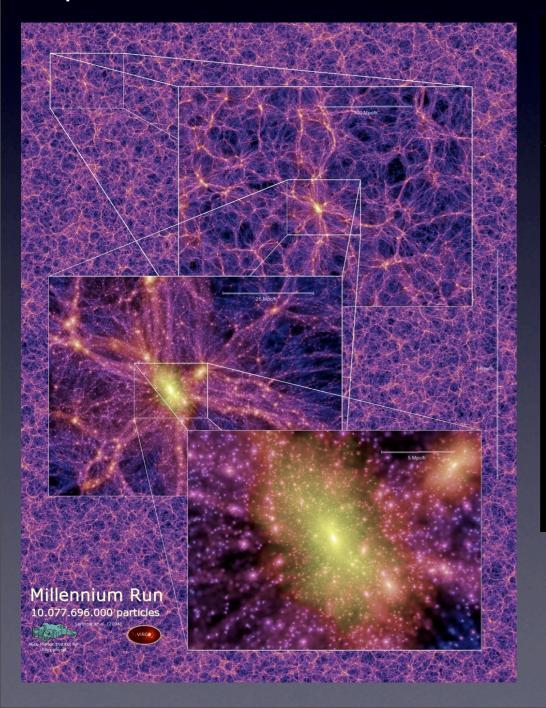
## Seeing [in] the Dark: Cosmic Shear from SDSS

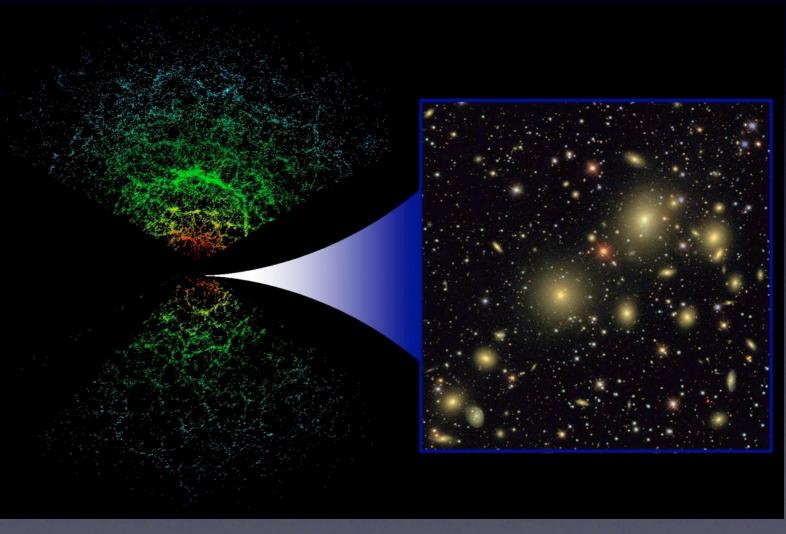
Eric Huff,

Tim Eifler, Chris Hirata, Rachel Mandelbaum, David Schlegel, Uros Seljak

### Why we want to measure weak lensing:

We can make very precise predictions about the invisible...





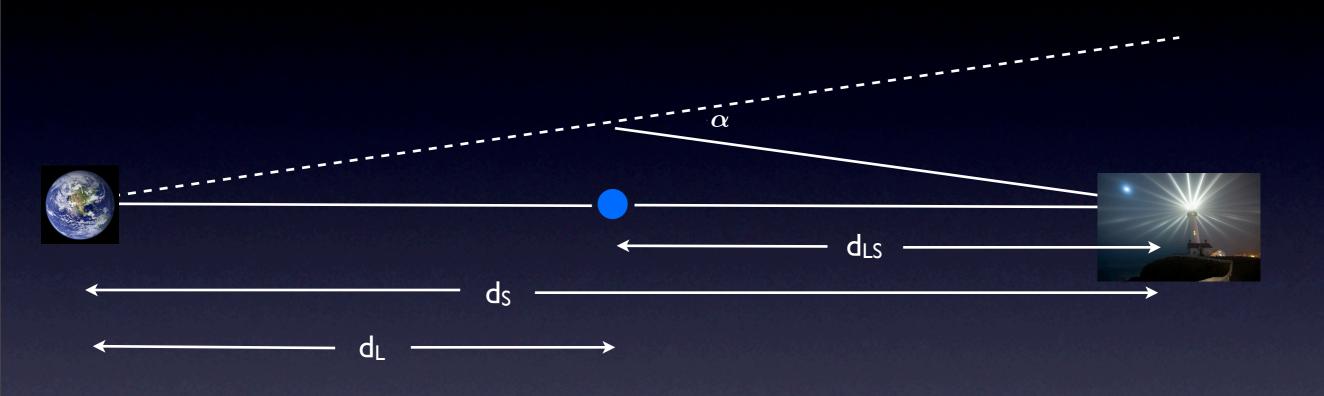
But the visible (baryonic) component is not easy to model.

Weak Lensing Refresher

• Part I: A Cosmic Shear Measurement

Part II: A Novel Magnification Measurement

# Lensing is the distortion of background images by foreground mass:



$$\kappa\left(\bar{\theta}\right) = \int_{0}^{z_{s}} \frac{c}{H_{0}} \frac{dz}{a} \frac{\rho_{m}\left(z\right)}{\Sigma_{\text{crit}}}$$

$$\Sigma_{\rm crit} = \frac{3H_0^2}{8\pi G} \frac{d_S}{d_L d_{LS}}$$

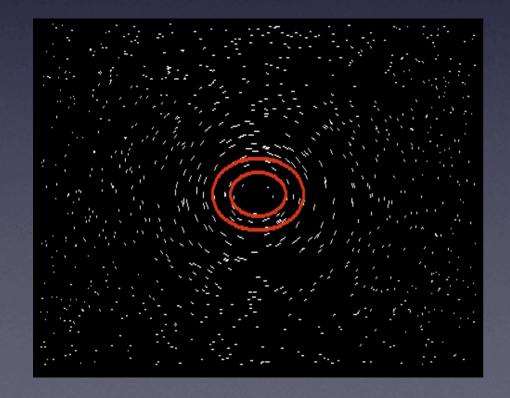
Sensitive to geometry and lens mass

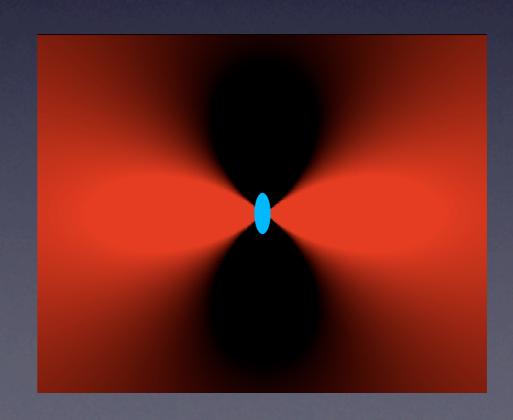
# This distortion can be decomposed into shear and magnification:

$$A_{ij} = \frac{\partial \theta_i^I}{\partial \theta_j^S} = \begin{pmatrix} 1 + \kappa + \gamma_1 & \gamma_2 \\ \gamma_2 & 1 + \kappa - \gamma_1 \end{pmatrix}.$$

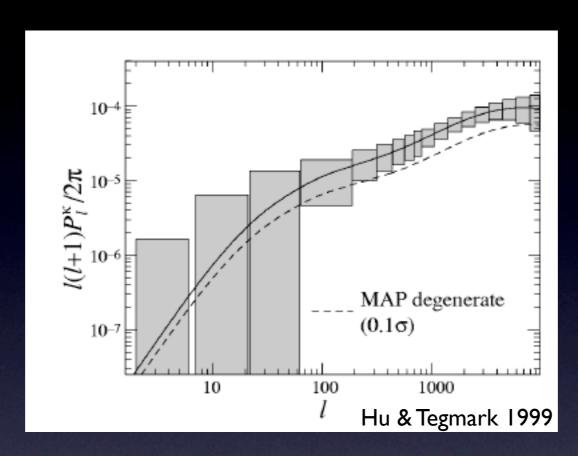
Let's focus on the shear distortion:

$$[\gamma_1 + i\gamma_2](\bar{\theta}) = -\frac{1}{\pi} \int d^2\bar{\phi} \frac{\kappa(\bar{\theta} - \bar{\phi})}{|\bar{\theta} - \bar{\phi}|^2} e^{2i\beta}$$





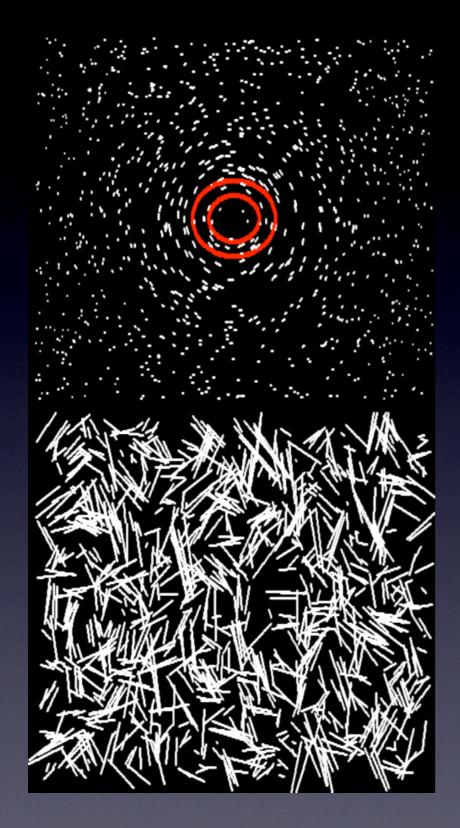
## The two-point statistics of the shear are related to the two-point statistics of the matter.



$$\xi_{\gamma\pm}(\theta) = <\gamma_t\gamma_t> \pm <\gamma_x\gamma_x> = \frac{1}{2\pi} \int_0^\infty d\ell \,\ell P_\kappa(\ell) J_{0,4}(\ell\theta)$$

$$P_{\kappa}(\ell) = \left(\frac{3}{2} \frac{\Omega_m}{d_H^2}\right)^2 \int_0^{\infty} \frac{d\chi}{a(\chi)^2} P_{\delta}\left(\frac{\ell}{d(\chi)}\chi\right) \left[\int_{\chi}^{\infty} d\chi' n(\chi') \frac{d(\chi' - \chi)}{d(\chi')}\right]^2$$

By measuring galaxy shapes, we can get a bias-free measurement of the clustering of matter.



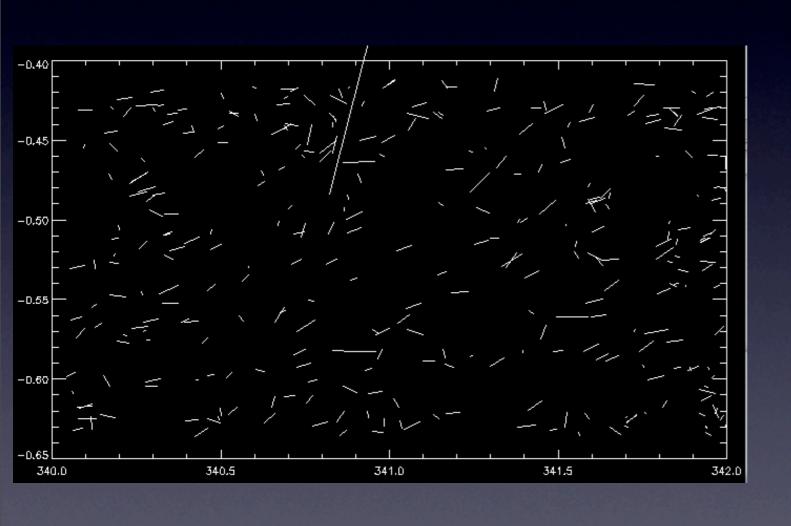
But the signal to noise is very small.

$$<\gamma>\sim 0.001-0.01$$

Especially compared to the random noise from galaxy shapes

$$\sigma_{\gamma} = 0.3$$

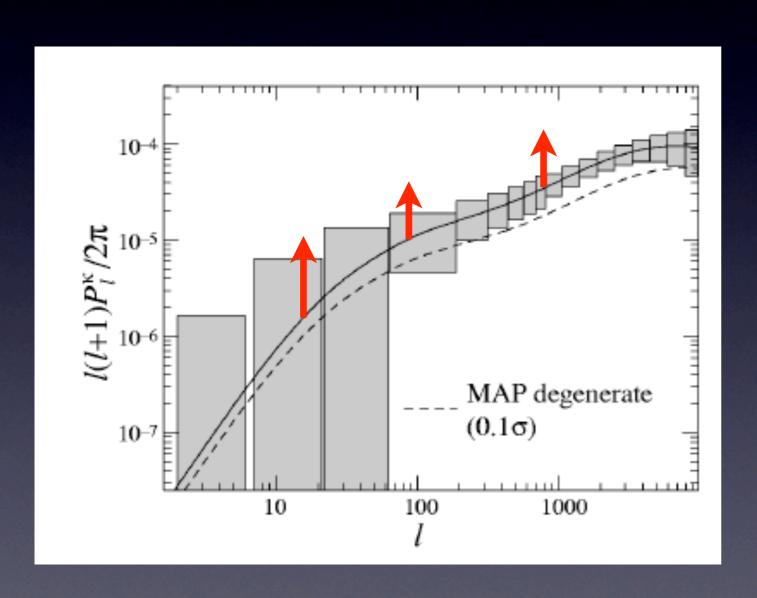
Few data sets exist that can measure this at high signal.



psf ellipticity

photometric redshifts

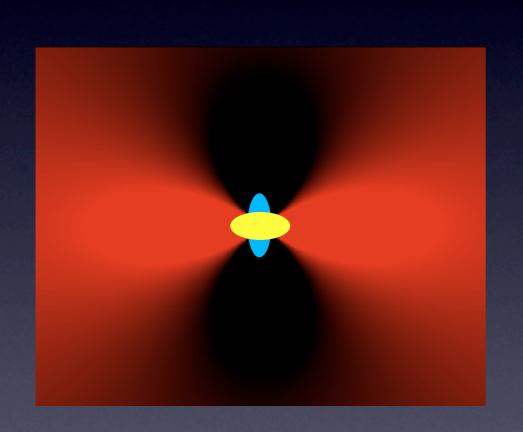
intrinsic alignments



psf ellipticity

photometric redshifts

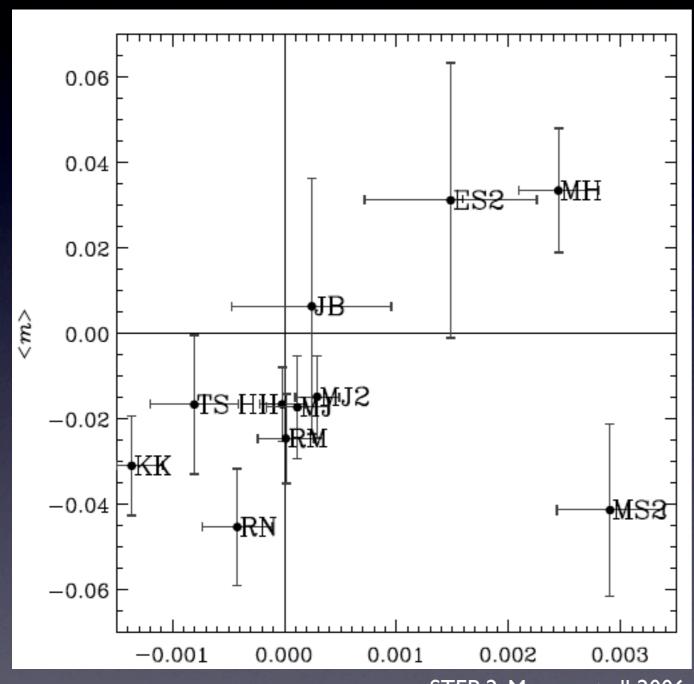
intrinsic alignments



psf ellipticity

photometric redshifts

intrinsic alignments



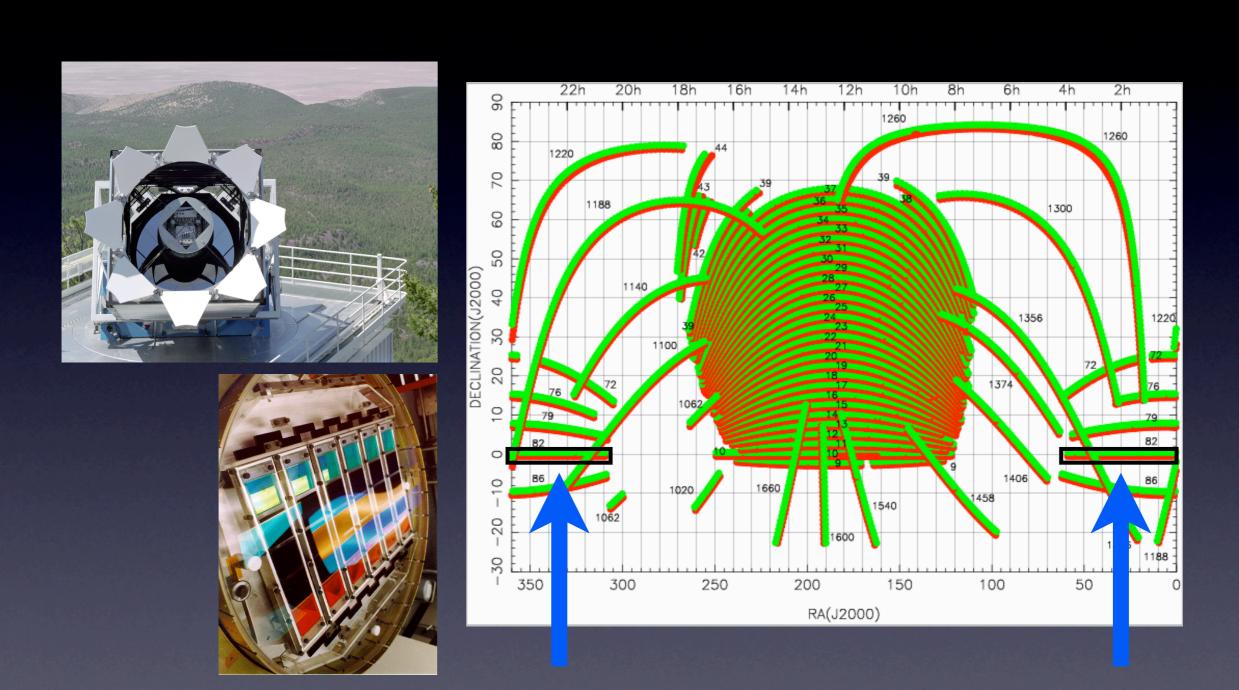
STEP 2: Massey et all 2006

psf ellipticity

photometric redshifts

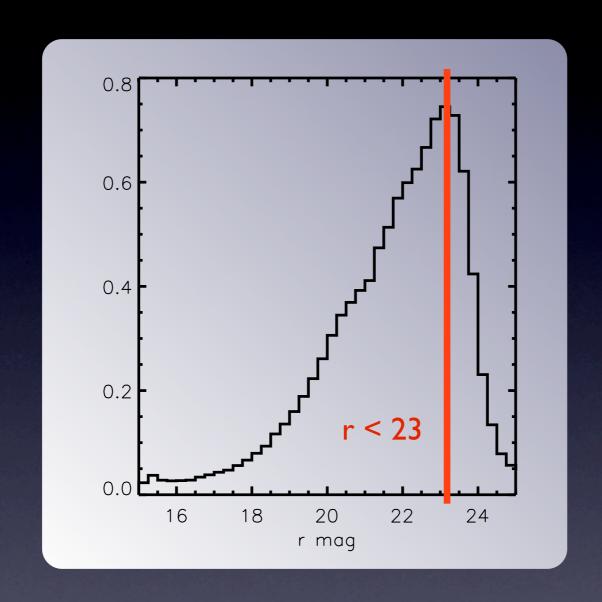
intrinsic alignments

# We've built a catalog from SDSS data that addresses these problems



We set out to coadd the 80+ epochs of Stripe 82 data

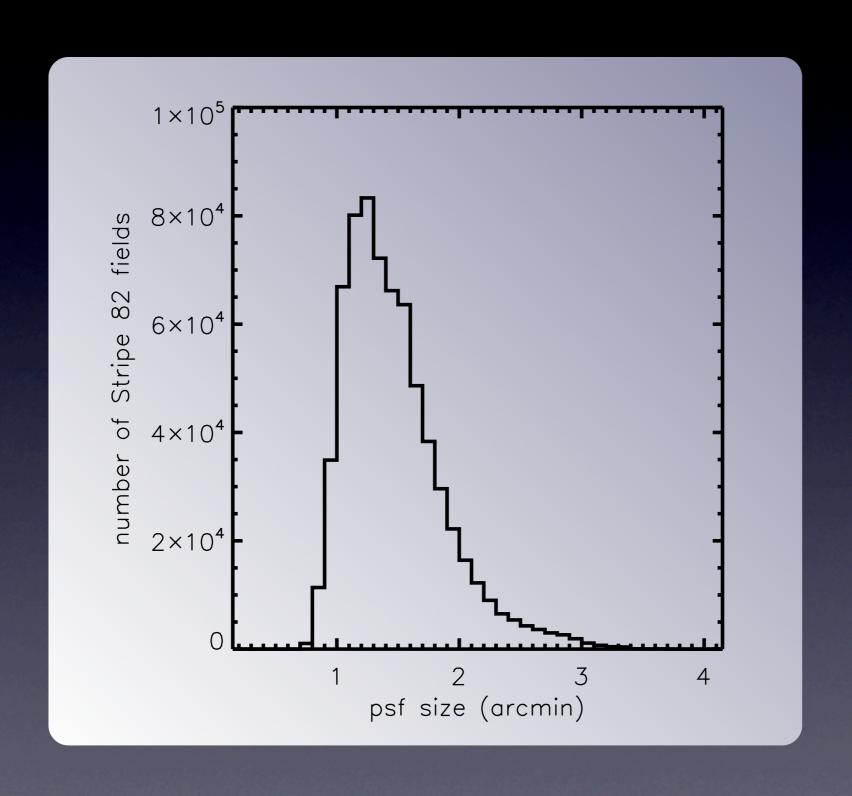
# SDSS is not the ideal survey for measuring cosmic shear



Even the SDSS coadds are fairly shallow.

By comparison, CFHTLS has i < 24.5

# SDSS is not the ideal survey for measuring cosmic shear



# Stripe 82 covers more area than any other cosmic shear survey.

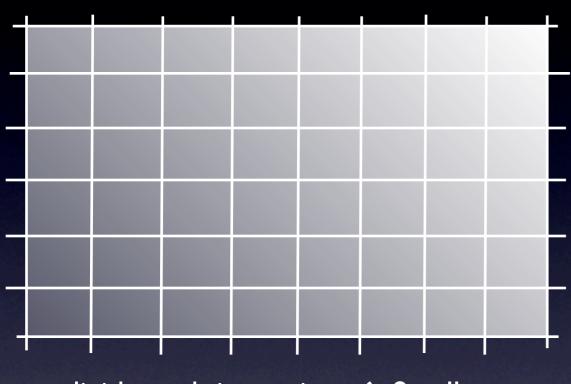
These are some of the surveys that have been done so far.

COSMOS

• CFHTLS

This Work

# We use a rounding kernel to homogenize the psf in each image.

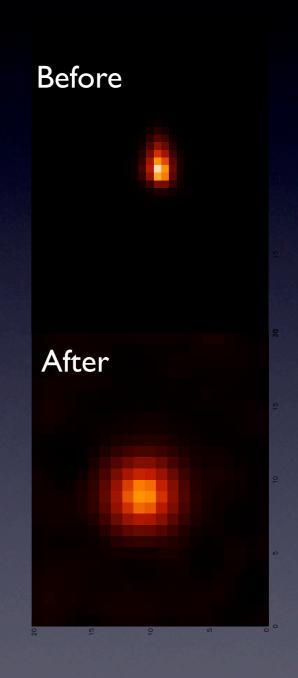


divide each image into 6x8 cells

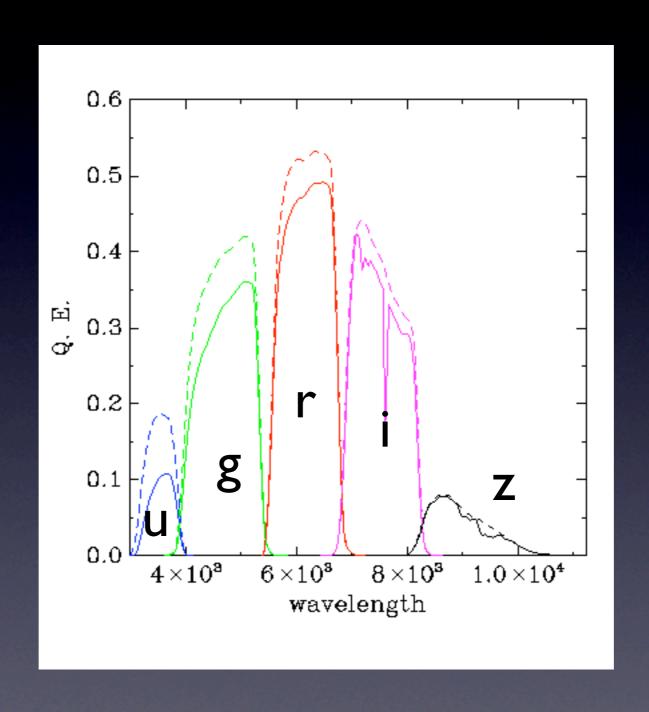
reconstruct SDSS psf in each cell

calculate convolution kernel in each cell

$$\hat{R} = \frac{\hat{G}\hat{T}}{\hat{G}^2 + \lambda}$$

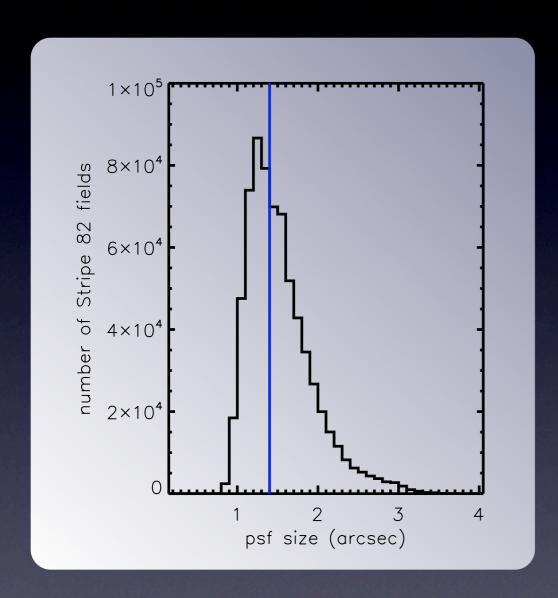


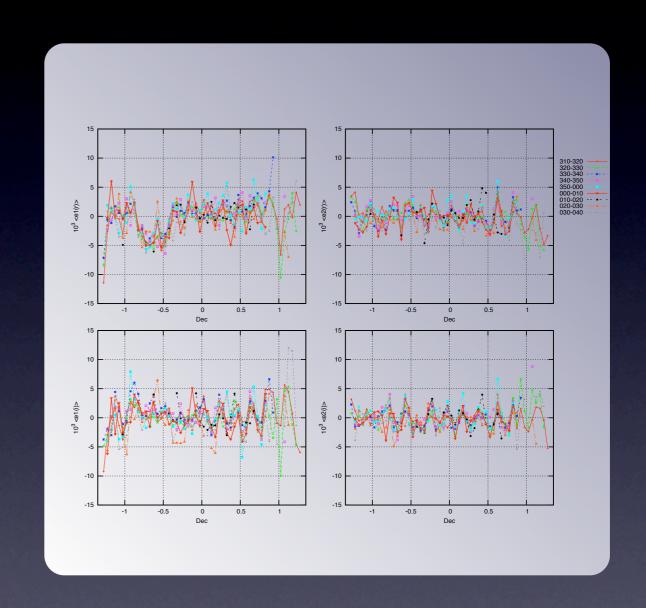
# We perform shape measurements in the SDSS r and i bands.



These are the deepest.

### We made many cuts on the data products

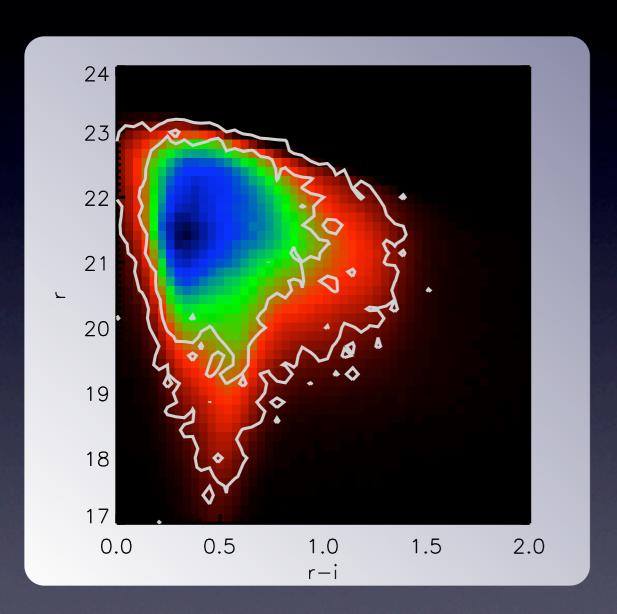


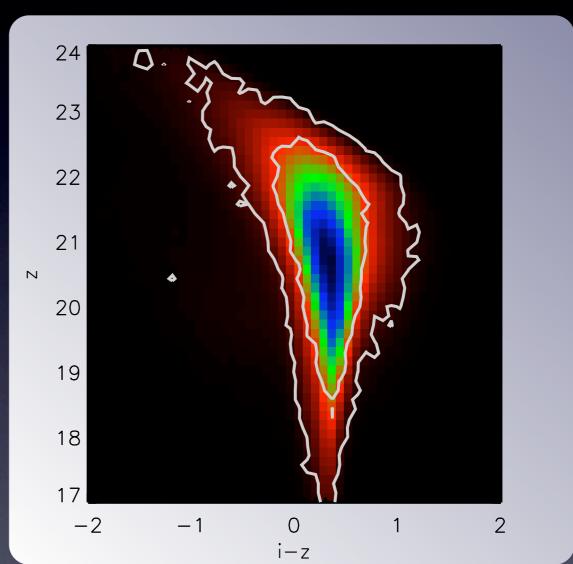


seeing < 1.4 arcsec

exclude r-band camcol 2

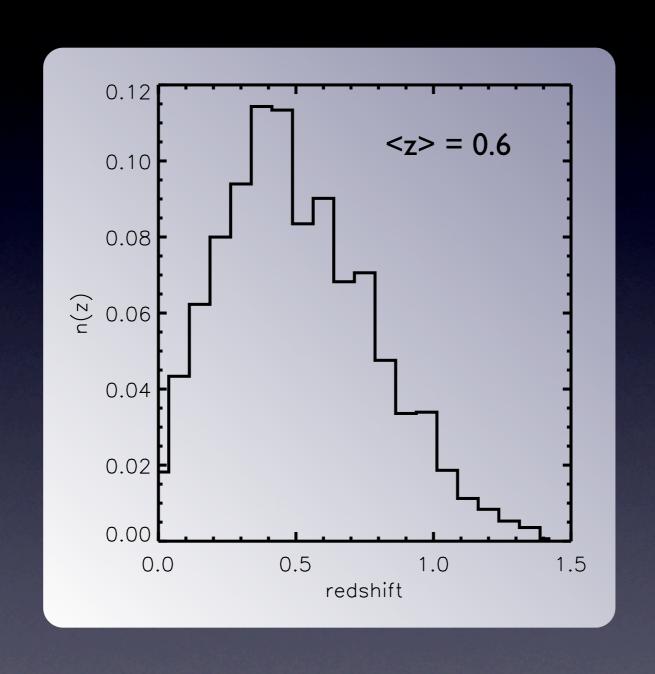
#### To estimate the redshift distribution:



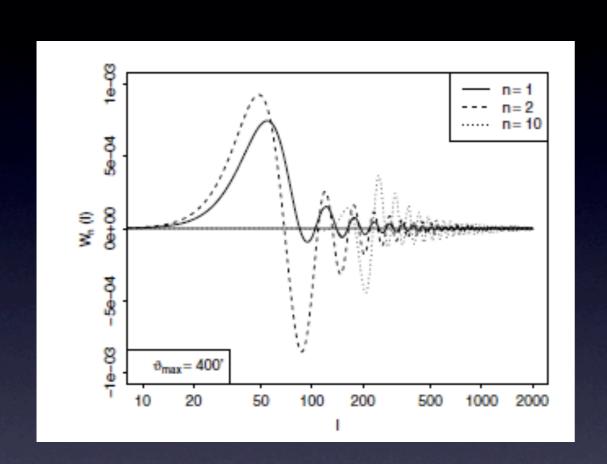


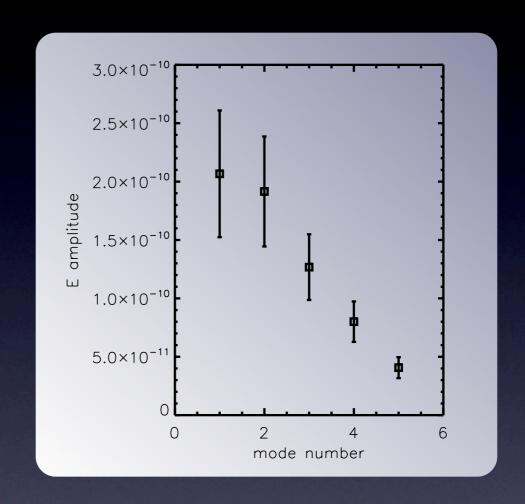
We re-weight the combined calibration sample to match 5-band magnitude distribution

### Our redshift distribution:



## We use the COSEBI basis to decompose into E and B modes.



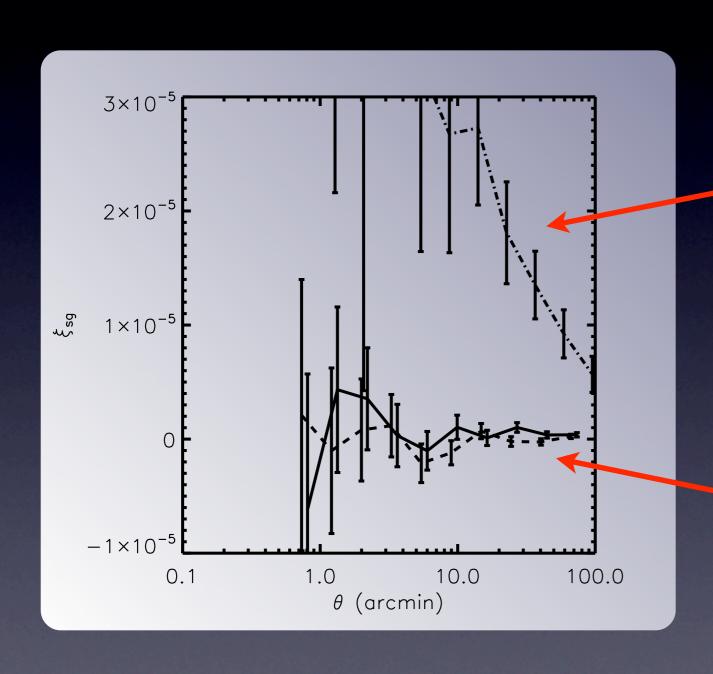


#### Advantages:

- I. small number of bins -- stable inverse covariance matrix
- 2. Virtually same information content as correlation function
- 3. Clean E/B decomposition -- removes ambiguous modes

### Systematics tests

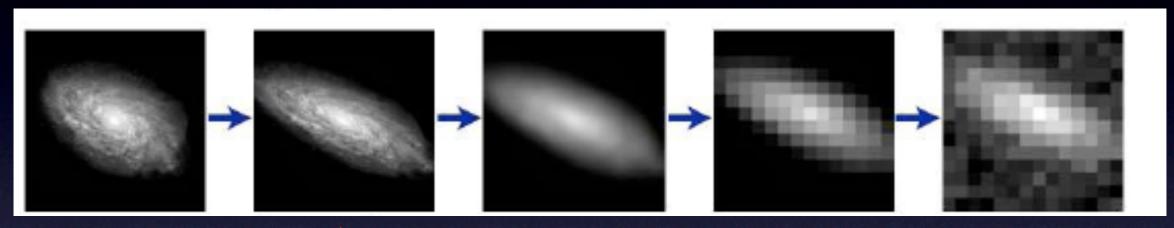
### These coadds have minimal psf shear.



Predicted Signal

Spurious psf shear

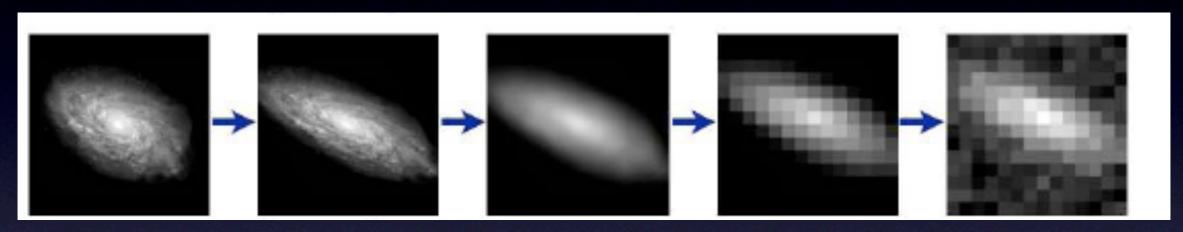
# We have tackled shear calibration and photoz errors using simulations.



Credit: Bridle et al, Great08 handbook

Apply a synthetic shear to COSMOS image

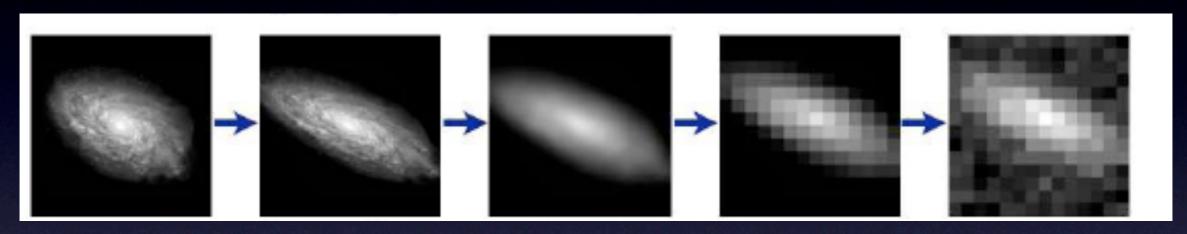
# We have tackled shear calibration and photoz errors using simulations.



Credit: Bridle et al, Great08 handbook

re-convolve with SDSS PSF

# We have tackled shear calibration and photoz errors using simulations

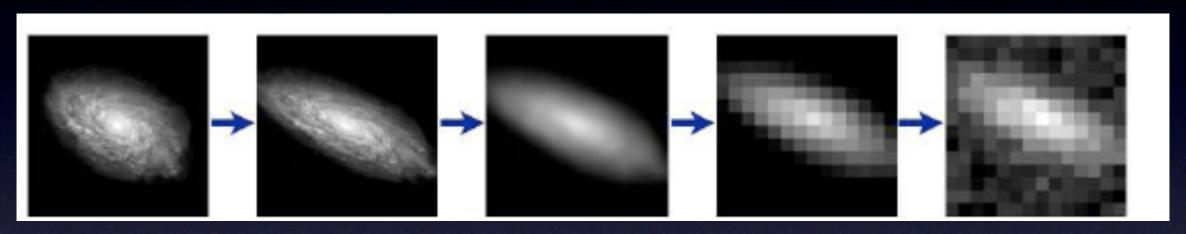


Credit: Bridle et al, Great08 handbook



insert galaxy into coadd image

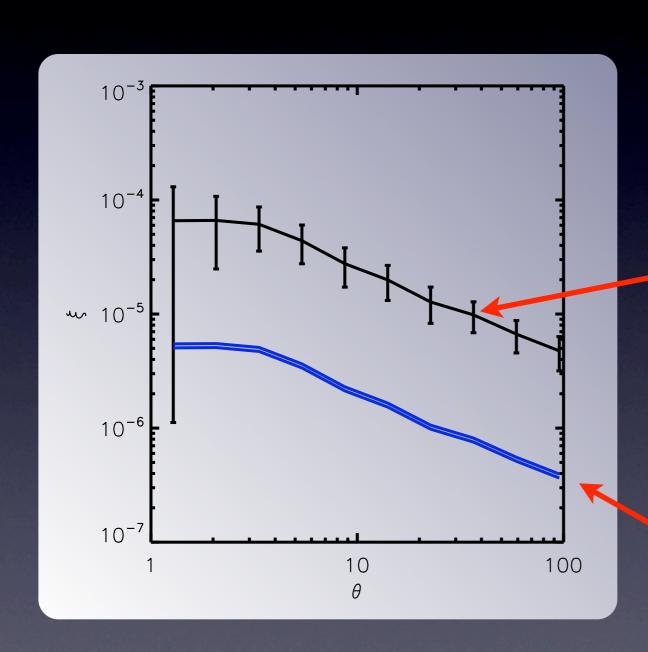
# We have tackled shear calibration and photoz errors using simulations



Credit: Bridle et al, Great08 handbook

rerun pipeline, measure shape and selection function

## Existing limits on intrinsic alignments add little uncertainty to this measurement

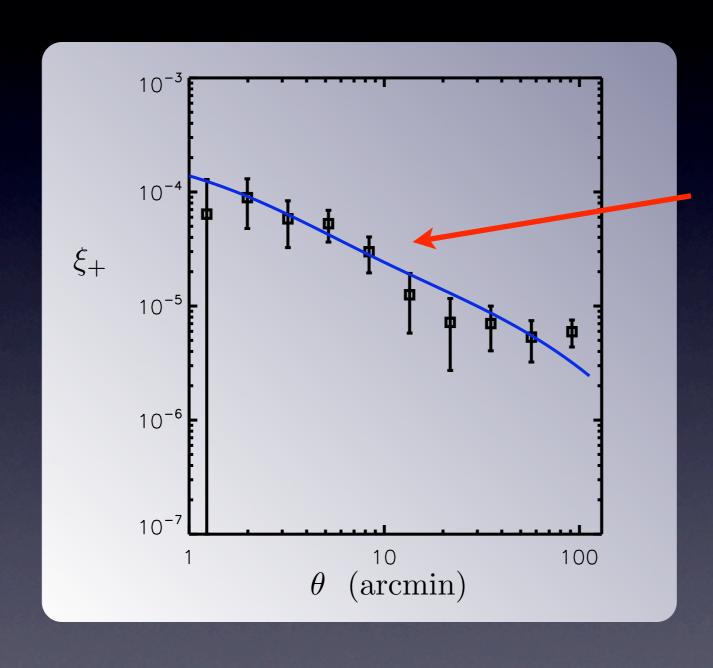


Predicted Signal

Intrinsic alignment from SDSS MegaZ-LRG (Joachimi et al 2011)

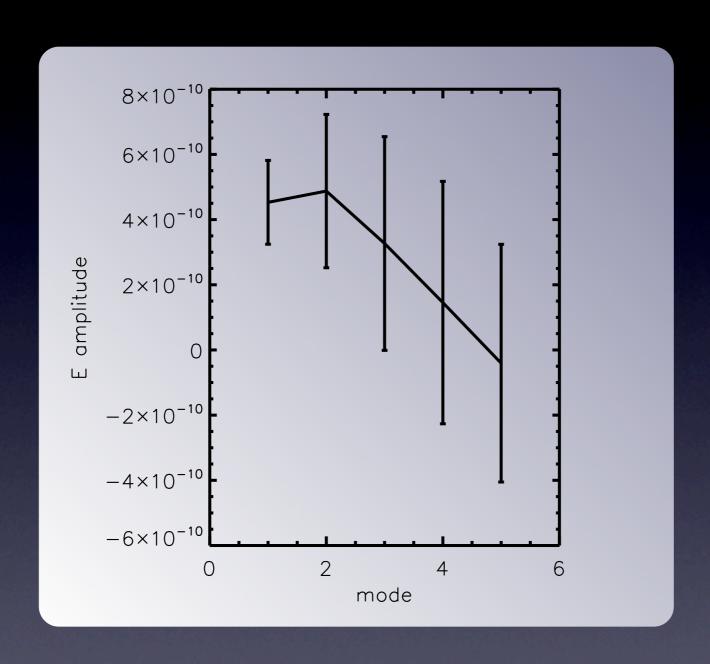
### Results

### We do detect a cosmic shear signal



WMAP7 prediction

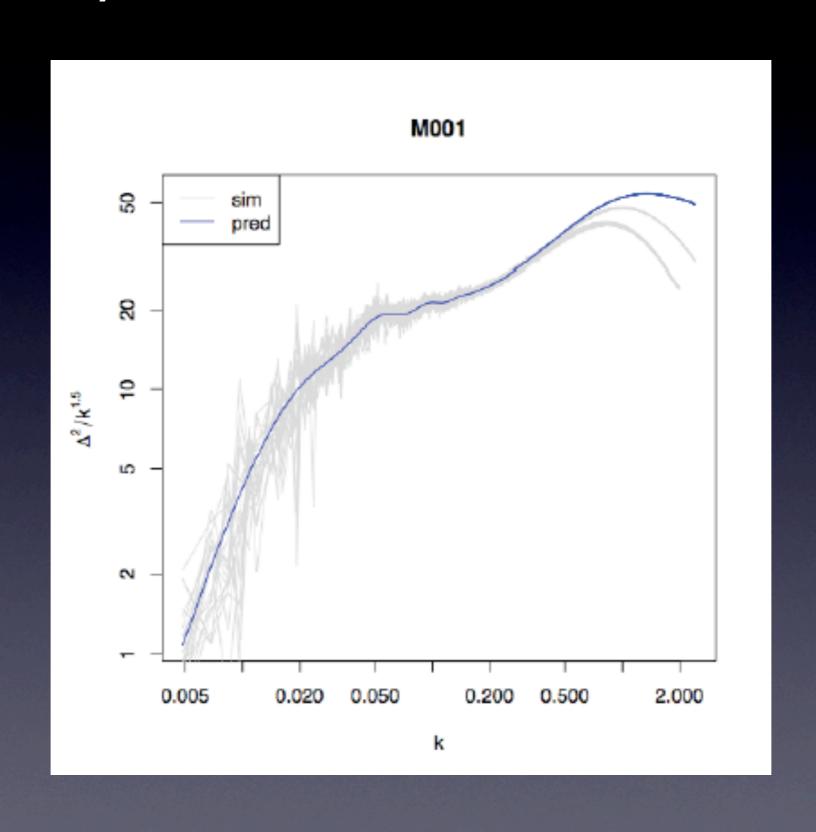
### We do detect a cosmic shear signal



We do not detect any B modes.

$$\chi^2 (B=0) = 1.05$$

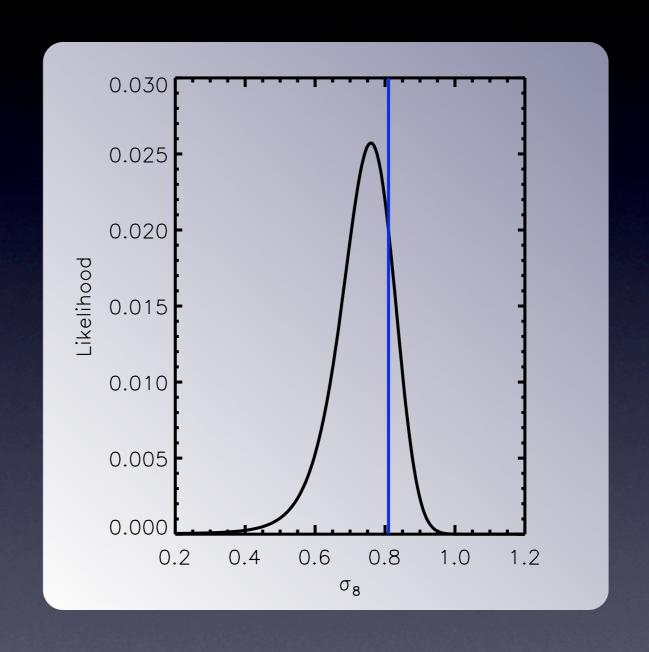
# To interpret this, we use the Coyote Universe simulation emulator



### We have an interesting constraint on structure.

 $\sigma_8 = 0.7578^{+0.070}_{-0.080}$ 

(fixing all other parameters)



#### Lessons Learned:

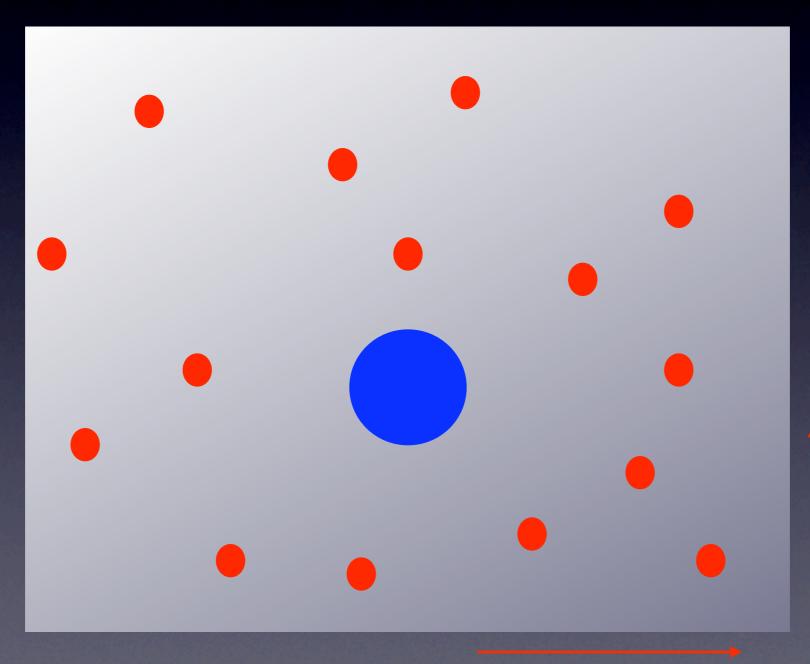
- 0. It is possible to measure cosmic shear from the ground.
- I. The rounding kernel method is effective if you have the right psf model.
- 2. Getting the right psf model is not easy.
- 3. Taking good care of systematics entails throwing away lots of data.
- 4. Low signal-to-noise is a big problem.

Is there an easier way?

Maybe.

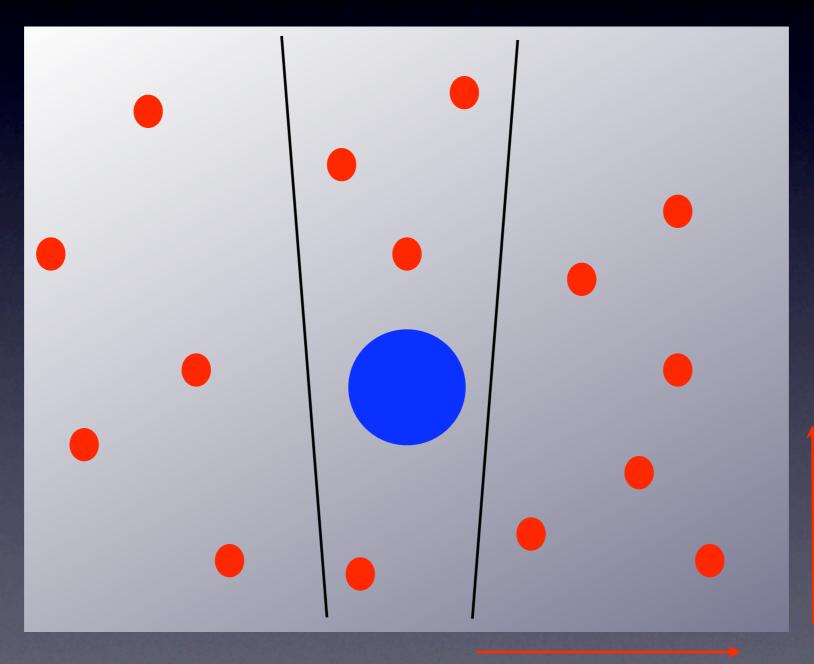
There are other components to the distortion tensor.

# The Effect of Magnification on galaxy sizes and luminosities

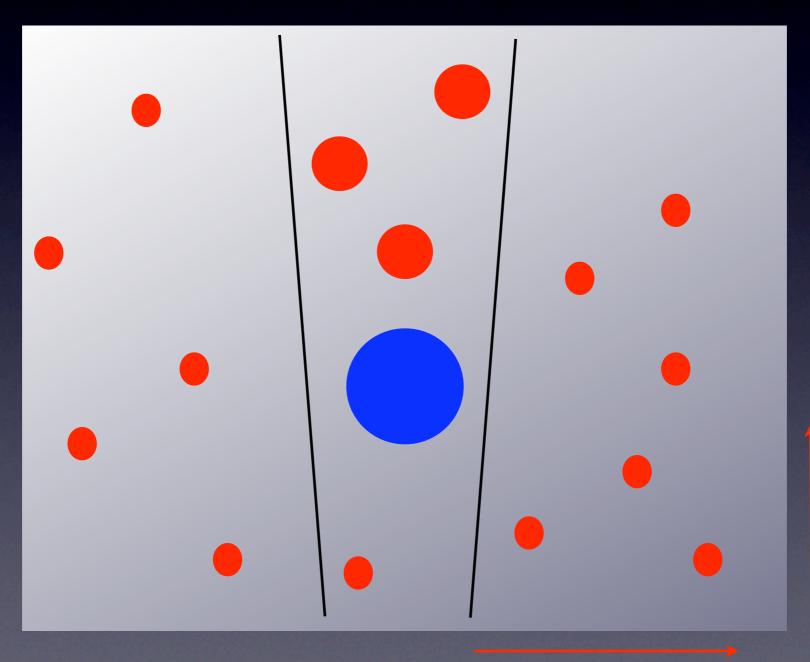


redshift

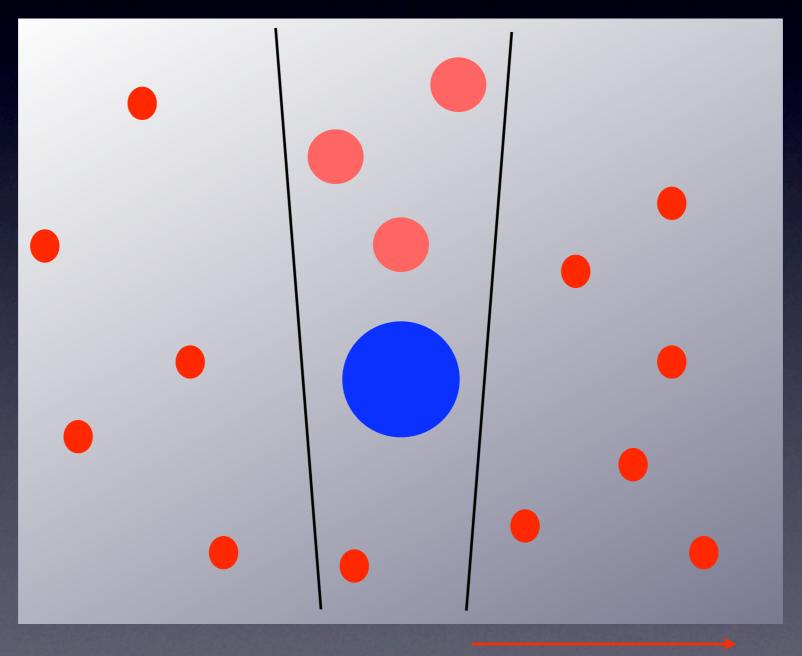
# The Effect of Magnification on galaxy sizes and luminosities



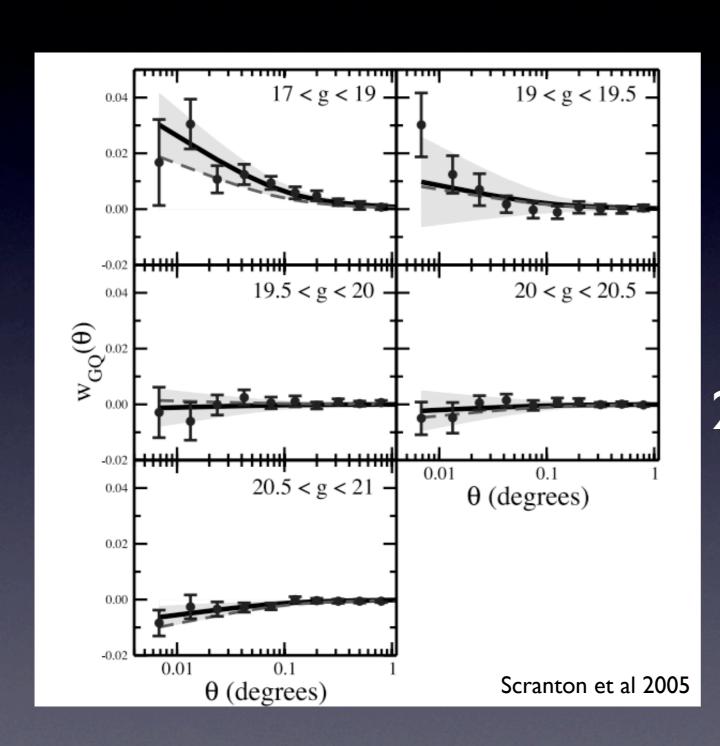
## The Effect of Magnification on galaxy sizes and luminosities



# The Effect of Magnification on galaxy sizes and luminosities

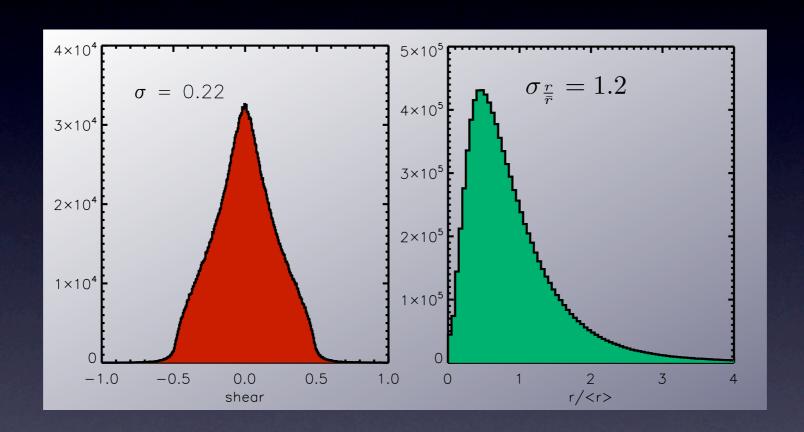


## The effect of Magnification on Luminosities



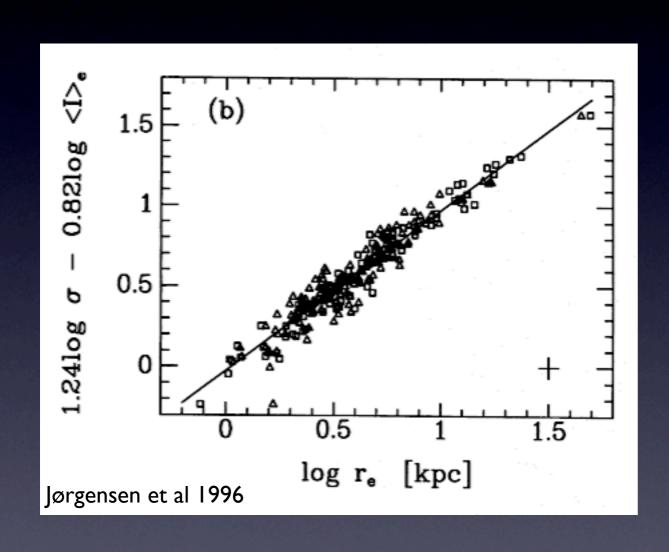
A heroic effort: 13.5 million galaxy lenses 225,000 quasar sources

#### Why shear is still much better than the alternatives:



We want a way to reduce the intrinsic scatter.

#### The Fundamental Plane of Early Type Galaxies

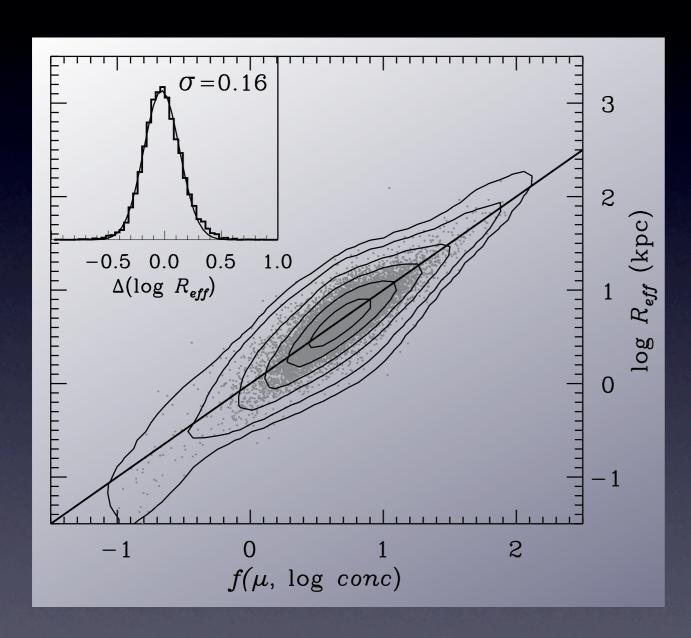


~15% intrinsic scatter

no detected variation with environment

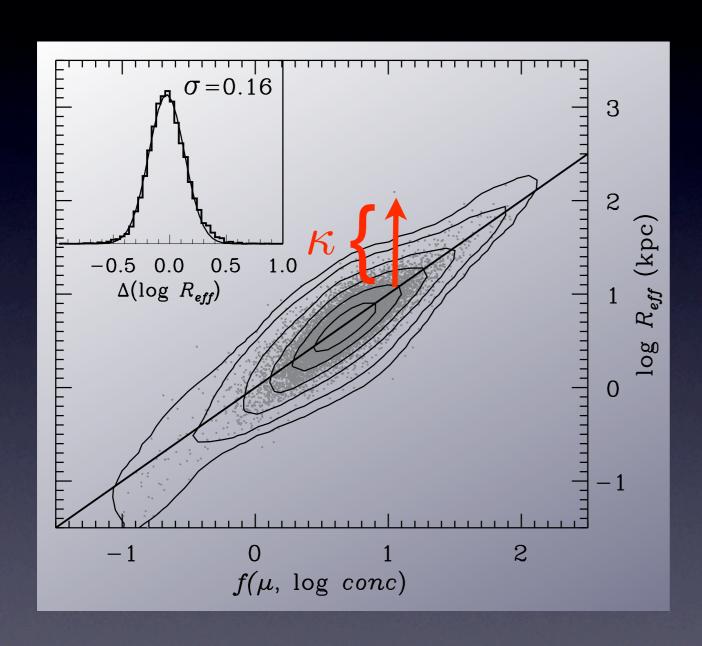
a photometric analogue exists

### The Effect of Magnification on the Photometric Plane



at fixed mass, concentration and effective radius are inversely correlated

## The Effect of Magnification on the Photometric Plane



$$\kappa = \log (R_{eff}) - f(\mu, \log conc)$$

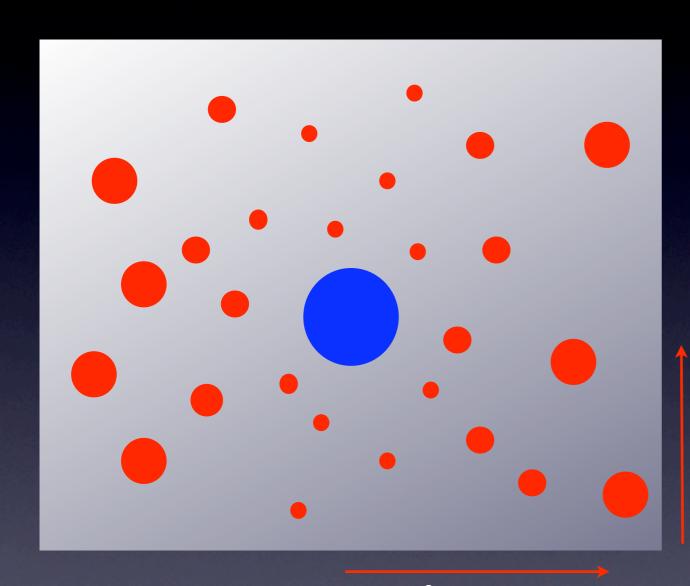
#### Constructing a Sample using SDSS

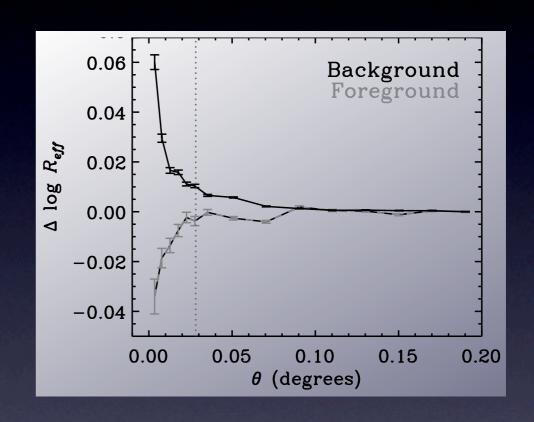
60,000 Lenses:
log (stellar mass) > 11.0
0.2 < z < 0.4

10 million Sources:
resolved galaxies
early-type SEDs (35%)

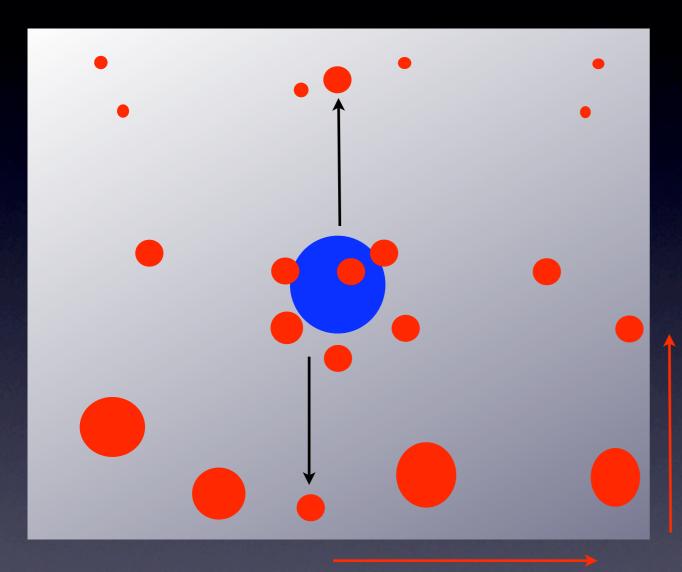


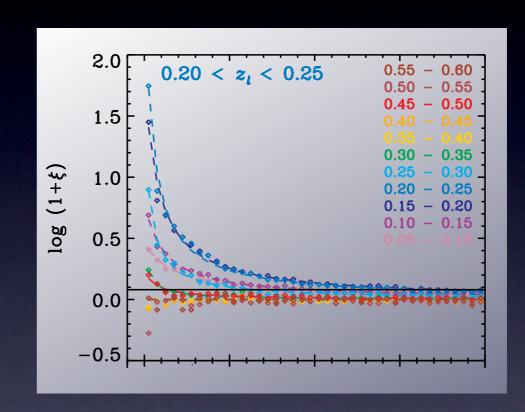
### Systematics: Sky Subtraction



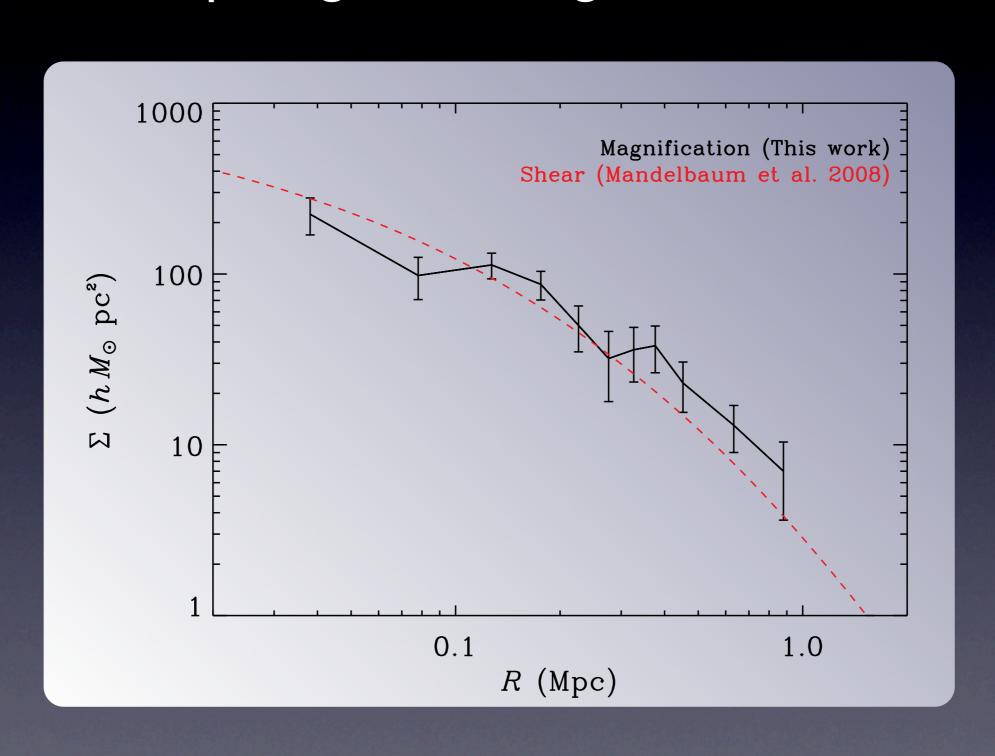


## Systematics: Source Clustering with Photo-z's





## Lensing Detection: Comparing to Existing Measurements



 Currently: we can control systematics in ground-based cosmic shear

 Galaxy scaling relations can yield much more weak lensing signal

