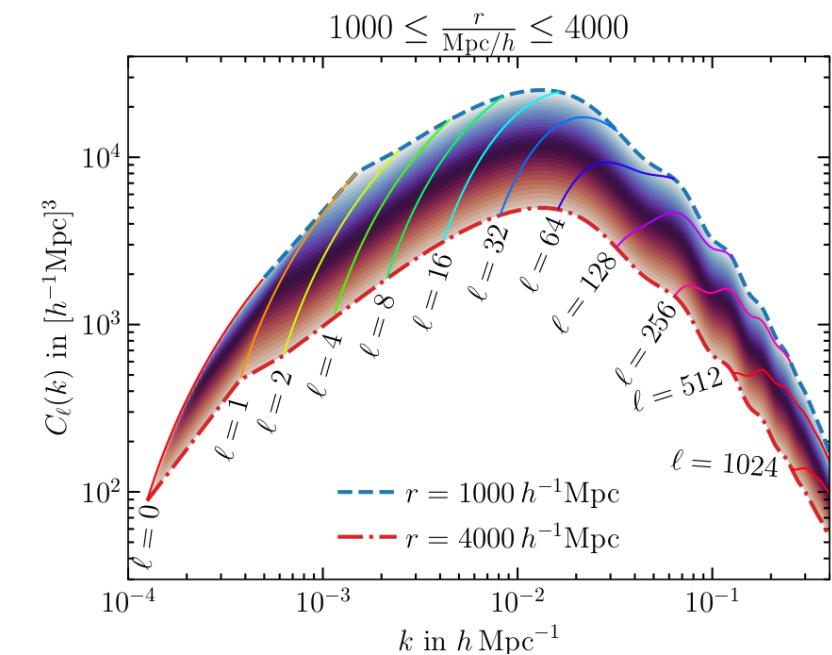
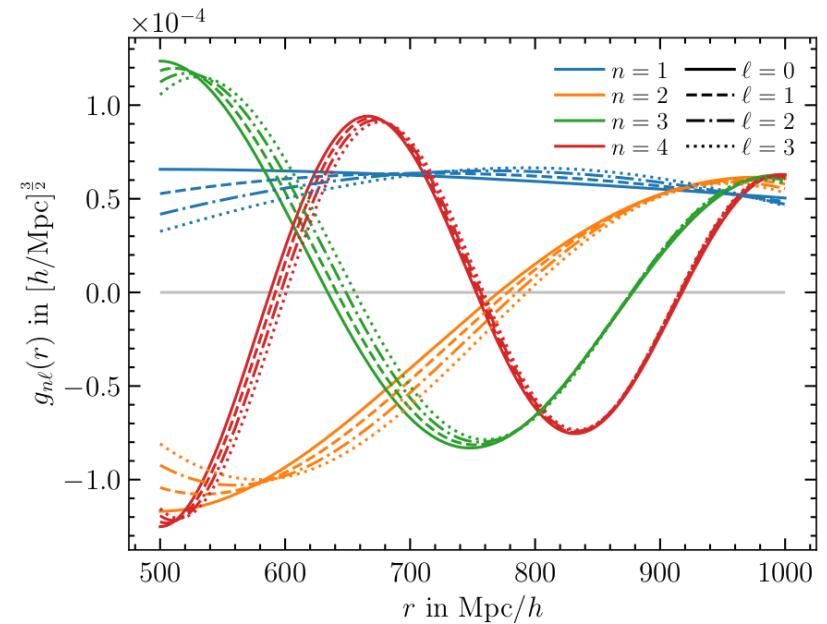


# SuperFaB: a fabulous code for Spherical Fourier-Bessel decomposition



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located at  
Jet Propulsion Laboratory,  
California Institute of Technology

[arXiv:2102.10079](https://arxiv.org/abs/2102.10079)



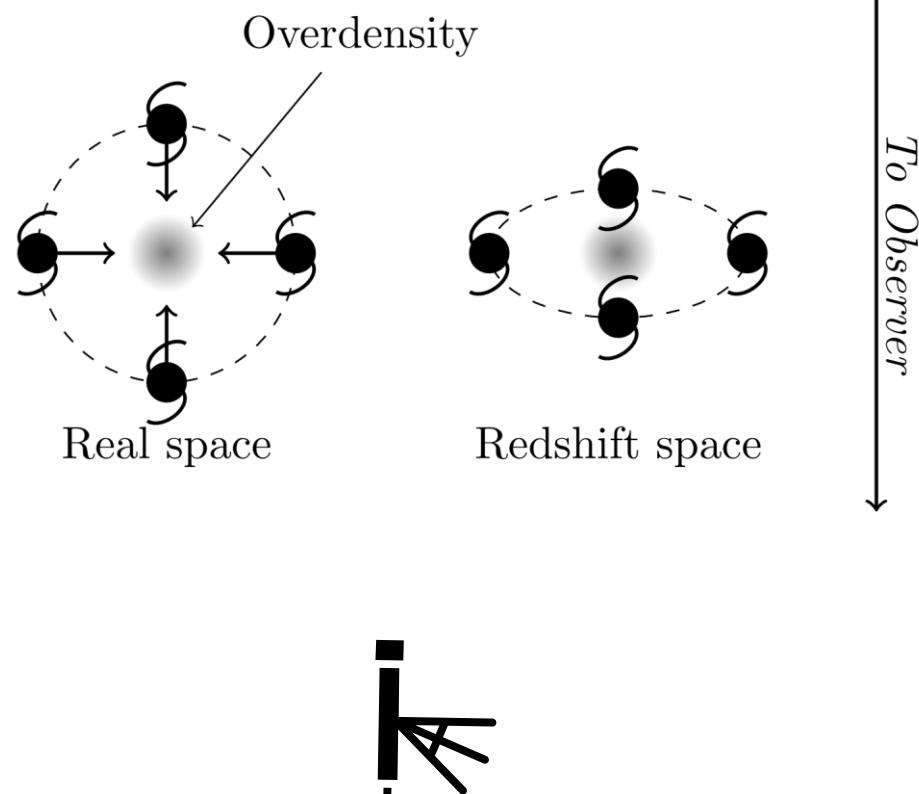
# Outline

- Motivation: Cosmology with the Power Spectrum
  - RSD
  - Redshift-evolution
- Motivation: Spherical Fourier-Bessel (SFB) vs. Fourier
- Intuition for the SFB power spectrum
- Measuring SFB modes
- Results: *Roman*-, *SPHEREx*-, *Euclid*-like

# Motivation: Roman, SPHEREx, Euclid, PFS, DESI, ...

- Galaxy surveys bridge the gap between low-redshift supernova Ia and high-redshift CMB measurements. (e.g. BAO)
- Redshift-space distortions (RSD) probe the growth of structure over cosmic time. → Tests for modified gravity.
- Large scales especially important for non-Gaussianity.

# Redshift-space distortions I: Kaiser effect



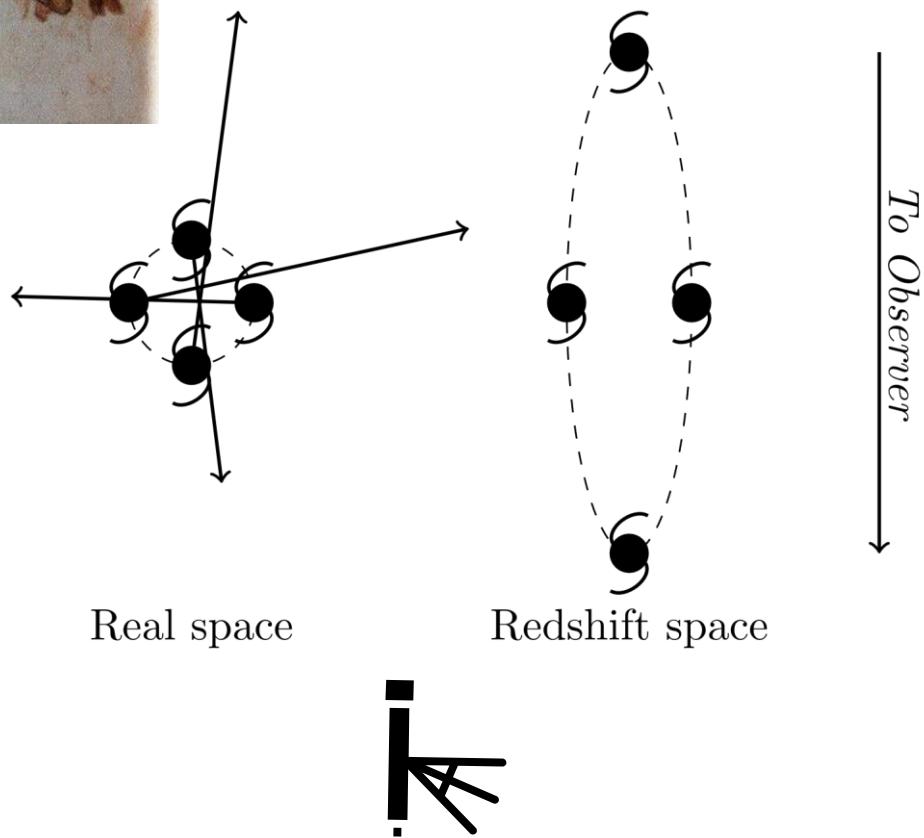
Coherent flow of galaxies towards overdensities leads to **increased clustering** along the line of sight direction.

→ Non-zero quadrupole!

$$s = r + \frac{v_{\parallel}}{aH}$$

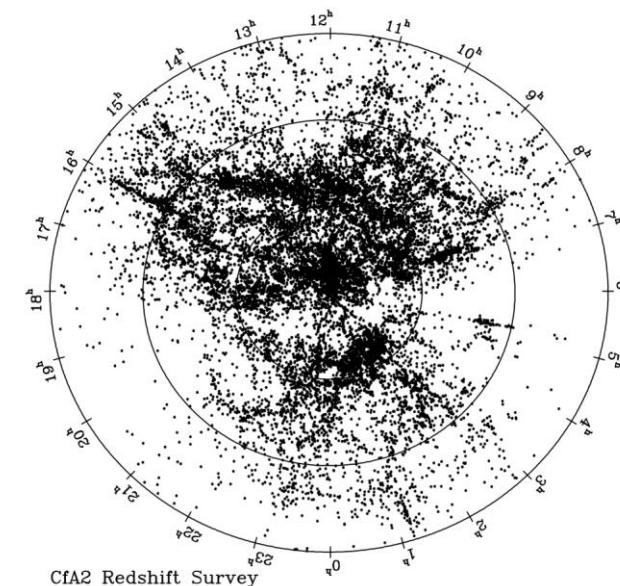


# Redshift-space distortions II: fingers of God

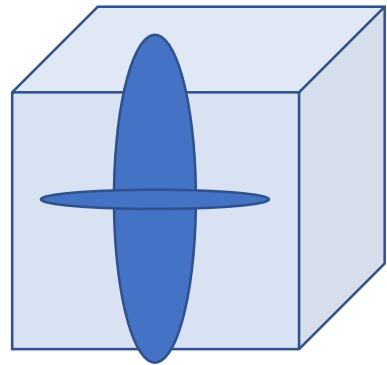
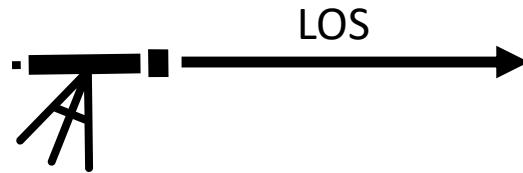


Stochastic motions of galaxies lead to a **suppression** of clustering along line-of-sight direction.

→ Small-scale suppression of power.



# Homogeneity and Isotropy with Redshift-space Distortions (RSD)



$$\delta(\vec{r}) = \frac{n(\vec{r}) - \bar{n}}{\bar{n}}$$

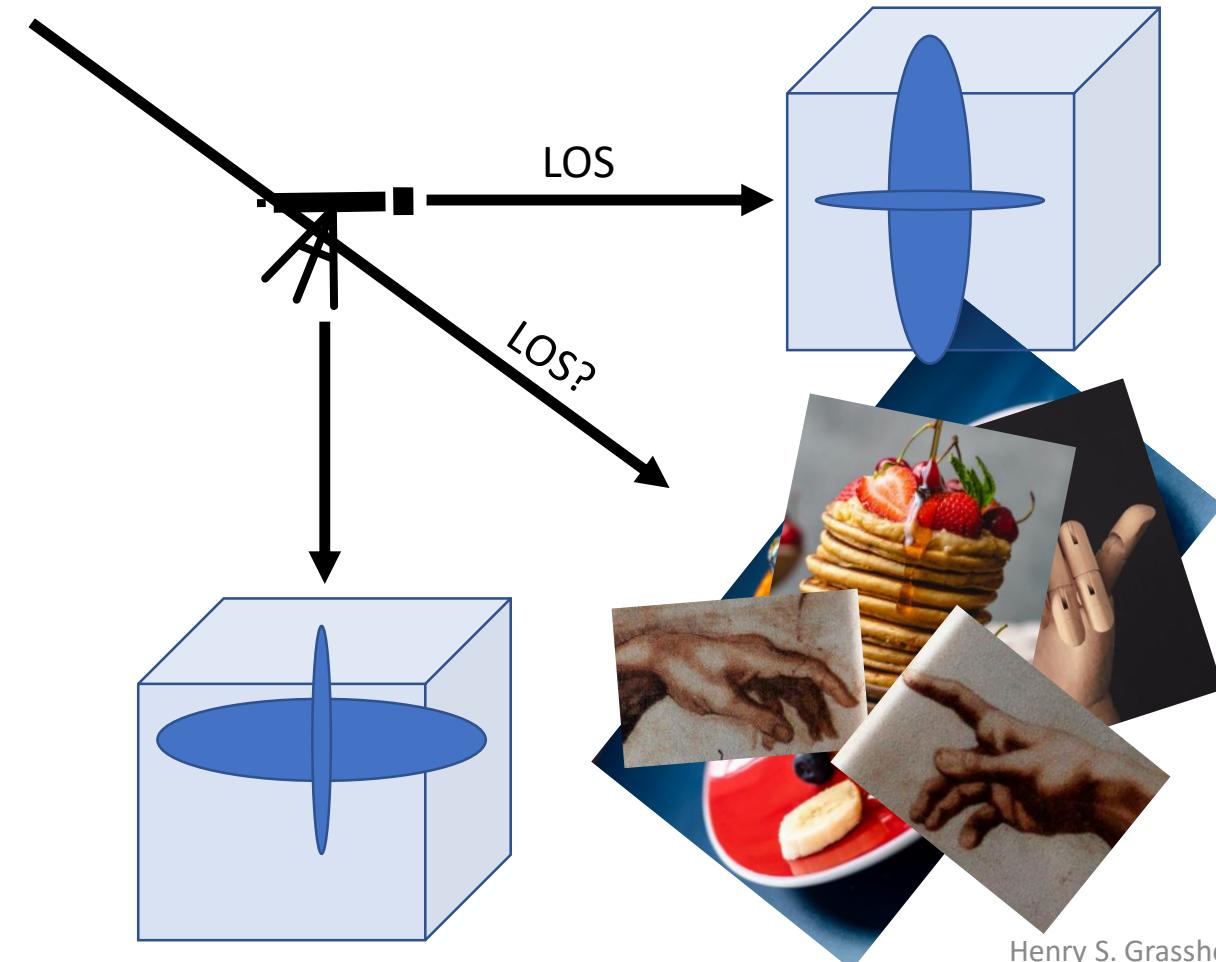
$$\delta(\vec{k}) = \int d^3 r e^{-i\vec{k}\cdot\vec{r}} \delta(\vec{r})$$

$$\langle \delta(\vec{k}) \delta^*(\vec{k}') \rangle = (2\pi)^3 \delta^D(\vec{k}' - \vec{k}) P(\vec{k})$$

Isotropy

Nope:  $P(k_\perp, k_\parallel) = (b + f\mu^2)^2 P(k)$

# ~~Homogeneity~~ and ~~Isotropy~~ with Redshift-space Distortions (RSD)



$$\delta(\vec{r}) = \frac{n(\vec{r}) - \bar{n}}{\bar{n}}$$

$$\delta(\vec{k}) = \int d^3r e^{-i\vec{k}\cdot\vec{r}} \delta(\vec{r})$$

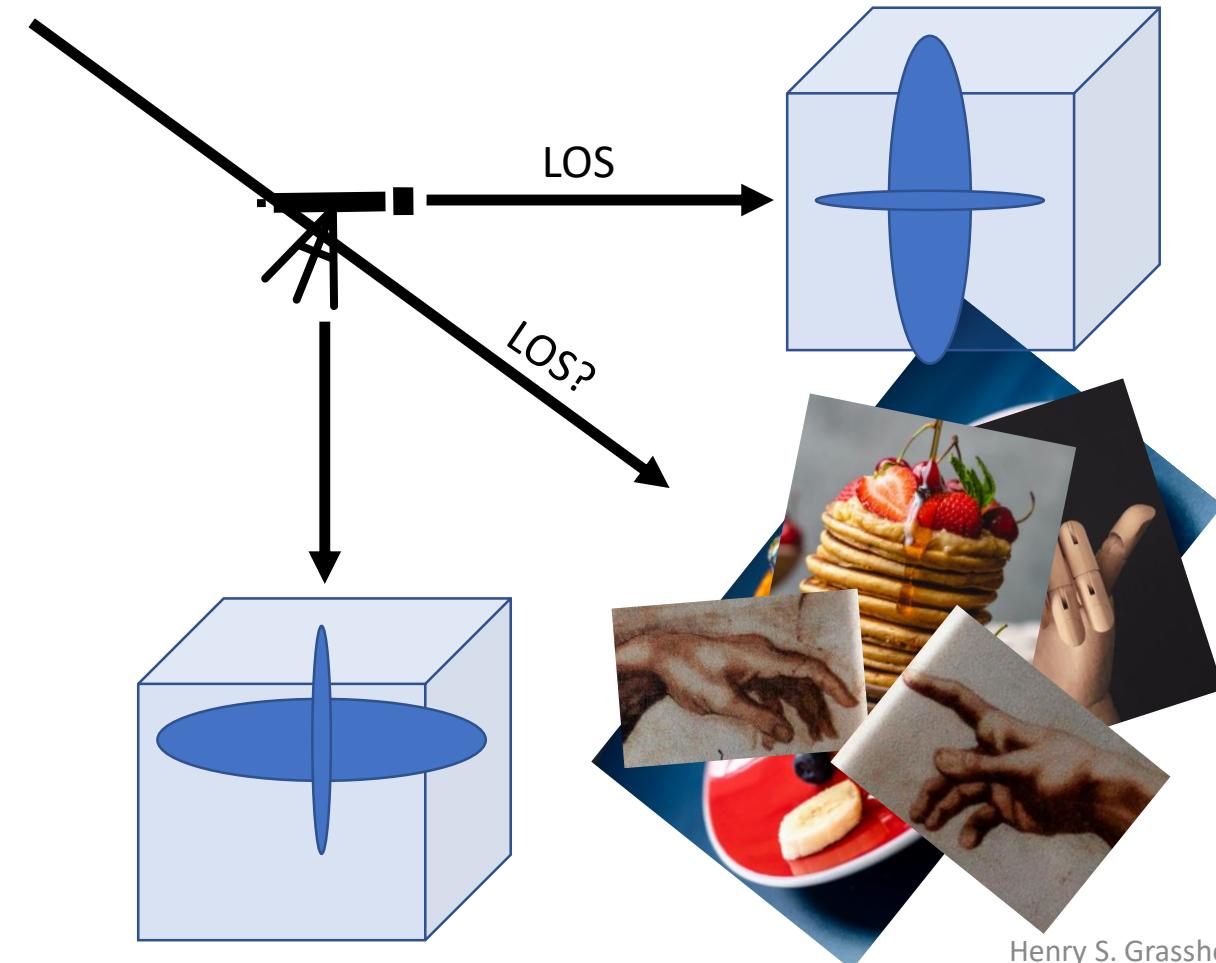
$$\langle \delta(\vec{k}) \delta^*(\vec{k}') \rangle = (2\pi)^3 \delta^D(\vec{k}' - \vec{k}) P(\vec{k})$$

Homogeneity      Isotropy

Nope:  $P(k_\perp, k_\parallel) = (b + f\mu^2)^2 P(k)$

Nope: Wide angles

# ~~Homogeneity~~ and ~~Isotropy~~ with Redshift-space Distortions (RSD)



$$\delta(\vec{r}) = \frac{n(\vec{r}) - \bar{n}}{\bar{n}}$$

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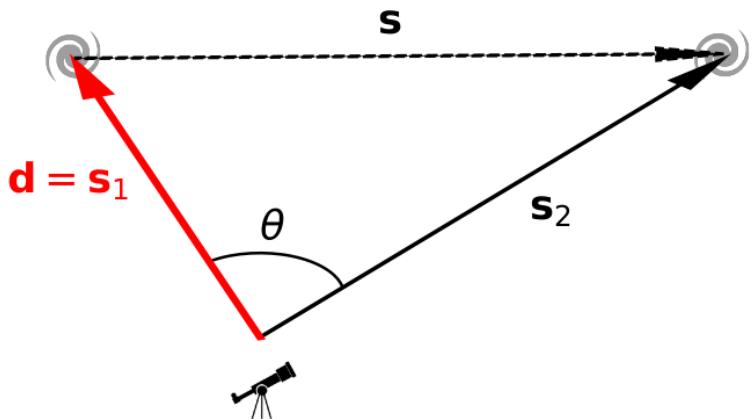
Homogeneity      Isotropy

Nope:  $P(k_\perp, k_\parallel) = (b + f\mu^2)^2 P(k)$

Nope: Wide angles



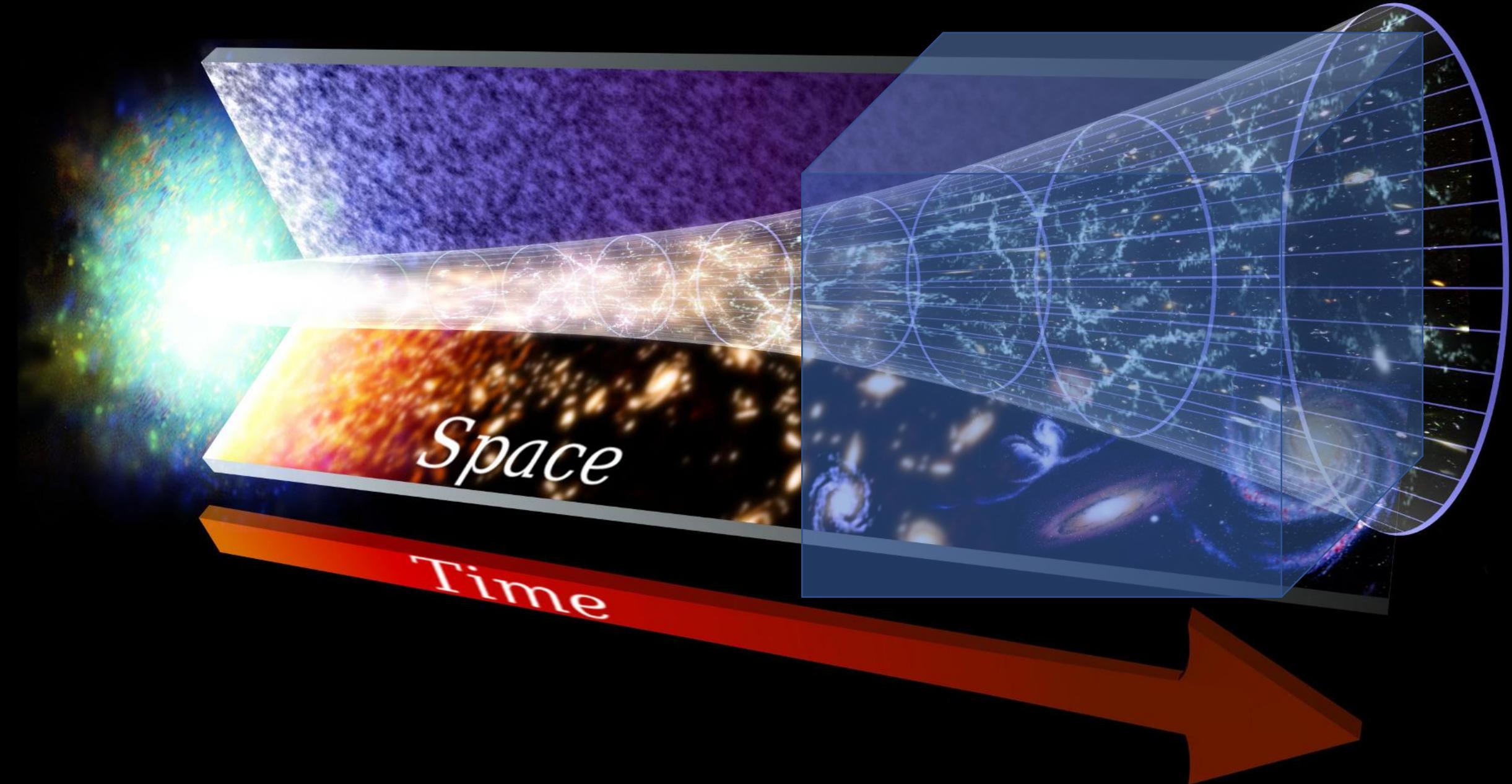
# Yamamoto to the rescue: Great for small angles!



- Choose one LOS per galaxy pair.
- Perturbative expansion for modeling.
- Suboptimal for the largest angles, because RSD can cancel.

(Figure from Beutler, Castorina, Zhang 2019)

Redshift-evolution: Our past lightcone is **not** homogeneous.



# What to do?

Back to Basics: Choose a coordinate system that is better suited for the fixed position of the observer than a Cartesian.

Laplacian eigenfunctions:  $\nabla^2 f(\vec{r}) = -k^2 f(\vec{r})$

Cartesian Coordinates:

$$f(\vec{r}) = e^{-i\vec{k}\cdot\vec{r}}$$

Fourier transform:

$$\delta(\vec{k}) = \int d^3r e^{-i\vec{k}\cdot\vec{r}} \delta(\vec{r})$$

Spherical Coordinates:

$$f(\vec{r}) = j_\ell(kr) Y_{\ell m}(\hat{r})$$

Spherical Fourier-Bessel transform:

$$\delta_{\ell m}(k) = \int d^3r j_\ell(kr) Y_{\ell m}^*(\hat{r}) \delta(\vec{r})$$

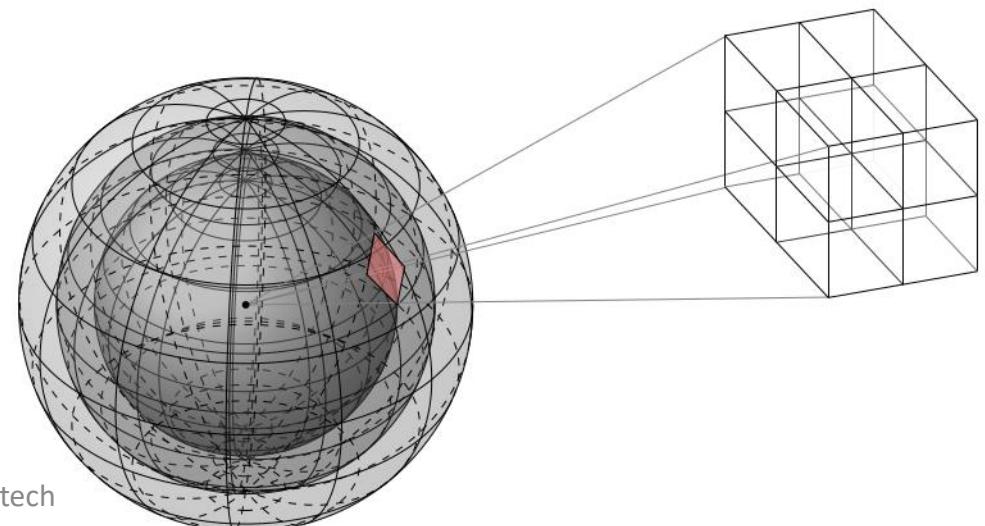
# SFB power spectrum

$$\langle \delta_{\ell m}(k) \delta_{\ell' m'}(k') \rangle = \underbrace{\delta^K_{\ell\ell'} \delta^K_{mm'}}_{\text{Isotropy}} C_\ell(k, k')$$

$$C_\ell(k, k') = \underbrace{\delta^D(k' - k)}_{\text{Homogeneity}} P(k)$$

# The SFB power spectrum is ideally suited to deep and wide galaxy surveys

- Natural separation between angular and radial coordinates
  - Individual line of sights for each galaxy
    - Maximal information from RSD
  - All wide-angle effects
  - Redshift-evolution (e.g., growth of the non-linear power spectrum)
- Nearly diagonal covariance matrix



# Limber's Approximation

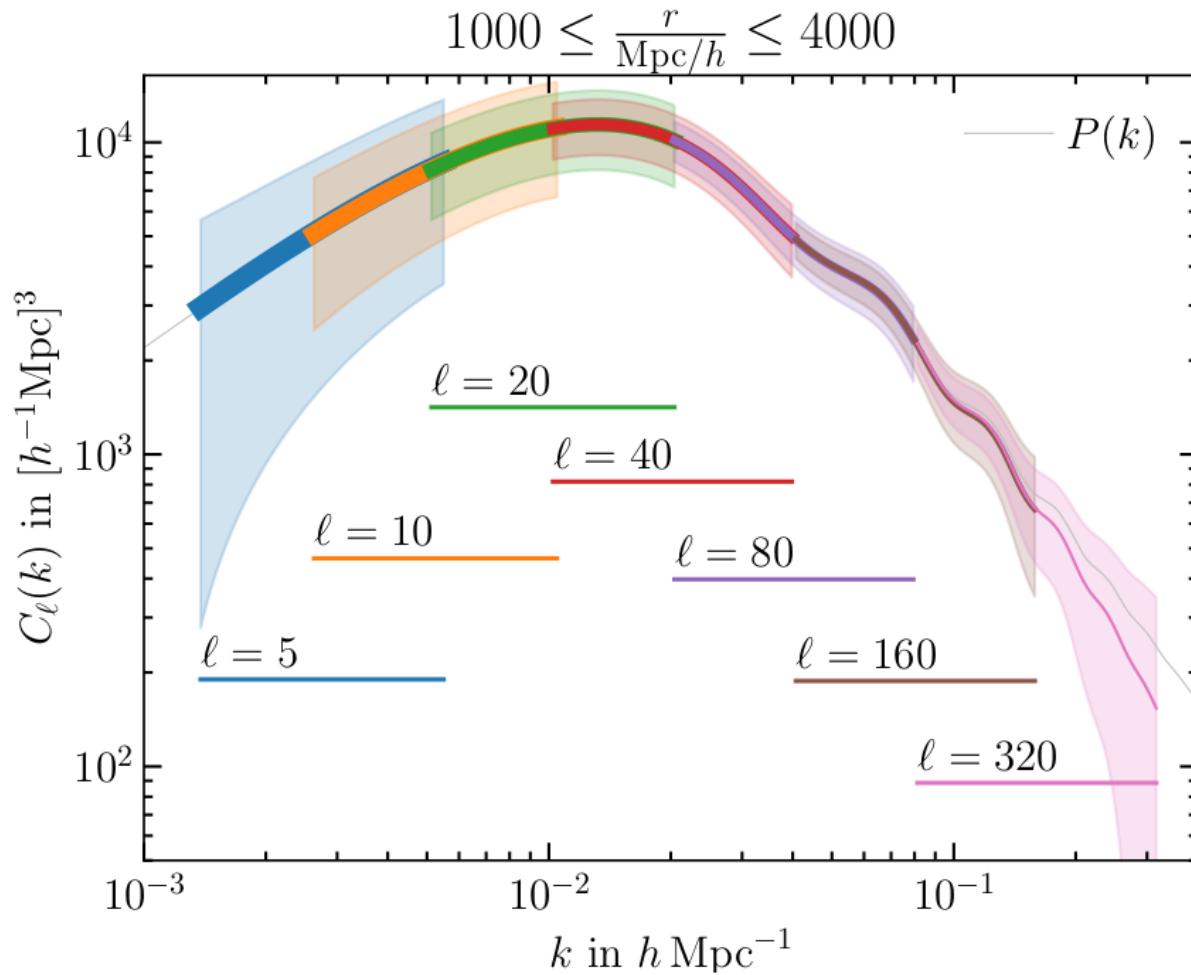
$$C_\ell(k, k') = P(k) e^{-\sigma_u^2 k^2} \delta^D(k - k')$$
$$\times \phi^2\left(\frac{\ell + \frac{1}{2}}{k}\right) D^2\left(\frac{\ell + \frac{1}{2}}{k}\right) b^2\left(\frac{\ell + \frac{1}{2}}{k}, k\right)$$
$$\times [1 - \beta(f_{-2}^\ell + f_0^\ell + f_2^\ell)]^2.$$

Diagram illustrating the components of Limber's Approximation:

- Power spectrum**:  $P(k)$  (blue circle)
- Fingers of God**:  $e^{-\sigma_u^2 k^2}$  (green oval)
- Redshift evolution**:  $\phi^2\left(\frac{\ell + \frac{1}{2}}{k}\right)$  (orange oval)
- Scale-dependent bias**:  $b^2\left(\frac{\ell + \frac{1}{2}}{k}, k\right)$  (red oval)
- Linear Kaiser effect**:  $[1 - \beta(f_{-2}^\ell + f_0^\ell + f_2^\ell)]^2$  (green oval)
- SFB power spectrum**:  $C_\ell(k, k')$  (overall expression)

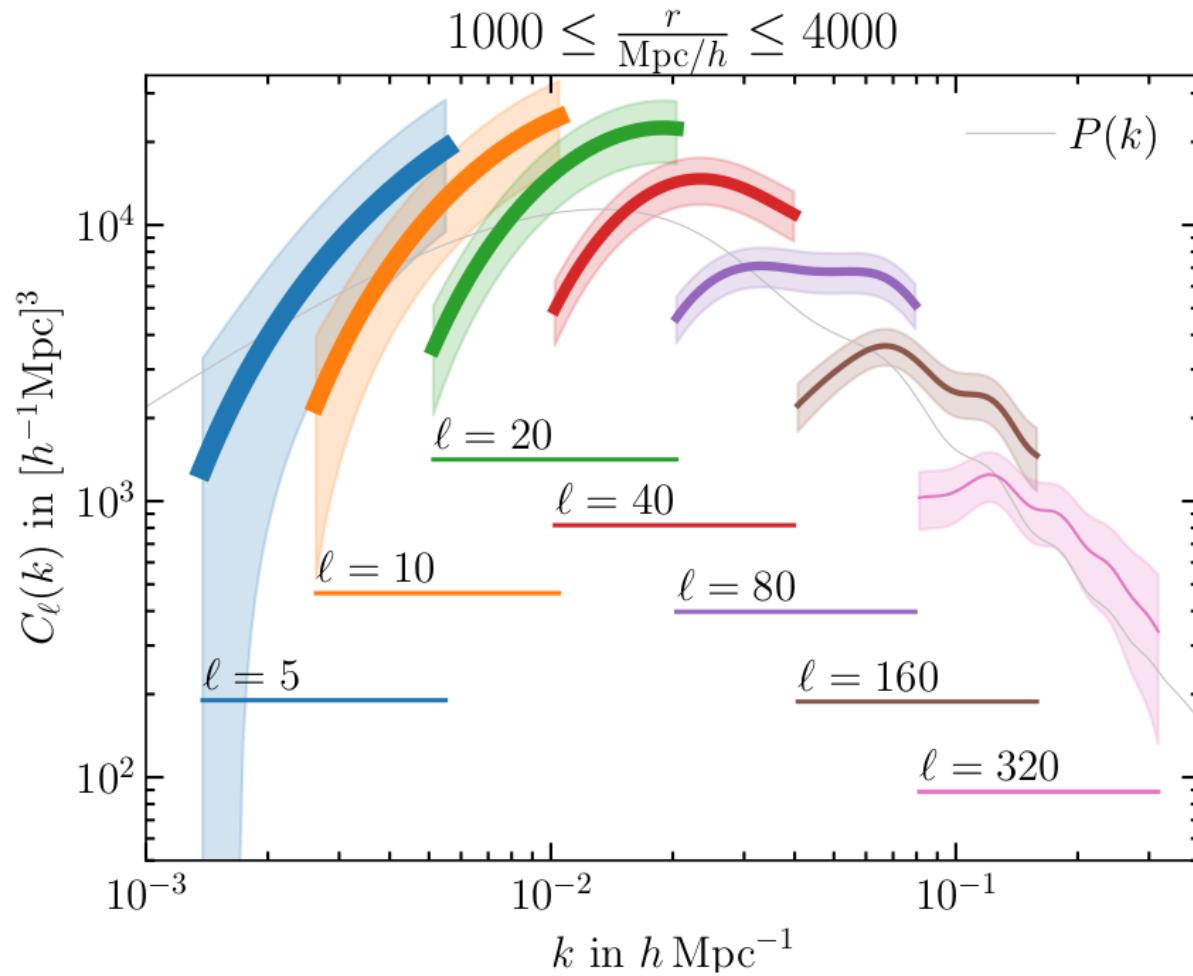
# The SFB power spectrum $C_\ell(k, k')$

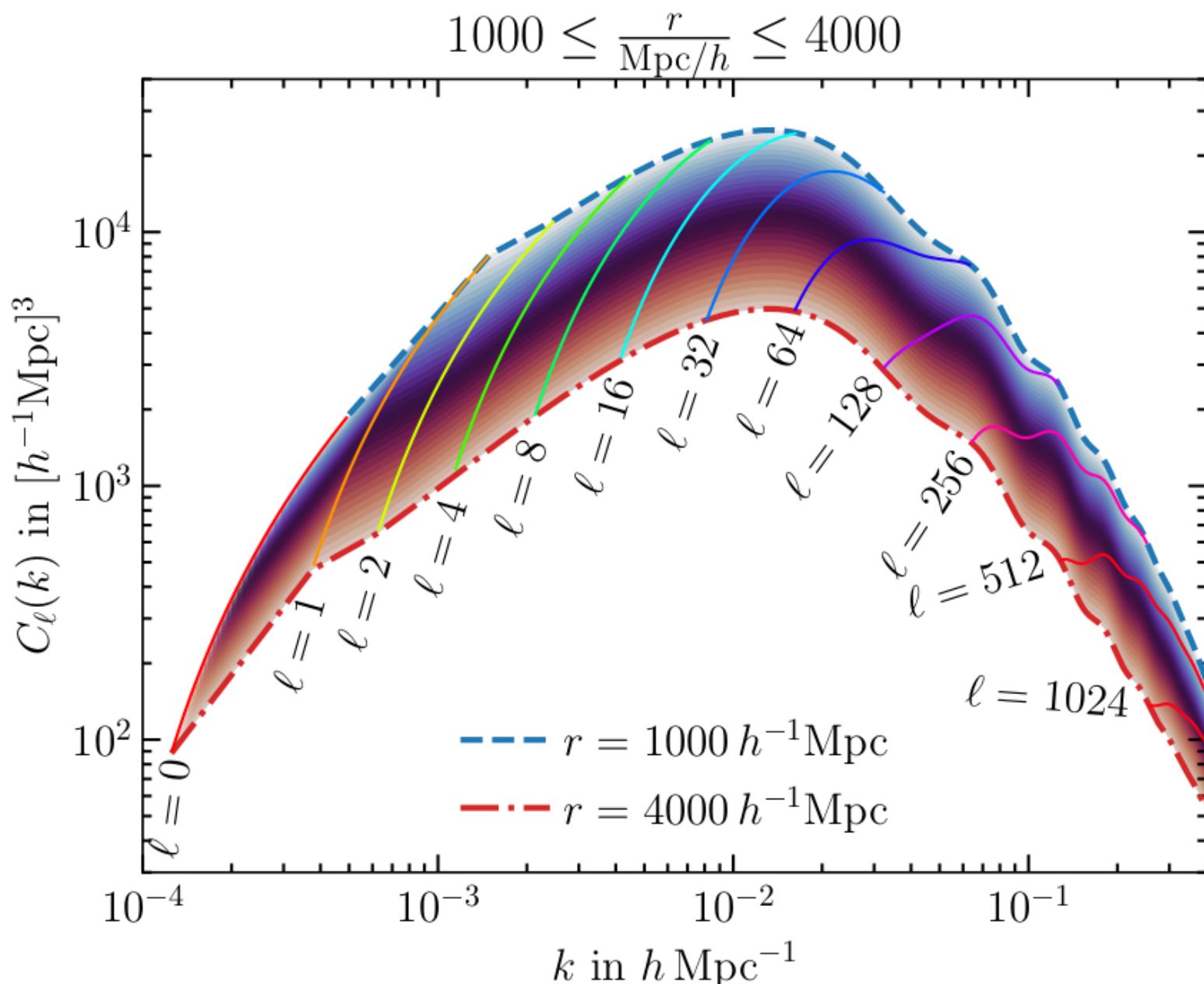
- $b(z) \propto \frac{1}{D(z)}$
- Limber distance:  
$$r = \frac{\ell + \frac{1}{2}}{k}$$



# The SFB power spectrum $C_\ell(k, k')$

- $b(z) = \text{const}$
- Limber distance:  
$$r = \frac{\ell + \frac{1}{2}}{k}$$





# Why a new code?

- We follow Samushia (2019) to introduce boundary conditions at both  $r_{min}$  and  $r_{max}$

$$\nabla^2 f = -k^2 f$$

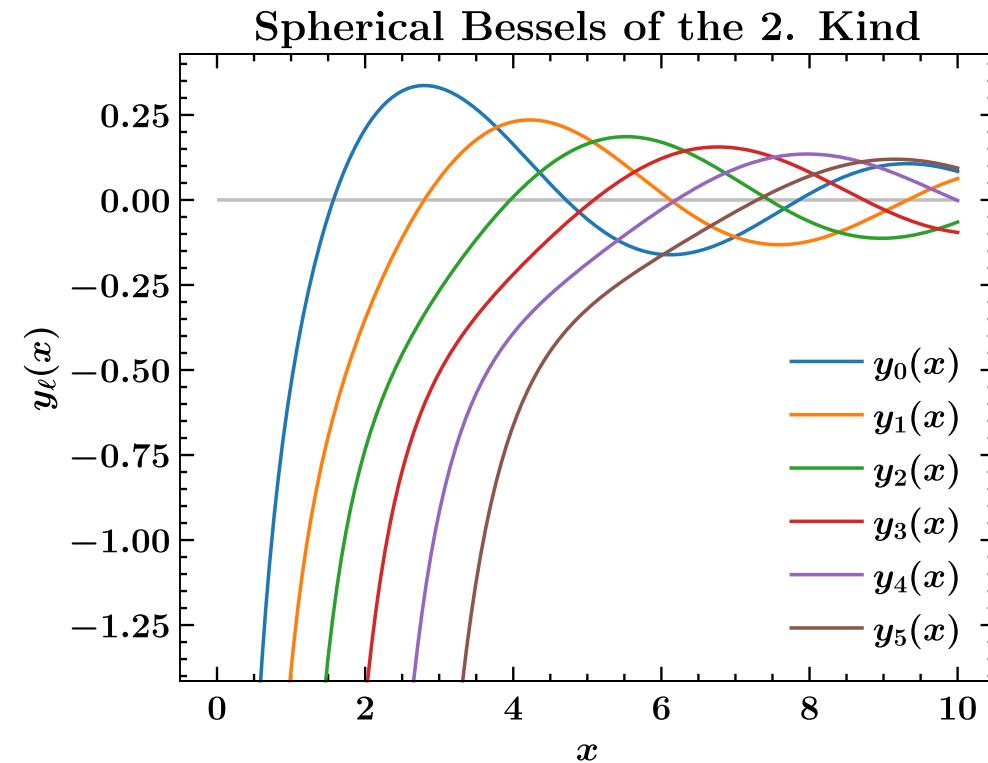
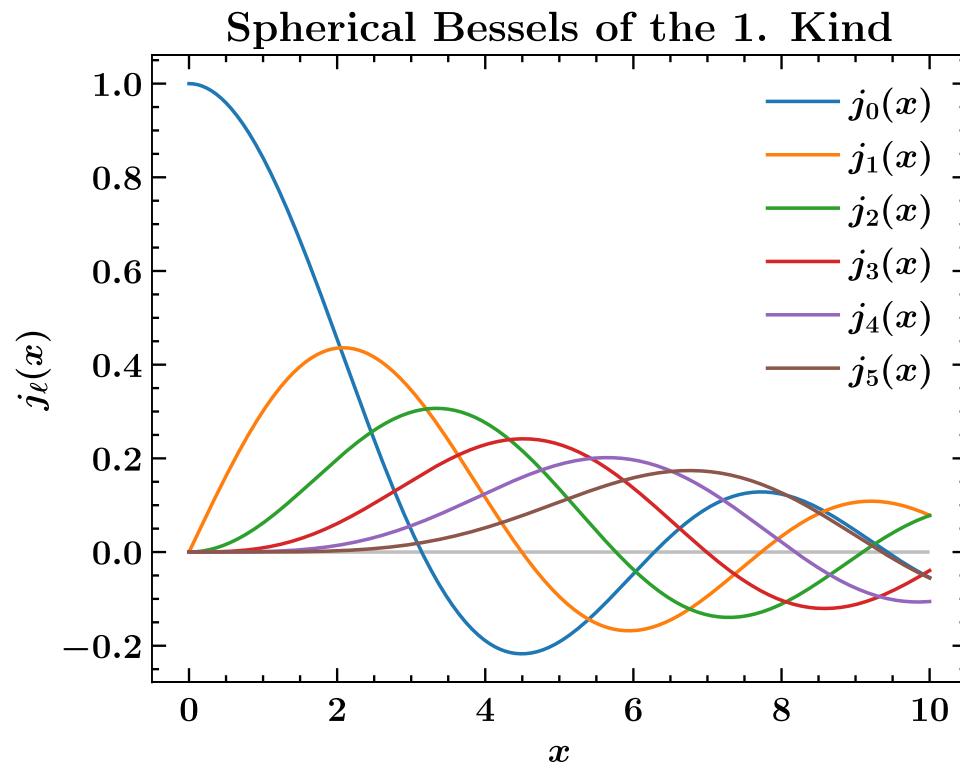
Spherical Bessels of 1. and 2. kind

$$f_{\ell\mu}(k; r, \theta, \phi) = [c_j j_\ell(kr) + c_y y_\ell(kr)] \times [c_p P_\ell^\mu(\cos \theta) + c_q Q_\ell^\mu(\cos \theta)] \times [c_+ e^{i\mu\phi} + c_- e^{-i\mu\phi}],$$

Spherical harmonics  $Y_{\ell m}$

- Numerically more stable
- We use potential boundary conditions

# Spherical Bessel and Neumann functions



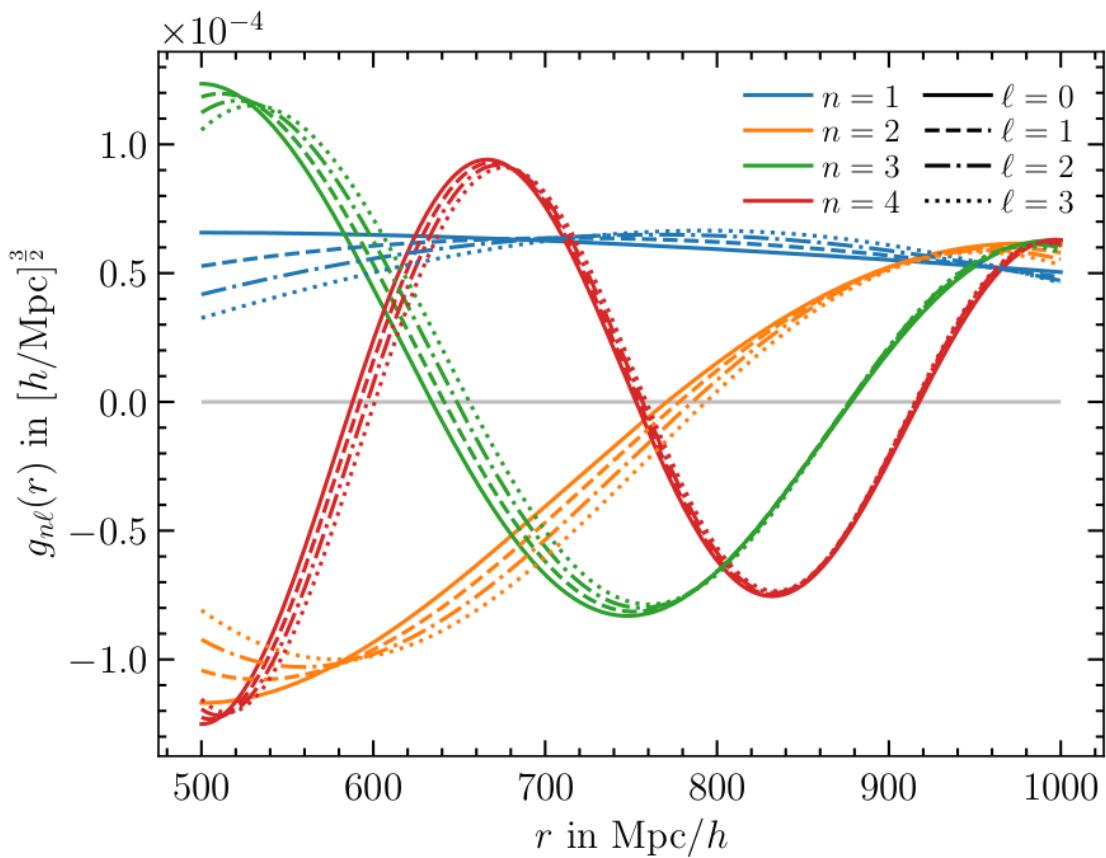
$$\begin{aligned} f_{\ell\mu}(k; r, \theta, \phi) &= [c_j j_\ell(kr) + c_y y_\ell(kr)] \\ &\times [c_p P_\ell^\mu(\cos \theta) + c_q Q_\ell^\mu(\cos \theta)] \\ &\times [c_+ e^{i\mu\phi} + c_- e^{-i\mu\phi}], \end{aligned}$$

# Radial basis functions

- Boundary conditions at both  $r_{min}$  and  $r_{max}$
- Spherical Bessels of the second kind

$$g_{n\ell}(r) = c_{n\ell} j_\ell(k_{n\ell} r) + d_{n\ell} y_\ell(k_{n\ell} r)$$

➤ Orthogonal basis where the survey is sensitive (Samushia 2019)

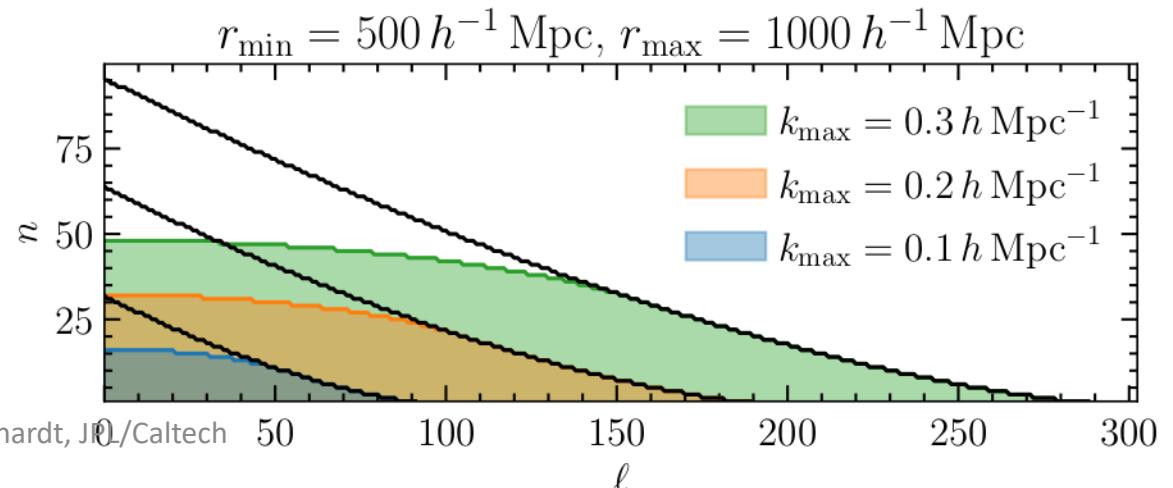


# Basic Algorithm

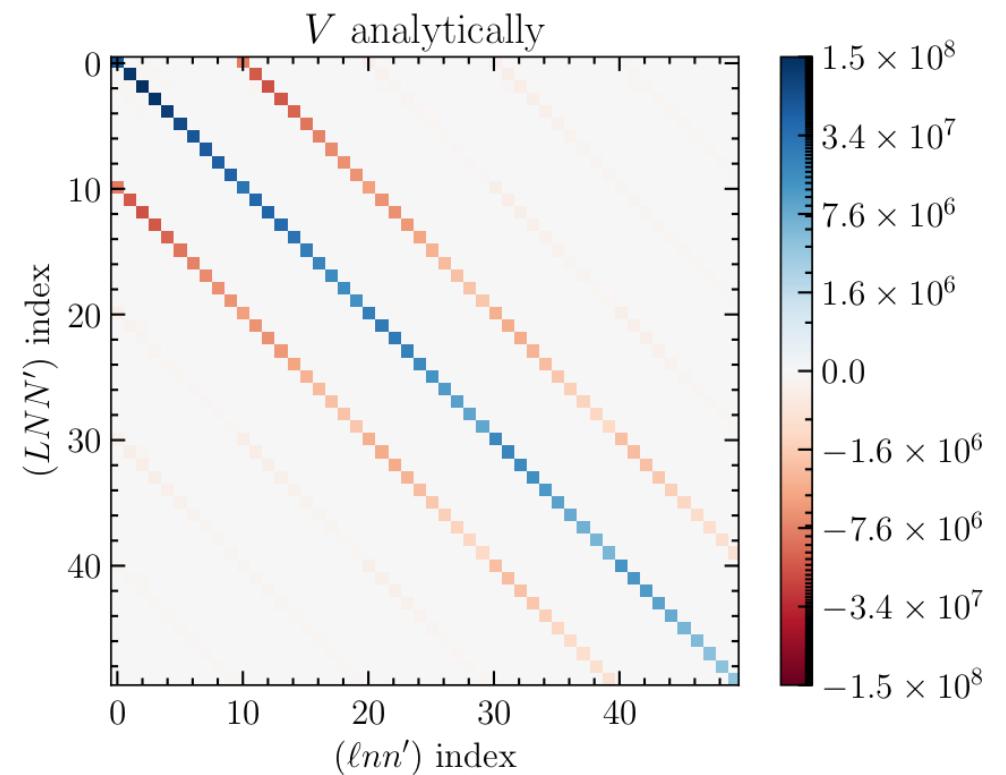
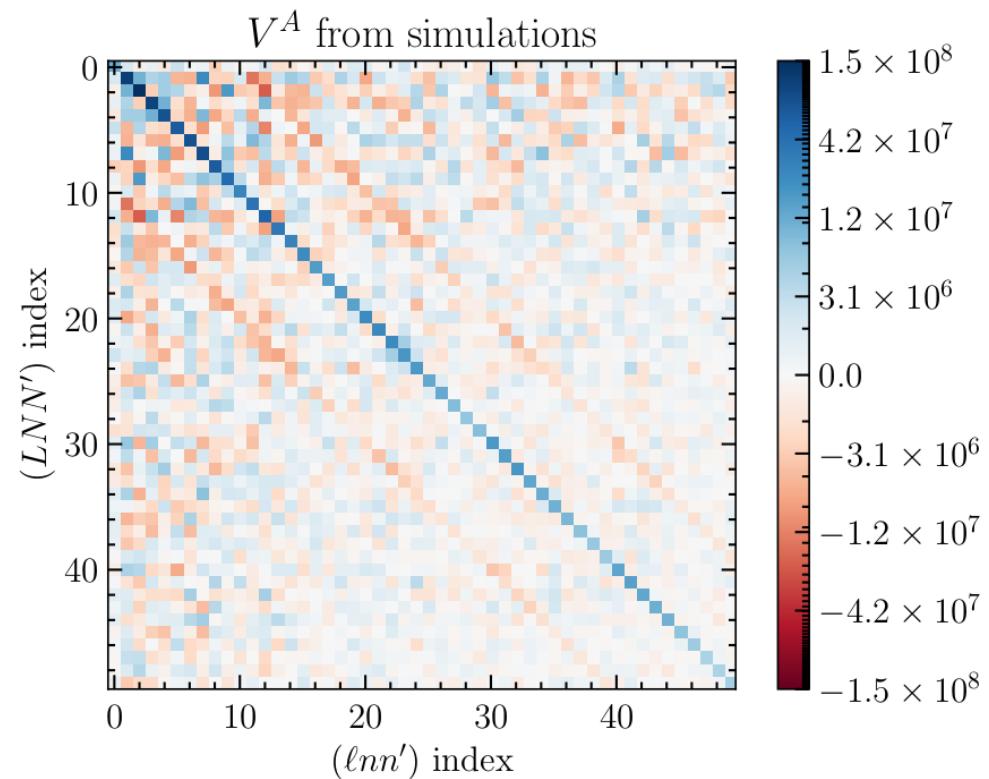
1. Radial Bessel transform with individual galaxies (Leistedt *et al* 2012)
2. Spherical harmonic transform with HEALPix (Gorski *et al* 2005)
3. Pseudo-SFB power spectrum bandpowers with window corrections (Hivon *et al* 2002)

$$\hat{C}_{\ell nn'}^{\text{obs}} = \frac{1}{2\ell + 1} \sum_m \delta_{n\ell m}^{\text{obs}} \delta_{n'\ell m}^{\text{obs},*}$$

4. Subtract shot noise

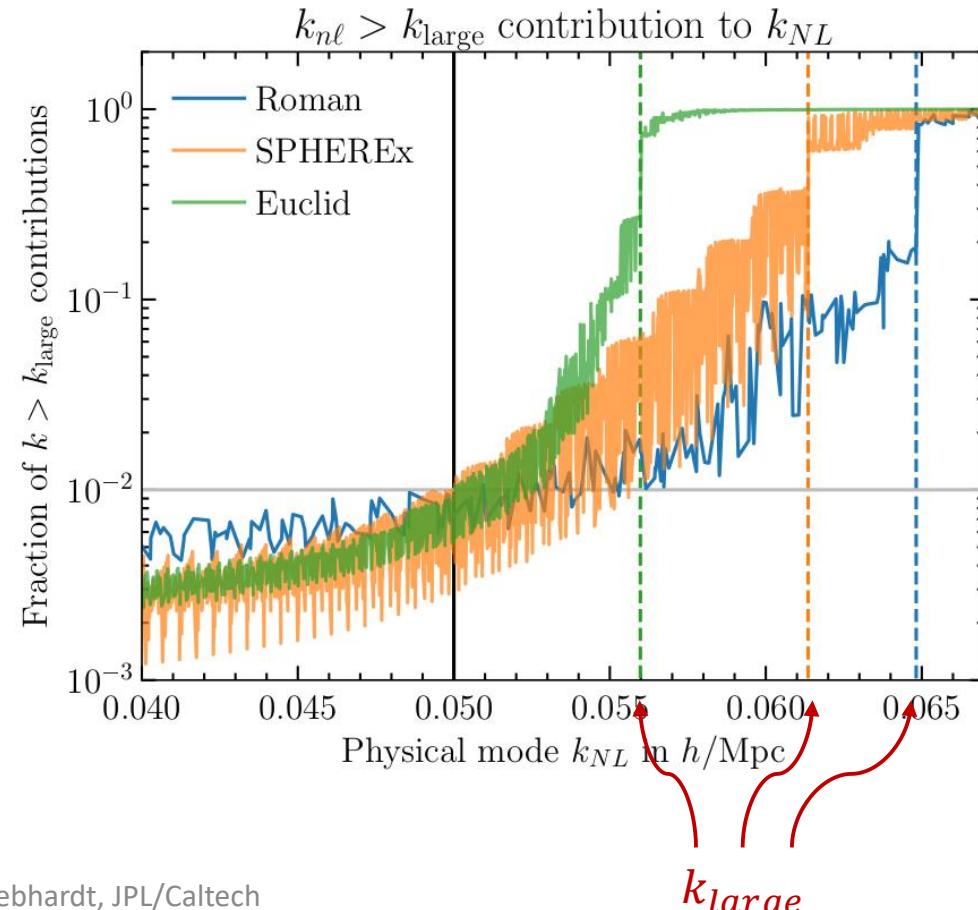
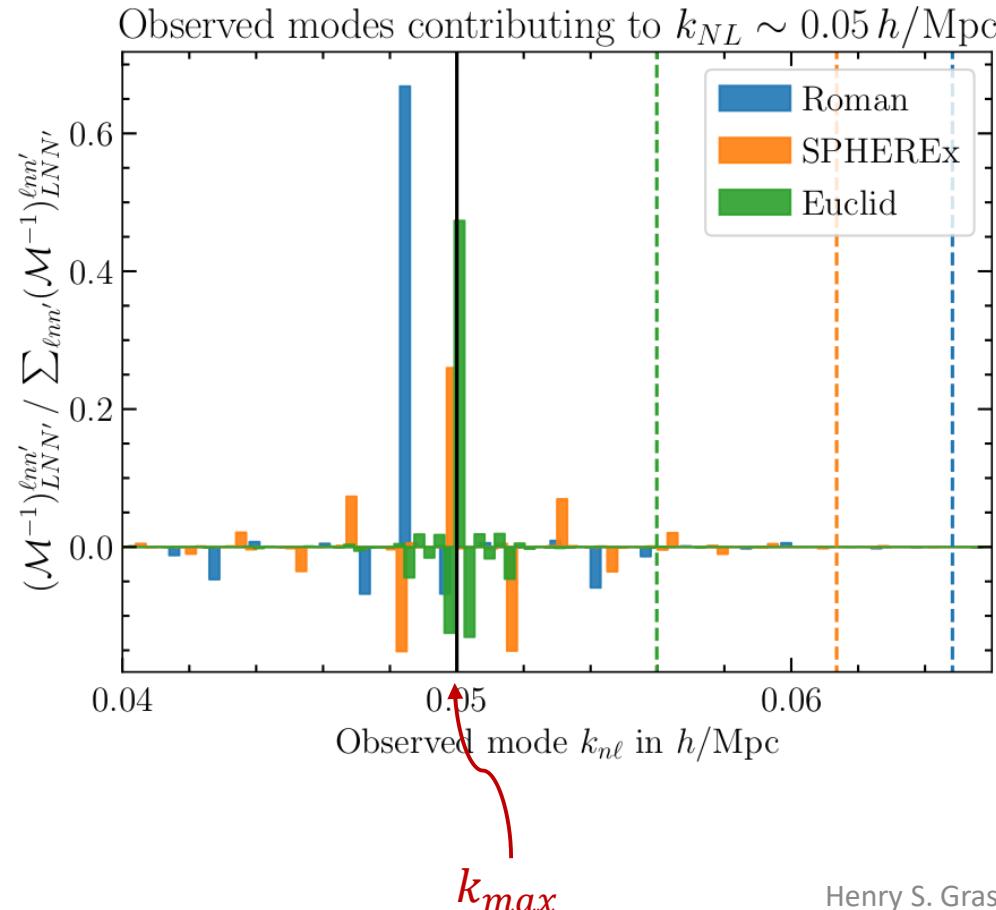


# Covariance matrix $V$

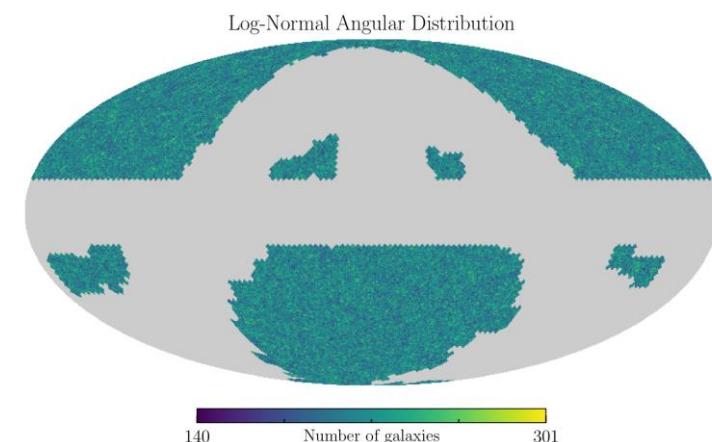
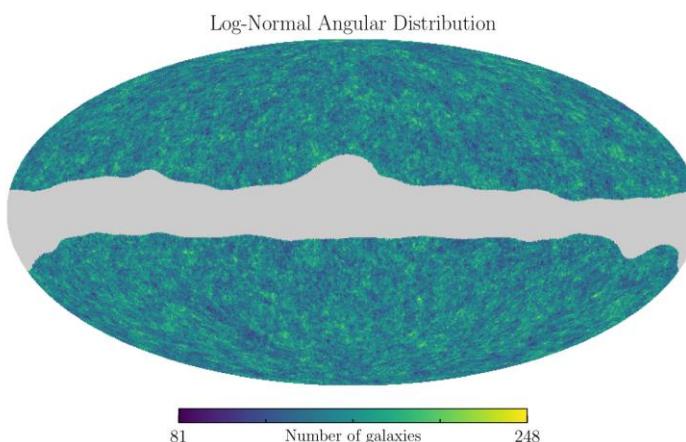
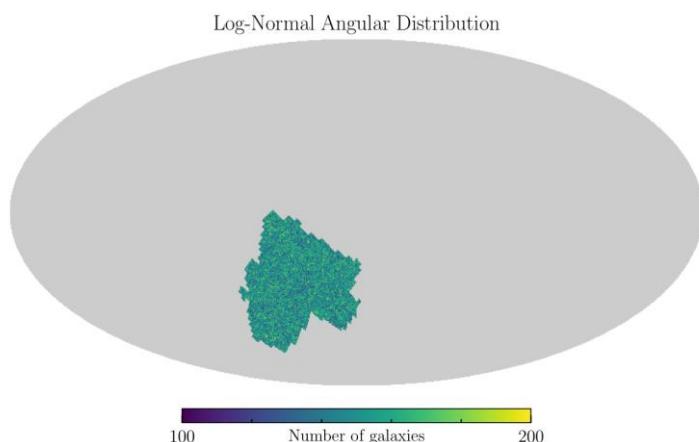
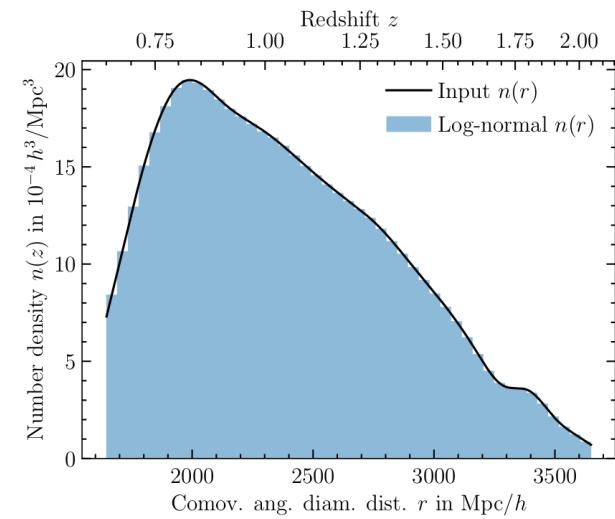
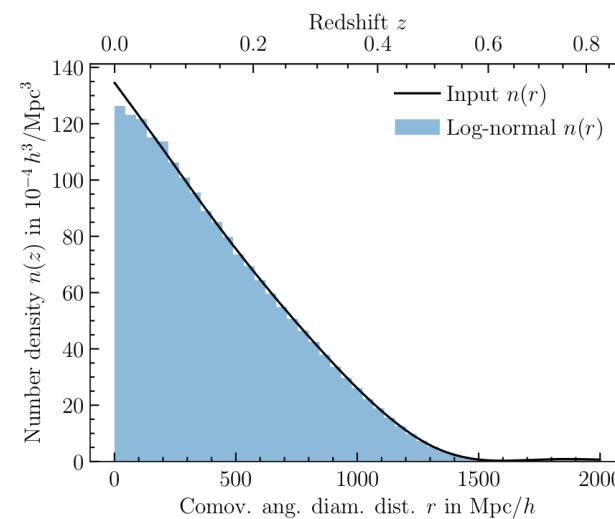
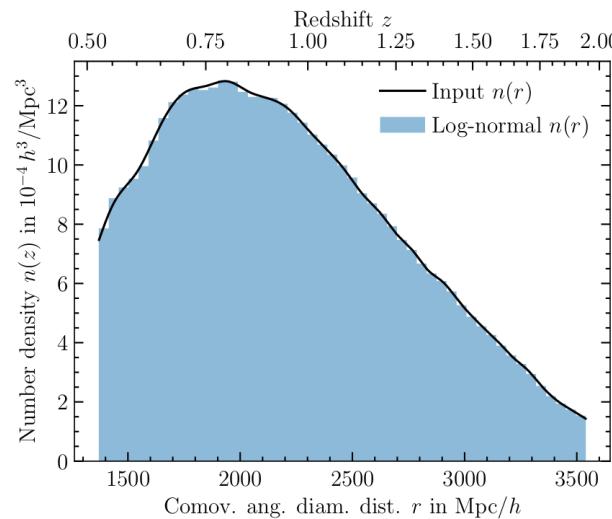


# Full window deconvolution requires modes above $k_{max}$

$$\hat{C}_{\ell nn'}^{\text{obs}} = \sum_{LNN'} \mathcal{M}_{\ell nn'}^{LNN'} \hat{C}_{LNN'}^A$$



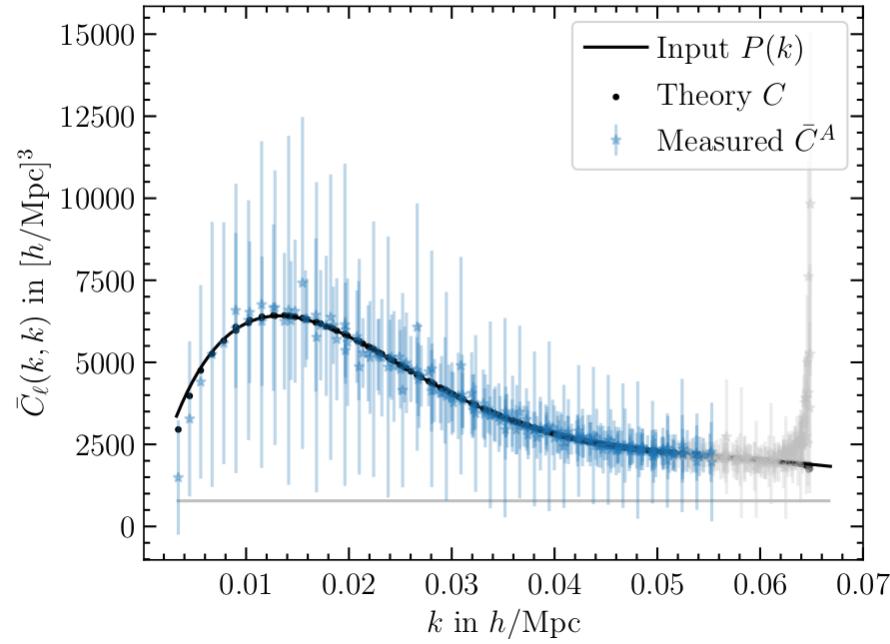
# Examples: *Roman*-, *SPHEREx*-, *Euclid*-like



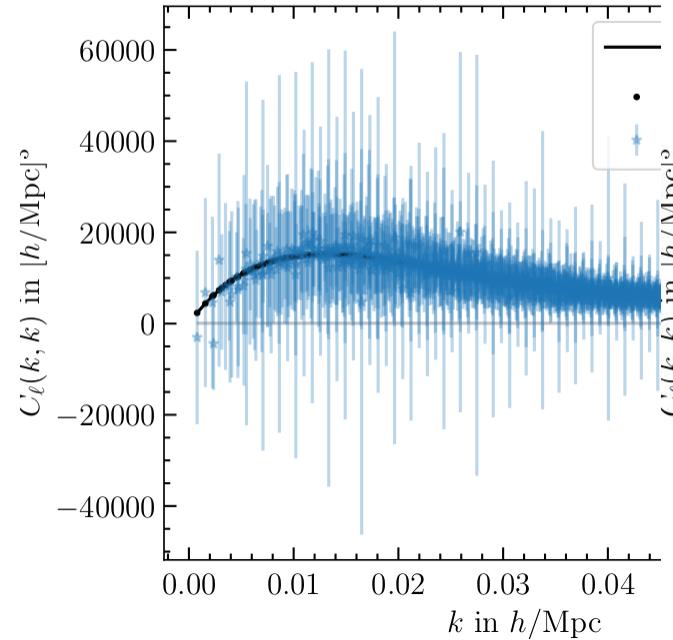
(Special thanks to Katarina Marković)

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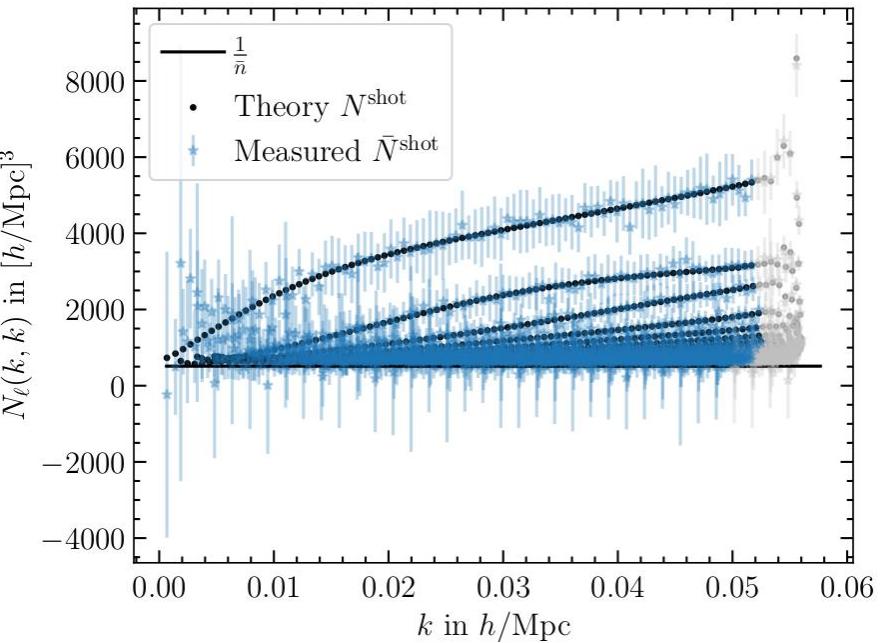
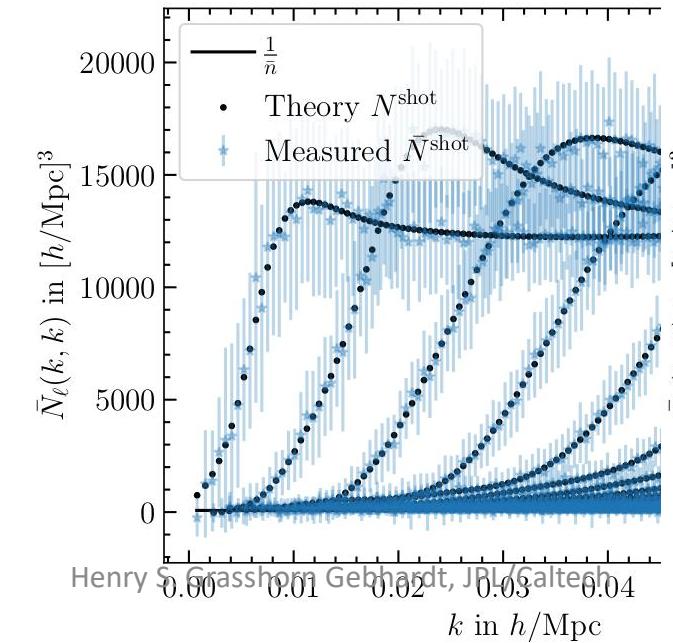
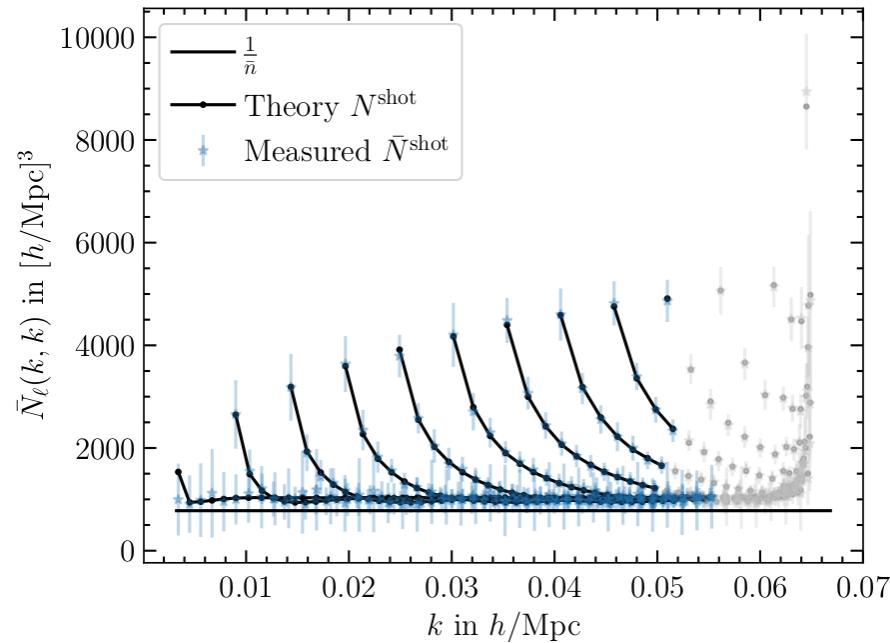
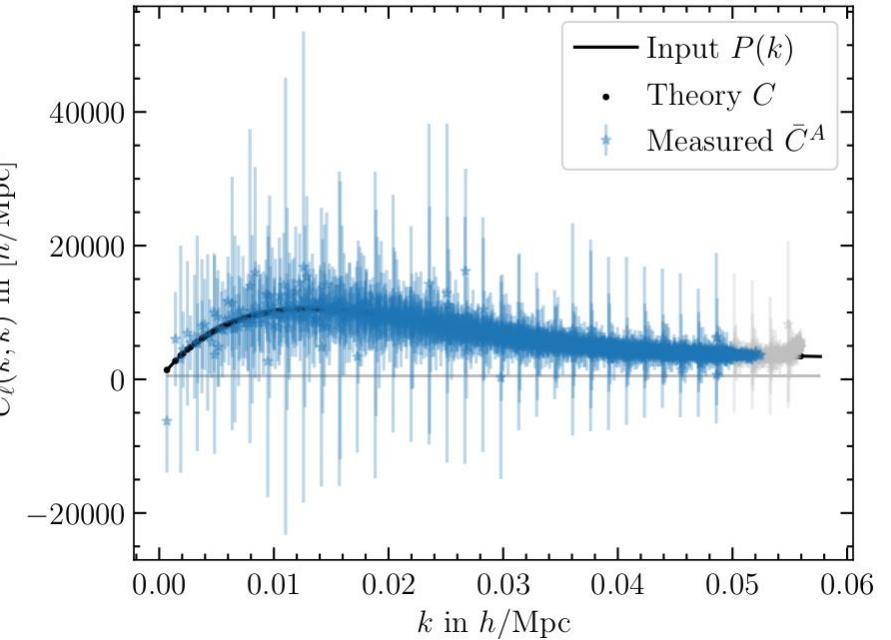
Roman-like



SPHEREx-like



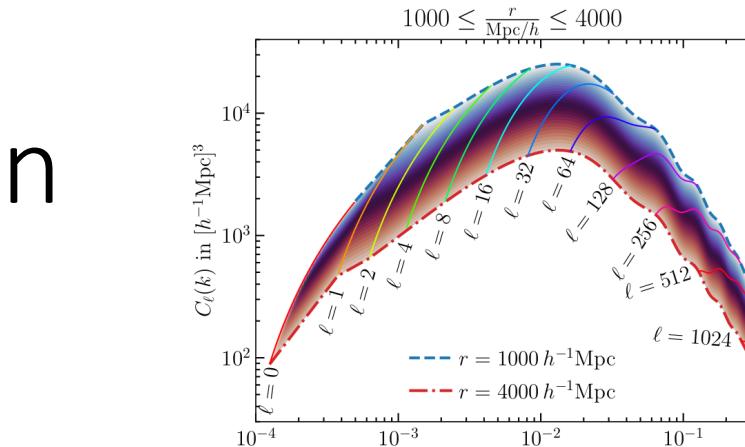
Euclid-like



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# Summary & Conclusion

- SFB is theoretically optimal for **deep and wide surveys**
- Boundary conditions at  $r_{min}$  and  $r_{max}$  make SFB analysis more **numerically stable**
- SuperFaB is feasible for large surveys
- Future:
  - Need fast  $P(k) \rightarrow C_\ell(k, k')$  code
  - Apply to data!



Log-Normal Angular Distribution

